# SINGLE TOP STUDIES WITH MCFM

### F. Tramontano

Università di Napoli "Federico II" and INFN sezione di Napoli

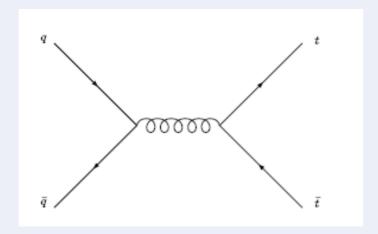
in collaboration with J. Campbell and R.K. Ellis

Munich

**DIS 2007** 

### Introduction

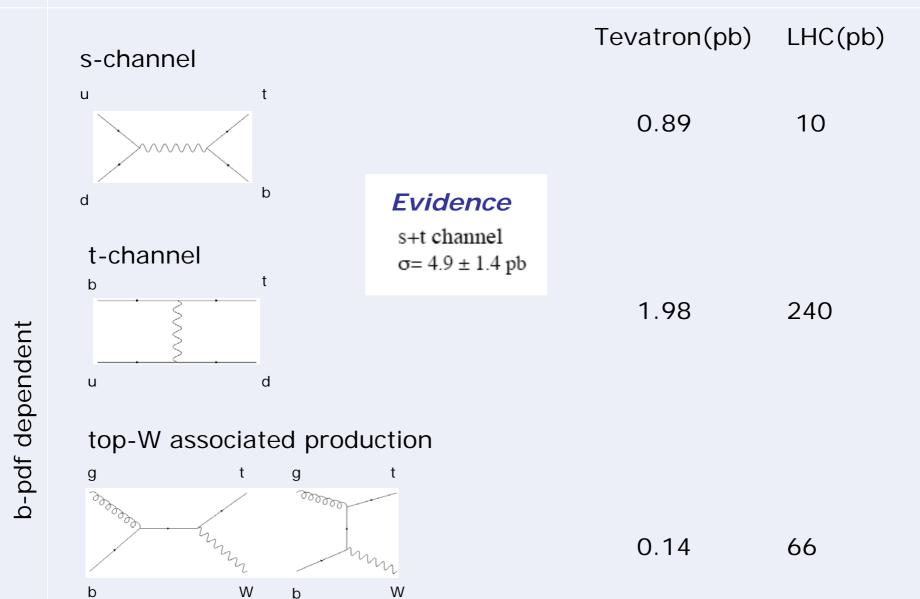
☐ The top quark has been discovered at Tevatron through the strong process of pair creation



