

Single cut spinor integration and tadpole coefficients

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In collaboration with R. Britto

Outline

- Introduction
- Double cut & Tadpoles
- Single cut & Tadpoles
- Conclusions

Introduction

$$\mathcal{A}_{1\text{-loop}} = \sum_i \left(\text{Diagram } i \right) d_i + \sum_j \left(\text{Diagram } j \right) c_j + \sum_k \left(\text{Diagram } k \right) b_k + \sum_m \left(\text{Diagram } m \right) a_m + \mathcal{R}$$

Computation of the coefficients & rational part

- Passarino-Veltman reduction, [Passarino, Veltman '79]
- singular limits of $\mathcal{A}_{1\text{-loop}}$ + on-shell formalism, [Bern, Dixon, Dunbar, Kosower '94,'95]
 - Generalized unitarity multiple-cuts
 - [Bern, Dixon, Kosower '98] [Bern *et al.* '05][Britto, Cachazo, Feng '05] [Mastrolia, '07] [Forde' 07] [Kilgore '08]
 - [Ellis, Giele, Kunszt '08] [Badger '09] [Ellis *et al.* '08] [Berger *et al.* '08] [Ossola, Papadopoulos, Pittau '07]
 - Double cut & spinor integration
 - [Britto *et al.* '07] [Britto, Feng, Mastrolia '06][Anastasiou *et al.* '07,'07]
 - ↪ closed formulae for d' s c' s and b' s
 - ↪ cannot compute a' s

Introduction

$$\mathcal{A}_{1\text{-loop}} = \sum_i \left[\text{Diagram } d_i \right] d_i + \sum_j \left[\text{Diagram } c_j \right] c_j + \sum_k \left[\text{Diagram } b_k \right] b_k + \sum_m \left[\text{Diagram } a_m \right] a_m + \mathcal{R}$$

This talk: tadpole coefficients using single cut & spinor integration

● Tricky business ...

- n -leg $\mathcal{A}_{1\text{-loop}} \leftrightarrow (n + 2)$ -leg tree level amplitude,
- tree-level recursion relations break down, [Hall '07][Schwinn, Weinzierl '07]
- gauge dependence arises, [Ellis *et al.* '08]
- connection tadpoles & massless bubbles.

● ... But promising!

- single cut to fully reconstruct $\mathcal{A}_{1\text{-loop}}$. [Catani *et al.* '08] [Glover, Williams '08]

Alternative approach: find a ' s from the UV & IR structure of $\mathcal{A}_{1\text{-loop}}$.

[Bern, Morgan '96] [Badger '08][Badger, Sattler, Yundin '10]

Auxiliary propagator method

Tadpole coefficients from the double cut

Auxiliary propagator method

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

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We want $a(0)$, the coefficient of $A_0(m_0^2)$.

General idea: [Britto, Feng '09] 1) apply the OPP decomposition,

$$\begin{aligned} \mathcal{I}_{2,2} &= \frac{a(0)}{D_0} + \frac{\tilde{a}(q; 0)}{D_0} + \frac{b(0, 1)}{D_0 D_1} \\ &+ \frac{\tilde{b}(q; 0, 1)}{D_0 D_1} + \dots \end{aligned}$$

Auxiliary propagator method

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

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$$\begin{aligned} \mathcal{I}_{2,2} \frac{1}{D_K} &= \frac{a(0)}{D_0 D_K} + \frac{\tilde{a}(q; 0)}{D_0 D_K} + \frac{b(0, 1)}{D_0 D_1 D_K} \\ &+ \frac{\tilde{b}(q; 0, 1)}{D_0 D_1 D_K} + \dots \end{aligned}$$

$$\mathcal{I}_{2,2} \frac{1}{D_K} = \frac{a_K(0)}{D_0} + \frac{b_K(0, K)}{D_0 D_K} + \frac{c_K(0, 1, K)}{D_0 D_1 D_K} + \dots$$

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$$\Delta_2 \left[\mathcal{I}_{2,2} \frac{1}{D_K} \right] = \Delta_2 \left[\frac{a_K(0)}{D_0} \right] + \Delta_2 \left[\frac{b_K(0, K)}{D_0 D_K} \right] + \Delta_2 \left[\frac{c_K(0, 1, K)}{D_0 D_1 D_K} \right] + \dots$$

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$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

