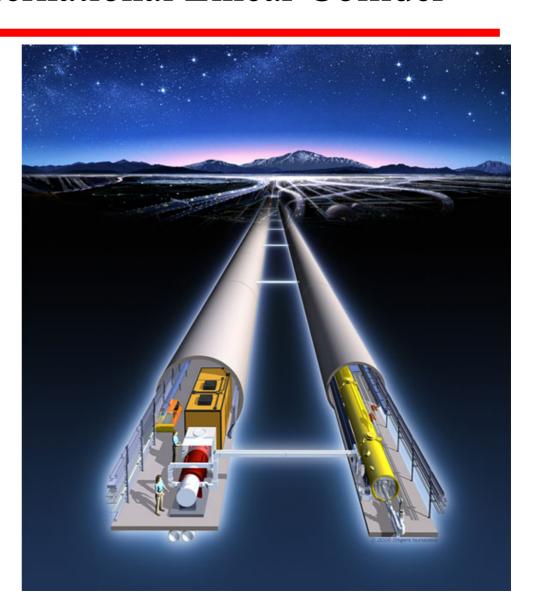
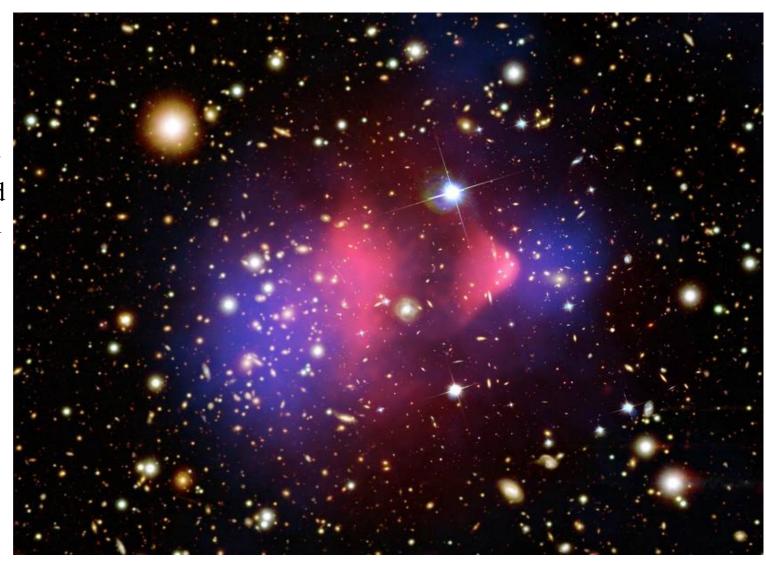
Physics with Flavour at the International Linear Collider

- Introduction
- Progress with the ILC
- ILC detectors
- Flavour identification at the ILC
- Some physics of, and with, flavour
- Summary



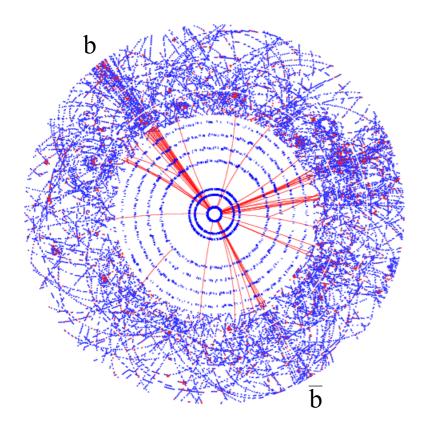
Why new colliders? There is more to heaven and earth...

- collision of two galaxy clusters seen using the Chandra X-ray Observatory, Hubble, ESO's Very Large Telescope and the Magellan optical telescopes.
- "Direct empirical proof of the existence of dark matter."
- Now we must study dark matter in the laboratory.

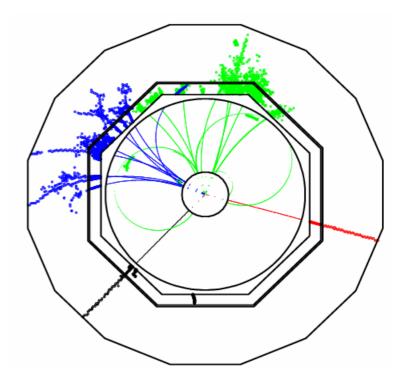


Why the ILC? Electron-positron collisions complement pp

pp → HX as expected in ATLAS detector at LHC:

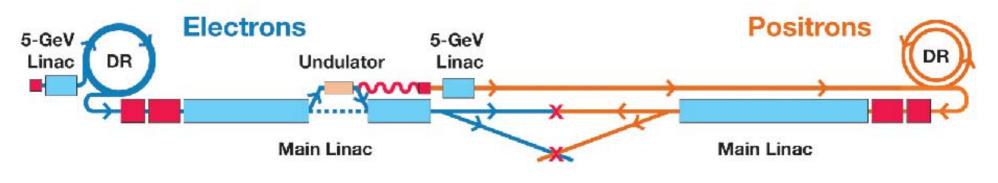


e⁺e⁻ → HZ as expected in LDC detector at ILC:

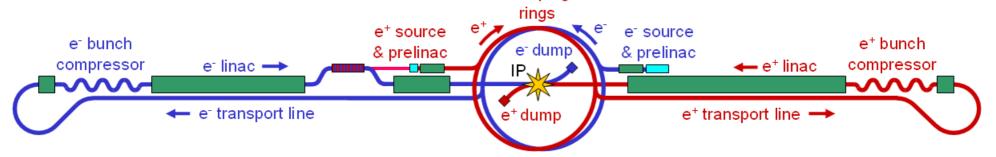


The International Linear Collider

■ ILC design July 2006 (Vancouver LCWS)



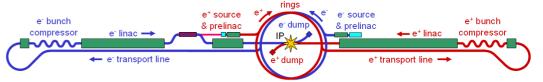
■ ILC now (Reference Design Report).



