

# Parton densities: progress and challenges

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Andreas Vogt

University of Liverpool

- **Introduction: partons in hard processes, LHC Higgs production**
- **Higher-order effects in the parton evolution, small- $x$  logarithms**
- **Parton evolution in practice: evolution codes and benchmarks**
- **Input shapes and factorization schemes, heavy-quark densities**
- **Highlights of recent parton analyses, future constraints at LHC**
- **Instead of a summary:  $W$ -mass at the LHC**

DIS 2007, Munich, 16-04-07

# W.L. van Neerven, 1947 – 2007

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Photo courtesy of DESY

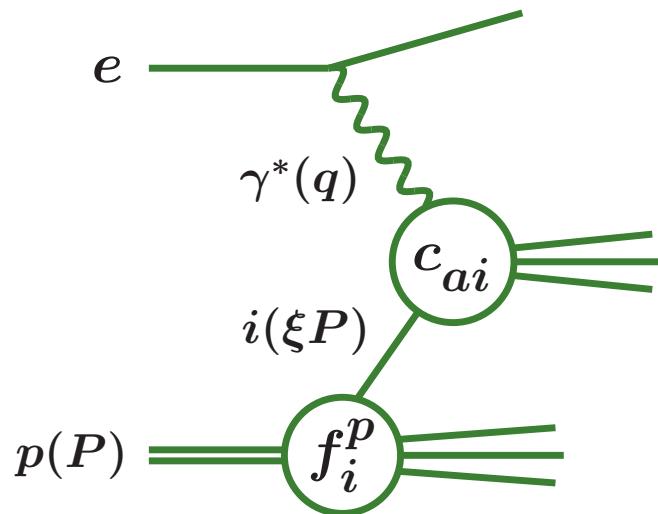
One of the greatest theorists  
of deep-inelastic processes

- ...
- Electro-weak corrections for LEP
- 2<sup>nd</sup> order QCD corr's for DY process
- NLO heavy-quark production in  $pp$
- $O(\alpha_s^2)$  coefficient functions in DIS
- Heavy quarks in structure functions
- Polarized NLO splitting functions
- 2<sup>nd</sup> order corr's for fragmentation
- Variable-flavour partons and DIS
- QCD corr's to LHC Higgs production

# Parton distributions and hard processes (I)

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Example: inclusive photon-exchange deep-inelastic scattering (DIS)



Hard scale, Bjorken variable

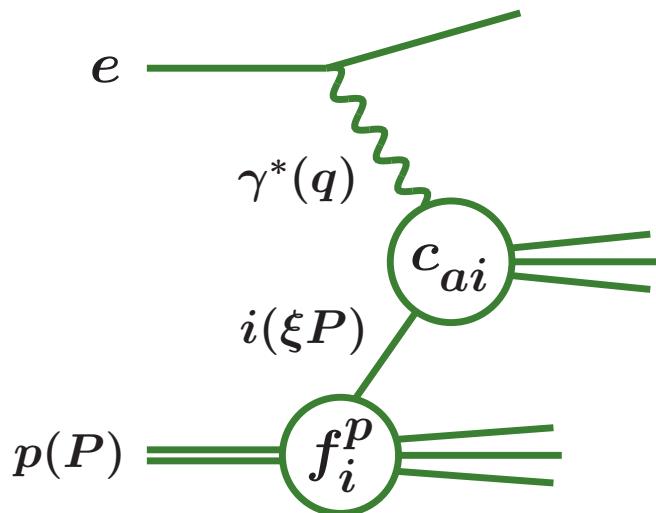
$$Q^2 = -q^2$$

$$x = Q^2/(2P \cdot q)$$

Lowest order, quarks:  $x = \xi$

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Lowest order, quarks:  $x = \xi$

Structure functions  $F_{2,L}$  [ up to  $1/Q^2$  terms – cut data to suppress ]

$$x^{-1} F_a^p(x, Q^2) = \sum_i \int_x^1 \frac{d\xi}{\xi} \, c_{a,i} \left( \frac{x}{\xi}, \alpha_s(\mu^2), \frac{\mu^2}{Q^2} \right) f_i^p(\xi, \mu^2)$$

Coefficient functions: scheme, scale  $\mu = \mathcal{O}(Q)$ , Mellin convolution

# Parton distributions and hard processes (II)

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Parton distributions  $f_i$ : evolution equations ( $\otimes$  = Mellin conv.)

$$\frac{d}{d \ln \mu^2} f_i(\xi, \mu^2) = \sum_k [P_{ik}(\alpha_s(\mu^2)) \otimes f_k(\mu^2)](\xi)$$

Initial conditions incalculable in pQCD. Lattice: low moments only

⇒ predictions: fits of suitable reference processes, universality

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Expansion in  $\alpha_s$ : splitting functions  $P$ , coefficient functions  $c_a$

$$P = \alpha_s P^{(0)} + \alpha_s^2 P^{(1)} + \alpha_s^3 P^{(2)} + \dots$$

$$c_a = \underbrace{\alpha_s^{n_a} \left[ c_a^{(0)} + \alpha_s c_a^{(1)} + \alpha_s^2 c_a^{(2)} + \dots \right]}_{}$$

LO: approximate shape, rough estimate of rate

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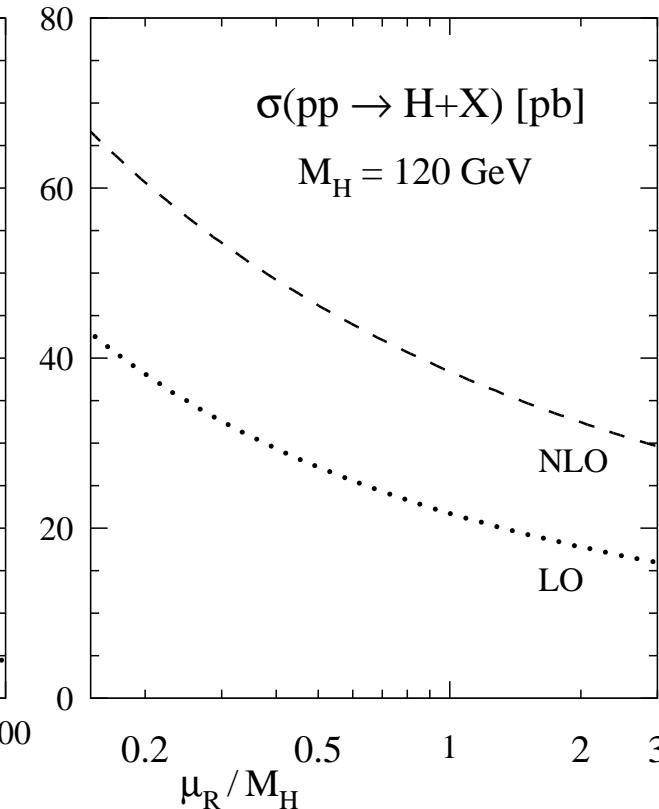
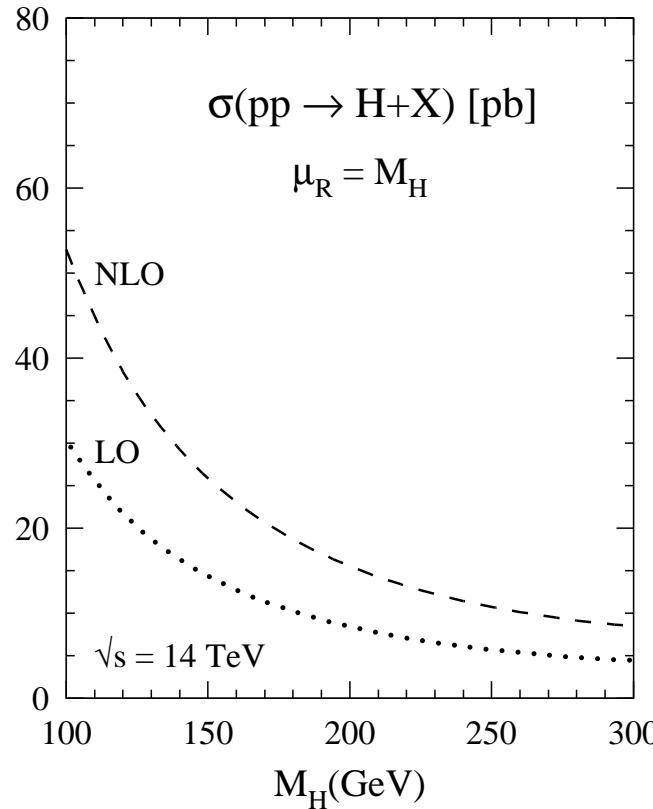
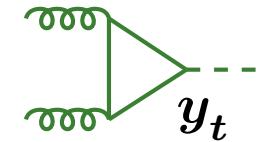
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NLO: first real prediction of size of cross sections

NNLO,  $P^{(2)}$ ,  $c_a^{(2)}$ : first serious error estimate of pQCD predictions

# Higgs boson production at the LHC (I)

Dominant channel:  $gg \rightarrow H + X$  via top-quark loop



Spira et al. (95)

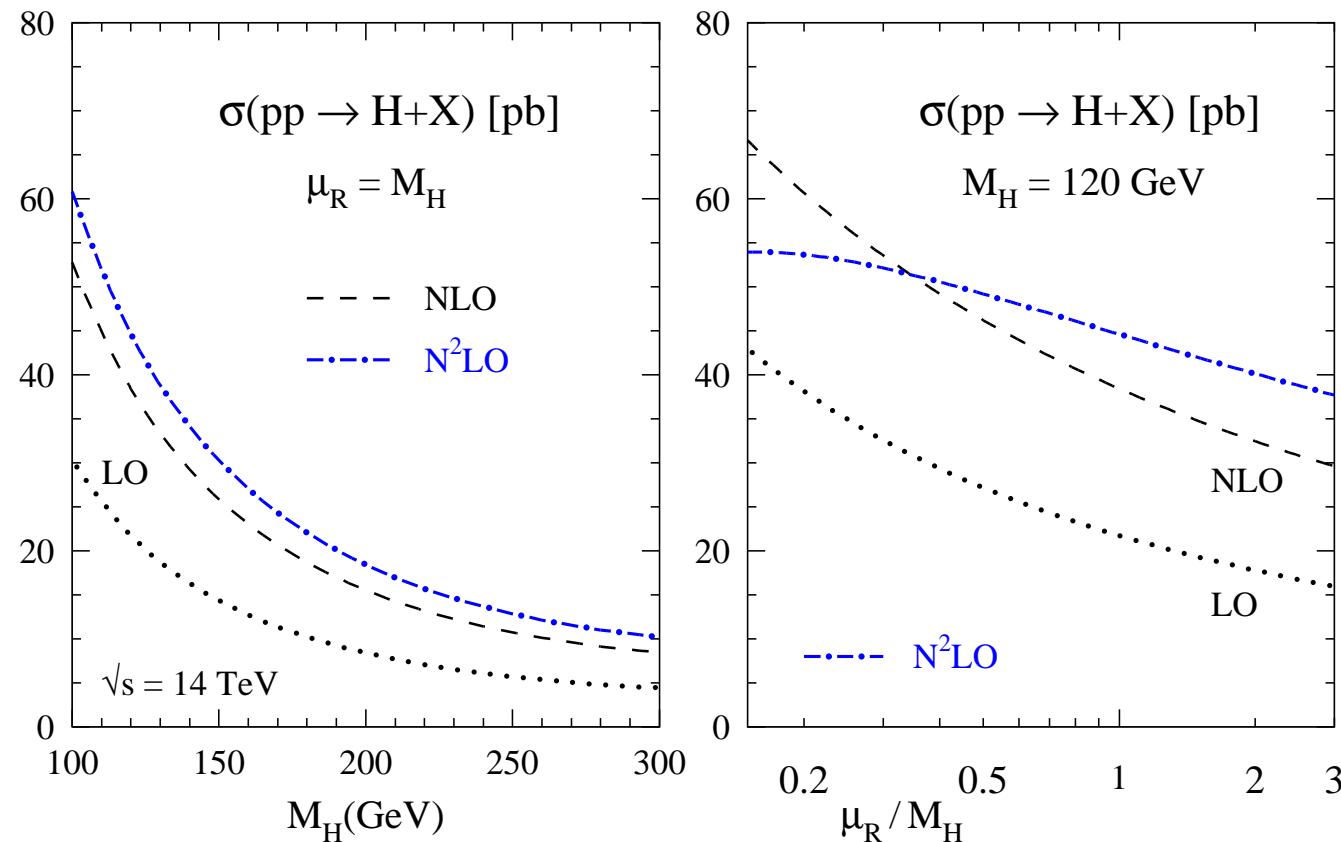
Error guess: apparent convergence, variation of scale  $\mu$  (how far?)

