Parton showers for non-global QCD observables

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Work in collaboration with A. Banfi and G. Corcella

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- How good is a given parton shower?
- Non-global observables and loss of angular-ordering (AO).
- Mismatch between AO and full leading-logarithmic calculations.
- HERWIG vs PYTHIA vs resummed perturbation theory.
- Concluding remarks.



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- Perturbative aspects parton shower (PS), ME corrections and matching.
- Non-perturbative aspects hadronisation and underlying event models.
- Tuning of event generators.

A lot of attention being paid to all aspects but are some details slipping the net? We shall address the first issue

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- Developments in all-order resummations have challenged understanding of soft radiation at large angles (angular ordering) even at leading (logarithmic) accuracy.
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- Observables sensitive to soft emission in limited regions include energy flow distributions, event shapes, jet distributions and many others.
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Non global observables and loss of AO





Examine differential E_t flow

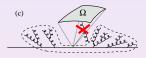
$$\frac{1}{\sigma} \frac{d\sigma}{dE_t}$$

Here $E_t = \sum_{i \in \Omega} E_{t,i}$ and Ω is interjet region e.g rapidity slice.

Leading logs are $\alpha_s^n L^n$ where $L = \ln Q/E_t$. Originate from $\omega_1 \gg \omega_2 \gg \omega_3 \cdots$ without angular-ordering. But AO a feature of MC's e.g HERWIG and old PYTHIA (before v 6.3).

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Resummation and angular ordering

We take two approaches to the problem.

- Take resummation program and force angular ordering toy model of AO.
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The leading logarithms resummed in large N_c limit by a dipole evolution Monte-Carlo.

Dasgupta and Salam 2001

Add soft gluon at scale $L' = \ln Q/\omega$ to dipole configuration C at scale L to get C'

$$P_{C'}(L') = \bar{\alpha_s}(L')\Delta_C(L, L')F_C(\theta', \phi')P_C(L)$$

$$F_c(\theta_k, \phi_k) = \sum_{\text{dipoles-ij}} \frac{2C_A (1 - \cos \theta_{ij})}{(1 - \cos \theta_{ik}) (1 - \cos \theta_{jk})}$$

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AO modification

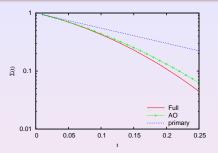
To force only configurations with AO:

$$F(\theta_{k},\phi_{k}) \rightarrow \frac{\theta\left(\cos\theta_{ik} - \cos\theta_{ij}\right)}{\left(1 - \cos\theta_{ik}\right)} + \frac{\theta\left(\cos\theta_{jk} - \cos\theta_{ij}\right)}{\left(1 - \cos\theta_{ik}\right)}$$

Then one gets $\Sigma_{AO}(\alpha_S L)$ as before.



Results

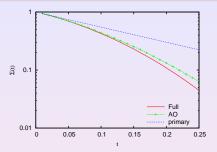


For unit rapidity slice in e^+e^- annihilation with

$$t \sim rac{lpha_{ extsf{S}}}{2\pi} \ln rac{ extsf{Q}}{ extsf{\textit{E}}_t}$$

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- HERWIG based on angular ordering, shd be close to full (large N_c) result.
- PYTHIA (old) ordering in m² and reject non AO configs, shd do worse.
- ARIADNE dipole phase space, shd have the full LL.
- PYTHIA (new) like ARIADNE?

Trick go to very high (10⁵ GeV) to kill subleading effects. Only interested in $t \sim \frac{\alpha_s}{2\pi} \ln \frac{Q}{F_s}$.



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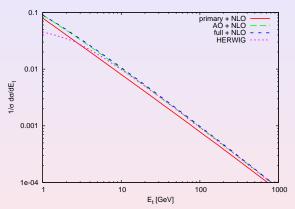
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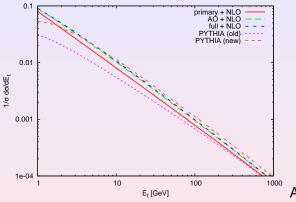
Comparison to HERWIG



effect at t = 0.15, $E_t = 10$ GeV.

Numerically 10%

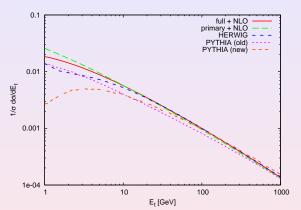
Comparison to PYTHIA



At t = 0.15

PYTHIA old deviates by 50 % from full. PYTHIA new only 7.5% off.

But for large gaps....



Problems seen with new PYTHIA at large rapidity intervals.....

- Angular ordering (implemented as in HERWIG) numerically includes a bulk of leading-log effects.
- The old PYTHIA versions (before 6.3) do not account for a large part of the leading perturbative logarithms for a number of observables.
- The new PYTHIA model works much better but we note problems at large rapidity intervals for energy flow between jets.
- Further studies are needed to understand the behaviour of the various parton showers in a quantitative fashion.
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