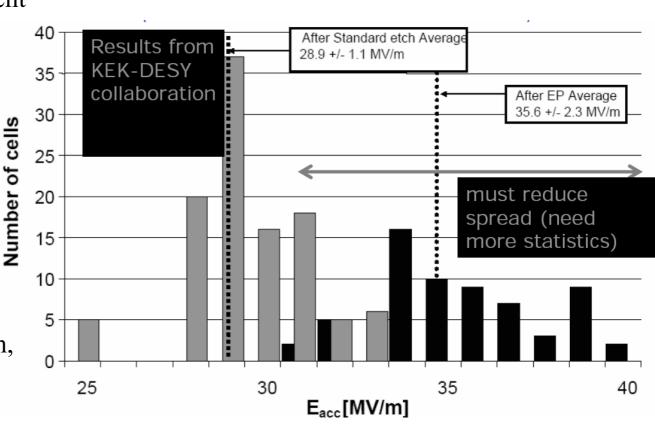
damping

- $\sqrt{s} = 200...500 \text{ GeV}$, upgrade to 1 TeV.
- $\mathcal{L} \sim 2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$, i.e. 500 fb⁻¹ in first 4 years.
- Energy stability better than 0.1%.
- Electron polarisation 80% (e⁺ 30%).
- Cost \$4.9B + \$1.8B + 13k person-years.

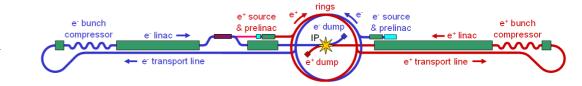
Superconducting cavities



- ILC relies on industrial production of high gradient SC cavities.
- Need peak gradient of 35 MV/m for $\sqrt{s} = 500 \text{ GeV}$.
- Material of choice niobium.
- Surface smoothness critical.
- Average gradient after standard etch ~ 29 MV/m, after electro-polishing ~ 36 MV/m.
- Single crystal cavity up to 45 MV/m.



Beam Delivery System



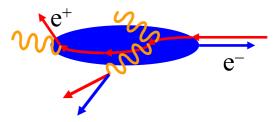
- Final focus.
- Luminosity given by:

$$\mathcal{L} = \frac{n_b N^2 f_{rep}}{A} H_D,$$

where:

- \bullet n_b, number of bunches in train.
- N, number of particles per bunch.
- \bullet f_{rep}, bunch train frequency.
- ♦ A, area of bunch at IP.
- → H_D, beam-beam enhancement factor.
- Need smallest possible crosssectional beam areas.

Particles pass through intense field of opposing beam, radiate photons.

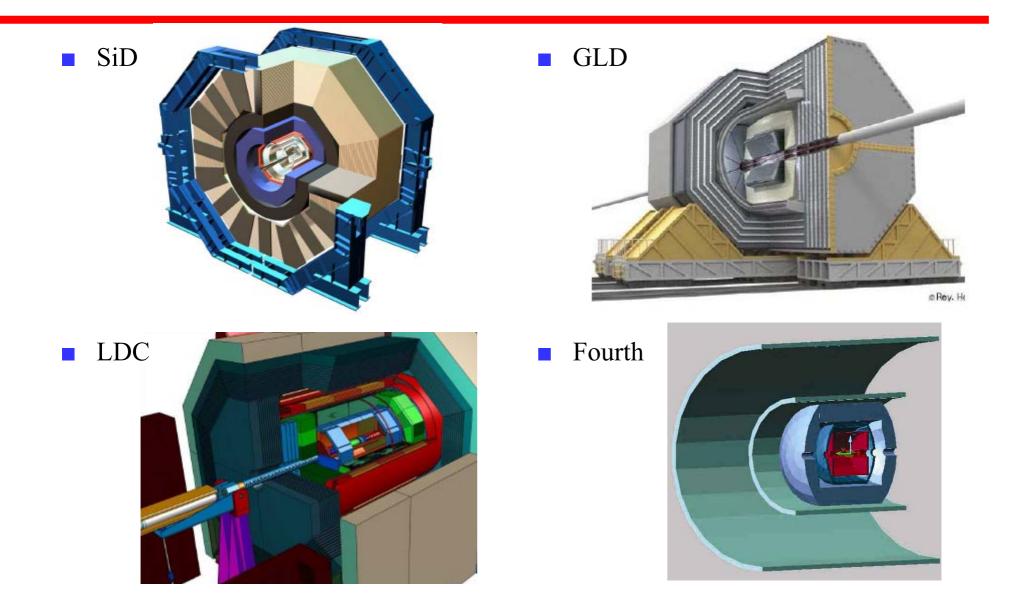


- These beamstrahlung photons interact with field of bunches, and generate e⁺e⁻ pairs.
- Beam-beam effects characterized by disruption parameter:

