

# The Double Spin Asymmetry in Exclusive $\pi^+$ Electro-production with CLAS



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DIS2007  
April 19, 2007

# Single $\pi^+$ electro-production

- Polarization dependence of the exclusive reaction

- by measuring the asymmetry

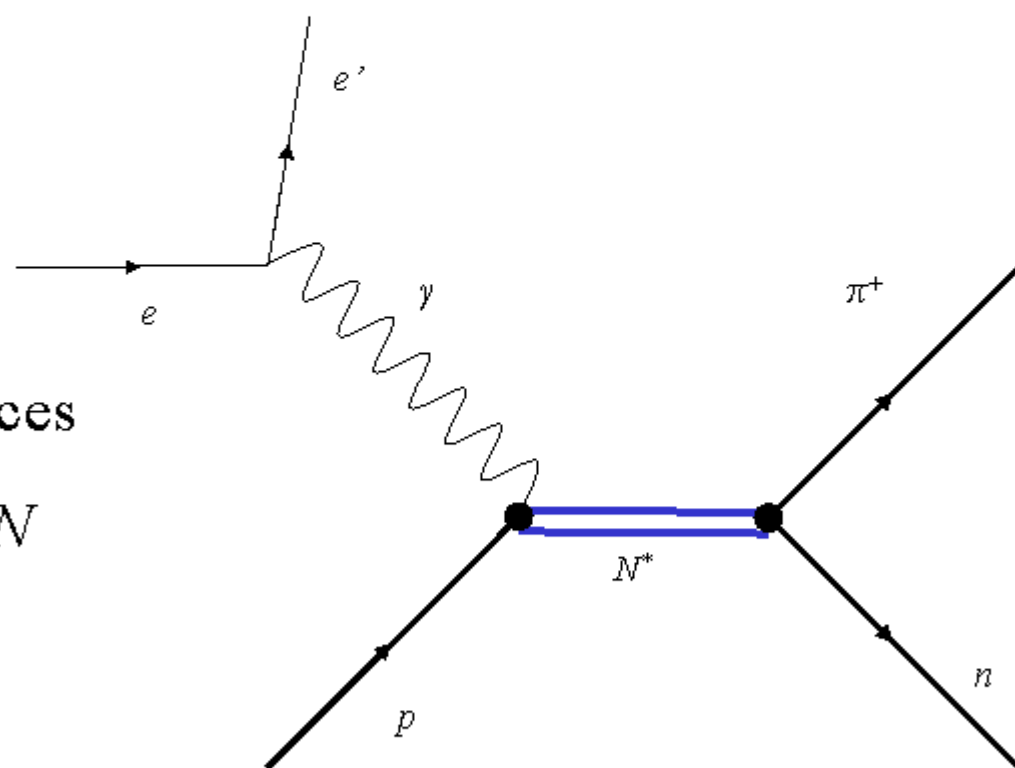
- Measured as a function of four variables.

- Can give insight into the spin structure of the resonances

- large branching ratios into  $\pi N$

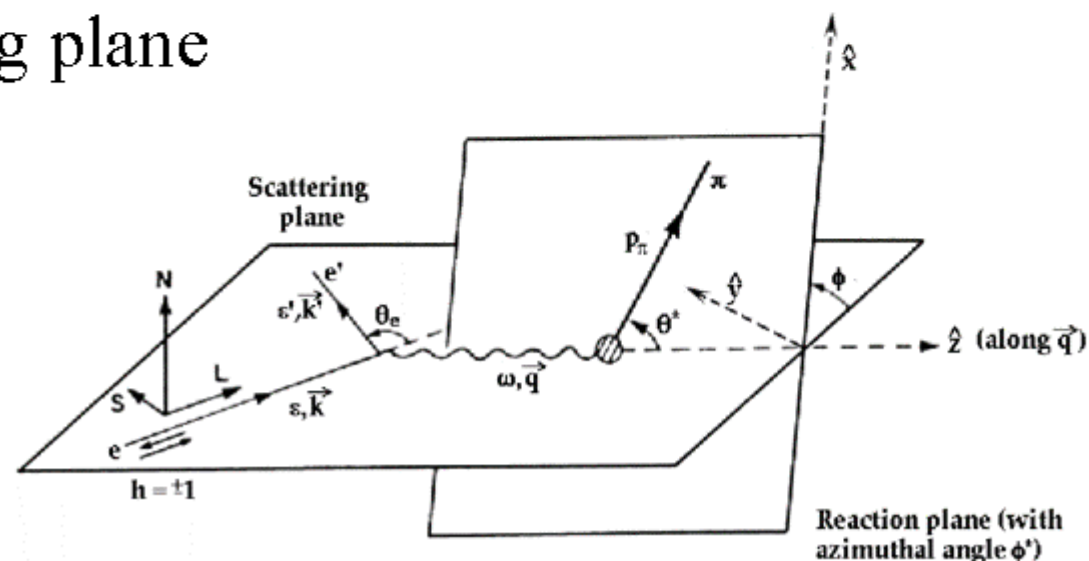
- $P_{11}(1440)$
    - $D_{13}(1520)$
    - $F_{15}(1680)$
    - others

$$ep \rightarrow e\pi^+ N$$



# Kinematic Variables

- $W$ ,  $Q^2$ ,  $\theta^*$ , and  $\phi^*$ 
  - $\theta^*$  and  $\phi^*$  are calculated in the rest frame of the pion/neutron system
  - $\theta^*$  is the angle between the pion production angle and the four momentum transfer  $\vec{q}$
  - $\phi^*$  is the angle between the lepton scattering plane and the hadron scattering plane



# Cross Section

$$d\sigma = \frac{\varepsilon_i}{k_i} \frac{m_e}{\varepsilon_i} \frac{m_N}{E_i} \frac{m_e}{\varepsilon_f} \frac{d^3\mathbf{k}_f}{(2\pi)^3} \frac{1}{2E_\pi} \frac{d^3\mathbf{p}_\pi}{(2\pi)^3} \frac{d^3\mathbf{p}_f}{(2\pi)^3} *$$

$$2\pi\delta^{(4)}(p_i + q - p_\pi - p_f) |\langle p_f, p_\pi | J^\mu | p_i \rangle q^{-2} \langle k_f | j_\mu | k_i \rangle|^2$$

$$\frac{d\sigma}{dE' d\Omega_e d\Omega_\pi} = \Gamma \frac{d\sigma_\nu}{d\Omega_\pi} \leftarrow \begin{array}{|l|} \hline \text{Virtual Photon} \\ \text{Cross Section} \\ \hline \end{array}$$

$$\begin{array}{|l|} \hline \text{Virtual Photon Flux} \\ \hline \end{array} \rightarrow \Gamma = \frac{\alpha}{2\pi^2} \frac{W^2 - m_p^2}{2m_p Q^2} \frac{E'}{E_0}$$