Next-to-leading order (NLO) insufficient for quantitative prediction

# Higgs boson production at the LHC (II)

$\hat{\sigma}_{\text{NNLO}}$  : Harlander, Kilgore (02); Anastasiou, Melnikov (02, 05 [ $\sigma_{\text{diff}}$  ])

Partons including NNLO: MRST (plot), Alekhin.

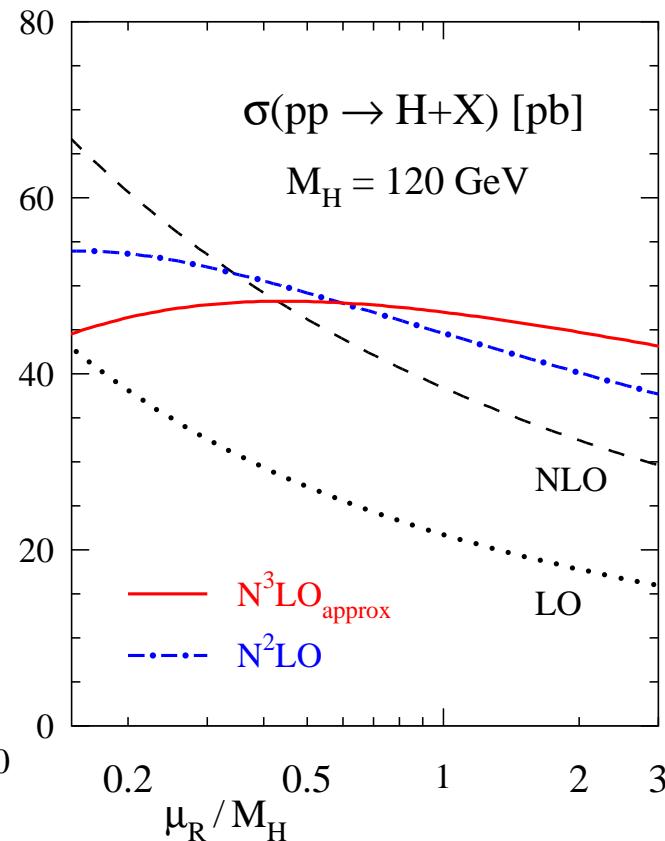
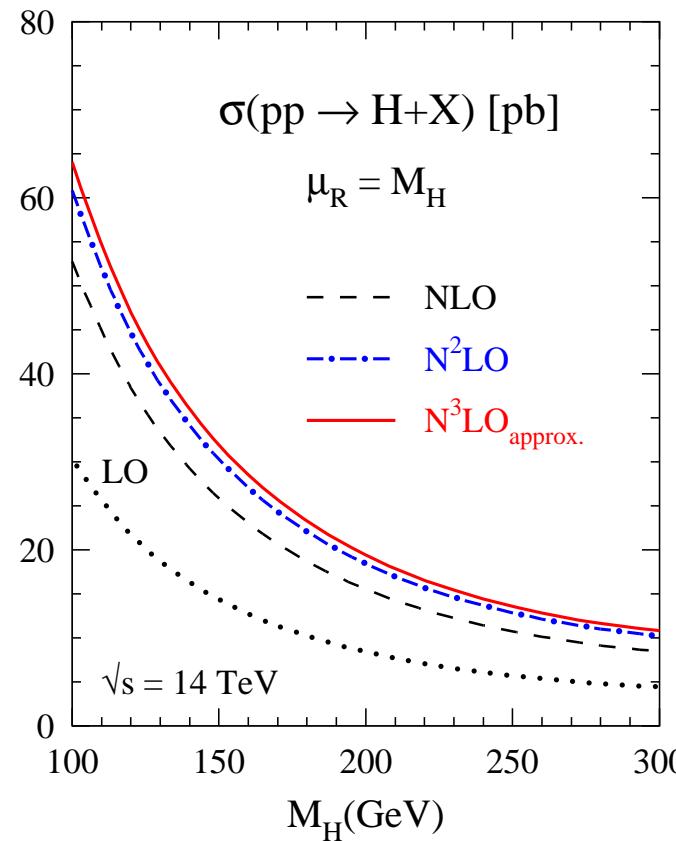


Higher-order uncertainties:  $\sim 15\%$  at NNLO

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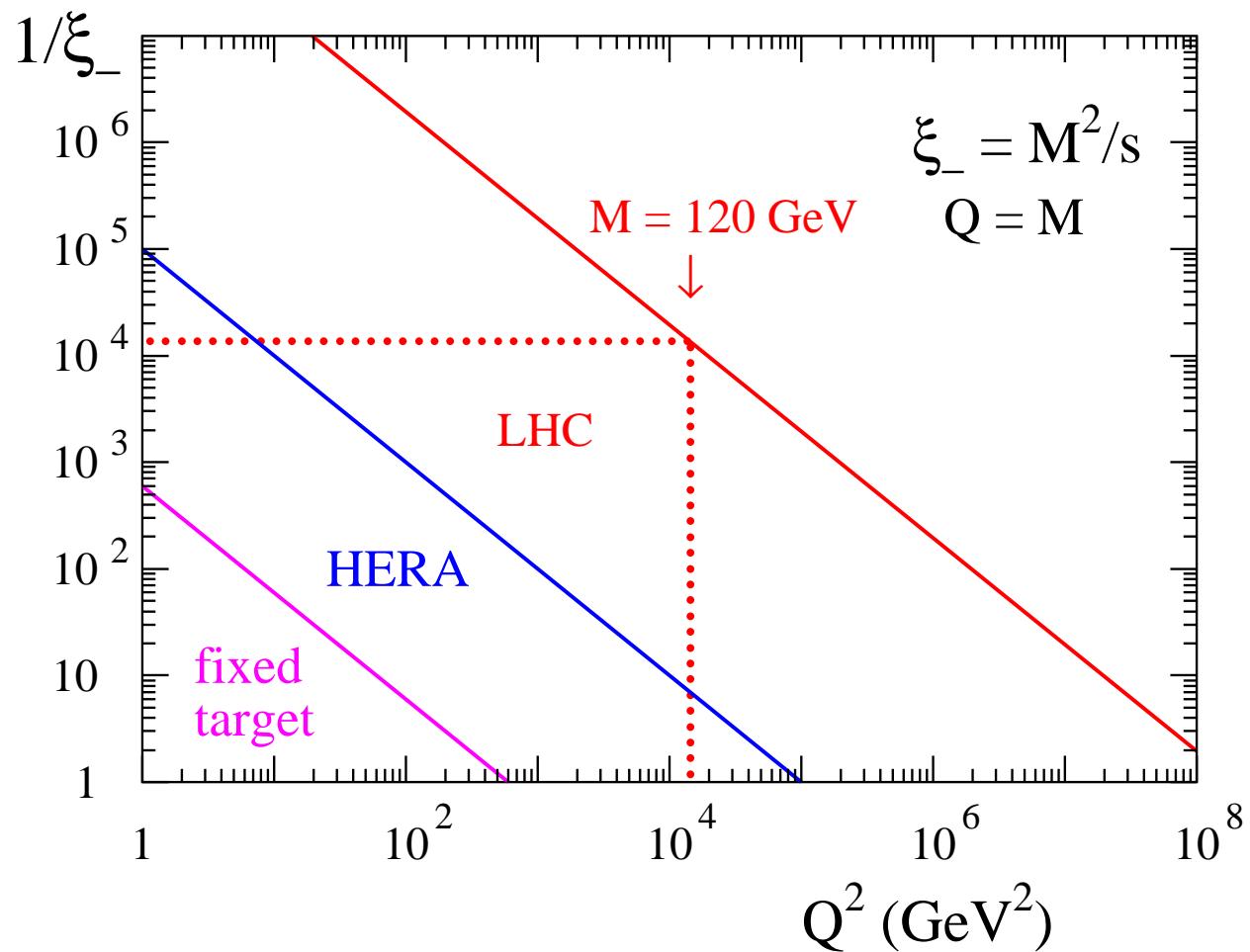
Moch, A.V. (05)

Higher-order uncertainties:  $\sim 15\%$  at NNLO,  $\sim 5\%$  at approx.  $N^3\text{LO}$

# Parton evolution from HERA to LHC

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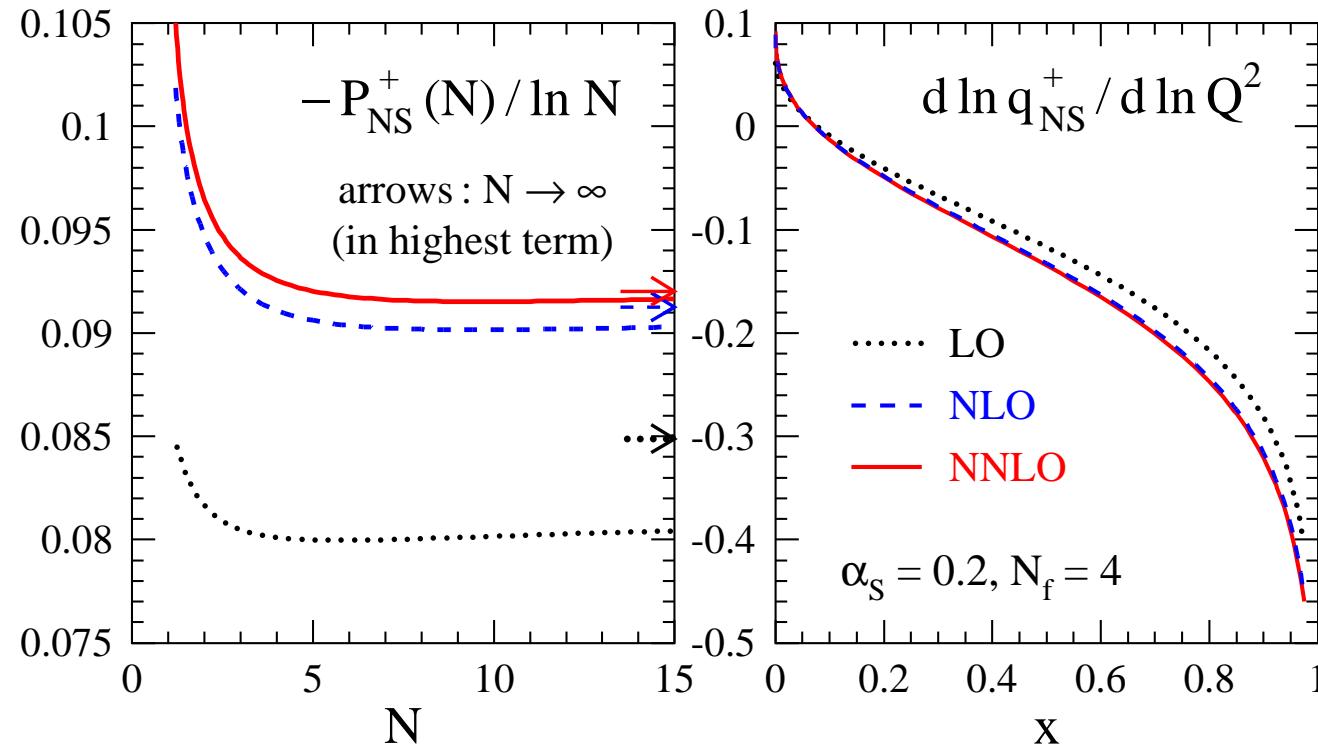
**Kinematics: partons with momentum fractions  $\xi_- < \xi < 1$  contribute**



