Full NLO predictions for ZZ scattering and its irreducible background at the LHC


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20 May 2021

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Introduction

## Characteristics of vector-boson scattering (VBS)



- Sensitivity to scalar sector of the SM and non-Abelian triple and quartic couplings.
- NLO EW corrections in VBS typically of the size of $\sim-15 \%$.
- VBS ZZ signature:
$\square 60$ partonic quark-induced channels compared to 40 for WZ and 12 for same-sign W.
$\square$ additional s-channel Higgs contribution, but in our setup the Higgs is cut out by invariant mass cut on the four leptons.


## Previous theoretical work in VBS

- NLO QCD corrections in VBS approximation for all VBS processes [Baglio (2014), Rauch (2017)].
- NLO EW corrections for same-sign W and WZ [Biedermann, Denner, Pellen (2016, 2017), Denner, Dittmaier, Maierhöfer, Pellen, Schwan (2019)].
- NLO EW + QCD corrections in same-sign W with event generator [Chiesa, Denner, Lang, Pellen (2019)].
- VBS ZZ:
$\square$ NLO QCD and EW corrections in VBS approximation [Jäger, Oleari, Zeppenfeld (2006)] and matched to QCD parton shower [Jäger, Karlberg, Zanderighi (2014)].
$\square$ NLO QCD corrections to QCD-induced process [Campanario, Kerner, Ninh, Zeppenfeld (2014)].
■ loop-induced ZZ $(+2 \mathrm{j})$ production matched to parton shower [Li, An, Charlot, Covarelli, Guan, Li (2020)].


## LO contributions



- EW contributions © $\mathcal{O}\left(\alpha^{6}\right)$
- interference contributions @ $\mathcal{O}\left(\alpha_{s} \alpha^{5}\right)$
- QCD contributions @ $\mathcal{O}\left(\alpha_{s}^{2} \alpha^{4}\right)$
- gluon-induced contributions @ $\mathcal{O}\left(\alpha_{s}^{4} \alpha^{4}\right)$



## NLO corrections (I)



- virtual corrections:

- real corrections:



## NLO corrections (II)

## Technical setup:

$\square$ Results at LO and NLO obtained with MoCaNLO + Recola [Actis, Denner, Hofer, Lang, Scharf, Uccirati $(2013,2017)$ )].
$\square$ NLO virtual corrections:
■ Evaluation of tensor and scalar integrals done by COLLIER [Denner, Dittmaier, Hofer (2016))].
$\square$ NLO real corrections:

- IR singularities handeled by Catani-Seymour dipole subtraction in QCD [Catani, Seymour $(1997,1998)$ )] and its equivalent in QED [Dittmaier (2000), Dittmaier, Kabelschacht, Kasprzik (2008)].
- collinear singularities also arises from diagrams with a low-virtual photon $\rightarrow$ photon-to-jet-conversion function [Denner, Dittmaier, Pellen, Schwan (2019)].
- For real EW corrections @ $\mathcal{O}\left(\alpha_{s}^{2} \alpha^{5}\right)$ recombination of hard photon and soft gluon can lead to infrared QCD singularity $\rightarrow$ fragmentation function [Denner, Hofer, Scharf, Uccirati (2015)].
- Comparisons:
$\square$ Independent checks in VBS between MoCaNLO and BBMC in the same sign W and MoCaNLO and BONSAY in WZ signature.VBS ZZ:
- virtuals checked with different implementations in COLLIER (COLI and DD mode).
- reals and integrated dipole contributions checked via different $\alpha$-parameters ( $\alpha=10^{-2}$ and $\alpha=1$ ).
- additional checks of a subset of channels between BBMC and MoCaNLO at all orders.


## Input parameters

- cms energy $\sqrt{s}=13 \mathrm{TeV}$.
- PDF set: NLO NNPDF-3.1 Lux QED (for both LO and NLO).
- strong coupling: $\alpha_{s}\left(M_{Z}\right)=0.118$.

