



TWO-LOOP QCD AMPLITUDES FOR $2 \rightarrow 3$ PROCESSES

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In collaboration with Michal Czakon, Alex Mitov, and Rene Poncelet

Based on arXiv:[1911.00479](https://arxiv.org/abs/1911.00479), [2012.13553](https://arxiv.org/abs/2012.13553), [2103.04319](https://arxiv.org/abs/2103.04319), [2105.06940](https://arxiv.org/abs/2105.06940)

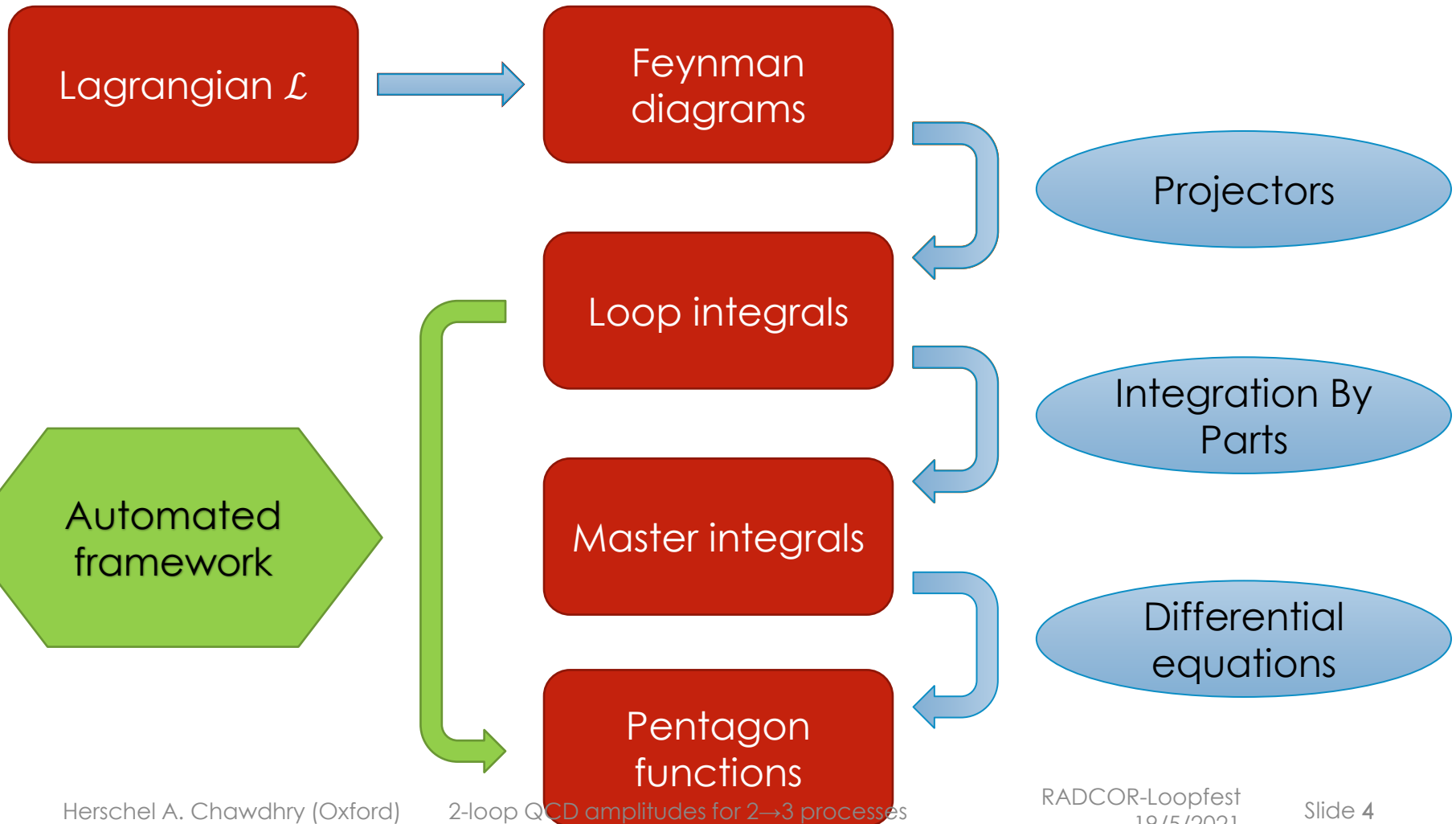
MOTIVATION

- Precise predictions for collider observables require higher-order QCD calculations
- NNLO QCD: $2 \rightarrow 1$ and $2 \rightarrow 2$ processes mostly done
 - Consistent with LHC measurements
- $2 \rightarrow 3$ is next frontier in NNLO QCD calculations
- 2-loop amplitudes are a key ingredient
 - Lots of work by many research groups (see summary slide by Vasily Sotnikov)
- $2 \rightarrow 3$ cross-sections at NNLO QCD
 - $pp \rightarrow \gamma\gamma$
 - [1911.00479](#) (HAC, Czakon, Mitov, Poncelet)
 - [2010.04681](#) (Kallweit, Sotnikov, Wiesemann)
 - $pp \rightarrow \gamma\gamma + j$
 - [2105.06940](#) (HAC, Czakon, Mitov, Poncelet)

OUTLINE

- Introduction
- Techniques
 - polarisation projectors
 - improved IBPs
 - automated framework for assembling amplitudes
- Results:
 - amplitudes for 3-photon production
 - cross-sections: 3-photon production at NNLO in QCD
 - amplitudes for 2-photon + jet production
 - cross-sections: 2-photon + jet production at NNLO in QCD
- Conclusions

OUR MULTI-LOOP TOOLBOX



POLARISATION PROJECTORS

[1904.00705](#) (Chen)

$$\mathcal{M}^{\bar{h}} = \varepsilon_{3,h_3}^\mu \varepsilon_{4,h_4}^\nu \varepsilon_{5,h_5}^\rho \text{Tr} \left\{ (u(h_1) \otimes \bar{v}(h_2)) \Gamma_{\mu\nu\rho} \right\}$$

