

New approach to DM searches with mono-photons

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CLIC workshop

02.10.2020

- Motivation
- Simulating mono-photon events with WHIZARD
 - new procedure for merging ISR structure function with ME
 - ArXiv: 2004.14486, **Eur. Phys. J. C (2020) 80: 634**
- New approach to DM searches with mono-photon signature
- Preliminary results
- Prospects

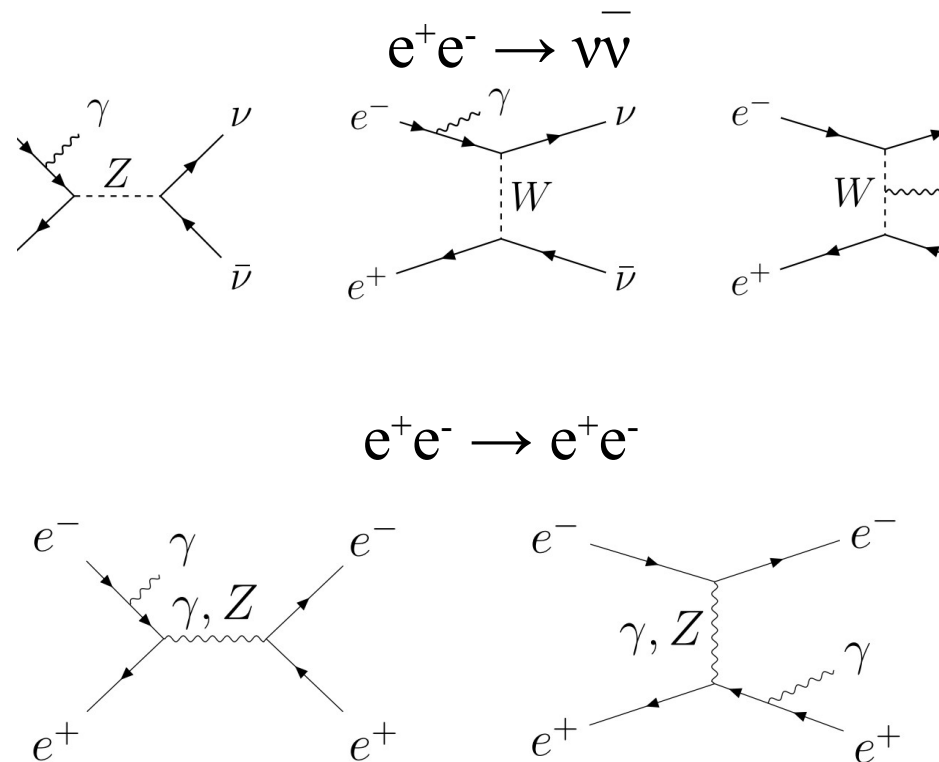
Motivation

- Initial aim: preparation of background and signal samples for **Dark Matter analysis using a mono-photon signature**. Following issues would need to be addressed:

- To avoid double counting of photons from the ISR (γ^{ISR}) and those from matrix element (γ^{ME})
- To have the same procedure for all backgrounds and DM signal, successful generation of all processes in the first place: cutting out all of the singularities, stable cross sections calculations etc

- Outcome: procedure general enough for any BSM searches with neutrino and Bhabha mono-photon events**

- **NOTE: one cannot properly simulate all matrix element photons: t-channel W radiation events (mostly with high photon multiplicity) are not included in WHIZARD description of the ISR**



Simulating mono-photon events with WHIZARD

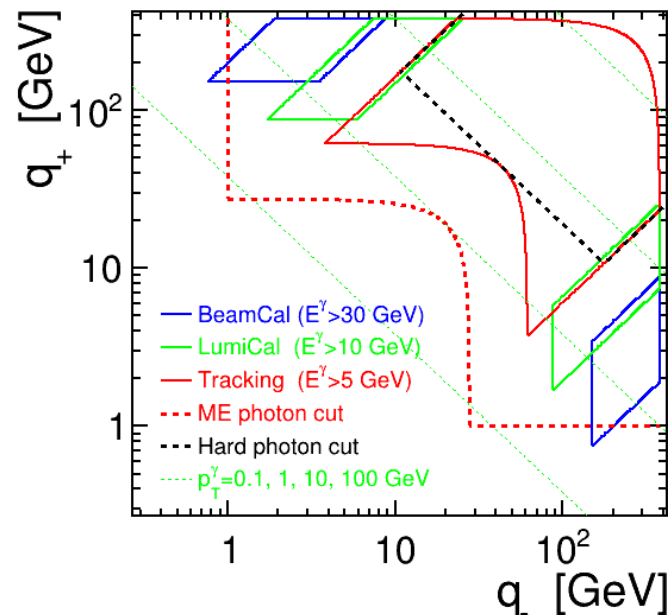
WHIZARD mono-photons

- Lets define variables q^+ and q^- calculated for each photon in the event. e^-/e^+ fixed as beam energies:

$$q_- = 2\sqrt{(E_{e^-} E_\gamma) \sin(\frac{\theta_\gamma}{2})}$$

$$q_+ = 2\sqrt{(E_{e^+} E_\gamma) \cos(\frac{\theta_\gamma}{2})}$$

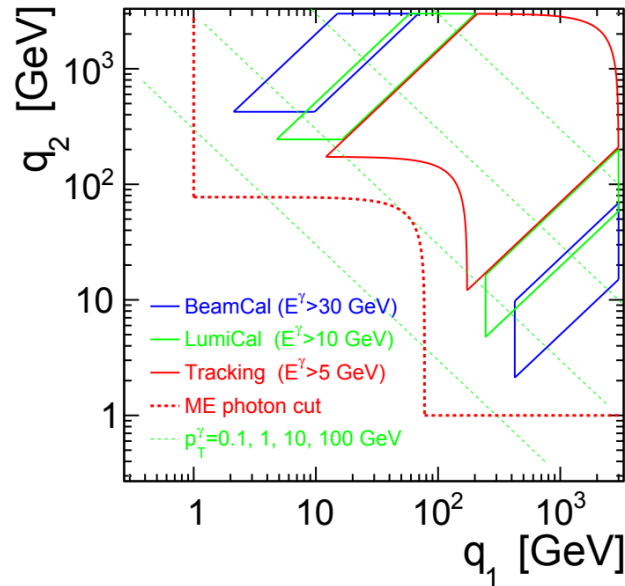
- Merging would be done for energy E^{\min} and q^{\min} : generated γ^{ME} should have q^+ and q^- above q^{\min} and energy above E^{\min}
- ISR photons accepted only below given E^{\min} and q^{\min} values
- Additional ‘hard photon’ selection: events with at least one γ having $p_T > 5 \text{ GeV}$ and $7^\circ < \theta < 173^\circ$



q^+/q^- phase space for \sqrt{s} of 380 GeV
(more cases in backup)

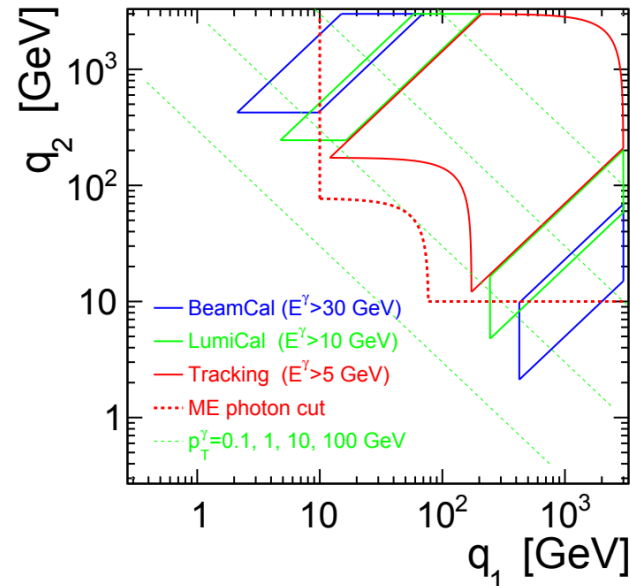
Merging ISR and ME photons

Using CLICdet design one can see that preferred value of q^{merge} is ~ 1 GeV:
no ISR photons reach the BeamCal



$$q^{\text{min}} = 1 \text{ GeV}$$

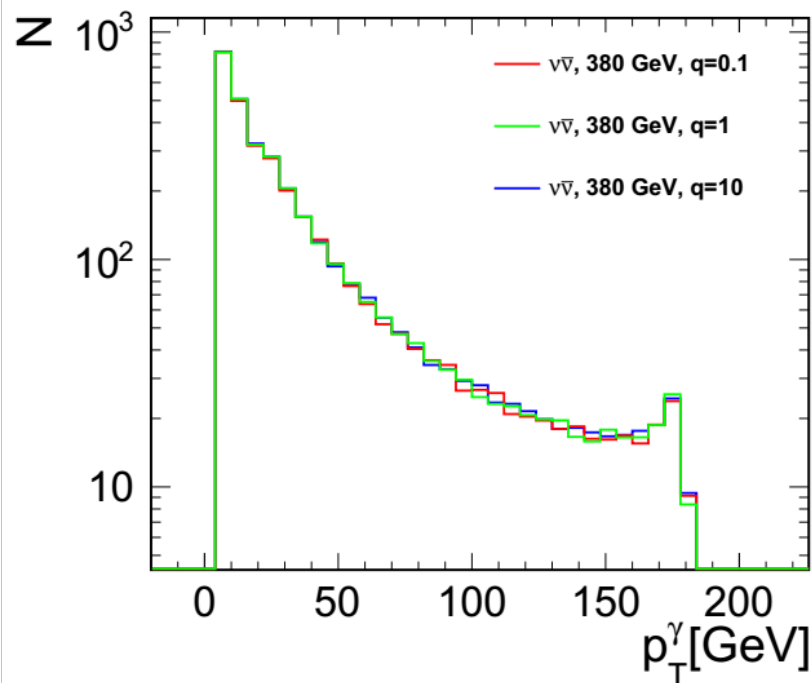
**Plots for
3 TeV
case**



$$q^{\text{min}} = 10 \text{ GeV}$$

Neutrino sample

$$e^+e^- \rightarrow \nu\bar{\nu} + N\gamma$$



Impact of q_{min} cut on cross section calculation: stable

$\sqrt{s}[\text{GeV}]$	$q_{min}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ [fb]				$\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ [fb] after ISR rejection
		$\nu\bar{\nu}$	$\nu\bar{\nu} + \gamma_{ME}$	$\nu\bar{\nu} + 2\gamma_{ME}$	$\nu\bar{\nu} + 3\gamma_{ME}$	
380	$q_{min}=0.1$	50900	16500	2600	220	50000
	$q_{min}=0.5$		12600	1600	110	50400
	$q_{min}=1$		10900	1200	55	50600
	$q_{min}=5$		7000	480	15	50700
	$q_{min}=10$		5300	270	7	50800

