

Can modifications of the standard Λ CDM model solve the Hubble tension?

NIKKI ARENDSE
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Beyond Standard Model: From Theory to Experiment (BSM-2021)

BSM

**March 29-
April 2, 2021**



Sponsored by Letters in High Energy Physics (LHEP)

One of the strongest hints for new physics: Hubble tension

HUBBLE TENSION



CMB

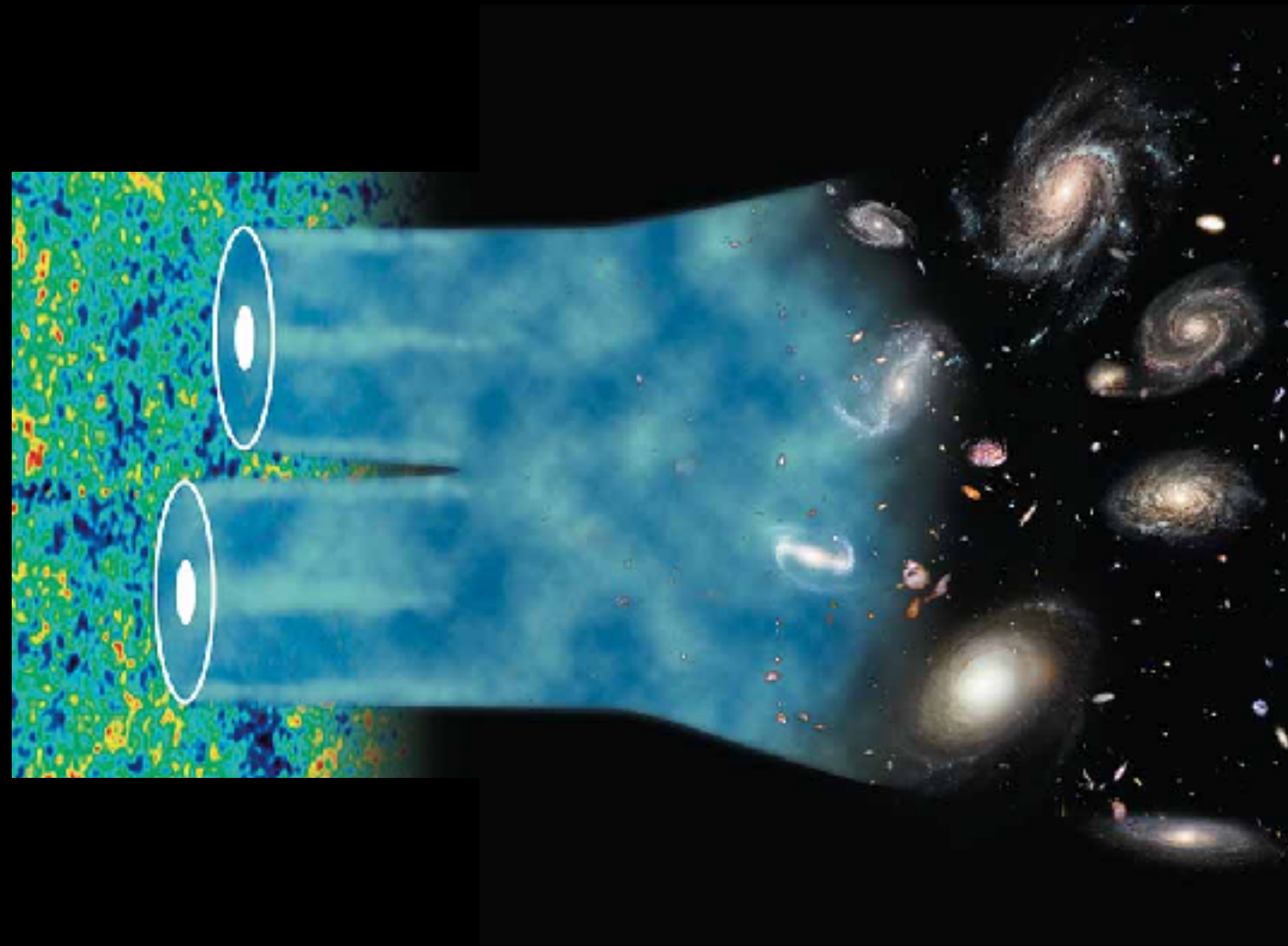
local

Expansion rate:

67.4

± 0.5 km/s/Mpc

from Planck



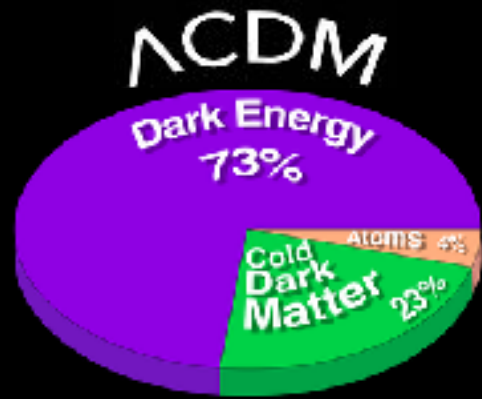
Expansion rate:

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from Cepheids (SH0ES)

HUBBLE TENSION



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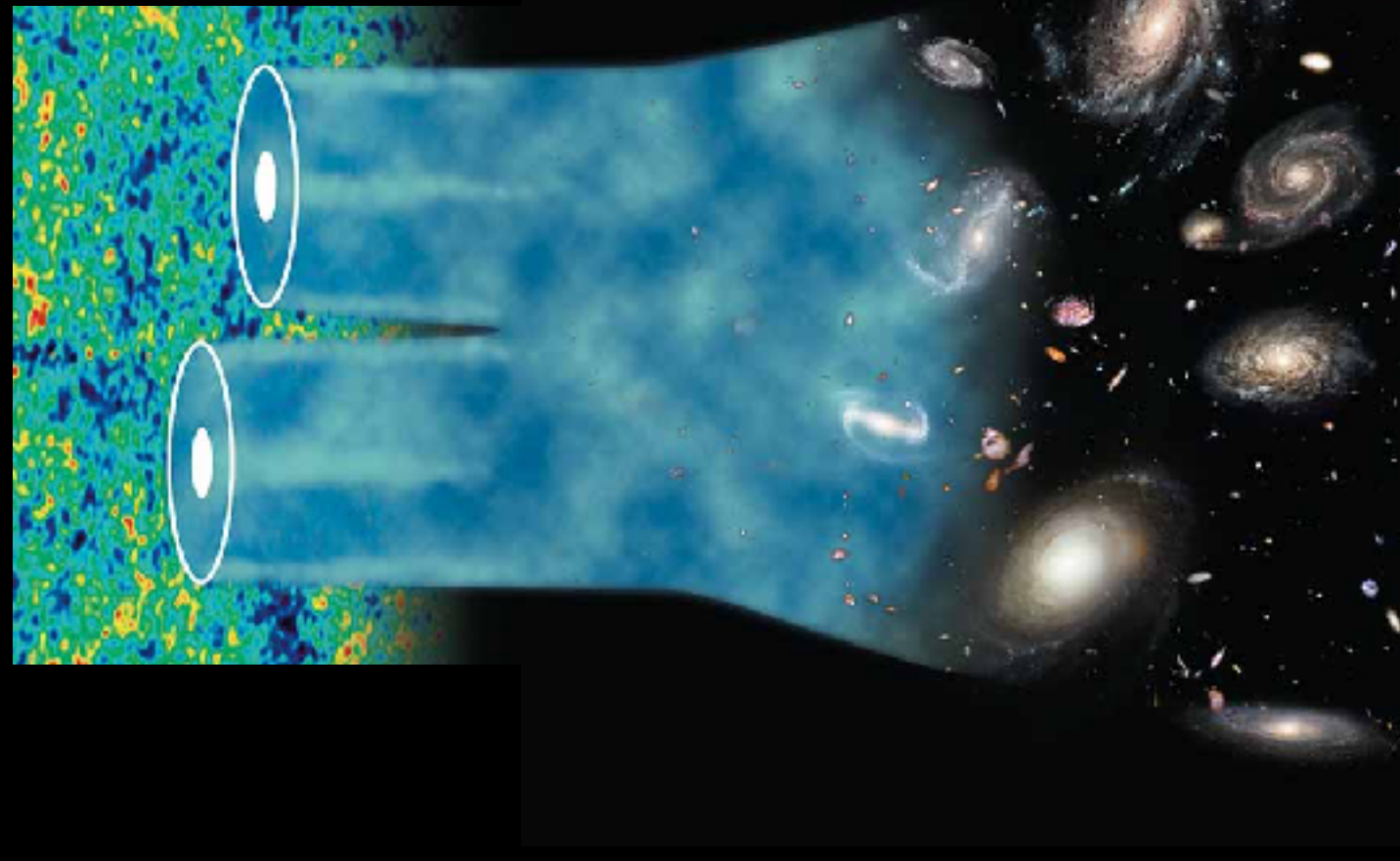
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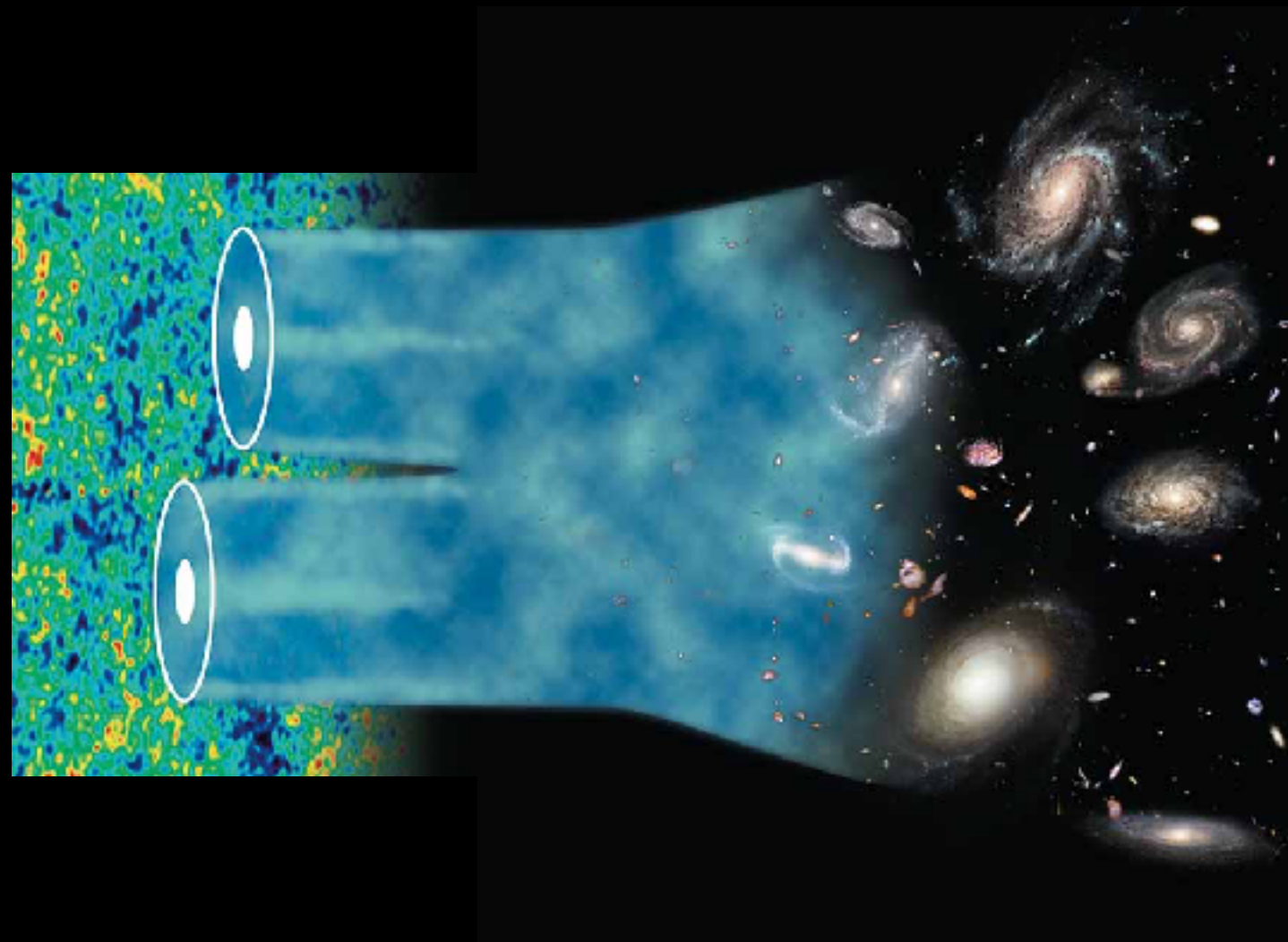
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Not just about H_0 !

