

Open Questions

G.G.Ross, DIS07, Munich, April07



Open Questions Beyond the Standard Model

- The origin of mass?

Due to a Higgs boson? Other physics?

Solution at energy $< 1\text{TeV}$

- Unification of fundamental forces?

At a very high scale?

- Quantum theory of gravity?

(Super)string theory, extra dimensions

- Origin of dark matter?

WIMP SUSY KK ...

LHC answers:

Electroweak spontaneous symmetry breaking $M_{W,Z}, m_{q_i}, m_{l_i}, V_{CKM}, V_{MNS}$

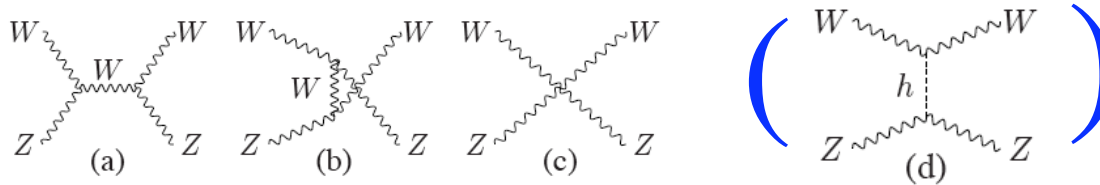
Higgs physics:

- Higgsless (or very heavy Higgs) $KK, \text{Technicolour} \dots$
- Light Higgs $SUSY, \text{eXtra Dimensions, little Higgs} \dots$

Higgsless (or very heavy Higgs)

A light Higgs scalar is needed for perturbative unification

● $W_L Z_L \rightarrow W_L Z_L$

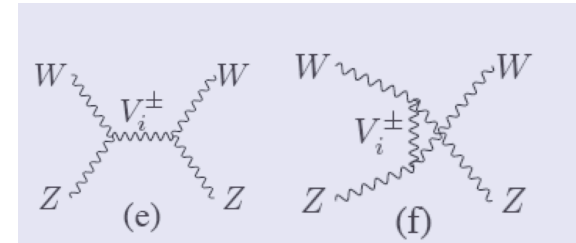
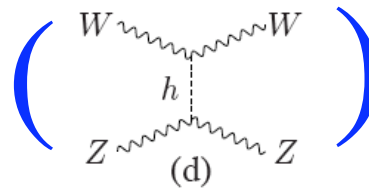
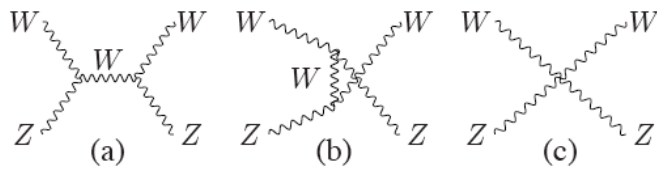


Unitarity violation: $\frac{g^2 \Lambda^2}{16\pi^2 M_W^2} = 1 \Rightarrow \Lambda \sim 1.8 TeV$

Higgsless (or very heavy Higgs)

A light Higgs scalar is needed for perturbative unification

● $W_L Z_L \rightarrow W_L Z_L$



Unitarity violation: $\frac{g^2 \Lambda^2}{16\pi^2 M_W^2} = 1 \Rightarrow \Lambda \sim 1.8 TeV$

$$\Lambda \sim \frac{3\pi^4}{g^2} \frac{M_W^2}{M} = (5-10) TeV$$

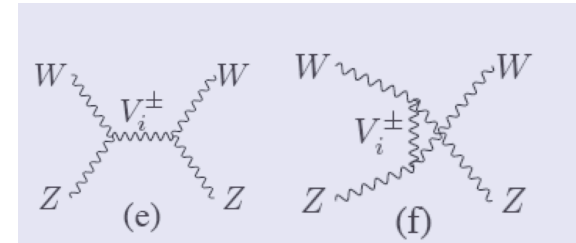
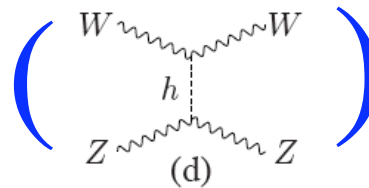
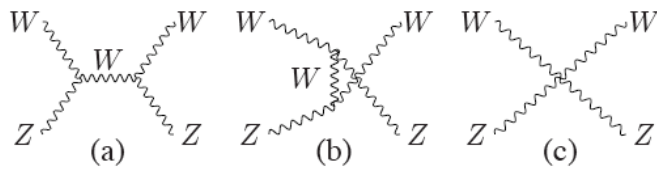
Higgsless

Chivukula, Discus, He

Higgsless (or very heavy Higgs)

A light Higgs scalar is needed for perturbative unification

● $W_L Z_L \rightarrow W_L Z_L$

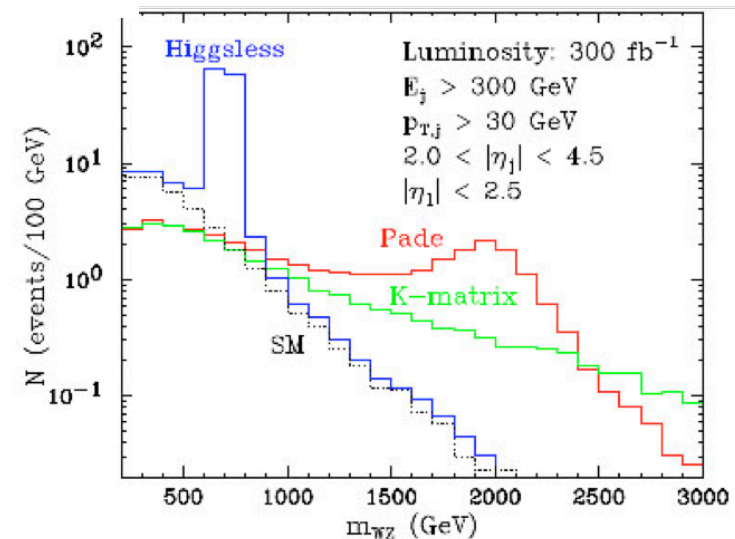
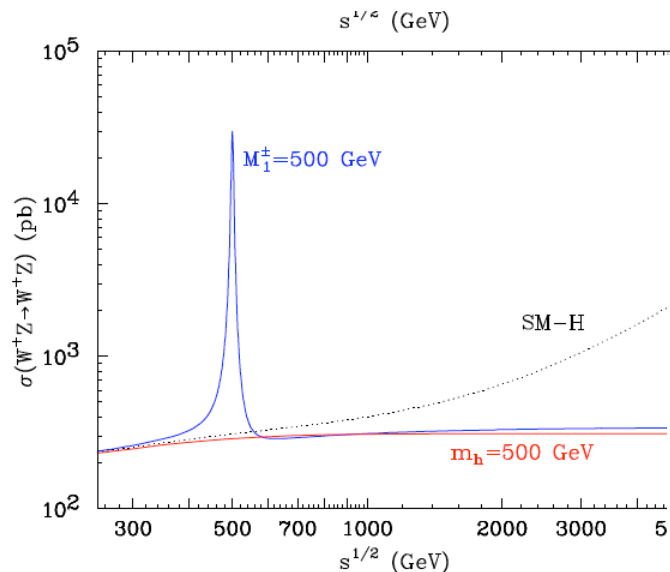


Unitarity violation: $\frac{g^2 \Lambda^2}{16\pi^2 M_W^2} = 1 \Rightarrow \Lambda \sim 1.8 TeV$

$\Lambda \sim \frac{3\pi^4}{g^2} \frac{M_W^2}{M} = (5-10) TeV$

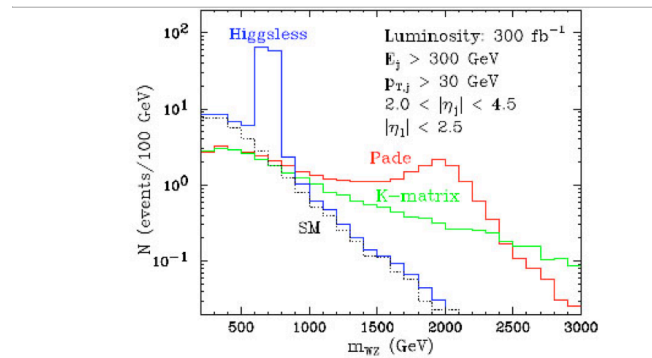
LHC:

W/Z bremsstrahlung off quarks \Rightarrow 2 forward jets + gauge boson pair \Rightarrow 2j+3l+ ~~E_T~~



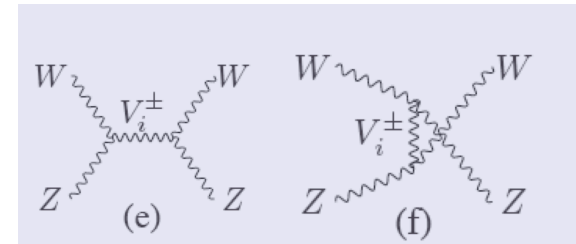
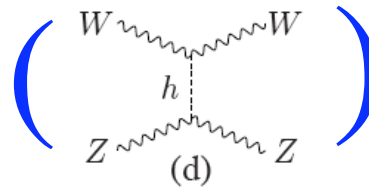
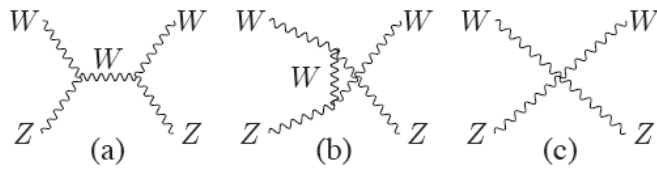
Birkedal,
Matchev,
Perelstein

