Dark photon oscillations in our inhomogeneous Universe





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Center for Cosmology and Particle Physics

BSM PANDEMIC October 13, 2020



Collab







Hongwan Liu (NYU/Princeton)

Papers

- 1. Caputo, Liu, SM, Ruderman, "Dark Photon Oscillations in Our Inhomogeneous Universe," PRL [arXiv:2002.05165]







Josh Ruderman (NYU)

Alfredo Urbano (INFN Trieste)

Maxim Pospelov (Minnesota)

2. Caputo, Liu, SM, Ruderman, "Modeling Dark Photon Oscillations in Our Inhomogeneous Universe," PRD [arXiv:2004.06733] 3. + Pospelov, Urbano, "Edges and Endpoints in 21-cm Observations from Resonant Photon Production," [arXiv:2009.03899]









Outline







Dark photons and resonant conversions

Dark photon oscillations with inhomogeneities

Dark photon signatures in 21-cm

Outline







Dark photons and resonant conversions

Dark photon oscillations with inhomogeneities

Dark photon signatures in 21-cm

Standard Model



Dark Sector











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(Some) Canonical portals:

Scalar
 Higgs portal

 $\lambda H^2 S^2 + \mu H^2 S$

• Fermion Neutrino portal

y(HL)N

Vector
 Kinetic mixing portal

 $\epsilon F^{\mu
u}F'_{\mu
u}$





(Some) Canonical portals:

• Scalar Higgs portal $\lambda H^2 S^2 + \mu H^2 S$

• Fermion Neutrino portal y(HL)N

• Vector Kinetic mixing portal

 $\epsilon F^{\mu
u}F'_{\mu
u}$



Kinetic mixing portal U(1)'

$$\Delta \mathscr{L} = -\frac{\epsilon}{2} F^{\mu\nu} F'_{\mu\nu} + \frac{m_{A'}^2}{2} A'^2$$

Holdom, PLB (1986)



Kinetic mixing portal U(1)'

$$\Delta \mathscr{L} = -\frac{\epsilon}{2} F^{\mu\nu} F'_{\mu\nu} + \frac{m_{A'}^2}{2} A'^2$$

$\gamma \sim \chi \sim \chi \sim \chi \sim \chi$

Holdom, PLB (1986)





 $\epsilon - m_{A'}$ plane









 $\epsilon - m_{A'}$ plane









 $\epsilon - m_{A'}$ plane: "ultralight" dark photons







 $\epsilon - m_{A'}$ plane: "ultralight" dark photons









Vacuum and resonant (in-medium) oscillations

Oscillations in vacuum $\omega^2 = k^2$



 M_{A}





Vacuum and resonant (in-medium) oscillations

Oscillations in vacuum $\omega^2 = k^2$



 M_{A}









$$m_{\gamma}^2(z) \approx \frac{4\pi\alpha n_{\rm e}(z)}{m_{\rm e}} \approx \frac{4\pi\alpha}{m_{\rm e}} x_{\rm e}(z) n_{\rm H}(z)$$

z) $n_{\rm H}(z)$

















Resonant oscillations in plasma: Landau-Zener formalism





Resonant oscillations in plasma: Landau-Zener formalism





$$\omega(z_{\rm res}) = \omega_{\rm obs}(1 + z_{\rm res})$$

$$\implies \text{Later resonances typically dominate}$$



Resonant oscillations in plasma: Landau-Zener formalism



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$$\omega(z_{\rm res}) = \omega_{\rm obs}(1 + z_{\rm res})$$

$$\implies \text{Later resonances typically dominate}$$

Similar formalism for neutrino oscillations (MSW effect)

Nonadiabatic Level Crossing in Resonant Neutrino Oscillations

Stephen J. Parke Fermi National Accelerator Laboratory, Batavia, Illinois 60510 (Received 27 May 1986)





Resonant oscillations in photon plasma





Resonant oscillations in photon plasma



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Resonant oscillations in photon plasma