**$W/Z, H, \text{top, new physics: } \xi_- \gtrsim 10^{-4}, \text{ can cut at } Q^2 \approx 10 \text{ GeV}^2$**

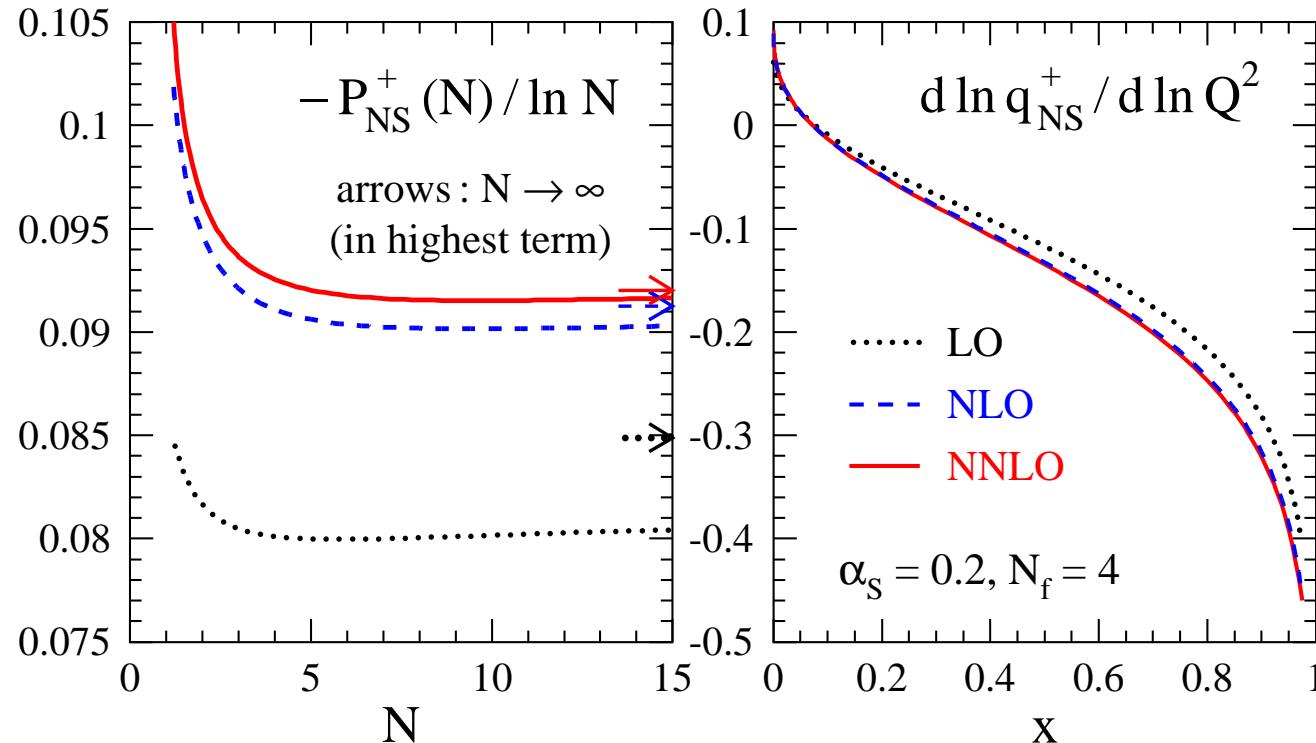
# Parton evolution at large $N$ / large $x$ in $\overline{\text{MS}}$

Recall  $A^N = \int_0^1 dx x^{N-1} A(x)$ . Non-singlet<sup>+</sup>:  $u + \bar{u} - (d + \bar{d})$  etc



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**$N^3\text{LO}$ :**  $P_{\text{ns}}^+$  now known for  $N=2$ ,  $n_f = 3$       **Baikov, Chetyrkin (06)**

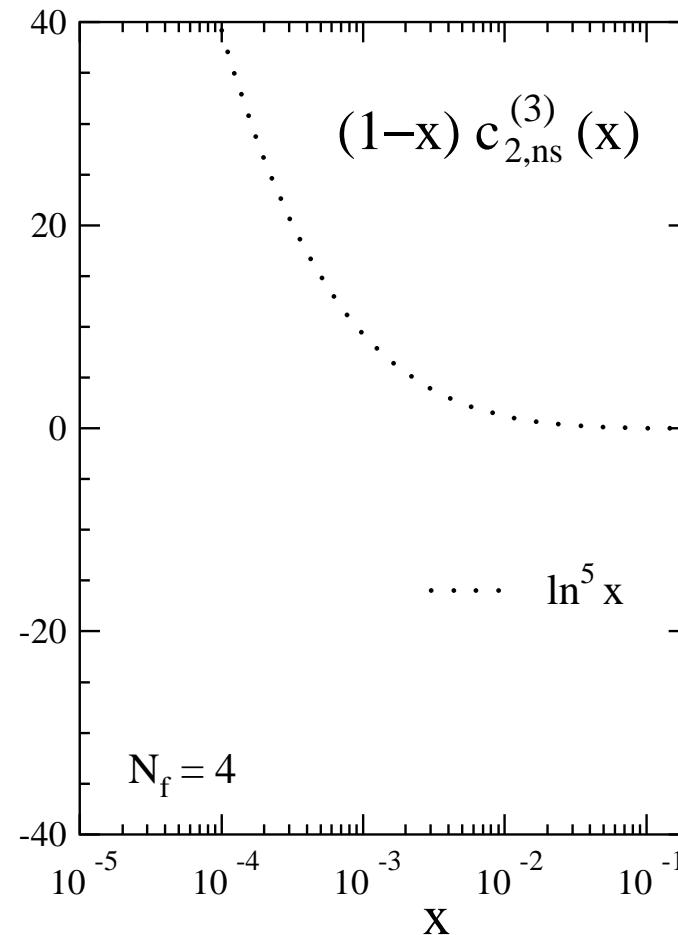
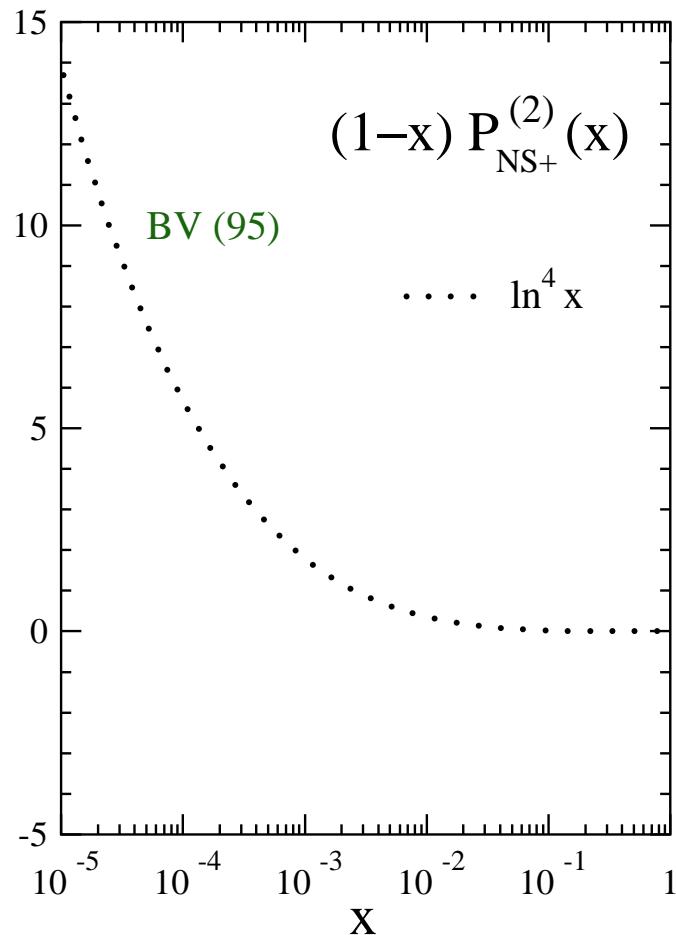
$$P_{\text{ns}}^+ = -0.283 \alpha_s [1 + 0.869 \alpha_s + 0.798 \alpha_s^2 + 0.926 \alpha_s^3 + \dots]$$