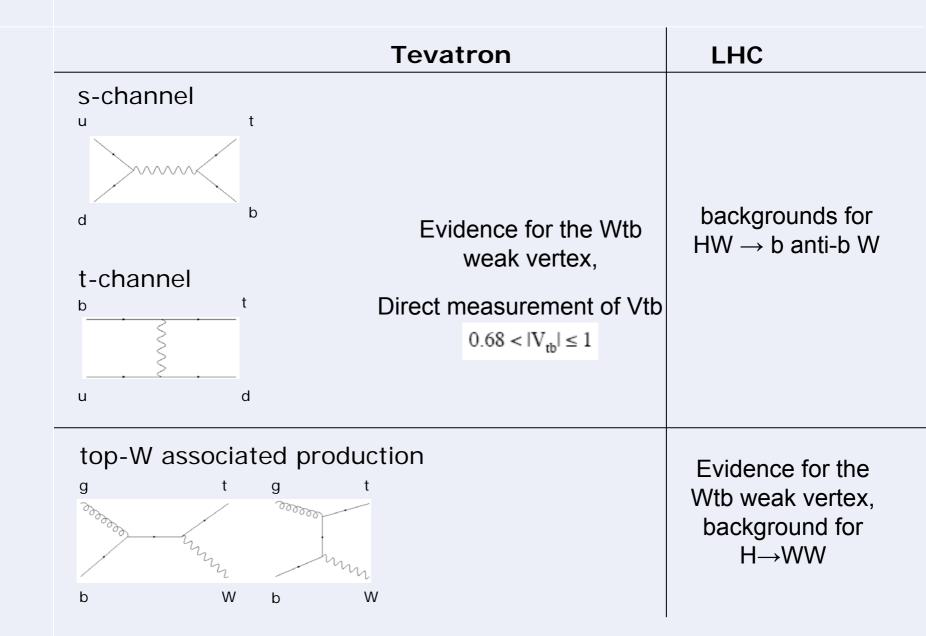
6pb at Tevatron 720pb at LHC

- ☐ Single top can be produced only weakly through the W-t-b vertex
- Evidence for single top allows new tests of the Standard Model
- ☐ The study of single top events passes from the search for their observation at the Tevatron to their study as a significant source of background events in new physics searches at the LHC

## Total cross-sections

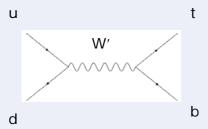


# Standard Model physics



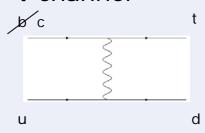
# Physics beyond the Standard Model

#### s-channel



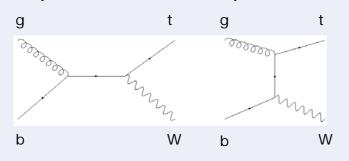
Sensitive to heavy W', charged H, new exotic charged bosons

#### t-channel



sensitive to FCNC

### top-W associated production

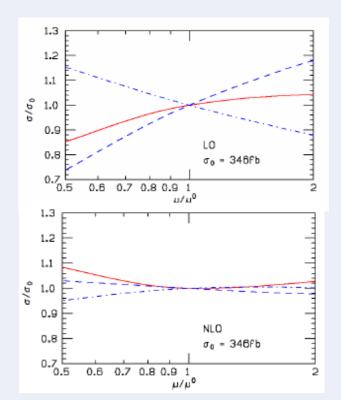


sensitive to anomalous Wt couplings

### Motivation for NLO calculation

- Although some information can be extracted from Born level calculations, the first serious approximation in QCD is obtained by including  $O(\alpha_s)$  radiative corrections
- ☐ Only at next to leading order we obtain accurate predictions of event rates that are sensitive to the structure of jets in the final state

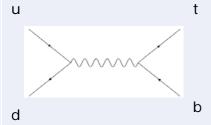
■ Next to leading order calculation give important information about the choice of factorization (dashed) and renormalization (dot-dashed) scales



Wt

# Previous single top NLO results

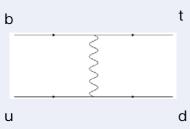
#### s-channel



#### Inclusive s-channel:

Smith and Willenbrock (1996).

#### t-channel



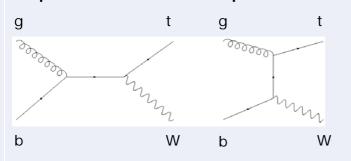
#### Inclusive t-channel:

Bordes, B. van Eijk (1995) Stelzer, Z. Sullivan, and S.Willenbrock (1997)

#### Differential:

Harris, Laenen, Phaf, Sullivan and Weinzierl (2002), Sullivan (2004)

#### top-W associated production



Study of the Wc channel at NLO (Giele, Keller, Laenen)

#### Inclusive Wt:

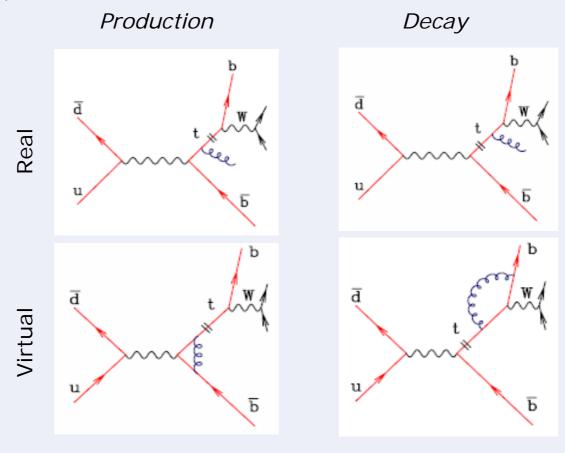
Zhu, 2002, no decay of W and top and no program available