CMB spectral distortions due to $\gamma \rightarrow A'$





$$I_{\omega} \left(m_{A'}, \epsilon; T_{\rm CMB} \right) = B_{\omega}$$

Blackbody spectrum
$$B_{\omega} = \frac{\omega^{3}}{2\pi^{2}} \left[\exp\left(\frac{\omega}{T_{\rm CMB}}\right) - 1 \right]^{-1}$$



CMB spectral distortions due to



$$\gamma \rightarrow A'$$

$$I_{\omega}\left(m_{A'}, \epsilon; T_{\rm CMB}\right) = B_{\omega}\left(1 - P_{\gamma \to A'}\right)$$

Blackbody spectrum γ disappearance probabil
$$B_{\omega} = \frac{\omega^{3}}{2\pi^{2}}\left[\exp\left(\frac{\omega}{T_{\rm CMB}}\right) - 1\right]^{-1} P_{\gamma \to A'} \simeq \frac{\pi\epsilon^{2}m_{A'}^{2}}{\omega\left(z_{\rm res}\right)}\left|\frac{d\ln m_{\gamma}^{2}(t_{\rm res})}{dt}\right|$$







CMB spectral distortions due to $\gamma \rightarrow A'$





$$I_{\omega}\left(m_{A'}, \boldsymbol{\epsilon}; T_{\text{CMB}}\right) = B_{\omega}\left(1 - P_{\gamma \to A'}\right)$$

Blackbody spectrum γ disappearance probabil
$$B_{\omega} = \frac{\omega^{3}}{2\pi^{2}}\left[\exp\left(\frac{\omega}{T_{\text{CMB}}}\right) - 1\right]^{-1} P_{\gamma \to A'} \simeq \frac{\pi \boldsymbol{\epsilon}^{2} m_{A'}^{2}}{\omega\left(z_{\text{res}}\right)}\left|\frac{d\ln m_{\gamma}^{2}(\boldsymbol{\epsilon})}{dt}\right|$$





$\epsilon - m_{A'}$ constraints from COBE/FIRAS



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Mirizzi, Redondo, Sigl [0901.0014]





Outline







Dark photons and resonant conversions

Dark photon oscillations with inhomogeneities

Dark photon signatures in 21-cm







$$m_{\gamma}^2(z) \approx \frac{4\pi \alpha \, \overline{n_{\rm e}(z)}}{m_{\rm e}}$$











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Computing conversion probability

Analytic approach

$$P_{\gamma \to A'} \simeq \pi \epsilon^2 m_{A'}^2 \left\langle \sum_i \frac{1}{\omega_i \left(z_{\text{res},i} \right)} \left| \frac{d \ln m_{\gamma}^2(t)}{dt} \right|_{z=z_{\text{res},i}}^{-1} \right\rangle$$

Compute average of stochastic process


Computing conversion probability





Computing conversion probability





Numerical approach



Bondarenko, Pradler, Sokolenko [2002.08942] Garcia et al [2003.10465]



1. Conversion probability along line of sight

$$\frac{\mathrm{d}P_{\gamma \to A'}}{\mathrm{d}t} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(t)} \delta_{\mathrm{D}} \left(m_{\gamma}^2(t) - m_{A'}^2 \right) m_{\gamma}^2(t)$$



1. Conversion probability along line of sight

$$\frac{\mathrm{d}P_{\gamma \to A'}}{\mathrm{d}t} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(t)} \delta_{\mathrm{D}} \left(m_{\gamma}^2(t) - m_{A'}^2 \right) m_{\gamma}^2(t)$$