# Response Functions

$$\begin{aligned}
 \frac{d\sigma_\nu}{d\Omega_\pi} = & \frac{|\vec{q}|}{q_\gamma^{CM}} [R_T + P_y R_T^y + \epsilon_L (R_L + P_y R_L^y) \\
 & + \sqrt{2\epsilon_L(1+\epsilon)} ((R_{LT} + P_y R_{LT}^y) \cos \phi^* + (P_x R_{LT}^x + P_z R_{LT}^z) \sin \phi^*) \\
 & + \epsilon ((R_{TT} + P_y R_{TT}^y) \cos 2\phi^* + (P_x R_{TT}^x + P_z R_{TT}^z) \sin 2\phi^*) \\
 & + h\sqrt{2\epsilon_L(1-\epsilon)} ((R_{LT'} + P_y R_{LT'}^y) \sin \phi^* + (P_x R_{LT'}^x + P_z R_{LT'}^z) \cos \phi^*) \\
 & + h\sqrt{1-\epsilon^2} (P_x R_{TT'}^x + P_z R_{TT'}^z) ]
 \end{aligned}$$

$$(P_x, P_y, P_z) = P_t (\sin \theta_\gamma \cos \phi^*, -\sin \theta_\gamma \sin \phi^*, \cos \theta_\gamma)$$

$$\begin{aligned}
 \frac{d\sigma_{et}}{d\sigma_0} = & \sqrt{2\epsilon_L(1-\epsilon)} [\sin \theta_\gamma (R_{LT'}^x \cos^2 \phi^* - R_{LT'}^y \sin^2 \phi^*) + R_{LT'}^z \cos \theta_\gamma \cos \phi^*] \\
 & + \sqrt{1-\epsilon^2} [R_{TT'}^x \sin \theta_\gamma \cos \phi^* + R_{TT'}^z \cos \theta_\gamma]
 \end{aligned}$$

# Cross Section (cont.)

$$\frac{d\sigma_v}{d\Omega_\pi^*} = \frac{|\mathbf{q}|}{q_\gamma^{CM}} \left\{ \frac{d\sigma_0}{d\Omega_\pi^*} + P_B \frac{d\sigma_e}{d\Omega_\pi^*} + P_T \frac{d\sigma_t}{d\Omega_\pi^*} - P_B P_T \frac{d\sigma_{et}}{d\Omega_\pi^*} \right\}$$

Polarization independent part  
of cross section

Part of the cross section that depends  
on both beam and target polarization

$$\frac{\sigma_{et}}{\sigma_0} = A_{et} = \frac{(\sigma(+,-) - \sigma(+,+)) + (\sigma(-,+) - \sigma(-,-))}{\sigma(+,+) + \sigma(-,-) + \sigma(+,-) + \sigma(-,+)}$$