V^\pm Resonance structure

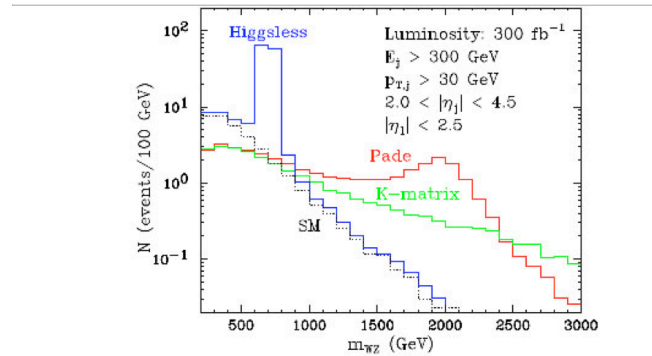


Narrow resonance
c.f. strongly bound state

● $W_L Z_L \rightarrow W_L Z_L$

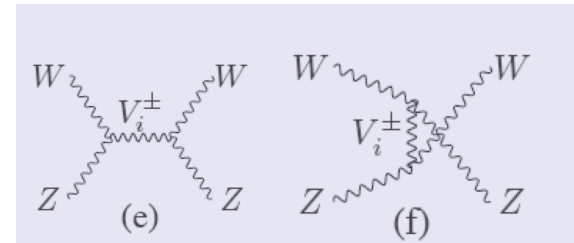
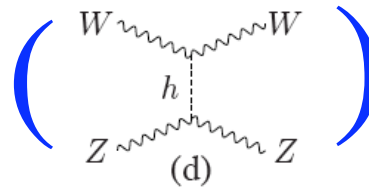
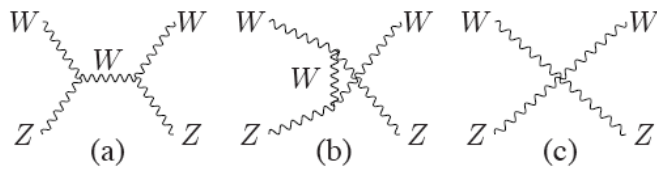


V^\pm Resonance structure



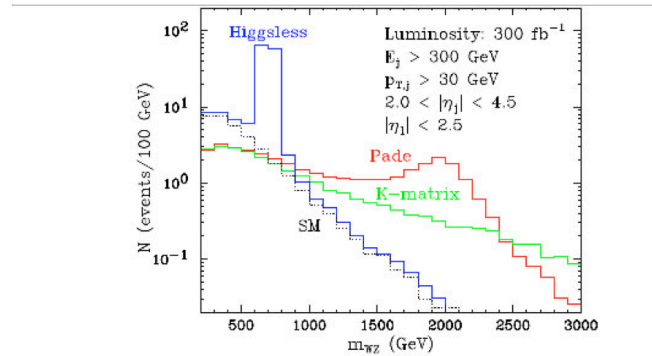
Narrow resonance
c.f. strongly bound state

● $W_L Z_L \rightarrow W_L Z_L$



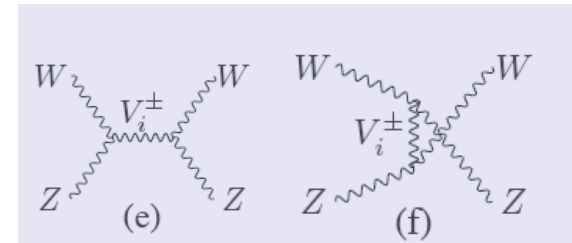
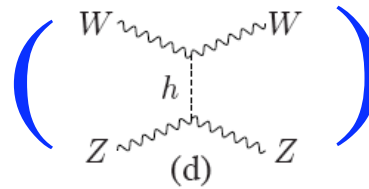
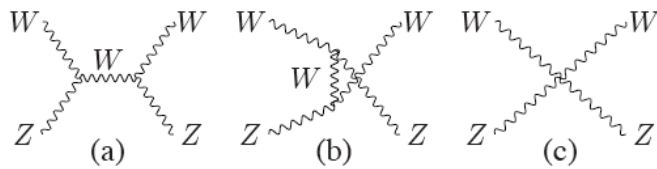
Model	$WW \rightarrow WW$	$WZ \rightarrow WZ$	$WW \rightarrow ZZ$
SM	Yes	No	Yes
Higgsless	Yes	Yes	No

V^\pm Resonance structure



Narrow resonance
c.f. strongly bound state

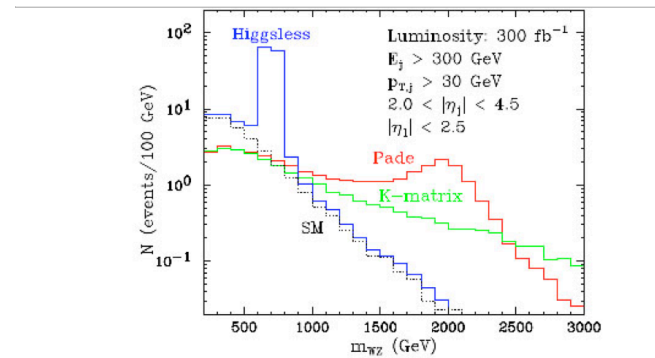
● $W_L Z_L \rightarrow W_L Z_L$



Difficult as large QCD backgrounds in multijet channel

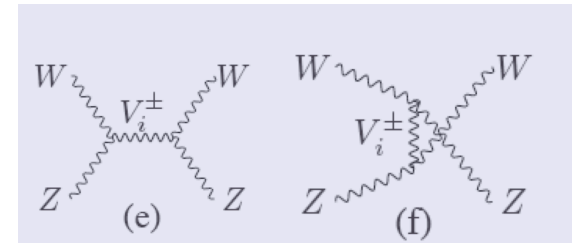
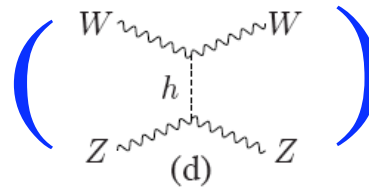
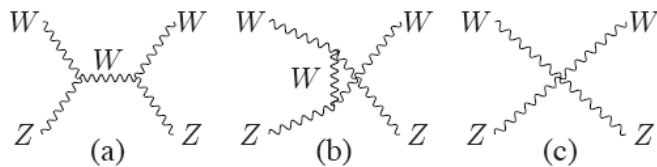
Model	$WW \rightarrow WW$	$WZ \rightarrow WZ$	$WW \rightarrow ZZ$
SM	Yes	No	Yes
Higgsless	Yes	Yes	No

V^\pm Resonance structure



Narrow resonance
c.f. strongly bound state

● $W_L Z_L \rightarrow W_L Z_L$



Difficult as large QCD backgrounds in multijet channel

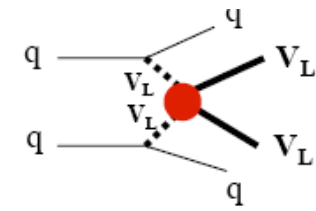
Model	$WW \rightarrow WW$	$WZ \rightarrow WZ$	$WW \rightarrow ZZ$
SM	Yes	No	Yes
Higgsless	Yes	Yes	No

Mass measurement can establish Higgsless models

$$g_{WWZZ} = g_{WWZ}^2 + \sum_i (g_{WZV}^{(i)})^2, \quad \text{Csaki et al}$$

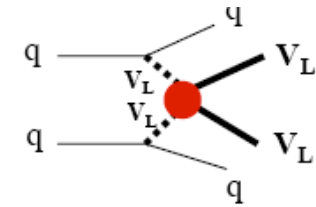
$$2(g_{WWZZ} - g_{WWZ}^2)(M_W^2 + M_Z^2) + g_{WWZ}^2 \frac{M_Z^4}{M_W^2} = \sum_i (g_{WZV}^{(i)})^2 \left[3(M_i^\pm)^2 - \frac{(M_Z^2 - M_W^2)^2}{(M_i^\pm)^2} \right]$$

Strongly coupled vector boson sector



Unitarity saturation signals strong interaction,-
observation of no excess means weakly coupled quanta below 1TeV