$$D_{x,y} = \frac{2r_e N\sigma_z}{\gamma \sigma_{x,y} (\sigma_x + \sigma_y)}.$$

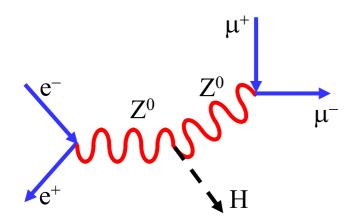
Flat beam, $\sigma_y < \sigma_x$, better than round: beam height ~5 nm, width ~500 nm.

Four detector concepts

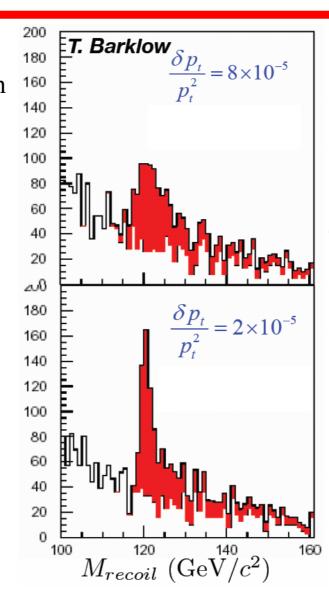


Detector requirements – tracking

Excellent momentum resolution to reconstruct "recoil" mass, e.g. when Higgs decays invisible in process:



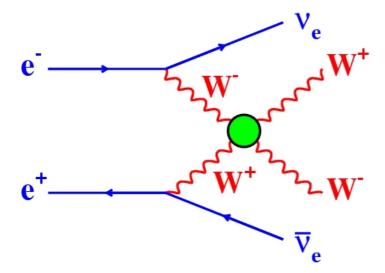
Must couple with large acceptance and robust pattern recognition capabilities to cope with multi-jet environment, e.g. six jets in $e^+e^- \rightarrow t \, \overline{t}$ events.



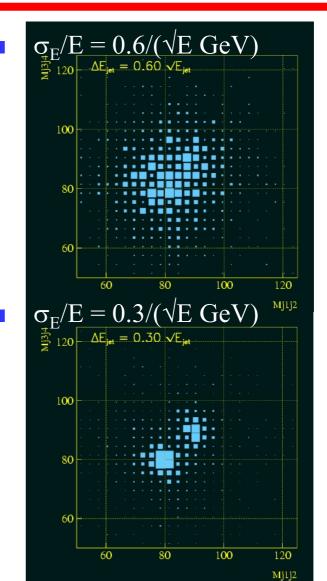
Target resolution: $\frac{\delta p_T}{^2} \sim 5 \times 10^{-5} \text{ GeV}^{-1}.$

Detector requirements – calorimetry

- Want to be able to separate final states $W \rightarrow q \overline{q}'$ and $Z \rightarrow q \overline{q}$.
- Allows e.g. study of processes:

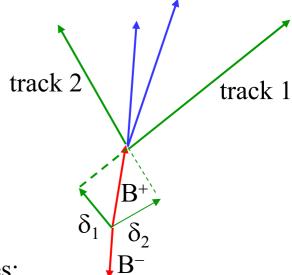


Good jet energy resolution required.



Detector requirements – vertexing

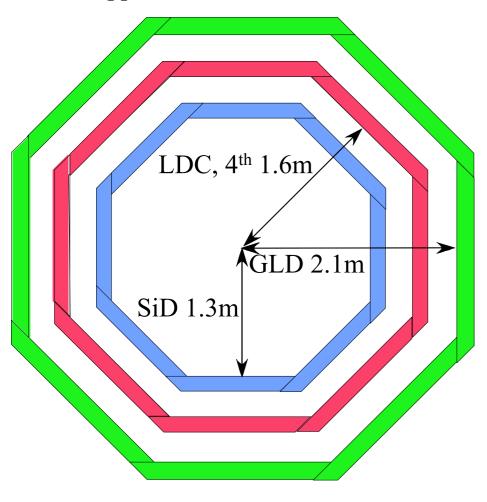
- Efficient identification required for τ leptons, c and b quarks.
- Average impact parameter δ of B decay products ~ 300 μm, of charmed particles less than 100 μm.
- Must resolve all tracks in dense jets.
- Cover large solid angle: forward/backward events are of particular significance for studies with polarised beams.
- Stand-alone reconstruction desirable.



- Implies:
 - Si pixels $\sim 20 \times 20 \ \mu m^2$ or smaller.
 - \bullet Hit resolution better than 5 μ m.
 - First measurement at $r \sim 15$ mm.
 - ♦ Five layers out to radius of about 60 mm, i.e. total ~ 10⁹ pixels
 - Material $\sim 0.1\% X_0$ per layer.
 - Detector covers $|\cos \theta| < 0.96$.