■ renormalization and factorization scale: $\mu_{\mathrm{ren}}=\mu_{\mathrm{fac}}=\sqrt{p_{\mathrm{T}, j_{1}} p_{\mathrm{T}, j_{2}}},\left(p_{\mathrm{T}, j_{1}}\right.$ and $p_{\mathrm{T}, j_{1}}$ are the two jets with the largest transverse momentum).

- scale variation: 7-point method $\left(\mu_{\text {ren }} / \mu_{0}, \mu_{\mathrm{fac}} / \mu_{0}\right)=(0.5,0.5),(0.5,1),(1,0.5),(1,1),(1,2),(2,1),(2,2)$.
- electromagnetic coupling: $G_{\mu}$ scheme.
- masses and widths:

$$
\begin{aligned}
m_{\mathrm{t}} & =173.0 \mathrm{GeV}, & \Gamma_{\mathrm{t}} & =0 \mathrm{GeV}, \\
M_{\mathrm{Z}}^{\mathrm{OS}} & =91.1876 \mathrm{GeV}, & \Gamma_{\mathrm{Z}}^{\mathrm{OS}} & =2.4952 \mathrm{GeV}, \\
M_{\mathrm{W}}^{\mathrm{OS}} & =80.379 \mathrm{GeV}, & \Gamma_{\mathrm{W}}^{\mathrm{OS}} & =2.085 \mathrm{GeV}, \\
M_{\mathrm{H}} & =125.0 \mathrm{GeV}, & \Gamma_{\mathrm{H}} & =4.07 \times 10^{-3} \mathrm{GeV} .
\end{aligned}
$$

- pole masses of vector bosons obtained from OS masses.


## Event selection

- combination of QCD partons into jets via anti- $k_{T}$ algorithm.
- cuts on leptons:

$$
p_{\mathrm{T}, \ell}>20 \mathrm{GeV}, \quad\left|\eta_{\ell}\right|<2.5, \quad \Delta R_{\ell \ell^{\prime}}>0.05, \quad M_{\ell^{+} \ell^{\prime}-}>4 \mathrm{GeV} .
$$

- cut on the invariant mass of two leptons:

$$
60 \mathrm{GeV}<M_{\ell^{+} \ell^{-}}<120 \mathrm{GeV}, \quad \ell=\mathrm{e}, \mu
$$

- cuts on the jets and (jets and leptons):

$$
p_{\mathrm{T}, \mathrm{j}}>30 \mathrm{GeV}, \quad\left|\eta_{\mathrm{j}}\right|<4.7, \quad \Delta R_{\mathrm{j} \ell}>0.4 .
$$

- cut on the invariant mass of the two $p_{T}$-hardest jets:

$$
M_{\mathrm{j}_{1} \mathrm{j}_{2}}>100 \mathrm{GeV} \quad \text { (inclusive setup), } \quad M_{\mathrm{j}_{1} \mathrm{j}_{2}}>500 \mathrm{GeV} \quad \text { (VBS setup) }
$$

- Results presented in this talk are based on the two setups for the invariant mass.
- Plots in this talk only in the inclusive setup.

Numerical Results

## Cross sections at LO

| Order | $\mathcal{O}\left(\alpha^{6}\right)$ | $\mathcal{O}\left(\alpha_{\mathrm{s}} \alpha^{5}\right)$ | $\mathcal{O}\left(\alpha_{\mathrm{s}}^{2} \alpha^{4}\right)$ | $\mathcal{O}\left(\alpha_{\mathrm{s}}^{4} \alpha^{4}\right)$ | Sum |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>100 \mathrm{GeV}$ |  |  |  |  |  |
| $\sigma_{\mathrm{LO}}[\mathrm{fb}]$ | $0.097683(2)$ | $0.008628(1)$ | $1.062478(48)$ | $0.12101(64)$ | $1.28980(64)$ |
| fraction[\%] | 7.57 | 0.67 | 82.38 | 9.38 | 100 |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>500 \mathrm{GeV}$ |  |  |  |  |  |
| $\sigma_{\mathrm{LO}}[\mathrm{fb}]$ | $0.073676(3)$ | $0.005567(1)$ | $0.136143(15)$ | $0.01345(29)$ | $0.22883(29)$ |
| fraction[\%] | 32.20 | 2.43 | 59.49 | 5.88 | 100 |

■ QCD contributions one order of magnitude larger as the EW contributions in contrast to same-sign W and WZ .