### Calculation

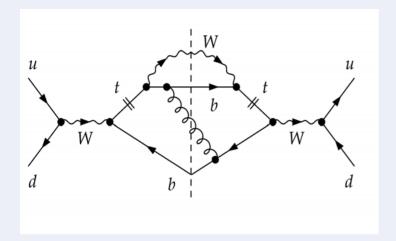
For all the 3 single top channels:

- We extend the previous program including the leptonic decay of the top quark with full spin correlation
- We also consider the effect of the gluon radiation in the decay of the top

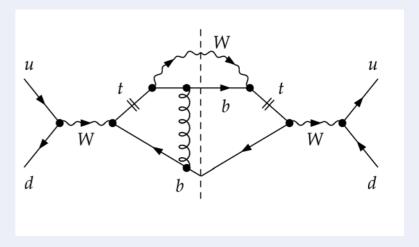
- We work in the on-shell approximation for the top quark: every diagram considered has 1 top quark exactly on its mass shell (double bars)
- $\hfill \Box$  Diagrams without an on-shell top quark are suppressed by  $\Gamma_t/m_t$  (~1%)



□ Neglect interference between radiation in the production and decay stages



Real interference



Virtual Interference

### Physical reason:

- □ Different characteristic time scales for Production (1/ $m_t$ ) and Decay (1/ $\Gamma_t$ )
- ☐ The two stages are separated by a large time
- ☐ The interference average to zero

#### Caveat

- Soft gluons potentially dangerous for this argument
  - time scale not confined to a time of order 1/m<sub>t</sub> -
- Soft radiation phase space is limited
  - resonant propagator and  $E_g\!<\!\Gamma_t$  -
- Because of the cancellation of real and virtual radiation the region of soft radiation is not especially privileged
- $\Box$  For infra-red safe variables these interference effects are expected to be of order α<sub>s</sub>Γ<sub>t</sub>/m<sub>t</sub> (few %)

Fadin, Khoze and Martin (1994), Melnikov and Yakovlev (1994)

#### Confirmation

- □ s-channel process, Pittau (1996)
- ightharpoonup e+e- ightharpoonup time including the sunsequent decay of the top quarks, Macesanu (2002)

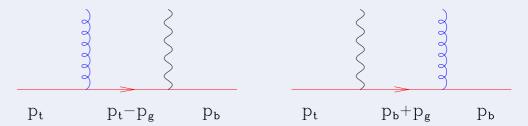
### In MCFM

### J. Campbell and R.K.Ellis

- □ Cancellation of soft and collinear divergences is performed using the subtraction method with massive quark Catani, Dittmaier, Seymour and Trócsányi (2002)
- ☐ Following Nagy and Trócsányi (massless case) we have introduced a tunable parameter (α) to control the size of the subtraction
- $\alpha = 1$  corresponds to the dipole formulas in Catani et al
- $\alpha\!<\!1$  corresponds to calculate the subtractions into a smaller phase space region around the singularities
  - 1. quicker calculation (subtractions are calculated only where they are needed)
  - 2. independence from  $\alpha$  is a valuable check of the implementation
  - 3. reduced mismach between reals and subtractions
- An extention of the method to all the massive dipole formulas has been performed

### Subtraction for radiation in top decay

Real corrections:



■ Require a counterterm that has the same soft and collinear singularities as the full matrix element

$$|\mathcal{M}(\cdots p_t, p_W, p_b, p_g)|^2 \to |\mathcal{M}_0(\cdots p_t, \tilde{p}_W, \tilde{p}_b)|^2 D(p_t, p_g, p_b, p_g, m_t^2, M_W^2)$$