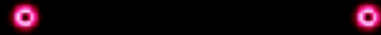
Another important cosmological parameter is the

SOUND HORIZON SCALE (R_s)

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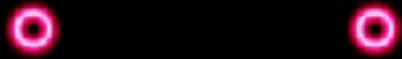


Over-densities

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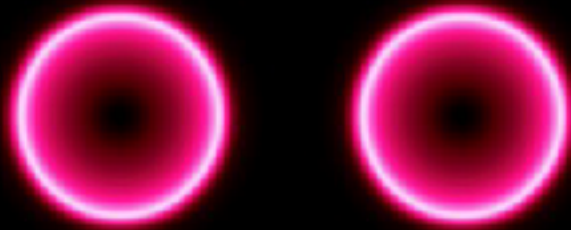


Sound waves

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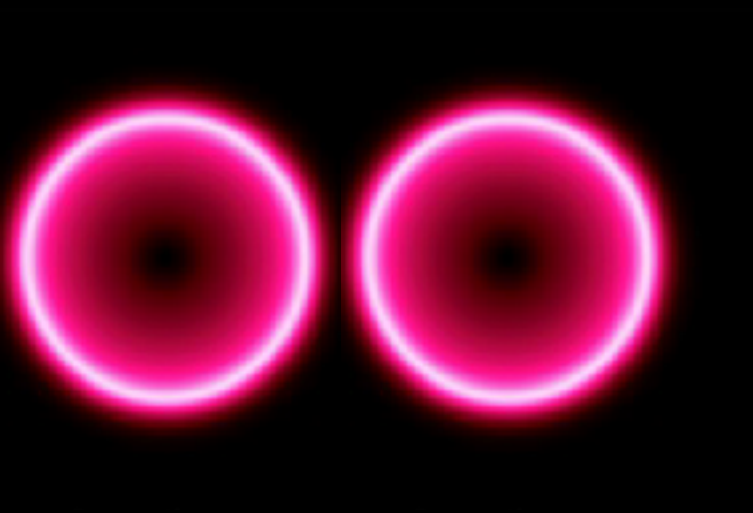


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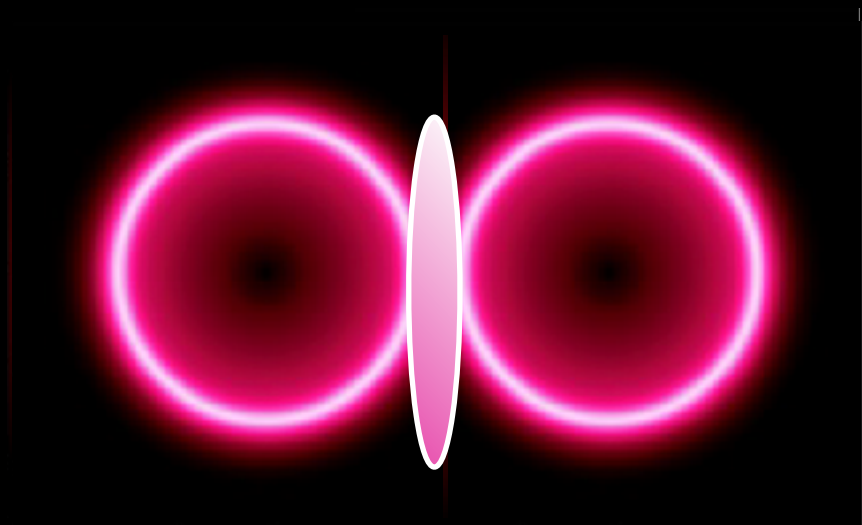


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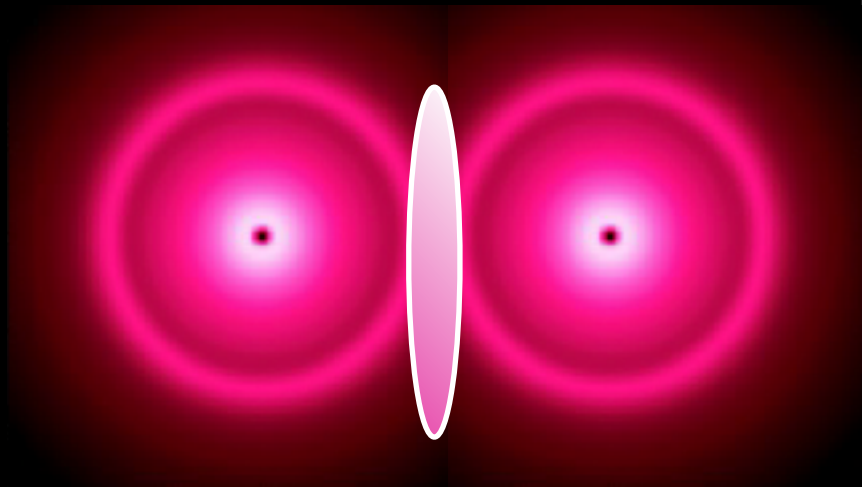


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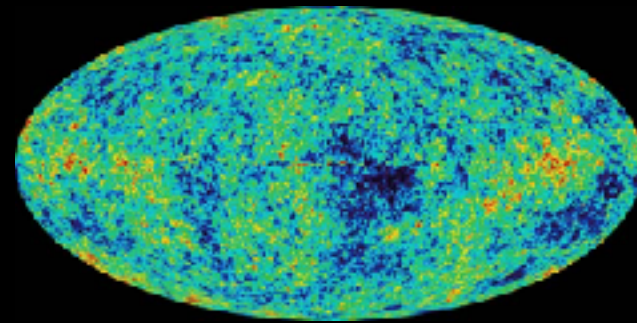
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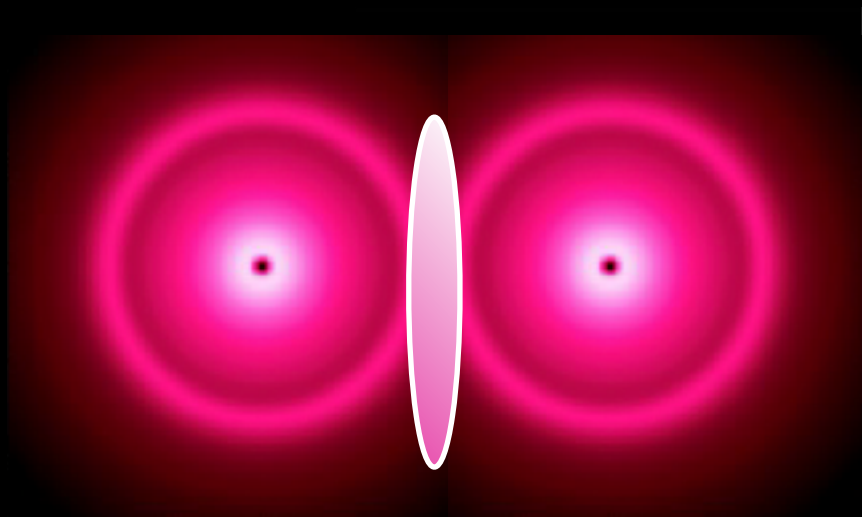