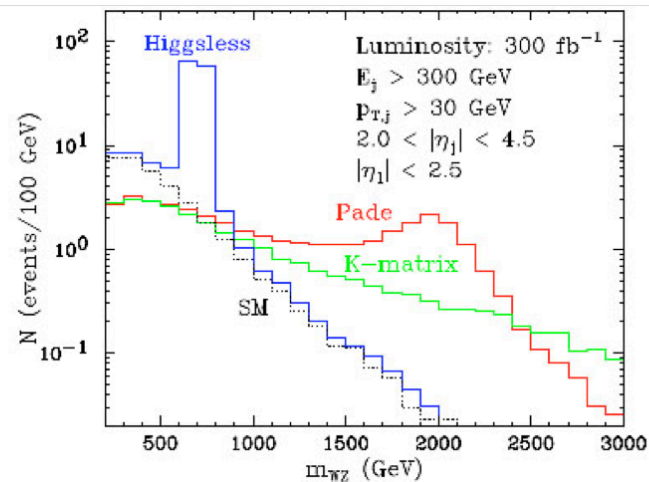
Strongly coupled vector boson sector



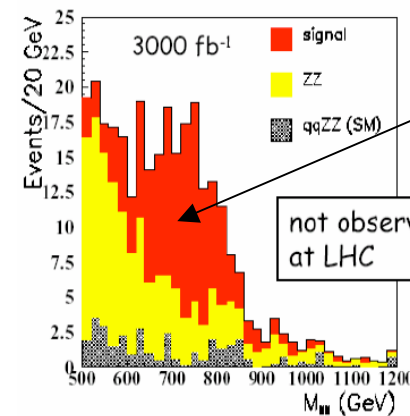
Unitarity saturation signals strong interaction,-
observation of no excess means weakly coupled quanta below 1TeV

Chanowitz “No-lose” theorem

● Unfortunately



Scalar resonance $Z_L Z_L \rightarrow 4\ell$

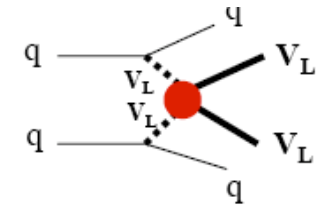


Need v accurate q dist fns..

not observable
at LHC

De Rocek

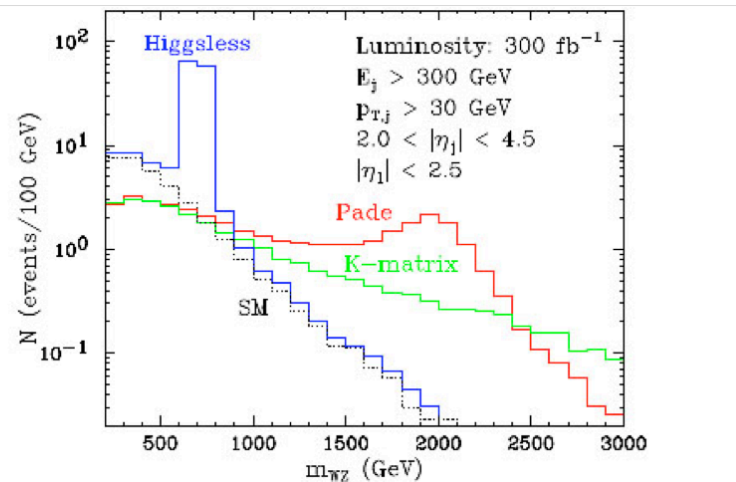
Strongly coupled vector boson sector



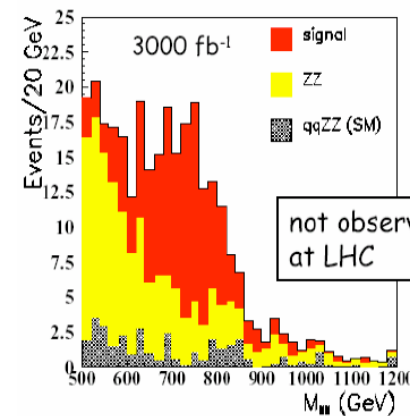
Unitarity saturation signals strong interaction,-
observation of no excess means weakly coupled quanta below 1TeV