2. Average conversion probability along l.o.s., weighted by PDF $f(m_{\gamma}^2)$

$$\frac{\mathrm{d}\left\langle P_{\gamma \to A'} \right\rangle}{\mathrm{d}z} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(z)} \left| \frac{\mathrm{d}t}{\mathrm{d}z} \right| \int \mathrm{d}m_{\gamma}^2 f\left(m_{\gamma}^2\right) \delta_{\mathrm{D}}\left(m_{\gamma}^2 - m_{A'}^2\right) dz$$

 m_{γ}^2



1. Conversion probability along line of sight

$$\frac{\mathrm{d}P_{\gamma \to A'}}{\mathrm{d}t} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(t)} \delta_{\mathrm{D}} \left(m_{\gamma}^2(t) - m_{A'}^2 \right) m_{\gamma}^2(t)$$

2. Average conversion probability along l.o.s., weighted by PDF $f(m_{\gamma}^2)$

$$\frac{\mathrm{d}\left\langle P_{\gamma \to A'} \right\rangle}{\mathrm{d}z} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(z)} \left| \frac{\mathrm{d}t}{\mathrm{d}z} \right| \int \mathrm{d}m_{\gamma}^2 f\left(m_{\gamma}^2\right) \delta_{\mathrm{D}}\left(m_{\gamma}^2 - m_{A'}^2\right)$$

3. Enforce resonance condition $m_{\gamma}^2 = m_{A'}^2$

$$\frac{\mathrm{d}\left\langle P_{\gamma \to A'} \right\rangle}{\mathrm{d}z} = \frac{\pi m_{A'}^4 \epsilon^2}{\omega(z)} \left| \frac{\mathrm{d}t}{\mathrm{d}z} \right| f\left(m_{\gamma}^2; z\right)$$

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 m_{γ}^2



1. Conversion probability along line of sight

$$\frac{\mathrm{d}P_{\gamma \to A'}}{\mathrm{d}t} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(t)} \delta_{\mathrm{D}} \left(m_{\gamma}^2(t) - m_{A'}^2 \right) m_{\gamma}^2(t)$$

2. Average conversion probability along l.o.s., weighted by PDF $f(m_{\gamma}^2)$

$$\frac{\mathrm{d}\left\langle P_{\gamma \to A'} \right\rangle}{\mathrm{d}z} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(z)} \left| \frac{\mathrm{d}t}{\mathrm{d}z} \right| \int \mathrm{d}m_{\gamma}^2 f\left(m_{\gamma}^2\right) \delta_{\mathrm{D}}\left(m_{\gamma}^2 - m_{A'}^2\right)$$

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m_{ν}^2

Rice's Formula (1944)

Mathematical Analysis of Random Noise By S. O. RICE

INTRODUCTION

HIS paper deals with the mathematical analysis of noise obtained by passing random noise through physical devices. The random noise





PDF of plasma mass fluctuations

 $f(m_{\gamma}^2) = \frac{\mathrm{d}\delta_{\mathrm{e}}}{\mathrm{d}m_{\gamma}^2}\mathscr{P}$

$$m_{\gamma}^{2}(\overrightarrow{x},z) = \overline{m_{\gamma}^{2}}(z) \left(1 + \delta_{e}(\overrightarrow{x},z)\right)$$

$$\mathcal{P}(\delta_{\rm e}) = \frac{\mathcal{P}(\delta_{\rm e})}{\overline{m_{\gamma}^2}}$$



PDF of plasma mass fluctuations

$$f(m_{\gamma}^2) = \frac{\mathrm{d}\delta_{\mathrm{e}}}{\mathrm{d}m_{\gamma}^2} \mathscr{P}(\delta_{\mathrm{e}}) = \frac{\mathscr{P}(\delta_{\mathrm{e}})}{\overline{m_{\gamma}^2}}$$

Electron and baryon fluctuations

$$\bar{n}_{e}(1+\delta_{e}) = \bar{x}_{e}(1+\delta_{x_{e}}) \bar{n}_{H}(1+\delta_{b})$$
$$\implies \delta_{e} = \delta_{b} + \delta_{x_{e}} + \delta_{x_{e}}\delta_{b}$$