Universe
becomes
transparent

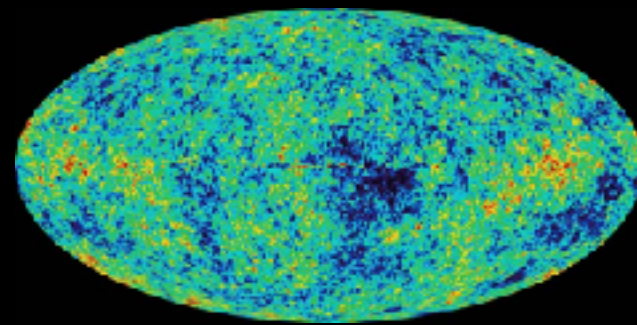
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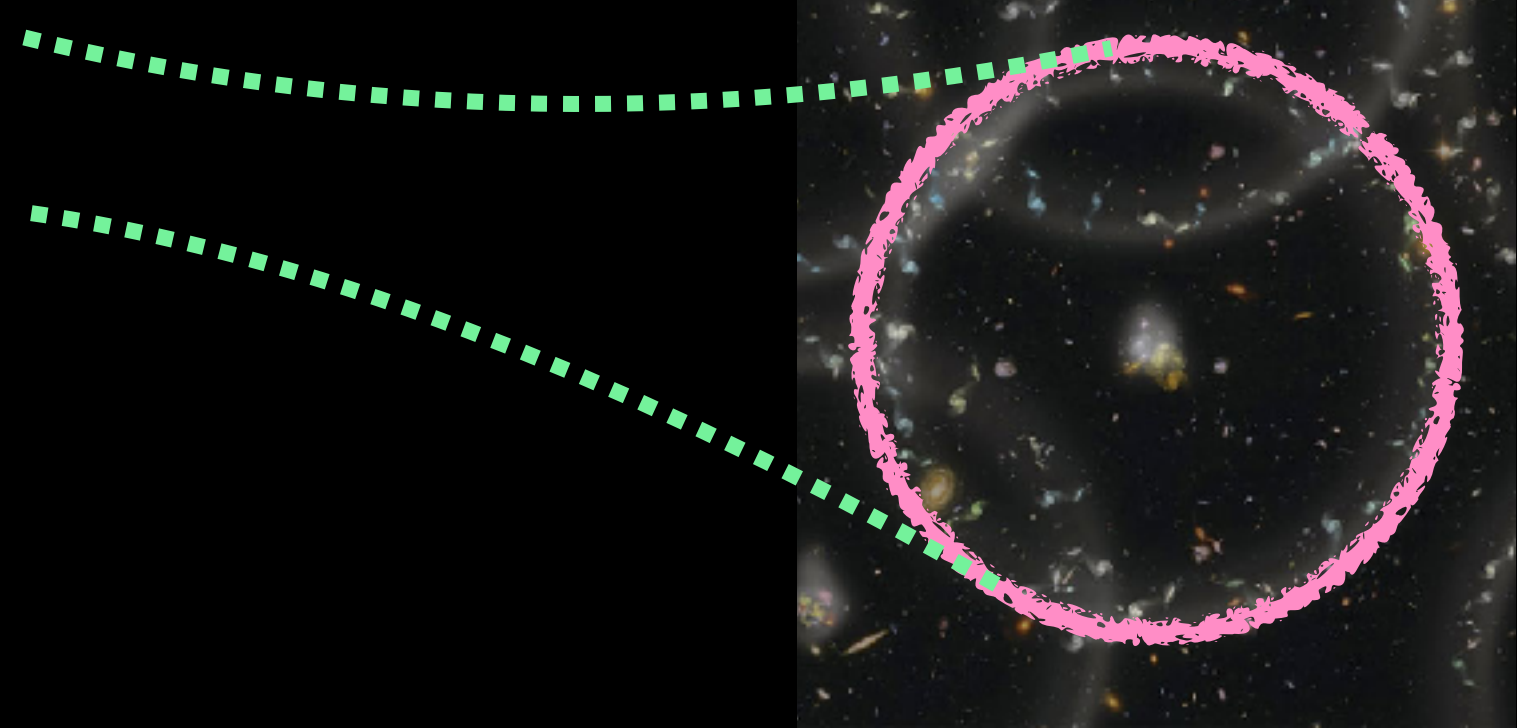
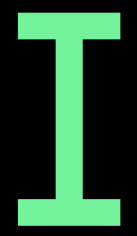
Can be seen in
the clustering of
galaxies (**BAO**)

Comparing r_s at recombination (drag epoch) to r_s at different redshifts tells us about the expansion of the Universe.



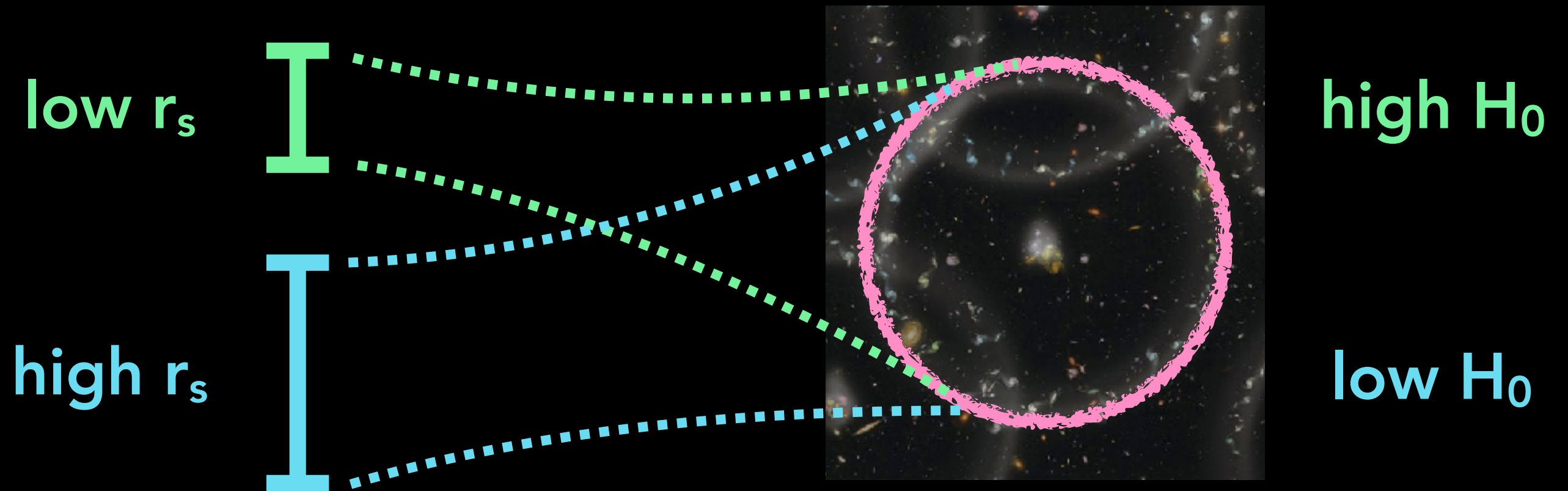
Comparing r_s at recombination (drag epoch) to r_s at different redshifts tells us about the expansion of the Universe.