 $\square$   $\widetilde{p}_b$  and  $\widetilde{p}_w$  are defined in such a way to absorb the four-momentum carried away by the gluon

$$p_{t} = p_{b} + p_{W} + p_{g} = \widetilde{p}_{b} + \widetilde{p}_{W} \qquad \qquad \widetilde{p}_{b}^{2} = 0 \quad \widetilde{p}_{W}^{2} = p_{W}^{2}$$

We define  $\widetilde{P}_W$  by a Lorentz boost so that the phase space for the subsequent decay of the W is unchanged, and obtain  $\widetilde{P}_b$  from top momentum conservation above

Usefull for tt also

$$\tilde{p}_{W} = \alpha \left( p_{W} - \frac{p_{t} \cdot p_{W}}{p_{t}^{2}} p_{t} \right) + \beta p_{t}$$

$$\alpha = \frac{p_{t}^{2} - p_{W}^{2}}{2\sqrt{(p_{t} \cdot p_{W})^{2} - p_{W}^{2} p_{t}^{2}}} \quad \beta = \frac{p_{t}^{2} + p_{W}^{2}}{2p_{t}^{2}}$$

# S/B analisys with MCFM

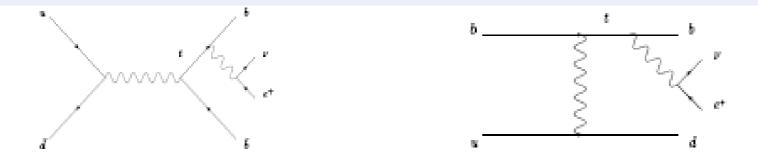
☐ **Signal** at Tevatron calculated at  $\sqrt{s} = 1.98 TeV$ 

By assuming:

$$p_T^e > 20 \,\text{GeV}, \qquad \left| \eta^e \right| < 1.1, \qquad E_T > 20 \,\text{GeV}, \qquad \Delta R^{jj} = 1.0,$$

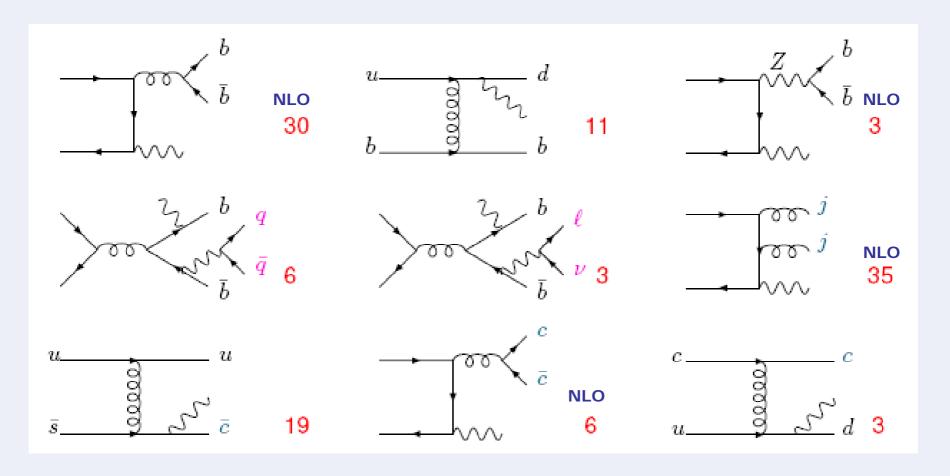
$$p_T^{jet} > 15 \,\text{GeV}, \qquad \left| \eta^{jet} \right| < 2.8, \qquad 140 \,\text{GeV} < m_{blv} < 210 \,\text{GeV}$$

Signal is 18fb from s-channel (7fb) and t-channel (11fb)



# **Backgrounds** at Tevatron calculated with MCFM at $\sqrt{s} = 1.98 TeV$

Including nominal tagging efficiencies and fake rates (fb)