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- 2) apply the double cut,
- 3) apply the decoupling conditions:
 $K \cdot K_1 = 0, \quad K^2 + m_K^2 - m_0^2 = 0.$

$$\begin{aligned} \Delta_2 \left[\mathcal{I}_{2,2} \frac{1}{D_K} \right] &= \Delta_2 \left[\frac{a(0)}{D_0 D_K} \right] + \Delta_2 \left[\frac{\tilde{a}(q; 0)}{D_0 D_K} \right] + \Delta_2 \left[\frac{b(0, 1)}{D_0 D_1 D_K} \right] \\ &+ \frac{f_1}{12} \Delta_2 \left[\frac{\tilde{b}_{00}(0, 1)}{D_0 D_K} \right] + \dots \end{aligned}$$

$$\Delta_2 \left[\mathcal{I}_{2,2} \frac{1}{D_K} \right] = \Delta_2 \left[\frac{a_K(0)}{D_0} \right] + \Delta_2 \left[\frac{b_K(0, K)}{D_0 D_K} \right] + \Delta_2 \left[\frac{c_K(0, 1, K)}{D_0 D_1 D_K} \right] + \dots$$

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$$a(0) = b_K(0, K) - \frac{f_1}{12} \tilde{b}_{00}(0, 1)$$

$$\tilde{b}_{00}(0, 1) = \frac{3}{K_1^2} \left[b(0, 1)|_{\mu^2} - c_K(0, 1, K)|_{\mu^2} \right]$$

$$\begin{cases} b(0, 1)|_{\mu^2} \\ c_K(0, 1, K)|_{\mu^2} \end{cases} \text{ from } \begin{cases} \mathcal{I}_{2,2} \\ \mathcal{I}_{2,2}/D_K \end{cases} \text{ using double cut in D-dimensions.}$$

Auxiliary propagator method

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

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$$a(0) = \frac{f_1}{3(K_1^2)^2} \left[K_1^2 (R_1 \cdot R_2) - 4(K_1 \cdot R_1)(K_1 \cdot R_2) \right] \quad [\text{OK with PV}]$$

Auxiliary propagator method

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

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- The auxiliary propagator method:
 - is valid in general,
 - requires few spurious terms,
 - ↪ traded with non-spurious ones . . .
 - ↪ . . . double cut in D-dimension sufficient
 - introduces extra D_K ,
 - ↪ computation is harder
 - assumes $K_1^2 \neq 0$.

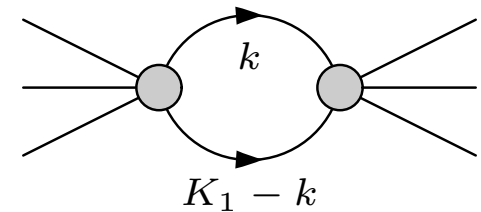
Evaluation of the double cut

How to get the double cut

Evaluation of the double cut

$$\Delta_2 [\mathcal{I}] \equiv \int d^4 k [\mathcal{I} D_0 D_1] \delta^+ (D_0) \delta^+ (D_1)$$

Definitions $D_0 = k^2 - m_0^2$; $D_1 = (k - K_1)^2 - m_1^2$;



Evaluation of the double cut

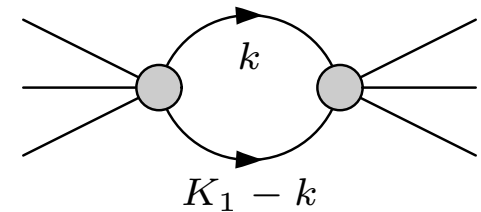
$$\Delta_2 [\mathcal{I}] = \int dt \int dz d\bar{z} \frac{[\mathcal{I} D_0 D_1]}{(1 + z\bar{z})^2}$$

Definitions $D_0 = k^2 - m_0^2$; $D_1 = (k - K_1)^2 - m_1^2$;

1) Reparametrizations

i) $k = \tilde{\ell} + \xi K_1$ [Anastasiou *et al.* '07]

ii) $\tilde{\ell}^\mu = t(p^\mu + \alpha q^\mu + z\epsilon_1^\mu + \bar{z}\epsilon_2^\mu)$ [Mastrolia '09]



