

# SINGLE TOP PRODUCTION AND DECAY IN MC@NLO

Patrick Motylinski

NIKHEF

DIS Munich, April 18th, 2007

# OUTLINE

- 1 INTRODUCTION
- 2 SINGLE TOP PRODUCTION
- 3 TOP DECAY AND SPIN CORRELATIONS

# MOTIVATIONS FOR SINGLE-TOP

- Direct Measurement of  $V_{tb}$
- Test the charged current coupling of the top quark
- Determine  $b$  quark density
- Single-top processes are important backgrounds for new physics ( $W'$ ,  $Z'$ , ...)

There are already several NLO calculations for single-top processes:

- B.W. Harris, E. Laenen, L. Phaf, Z. Sullivan, S. Weinzierl
- Z. Sullivan (ZTOP)
- J. Campbell, R.K. Ellis, F. Tramontano (MCFM)
- Q.-H. Cao, R. Schwienhorst, C.P. Yuan

# SINGLE TOP IN MC@NLO

(S. FRIXIONE, E. LAENEN, P. MOTYLINSKI, B. WEBBER: *hep-ph/0512250*)

- Matching NLO with Parton Showers (gives better final state description)
- MC@NLO is an event generator (!)
- Hardest  $p_T$  emission is computed exactly in agreement with NLO calculation

What had to be done:

- Compute NLO using FKS subtraction method
- Construct MC counter terms
- A LOT of testing...

# THE FKS METHOD

(S. FRIXIONE, Z. KUNSZT, A. SIGNER)

- Soft and/or collinear divergencies are dealt with through subtraction method
- Essentially partitioning phase space according to:
  - No overlapping regions
  - At most *one* collinear and *one* soft singularity in each region

How is this done for the single top process in MC@NLO?

Introducing functions with the properties:

$$1 = \mathcal{S}_i^{(\text{IN})} + \mathcal{S}_i^{(\text{OUT})} \quad \forall i$$

$$d\sigma = \int d\text{PS} |\mathcal{M}|^2 \left( \mathcal{S}_i^{(\text{IN})} + \mathcal{S}_i^{(\text{OUT})} \right)$$

In case of single top: we use  $i = 3$ .

$\mathcal{S}^{(\text{IN})}_i = 1$  where particle  $i$  becomes collinear with one of the incoming particles.  
 Similarly  $\mathcal{S}^{(\text{OUT})}_i = 1$  where particle  $i$  becomes collinear with one of the outgoing particles.

$$\mathcal{S}_3^{(\text{IN})} = \frac{(k_3 \cdot k_1)^a (k_3 \cdot k_2)^a}{(k_3 \cdot k_1)^a (k_3 \cdot k_2)^a + (k_3 \cdot p_1)^a (k_3 \cdot p_2)^a}$$

$$\mathcal{S}_3^{(\text{OUT})} = \frac{(k_3 \cdot p_1)^a (k_3 \cdot p_2)^a}{(k_3 \cdot k_1)^a (k_3 \cdot k_2)^a + (k_3 \cdot p_1)^a (k_3 \cdot p_2)^a}$$

Thus in each region of phase space the corresponding part of the cross section is treated separately.  $\Rightarrow$  In case of singularity: subtraction method applied to this part of the cross section.

## MC COUNTERTERM

- Analogous to  $BQ(x)/x$  (see next slide)
- Needed in final state (recoil jet)  
Special in the single-top case:
  - first inclusion of final state jet in MC@NLO
- Easy colour structure

How MC@NLO works:

$$\begin{aligned}
 \left( \frac{d\sigma}{d\mathcal{O}} \right) = \int_0^1 dx & \left[ \underbrace{I_{MC}(\mathcal{O}, x_M(x)) \frac{\alpha(R(x) - \overbrace{BQ(x)})}{x}}_{\text{Real emission + MC shower}} \right. \\
 & \left. + \underbrace{I_{MC}(\mathcal{O}, 1) \left( B + \alpha V - \frac{\alpha B(Q(x) - 1)}{x} \right)}_{\text{(LO \& virtual) + MC shower}} \right]
 \end{aligned}$$

parton shower @  $\mathcal{O}(\alpha)$

where

$B$  is the Born cross section

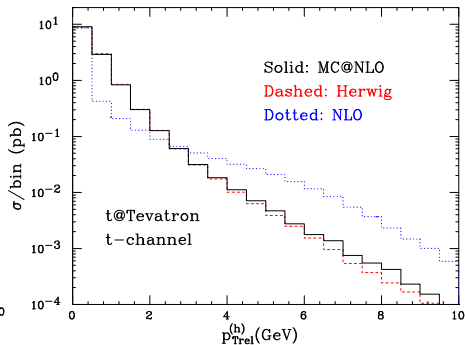
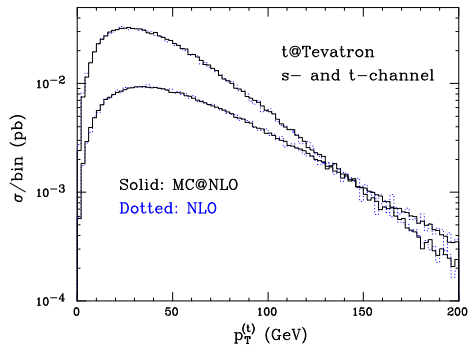
$V$  is the virtual correction