With ‘hard photon’ selection:

$\sqrt{s}[\text{GeV}]$	$q_{min}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ [fb]			$\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ [fb] after ISR rejection
		$\nu\bar{\nu} + \gamma_{ME}$	$\nu\bar{\nu} + 2\gamma_{ME}$	$\nu\bar{\nu} + 3\gamma_{ME}$	
380	$q_{min}=0.1$	3200	910	120	3000
	$q_{min}=0.5$	3200	670	55	3000
	$q_{min}=1$	3200	570	43	3100
	$q_{min}=5$	3200	340	10	3100
	$q_{min}=10$	3200	230	6	3100

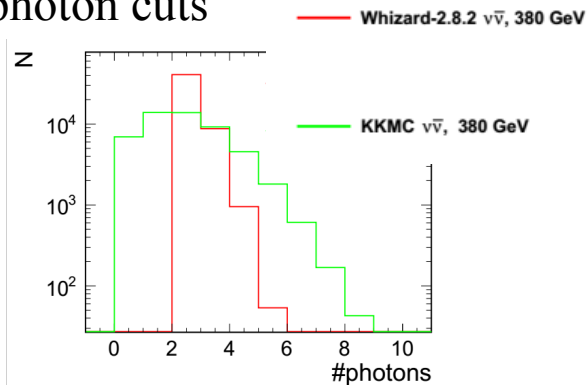
$\nu\bar{\nu}$ sample: WHIZARD vs KKMC



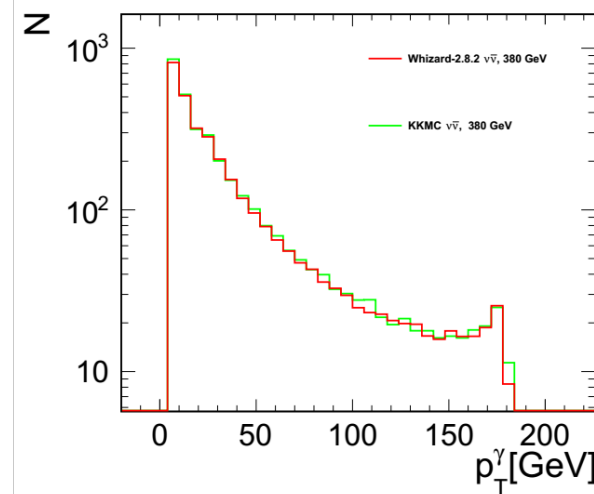
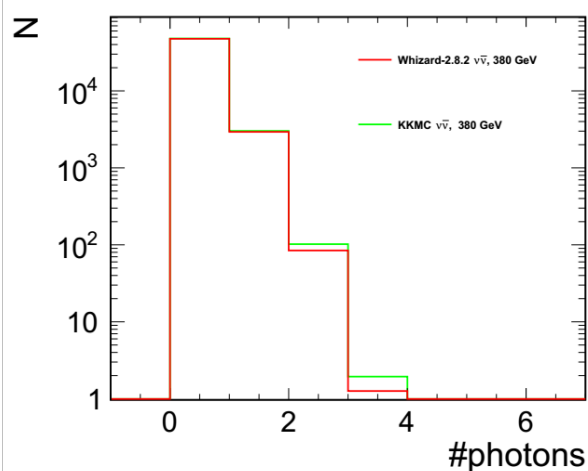
Comparison of WHIZARD and KKMC generators for neutrino sample at $\sqrt{s} = 380$ GeV

KKMC generator:
S. Jadach, B.F.L. Ward, Z. Was
Computer Physics Communications 130 (2000) 260–325

- Using generators with different ISR/FSR descriptions
- no photon cuts



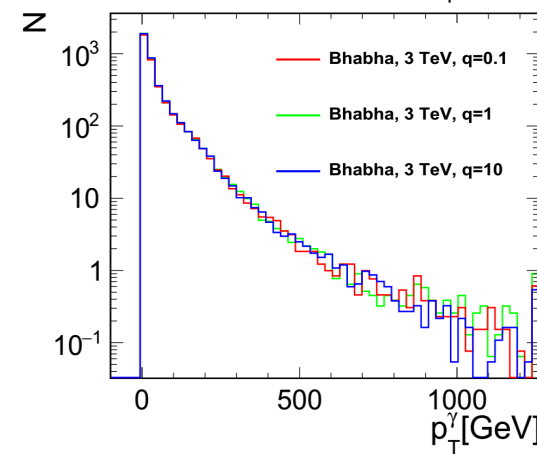
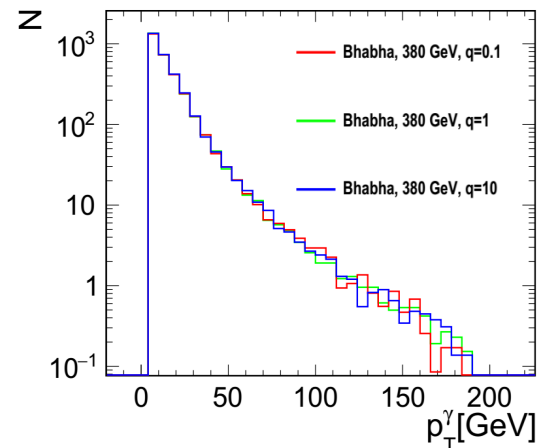
The same 'hard photon' selection imposed:



q^{\min} set to 1 GeV

Bhabha sample and ISR/ γ^{ME} merging

- Comparison of p_T spectra for energies of 380 GeV and 3 TeV for Bhabha background as a function of different merging parameters q^{min}
- Distributions show small impact on the shapes of the given spectra
- For further studies q^{min} value set to 1 GeV



New approach to DM searches with mono-photon signature

New approach to DM searches

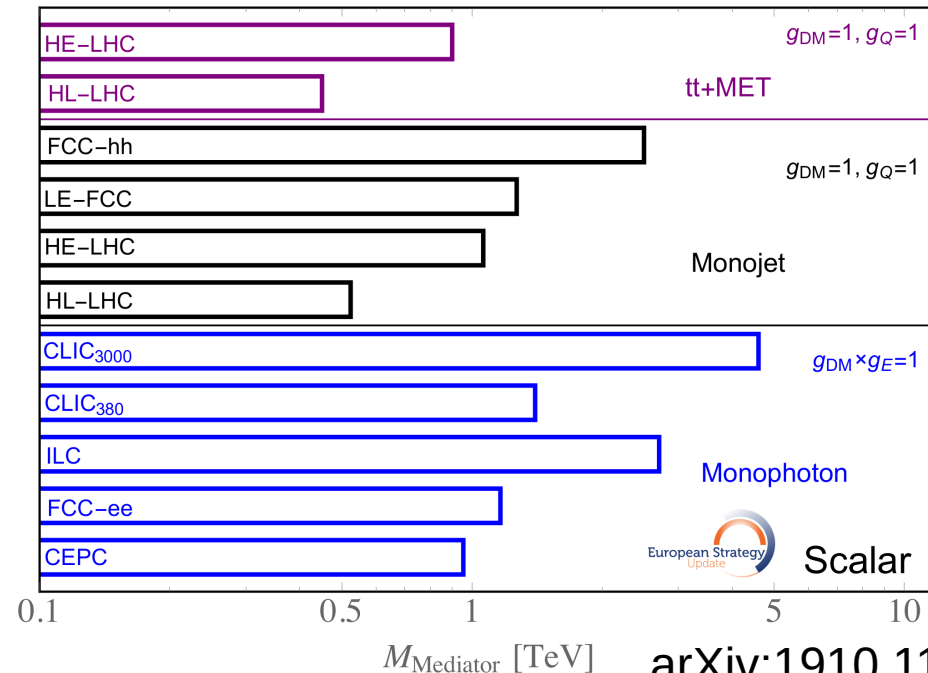
- Common approach:
setting limits on mediator mass for given coupling structure/values

- This approach is suitable,
if we can assume

$$M_{\text{med}} \gg \sqrt{s}$$

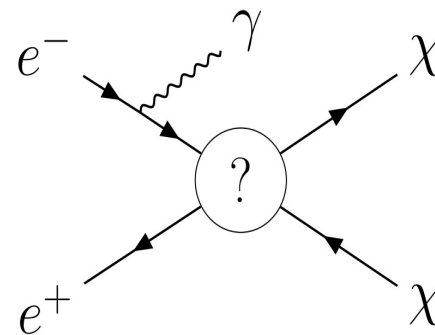
but this does not need
to be the case...