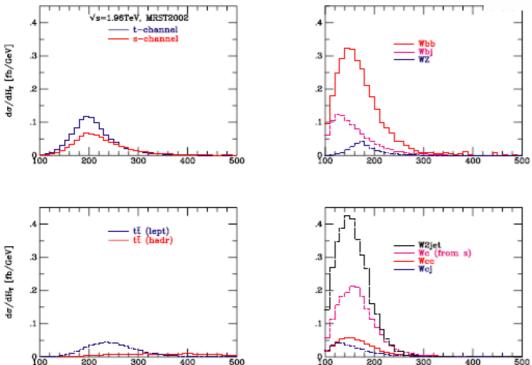
S:B 1:6 with our nominal efficiencies

# $H_T$ distrbution:

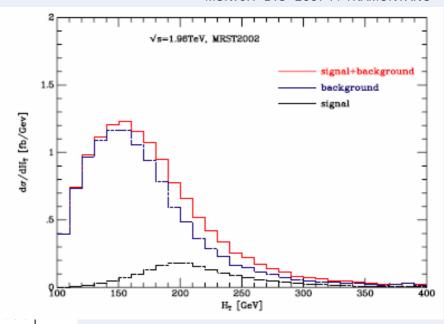
H<sub>T</sub> [GeV]

Scalar sum of jet, lepton and missing  $E_{\mathcal{T}}$ 

$$H_T = |p_T(\text{lepton})| + |\not \!\!E_T| + \sum |p_T(\text{jet})|$$



H<sub>7</sub> [GeV]

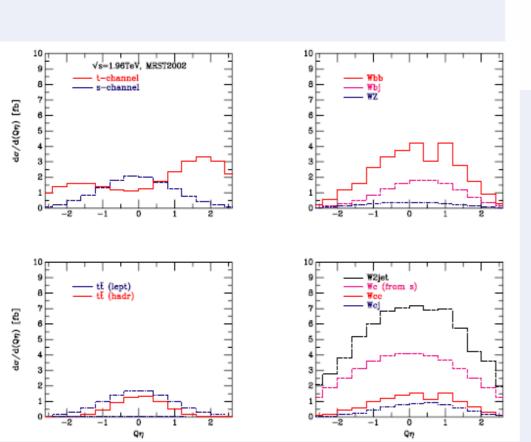


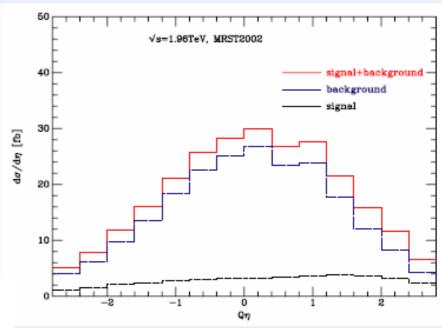
 $\boldsymbol{H}_T$  distribution can only be calculated when the momenta of the decay products are known

**Tevatron** 

# $Q\eta$ distrbution:

Lepton charge in unit of positron X pseudorapidity of the untagged jet





**Tevatron** 

# Wt production and decay

- The last single top channel has also been implemented in MCFM
- We used helicity amplitudes decomposing the top momentum along the initial gluon and an auxiliary massless momentum t<sub>0</sub>

$$u(t)_{\uparrow} = \frac{(f+m_t)}{[t_0g]}|g,-1\rangle, \qquad \bar{u}(t)_{\uparrow} = \langle g,-1|\frac{(f+m_t)}{\langle gt_0\rangle}$$

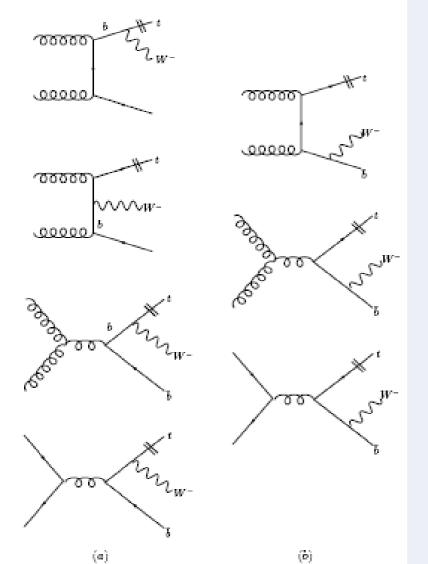
$$u(t)_{\downarrow} = \frac{(f+m_t)}{\langle t_0g\rangle}|g,+1\rangle, \qquad \bar{u}(t)_{\downarrow} = \langle g,+1|\frac{(f+m_t)}{[gt_0]}$$