Chanowitz “No-lose” theorem

- Unfortunately

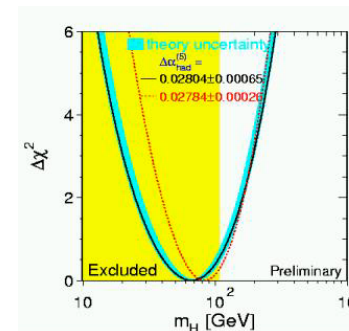


Scalar resonance $Z_L Z_L \rightarrow 4\ell$

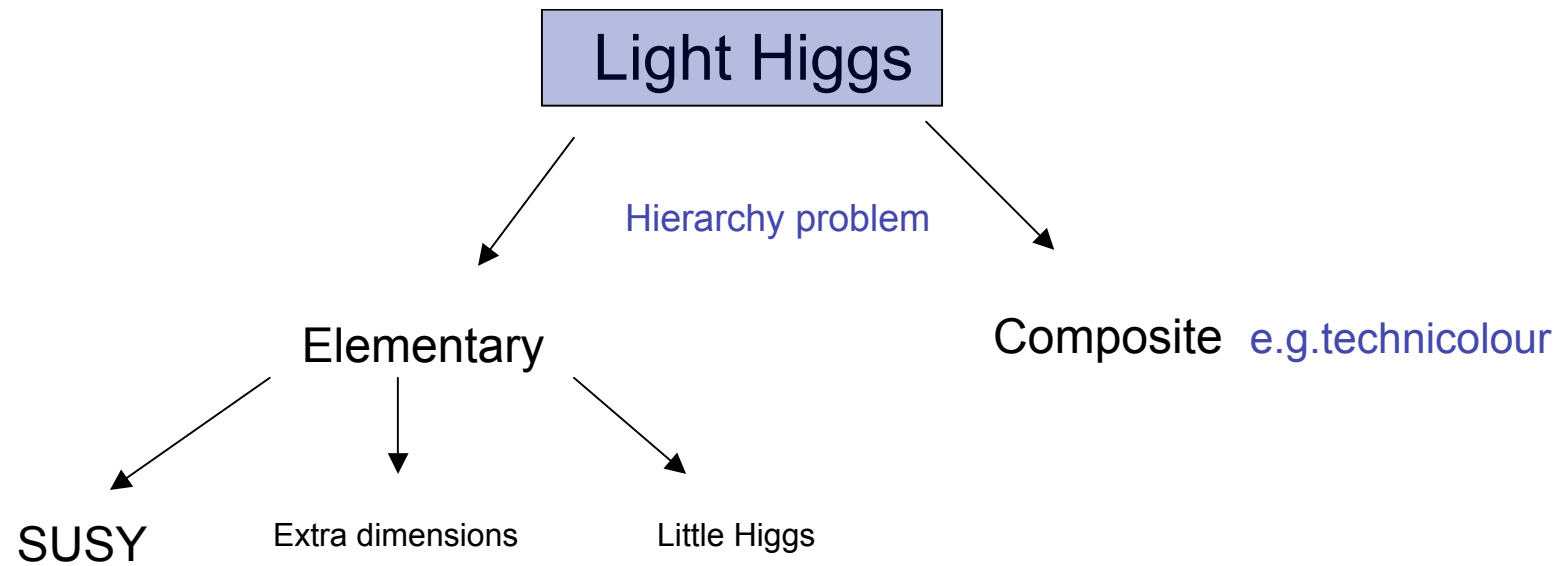


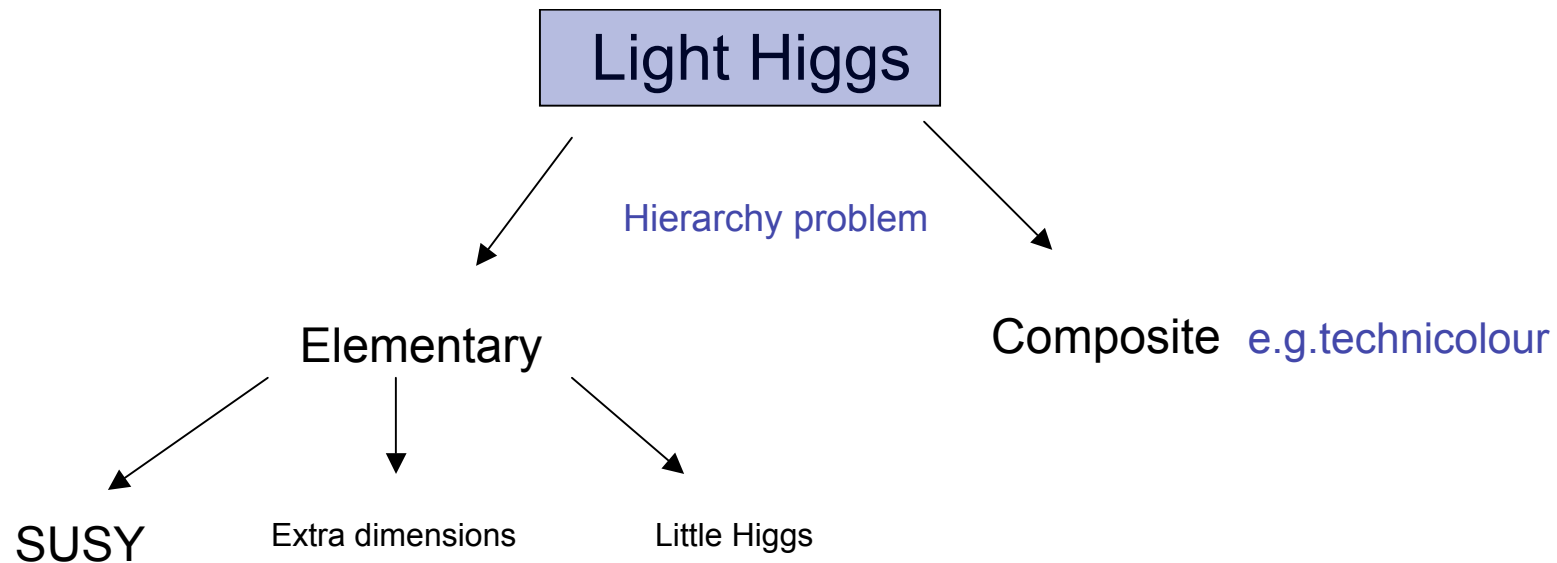
De Rocek

- But precision tests suggest strong interaction cannot be all

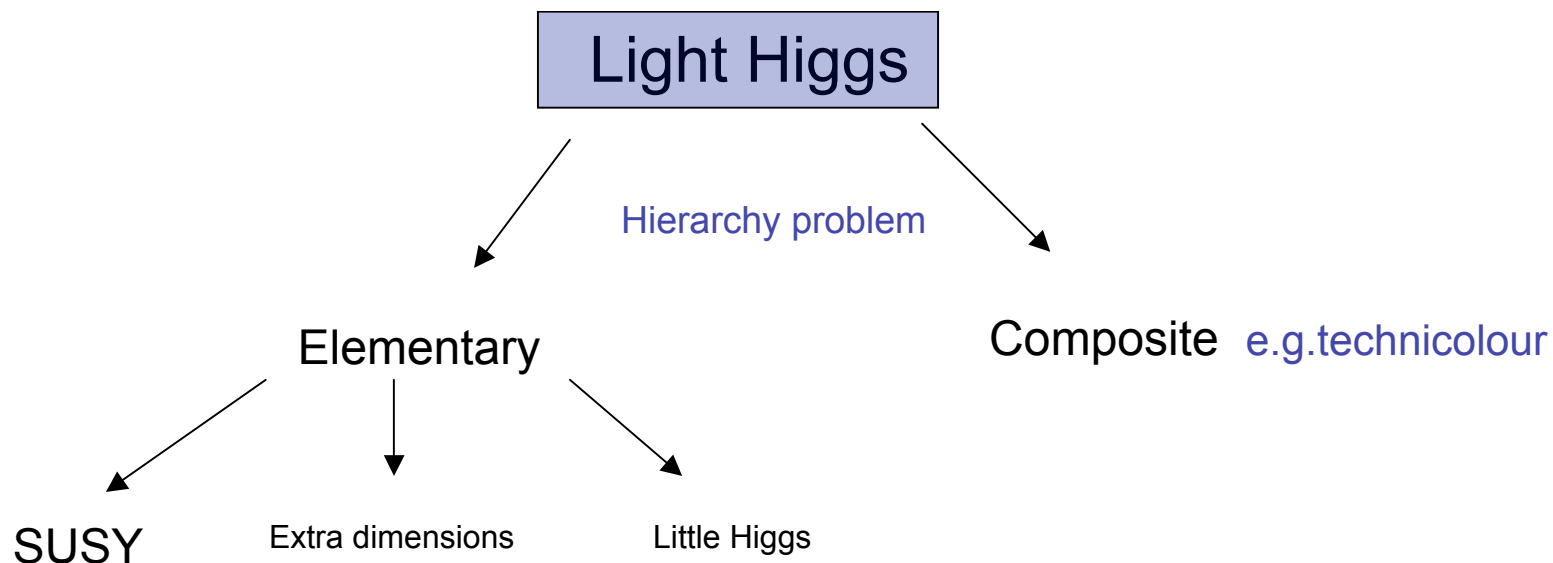


Light Higgs



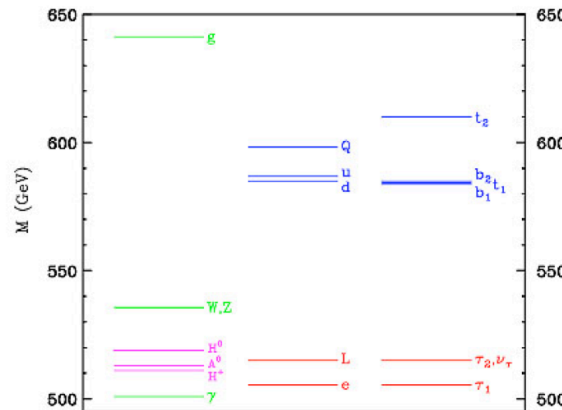
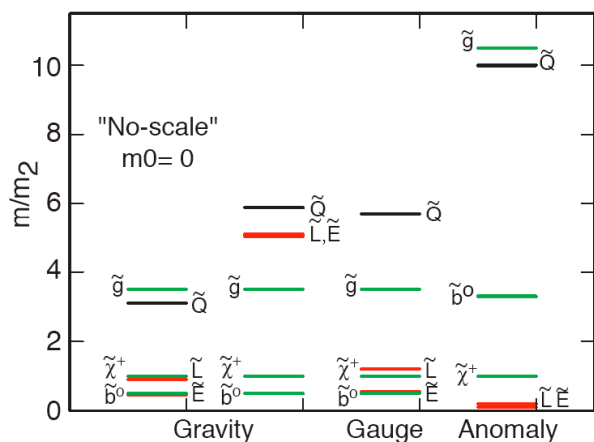


How will we distinguish these possibilities?



How will we distinguish them?

e.g. SUSY and UED have similar level structure

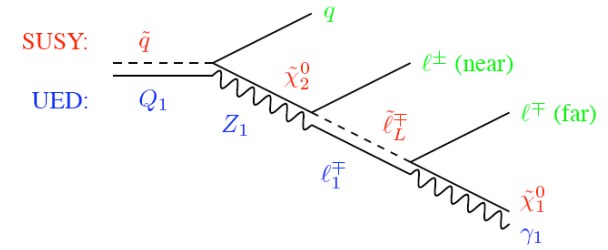


Similar radiative corrections

Cheng
et al

Also the couplings of both the SUSY and KK states are the same as their SM partners

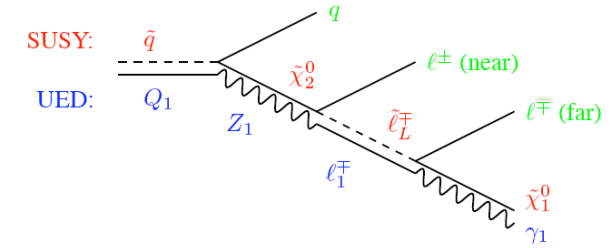
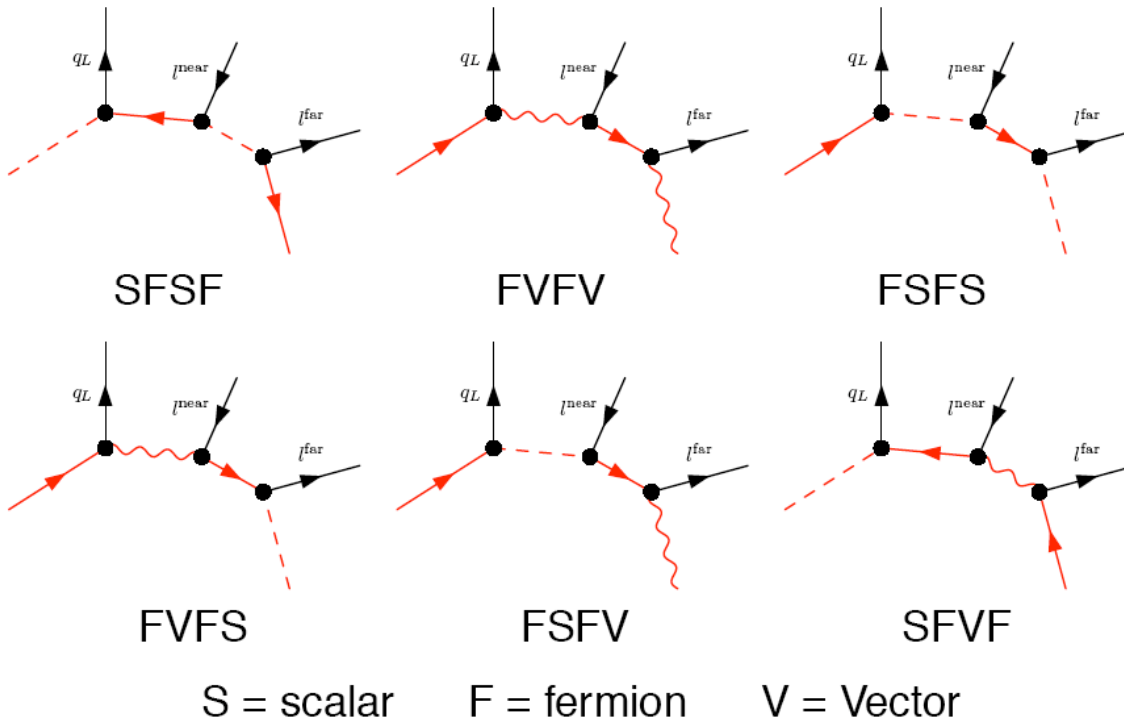
Spin discriminant