Evaluation of the double cut

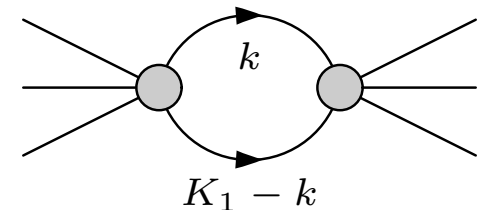
$$\Delta_2 [\mathcal{I}] = \int dt \left\{ \lim_{\Lambda \rightarrow \infty} \frac{1}{2} \int_0^{2\pi} d\alpha \Lambda e^{i\alpha} \mathcal{G}(\Lambda e^{i\alpha}, \Lambda e^{-i\alpha}) - \pi \sum_{\text{poles } z_j} \text{Res}\{\mathcal{G}(z, \bar{z}), z_j\} \right\}$$

Definitions $D_0 = k^2 - m_0^2$; $D_1 = (k - K_1)^2 - m_1^2$; $\mathcal{G}(z, \bar{z}) = \int d\bar{z} \frac{[\mathcal{I} D_0 D_1]}{(1 + z\bar{z})^2}$

1) Reparametrizations

i) $k = \tilde{\ell} + \xi K_1$ [Anastasiou et al. '07]

ii) $\tilde{\ell}^\mu = t(p^\mu + \alpha q^\mu + z\epsilon_1^\mu + \bar{z}\epsilon_2^\mu)$ [Mastrolia '09]



2) Generalized Cauchy formula

i) $z = \Lambda e^{i\alpha}$ $\bar{z} = \Lambda e^{-i\alpha}$

ii) the double cut given by the residues of \mathcal{G}

Evaluation of the single cut

What about the single cuts ?

Evaluation of the single cut

$$\Delta_1 [\mathcal{I}] \equiv \int d^4 k [\mathcal{I} D_0] \delta^+ (D_0)$$

Definitions $D_0 = k^2 - m_0^2$;

Evaluation of the single cut

$$\Delta_1 [\mathcal{I}] = \int dt \left\{ \lim_{\Lambda \rightarrow \infty} \frac{1}{2} \int_0^{2\pi} d\alpha \Lambda e^{i\alpha} \mathcal{G}(\Lambda e^{i\alpha}, \Lambda e^{-i\alpha}) - \pi \sum_{\text{poles } z_j} \text{Res}\{\mathcal{G}(z, \bar{z}), z_j\} \right\}$$

Definitions $D_0 = k^2 - m_0^2$; $\mathcal{G}(z, \bar{z}) = \frac{2K^2 t}{4} \int d\bar{z} [\mathcal{I} D_0]$;

1) Reparametrizations

i) $k = \tilde{\ell} + \xi K$ (K arbitrary)

ii) $\tilde{\ell}^\mu = t(p^\mu + \alpha q^\mu + z\epsilon_1^\mu + \bar{z}\epsilon_2^\mu)$ [Mastrolia '09]

iii) let $K^2 \rightarrow \infty$

2) Generalized Cauchy formula

i) $z = \Lambda e^{i\alpha}$ $\bar{z} = \Lambda e^{-i\alpha}$

Evaluation of the single cut

$$\Delta_1 \left[\frac{1}{D_0} \right] = \int dt \left\{ \lim_{\Lambda \rightarrow \infty} \frac{1}{2} \int_0^{2\pi} d\alpha \frac{2K^2 t}{4} \Lambda^2 - \pi \sum_{\text{poles } z_j} \text{Res} \left\{ \frac{2K^2 t}{4} \bar{z}, z_j \right\} \right\}$$

Definitions $D_0 = k^2 - m_0^2$; $\mathcal{G}(z, \bar{z}) = \frac{2K^2 t}{4} \bar{z}$;

Tadpole integral $\int d^4 k \frac{1}{D_0} \longrightarrow \mathcal{I} = \frac{1}{D_0}$

Evaluation of the single cut

$$\Delta_1 \left[\frac{1}{D_0} \right] = \lim_{\Lambda \rightarrow \infty} \frac{\pi K^2}{2} \Lambda^2 \left(\int dt t \right)$$

Definitions $D_0 = k^2 - m_0^2$; $\mathcal{G}(z, \bar{z}) = \frac{2K^2 t}{4} \bar{z}$;

Tadpole integral $\int d^4 k \frac{1}{D_0} \longrightarrow \mathcal{I} = \frac{1}{D_0}$