- Different approach needed
for mediator masses $\sim \sqrt{s}$



New approach to DM searches

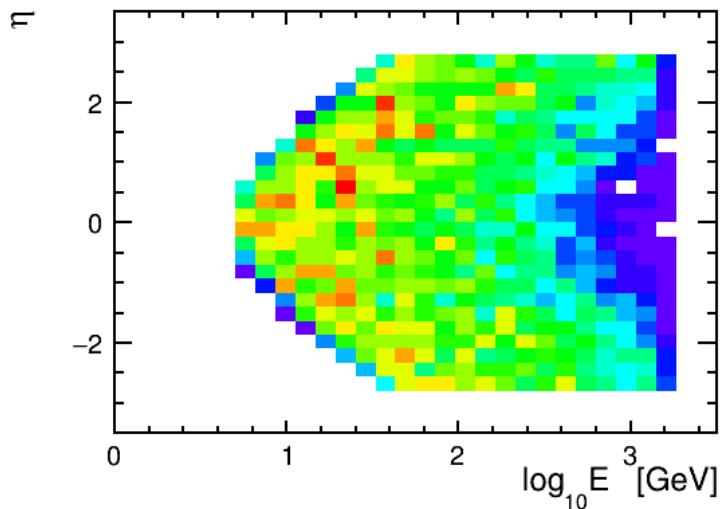
- Experimental approach:
set limits on the DM production cross section
 - as a function of the assumed **mediator mass and mediator width**
DM mass not relevant, if sufficiently light
 - depends on the assumed **coupling structure**,
but coupling values not relevant
 - most model independent limits
can be then interpreted in different models



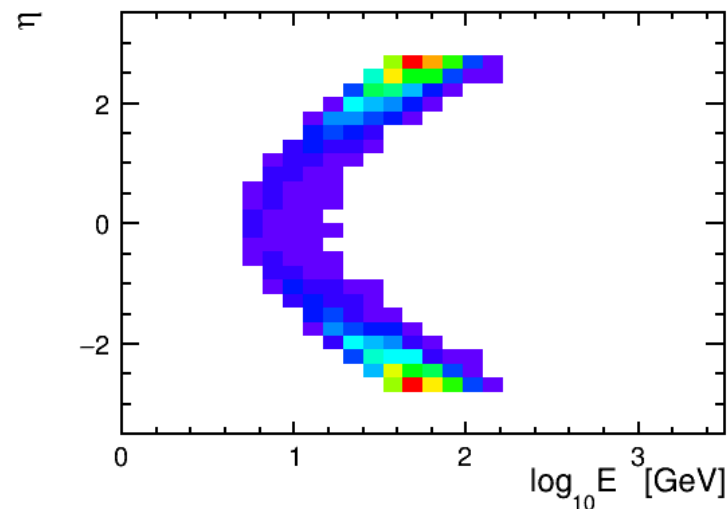
New approach to DM searches

- Experimental approach:
consider 2-D distribution of mono-photon event
- Main backgrounds (generator level cuts only) @ 3TeV CLIC:

$$e^+e^- \rightarrow \nu\nu$$



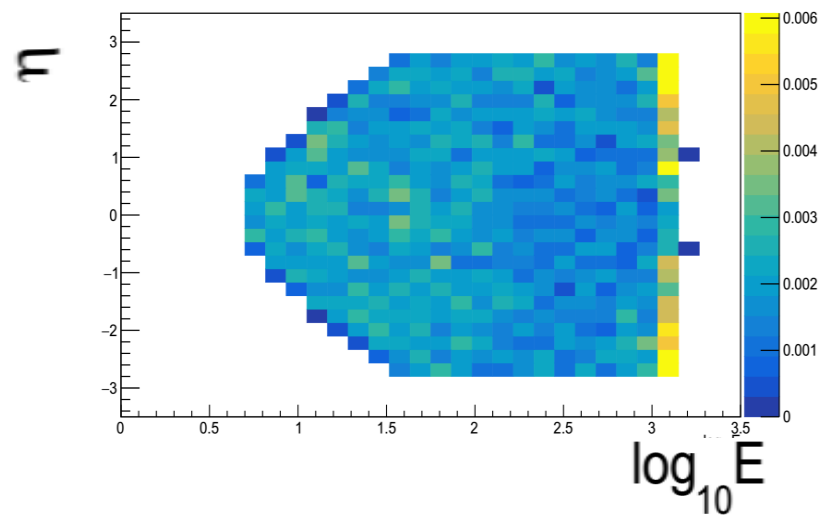
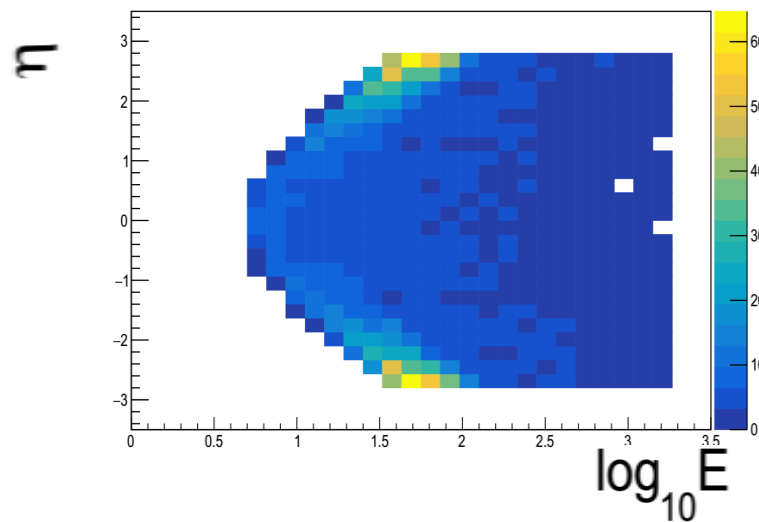
$$e^+e^- \rightarrow e^+e^-$$



New approach to DM searches

- Experimental approach:
consider 2-D distribution of mono-photon event

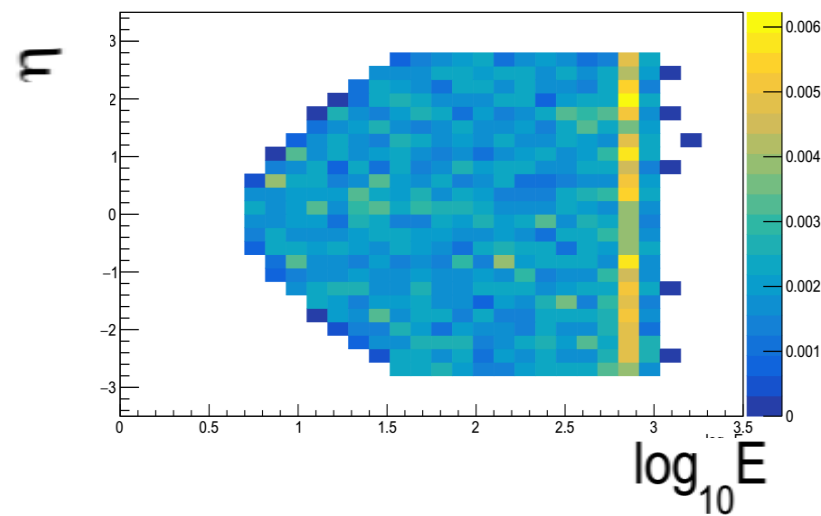
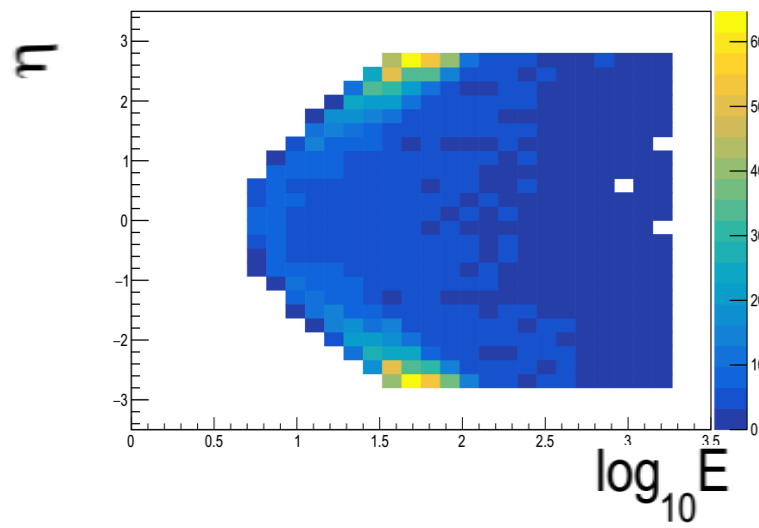
- Main backgrounds vs signal @ 3TeV CLIC
 $e^+e^- \rightarrow \nu\bar{\nu}$ & $e^+e^- \rightarrow e^+e^-$ $M_{\text{DM}}=50\text{GeV}$ $M_{\text{med}}=1\text{TeV}$ $\Gamma_{\text{med}}=50\text{GeV}$



New approach to DM searches

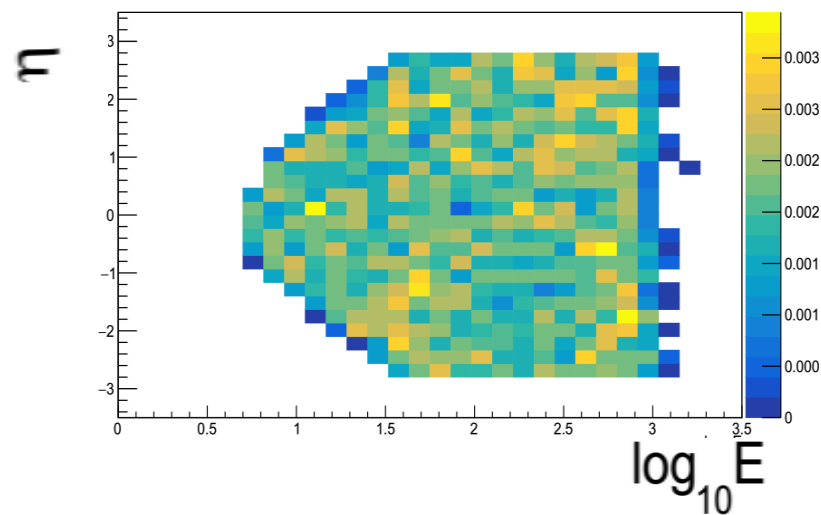
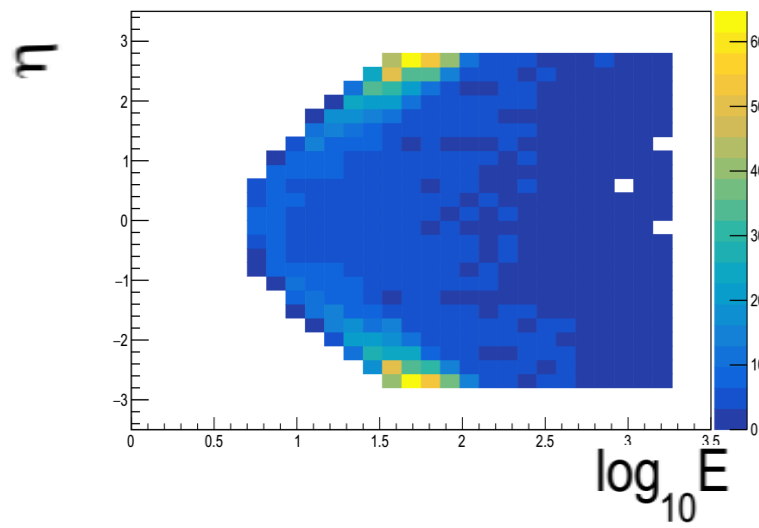
- Experimental approach:
consider 2-D distribution of mono-photon event