Spin discriminant

Cascade Decay Chains

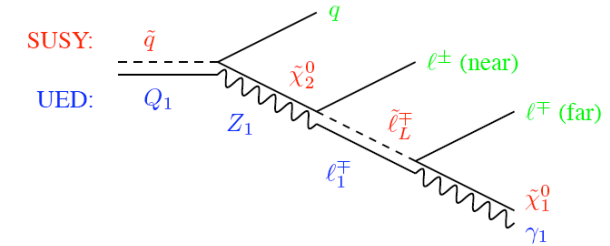
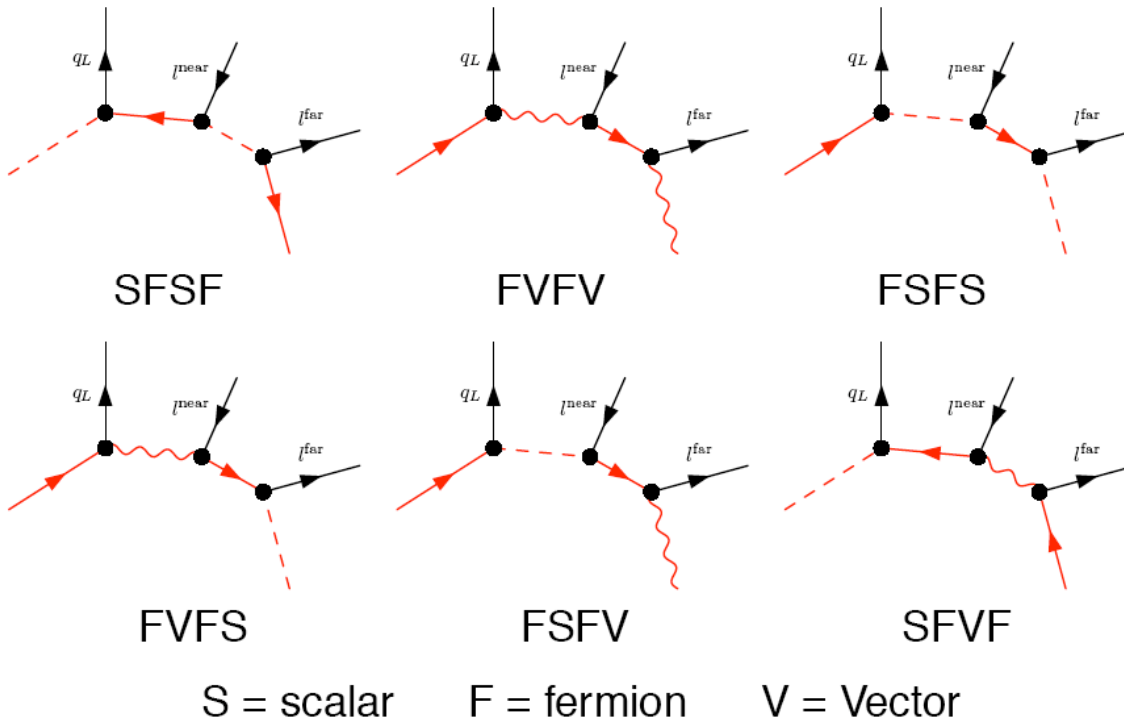
There are 6 possible spin assignments:



Spin discriminant

Cascade Decay Chains

There are 6 possible spin assignments:



Discriminant:

$$\frac{dP}{dm} = \frac{1}{\Gamma} \frac{d\Gamma}{dm}$$

Γ total decay rate of chain

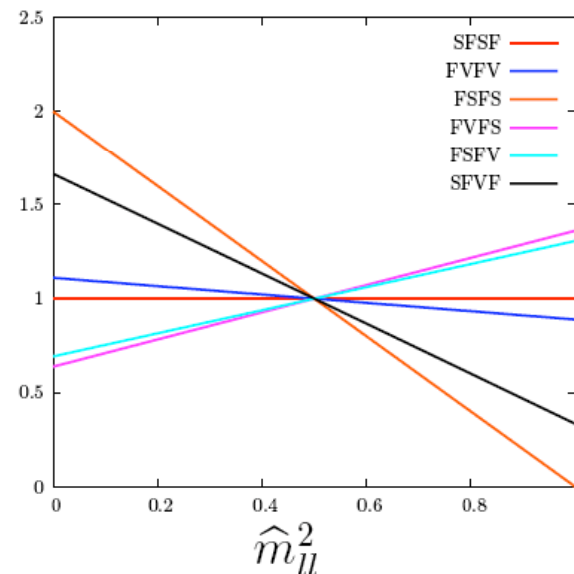
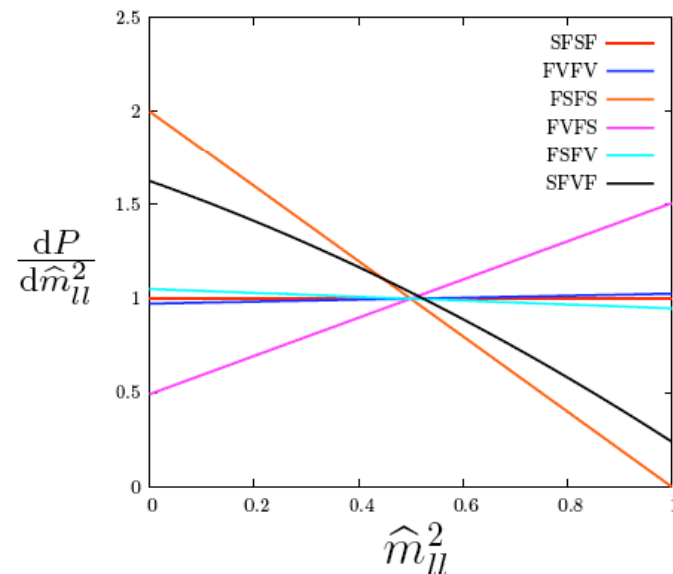
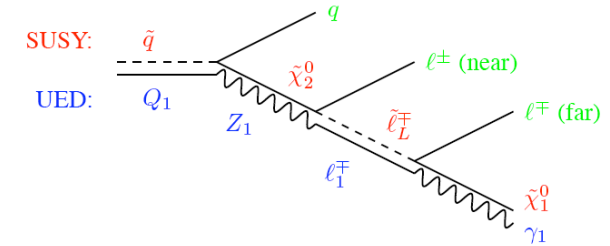
$$m_{ab}^2 = (p_a + p_b)^2$$

Barr

Athanasίου, Lester, Smillie, Webber

For Example, l^+l^-

The m_{ll}^2 distributions for SPS 1a masses and UE ($R^{-1} = 800\text{GeV}, \Lambda R = 20$) are:



$$\frac{dP}{dm} = \frac{1}{\Gamma} \frac{d\Gamma}{dm}$$

Γ total decay rate of chain

Discrimination

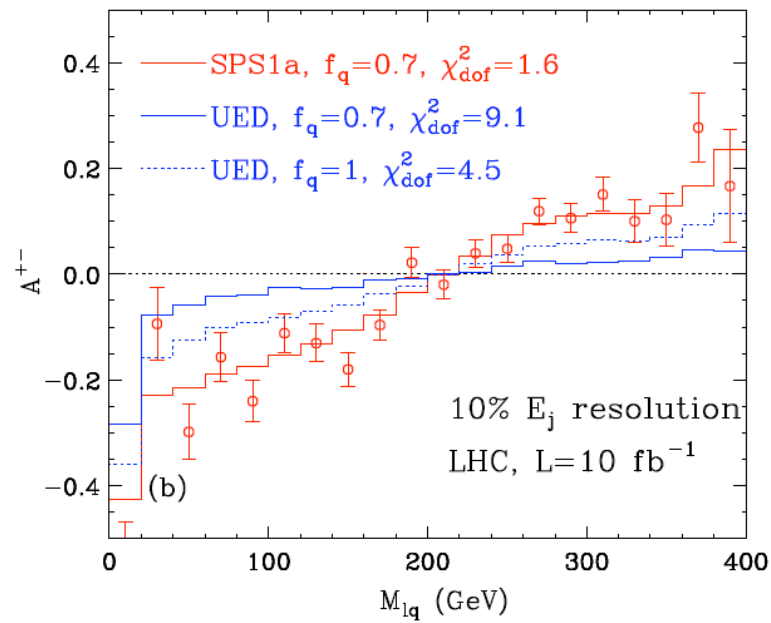
$T \downarrow S \rightarrow$	SFSF	FVFS	FSFS	FVFS	FSFV	SFVF
SFSF	∞	60486	23	148	15608	66
FVFS	60622	∞	22	164	6866	62
FSFS	36	34	∞	16	39	266
FVFS	156	173	11	∞	130	24
FSFV	15600	6864	25	122	∞	76
SFVF	78	73	187	27	90	∞

Number of events, assuming FSFS is true, such that FSFS is 1000 times more likely than other model.

Athanasiou, Lester, Smillie, Webber

Caveats: jet identification, mass measurements uncertain, charge asymmetry only comes from quark production...