- Construct explicit representations of the external wavefunctions for a given helicity or polarisation

$$\varepsilon_{i,X}^\mu = c_{i,1}^X p_1^\mu + c_{i,2}^X p_2^\mu + c_{i,3}^X p_i^\mu$$

$$(\varepsilon_{i,X})^2 = -1 \quad , \quad \varepsilon_{i,X} \cdot q = 0 \quad , \quad \varepsilon_{i,X} \cdot p_i = 0$$

$$\varepsilon_{i,h}^\mu = \frac{1}{\sqrt{2}} (\varepsilon_{i,X}^\mu + h i \varepsilon_{i,Y}^\mu)$$

$$(u \otimes \bar{v})_{\alpha\beta} = \frac{\bar{u} N v}{\bar{u} N v} (u \otimes \bar{v})_{\alpha\beta} = \frac{1}{\bar{u} N v} (u \otimes \bar{u})_{\alpha\gamma} N_{\gamma\delta} (v \otimes \bar{v})_{\delta\beta} = \frac{1}{\mathcal{N}} [(u \otimes \bar{u}) N (v \otimes \bar{v})]_{\alpha\beta}$$

- Construction is in $d = 4$
 - Bare amplitudes are scheme-dependent
 - Dependence cancels in finite remainder

IBPS: A NEAT STRATEGY

[1805.09182](#) (HAC, Lim, Mitov). See also [1812.01491](#) (Maierhöfer, Usovitsch)

- Loop integrals live in a vector space over $\mathbb{Q}(\{s_{ij}\}, d)$
- Master integrals $\{M_n\}$ form a basis for this space
- Can therefore simplify the IBP equations by *taking coefficients w.r.t. the basis vectors*
- In practice:
 - Set $M_n = 0$ for all $n \neq 1$
 - Solve simplified IBPs
 - Repeat in turn for $n = 2, n = 3, \dots$ to obtain complete IBP solution
- Correct because in any vector space, a vector can be uniquely written as a vector superposition of a basis
 - Equivalently, correct because IBP equations - and solutions - are linear (in the integrals) and homogeneous
- Advantage: one simplified IBP system per M_j
 - Independent \rightarrow can parallelise
 - Simpler \rightarrow faster to solve, requires less RAM