$N > 2$ : similar/smaller (as  $\ln N$  coeff's).  $n_f > 3$ : HO corr's smaller

# Non-singlet three-loop quantities at small $x$

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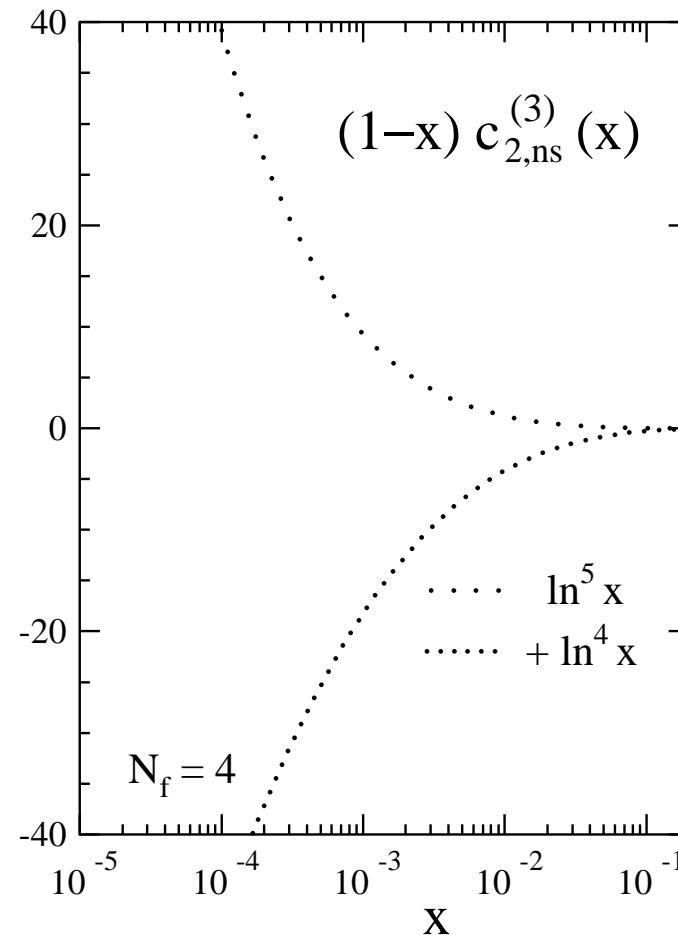
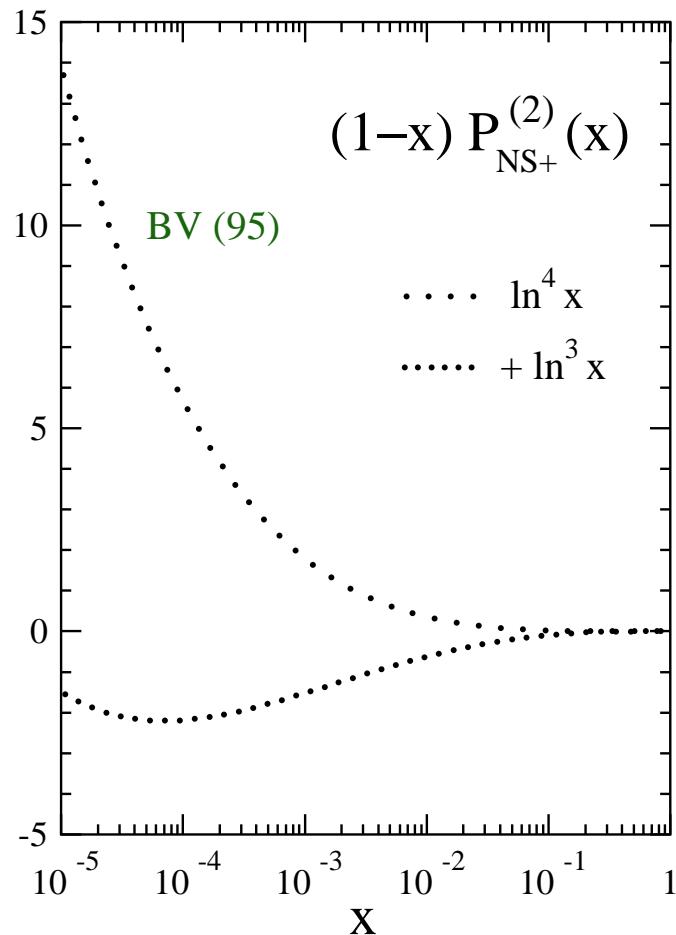
Order  $\alpha_s^n$ : small- $x$  terms  $\ln^k x$  with  $k$  up to  $2n-2$  ( $2n-1$ ) in  $P(c)$



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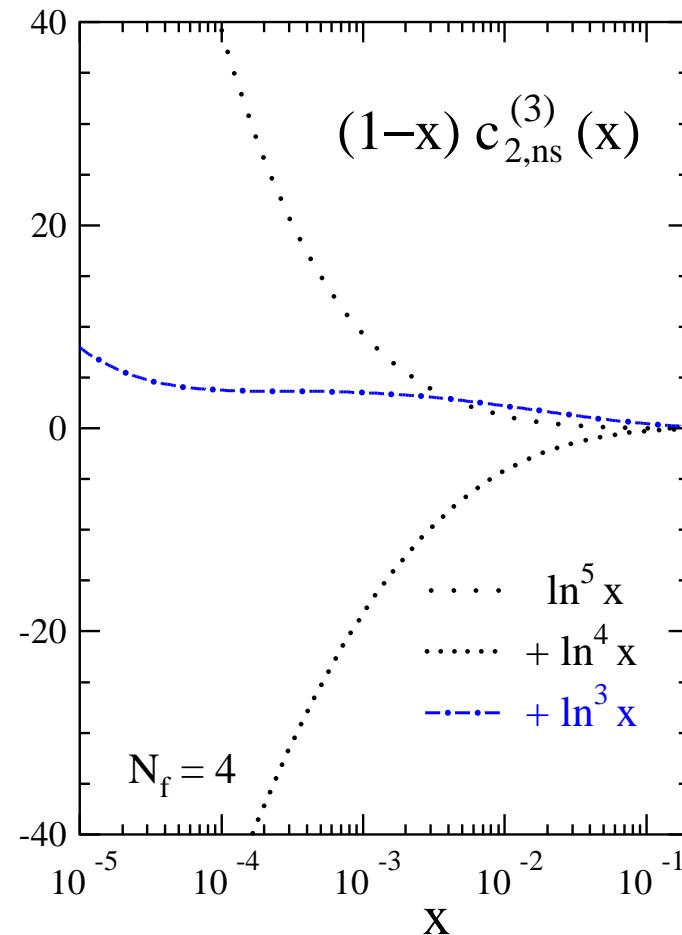
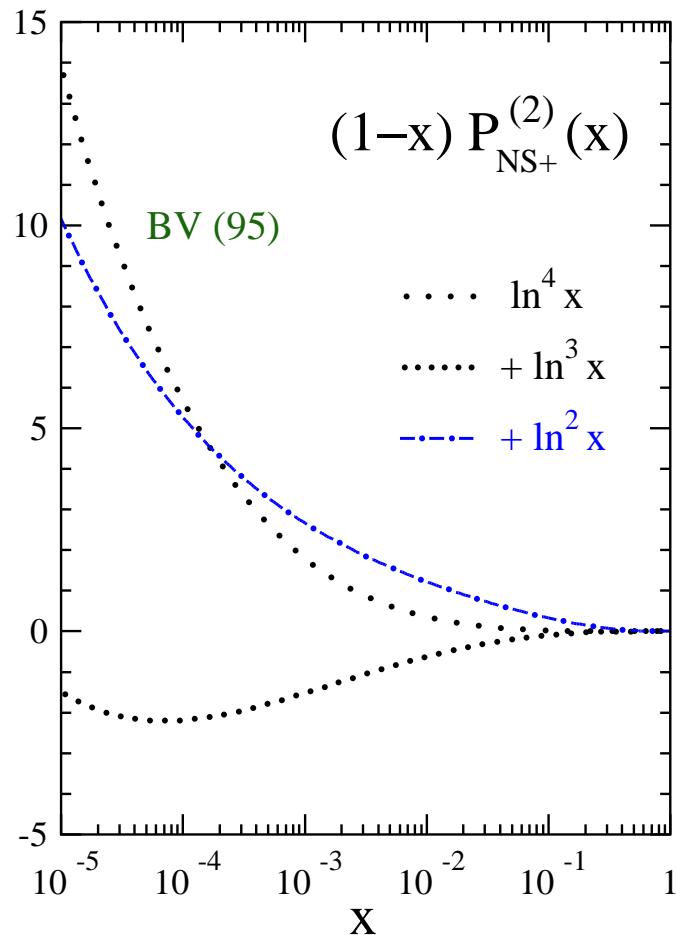
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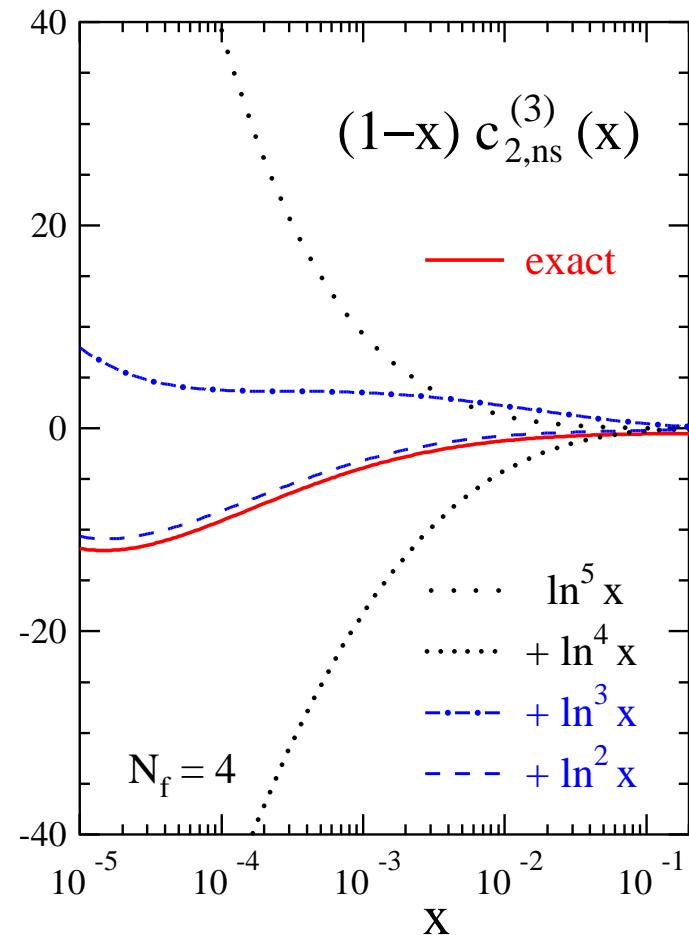
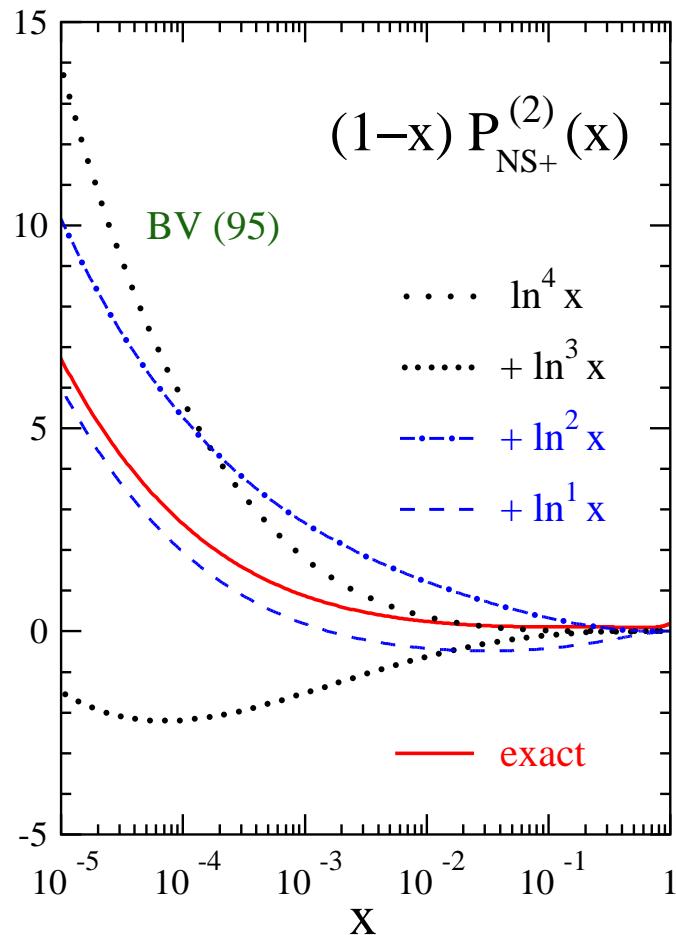
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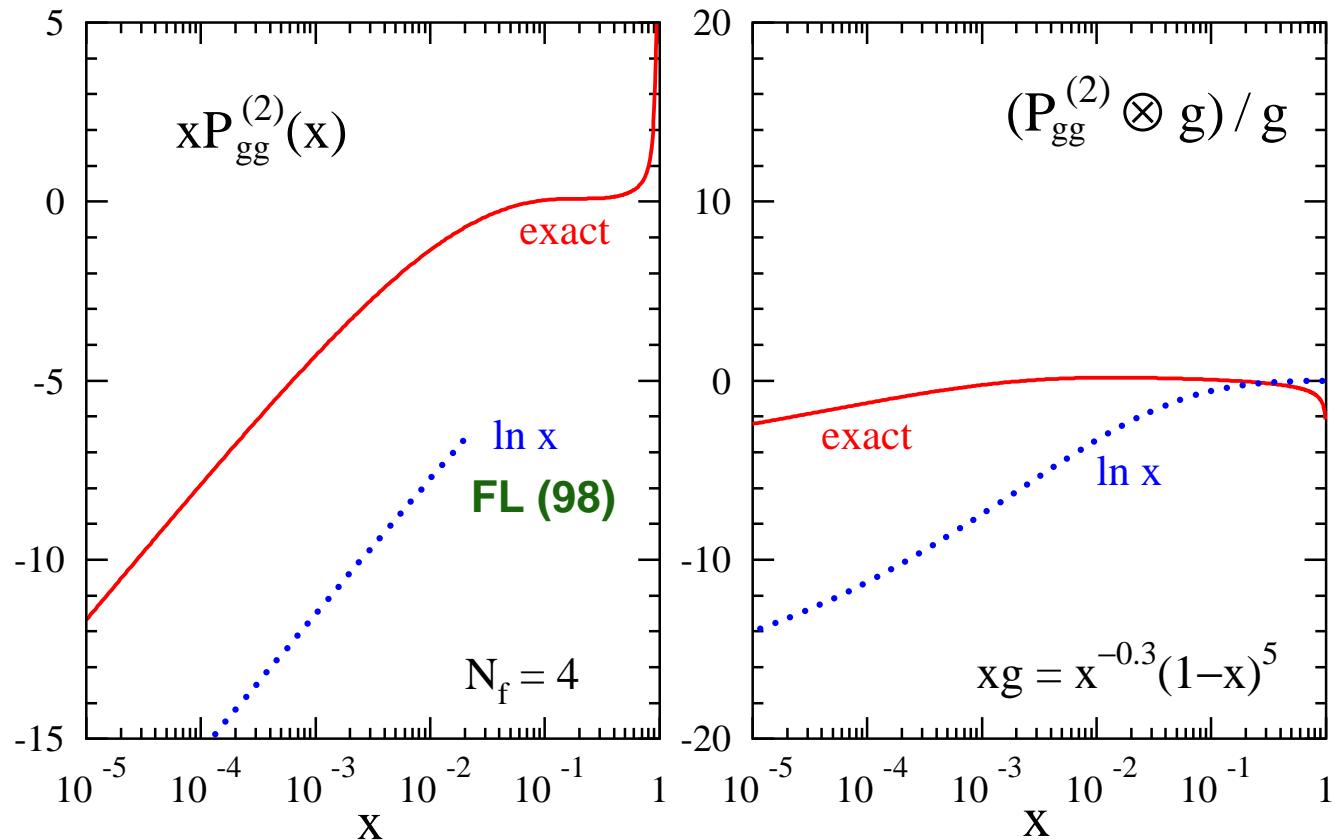