The four detector concepts

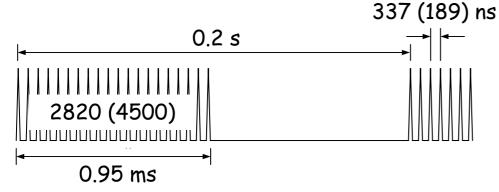
Approximate relative sizes:



- ECal and HCal inside coil.
- SiD:
 - ♦ W/Si ECal, Fe/RPC HCal.
 - ◆ All silicon tracking + VXD.
- LDC:
 - W/Si ECal, Fe/Scint or Fe/RPC HCal.
 - ◆ TPC + silicon tracking + VXD.
- **GLD**:
 - W/Scint Ecal, Pb/Scint Hcal.
 - ◆ TPC + silicon tracking +VXD.
- **Fourth:**
 - Cystal ECal, W/multi-fibre HCal.
 - ◆ TPC + silicon tracking +VXD.

Vertex detectors – constraints due to machine

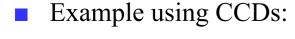
- Minimum beam pipe radius ~14 mm.
- Pair background at this radius in ~4T field causes ~0.03 (0.05) hits per BC and mm² at $\sqrt{s} = 500$ (800) GeV.
- Bunch train structure:

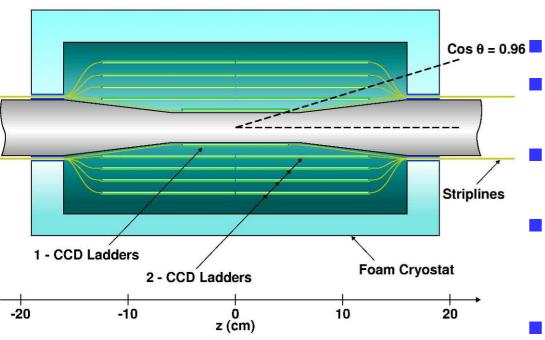


For 10^9 pixels of size $20 \times 20 \mu m^2$, implies readout or storage of signals ~ 20 times during bunch train to obtain occupancy less than $\sim 0.3\%$ (0.9%).

- Must withstand:
 - ◆ Radiation dose of ~50 krad p.a.
 - ◆ Annual dose of neutrons from beam and beamstrahlung dumps
 ~1 x 10⁹ 1 MeV equiv. n/cm².
- Must cope with operation in magnetic field of up to 5 T.
- Must be robust against beam-related RF pickup and noise from other detectors.
- None of available sensor technologies yet satisfies all these requirements.

Conceptual vertex detector design

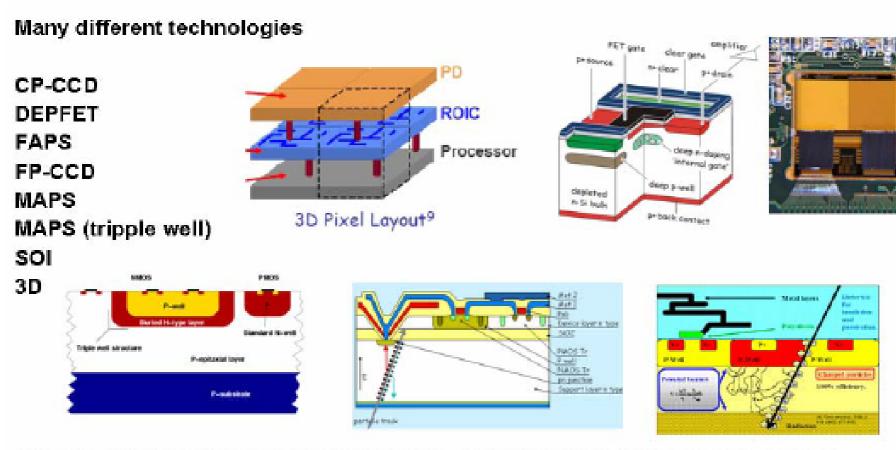




- Surrounded by ~2 mm thick Be support cylinder.
- Allows Be beam pipe to be of thickness of ~0.25 mm.

- Pixel size $20 \times 20 \mu m^2$, 8×10^8 pixels in total.
- 50 MHz readout of inner layer.
- Standalone tracking using outer 4 layers.
- Hits in first layer improve extrapolation of tracks to IP.
- Sensor operation at ~220 K, gas cooling, additional evaporative cooling for electronics if needed.
- Readout and drive connections routed along BP.
- Important that access to vertex detector possible, "roll" outer tracker along BP as done at SLD.

Sensors for the vertex detector



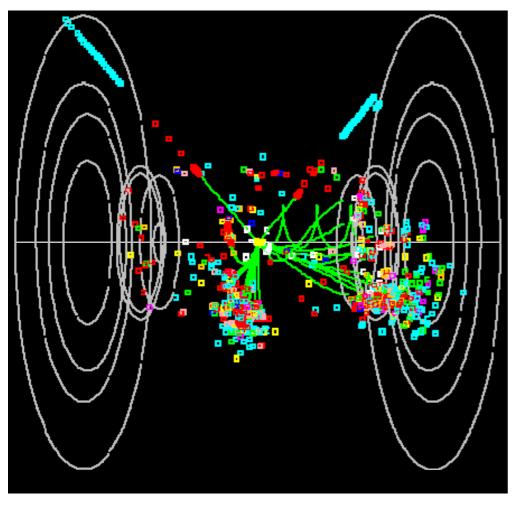
Excellent reviews by Konstantin Stefanov, Joel Goldstein, Rainer Richter, Hans Krueger, Marc Winter, Devis Contarato, Valerio Re, Tadashi Nagamine, Ray Yarema, Wojciech Kucewicz