- with stronger cut on $M_{j_{1} j_{2}}$ QCD background significantly decreases while EW corrections only moderately decrease.
■ Loop-induced process contributes at $\sim 10 \%$. Contribution decreases with stronger cut on the two hardest jets.


## LO distributions (I)

- Rapidity difference (left) and invariant mass (right) of the two hardest jets $j_{1}$ and $j_{2}$.

- Rapidity difference of the two hardest jets enhance the EW contributions only for values $\left|\Delta y_{j_{1} j_{2}}\right|>5$.
- Invariant mass distribution of the two hardest jets for the EW contributions enhanced for values $M_{j_{1} j_{2}}>1200 \mathrm{GeV}$, while the QCD contribution exceed the latter in the region below.


## Cross sections at NLO (I)

| Contribution | $\sigma_{\alpha^{6}}[\mathrm{ab}]$ | $\Delta \sigma_{\alpha^{7}}[\mathrm{ab}]$ | $\Delta \sigma_{\alpha^{7}} / \sigma_{\alpha^{6}}[\%]$ | $\Delta \sigma_{\alpha_{\mathrm{s}} \alpha^{6}}[\mathrm{ab}]$ | $\Delta \sigma_{\alpha_{\mathrm{s}} \alpha^{6}} / \sigma_{\alpha^{6}}[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>100 \mathrm{GeV}$ |  |  |  |  |  |
| all | 97.683(2) | -15.55(5) | -15.9 | 23.10(11) | 23.6 |
| VBS-WW | 95.237(2) | -15.28(5) | -16.0 | 1.33(11) | 1.4 |
| VBS-ZZ | 1.9463(2) | -0.1979(6) | -10.2 | 3.892(4) | 200 |
| WZZ | 0.1361(1) | -0.0142(1) | -10.5 | 13.850(4) | 10174 |
| ZZZ | 0.3629(1) | -0.0542(6) | -14.9 | 4.029(3) | 1110 |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>500 \mathrm{GeV}$ |  |  |  |  |  |
| all | 73.679(2) | -13.01(4) | -17.7 | 0.07(25) | 0.10 |
| VBS-WW | 72.846(2) | -12.91(4) | -17.7 | -2.73(25) | -3.7 |
| VBS-ZZ | 0.8096(2) | -0.0986(3) | -12.2 | 0.486(6) | 60.1 |
| WZZ | 0.00471(2) | -0.00085(1) | -18.1 | 1.849 (5) | 39258 |
| ZZZ | 0.01887(1) | -0.00529(2) | -28.0 | 0.470(1) | 2488 |

- 16 partonic channels with subprocess WW-ZZ dominate LO and NLO EW corrections at $\mathcal{O}\left(\alpha^{7}\right)$.
- $\mathcal{O}\left(\alpha_{s} \alpha^{6}\right)$ corrections dominated by triple-vector-boson production in the inclusive setup.
- Stronger cut on $M_{j_{1} j_{2}}>500 \mathrm{GeV}$ reduces contributions of triple-vector-boson channels.


## NLO distributions (I)

■ invariant mass of the two hardest jets (left) and rapidity separation of the two hardest jets (right):




- Plots are normalized to the order $\mathcal{O}\left(\alpha^{6}\right)$ Born cross section.
- Bulk of NLO QCD corrections is in low di-jet invariant mass regime or regions with low rapidity difference of the two hardest jets.


## Large QCD corrections in the inclusive setup

- In case of real radiation invariant mass cut of 100 GeV does not necessarily apply to the two quarks coming from vector-boson-decay.
- Extra jet can invoke kinematical configurations, where $M_{j_{1} j_{2}}>100$.