$x$ -values for colliders: not even shape guaranteed by leading logs

# Singlet splitting and evolution at small $x$

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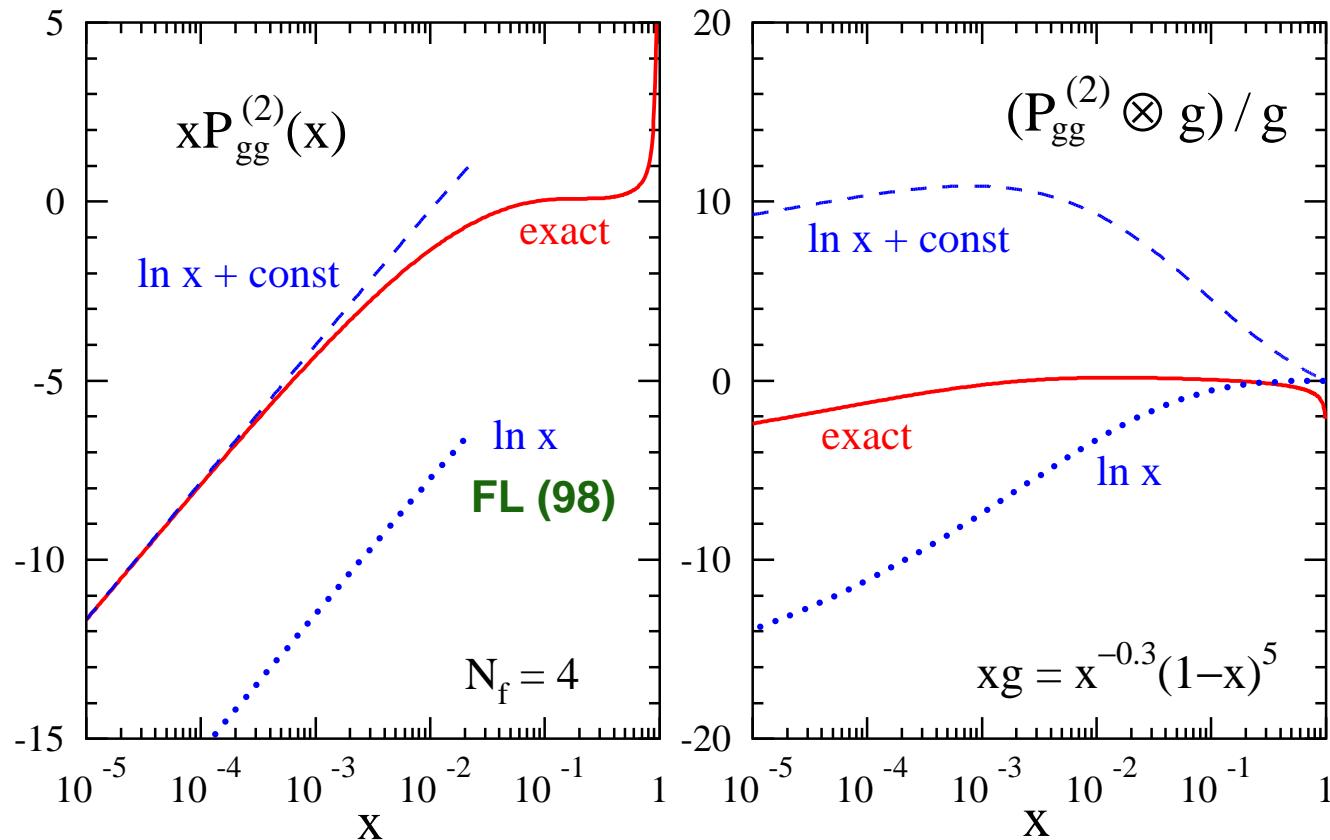
Splitting functions  $\rightarrow$  observables: convolutions,  $\int_x^1 \frac{dy}{y} P(y) f\left(\frac{x}{y}\right)$



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**Splitting functions → observables:** convolutions,  $\int_x^1 \frac{dy}{y} P(y) f\left(\frac{x}{y}\right)$

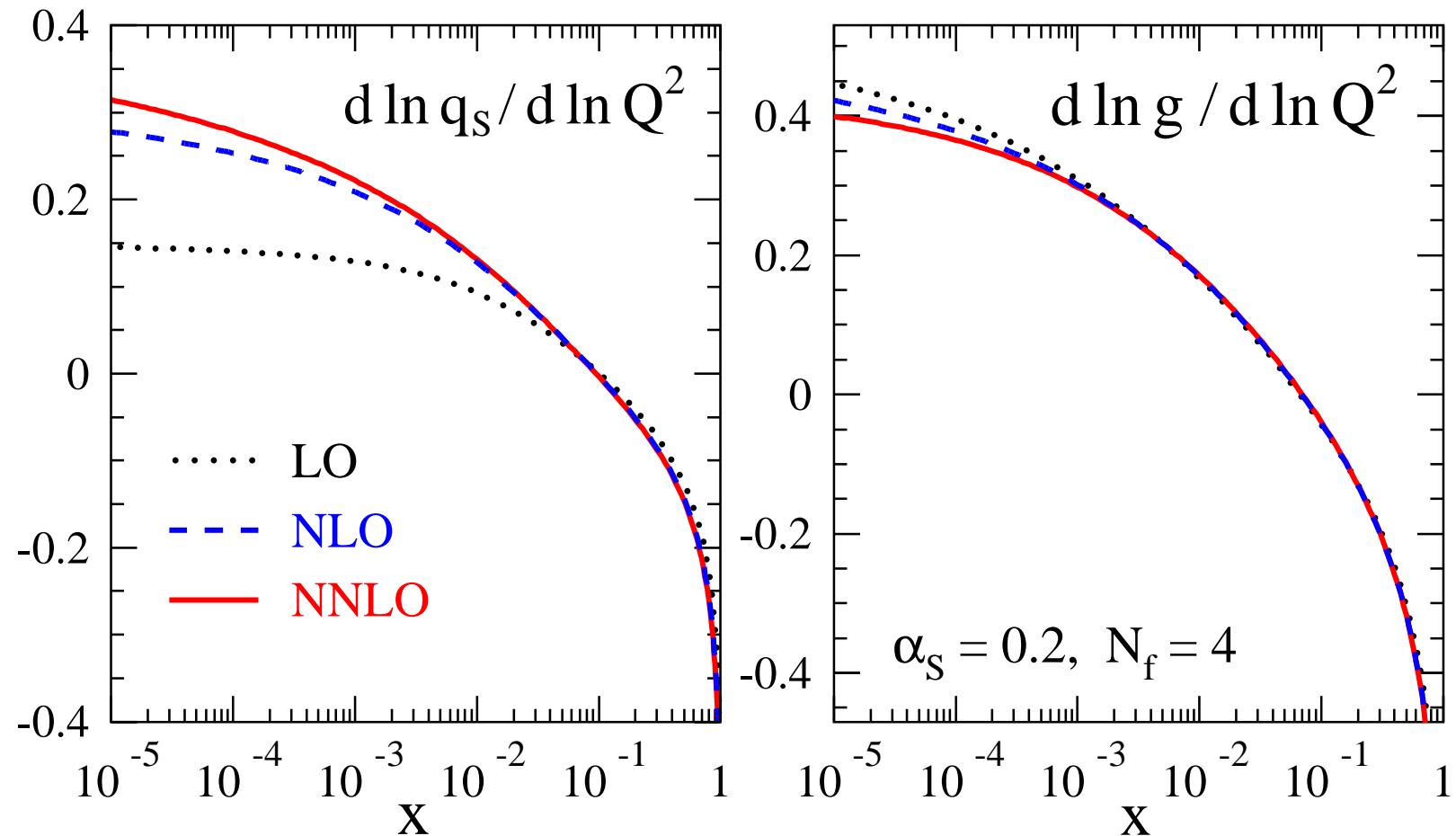


**General:** small- $x$  limits of splitting fct's insufficient in convolutions

**First estimate of corr's beyond NNLO:** future moments at order  $\alpha_s^4$

# Evolution of singlet parton distributions

Example: quark and gluon scale derivatives at scale  $Q^2 \approx 30 \text{ GeV}^2$