● Single cuts of a tadpole $\sim \Lambda^2$

● bubble, triangle, box, behave differently

e.g. for a bubble $\Delta_1[1/(D_0 D_1)] \sim \log(\alpha \Lambda^2)$

● Strategy

Find the tadpole coefficient of $\int \mathcal{I}$ selecting the $\Lambda^2 t$ -terms of $\Delta_1[\mathcal{I}]$

● Caveat: spurious terms enter!

i.e. $\int \mathcal{I} = 0 \not\Rightarrow \Delta_1[\mathcal{I}] = 0$

Single cut & Tadpoles – I

Tadpole coefficients, (next-to-) simplest example

Single cut & Tadpoles – I

$$\mathcal{I}_{2,1} = \frac{(2R \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

We want $a(0)$, the coefficient of $A_0(m_0^2)$.

1) OPP decomposition ($K_1 \perp$ to n, ℓ_7, ℓ_8)

$$\mathcal{I}_{2,1} = \frac{1}{D_0} a(0) + \dots$$

Single cut & Tadpoles – I

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- 1) OPP decomposition ($K_1 \perp$ to n, ℓ_7, ℓ_8)
- 2) Take the $\Lambda^2 t$ part of the cut (recall the strategy!)

$$\Delta_1 [\mathcal{I}_{2,1}] = \Delta_1 \left[\frac{1}{D_0} \right] a(0) + \Delta_1 [\dots]$$

Single cut & Tadpoles – I

$$\mathcal{I}_{2,1} = \frac{(2R \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

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- 1) OPP decomposition ($K_1 \perp$ to n, ℓ_7, ℓ_8)
- 2) Take the $\Lambda^2 t$ part of the cut (recall the strategy!)
- 3) Compute the single cut

$$0 = \frac{q^\mu}{(q \cdot K_1)} \left[K_{1\mu} a(0) + R_{1\mu} + (\text{terms } \perp \text{ to } K_1^\mu) \right]$$

Single cut & Tadpoles – I

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We want $a(0)$, the coefficient of $A_0(m_0^2)$.

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- 2) Take the $\Lambda^2 t$ part of the cut (recall the strategy!)
- 3) Compute the single cut
- 4) decompose $R^\mu = \alpha_1 K_1^\mu + \alpha_n n^\mu + \alpha_7 \ell_7^\mu + \alpha_8 \ell_8^\mu$

$$0 = \frac{q^\mu}{(q \cdot K_1)} \left[K_{1\mu} (a(0) + \alpha_1) \quad + \quad (\text{terms } \perp \text{ to } K_1^\mu) \right]$$

Single cut & Tadpoles – I

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$$0 = \frac{q^\mu}{(q \cdot K_1)} \left[K_{1\mu} (a(0) + \alpha_1) \quad + \quad (\text{terms } \perp \text{ to } K_1^\mu) \right]$$

$$a(0) = -\alpha_1 = -\frac{2K_1 \cdot R_1}{K_1^2}$$

Single cut & Tadpoles – I

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m_0^2, \quad D_1 = (k - K_1)^2 - m_1^2, \quad f_1 = K_1^2 + m_0^2 - m_1^2]$$

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- 3) Compute the single cut
- 4) decompose $R_1^\mu R_2^\nu = \alpha_1 K_1^\mu K_2^\mu + \alpha_n n^\mu n^\nu + \dots$

$$0 = \frac{q^\mu q^\nu}{(q \cdot K_1)^2} \left[\left(\alpha_1 + \frac{a(0)}{f_1} + \frac{\tilde{b}_{00}}{3} \right) K_{1\mu} K_{1\nu} - \left(\tilde{b}_{00} - \alpha_n \right) n_\mu n_\nu + \dots \right]$$

$$\alpha_1 + \frac{a(0)}{f_1} + \frac{\tilde{b}_{00}}{3} = 0, \quad \tilde{b}_{00} - \alpha_n = 0$$

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$$0 = \frac{q^\mu q^\nu}{(q \cdot K_1)^2} \left[\left(\alpha_1 + \frac{a(0)}{f_1} + \frac{\tilde{b}_{00}}{3} \right) K_{1\mu} K_{1\nu} - \left(\tilde{b}_{00} - \alpha_n \right) n_\mu n_\nu + \dots \right]$$

$$\alpha_1 + \frac{a(0)}{f_1} + \frac{\tilde{b}_{00}}{3} = 0, \quad \tilde{b}_{00} - \alpha_n = 0$$

$$a(0) = -f_1 \left(\alpha_1 + \frac{\alpha_n}{3} \right) = \frac{f_1}{3(K_1^2)^2} \left[K_1^2 (R_1 \cdot R_2) - 4(K_1 \cdot R_1)(K_1 \cdot R_2) \right] \quad \text{[OK!]}$$