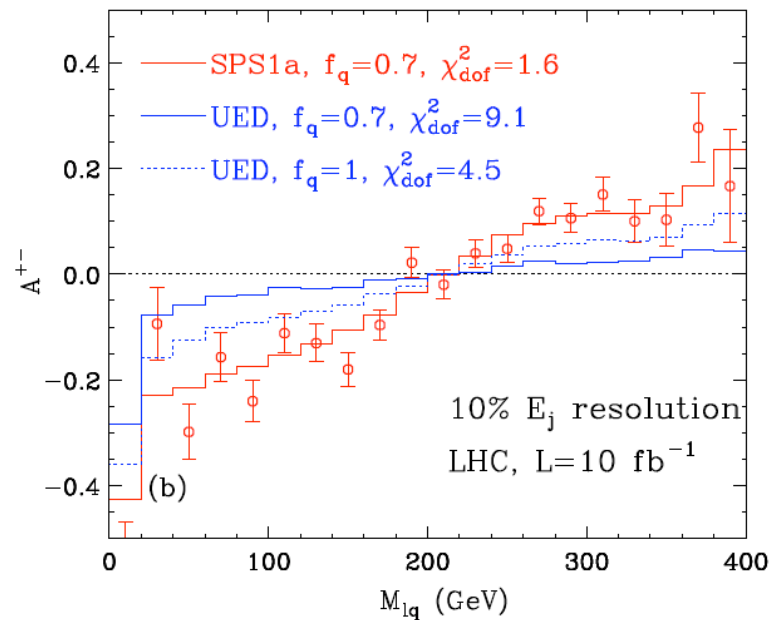
e.g. $\frac{dP}{dm_{ll}^2}$ - Difference between SUSY and UED $\propto \beta(m_i)$, small in favoured parameter space



Kong, Matchev

Caveats: jet identification, mass measurements uncertain, charge asymmetry only comes from quark production...

e.g. $\frac{dP}{dm_{ll}^2}$ - Difference between SUSY and UED $\propto \beta(m_i)$, small in favoured parameter space



Alternative methods being developed may offer complementary information

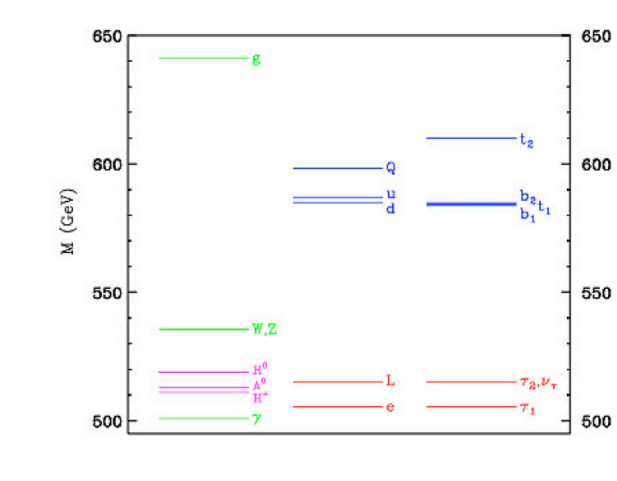
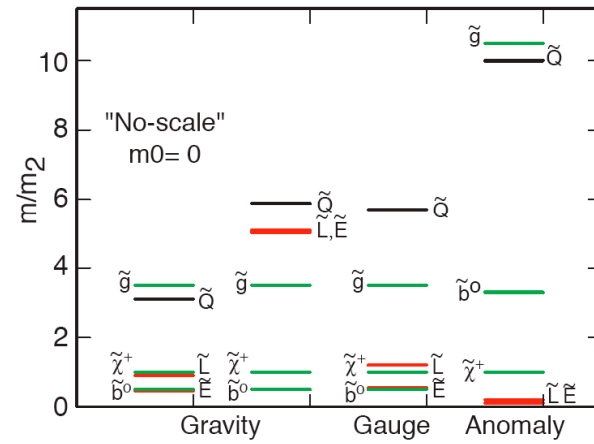
e.g. $\left(\frac{d\sigma}{d\cos\theta^*} \right)_i$

θ^* angle between incoming quark and lepton

$q\bar{q} \rightarrow Z^0/\gamma \rightarrow \tilde{\ell}^+\tilde{\ell}^- \rightarrow \tilde{\chi}_1^0\ell^+ \tilde{\chi}_1^0\ell^-$,
 $q\bar{q} \rightarrow Z^0/\gamma \rightarrow \ell_1^+\ell_1^- \rightarrow \gamma_1 \ell^+\gamma_1 \ell^-$.

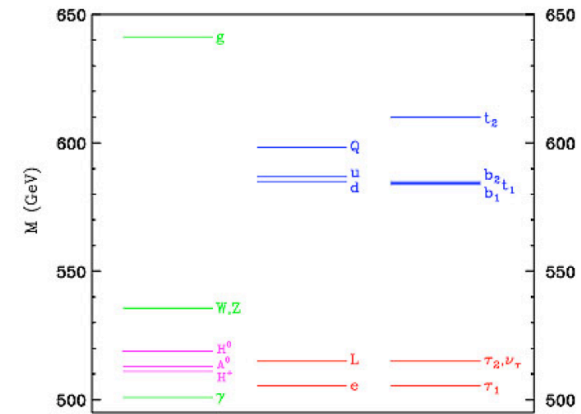
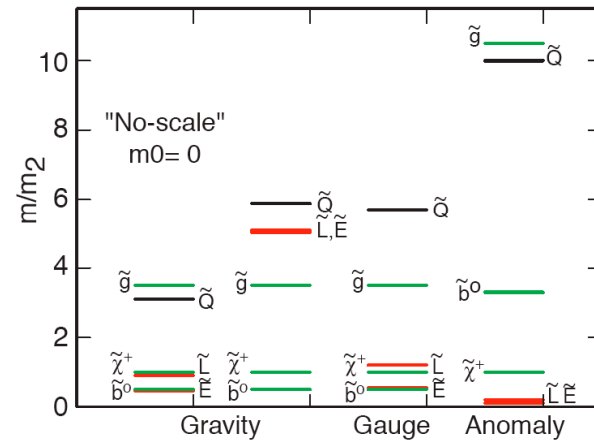
Barr

Spectrum



What do we learn from measuring the spectrum?

Spectrum



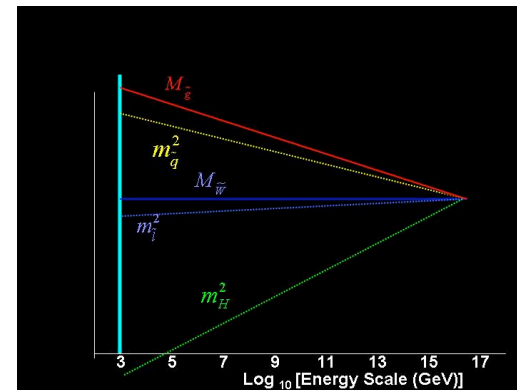
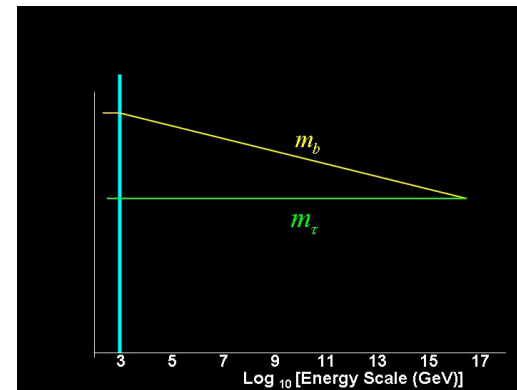
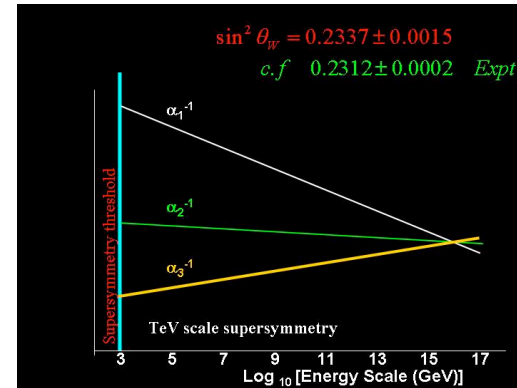
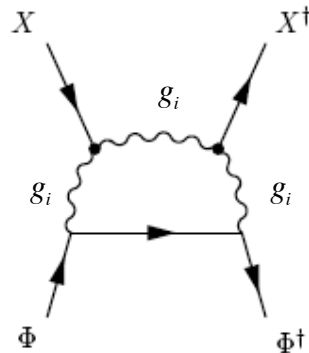
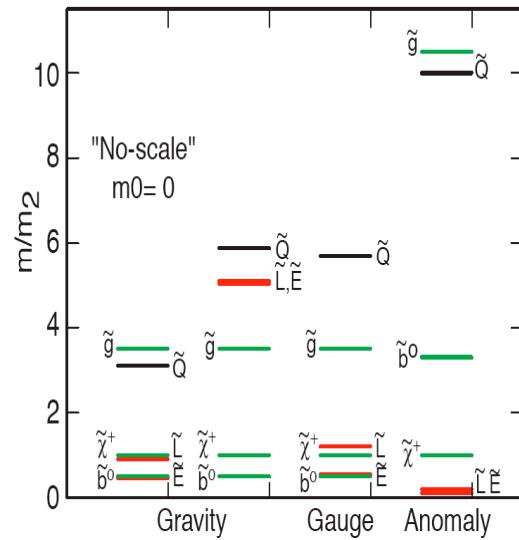
What do we learn from measuring the spectrum?