Expansion very stable for main LHC momentum fractions  $x \gtrsim 10^{-4}$

# Available evolution codes including NNLO

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**$x$ -space: discretization in  $x, \mu_f$  of coupled integro-differential eqs.**

HOPPET (G. Salam, from 2001), <http://hepforge.cedar.ac.uk/hoppet/>

QCDNUM (M. Botje, new: v. 17), <http://www.nikhef.nl/~h24/qcdnum/>

**$N$ -space: ordinary diff. eqs., time-ordered exp.,  $N \rightarrow x$  numerical**

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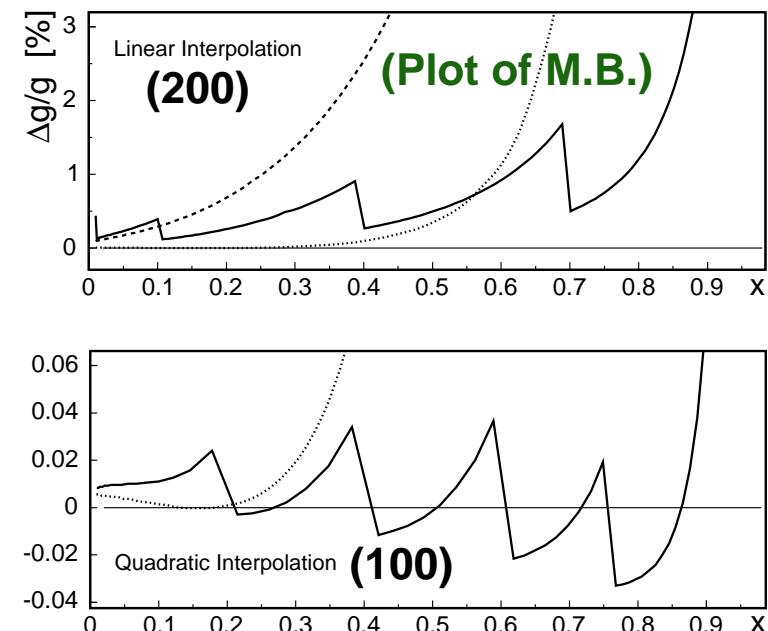
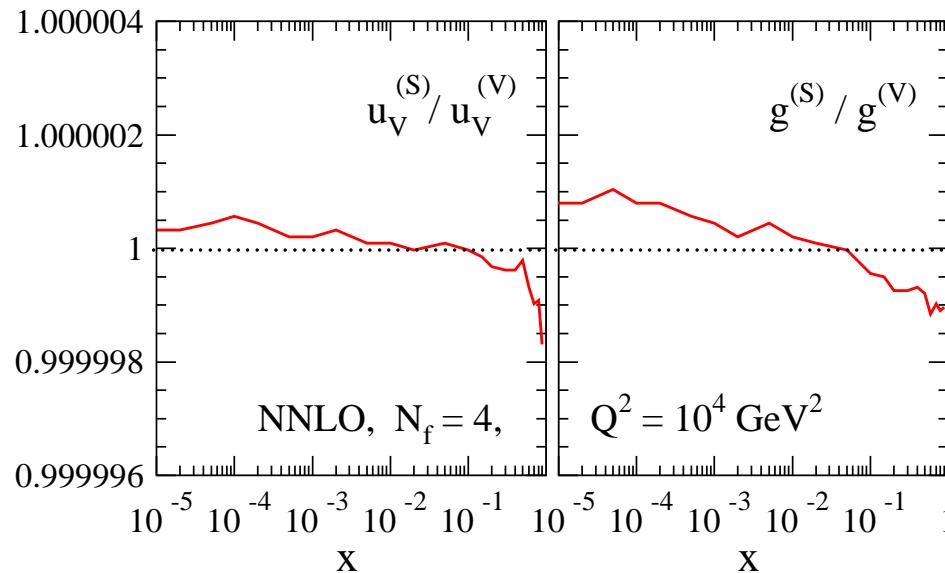
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## Sample comparisons



# Benchmark tables for the parton evolution

---

Evolution of Les Houches (2001) reference input at  $\mu_{f,0}^2 = 2 \text{ GeV}^2$

$$xu_v(x, \mu_{f,0}^2) = 5.1072 x^{0.8} (1-x)^3 , \dots$$
$$xg(x, \mu_{f,0}^2) = 1.7000 x^{-0.1} (1-x)^5$$

with

$$\alpha_s(\mu_r^2 = 2 \text{ GeV}^2) = 0.35$$

at LO, NLO, NNLO, for  $\mu_r = \{0.5, 1, 2\} \mu_f$ , with fixed / variable  $N_f$

Use of two completely different codes.

G. Salam, A.V. (2002, 05)

Five-digit agreement over wide range in  $x, \mu_f^2 \Rightarrow$  reference tables

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Example: (iterated) NNLO,  $\mu_r = 2\mu_f$ ,  $N_f = 4$  at  $\mu_f^2 = 10^4 \text{ GeV}^2$

$$x = 10^{-5}, \quad xu_v = 2.9032 \cdot 10^{-3}, \quad \dots, \quad xg = 2.2307 \cdot 10^2$$

...

$$x = 0.9, \quad xu_v = 3.6527 \cdot 10^{-4}, \quad \dots, \quad xg = 1.2489 \cdot 10^{-6}$$

# Input shapes and factorization schemes

---

**MS evolution:** convenient, stable (see above) – but pdf's not physical

Shapes for  $x \rightarrow 1$ , e.g., more natural in other scheme? Positivity?

NLO partons (which scheme?) best choice with LO Monte-Carlos?

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Traditional alternative: **DIS scheme, quarks physical via  $F_2$ .** Singlet:

$$q_S^{\text{DIS}} = q_S + \alpha_s [ c_{2,q}^{(1)} \otimes q_S + c_{2,g}^{(1)} \otimes g ] + \dots$$

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Interesting old concept: **DIS<sub>φ</sub> scheme**      Furmanski, Petronzio (81)

Also gluon shape physical via structure fct.  $F_\phi$  of scalar coupling to  $G^{\mu\nu}G_{\mu\nu}$  ( $\leftrightarrow$  Higgs in large- $m_t$  limit) –  $c_{\phi,i}^{(1,2,3)}(x)$  now computed

⇒ End-point constraints from positivity of  $F_\phi$ , ...

# Heavy quarks in parton evolution

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Below: disregard ‘intrinsic charm’ – can be relevant at large  $x$

Pumplin, Lai, Tung (2007)

**$\overline{\text{MS}}$  evolution of partons and  $\alpha_s$  with variable number of flavours:  
matching of effective theories. For partons at  $\mu_F = m_h$  (pole mass):**

$$\begin{aligned} l_i^{(N_f+1)} &= l_i^{(N_f)} + \delta_{m2} a_s^2 A_{qq,h}^{\text{NS},(2)} \otimes l_i^{(N_f)} \\ g^{(N_f+1)} &= g^{(N_f)} + \delta_{m2} a_s^2 [A_{gq,h}^{\text{S},(2)} \otimes q_S^{(N_f)} + A_{gg,h}^{\text{S},(2)} \otimes g^{(N_f)}] \\ (h + \bar{h})^{(N_f+1)} &= \delta_{m2} a_s^2 [A_{hq}^{\text{S},(2)} \otimes q_S^{(N_f)} + A_{hg}^{\text{S},(2)} \otimes g^{(N_f)}] \end{aligned}$$

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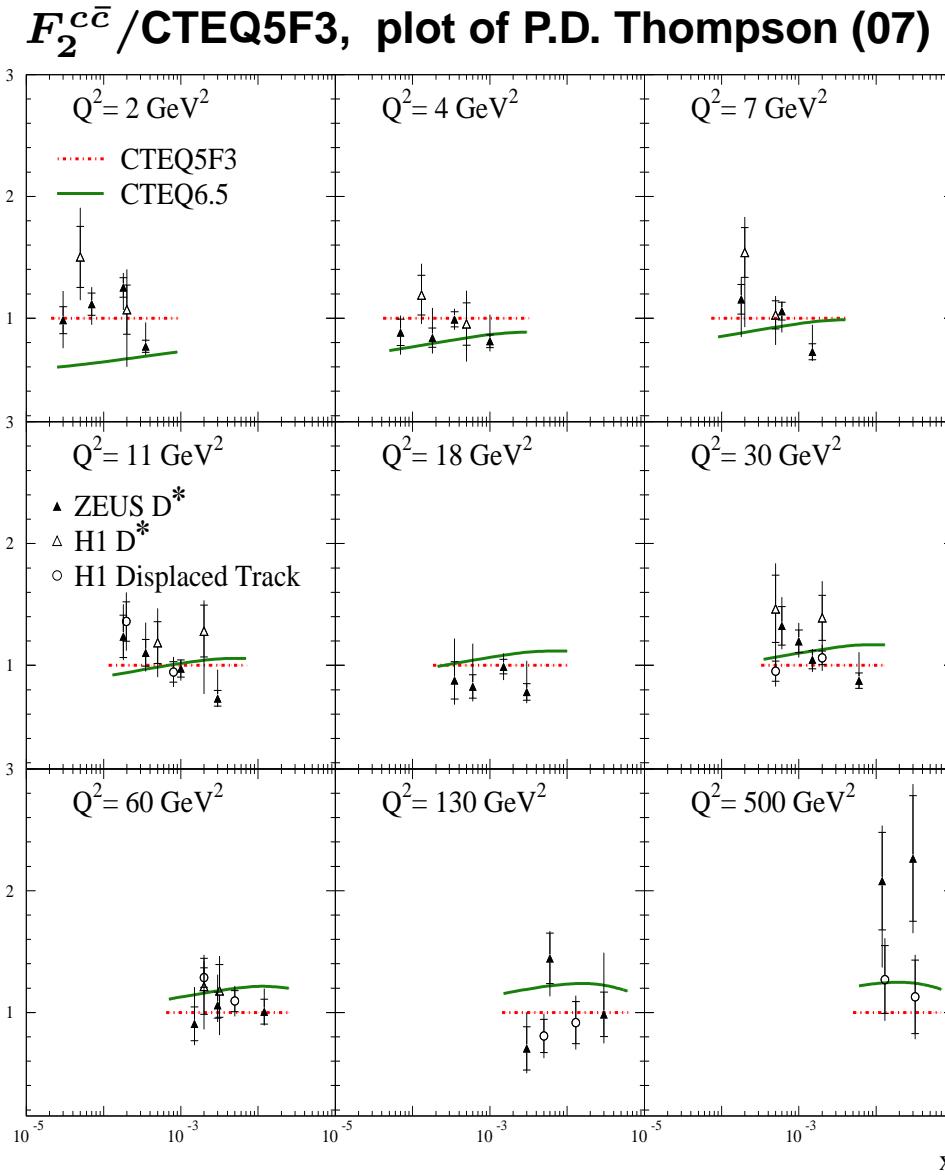
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Included in the evolution codes and benchmarks discussed above

Ignored in NNLO MRST partons, now implemented for MSTW (07) sets  
+ more improvements, e.g., fastNLO (Kluge, Rabberts, Wobisch) instead of  $K$ -factors