low r_s



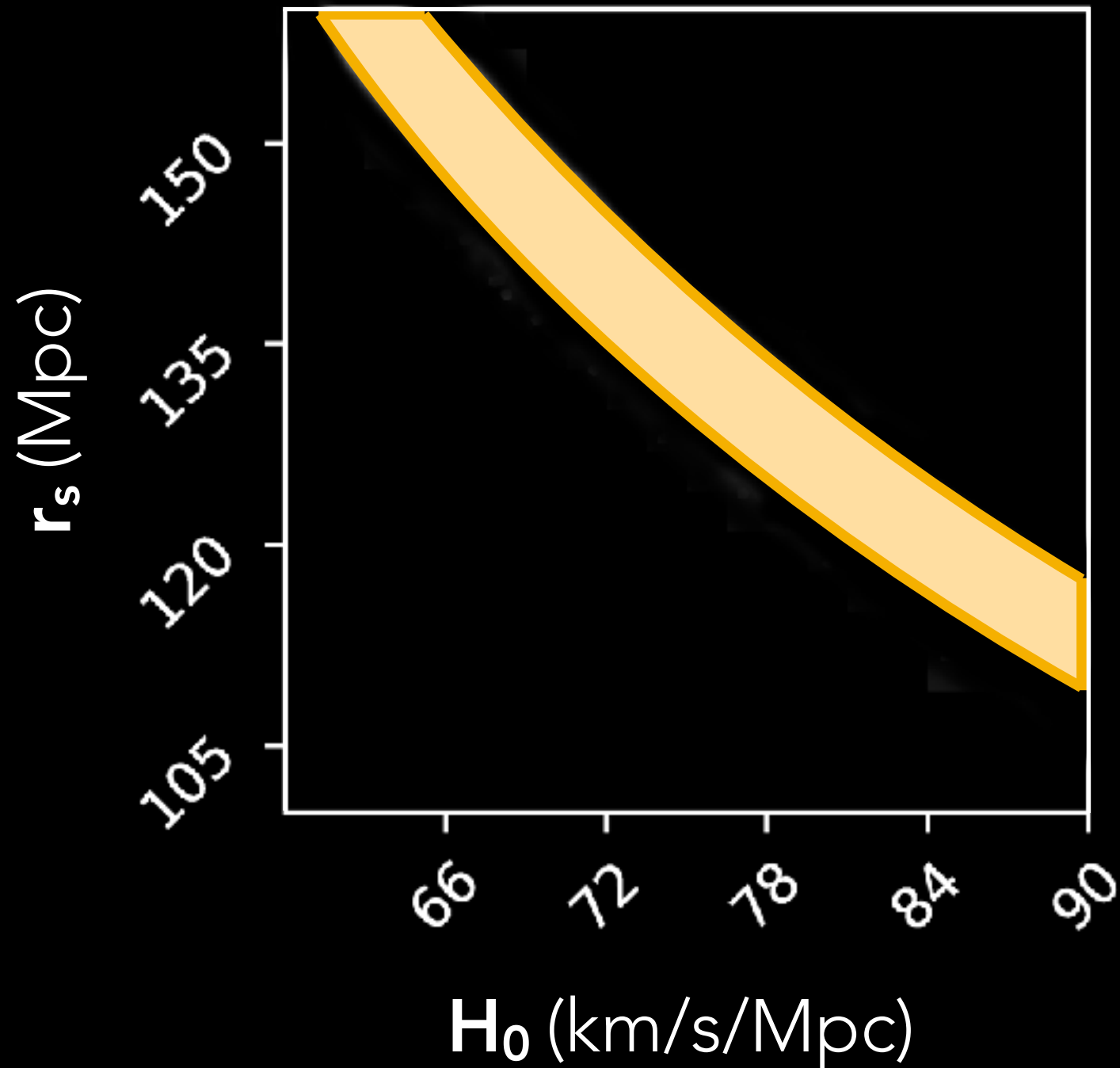
high H_0

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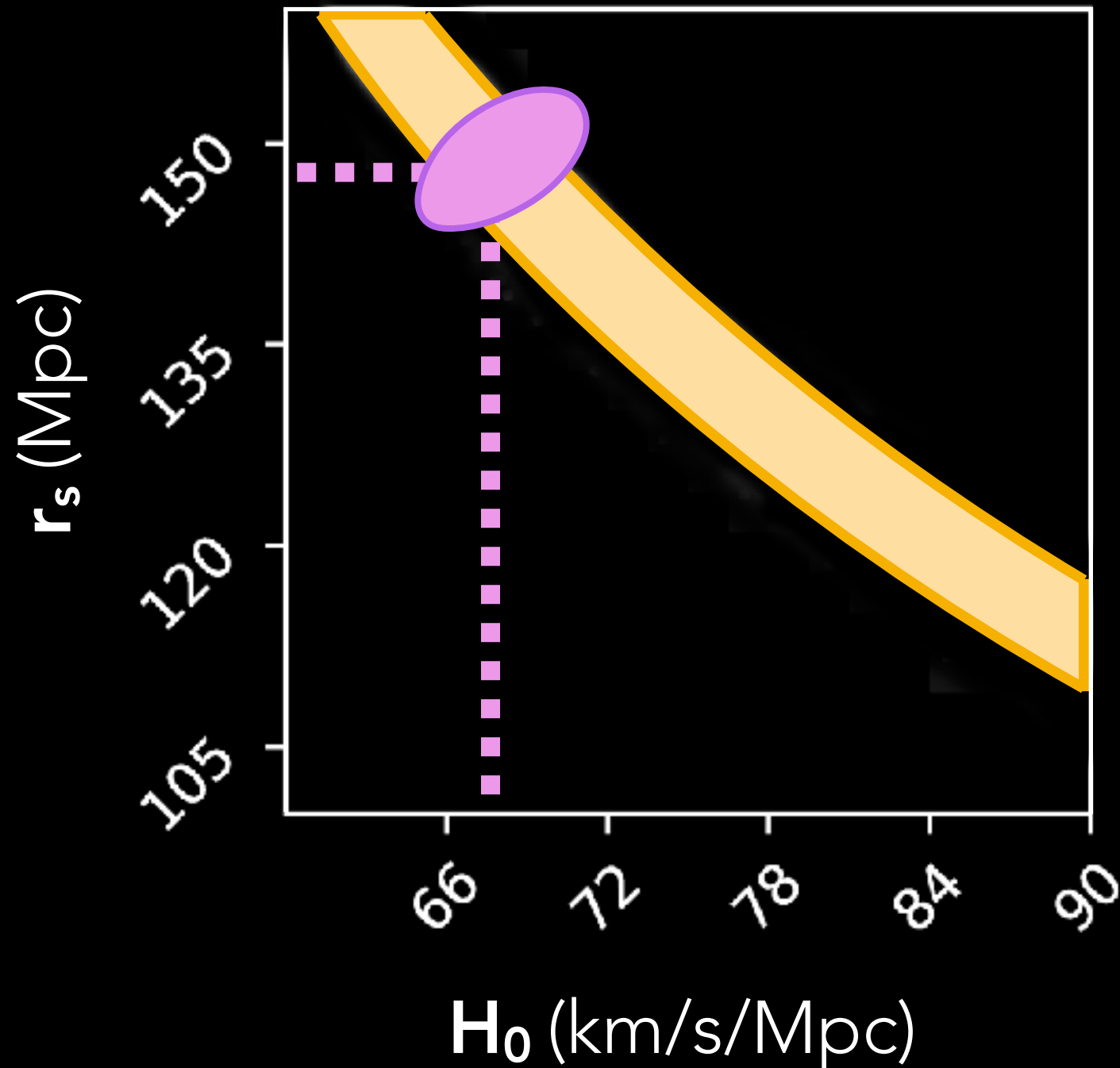


In this way, BAO constrains the product $H_0 r_s$.

The tension is a problem in both parameters,
they are degenerate with each other.

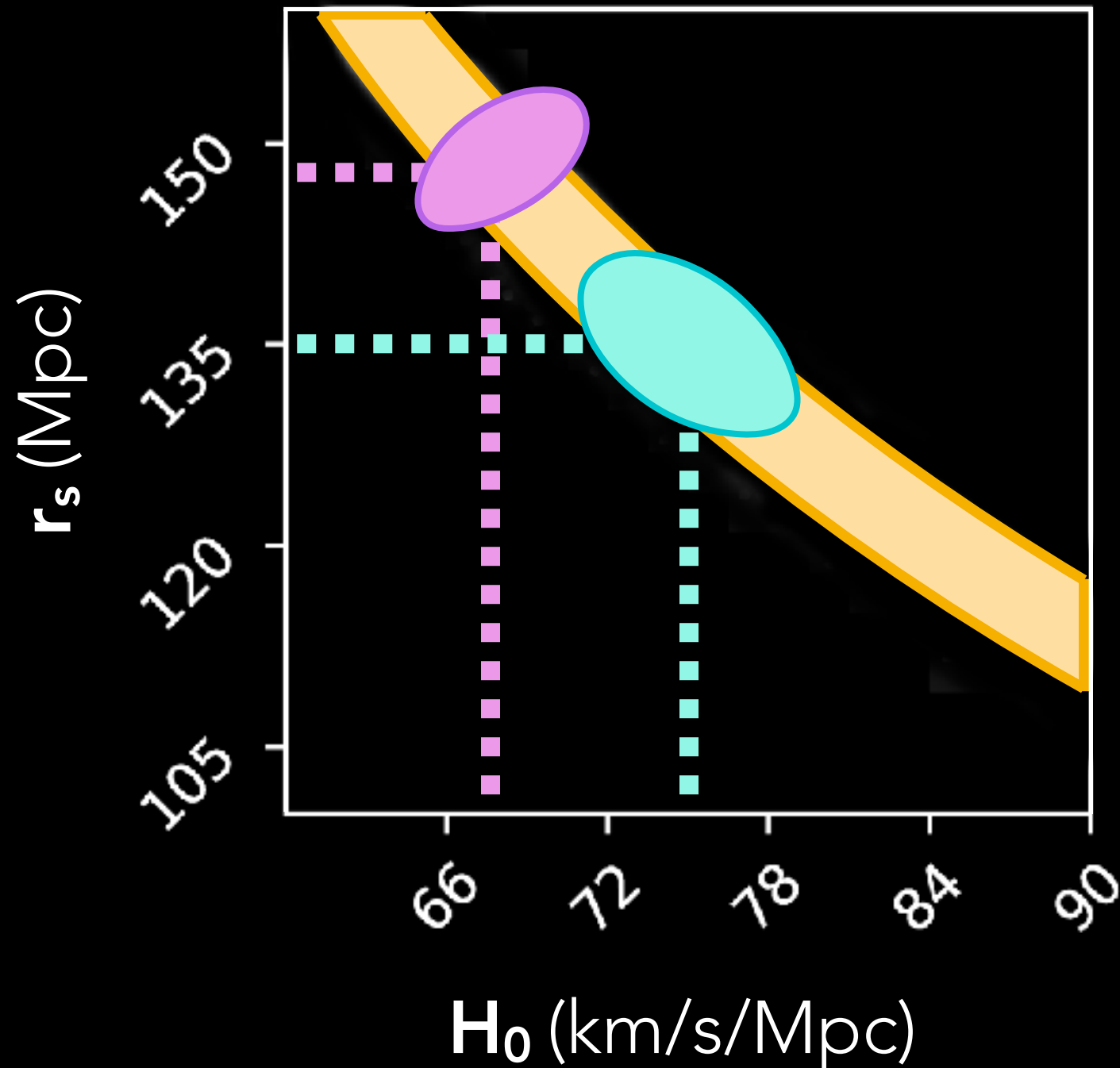


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Planck measures a **high** r_s and
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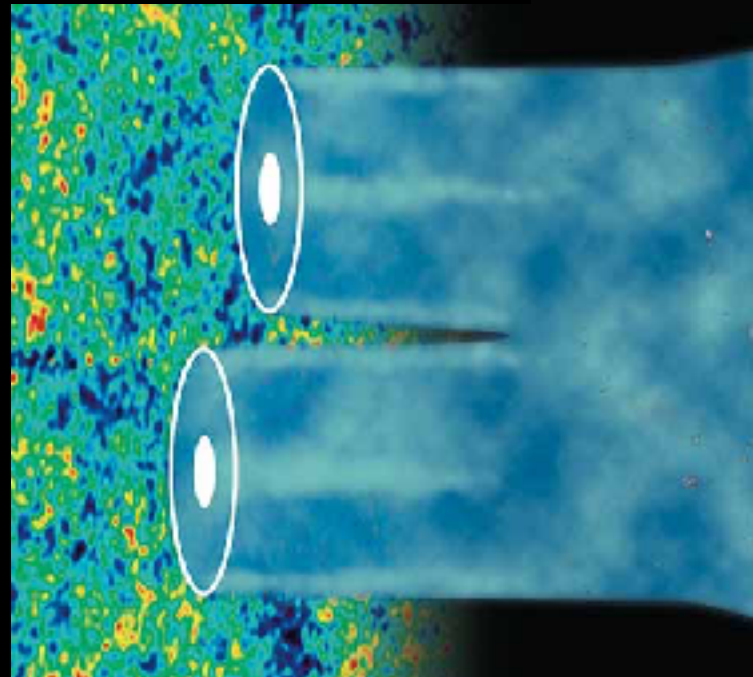