May 31, 2006

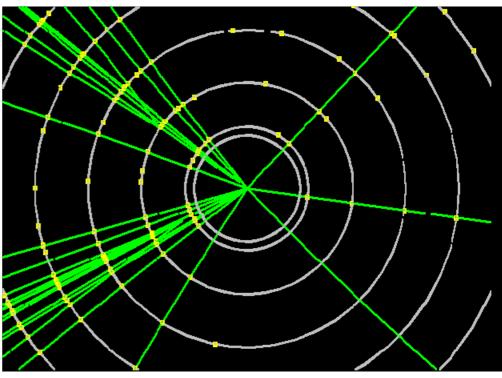
H-G Moser MPI fuer Physik/HLL

Finding decay vertices using the vertex detector

■ E.g. $e^+e^- \rightarrow ZH \rightarrow \mu^+\mu^-$ jet jet.



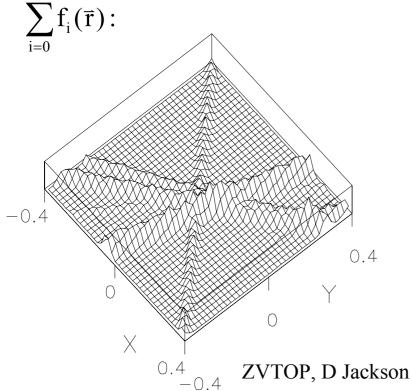
Vertex region:



Challenge is to associate all charged tracks with correct vertex in high track density environment.

Topological vertex reconstruction

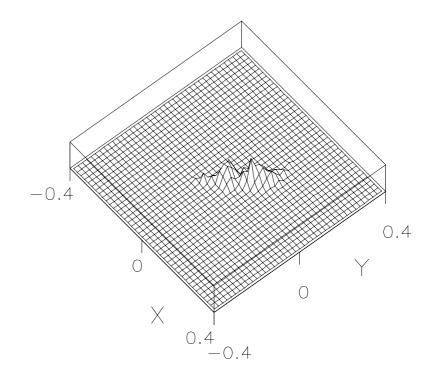
- Using VXD hits, tracks approximated as probability tubes, $f_i(\vec{r})$, beam spot as ellipsoid, $f_0(\vec{r})$.
- E.g. N tracks, integrating over z,



From these define vertex function:

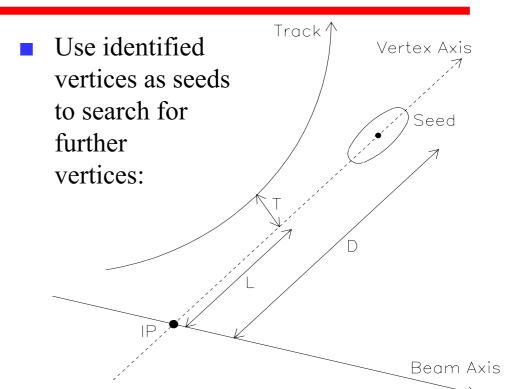
$$V(\vec{r}) = \sum_{i=0}^{N} f_i(\vec{r}) - \frac{\sum_{i=0}^{N} f_i^2(\vec{r})}{\sum_{i=0}^{N} f_i(\vec{r})}$$

Returning to e.g.:



Topological vertex reconstruction

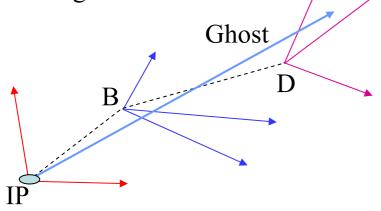
- Find "seed" maxima in V(r̄) for track-track and vertex-track pairs which are "resolvable" and for which χ² good.
- Search around seeds for true maxima in $V(\bar{r})$, including all tracks and vertex.
- Spatially resolved maxima form candidate vertices if associated with 2 or more tracks.
- Track/vertex association ambiguities are decided according to largest $V(\bar{r})$ after quality cuts.
- Vertex that includes the IP ellipsoid is primary vertex.



- Tune cut on L/D to find new vertices, possibly "one prong".
- (Additional weighting factors can be applied to favour vertices near jet core, to suppress vertices v. close to IP etc.)

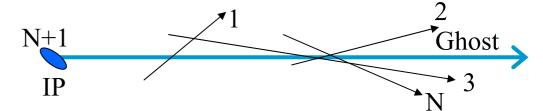
Ghost track algorithm

Complementary approach to vertex finding.

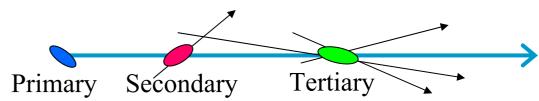


- Start with straight "ghost" track at IP along jet axis direction.
- Calculate χ^2 of DCA each track in jet to ghost track.
- Swivel ghost track in θ and φ to minimise $\sum \chi^2$ for all tracks in jet.