- Main backgrounds vs signal @ 3TeV CLIC
 $e^+e^- \rightarrow \nu\bar{\nu}$ & $e^+e^- \rightarrow e^+e^-$ $M_{\text{DM}}=50\text{GeV}$ $M_{\text{med}}=2\text{TeV}$ $\Gamma_{\text{med}}=50\text{GeV}$



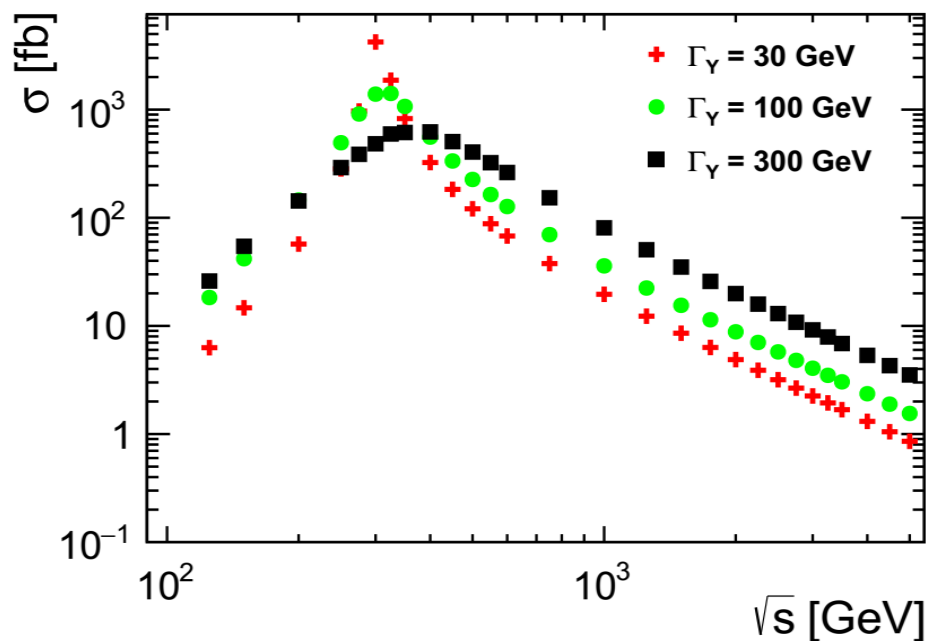
New approach to DM searches

- Experimental approach:
consider 2-D distribution of mono-photon event
- Main backgrounds vs signal @ 3TeV CLIC
 $e^+e^- \rightarrow \nu\bar{\nu}$ & $e^+e^- \rightarrow e^+e^-$ $M_{\text{DM}}=50\text{GeV}$ $M_{\text{med}}=2\text{TeV}$ $\Gamma_{\text{med}}=500\text{GeV}$



Cross sections

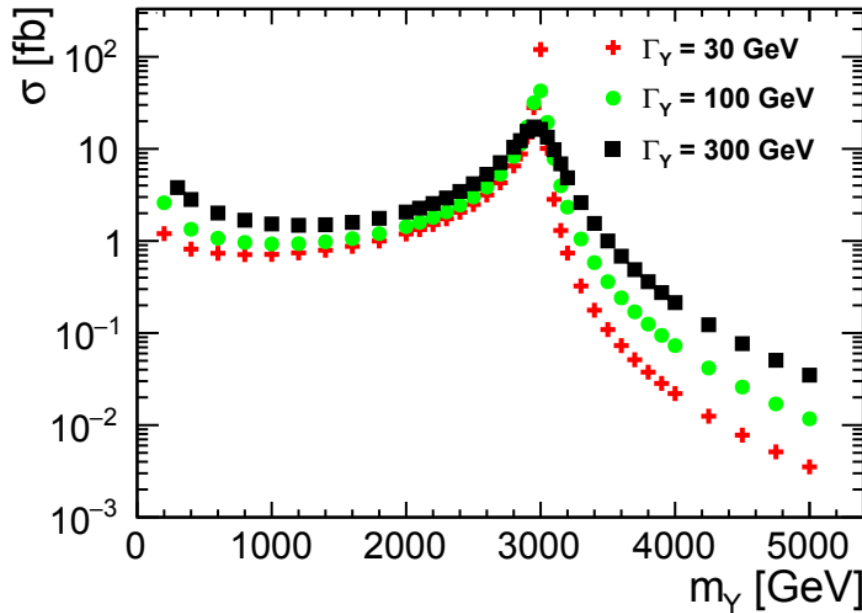
- For now one scenario studied
 - Dirac Dark Matter of 50 GeV
 - Vector mediator with different masses and widths
- Example of cross section behaviour for different collision energies



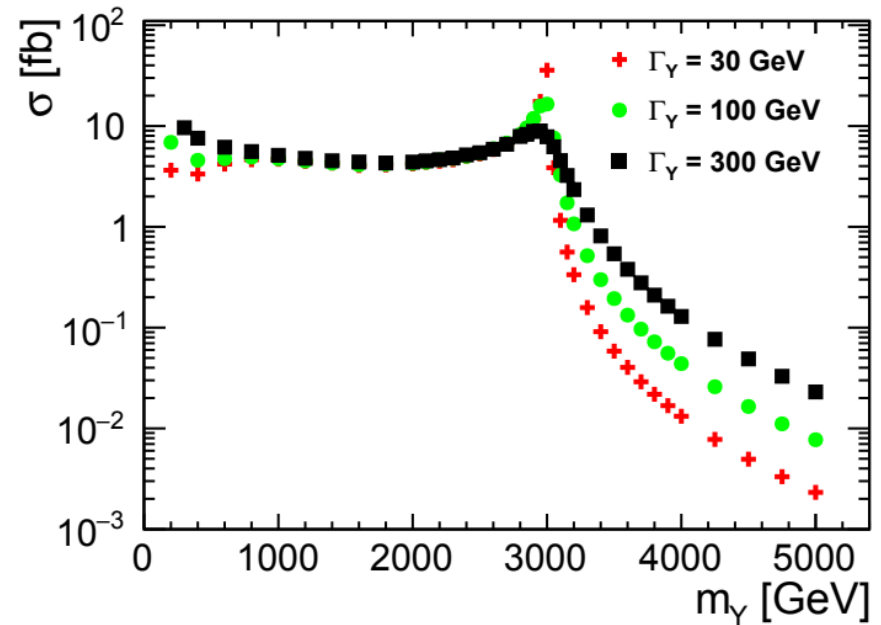
Assumed mediator mass
of 300 GeV

Cross sections as a function of mediator mass

Without luminosity spectra



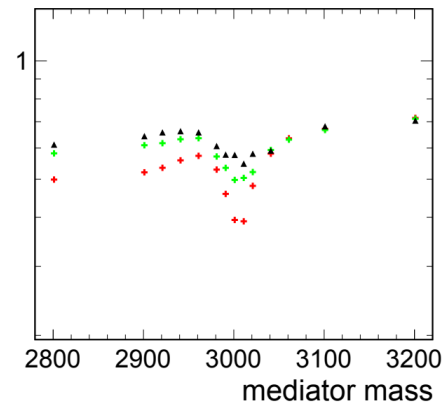
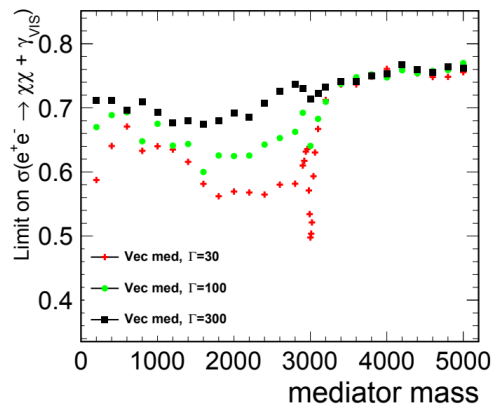
With luminosity spectra



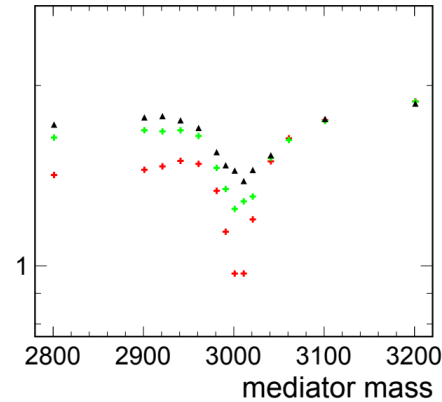
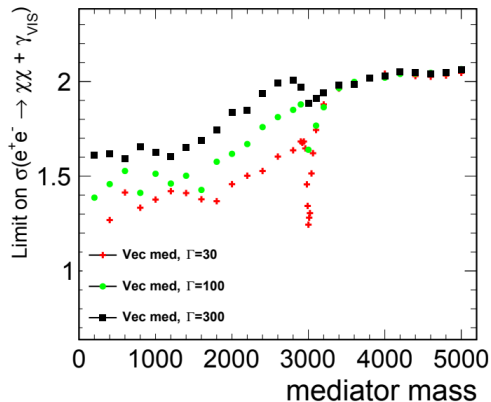
For 3 TeV collision energy @ CLIC