If
$$\delta_{x_e} \ll \delta_b \implies \delta_e \approx \delta_b$$

$$m_{\gamma}^{2}(\overrightarrow{x},z) = \overline{m_{\gamma}^{2}}(z) \left(1 + \delta_{e}(\overrightarrow{x},z)\right)$$



PDF of plasma mass fluctuations

$$f(m_{\gamma}^2) = \frac{\mathrm{d}\delta_{\mathrm{e}}}{\mathrm{d}m_{\gamma}^2}\mathcal{G}$$

 $\mathscr{P}\left(\delta_{\mathsf{b}} \right)$

Electron and baryon fluctuations

$$\bar{n}_{e}(1+\delta_{e}) = \bar{x}_{e}(1+\delta_{x_{e}}) \bar{n}_{H}(1+\delta_{b})$$
$$\implies \delta_{e} = \delta_{b} + \delta_{x_{e}} + \delta_{x_{e}}\delta_{b}$$
$$\text{If } \delta_{x_{e}} \ll \delta_{b} \implies \delta_{e} \approx \delta_{b}$$

$$m_{\gamma}^{2}(\overrightarrow{x},z) = \overline{m_{\gamma}^{2}}(z) \left(1 + \delta_{e}(\overrightarrow{x},z)\right)$$







PDF of plasma mass fluctuations: Gaussian toy example





PDF of plasma mass fluctuations: Gaussian toy example





PDF of plasma mass fluctuations: Gaussian toy example



Fluctuations at late times highly <u>non-Gaussian</u>, <u>non-linear</u>



Log-normal PDF

$$\mathscr{P}_{\rm LN}\left(\boldsymbol{\delta}_{\rm b};z\right) = \frac{\left(1+\boldsymbol{\delta}_{\rm b}\right)^{-1}}{\sqrt{2\pi\Sigma^{2}(z)}} \exp\left(-\frac{\left[\ln\left(1+\boldsymbol{\delta}_{\rm b}\right)+\Sigma^{2}(z)/2\right]^{2}}{2\Sigma^{2}(z)}\right)$$



Log-normal PDF

$$\mathscr{P}_{\rm LN}\left(\boldsymbol{\delta}_{\rm b};z\right) = \frac{\left(1+\boldsymbol{\delta}_{\rm b}\right)^{-1}}{\sqrt{2\pi\Sigma^{2}(z)}} \exp\left(-\frac{\left[\ln\left(1+\boldsymbol{\delta}_{\rm b}\right)+\Sigma^{2}(z)/2\right]^{2}}{2\Sigma^{2}(z)}\right)$$





The non-linear baryonic power spectrum





The non-linear baryonic power spectrum



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Non-linear baryonic power spectra from hydrodynamical simulations:

- Illustris
- IllustrisTNG
- BAHAMAS
- EAGLE

Foreman et al [1910.03597] van Daalen et al [1906.00968]



Alternative PDF prescriptions

Log-normal PDF

Log-normal PDF with nonlinear baryon power spectrum

$$\mathscr{P}_{\rm LN}\left(\boldsymbol{\delta}_{\rm b};z\right) = \frac{\left(1+\boldsymbol{\delta}_{\rm b}\right)^{-1}}{\sqrt{2\pi\Sigma^{2}(z)}} \exp\left(-\frac{\left[\ln\left(1+\boldsymbol{\delta}_{\rm b}\right)+\Sigma^{2}(z)/2\right]^{2}}{2\Sigma^{2}(z)}\right)$$



$\mathscr{P}_{\rm LN}\left(\boldsymbol{\delta}_{\rm b}; z\right) = \frac{\left(1 + \boldsymbol{\delta}_{\rm b}\right)^{-1}}{\sqrt{2\pi \Sigma^2(z)}} \exp_{\boldsymbol{\lambda}}$