Now have N (tracks) + 1 (IP) initial vertex candidates.



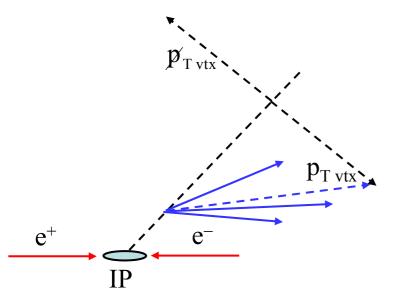
- Calculate fit probability for all trackghost track or track-IP combinations.
- If prob. high, combine objects to form vertex.
- Iterate.



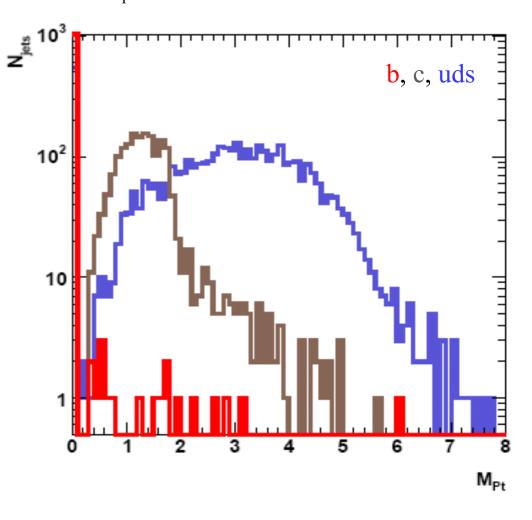
Flavour identification

- Identify variables that provide discrimination between uds/gluon and c and b jets.
- E.g., if vertex found, p_T corrected mass, m_{pT} , is useful variable,

$$m_{pT} = \sqrt{m_{vtx}^{2} + p_{Tvtx}^{2}} + p_{Tvtx}.$$



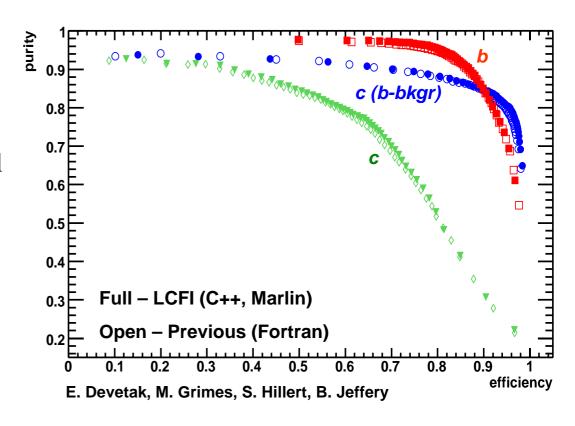
 \mathbf{m}_{pT} distributions:



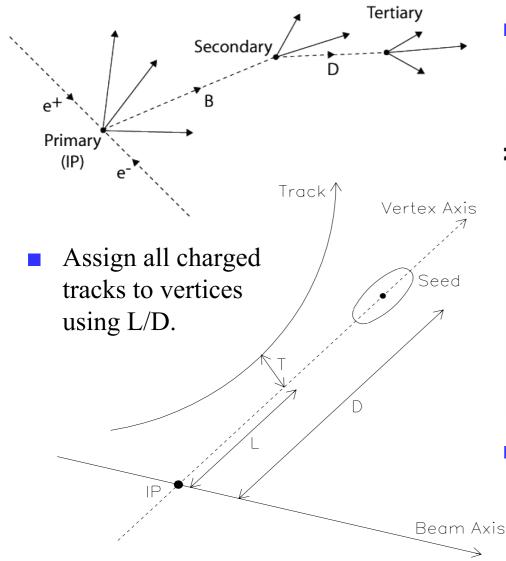
Flavour identification

- Combine variables using neural net.
- Input variables:
 - Momentum and impact parameter significance in r-φ and z of two most significant tracks.
 - Probability in r-φ and z that all tracks originate from IP.
- Additionally, if secondary vertex found:
 - \bullet m_{pT}.
 - Momentum associated with vertex.
 - Decay length.
 - Decay length significance.

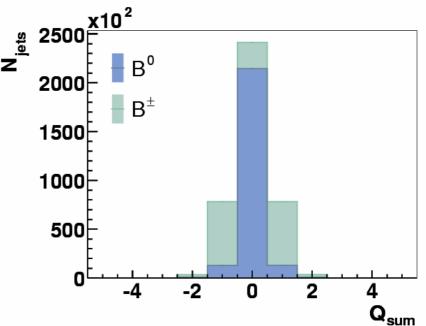
Resulting flavour identification performance (here at $\sqrt{s} = m_7$).