Prospects on DM searches

- Limits calculated with CLs method from 2D photon distributions
 - Assumed luminosity: 1 ab^{-1}
 - 3 TeV CLIC
 - Require only one photon in the detector (LumiCal and BeamCal veto)
 - Consider electron beam polarisation
 - Limits not corrected for acceptance
- Stable cross section limits for a wide range of mediator masses
- Limits improve for narrow resonances



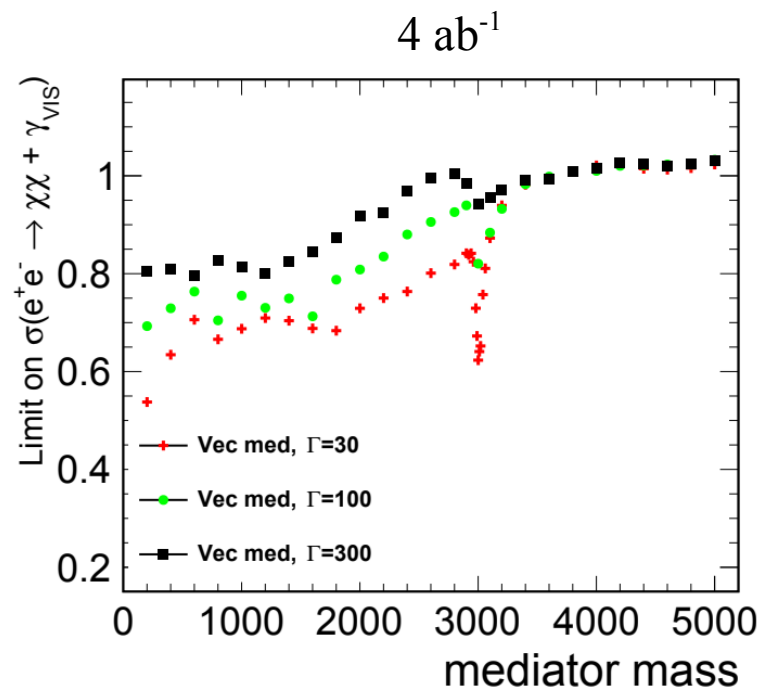
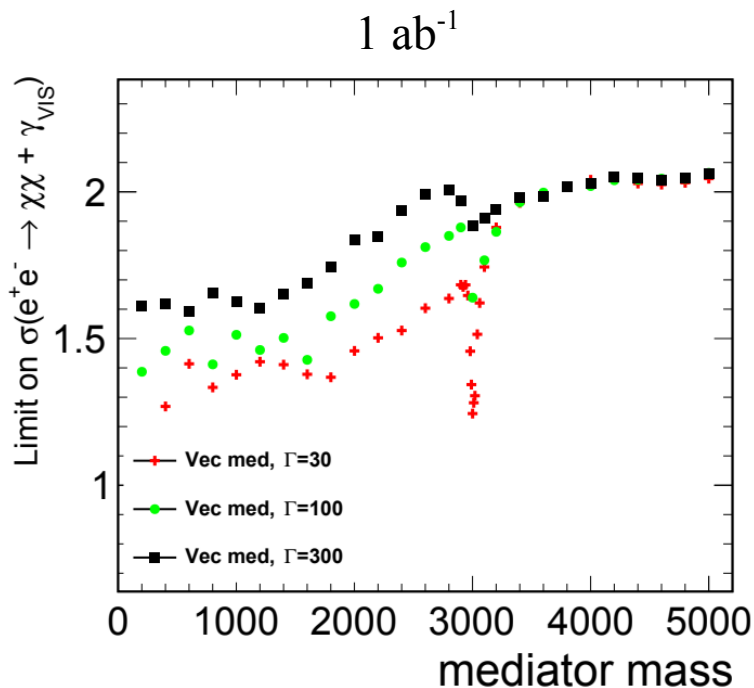
Positive
pol.



Negative
pol.

Prospects on DM searches

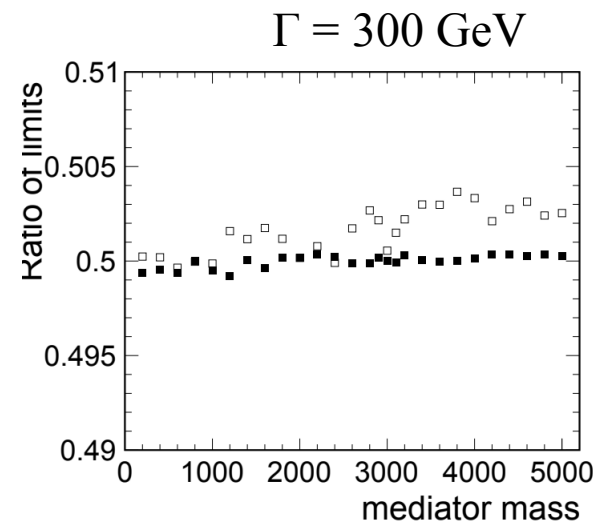
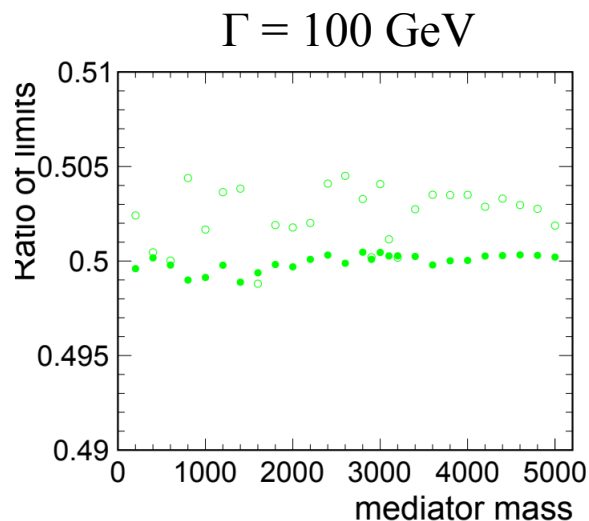
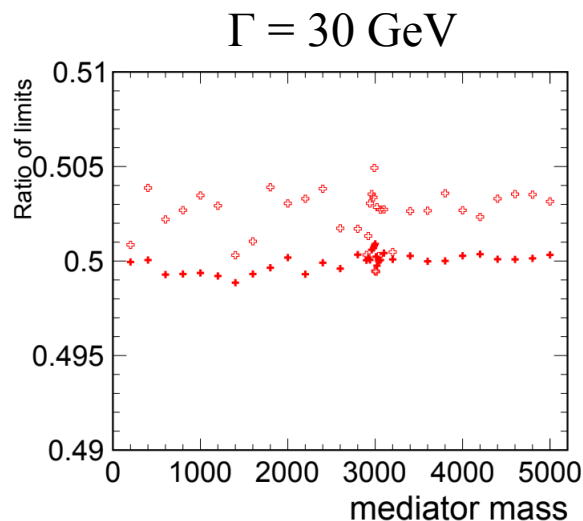
Impact of assumed luminosity Negative polarisation



Ratios of limits

Ratios of limits calculated for assumed luminosity of 4 ab^{-1} with respect to 1 ab^{-1}

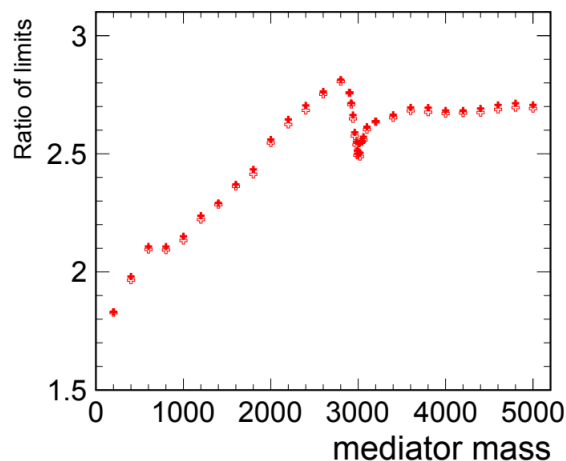
positive (open symbols) and negative (full) polarisation



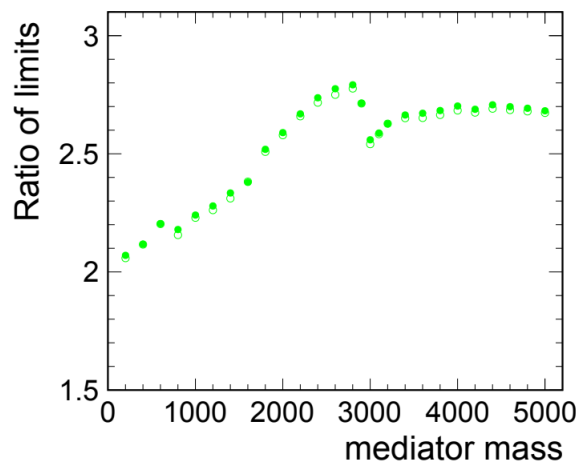
Ratios of limits

Ratios of limits calculated for negative polarisations to positive polarisation limits, for assumed luminosity of 1 ab^{-1}