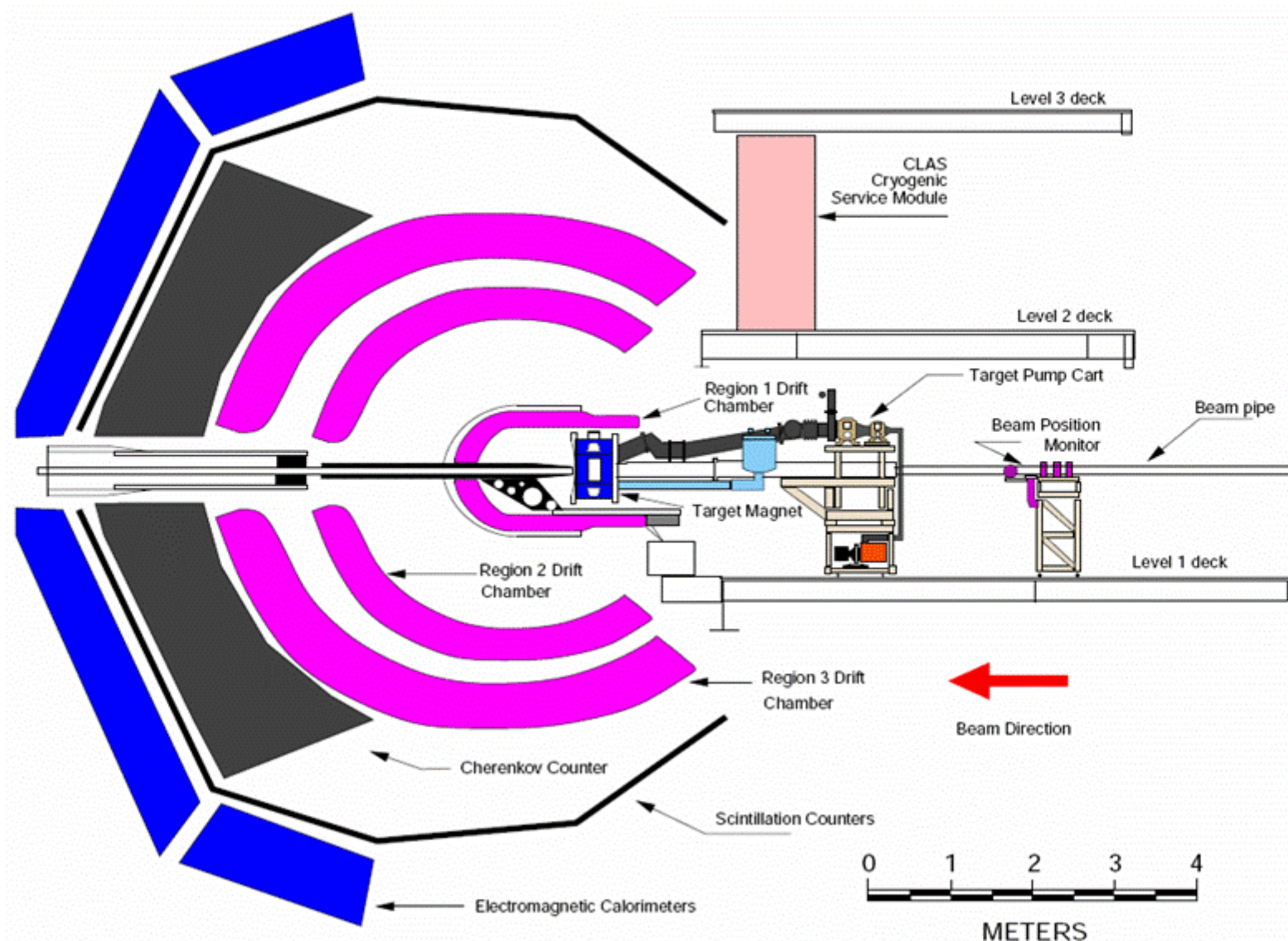
Double spin Asymmetry

$$A_{et} = \frac{1}{f_D P_B^\uparrow P_T^\downarrow} \frac{N'_{\downarrow\uparrow} - N'_{\uparrow\uparrow} + r_B (N'_{\uparrow\downarrow} - N'_{\downarrow\downarrow})}{N'_{\downarrow\uparrow} + N'_{\uparrow\uparrow} + r_T (N'_{\uparrow\downarrow} + N'_{\downarrow\downarrow})}$$

# eg1b aka EG2000

- Hall B at Jefferson Lab (CLAS)
- Polarized electron beam
  - Four different beam energies (1.6, 2.5, 4.2, 5.6 GeV)
- Polarized nuclear target
  - Polarized  $\text{NH}_3$  as a proton target
  - Polarized  $\text{ND}_3$  as an effective neutron target
  - Carbon and Nitrogen targets for background studies

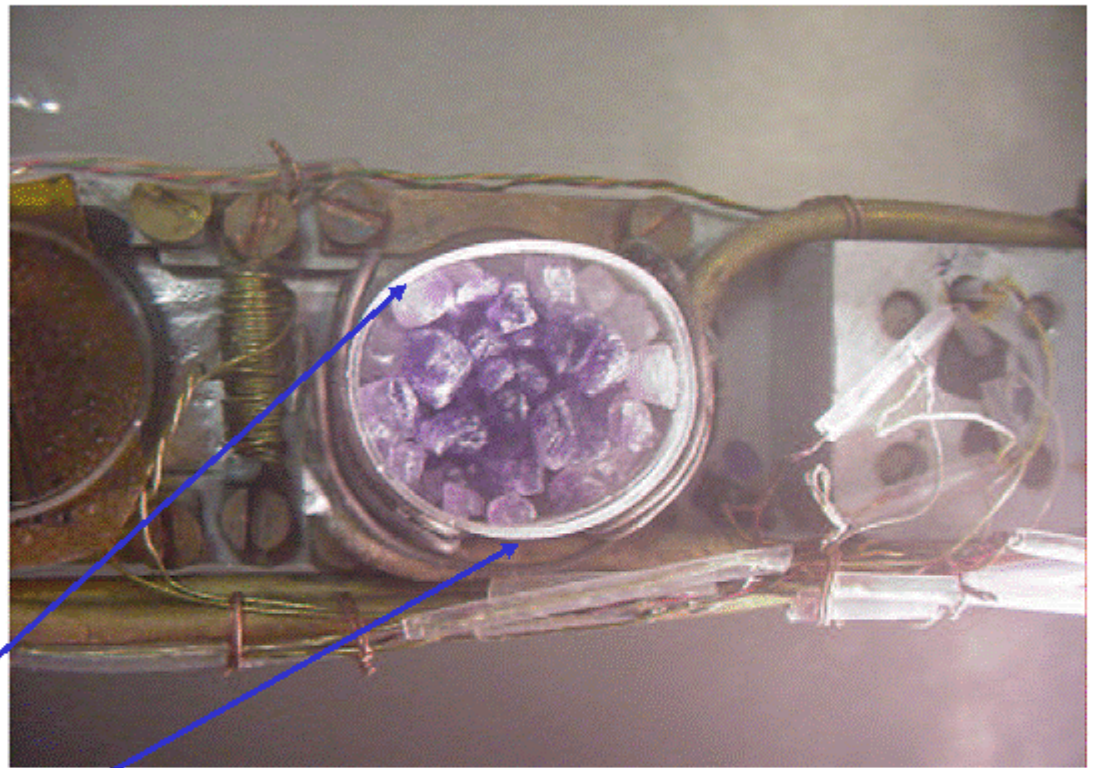
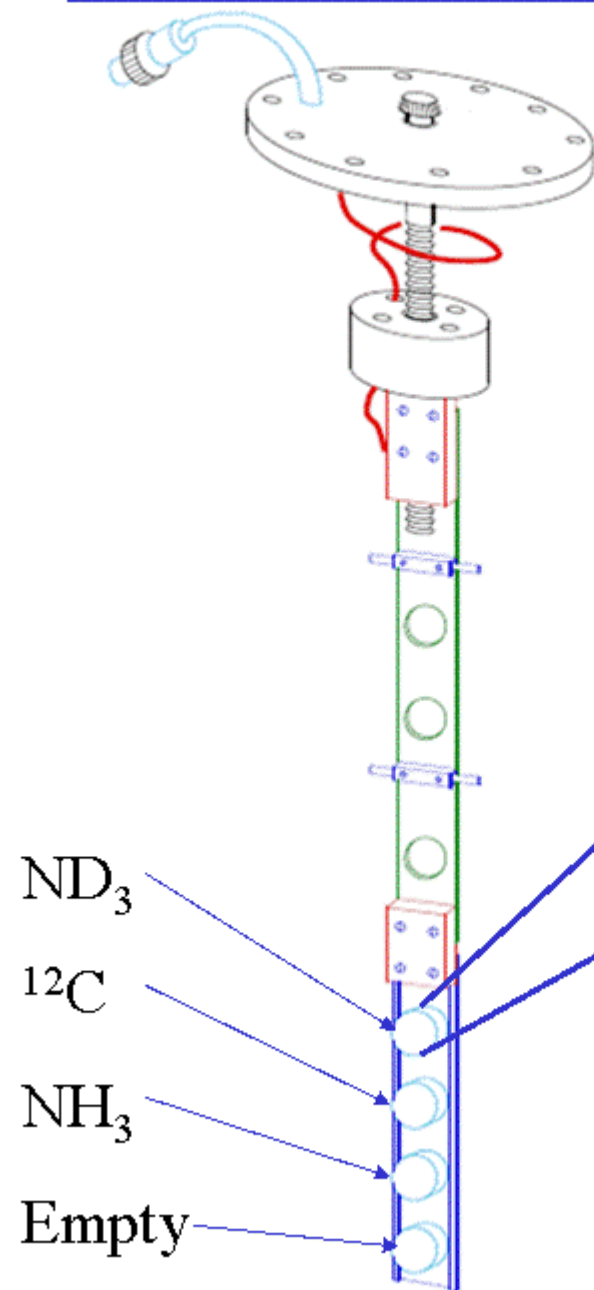
# EG1 Target in CLAS