SUSY: SUSY breaking mechanism, hidden sector, unification

XD: Compactification (need higher levels)

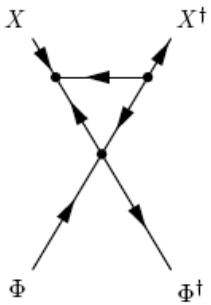
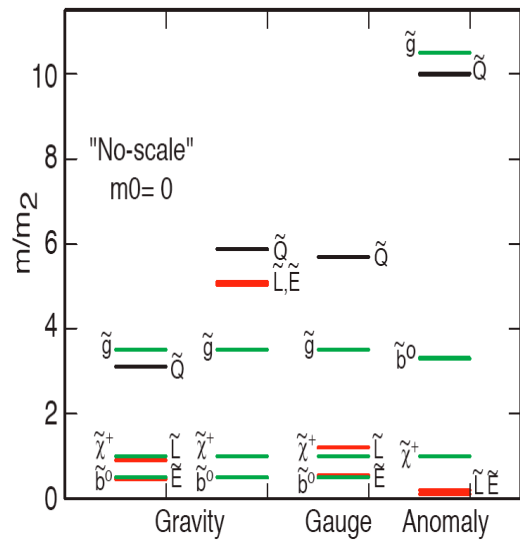
Dark matter abundance

SUSY: Soft mass unification

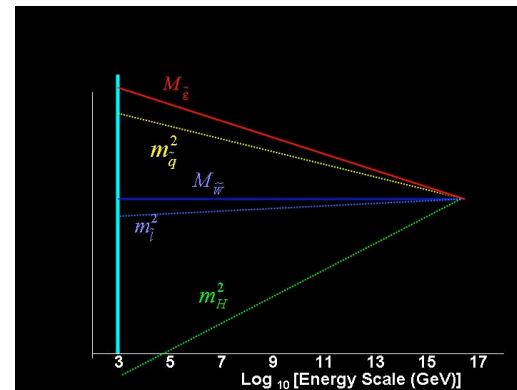
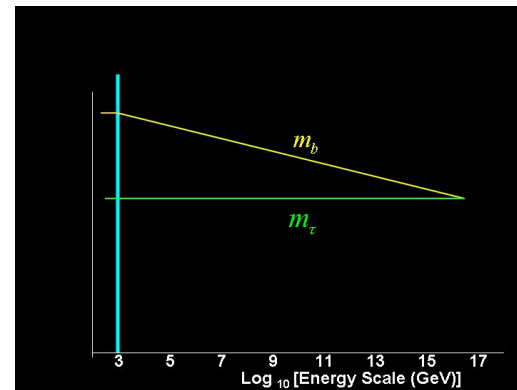
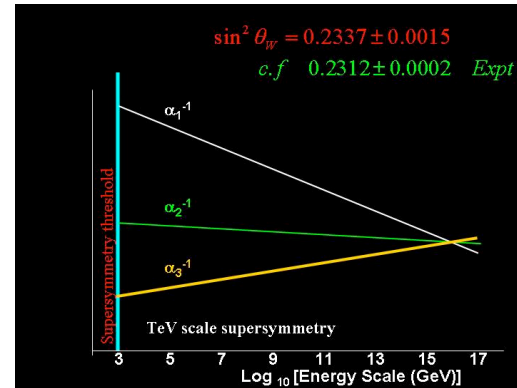
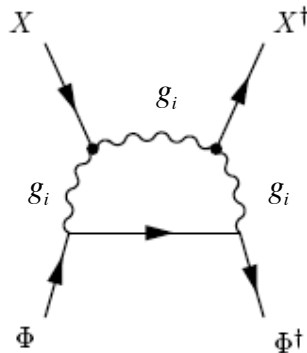


$SU(3) \otimes SU(2) \otimes U(1) \rightarrow SU(3) \otimes U(1)_{EM}$

SUSY: Soft mass unification

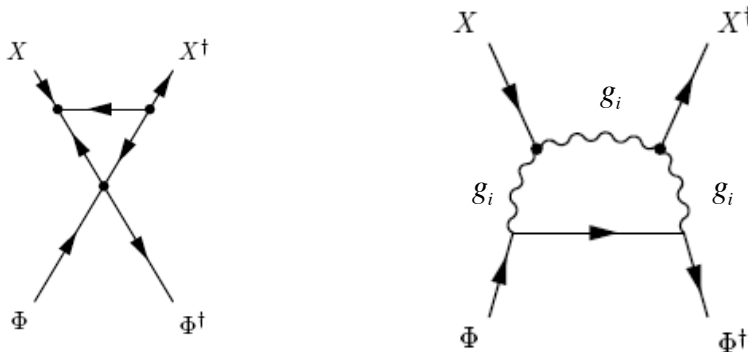
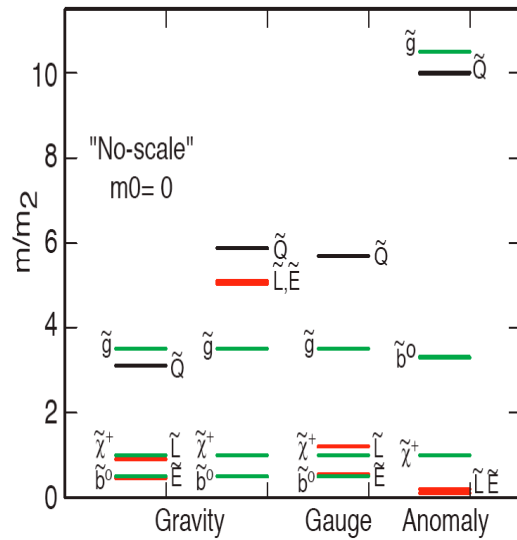


Cohen, Roy, Schmaltz

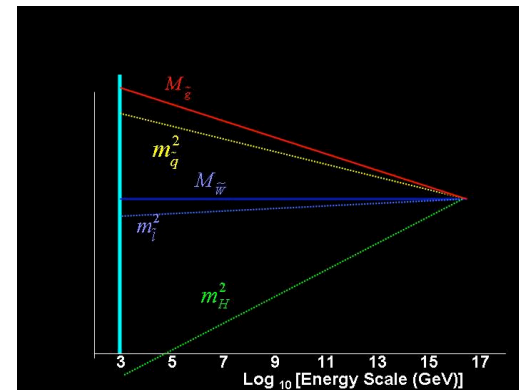
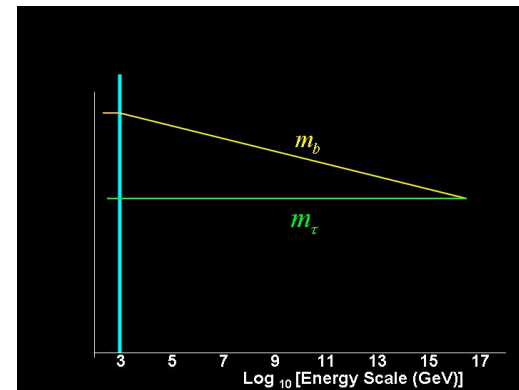
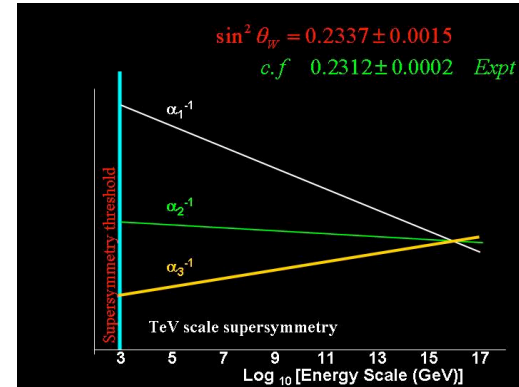


$$SU(3) \otimes SU(2) \otimes U(1) \rightarrow SU(3) \otimes U(1)_{EM}$$

SUSY: Soft mass unification



- Gaugino mass unification unaffected
- Soft mass unification (usually) unchanged

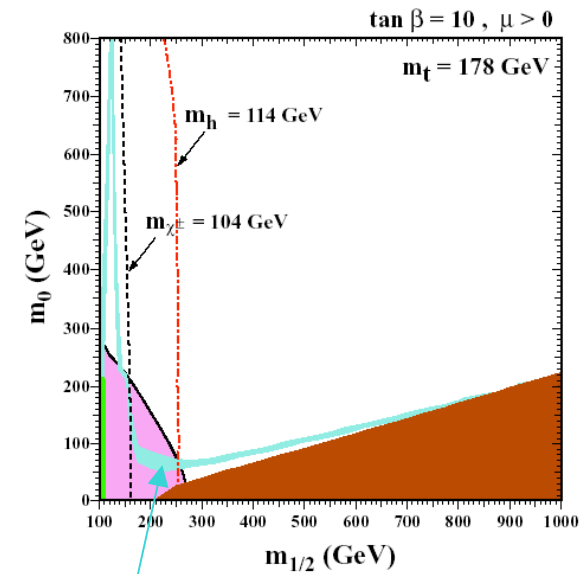


$$SU(3) \otimes SU(2) \otimes U(1) \rightarrow SU(3) \otimes U(1)_{EM}$$

Neutralino LSP

Dark matter abundance

Dark matter abundance in SUSY LSP very sensitive to slepton mass(es)



WMAP constraint on relic density

Ellis, Olive, Santoso, Spanos

Dark matter abundance

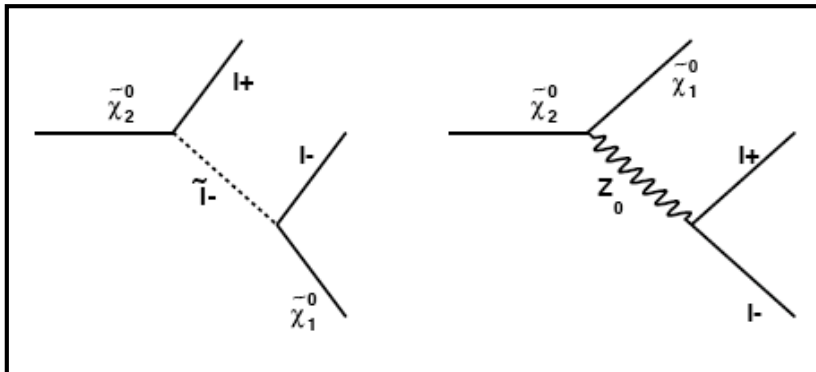
Dark matter abundance in SUSY LSP very sensitive to slepton mass(es)