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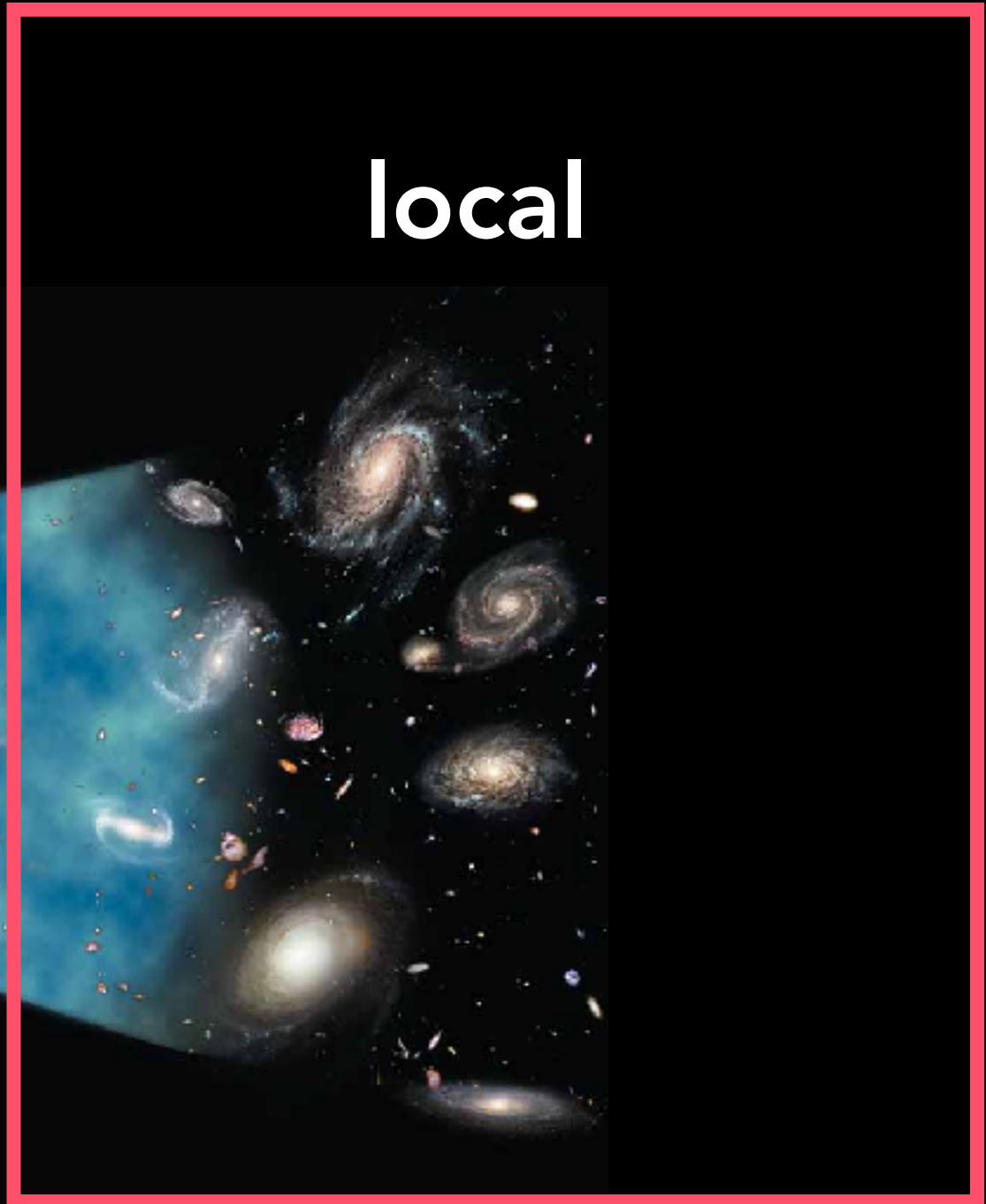
While **local measurements**
get a **low** r_s and a **high** H_0 .

LOCAL DETERMINATION OF H_0 AND R_S

CMB

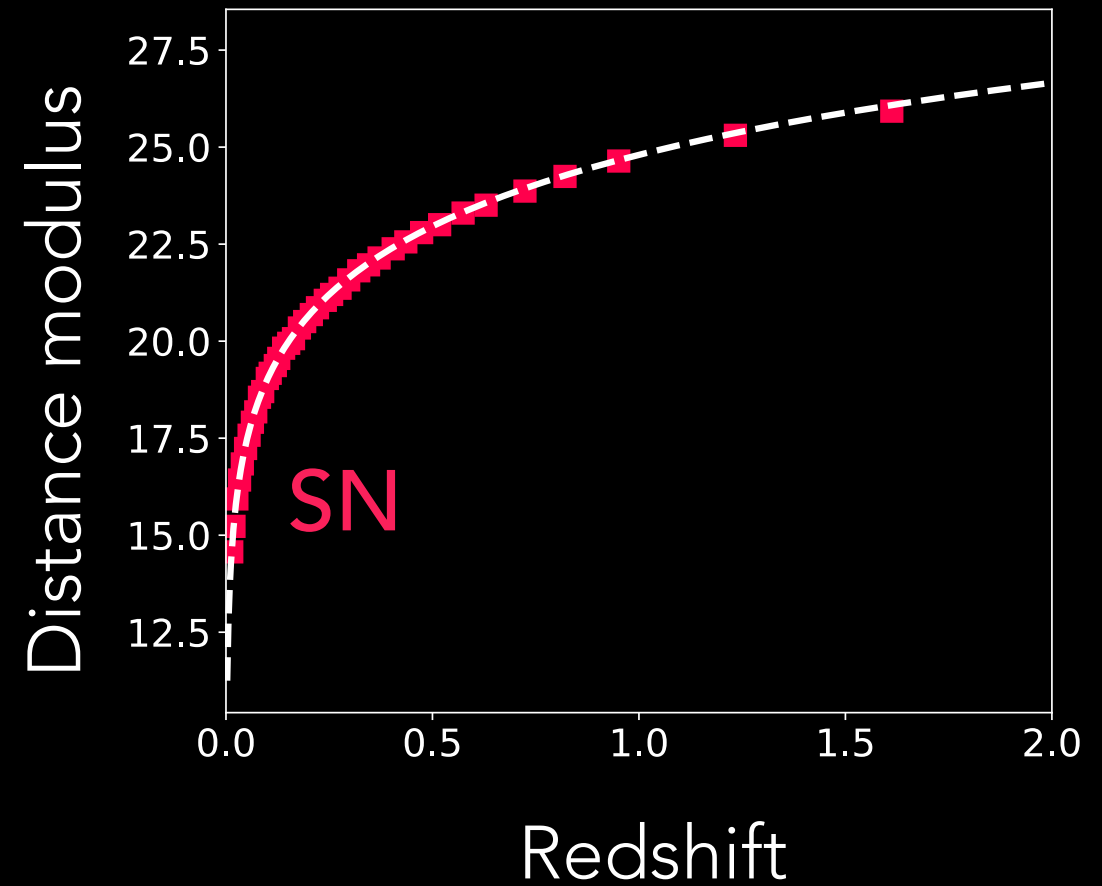
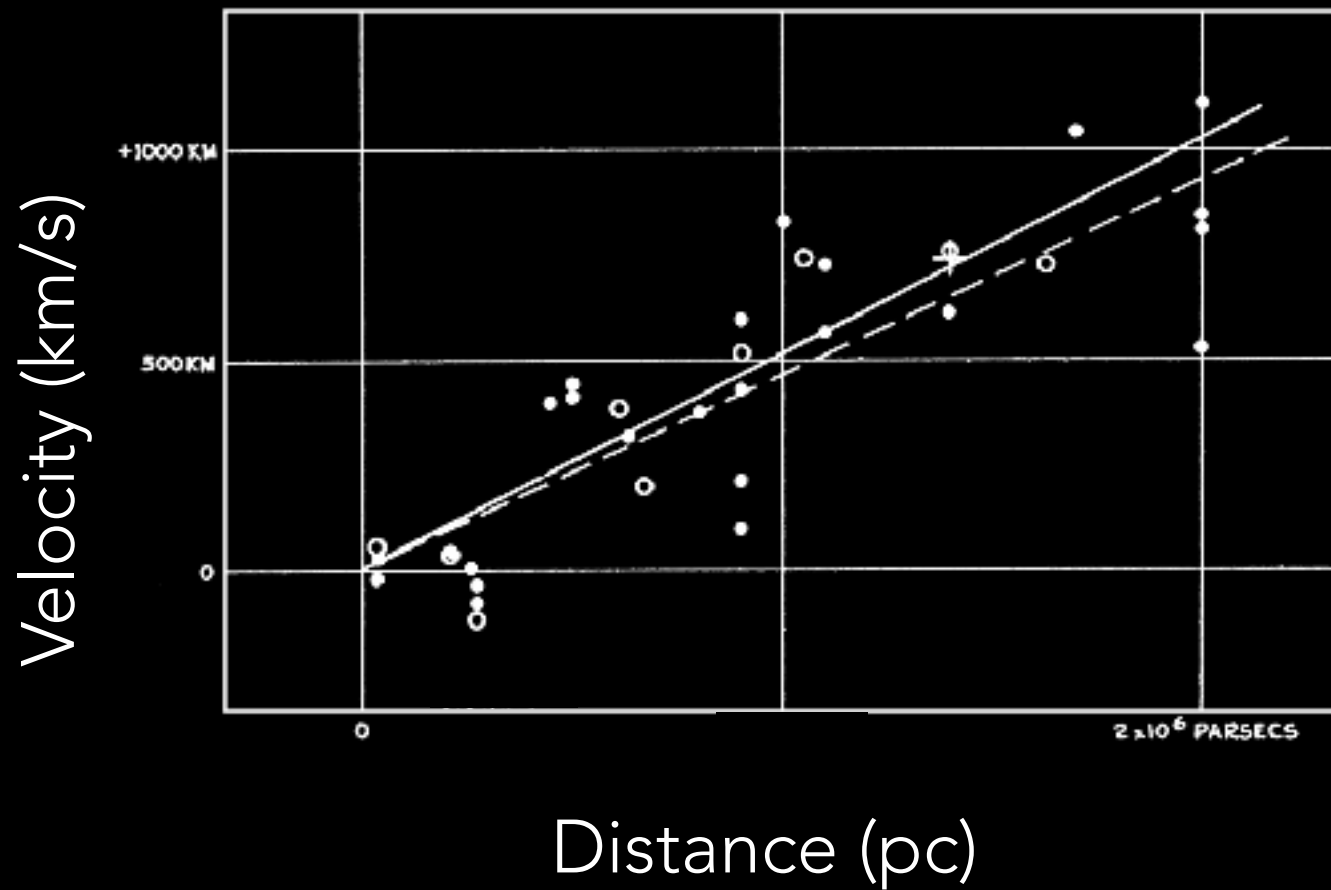