Alternative PDF prescriptions

Log-normal PDF Log-normal PDF with nonlinear

baryon power spectrum

"Analytic" PDF Non-linear spherical collapse of linear matter field

Ivanov, Kaurov, Sibiryakov [1811.07913]

$$\mathscr{P}_{\rm an}\left(\boldsymbol{\delta}_{\rm b};z\right) = \frac{\hat{C}\left(\boldsymbol{\delta}_{\rm b}\right)}{\sqrt{2\pi\sigma_{R_{\rm J}}^2(z)}} \exp\left[\frac{-\frac{F^2\left(\boldsymbol{\delta}_{\rm b}\right)}{2\sigma_{R_{\rm J}}^2(z)}\right]$$

$$\operatorname{xp}\left(\frac{\left[\ln\left(1+\delta_{b}\right)+\Sigma^{2}(z)/2\right]^{2}}{2\Sigma^{2}(z)}\right)$$



Cosmic voids PDF

Log-normal PDF

"Analytic" PDF

linear matter field

baryon power spectrum

Log-normal PDF with nonlinear

PDF of matter underdensities

Ivanov, Kaurov, Sibiryakov [1811.07913]

Adermann et al [1703.04885, 1807.02938]

 $\mathscr{P}_{\text{voids}}\left(\delta_{\mathbf{b}};z\right) \sim \text{from simulations}$

"Analytic" PDF Non-linear spherical collapse of $\mathscr{P}_{an}\left(\delta_{b};z\right) = \frac{\widehat{C}\left(\delta_{b}\right)}{\sqrt{2\pi\sigma_{R_{J}}^{2}(z)}} \exp \frac{1}{\sqrt{2\pi\sigma_{R_{J}}^{2}(z)}}$

$\mathscr{P}_{\rm LN}\left(\delta_{\rm b};z\right) = \frac{\left(1+\delta_{\rm b}\right)^{-1}}{\sqrt{2\pi\Sigma^2(z)}}\,{\rm ex}$

Alternative PDF prescriptions

$$xp\left(\frac{\left[\ln\left(1+\delta_{b}\right)+\Sigma^{2}(z)/2\right]^{2}}{2\Sigma^{2}(z)}\right)$$

$$p\left[-\frac{F^2\left(\delta_{\rm b}\right)}{2\sigma_{R_{\rm J}}^2(z)}\right]$$



Cosmic voids PDF

Log-normal PDF

"Analytic" PDF

linear matter field

baryon power spectrum

Log-normal PDF with nonlinear

Non-linear spherical collapse of

Ivanov, Kaurov, Sibiryakov [1811.07913]

PDF of matter underdensities

Adermann et al [1703.04885, 1807.02938]

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 $\mathscr{P}_{\text{voids}}\left(\delta_{\rm b};z\right) \sim \text{from simulations}$

$\mathscr{P}_{\rm an}\left(\boldsymbol{\delta}_{\rm b};z\right) = \frac{\widetilde{C}\left(\boldsymbol{\delta}_{\rm b}\right)}{\sqrt{2\pi\sigma_{R_{\rm J}}^2(z)}} \exp\left[-\frac{F^2\left(\boldsymbol{\delta}_{\rm b}\right)}{2\sigma_{R_{\rm J}}^2(z)}\right]$

$\mathscr{P}_{\rm LN}\left(\boldsymbol{\delta}_{\rm b};z\right) = \frac{\left(1+\boldsymbol{\delta}_{\rm b}\right)^{-1}}{\sqrt{2\pi\Sigma^2(z)}}\,{\rm ex}$

Alternative PDF prescriptions

$$xp\left(\frac{\left[\ln\left(1+\delta_{b}\right)+\Sigma^{2}(z)/2\right]^{2}}{2\Sigma^{2}(z)}\right)$$









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 $\gamma \rightarrow A'$

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 $\gamma \rightarrow A'$





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$\gamma \rightarrow A'$





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 $\gamma \to A'$

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Note: Additional constraints apply when A' is the DM; see papers, also Witte et al [2003.13698]