- The inclusive setup should therefore be avoided when using VBS approximations neglecting diagrams containing tri-boson contributions.
- Subtraction of tri-boson contributions generally not favourable due to possible gauge-invariance violation.


## Cross sections at NLO (II)

| Part. channel | $\sigma_{\alpha^{6}}[\mathrm{ab}]$ | $\delta_{\alpha^{7}}[\%]$ | $\delta_{\text {LL }}[\%]$ | $\delta_{\text {LL+SSC }}[\%]$ | subprocesses |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{ud} \rightarrow \mathrm{e}^{+} \mathrm{e}^{-} \mu^{+} \mu^{-} \mathrm{ud}$ | $51.537(2)$ | $-17.3(1)$ | -16.4 | -14.6 | VBS-WW/VBS-ZZ |
| us $\rightarrow \mathrm{e}^{+} \mathrm{e}^{-} \mu^{+} \mu^{-} \mathrm{dc}$ | $12.769(1)$ | $-15.1(1)$ | -14.2 | -12.6 | VBS-WW |
| $\mathrm{u} \overline{\mathrm{u}} \rightarrow \mathrm{e}^{+} \mathrm{e}^{-} \mu^{+} \mu^{-} \mathrm{d} \overline{\mathrm{d}}$ | $10.666(1)$ | $-15.0(1)$ | -13.6 | -10.1 | VBS-WW/ZZZ |
| $\mathrm{uu} \rightarrow \mathrm{e}^{+} \mathrm{e}^{-} \mu^{+} \mu^{-} \mathrm{uu}$ | $0.37718(5)$ | $-11.8(1)$ | - | - | VBS-ZZ |
| $\mathrm{u} \overline{\mathrm{d}} \rightarrow \mathrm{e}^{+} \mathrm{e}^{-} \mu^{+} \mu^{-} \mathrm{u} \overline{\mathrm{d}}$ | $0.24011(5)$ | $-10.2(1)$ | - | - | WZZ |
| $\mathrm{u} \overline{\mathrm{u}} \rightarrow \mathrm{e}^{+} \mathrm{e}^{-} \mu^{+} \mu^{-} \mathrm{u} \overline{\mathrm{u}}$ | $0.15878(4)$ | $-11.6(1)$ | - | - | VBS-ZZ/ZZZ |
| $\mathrm{d} \overline{\mathrm{d}} \rightarrow \mathrm{e}^{+} \mathrm{e}^{-} \mu^{+} \mu^{-} \mathrm{s} \overline{\mathrm{s}}$ | $0.11638(3)$ | $-11.0(1)$ | - | - | ZZZ |

- The $\mathcal{O}\left(\alpha^{7}\right)$ contributions dominated by large Sudakov logarithms.
- Sudakov approximation for VBS-WW channels in the LL and SSC [Denner, Pozzorini (2001), Accomando, Denner, Pozzorini (2007)]:
$\square$ Universal and for all VBS processes:

$$
\delta_{\mathrm{LL}}=\frac{\alpha}{4 \pi}\left\{-4 C_{W}^{\mathrm{EW}} \log ^{2}\left(\frac{Q^{2}}{M_{\mathrm{W}}^{2}}\right)+2 b_{W}^{\mathrm{EW}} \log \left(\frac{Q^{2}}{M_{\mathrm{W}}^{2}}\right)\right\} .
$$

$\square$ Process dependent (WW $\rightarrow$ ZZ):

$$
\delta_{\mathrm{SSC}}=\frac{\alpha}{\pi s_{\mathrm{w}}^{2}} 2 \ln \left(\frac{Q^{2}}{M_{\mathrm{W}}^{2}}\right)\left[-\ln \frac{s_{12}}{Q^{2}}+\frac{s_{23}}{s_{12}} \ln \frac{s_{13}}{Q^{2}}+\frac{s_{13}}{s_{12}} \ln \frac{s_{23}}{Q^{2}}\right] .
$$

- Agreement of approximation and full results within $2 \%$.