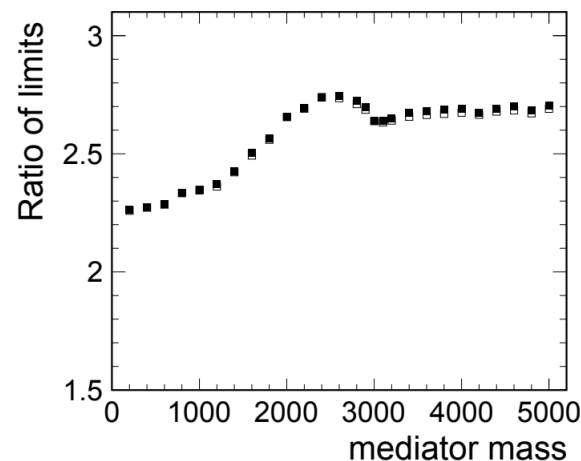
$\Gamma = 30 \text{ GeV}$



$\Gamma = 100 \text{ GeV}$



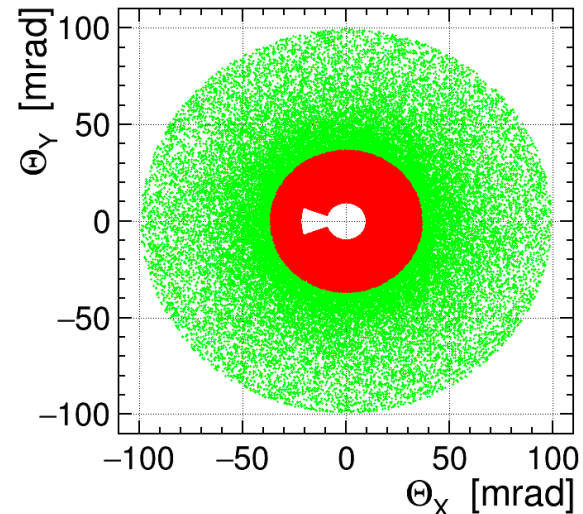
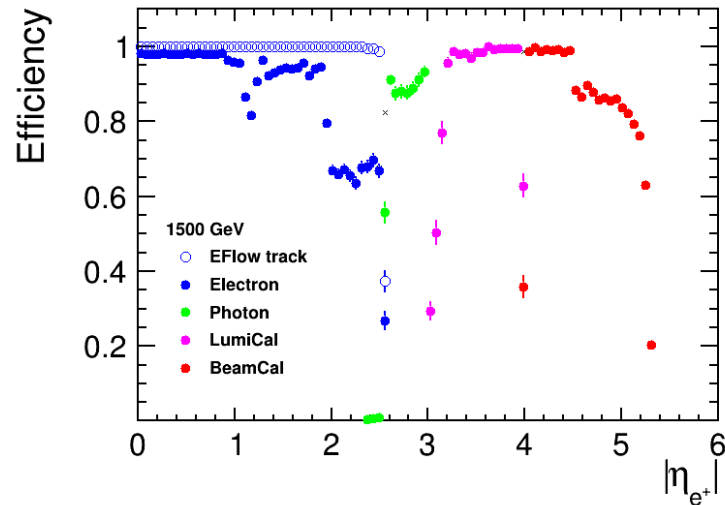
$\Gamma = 300 \text{ GeV}$



For the same luminosity, positive electron beam polarisation improves the sensitivity to DM production by about a factor of two

Prospects on DM searches

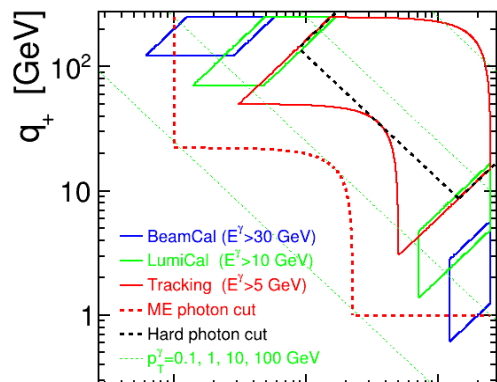
- All results so far based on generator level cuts
- Delphes model of CLICdet has been modified to allow for better simulation of forward calorimeter acceptance



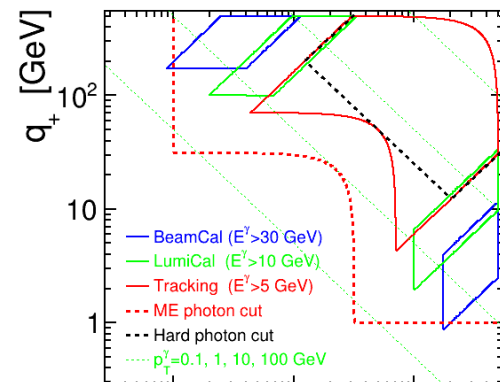
- Presented merging procedure allows to treat two main backgrounds of mono-photon analysis with the same approach and obtain stable cross section calculations also for the Bhabha background. **This procedure is general enough to be used for arbitrary BSM scenarios with mono-photon signature**
- Working merging procedure allows to obtain model independent limits on Dark Matter production cross sections for different experimental setups
 - **Only coupling structure needed**
 - **Impact of the mediator width on signal-to-background separation**
- Scans of upper limits for a single DM scenario
 - **Light mediators clearly distinguishable**
 - **Positive polarisation beam setup gives better limits**
- Preliminary results based on generator level cuts, study based on Delphes detector model will follow

Backup

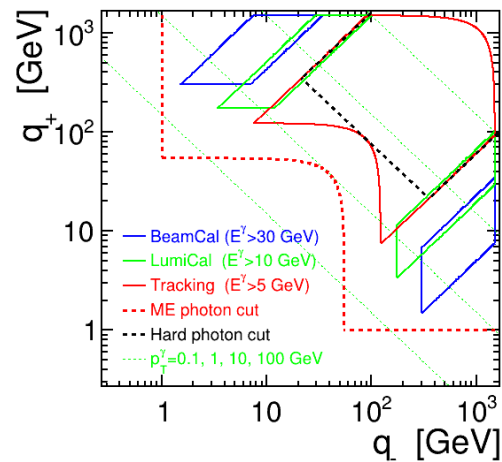
$\sqrt{s} = 250 \text{ GeV}$



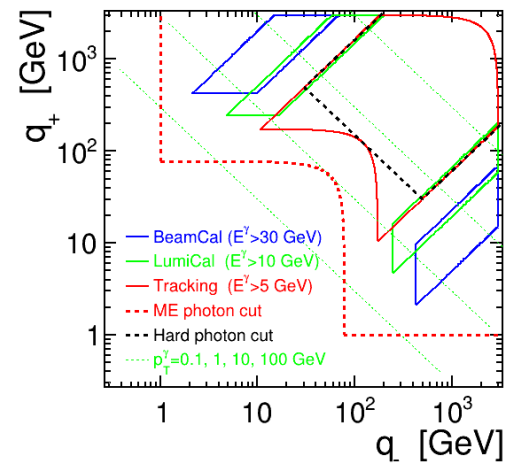
$\sqrt{s} = 500 \text{ GeV}$



$\sqrt{s} = 1.5 \text{ TeV}$



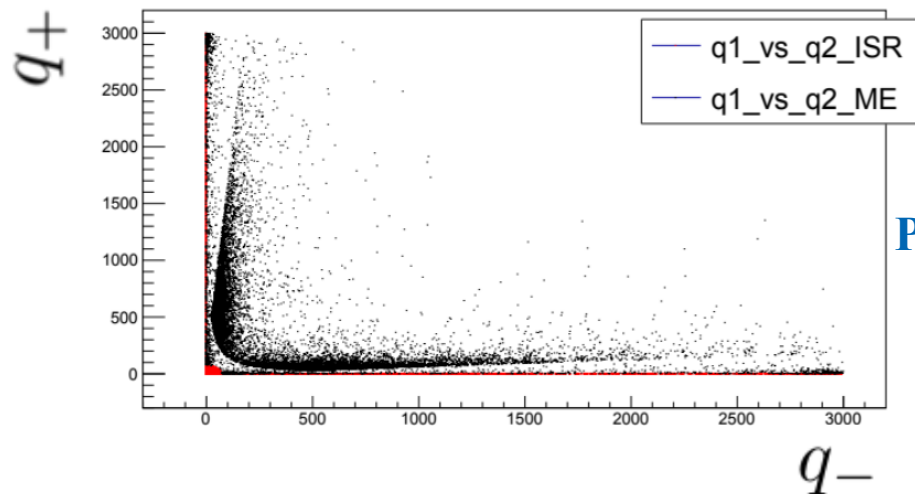
$\sqrt{s} = 3 \text{ TeV}$



Merging ISR and γ^{ME} for sample with $P_t > 5 \text{ GeV}$ cut

$$q^{\text{merge}} = 1 \text{ GeV}$$

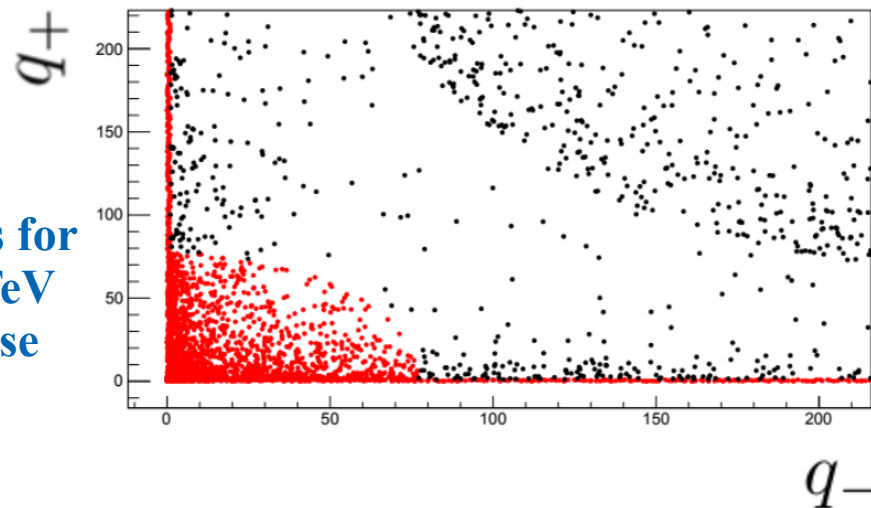
Bhabha sample



Plots for
3 TeV
case

Red – γ ISR
Black – γ^{ME}

Zoom



$$q^{\text{merge}} = 1 \text{ GeV}$$

Merging ISR and γ^{ME}



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Dependency of the **neutrino** sample on the merging parameter q^{min}