local



LOCAL DETERMINATION OF H_0 AND R_s

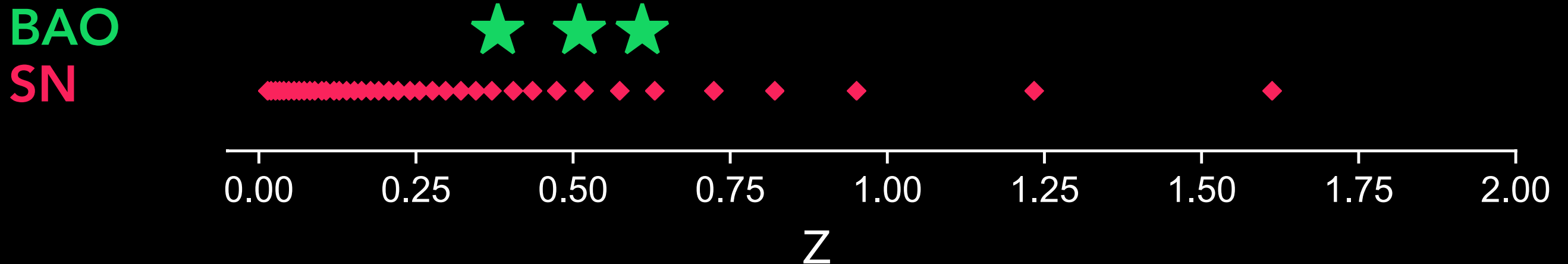
Modern & improved Hubble diagram



RELATIVE DISTANCE INDICATORS

Supernovae type Ia (Pantheon, Scolnic et al 2018) and baryon acoustic oscillations (**BAO**) (BOSS, Alam et al 2017) yield relative distance measurements.

BAO introduce the sound horizon as a free parameter.



CALIBRATE DATA

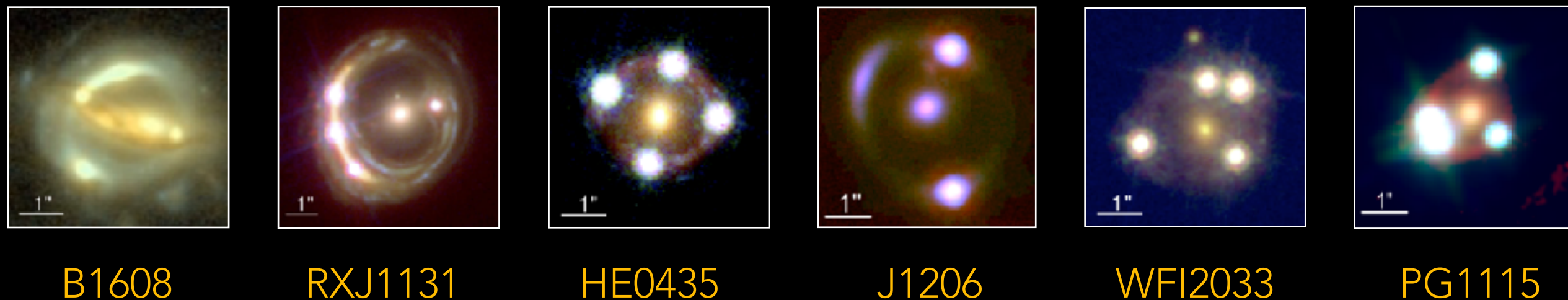
This is often done with **Cepheids**, stars at the Tip of the Red Giant Branch (**TRGB**) or Surface Brightness Fluctuations (**SBF**).

But those can only be measured at low redshifts.



H0LiCOW

H0 Lenses in COSMOSGRAB Wellspring



Lenses

SBF



Cepheids



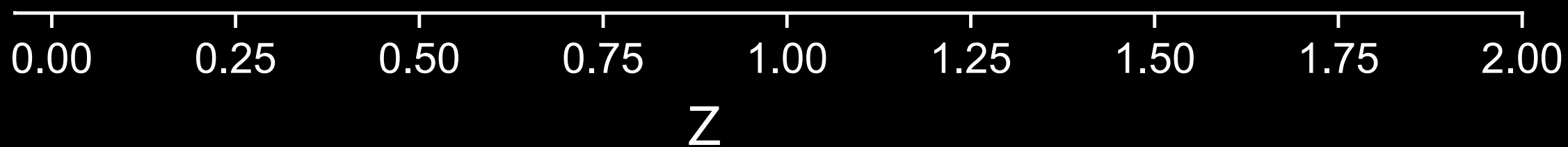
TRGB



BAO



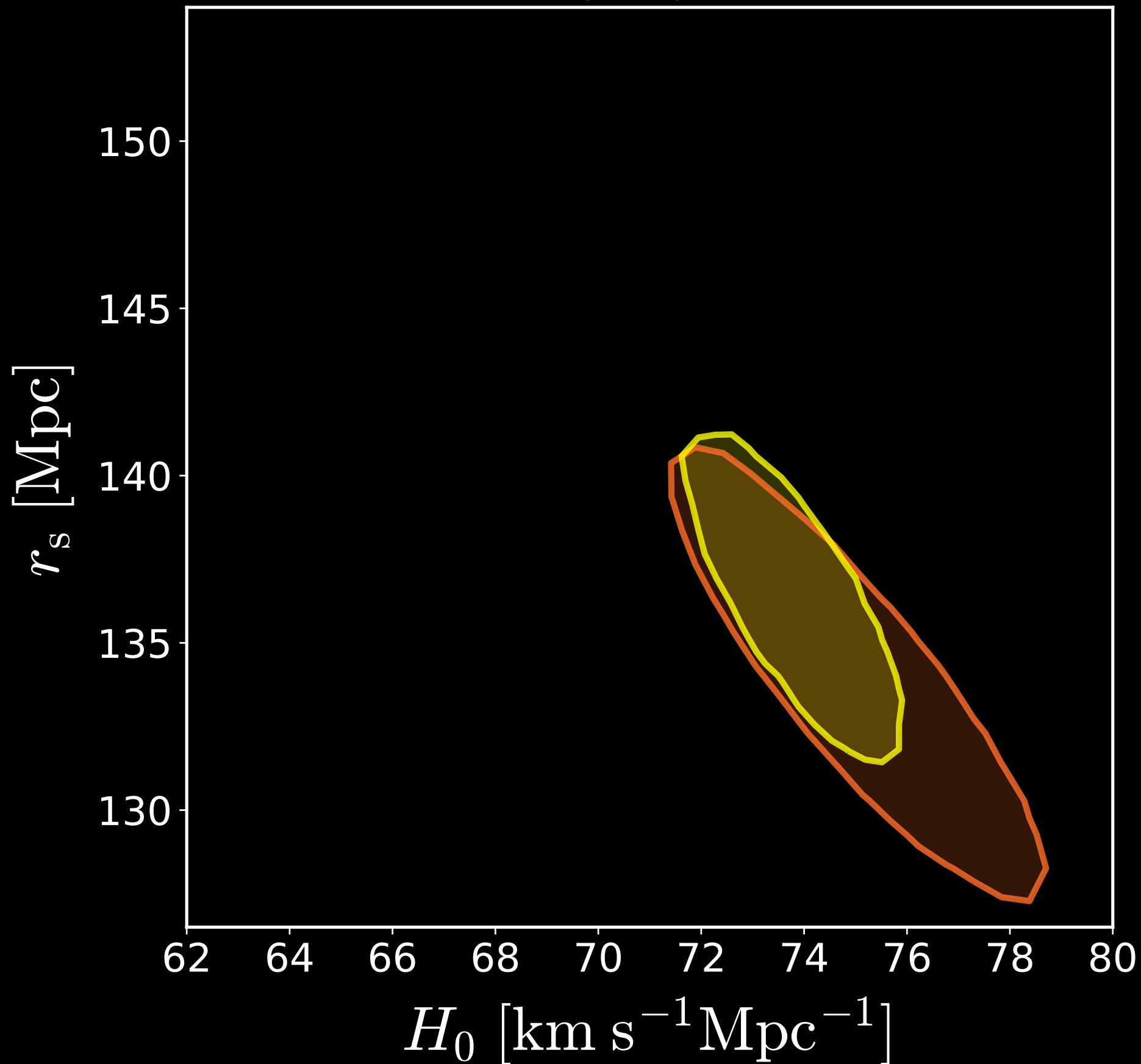
SN



RESULTS



Arendse, Wojtak, Agnello et al. (2020)