AUTOMATED ASSEMBLY OF THE AMPLITUDES

- Ingredients
 - Coefficients of unreduced integrals in amplitude
 - IBP solutions
 - master integral solutions
 - IR counterterms (if seeking finite remainder)
- All these ingredients are calculated analytically, but assembling them is non-trivial
- Our automated framework uses finite fields for the assembly
 - “Numerically” evaluate each of the ingredients
 - Calculate numerical value of the amplitude/finite remainder
 - Interpolate to obtain analytical amplitude
- Complications (see later)
 - Momentum crossings
 - Expansion in $\varepsilon \equiv (4 - d)/2$
 - $A^{(2\text{-loop})} = \sum_n \varepsilon^n \sum_k r_{n,k}(s_{ij}) g_k(s_{ij})$, where $g_k(s_{ij})$ are pentagon functions
- Simultaneously calculate multiple amplitude structures A_n by sharing finite-field samples of IBPs and masters
 - useful for processes with many partonic channels, and/or helicity amplitudes

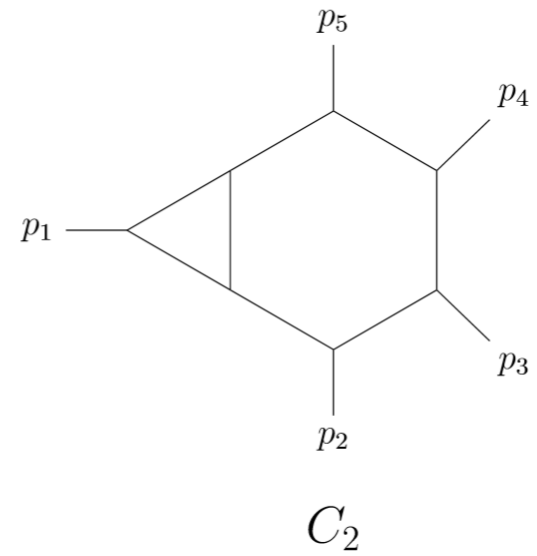
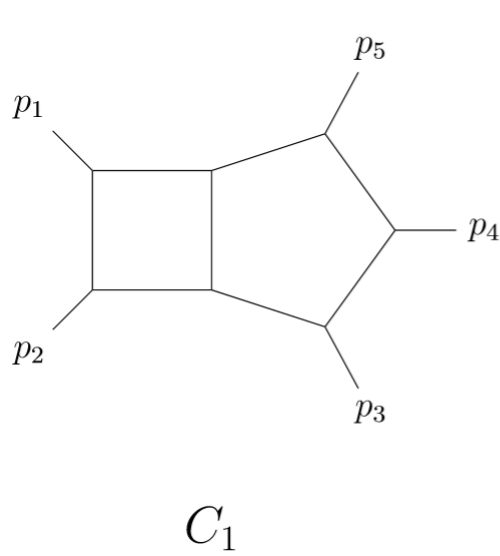
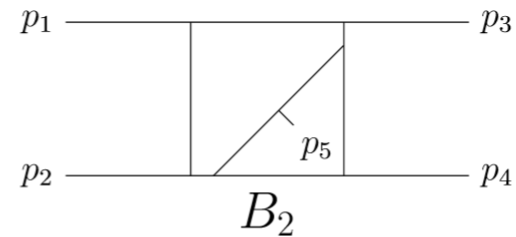
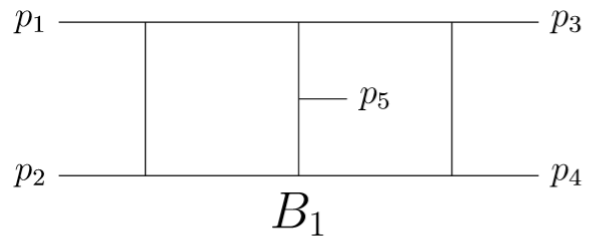
FINITE-FIELD METHODS