Quark charge identification



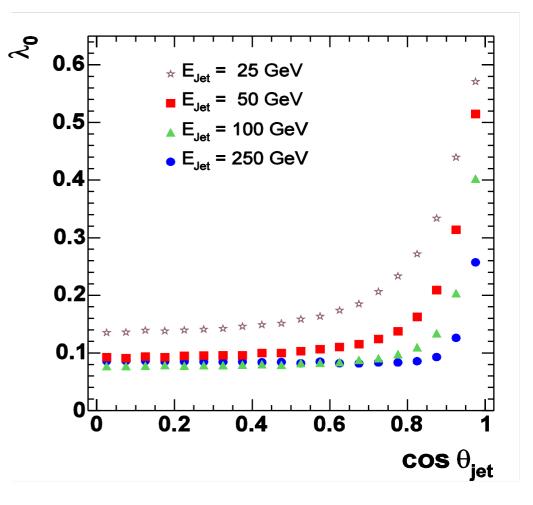
Sum charge of tracks at secondary vertex to form Q_{sum} :



Define $Q_{vtx} = sgn(Q_{sum})$.

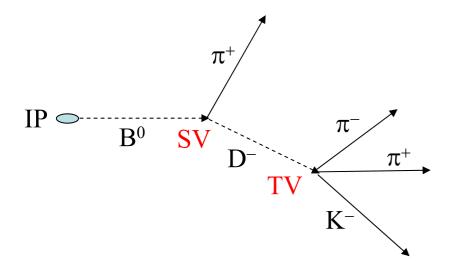
Quark charge identification

- Quantify performance in terms of λ_0 , probability of reconstructing neutral B hadron as charged.
- Degradation at large $\cos \theta$ caused by loss of tracks and multiple scattering.
- Effect stronger at lower jet energy (broader jets, lower momenta...)



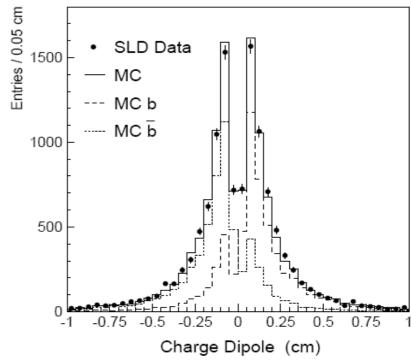
Identifying quark charge – dipole method

- Applicable when b fragments to \overline{B}^0 or \overline{b} to B^0 .
- Identify secondary and tertiary vertices and measure charge.



 $Q_{SV} = +1, Q_{TV} = -1, \Rightarrow B^{0};$ $Q_{SV} = -1, Q_{TV} = +1, \Rightarrow \overline{B}^{0}.$

Charge dipole distribution:

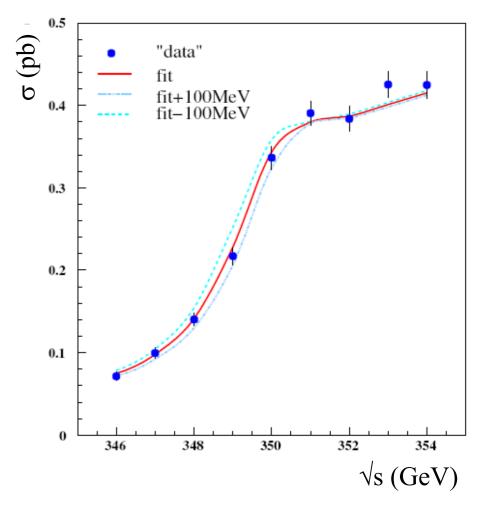


- Here particle ID can help with charge measurement.
- Of course, quark charge determination affected by mixing!

Flavour physics at the ILC – top mass determination

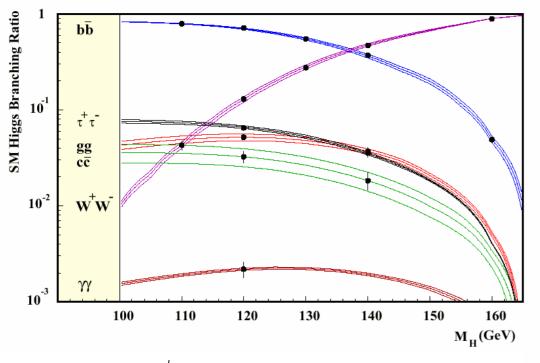
- LHC will provide m_t measurement with precision of ~ 1 GeV.
- Will remain large source of uncertainty for many SM calculations.
- Threshold scan at ILC allows determination of m_t with precision of 50...100 MeV.
- Top width also measured to 3...5%.

■ Top production cross section:



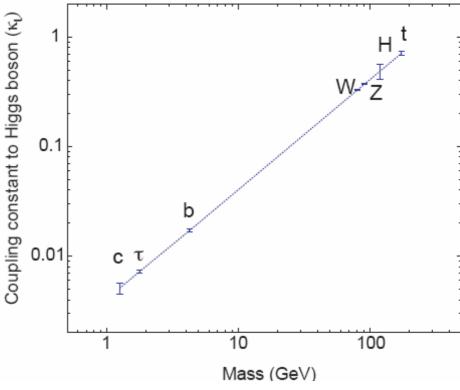
Higgs branching ratios and couplings

 Crucial to precisely measure branching ratios of Higgs boson.



Here $\sqrt{s} = 350$ GeV, integrated luminosity 500 fb⁻¹.