Single cut & Tadpoles – II

More difficult: massless momentum & degenerate masses

Single cut & Tadpoles – II

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m^2, \quad D_1 = (k - K_1)^2 - m^2, \quad f_1 = 0]$$

We want the coefficient of $A_0(m^2)$, a_{total}

● $K_1^2 = 0$ & $m = m_0 = m_1$, OPP & PV decomposition has to be changed

$$\mathcal{I}_{2,2} = a(0) \frac{1}{D_0} + a(1) \frac{1}{D_1} + b \frac{1}{D_0 D_1} + \hat{b}_{00} \frac{(2k \cdot v)^2}{D_0 D_1} + \text{spurious terms}$$

Single cut & Tadpoles – II

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m^2, \quad D_1 = (k - K_1)^2 - m^2, \quad f_1 = 0]$$

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$$\bullet \int \frac{(2k \cdot v)^2}{D_0 D_1} = \frac{(2K_1 \cdot v)^2}{3m^2} A_0(m^2) \rightsquigarrow \hat{b}_{00} \text{ needed}$$

$$\bullet \int \frac{1}{D_0} = \int \frac{1}{D_1} = A_0(m^2) \rightsquigarrow a(0) \text{ \& } a(1) \text{ needed}$$

$$\bullet \int \frac{1}{D_0 D_1} = \frac{A_0(m^2)}{m^2} \rightsquigarrow b \text{ needed}$$

Single cut & Tadpoles – II

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m^2, \quad D_1 = (k - K_1)^2 - m^2, \quad f_1 = 0]$$

We want the coefficient of $A_0(m^2)$, $a_{\text{total}} = a(0) + a(1) + \frac{b}{m^2} + \hat{b}_{00} \frac{4(K_1 \cdot v)^2}{3m^2}$

Strategy: 1) extract $a(0) + a(1)$ & \hat{b}_{00} using the $\Lambda^2 t$ terms of the cut
2) extract b selecting the logarithmic coefficients of the cut

$$\begin{aligned} a(0) + a(1) & : (R_1 \cdot R_2) \\ \frac{4(K_1 \cdot v)^2}{3m^2} \hat{b}_{00} & : \frac{4(K_1 \cdot R_1)(K_1 \cdot R_1)}{3m^2} \\ \frac{1}{m^2} b & : (R_1 \cdot R_2) \end{aligned}$$

Single cut & Tadpoles – II

$$\mathcal{I}_{2,2} = \frac{(2R_1 \cdot k)(2R_2 \cdot k)}{D_0 D_1}, \quad [D_0 = k^2 - m^2, \quad D_1 = (k - K_1)^2 - m^2, \quad f_1 = 0]$$

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Strategy: 1) extract $a(0) + a(1)$ & \hat{b}_{00} using the $\Lambda^2 t$ terms of the cut
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$$\begin{array}{rcl}
 a(0) + a(1) & : & (R_1 \cdot R_2) \quad + \\
 \frac{4(K_1 \cdot v)^2}{3m^2} \hat{b}_{00} & : & \frac{4(K_1 \cdot R_1)(K_1 \cdot R_1)}{3m^2} \quad + \\
 \frac{1}{m^2} b & : & (R_1 \cdot R_2) \quad = \\
 \hline
 a_{\text{total}} & : & 2R_1 \cdot R_2 + \frac{4}{3m^2} (K_1 \cdot R_1)(K_1 \cdot R_2) \quad \text{[OK with PV]}
 \end{array}$$

Conclusions & Outlook

Conclusions

- Auxiliary propagator method
 - exploits D-dimensional double cut
 - gives the tadpole coefficients
 - "free" from spurious terms
- Single cut
 - useful to compute the tadpole coefficients . . .
 - . . . and the bubble coefficients as well!
 - spurious terms enter
 - ok when the Gram determinant, Δ_G , vanishes

Outlook

- Study the auxiliary propagator method in the $\Delta_G = 0$ case
- Apply the single cut to get the (massless) bubble coefficients
- Apply both methods to physically relevant cases