$\sqrt{s}[\text{GeV}]$	$q_{\text{min}}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow \nu\bar{\nu})$				$\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ after ISR selection
		$\nu\bar{\nu}$	$\nu\bar{\nu} + \gamma_{\text{ME}}$	$\nu\bar{\nu} + 2\gamma_{\text{ME}}$	$\nu\bar{\nu} + 3\gamma_{\text{ME}}$	
380	$q_{\text{min}}=0.1$	50849	16500	2600	220	50000
	$q_{\text{min}}=0.5$		12600	1600	110	50400
	$q_{\text{min}}=1$		10900	1200	55	50600
	$q_{\text{min}}=5$		7000	480	15	50700
	$q_{\text{min}}=10$		5300	270	7	50800
	$q_{\text{min}}=50$		1400	21	0	50500
3000	$q_{\text{min}}=0.1$	53921	26200	6300	970	52200
	$q_{\text{min}}=0.5$		20700	4000	410	52300
	$q_{\text{min}}=1$		18500	3100	180	52400
	$q_{\text{min}}=5$		13000	1600	74	52700
	$q_{\text{min}}=10$		10700	1100	57	52600
	$q_{\text{min}}=50$		5200	260	2	52800

Merging ISR and γ^{ME}



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Dependency of the **neutrino** sample with a '**hard**' γ^{ME} on the merging parameter q^{min}

$\sqrt{s}[\text{GeV}]$	$q_{\text{min}}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow \nu\bar{\nu})$			$\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ after ISR selection
		$\nu\bar{\nu} + \gamma_{\text{ME}}$	$\nu\bar{\nu} + 2\gamma_{\text{ME}}$	$\nu\bar{\nu} + 3\gamma_{\text{ME}}$	
380	$q_{\text{min}}=0.1$	3200	910	120	3000
	$q_{\text{min}}=0.5$	3200	670	55	3000
	$q_{\text{min}}=1$	3200	570	43	3100
	$q_{\text{min}}=5$	3200	340	10	3100
	$q_{\text{min}}=10$	3200	230	6	3100
	$q_{\text{min}}=50$	1400	22	0	1400
3000	$q_{\text{min}}=0.1$	4200	1900	260	3700
	$q_{\text{min}}=0.5$	4200	1500	220	3900
	$q_{\text{min}}=1$	4200	1300	140	3900
	$q_{\text{min}}=5$	4200	850	66	3900
	$q_{\text{min}}=10$	4200	680	31	3900
	$q_{\text{min}}=50$	4000	240	2	3700

Merging ISR and γ^{ME}



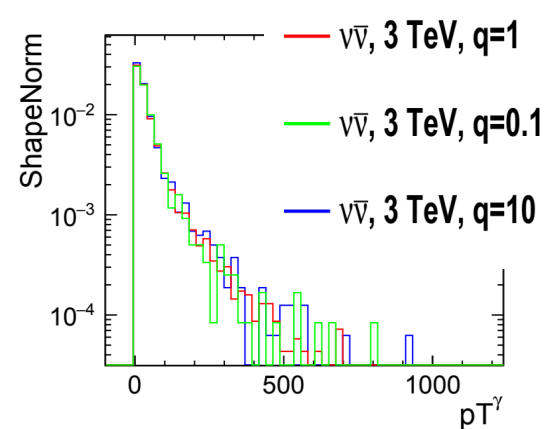
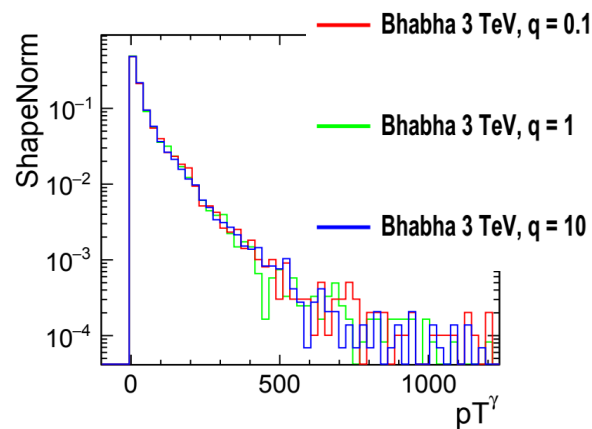
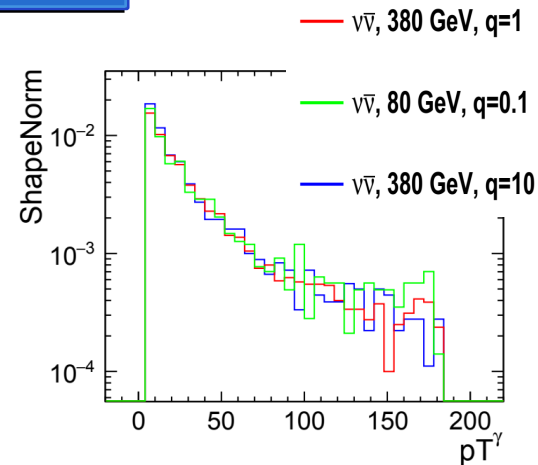
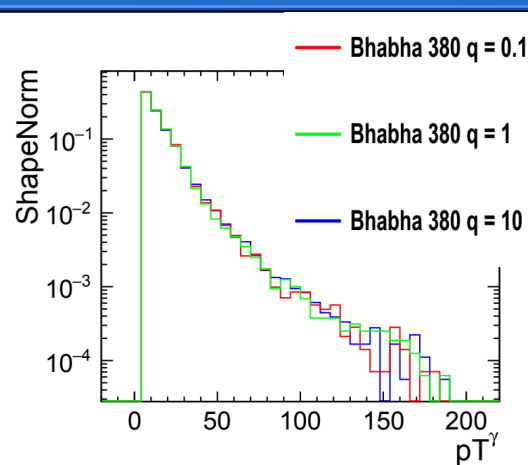
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Dependency of the **Bhabha** sample with a ‘hard’ γ^{ME} on the merging parameter q^{min}
Bhabha background requires ‘hard’-photon approach to have the cross section calculations
even converge

$\sqrt{s}[\text{GeV}]$	$q_{\text{min}}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow e^+e^- + N\gamma)$			$\sigma(e^+e^- \rightarrow e^+e^- + N\gamma)$ after ISR selection
		$e^+e^- + \gamma_{\text{ME}}$	$e^+e^- + 2\gamma_{\text{ME}}$	$e^+e^- + 3\gamma_{\text{ME}}$	
380	$q_{\text{min}}=0.1$	139500	29100	650	119800
	$q_{\text{min}}=0.5$	139200	20700	630	124900
	$q_{\text{min}}=1$	139300	17200	540	125600
	$q_{\text{min}}=5$	139700	9200	310	129400
	$q_{\text{min}}=10$	139400	4900	140	130000
	$q_{\text{min}}=50$	11200	150	2	11000
3000	$q_{\text{min}}=0.1$	8800	3600	74	6200
	$q_{\text{min}}=0.5$	9000	2800	73	6800
	$q_{\text{min}}=1$	8800	2700	33	7000
	$q_{\text{min}}=5$	8900	1800	30	7500
	$q_{\text{min}}=10$	8800	1200	22	7300
	$q_{\text{min}}=50$	7000	280	7	6000

Merging ISR/ γ ^{ME}

- Comparison of pT spectra for energies of 380 GeV and 3 TeV for neutrino and Bhabha backgrounds as a function of different merging parameters q^{\min}
- Distributions show small impact on the shapes of the given spectra
- For further studies q^{\min} value set to 1 GeV



More on σ and collision energy



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Backgrounds σ as a function of collision energy

Bhabha samples

$\sqrt{s}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow e^+e^- + N\gamma)$			$\sigma(e^+e^- \rightarrow e^+e^- + N\gamma)$ after ISR selection
	$e^+e^- + \gamma_{ME}$	$e^+e^- + 2\gamma_{ME}$	$e^+e^- + 3\gamma_{ME}$	
240	233400	24900	480	216000
250	223600	24800	730	208100
380	139300	17200	540	125600
500	99300	14200	330	88800
1000	39900	7200	270	34000
1500	22900	4700	140	18800
3000	8800	2700	33	7000

neutrino samples

$\sqrt{s}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow \nu\bar{\nu})$			$\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ after ISR selection
	$\nu\bar{\nu} + \gamma_{ME}$	$\nu\bar{\nu} + 2\gamma_{ME}$	$\nu\bar{\nu} + 3\gamma_{ME}$	
240	4200	730	37	4100
250	4100	700	48	3900
380	3200	570	43	3100
500	3100	590	48	3000
1000	3500	780	74	3300
1500	3800	940	98	3600
3000	4200	1300	140	3900

$\sqrt{s}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow \nu\bar{\nu})$				$\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ after ISR selection
	$\nu\bar{\nu}$	$\nu\bar{\nu} + \gamma_{ME}$	$\nu\bar{\nu} + 2\gamma_{ME}$	$\nu\bar{\nu} + 3\gamma_{ME}$	
240	53900	12600	1300	64	53800
250	53300	12200	1300	60	53000
380	50900	10900	1200	55	50600
500	51200	11300	1200	75	51200
1000	52800	13600	1700	120	52200
1500	53300	15400	2200	170	52100
3000	53900	18500	3100	180	52400

Merging ISR/ γ^{ME} and beamstrahlung



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Note: values in [fb], rounded differently for each order of mag.