- Long-used in computer algebra (e.g. Mathematica), now also used by physicists
 - e.g. [[1406.4513](#) – Manteuffel, Schabinger], [[1608.01902](#) – Peraro]
- Speed largely determined by the polynomial degrees of the expressions in the final result
- Reconstruct analytical results using interpolation and Chinese remainder theorem
 - we use FireFly library [[1904.00009](#) , [2004.01463](#) – Klappert, Klein, Lange]

IBP SOLUTIONS

- Using our analytical solutions from [1805.09182](#) (HAC, Lim, Mitov)
- Preparatory step: analytically expand IBPs to 5 orders in $\varepsilon \equiv (4 - d)/2$
 - Insert book-keeping terms to account for truncated higher-orders in ε
- Analytical momentum crossings not required
 - During finite-field sampling, apply momentum crossings numerically to the finite-field values of $\{s_{ij}\}$
- Time needed to evaluate all required IBPs: <10s per crossing (without any optimisations)

INTEGRAL FAMILIES



MASTER INTEGRALS

- Solved several times:
 - [1511.09404](#) (Papadopoulos, Tommasini, Wever)
 - [1807.09812](#) (Gehrmann, Henn, Lo Presti)
 - [2009.07803](#) (Chicherin, Sotnikov)
- C1 masters \rightarrow UT masters \rightarrow pentagon functions
 - produced analytically for each momentum crossing
 - minimal basis of pentagon functions
 - analytically expanded to 5 orders in ε
 - inserted book-keeping terms (as with IBPs)
- mapped C2 masters \rightarrow C1 masters
- Finite-field evaluation time < 1 s per crossing

SIMPLIFYING THE FINAL RESULTS

- \mathbb{Q} -linear relations between the coefficients r_k of pentagon functions g_k
 - $\sum_k \lambda_k r_k = 0$, with $\lambda_k \in \mathbb{Q}$ recall: $A = \sum r_k g_k$, with $r_k \in \mathbb{Q}(\{s_{ij}\}, d)$
 - Greatly reduces number - and total size - of coefficients
 - Used by e.g.
 - [\[1904.00945\]](#) – Abreu, Dormans, Febres Cordero, Ita, Page, Sotnikov]
 - [\[1911.00479\]](#) – HAC, Czakon, Mitov, Poncelet]
 - [\[2102.01820\]](#) – Agarwal, Buccioni, von Manteuffel, Tancredi]
 - [\[2102.02516\]](#) – Badger, Hartanto, Zoia]
 - [\[2102.13609\]](#) – Abreu, Febres Cordero, Ita, Page, Sotnikov]
 - Exploration of origin and further applications?
- partial fractioning
 - e.g. MultivariateApart [\[2101.08283\]](#) – Heller, von Manteuffel]
 - Further reduces sizes of coefficients by factor of $\sim 2-4$

RESULTS:

LEADING-COLOUR $q\bar{q} \rightarrow \gamma\gamma\gamma$

- ✓ Spin-averaged $2\text{Re}\langle A^{(0)} | A^{(2)} \rangle$ [[1911.00479](#)] (HAC, Czakon, Mitov, Poncelet)
 - Automated results checked against by-hand calculation
 - Pentagon functions' coefficients were large (but fast)
 - For pheno, evaluate coefficients using exact rational arithmetic
 - Takes 1-2 seconds per kinematic point (full set of coefficients)
 - Pentagon functions themselves are much slower
 - 10-50 min per kinematic point (average 17 min) to evaluate functions from [1807.09812](#) (Gehrmann, Henn, Lo Presti)
 - [2009.07803](#) (Chicherin, Sotnikov) offers large improvements
- ✓ Full set of 16 linearly-polarised amplitudes
 - Checked against unpolarised $2\text{Re}\langle A^{(0)} | A^{(2)} \rangle$
- ✓ Full set of helicity amplitudes [[2012.13553](#)] (HAC, Czakon, Mitov, Poncelet)
 - Much lower polynomial degrees than in linearly-polarised
 - Raw results ~100 MB (reduces to 1.8 MB using \mathbb{Q} -linear relations and partial fractioning)

RESULTS:

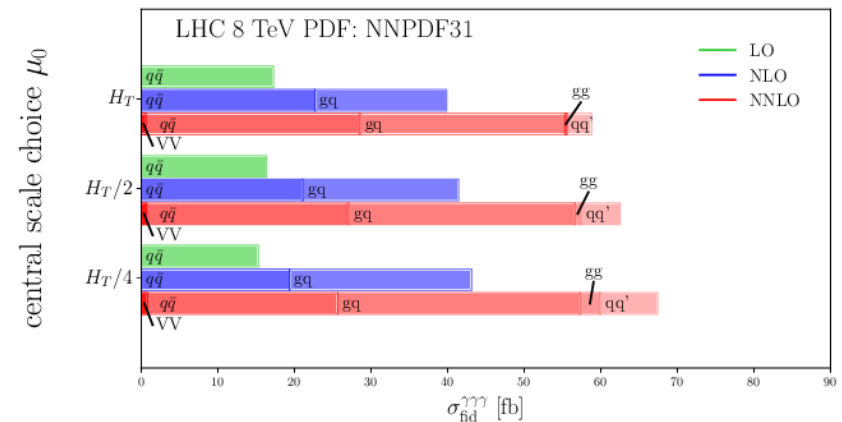
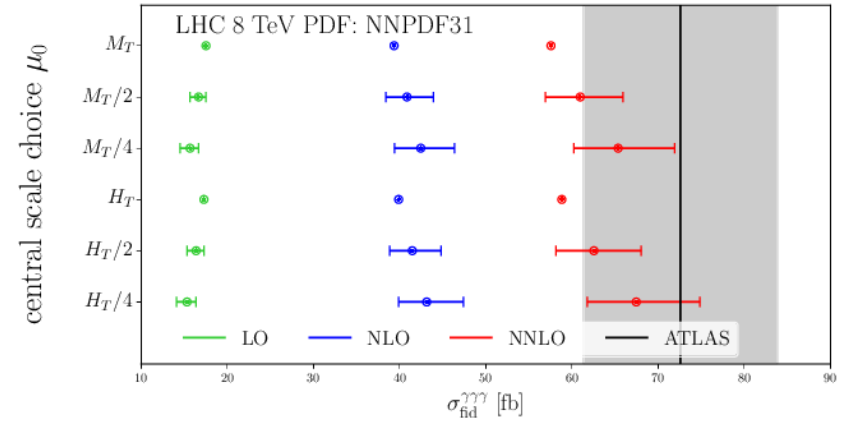
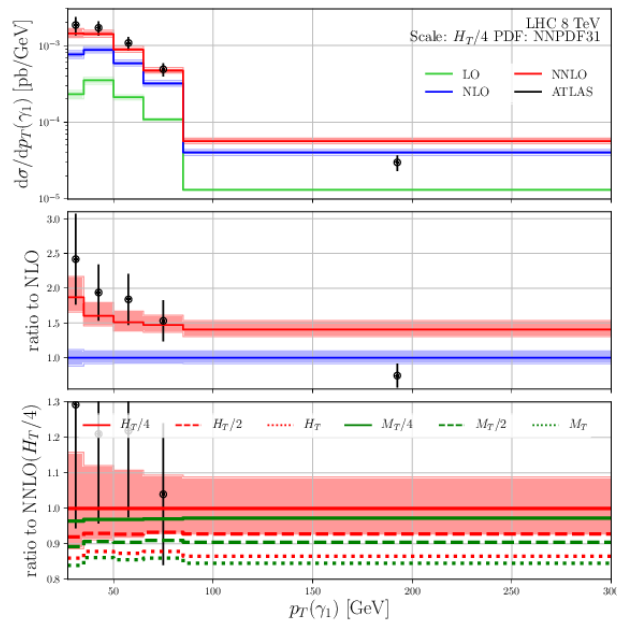
LEADING-COLOUR $\gamma\gamma + j$

- Full set of helicity amplitudes [\[2103.04319\]](#) (HAC, Czakon, Mitov, Poncelet)
- 2 partonic channels, 9 helicity amplitudes
 - $q\bar{q} \rightarrow g\gamma\gamma$ (3 independent helicities)
 - $qg \rightarrow q\gamma\gamma$ (6 independent helicities)
- Automated calculation took $O(1)$ day
- Raw results \sim a few hundred MB (reduces to ~ 4 MB after \mathbb{Q} -linear relations and partial fractioning)
- Obtained the first NNLO QCD cross-sections for $\gamma\gamma + j$ production at the LHC (see later) [\[2105.06940\]](#) (HAC, Czakon, Mitov, Poncelet)