CEBAF Large Acceptance Spectrometer



# Polarized Target



- DNP used to polarize target
- Super-conducting magnet (5T)
- Microwaves (140 GHz)
- He evaporation refrigerator (1K)

# Asymmetry Instead of Cross Sections

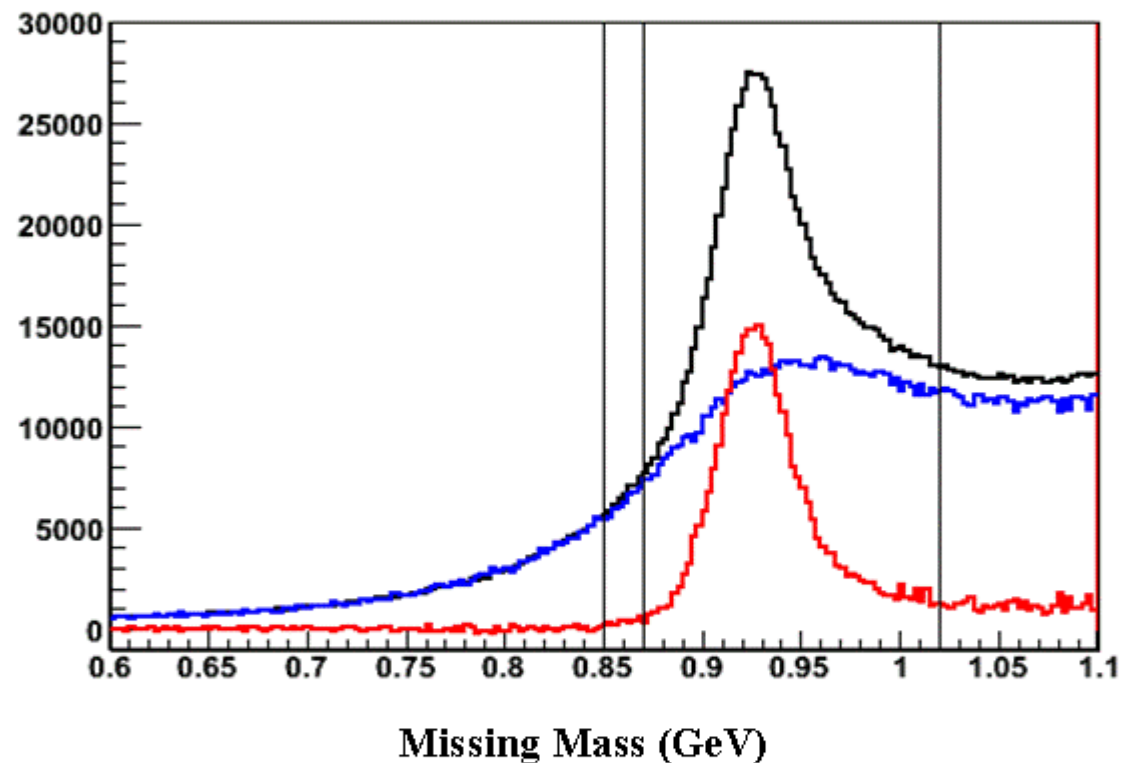
- Insensitive to acceptance
  - The acceptance cancels in the asymmetry
- Only relative luminosities are important
- Limited background contribution
  - Spin independent backgrounds enter as a dilution
- Interference terms allow relatively small quantities to be measured more easily

# Asymmetry

- Fast flipping of electron helicity (pseudo-random sequence at 1 Hz)
  - Polarization is nearly identical for both helicities ( $\sim 70\%$ )
- Slow flipping of proton helicity (once per  $E_0$ )
  - Large, systematic difference in polarizations for different helicities ( $\sim 70\%$  with fairly large variations for two helicities)
- Define two data groups
  - Data taken with positive target helicity and data taken with negative target helicity
  - Data from the second group normalized by the ratio of the polarizations from the two groups.