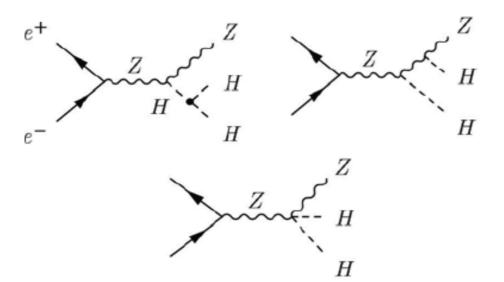
Extract couplings and check proportional to masses:



Here $m_H = 120 \text{ GeV}$, 500 Gev (700 GeV) running for HHH (t \overline{t} H) couplings.

Determining the Higgs self coupling and the top Yukawa coupling

- Quark flavour and charge identification important tool in many measurements.
- E.g. study of $e^+e^- \rightarrow HHZ$:

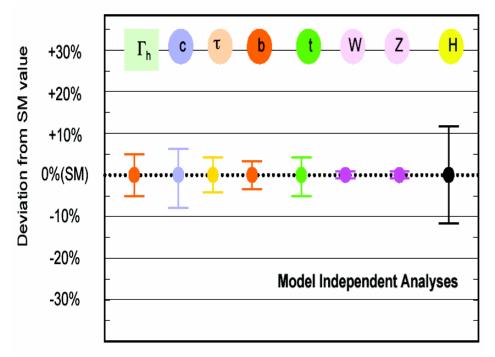


Identify four b quarks!

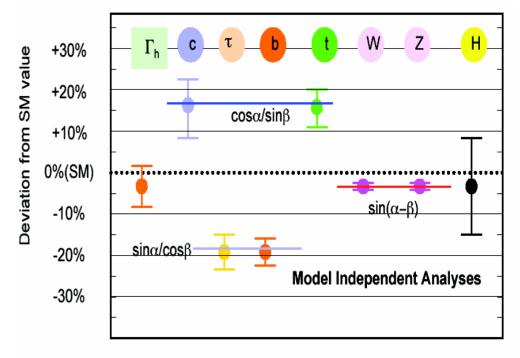
- Identification of b quarks and separation of b from \overline{b} reduces significantly combinatorial background.
- Study of this process allows determination of Higgs self-coupling.
- Similarly, b identification aids separation of e⁺e⁻ → t t H from background with typically >1000 times higher rate.
- Charge ID also helps through reduction of combinatorial BG.

Higgs and SUSY at the ILC

Such precision studies of Higgs are powerful means of checking whether Higgs properties are as expected in SM...

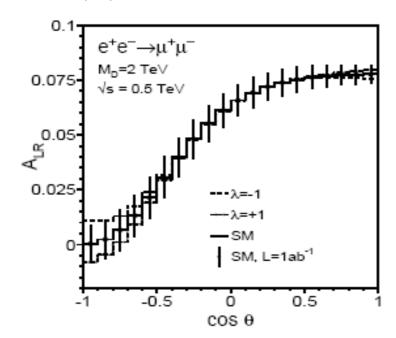


...or more like predictions from Minimal SUSY SM:

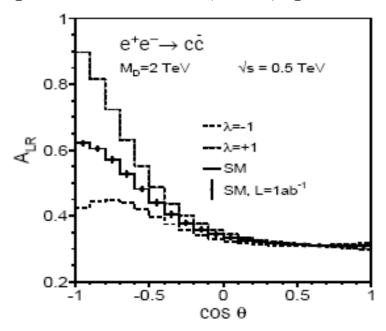


Quark charge identification

- Increases sensitivity to new physics.
- E.g. effects of large extra dimensions on $e^+e^- \rightarrow f \overline{f}$.
- Study $A_{LR} = (\sigma_L \sigma_R)/\sigma_{tot}$ as a function of $\cos \theta$.
- For $\mu^+\mu^-$, effects of ED not visible:



Changes from SM much more pronounced for c (and b) quarks:

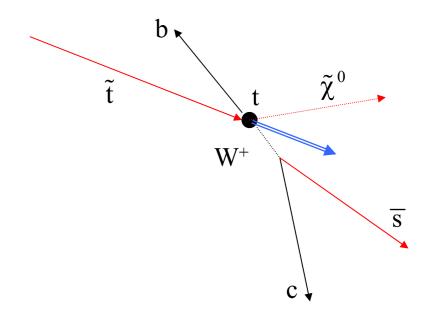


Efficient flavour identification and charge determination needed out to large $\cos \theta$.

Quark charge identification

- Provides new tools for physics studies.
- E.g. can measure top polarisation in decay $t \rightarrow W^+b$
- Top decays before hadronisation.
- Anti-strange jet has $1 \cos \theta$ distribution w.r.t. top polarisation direction.
- Distinguish between t and \overline{t} by tagging b and c jets.
- Determine quark charge for (at least) one of these jets.

Example of physics made accessible using this technique:



Determine tan β and tri-linear coupling A_t and A_b through measurements of top polarisation in t̃ and b̃ decays.

Summary

- The precision of the International Linear Collider provides an excellent complement to the discovery potential of the LHC.
- A feasible baseline ILC design has been produced.
- The challenges for experiments at the ILC are different to those at the LHC and are leading to novel approaches to detector design.
- These offer exciting opportunities for studying the physics of heavy flavours and for using heavy flavours as a tool in many investigations.
- Next steps include production of Technical Design Reports for ILC machine and detectors around 2010...
- ...hopefully leading to ILC physics results ~2020.