PHENO RESULTS: 3-PHOTON

[1911.00479](#) (HAC, Czakon, Mitov, Poncelet)

- NNLO gives good description of LHC data
- scale-independent leading-colour 2-loop finite remainder contribution is at % level



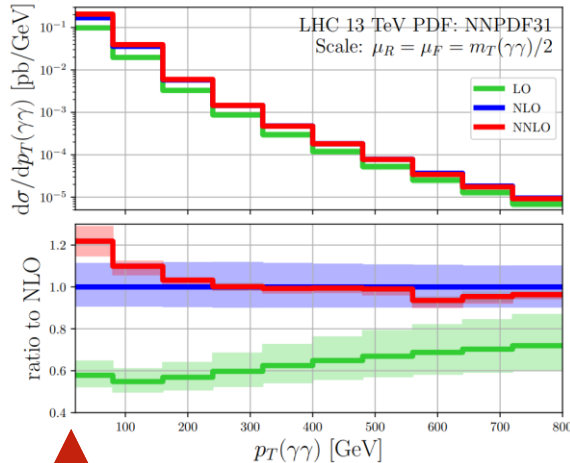
See talk by Rene Poncelet on Friday at 15:00 (Florida time)

PHENO RESULTS: 2-PHOTON + JET

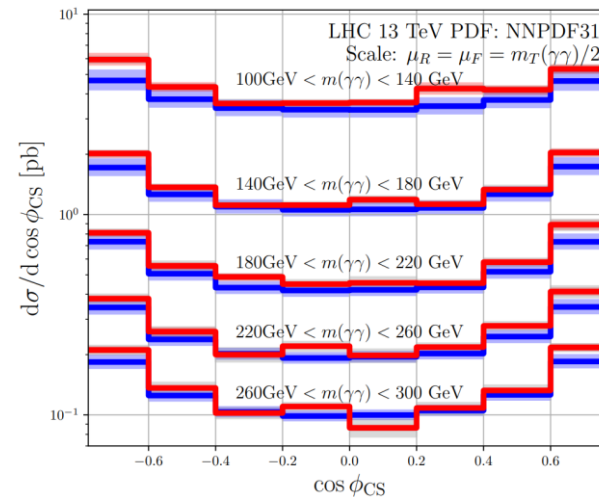
[2105.06940](#) (HAC, Czakon, Mitov, Poncelet)

See talk by Rene Poncelet on Friday at 15:00 (Florida time)

- First NNLO-accurate prediction for $p_T(\gamma\gamma)$ distribution
- Main background for Higgs production at high p_T



↑
Require $p_T(\gamma\gamma) > 20$ GeV



- 2-loop subleading-colour effects likely to be 1-2% at most
- Future work: calculate 2-loop amplitude for loop-induced process $gg \rightarrow g\gamma\gamma$ (gives a partial N³LO contribution)

CONCLUSIONS

- Fast automated framework for assembling amplitudes
 - well-suited for processes with several partonic channels and/or helicities and/or colour-structures
- Several 2-loop QCD amplitudes for 2 → 3 processes
 - $q\bar{q} \rightarrow \gamma\gamma\gamma$ (leading colour)
 - ✓ spin-averaged $2\text{Re}[\langle A^{(0)} | A^{(2)} \rangle]$
 - evaluation time dominated by pentagon functions, not their coefficients
 - ✓ complete set of 16 linearly-polarised amplitudes
 - ✓ complete set of helicity amplitudes
 - ✓ NNLO QCD cross-sections for 3-photon production
 - $\gamma\gamma$ +jet production
 - ✓ complete set of helicity amplitudes (leading colour)
 - ✓ NNLO QCD cross-sections for 2-photon+jet production
- More details: [[1911.00479](#), [2012.13553](#), [2103.04319](#), [2105.06940](#)]