# Event Selection

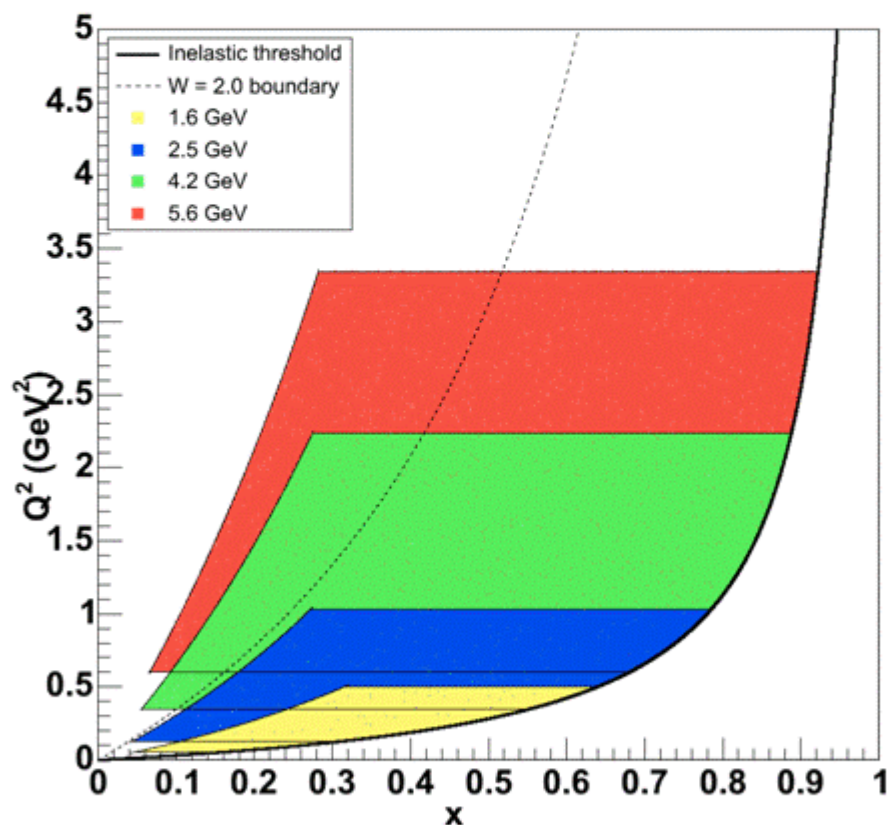
- Events require the detection of scattered electron and the pion
- Neutron not detected, but channel determined by a cut on missing mass of undetected particle



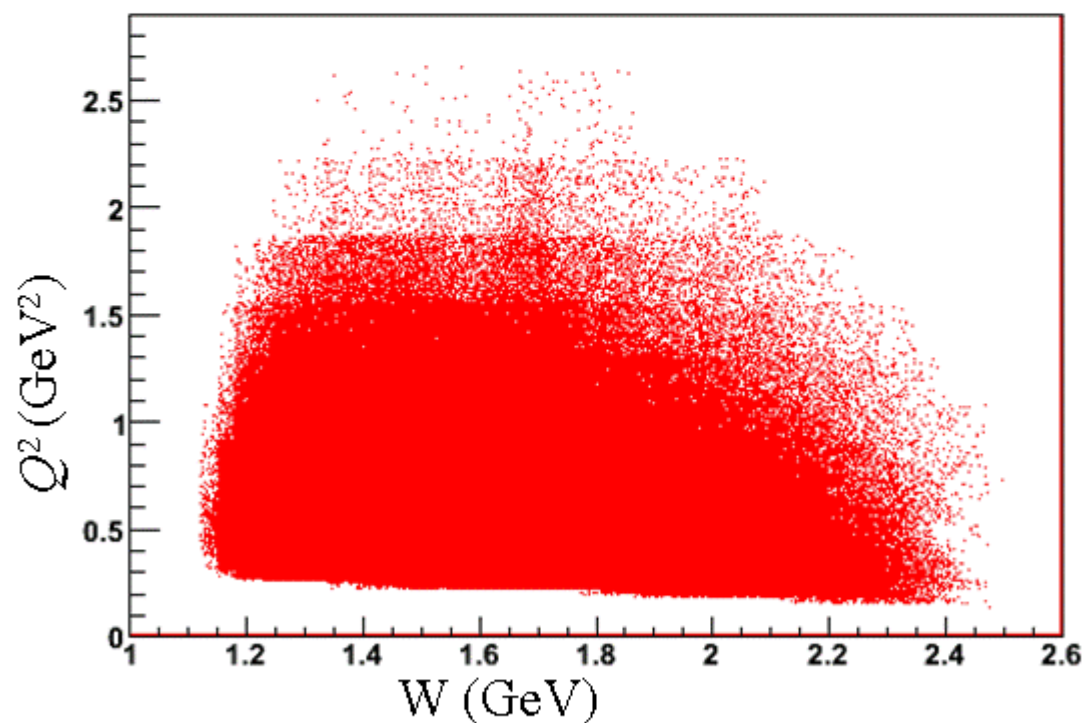
- Necessary due to low efficiency in detection of neutrons