$$u(t)_{\uparrow}\bar{u}(t)_{\uparrow} + u(t)_{\downarrow}\bar{u}(t)_{\downarrow} = f+m_t$$

$$t_0^{\mu} = t^{\mu} - \frac{m_t^2}{2t \cdot g}g^{\mu}$$

- We used dimensional regolarization
- ☐ Compact expressions are obtained introducing functions that are combination of the usual scalar integrals

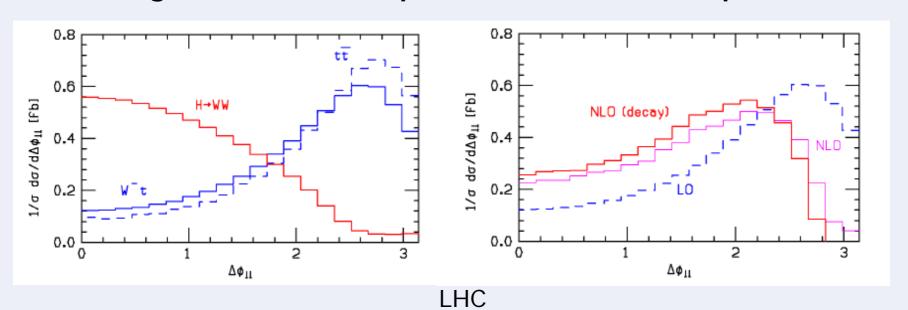
☐ For the Reals of the Wt calculation a complication arise due to the presence of double resonant top contribution in the Real corrections with W-t-b in the final state



- For small  $p_T$  of the b quark the contribution from diagrams (a) is in the b-quark pdf (up to  $\mu_F$  by pdf definition)
- $\checkmark$  μ<sub>F</sub> is limited by the validity of the collinear approximation  $μ_F \sim (m_t + m_W)/4 \sim 65 \text{ GeV}$
- ✓ If we consider the b-quark up to a  $p_T \sim \mu_F \sim 65$  GeV then tt associated production diagrams (b) give small interference and can be separeted
- ✓ We define single top events with a p<sub>T</sub> veto on the b-jet
- ✓ For higher p<sub>T</sub> the tt production gives a better description

- At the LHC the Wt channel has a total rate of 66pb that is larger then the s-channel (10pb)
- □ For 150<m<sub>H</sub><180 Wt associated production is an important source of background for H→WW\* discovery channel
  - WW from the continuum is the main background
  - > tt and Wt are large and comparable

#### Angle between the leptons in the transverse plane



NLO introduce important modification for this distribution

## **Approximations**

- $\rightarrow$  We set  $m_b = 0$ 
  - ~1% (from a LO calculation of the s-channel with massive b)
- We put the top quark on its mass shell
  - ~1% (from a LO calculation with a Breight-Wigner for the top)
- We neglect interference between radiation emitted in the production and decay stages

Expectet to give an  $O(\alpha_s \Gamma_t/m_t)$  correction

> We include no showering and hadronization

We have defined a large pseudocone, but a more detailed analysis should await the combination with a parton shower

# Conclusion

- Single top evidence at Tevatron makes possible new tests of the Standard Model and the evaluation of important backgrounds to the Higgs search at Tevatron and LHC
- All the 3 single top channels including top decay with full spin correlation have been implemented at NLO in MCFC
- ☐ The impelmentation of the subtraction specific for the top decay could also be applied to the decay of the top quarks in the tt associated production process
- Distributions can be generated for single top signal and backgrounds (most at NLO) using MCFM