# Heavy quarks in structure functions

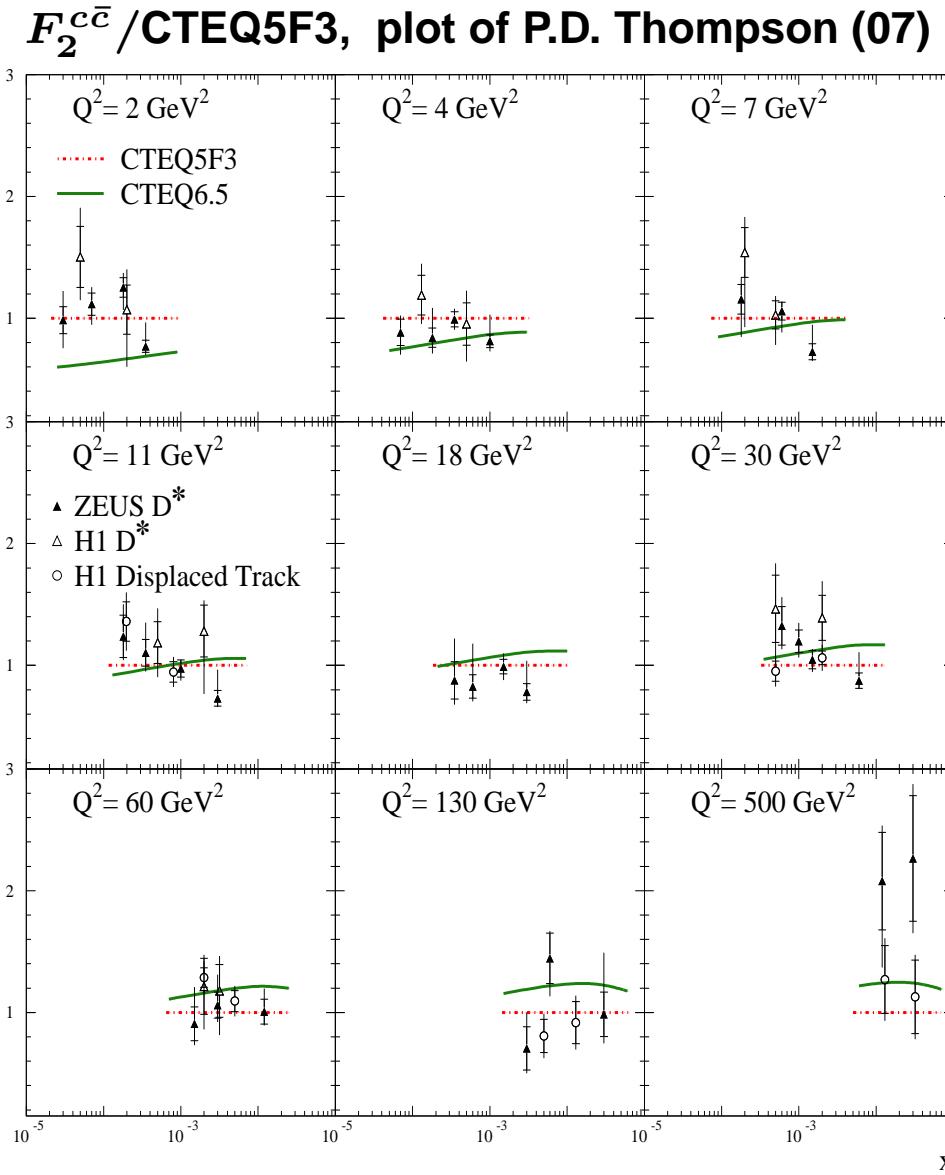


$Q \gg m_c : u, d, s, g$  partons  
+ massive  $c$  coeff. fct's, FFNS

NLO: Laenen et al. (92); HVQDIS (95)

$Q \ggg m_c$ : terms  $m_c/Q \rightarrow 0$   
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$Q \gg m_c$ : terms  $m_c/Q \neq 0$ ,  
but quasi-collinear logs large  
 $n_f = 4$  pdf's, 'interpolating'  
coeff. functions, (GM-) VFNS  
ACOT; BMSN/CSN; (Roberts-) Thorne

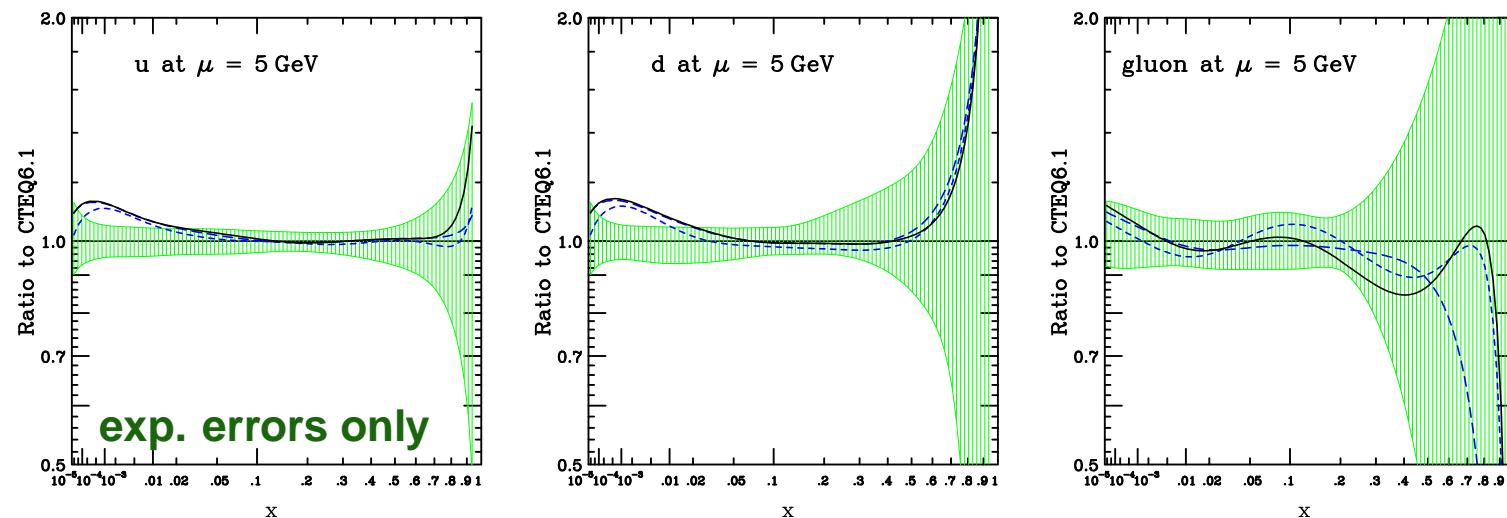
Transition process-dependent  
Exp.+ th.: " $Q \gg m_c$ " for HERA  $F_2^{c\bar{c}}$

# Highlights of recent parton analyses (I)

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CTEQ6.5, ..., W.K. Tung et al. (06)

Charm mass suppression in DIS finally implemented in fits  
Improved iterative treatment of inconsistent data under way, ...

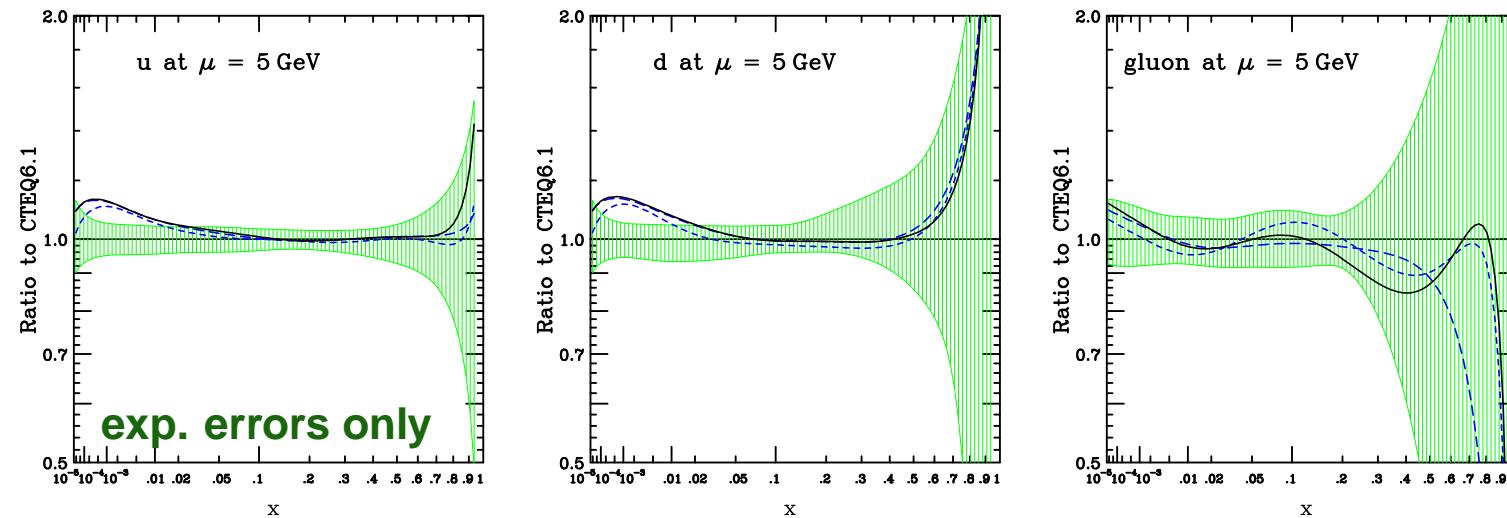


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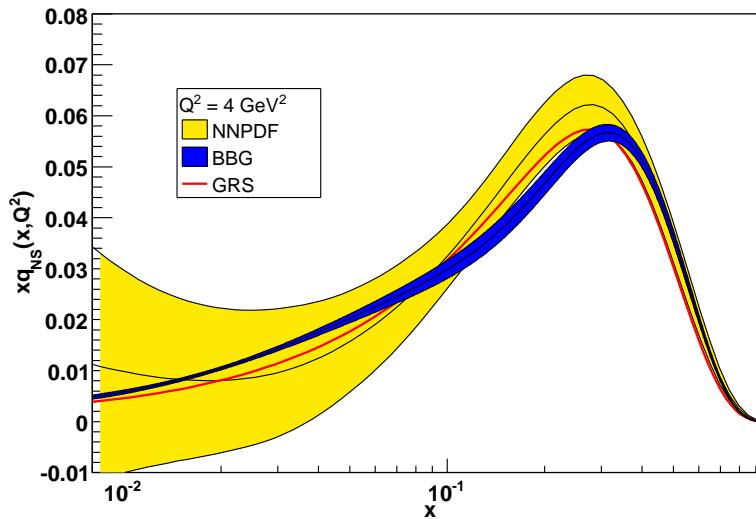
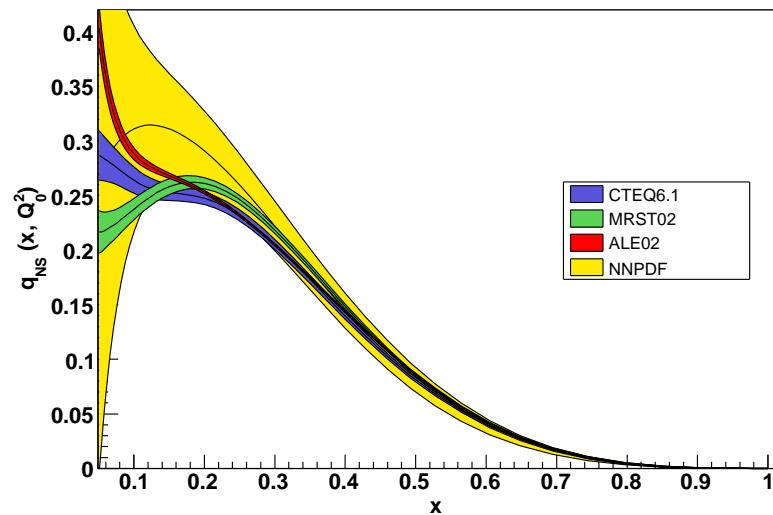
Extension of fits (S.A., 02: DIS only) to consistent subset of DY data  
NNLO corr's to Drell-Yan cross sections crucial; preferred  $\alpha_s$  low