# eg1b Kinematic Coverage

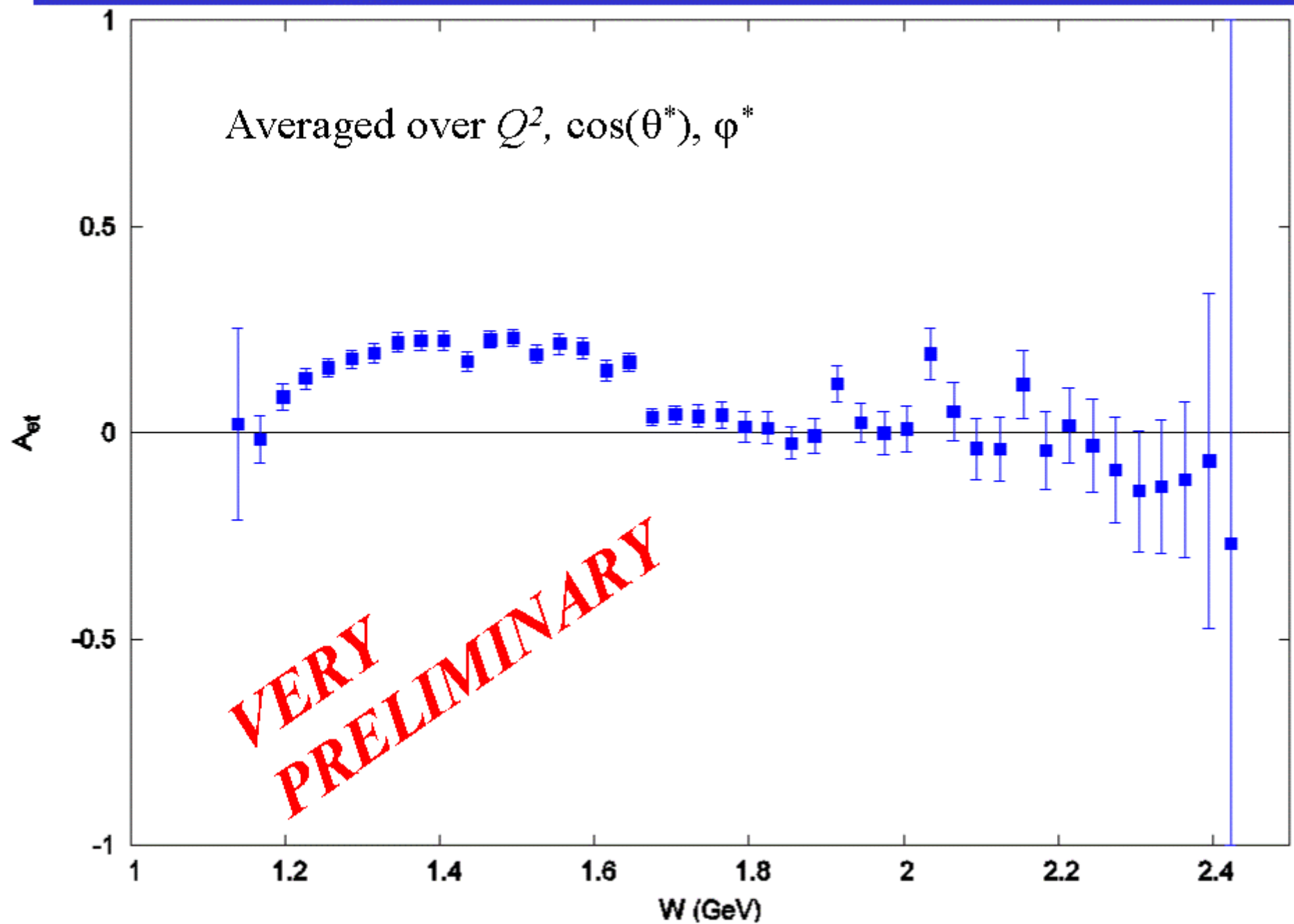
Coverage for inclusive events



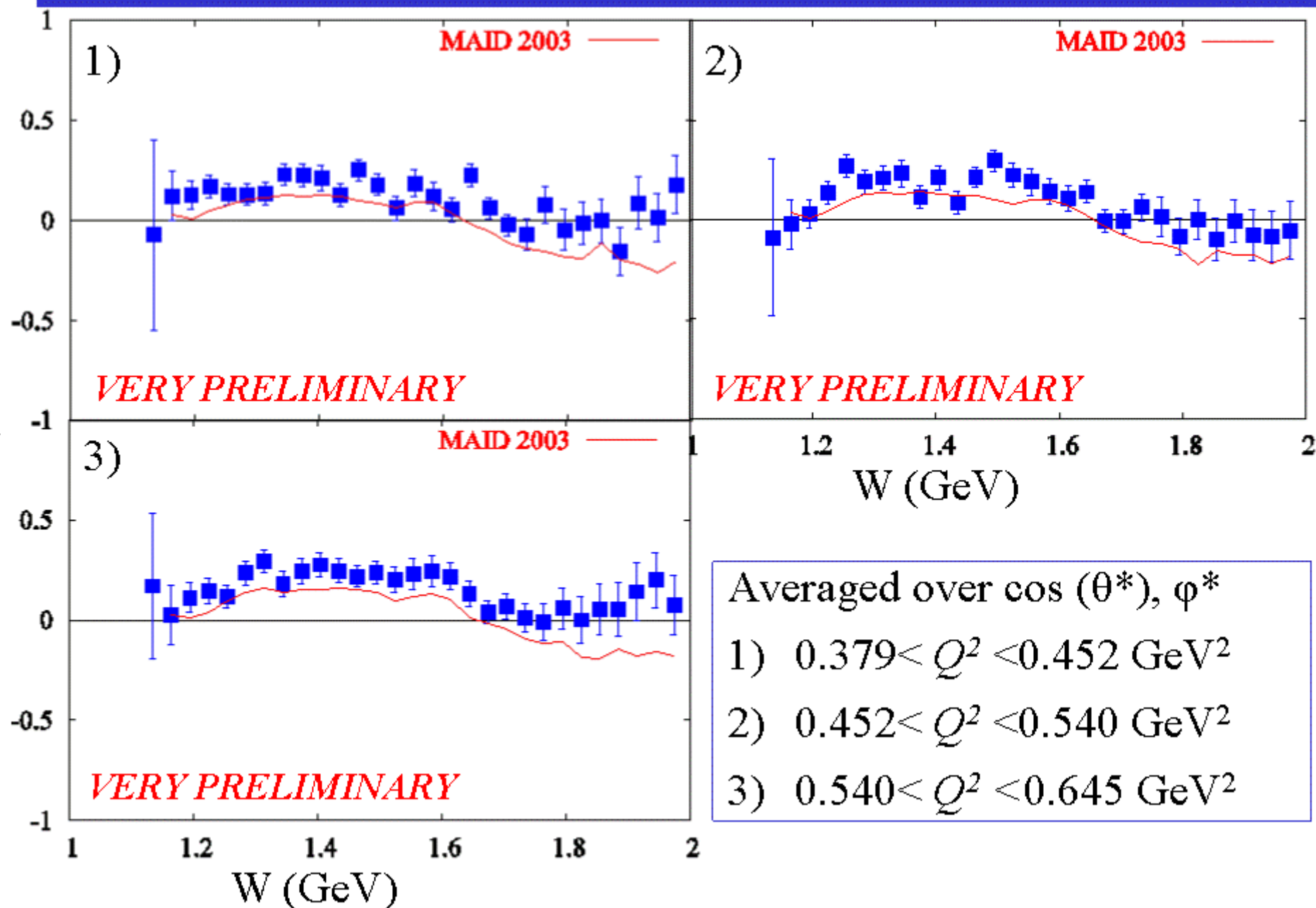
Coverage for exclusive  $\pi^+$  electro-production events.  
 $E_0 = 4.2 \text{ GeV}$