$\gamma \to A'$











































Stronger constraints are possible with a better understanding of larger over/under-densities

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Outline







Dark photons and resonant conversions

Dark photon oscillations with inhomogeneities

Dark photon signatures in 21-cm

Hyperfine splitting of hydrogen 1s:









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Neutral hydrogen (HI) In the intergalactic medium (IGM)



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 $\omega = \omega_{21}$

 T_{γ}

CMB Source of 21-cm photons



Spin temperature $T_{\rm s}$

$$\frac{n_{\rm t}}{n_{\rm s}} \equiv 3 \exp\left(-\frac{\omega_{21}}{T_{\rm s}}\right)$$



Observation



 $\checkmark T_{b}^{21}$

Neutral hydrogen (HI) In the intergalactic medium (IGM)



 $\omega_{\rm obs} = \frac{\omega_{21}}{1+z}$

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 $\omega = \omega_{21}$

 T_{γ}

CMB Source of 21-cm photons



Spin temperature T_s

$$\frac{n_{\rm t}}{n_{\rm s}} \equiv 3 \exp\left(-\frac{\omega_{21}}{T_{\rm s}}\right)$$



Observation



Neutral hydrogen (HI) In the intergalactic medium (IGM)



$$\omega_{\rm obs} = \frac{\omega_{21}}{1+z}$$

Brightness temperature ΔT_b^{21} $\Delta T_b^{21} \propto x_{\rm HI} \left(1 - \frac{T_{\gamma}}{T_{\rm s}} \right)$

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 $\omega = \omega_{21}$

 T_{γ}

CMB Source of 21-cm photons



Spin temperature T_s

$$\frac{n_{\rm t}}{n_{\rm s}} \equiv 3 \exp\left(-\frac{\omega_{21}}{T_{\rm s}}\right)$$



Observation

$\checkmark T_{b}^{21}$

Neutral hydrogen (HI) In the intergalactic medium (IGM)



$$\omega_{\rm obs} = \frac{\omega_{21}}{1+z}$$

Brightness temperature
$$\Delta T_b^{21}$$

$$\Delta T_b^{21} \propto x_{\rm HI} \left(1 - \frac{T_{\gamma}}{T_{\rm S}} \right)$$

 $T_{\rm s} > T_{\gamma} \implies \Delta T_{\rm b}^{21} > 0$ (emission) $T_{\rm s} < T_{\gamma} \implies \Delta T_{\rm b}^{21} < 0$ (absorption)

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 $\omega = \omega_{21}$

 T_{γ}

CMB Source of 21-cm photons



Spin temperature T_s

$$\frac{n_{\rm t}}{n_{\rm s}} \equiv 3 \exp\left(-\frac{\omega_{21}}{T_{\rm s}}\right)$$



Observations of the (global) 21-cm signal





Observations of the (global) 21-cm signal





EDGES 21-cm signal (Experiment to Detect the Global EoR Signature)





Bowman et al, Nature (2018)

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EDGES 21-cm signal (Experiment to Detect the Global EoR Signature)





Bowman et al, Nature (2018)

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Extra 21-cm absorption with new physics

$\Delta T_b^{21} \propto 1 - \frac{T_{\gamma}}{T_s}$

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Extra 21-cm absorption with new physics



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Cool baryons

Muñoz & Loeb [1802.10094] Falkowski & Petraki [1803.10096] Barkana [1803.06698] Barkana et al [1803.03091] Berlin et al [1803.02804] Liu et al [1908.06986]



Extra 21-cm absorption with new physics



Heat photons

Pospelov et al [1803.07048] Moroi, Nakayama, Tang [1804.10378] Choi, Seong, Yun [1911.00532]



Cool baryons

Muñoz & Loeb [1802.10094] Falkowski & Petraki [1803.10096] Barkana [1803.06698] Barkana et al [1803.03091] Berlin et al [1803.02804] Liu et al [1908.06986]