## Cross sections at NLO (III)

| Order | $\mathcal{O}\left(\alpha^{7}\right)$ | $\mathcal{O}\left(\alpha_{\mathrm{s}} \alpha^{6}\right)$ | $\mathcal{O}\left(\alpha_{\mathrm{s}}^{2} \alpha^{5}\right)$ | $\mathcal{O}\left(\alpha_{\mathrm{s}}^{3} \alpha^{4}\right)$ | Sum |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>100 \mathrm{GeV}$ |  |  |  |  |  |
| $\Delta \sigma_{\mathrm{NLO}}[\mathrm{fb}]$ | $-0.01557(4)$ | $0.0231(1)$ | $-0.0862(1)$ | $-0.0530(16)$ | $-0.1317(16)$ |
| $\Delta \sigma_{\mathrm{NLO}} / \sigma_{\mathrm{LO}}[\%]$ | $-1.33(1)$ | $1.98(1)$ | $-7.38(2)$ | $-4.54(14)$ | $-11.27(14)$ |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>500 \mathrm{GeV}$ |  |  |  |  |  |
| $\Delta \sigma_{\mathrm{NLO}}[\mathrm{fb}]$ | $-0.01299(5)$ | $0.00008(25)$ | $-0.0142(1)$ | $0.0058(11)$ | $-0.0214(11)$ |
| $\Delta \sigma_{\mathrm{NLO}} / \sigma_{\mathrm{LO}}[\%]$ | $-6.03(3)$ | $0.04(12)$ | $-6.60(5)$ | $2.67(50)$ | $-9.91(51)$ |

- For $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>100 \mathrm{GeV}$ corrections at the order $\mathcal{O}\left(\alpha^{7}\right)$ and $\mathcal{O}\left(\alpha_{s} \alpha^{6}\right)$ are small, varying between $1 \%$ and $2 \%$.
- In the VBS setup with $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>500 \mathrm{GeV}$ the largest contribution is the one at order $\mathcal{O}\left(\alpha_{s}^{2} \alpha^{5}\right)$, dominated by the EW corrections to the LO QCD contribution.


## NLO distributions (II)

■ invariant mass of the two hardest jets (left) and rapidity separation of the two hardest jets (right):




- Plots are normalized to the full LO cross section.
- Both distributions receive sizeable contributions in the large $M_{j_{1} j_{2}}$ and $\left|\Delta y_{j_{1} j_{2}}\right|$ regime at the order $\mathcal{O}\left(\alpha^{6}\right)$ while other distributions are dominated by the $\mathcal{O}\left(\alpha_{s}^{2} \alpha^{4}\right)$.
- Large Sudakov corrections visible both for small invariant mass or rapidity difference at the order $\mathcal{O}\left(\alpha_{s}^{2} \alpha^{5}\right)$ and for large values $M_{j_{1} j_{2}}$ and $\left|\Delta y_{j_{1} j_{2}}\right|$ at the order $\mathcal{O}\left(\alpha^{7}\right)$.


## NLO distributions (III)

- invariant mass (left) and transverse momentum of the four leptons (right).



- Plots are normalized to the full LO cross section.
- Both distributions display similar behaviour as they are correlated.
- The order $\mathcal{O}\left(\alpha_{s}^{2} \alpha^{5}\right)$ show typical Sudakov behaviour reaching up to $30 \%$ while the $\mathcal{O}\left(\alpha^{7}\right)$ corrections are damped due to the dominating LO QCD contributions.
- QCD corrections at the order $\mathcal{O}\left(\alpha_{s}^{3} \alpha^{4}\right)$ show stronger impact reaching $-50 \%$ for $M_{4 \ell}=2 \mathrm{TeV}$ or $-20 \%$ for high $p_{T}$.