# $A_{\text{et}}$ as a Function of $W$

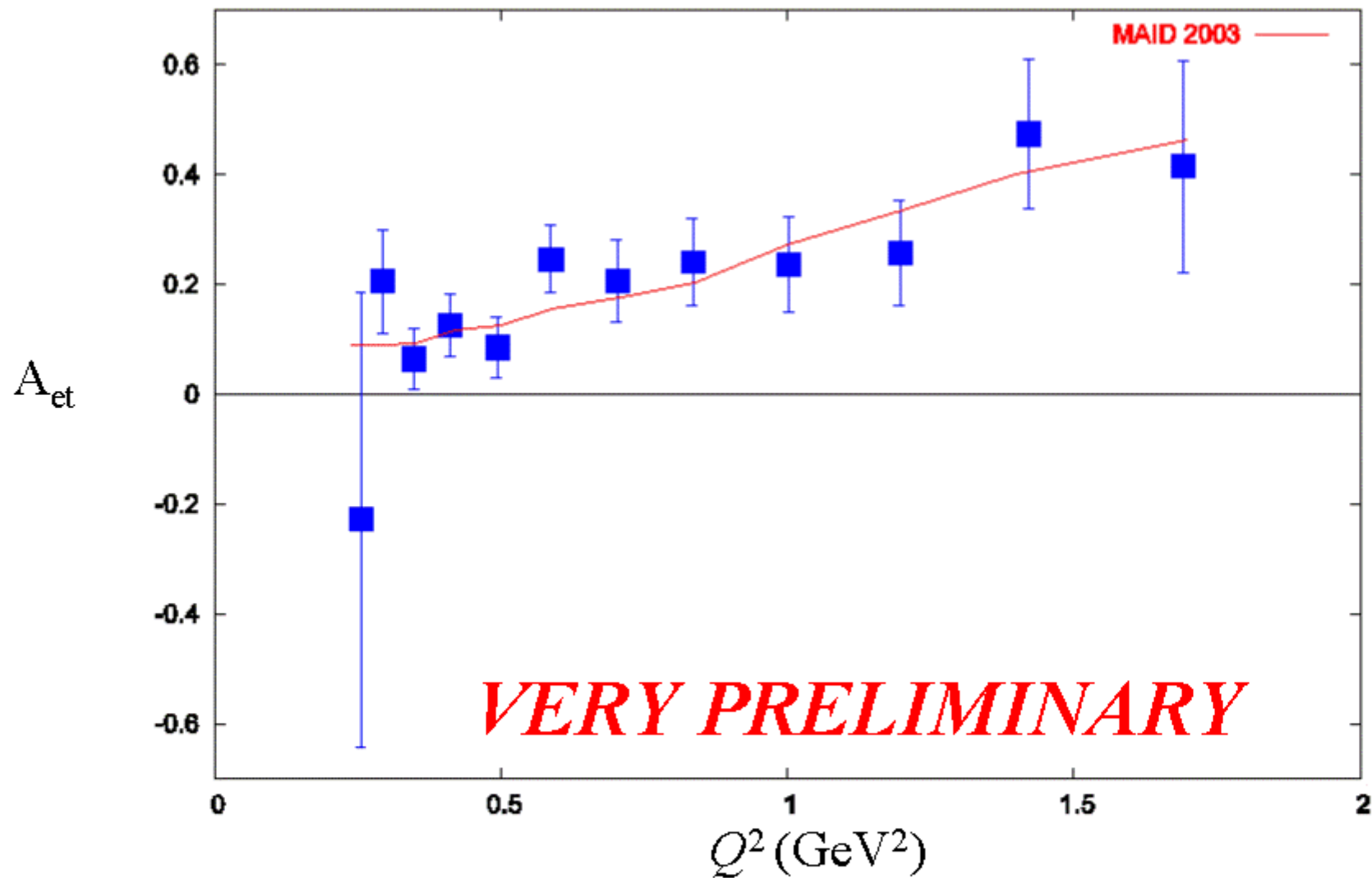


# $A_{\text{et}}$ as a Function of $W$ for fixed $Q^2$





# $A_{\text{et}}$ as a Function of $Q^2$ for fixed $W$

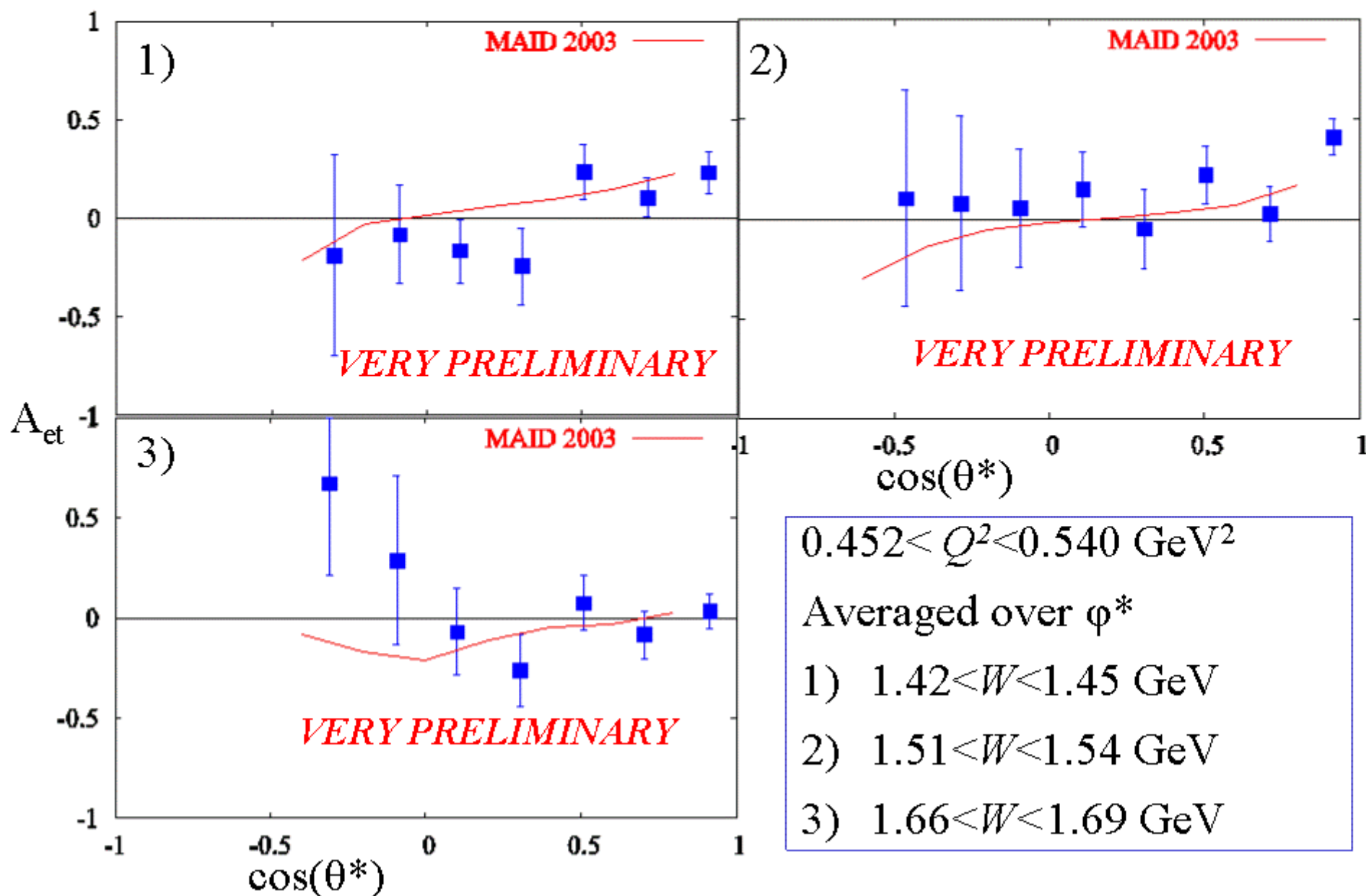


Averaged over  $\cos(\theta^*)$ ,  $\varphi^*$

$1.420 < W < 1.450 \text{ GeV}$



# $A_{\text{et}}$ as a Function of $\cos(\theta^*)$ for fixed $W$ and $Q^2$



$0.452 < Q^2 < 0.540$  GeV<sup>2</sup>

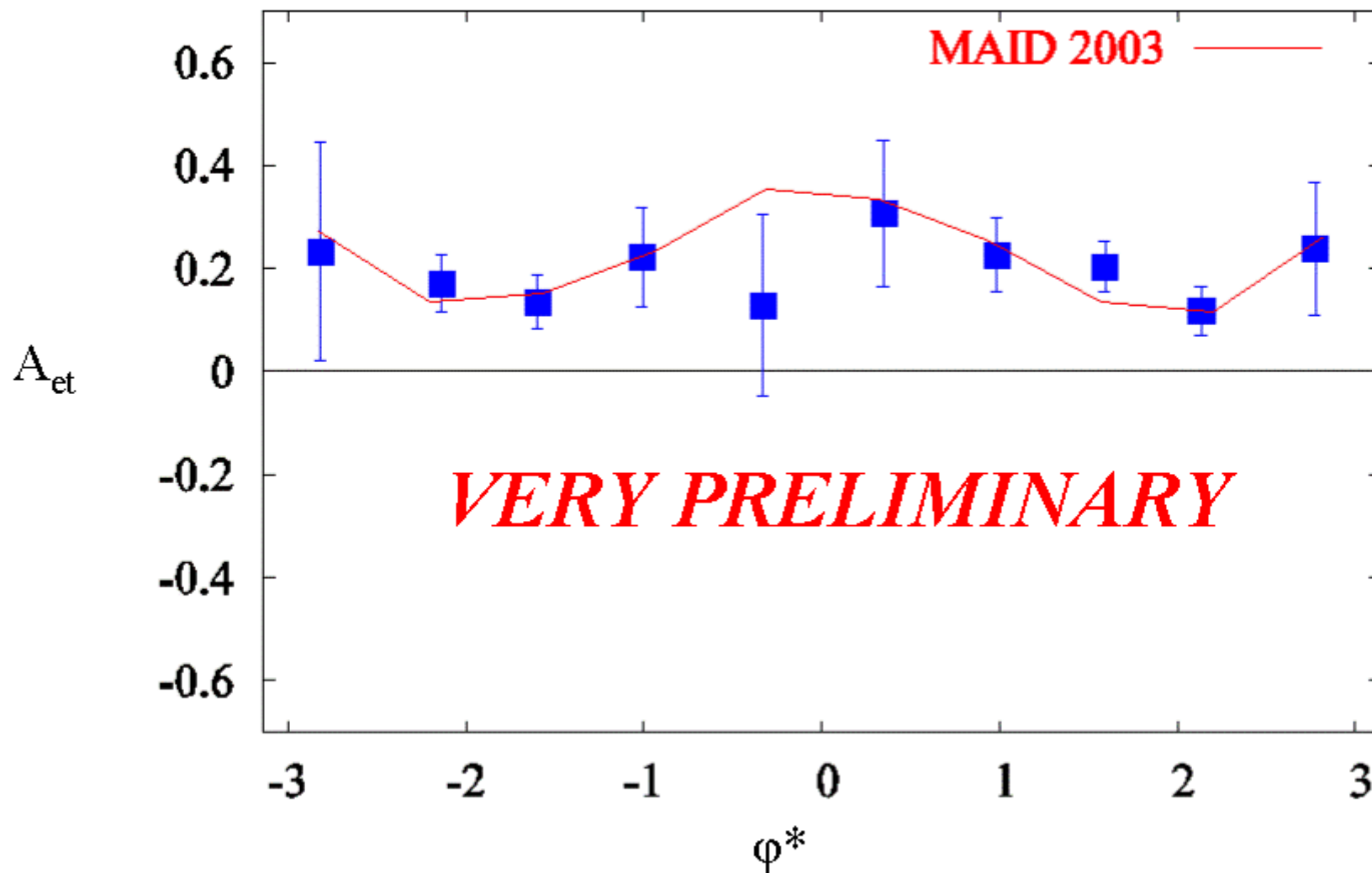
Averaged over  $\phi^*$

1)  $1.42 < W < 1.45$  GeV

2)  $1.51 < W < 1.54$  GeV

3)  $1.66 < W < 1.69$  GeV

# $A_{\text{et}}$ as a Function of $\varphi^*$ for fixed $W$



Averaged over  $Q^2$ ,  $\cos(\theta^*)$

$1.420 < W < 1.450$

# Summary

- Spin Dependence of  $ep \rightarrow e\pi^+ N$  is useful in determining the spin structure of the resonances
- The double spin asymmetry has been measured as a way to show this dependence
- The double spin asymmetry has been measured as a function of  $W$ ,  $Q^2$ ,  $\theta^*$ , and  $\phi^*$
- Analysis of single spin asymmetries are also underway
- More energy settings, much more data
- Analysis of other channels