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Heating the CMB with dark photon oscillations

Pospelov, Pradler, Ruderman, Urbano [1803.07048]

1. Decay of long-lived particle a (DM) to A'







Heating the CMB with dark photon oscillations

Pospelov, Pradler, Ruderman, Urbano [1803.07048]

1. Decay of long-lived particle a (DM) to A'



2. Resonant conversion of A' to γ

$$A' \sim 10^{-13} - 10^{-9} \,\mathrm{eV}$$





Features in 21-cm from dark photons









Features in 21-cm from dark photons









Features in 21-cm from dark photons









Benchmark 1: signal during cosmic dawn

$$m_{A'} = 10^{-11} \text{ eV}$$

$$m_a = 5 \times 10^{-4} \text{ eV}$$

$$z_{edge} \simeq 660$$

$$z_{end} \simeq 15$$

$$\epsilon = 5 \times 10^{-8}$$



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Benchmark 1: signal during cosmic dawn

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21-cm from the moon?



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Benchmark 2: signatures during dark ages

 $m_{A'} = 3 \times 10^{-13} \,\mathrm{eV}$ $z_{\rm edge} \simeq 95$ $m_a = 2 \times 10^{-5} \,\mathrm{eV}$ $z_{\rm end} \simeq 65$ $\epsilon = 5 \times 10^{-10}$







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Benchmark 2: signatures during dark ages

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Work in progress

Implications for axion-like particles



Mirizzi, Redondo, Sigl [0905.4865]



Work in progress

Implications for axion-like particles



Mirizzi, Redondo, Sigl [0905.4865]

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Effect on CMB and 21-cm anisotropy







Conclusions



$$\prime \rightarrow A'$$

Inhomogeneities can have significant observable effects for resonant photon-to-dark photon conversions



$$A' \to \gamma$$

Resonant dark photon-to-photon conversion can leave striking signatures in 21-cm observations

Conclusions



$$\prime \rightarrow A'$$

Inhomogeneities can have significant observable effects for resonant photon-to-dark photon conversions



$$A' \to \gamma$$

Resonant dark photon-to-photon conversion can leave striking signatures in 21-cm observations

More information

Papers:

- "Dark Photon Oscillations in Our Inhomogeneous" Universe," Caputo, Liu, SM, Ruderman [2002.05165]
- "Modeling Dark Photon Oscillations in Our Inhomogeneous Universe," Caputo, Liu, SM, Ruderman [2004.06733]
- "Edges and Endpoints in 21-cm Observations from Resonant Photon Production," + Pospelov, Urbano [2009.03899]

Codes:

- <u>https://github.com/smsharma/dark-photons-perturbations</u>
- https://github.com/smsharma/edges-endpoints-21cm
- <u>https://github.com/smsharma/twentyone-global</u>



Additional slides

$\epsilon - m_{A'}$ constraints on dark photon dark matter*

Additional constraints apply when the A' is the dark matter

McDermott & Witte [1911.05086]

- Anomalous heating of the IGM during He II reionization is constrained to be < 1 eV
- This constrains the energy injected due to $A' \rightarrow \gamma$ during $2 \leq z \leq 6$

See also Witte et al [2003.13698]



*Assumes energy is uniformly distributed among baryons







Local heating prescription for $A' \text{ DM} \rightarrow \gamma$





Comparison with numerical approach



Bondarenko, Pradler, Sokolenko [2002.08942] Garcia et al [2003.10465]



Resonant oscillations in plasma: Landau-Zener formalism

Non-adiabatic level crossings for two-level quantum system





Resonant oscillations in plasma: Landau-Zener formalism

Non-adiabatic level crossings for two-level quantum system



Similar formalism for neutrino oscillations (MSW effect)

Nonadiabatic Level Crossing in Resonant Neutrino Oscillations

Stephen J. Parke Fermi National Accelerator Laboratory, Batavia, Illinois 60510 (Received 27 May 1986)