## NLO distributions (IV)

- transverse momentum of the hardest jet (left) and of an electron-positron pair (right):



- Plots are normalized to the full LO cross section.
- NLO correction to the contribution of transverse momentum of the hardest jet reaches up to $25 \%$. Overall behaviour is directed to the order $\mathcal{O}\left(\alpha_{s}^{3} \alpha^{4}\right)$ corrections.
- Large negative corrections towards high energies reaching $-60 \%$ at 700 GeV . NLO corrections leave LO uncertainty band at about 200 GeV and NLO scale uncertainty increases significantly for high leptonic energies.


## Conclusions

- LO contributions for VBS ZZ at the orders $\mathcal{O}\left(\alpha^{6}\right), \mathcal{O}\left(\alpha_{s} \alpha^{5}\right), \mathcal{O}\left(\alpha_{s}^{2} \alpha^{4}\right)$ and $\mathcal{O}\left(\alpha_{s}^{4} \alpha^{4}\right)$.
- Gluon-induced contributions at LO are non-negligable.
- EW corrections large by about $-16 \%$ in agreement with previous VBS calculations. Can reach about $-40 \%$ in high energy tails of distributions.
- QCD corrections at the order $\mathcal{O}\left(\alpha_{s} \alpha^{6}\right)$ reach $20 \%$ in the inclusive setup due to tri-boson contributions. Effect can be decreased by choosing sensible cut on $M_{j_{1} j_{2}}$ or $\Delta y_{j_{1} j_{2}}$.
- $\mathcal{O}\left(\alpha_{s}^{2} \alpha^{5}\right)$ also show large Sudakov corrections.
- In high energy tails of distributions the $\mathcal{O}\left(\alpha_{s}^{3} \alpha^{4}\right)$ corrections are large.

Thank you for your attention

Backup slides

## Cross sections at NLO (backup)

| Order | $\mathcal{O}\left(\alpha^{6}\right)+\mathcal{O}\left(\alpha^{7}\right)$ | $\mathcal{O}\left(\alpha^{6}\right)+\mathcal{O}\left(\alpha_{\mathrm{s}} \alpha^{6}\right)$ | $\mathcal{O}\left(\alpha^{6}\right)+\mathcal{O}\left(\alpha^{7}\right)+\mathcal{O}\left(\alpha_{\mathrm{s}} \alpha^{6}\right)$ |
| :--- | :--- | :--- | :--- |


| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>100 \mathrm{GeV}$ |  |  |  |
| :--- | :---: | :---: | :---: |
| $\sigma_{\mathrm{NLO}}[\mathrm{fb}]$ | $0.08211(4)$ | $0.12078(11)$ | $0.10521(11)$ |
| $\sigma_{\mathrm{NLO}}^{\mathrm{max}}[\mathrm{fb}]$ | $0.08728(5)[+6.3 \%]$ | $0.12540(13)[+3.8 \%]$ | $0.10838(14)[+3.0 \%]$ |
| $\sigma_{\mathrm{NLO}}^{\mathrm{min}}[\mathrm{fb}]$ | $0.07749(4)[-5.6 \%]$ | $0.11656(9)[-3.5 \%]$ | $0.10225(9)[-2.8 \%]$ |
| $\delta[\%]$ | -15.9 | 23.6 | 7.7 |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>500 \mathrm{GeV}$ |  |  |  |
| $\sigma_{\mathrm{NLO}}[\mathrm{fb}]$ | $0.06069(4)$ | $0.07375(25)$ |  |
| $\sigma_{\mathrm{NLO}}^{\max }[\mathrm{fb}]$ | $0.06568(5)[+8.2 \%]$ | $0.07466(26)[+1.2 \%]$ | $0.06149(24)[+1.2 \%]$ |
| $\sigma_{\mathrm{NLO}}^{\min }[\mathrm{fb}]$ | $0.05636(4)[-7.1 \%]$ | $0.07282(21)[-1.3 \%]$ | $0.05977(30)[-1.6 \%]$ |
| $\delta[\%]$ | -17.6 | 0.1 | -17.5 |

- EW corrections range between $\sim 16-17 \%$ depending on the invariant mass cut.
- QCD corrections sizeable in the inclusive setup but drastically reduced for stronger invariant mass cut.