SNe Ia

Scolnic et al. (2018)

BAO

Alam et al. (2017)

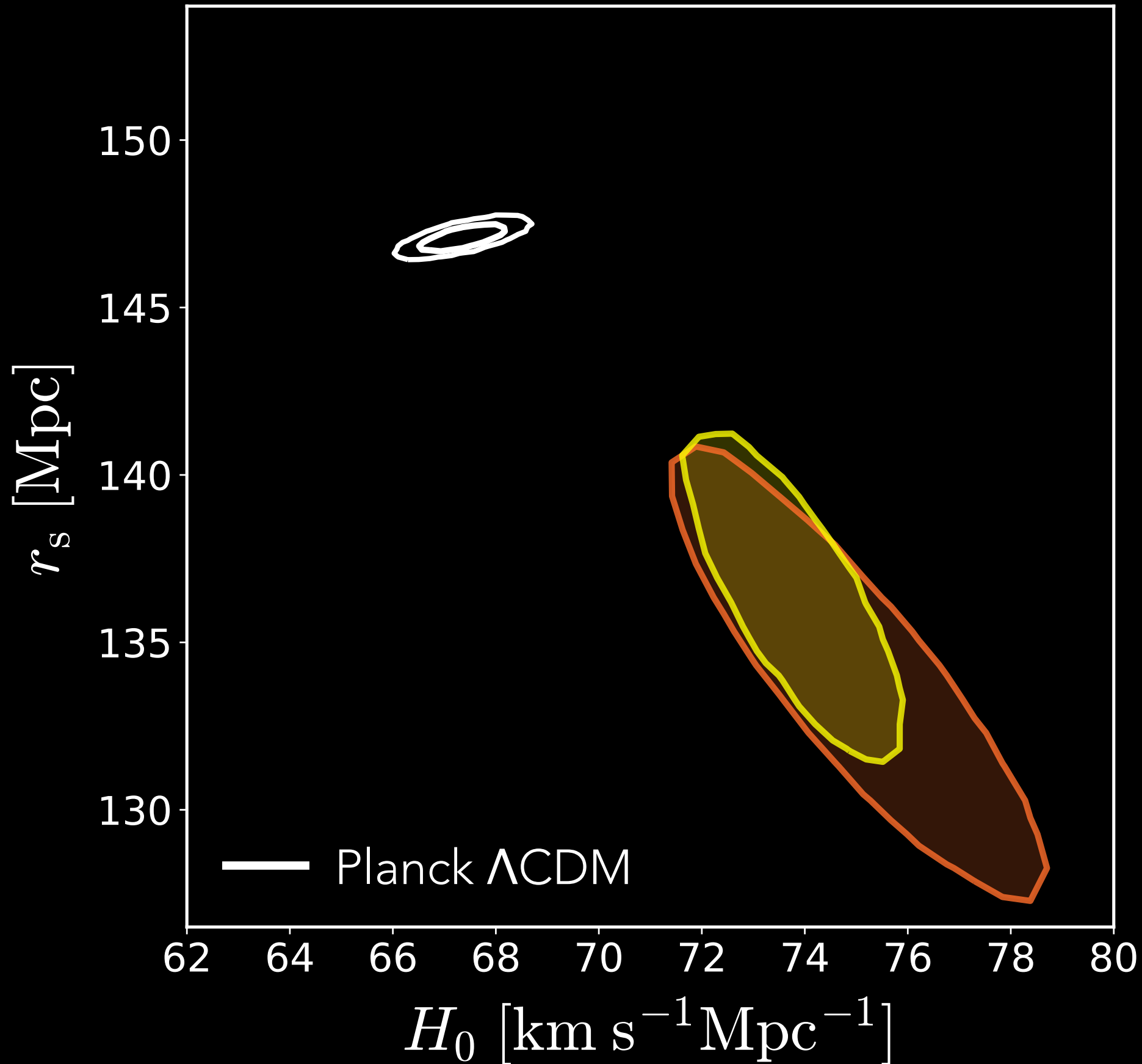
Cepheids (SH0ES)

Riess et al. (2020)

Lenses (H0LiCOW)