Resonant oscillations in plasma: Landau-Zener formalism







$P_{\gamma \rightarrow A'}$ in an inhomogeneous plasma

Conversion along a line of sight

$$P_{\gamma \to A'} \simeq \pi \epsilon^2 m_{A'}^2 \sum_{i} \frac{1}{\omega_i \left(z_{\text{res},i} \right)} \left| \frac{d \ln m_{\gamma}^2(t)}{dt} \right|_{z=z_{\text{res},i}}^{-1}$$

Perturbations in the Photon Plasma Mass



Averaged conversion $P_{\gamma \to A'} \simeq \pi \epsilon^2 m_{A'}^2 \left\langle \sum_i \frac{1}{\omega_i (z_{\text{res},i})} \right|^2$ $d\ln m_{\gamma}^2(t)$ d*t* $z = z_{\rm res.}$

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Resonant oscillations in photon plasma











Photon plasma mass variance



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PDF snapshots



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Baryon/electron fluctuations









Simulation-inferred baryon power spectra



Simulation vs analytics comparison





Conversion probabilities





Dependence of total probability on k_{\max}





Redshift and PDF systematics





Constraints on EDGES explanation of model

Pospelov, Pradler, Ruderman, Urbano [1803.07048]

1. Decay of long-lived particle a (DM) to A'





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Constraints on EDGES explanation of model

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1. Decay of long-lived particle a (DM) to A'



2. Resonant conversion of A' to γ

$$A' \sim X \sim Y \sim Y \sim E$$



Power-law injection of CMB photons



$$T_{\gamma}^{21} = T_{\text{CMB}}(1+z) \left[1 + f_{\text{r}}A_{\text{r}} \left(\frac{\nu_{21}/(1+z)}{78\text{MHz}} \right) \right]$$



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- Collisional coupling $T_{\rm k} \approx T_{\rm s}$
- Compton scattering effective, $T_{\rm k} \approx T_{\gamma}$
- $\Delta T_{\rm b}^{21} \approx 0$

- Gas adiabatically cools as $(1 + z)^{-2}$
- $T_{\rm s} < T_{\gamma} \implies \Delta T_{\rm b}^{21} < 0$ (absorption)
- Collisional coupling becomes ineffective

• $\Delta T_{\rm b}^{21} \approx 0$









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21-cm temperature evolution



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 $\Delta T_{\rm b}^{21} \propto x_{\rm HI} \left(1 - \frac{T_{\gamma}}{T_{\rm s}} \right)$

- First stars produce Ly- α photons, couples $T_{\rm s}$ to $T_{\rm k}$
- $\Delta T_{\rm b}^{21} < 0$ (absorption feature)

- X-ray sources (e.g. quasars) heat the gas
- $T_{\rm s} > T_{\gamma}$
- $\Delta T_{\rm b}^{21} > 0$ (emission feature)



21-cm temperature evolution



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21-cm temperature evolution under perfect W-F coupling







21-cm public code

https://github.com/smsharma/twentyone-global

- Lightweight code to model global 21-cm signal
- Simple models of astrophysical (UV/X-ray) emission
- Easy to add extra sources of photons

•	smsharma Updated readme	4485578 on Sep 9	🕚 14 commits
	data	Basic repo structure and code	3 months ago
	notebooks	Updated to arXiv version	2 months ago
	twentyone	Merge branch 'master' of https://github.com/smsharma/twentyone	2 months ago
D	.gitignore	Basic repo structure and code	3 months ago
D	LICENSE	Initial commit	3 months ago
D	README.md	Updated readme	2 months ago
ß	environment.yml	Updated to arXiv version	2 months ago

README.md

twentyone-global

Simplified framework for modeling the global 21-cm absorption signal, with a focus on studying the implications of non-standard 21-cm CMB temperature evolution. For details about the modeling, see 2009.03899.



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