Slepton mass measurement at the LHC

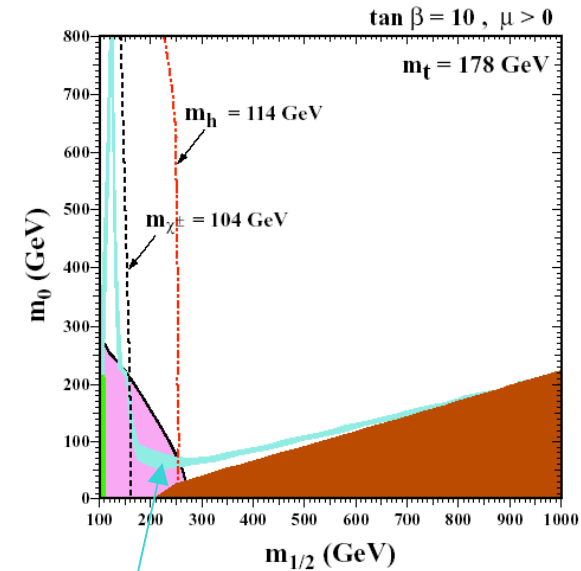
- Direct production channel has large WW, tt backgrounds
- Cascade decays promising - dilepton invariant mass distribution

e.g. $\tilde{\chi}_2^0 \rightarrow l^\pm l^\mp \tilde{\chi}_1^0$

Endpoint for virtual intermediate states: $m_{ll,\max} = m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$



Neutralino LSP



WMAP constraint on relic density

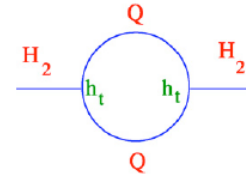
Ellis, Olive, Santoso, Spanos

Event rate for mSUGRA study points:

Point	M_0	$M_{\frac{1}{2}}$	$M_{\tilde{t}}$	σ	$N(10\text{fb}^{-1})$
A	40 GeV	189 GeV	92 GeV	170 pb	$1.7 \cdot 10^6$
B	150 GeV	187 GeV	96 GeV	150 pb	$1.5 \cdot 10^6$
C	3280 GeV	300 GeV	3277 GeV	4.4 pb	44,000
$t\bar{t}$ (SM background)	NA	NA	NA	425 pb	$4.25 \cdot 10^6$

Birkedal, Group, Matchev

Heavy SUSY - Little Higgs



Little hierarchy problem:

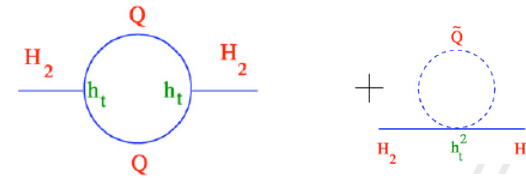
Standard Model

$$\delta m_H^2 \propto y_t^2 \Lambda_{UV}^2$$

Fine tuned

Heavy SUSY - Little Higgs

Little hierarchy problem:



Standard Model

$$\delta m_H^2 \propto y_t^2 \Lambda_{UV}^2$$

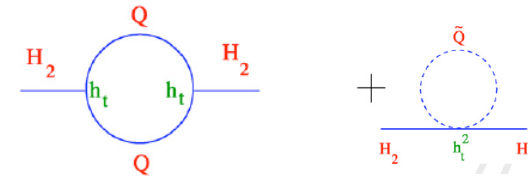
Fine tuned

SUSY

$$\delta m_H^2 \simeq -\frac{3y_t^2}{8\pi^2} m_{\tilde{t}}^2 \log \left(\frac{\Lambda_{UV}^2}{m_{\tilde{t}}^2} \right)$$

Fine tuned

Heavy SUSY - Little Higgs



Little hierarchy problem:

Standard Model

$$\delta m_H^2 \propto y_t^2 \Lambda_{UV}^2$$

Fine tuned

SUSY

$$\delta m_H^2 \simeq -\frac{3y_t^2}{8\pi^2} m_{\tilde{t}}^2 \log \left(\frac{\Lambda_{UV}^2}{m_{\tilde{t}}^2} \right)$$

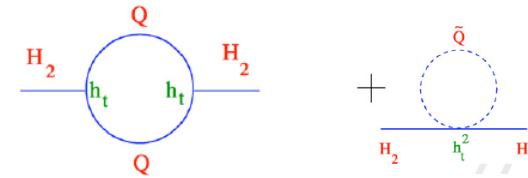
Fine tuned

Little Higgs

$$\delta m_H^2 \propto \frac{y_t^2}{16\pi^2} f^2$$

(Pseudo Goldstone boson)

Heavy SUSY - Little Higgs



Little hierarchy problem:

Standard Model

$$\delta m_H^2 \propto y_t^2 \Lambda_{UV}^2$$

Fine tuned

SUSY

$$\delta m_H^2 \simeq -\frac{3y_t^2}{8\pi^2} m_{\tilde{t}}^2 \log \left(\frac{\Lambda_{UV}^2}{m_{\tilde{t}}^2} \right)$$

Fine tuned

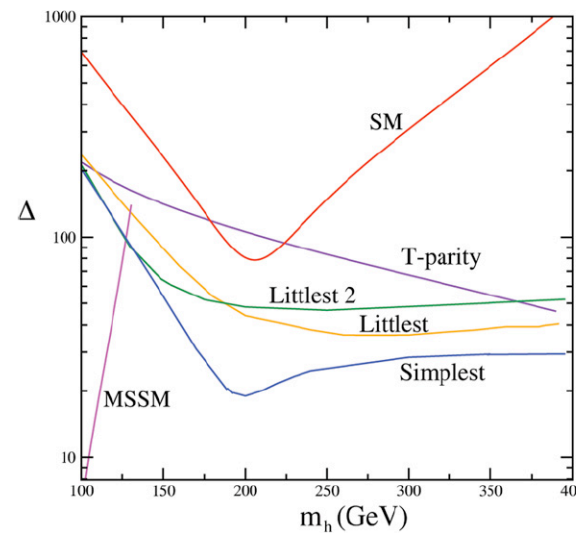
Little Higgs

$$\delta m_H^2 \propto \frac{y_t^2}{16\pi^2} f^2$$

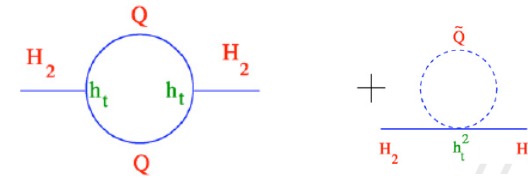
Fine tuned

(Pseudo Goldstone boson)

..but precision tests $\Rightarrow f = O(TeV)$



Heavy SUSY - Little Higgs



Little hierarchy problem:

Standard Model

$$\delta m_H^2 \propto y_t^2 \Lambda_{UV}^2$$

Fine tuned

SUSY

$$\delta m_H^2 \simeq -\frac{3y_t^2}{8\pi^2} m_{\tilde{t}}^2 \log \left(\frac{\Lambda_{UV}^2}{m_{\tilde{t}}^2} \right)$$

Fine tuned

Little Higgs

$$\delta m_H^2 \propto \frac{y_t^2}{16\pi^2} f^2$$

(Pseudo Goldstone boson)

Fine tuned

..but precision tests $\Rightarrow f = O(TeV)$

Double protection

$$\delta m_H^2 \approx -\frac{3y_t^2}{8\pi^2} [(m_{\tilde{t}}^2 + m_T^2) \log(m_{\tilde{t}}^2 + m_T^2) - m_{\tilde{t}}^2 \log(m_{\tilde{t}}^2) - m_T^2 \log(m_T^2)]$$

(SUSY+PG boson)

Vectorlike top quark

Falkowski et al; Csaki et al...

(Needed to cancel top quark contribution)

$$M_{SUSY} \rightarrow 10 TeV, \quad m_T \leq 1 TeV$$

New heavy quarks, new heavy gauge bosons $\sim TeV$

SUMMARY

LHC will probe new energy regime relevant to EW breaking.

Many possibilities identified - it will require extensive correlated information to distinguish between them. This will need:

- Control over SM backgrounds in a wide variety of (multiparticle) processes
- Higher order radiative corrections $(\sigma(gg \rightarrow H) \simeq \sigma_{LO}(1 + 0.7 + 0.3 + ..))$
- Develop techniques to measure spin and mass of new states

SUMMARY

LHC will probe new energy regime relevant to EW breaking.

Many possibilities identified - it will require extensive correlated information to distinguish between them. This will need:

- Control over SM backgrounds in a wide variety of (multiparticle) processes
- Higher order radiative corrections $(\sigma(gg \rightarrow H) \simeq \sigma_{LO}(1 + 0.7 + 0.3 + ..))$
- Develop techniques to measure spin and mass of new states

Much has been done but much still to do!