$R$  describes real corrections

$x_M$  is the maximal energy available for subsequent photon emission (via showers)

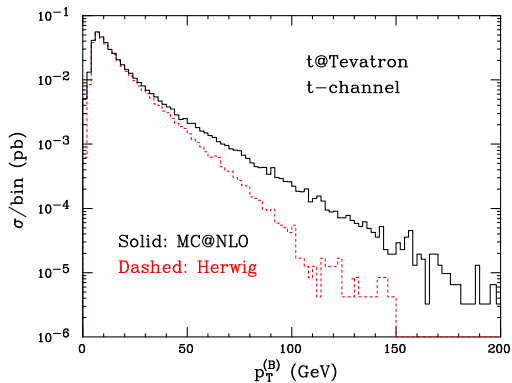
$I_{MC}$  is the interface-to-MC:  $I_{MC}(\mathcal{O}, x_M)$  is the diff. cross section obtained by running MC starting at scale  $x_M$ .





Left:  $p_T$  of the top quark

Right:  $p_T$  of particles included in the jet relative to the axis of the jet.



# TOP DECAY AND SPIN CORRELATIONS

(S. FRIXIONE, E. LAENEN, P. MOTYLINSKI, B. WEBBER: *hep-ph/0702198*)

The aim of studying spin correlations in (single) top decay is to determine the handed-ness of the EW coupling of the top. (Std. Model: purely left-handed)

Due to strong correlation with the decay products it will be possible to measure the coupling. (100% correlation between top-spin and lepton direction)

The full (leptonic) matrix element is bounded from above by an overall constant times the undecayed matrix elements, i.e.:

$$\frac{d\sigma_{l\nu b}}{dPS} \leq \frac{4g_W^4 |V_{tb}|^2 (r \cdot k_2)(p \cdot k_1)}{\left((q^2 - m_W^2)^2 + (m_W \Gamma_W)^2\right) \left((p^2 - m_t^2)^2 + (m_t \Gamma_t)^2\right)} d\sigma_t$$

How is this done in MC@NLO?

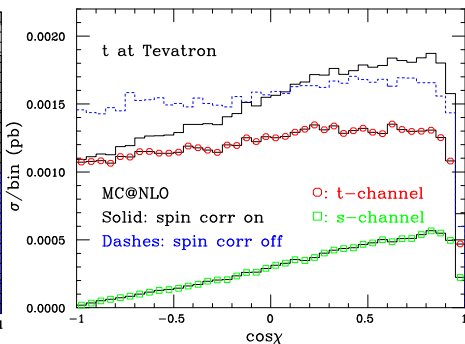
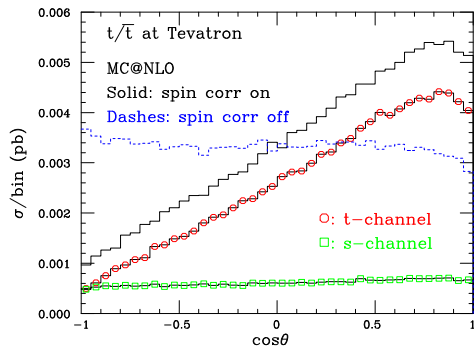
Make angular correlations precise to

- NLO only for real (hard) emissions
- LO for soft/collinear regions

In practice

- 1. Integrate undecayed matrix element via MC@NLO (creates set of events)
- 2. Generate hard events using result from step 1
- 3. For each hard event generate lepton and  $b$  quark momenta in the decay space of the top
- 4. Compute (full) lepton matrix element (using MadGraph) using momenta from 3, and compute undecayed matrix element using momenta from step 2.
- 5. Do hit-and-miss

## EFFECTS OF SPIN CORRELATIONS



where

$\theta$  is the angle between the direction of flight of the lepton (from top decay) and the hardest non- $b$  jet

$\chi$  is the angle between the direction of flight of the lepton (from top decay) and the anti-proton beam.

# “THE LAST SLIDE”

Angular correlations are now included for the single top process in MC@NLO ver. 3.3!

Angular correlations are included for

- $t\bar{t}$  production
- $VV$  production