# Highlights of recent parton analyses (II)

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NNPDF, L. Del Debbio et al (07)

Hybrid evolution combining advantages of  $N$ - and  $x$ -space  
Non-singlet (so far); neural networks to avoid parametrization bias

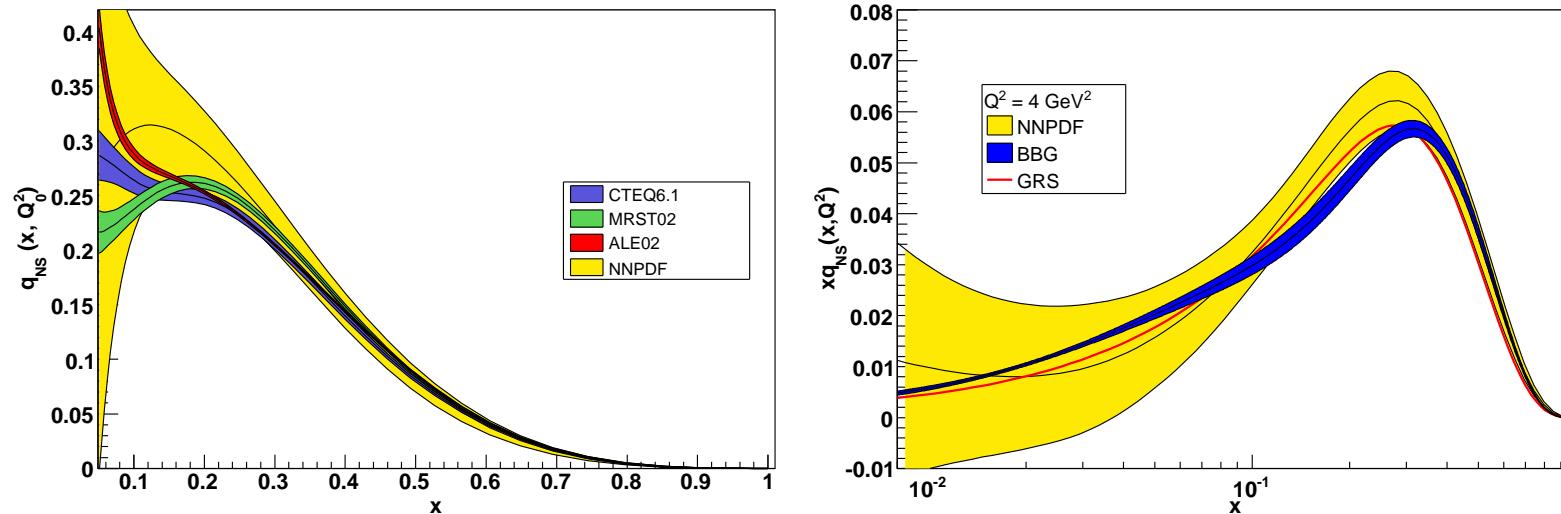


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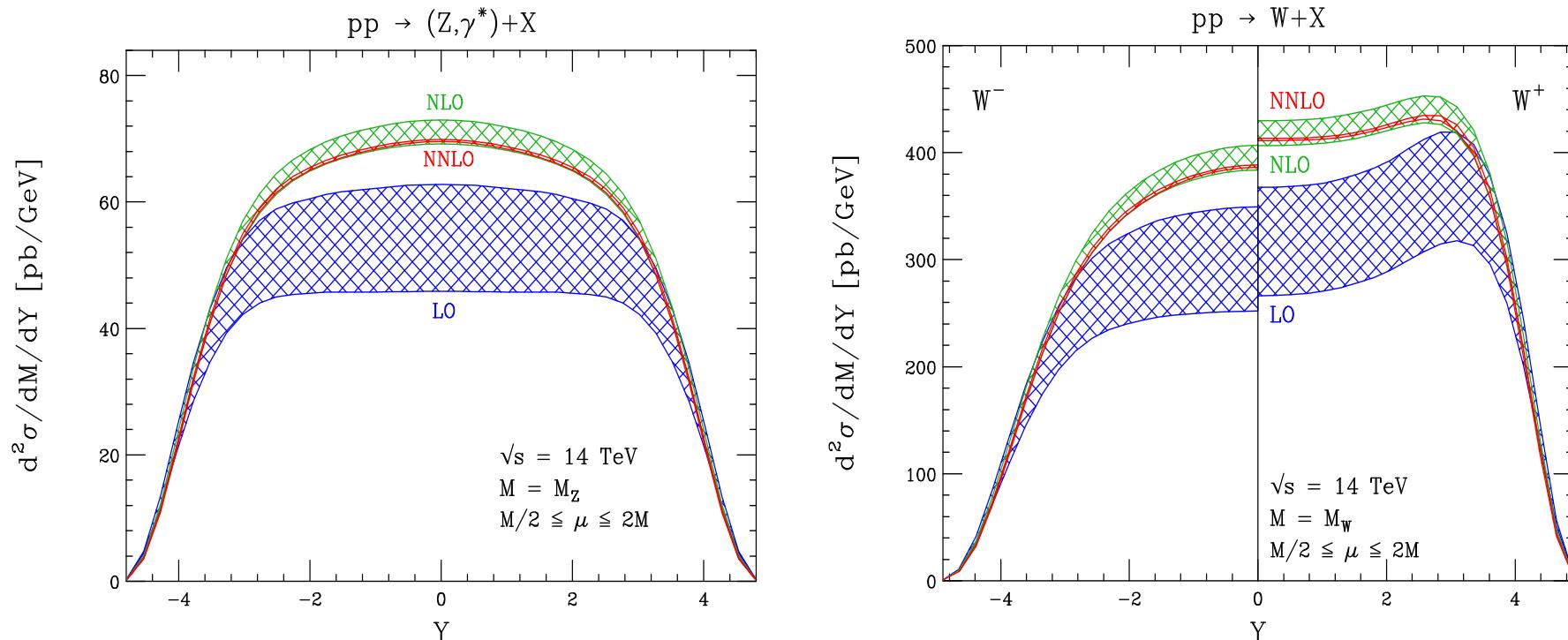
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Blümlein, Böttcher, Guffanti (06)

Global non-singlet DIS analysis with determination of  $\alpha_s$  to N<sup>3</sup>LO

$\alpha_s(M_Z) = 0.1134,41 \pm 0.002$  at N<sup>2,3</sup>LO, consistent with AMP (06)

# Constraints of LHC gauge-boson production



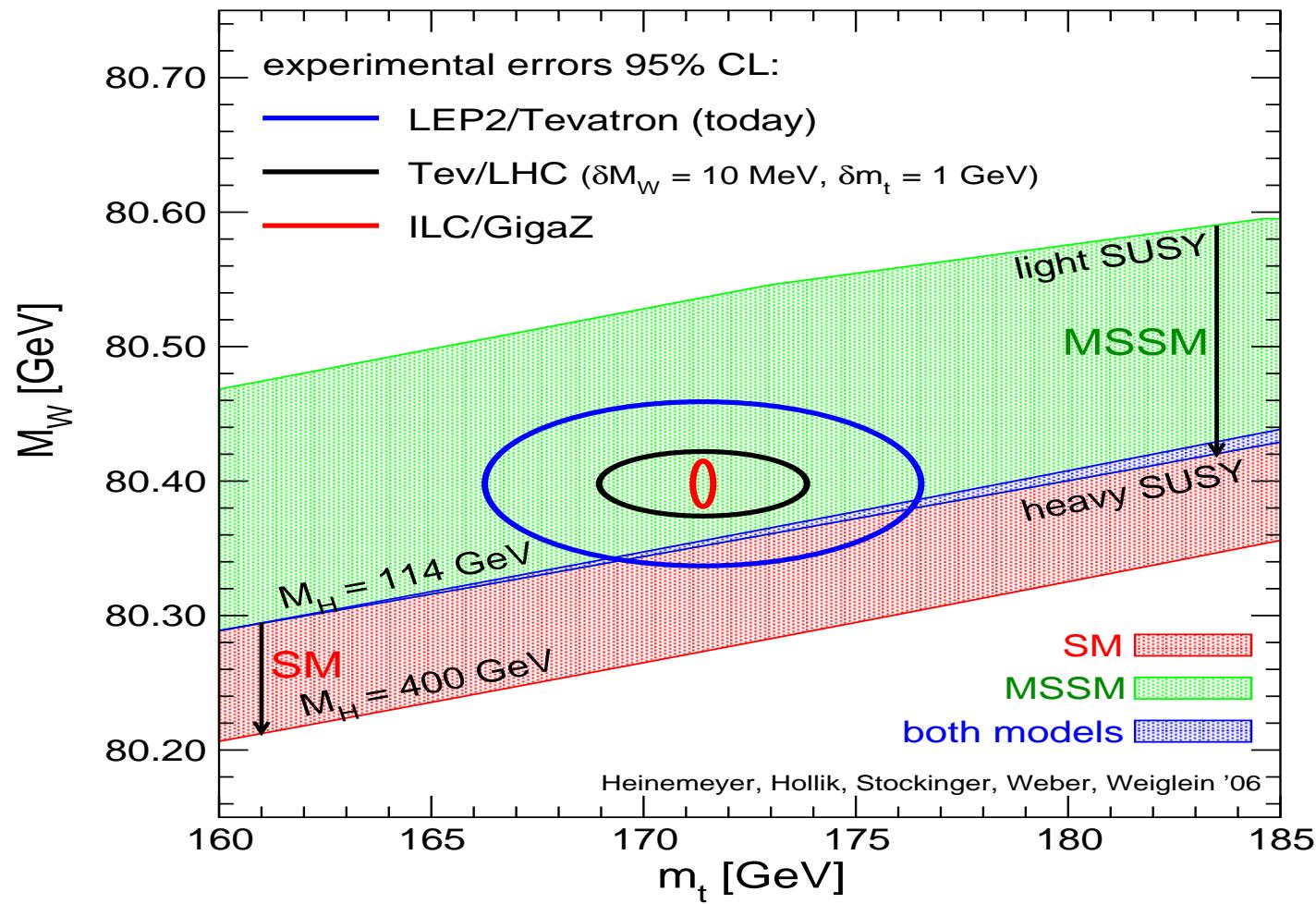
diff.  $\hat{\sigma}_{\text{NNLO}}$ : Anastasiou, Dixon, Melnikov, Petriello (03)

'Gold-plated' processes: NNLO perturbative accuracy better than 1%

⇒ use to improve upon pre-LHC determinations of parton densities  
A. Cooper-Sarkar, M. Klein, ... (this workshop)

# Instead of a summary: $W$ -mass at the LHC

$M_W$  as function of  $\tau_\mu$ ,  $M_Z$ , ... can discriminate between theories



Precision QCD and partons required to make the black ellipse happen