Impact of beamstrahlung (BS in tables) on backgrounds' cross sections:

Cross sections for **neutrino**
background samples are **not affected**
by switching on the beam spectra

Beam spectra crucial for 3TeV Bhabha background

Bhabha

$\sqrt{s}[\text{GeV}]$	$q_{\min}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow e^+e^- + N\gamma)$				$\sigma(e^+e^- \rightarrow e^+e^- + N\gamma)$ after ISR selection
		$e^+e^- + \gamma_{\text{ME}}$	$e^+e^- + 2\gamma_{\text{ME}}$	$e^+e^- + 3\gamma_{\text{ME}}$	$e^+e^- + 4\gamma_{\text{ME}}$	
380	noISR+noBS	128100	16700	1100	2	145900
	BS	407600	29600	640	4	437800
	ISR	139200	17600	520	3	125100
	BS+ISR	144900	18200	530	0	131700
3000	noISR+noBS	5800	2200	77	0	8100
	BS	221600	15100	88	2	236800
	ISR	8900	2500	74	0	6900
	BS+ISR	20400	3800	100	0	15900

Neutrinos

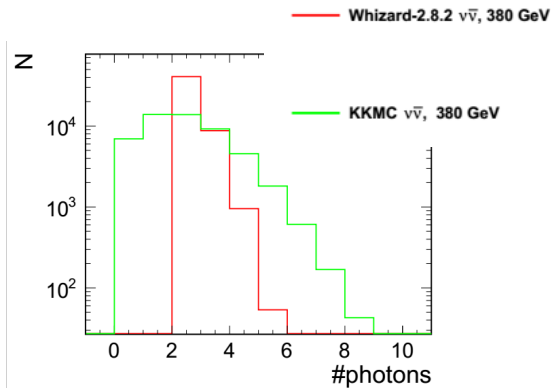
$\sqrt{s}[\text{GeV}]$	$q_{\min}[\text{GeV}]$	Whizard-2.8 $\sigma(e^+e^- \rightarrow \nu\bar{\nu})$				$\sigma(e^+e^- \rightarrow \nu\bar{\nu})$ after ISR selection
		$\nu\bar{\nu} + \gamma_{\text{ME}}$	$\nu\bar{\nu} + 2\gamma_{\text{ME}}$	$\nu\bar{\nu} + 3\gamma_{\text{ME}}$	$\nu\bar{\nu} + 4\gamma_{\text{ME}}$	
380	noISR+noBS	3000	550	48	1	3600
	BS	3000	560	45	1	3600
	ISR	3200	570	48	0	3100
	BS+ISR	3200	580	43	2	3100
3000	noISR+noBS	4300	1300	150	4	5700
	BS	4100	1200	130	1	5400
	ISR	4200	1300	140	2	3900
	BS+ISR	4000	1200	130	1	3800

$\nu\bar{\nu}$ sample: WHIZARD vs KKMC

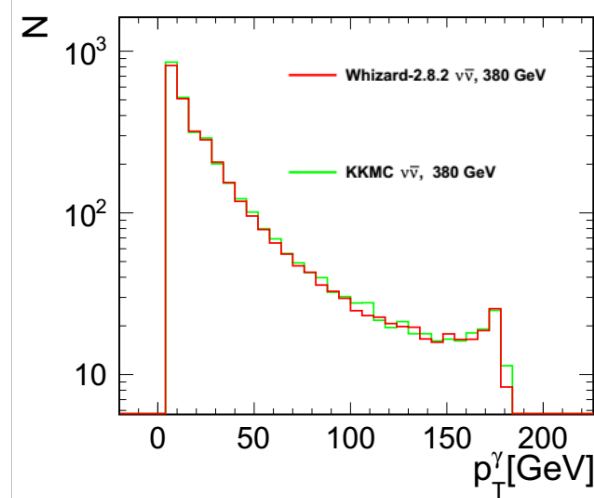
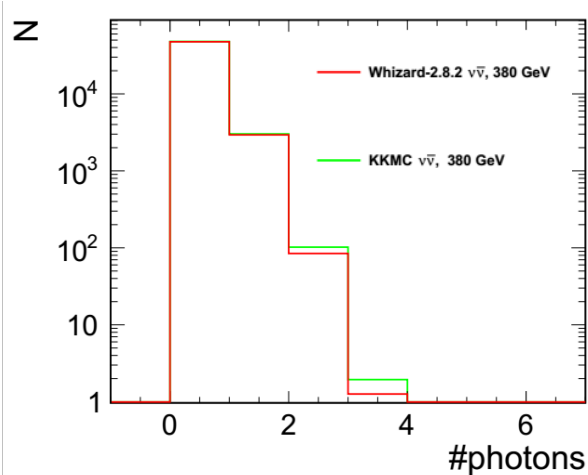
Comparison of WHIZARD and KKMC generators for neutrino sample at $\sqrt{s} = 380$ GeV

KKMC generator:
S. Jadach, B.F.L. Ward, Z. Was
Computer Physics Communications 130 (2000) 260–325

- Using generators with different ISR/FSR descriptions
- no photon cuts



The same ‘hard photon’ selection imposed:



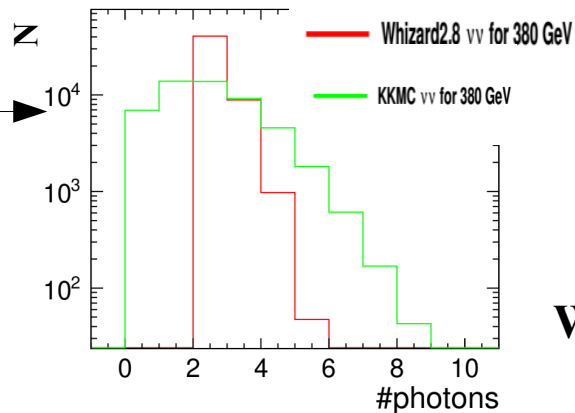
q^{\min} set to 1 GeV

Comparison with KKMC

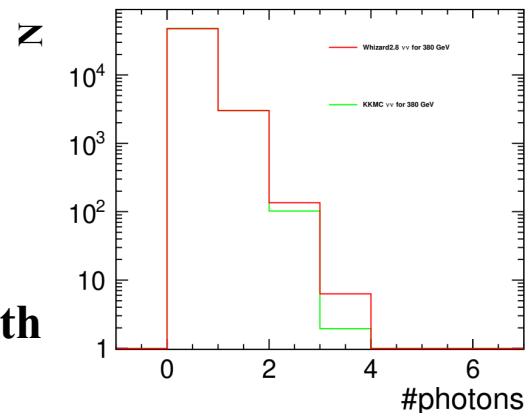


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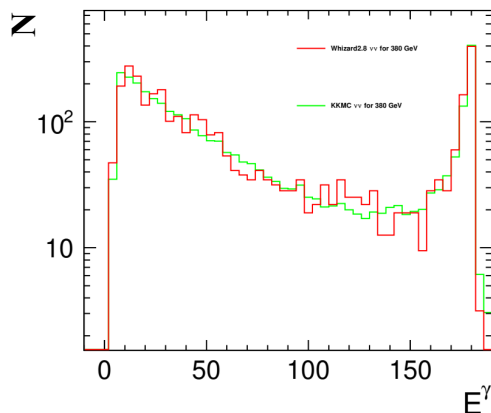
**Before
selection**



**Whizard sample with
basic merging**

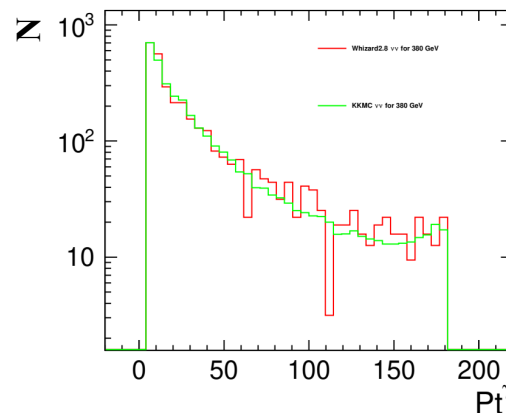


**After
selection**



Selecting photons:

$$p_T > 5 \text{ GeV} \\ 7^\circ < \theta < 173^\circ$$

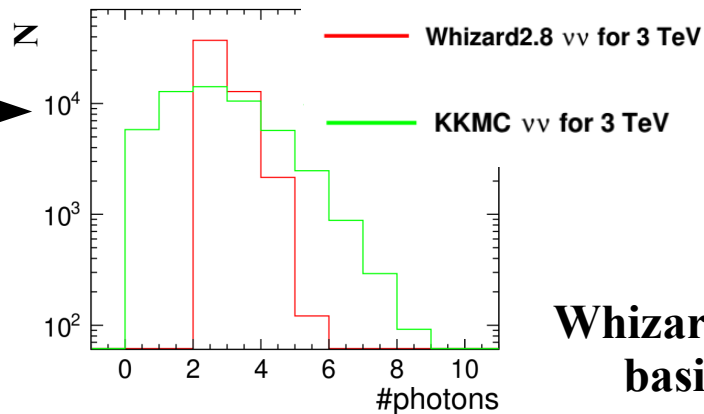


Comparison with KKMC

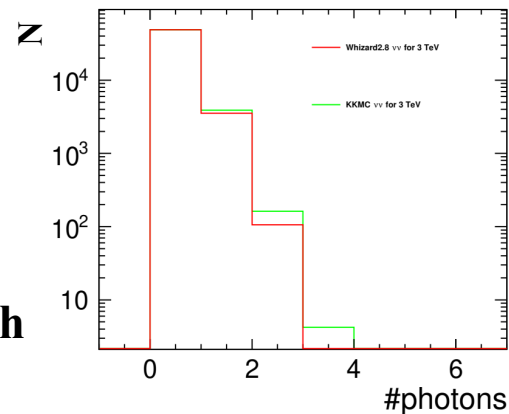


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WARSZAWSKI

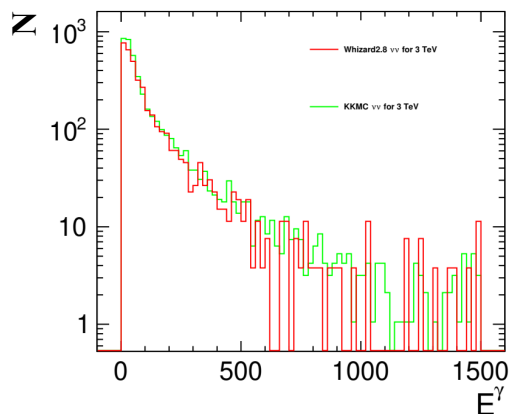
Before
selection



Whizard sample with
basic merging



After
selection



Selecting photons:

$$p_T > 5 \text{ GeV}$$
$$7^\circ < \theta < 173^\circ$$