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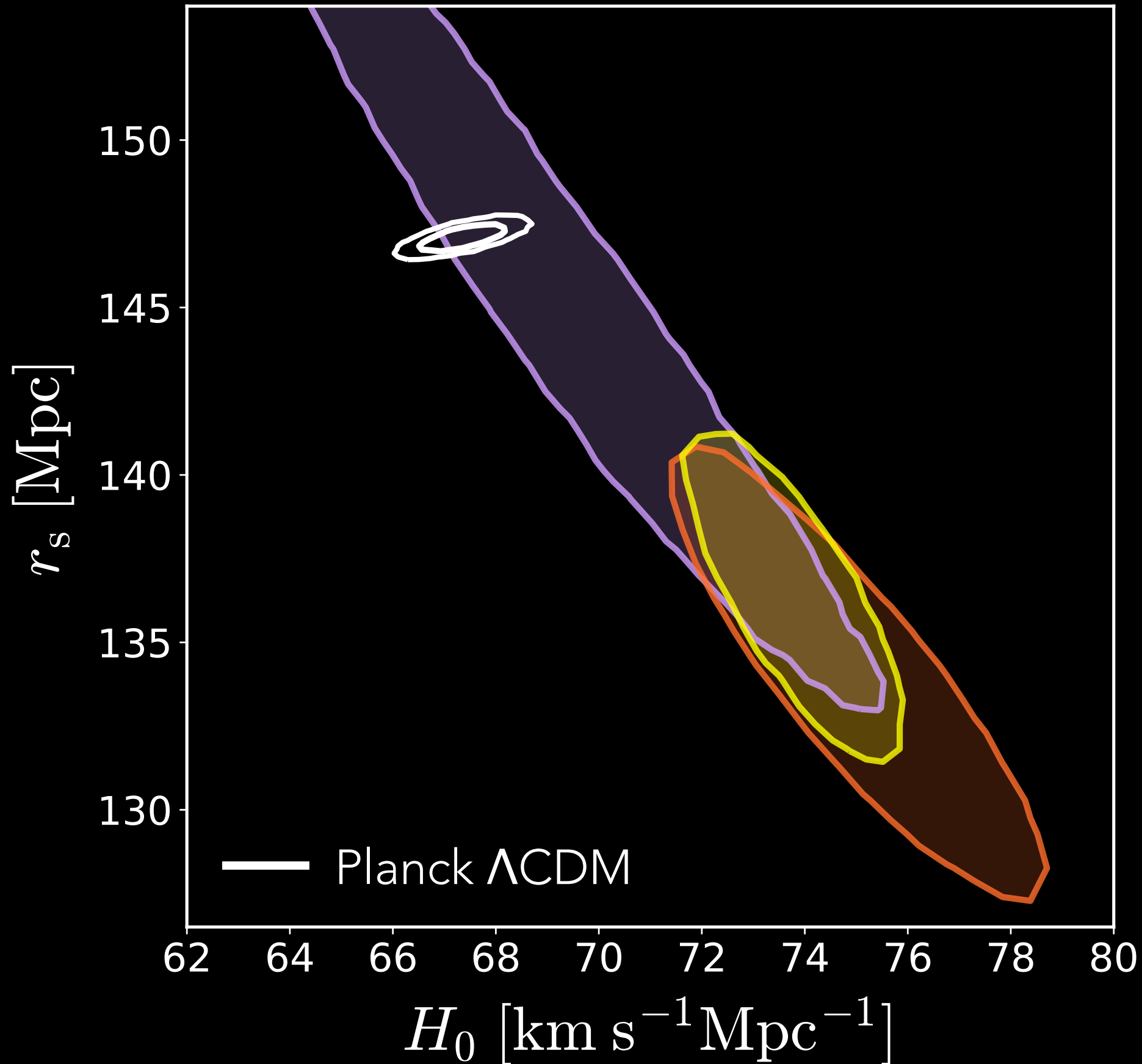
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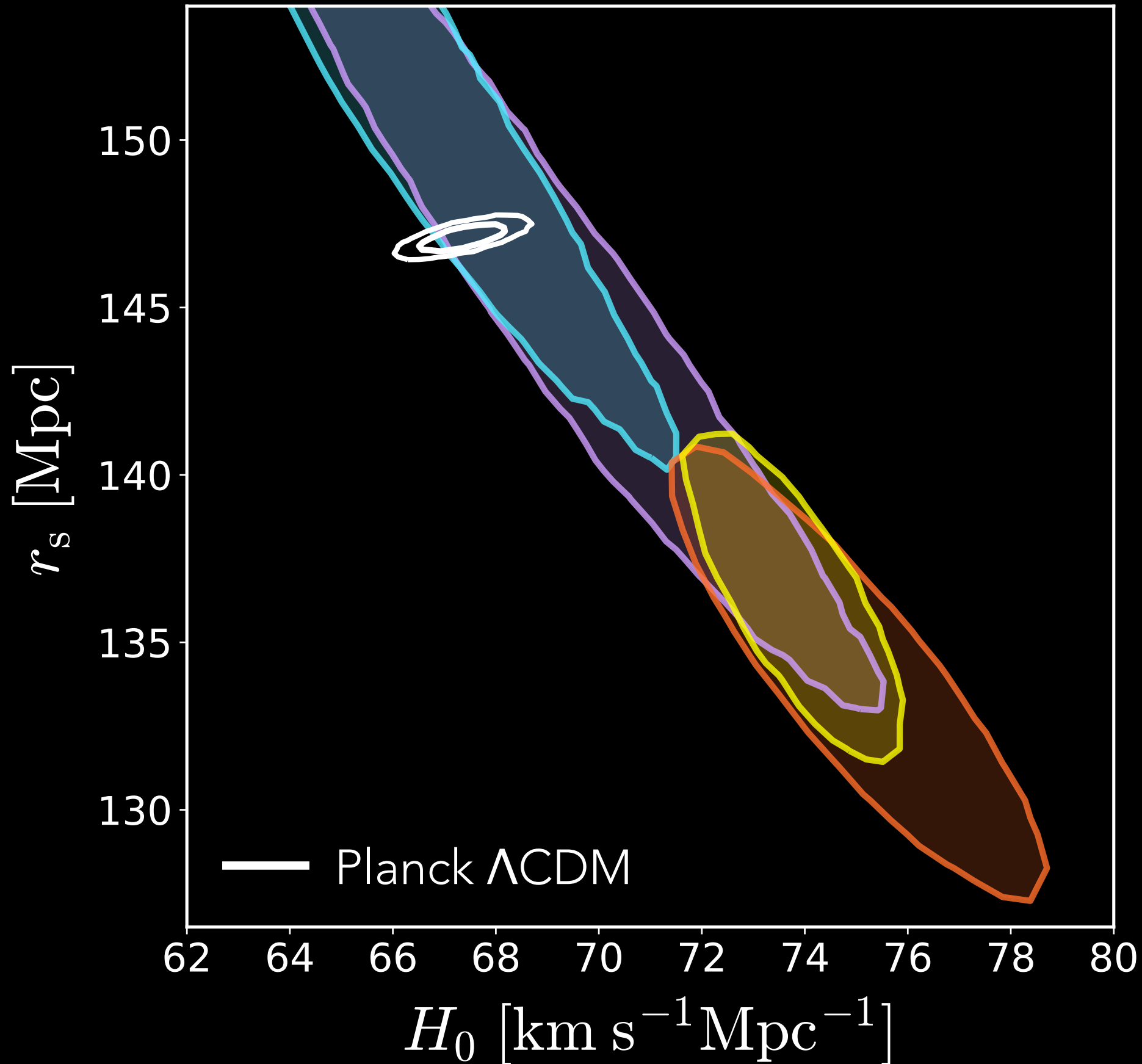
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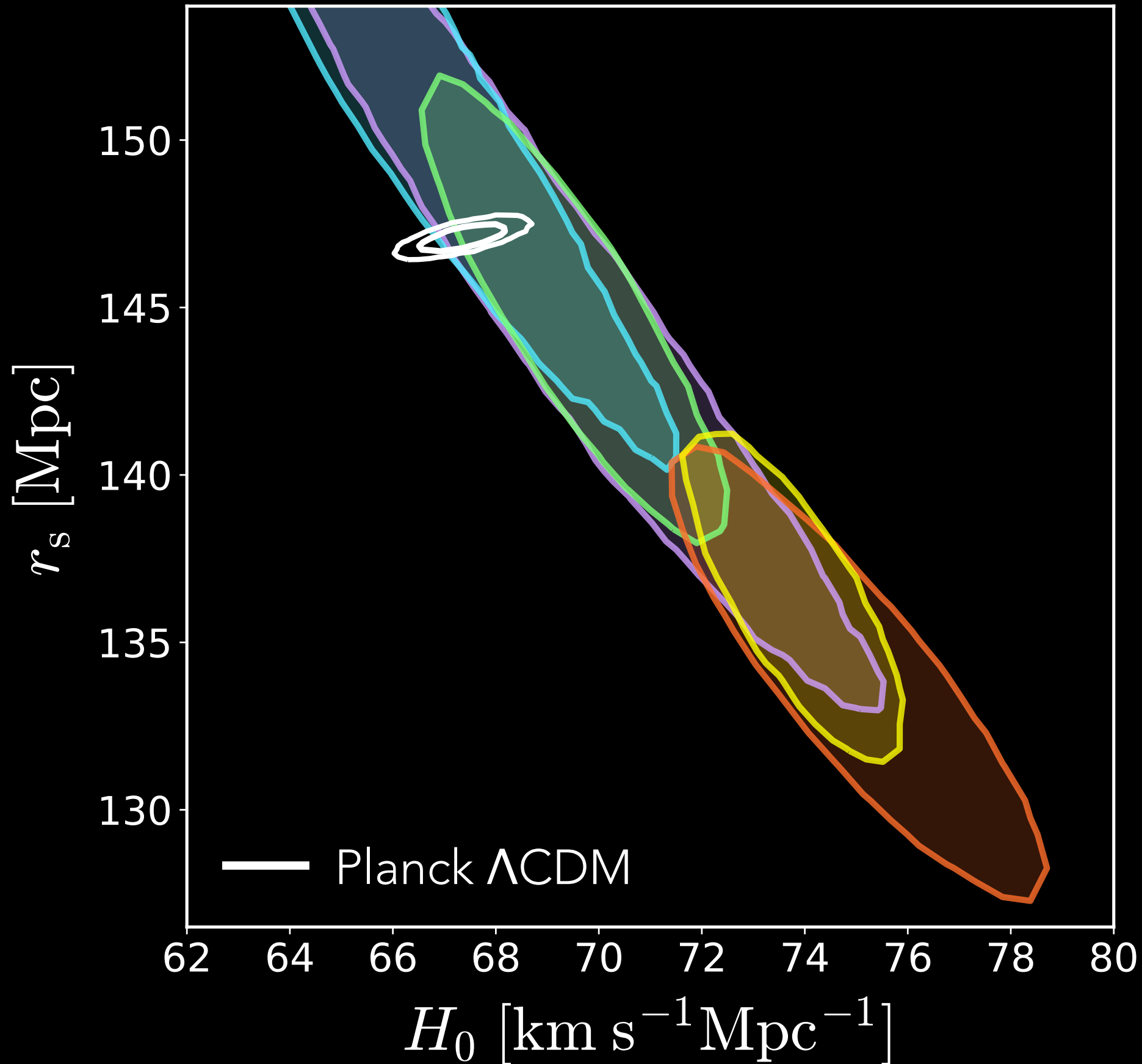
SBF

Khetan et al. (2020)

Lenses (TDCOSMO)

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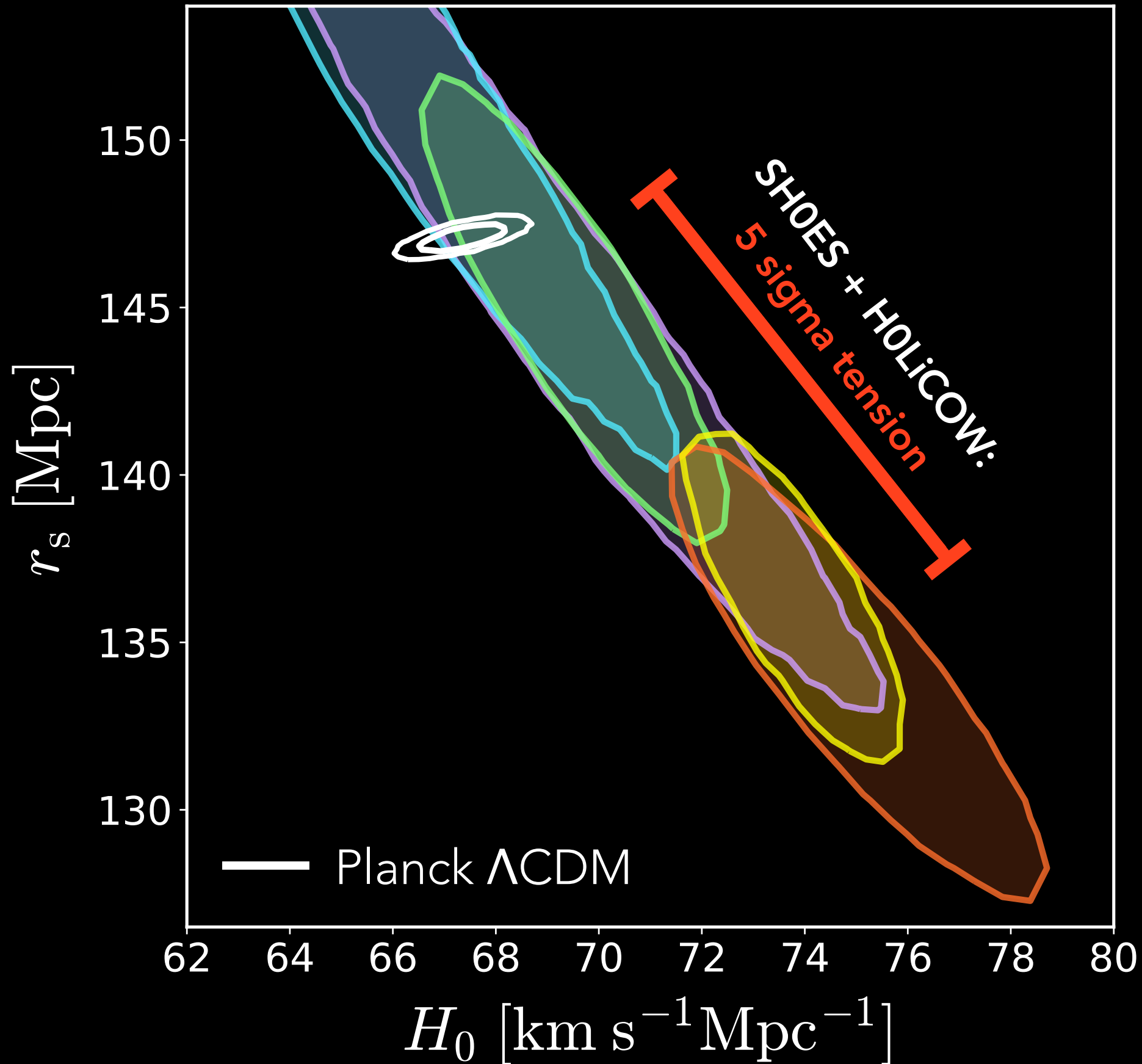
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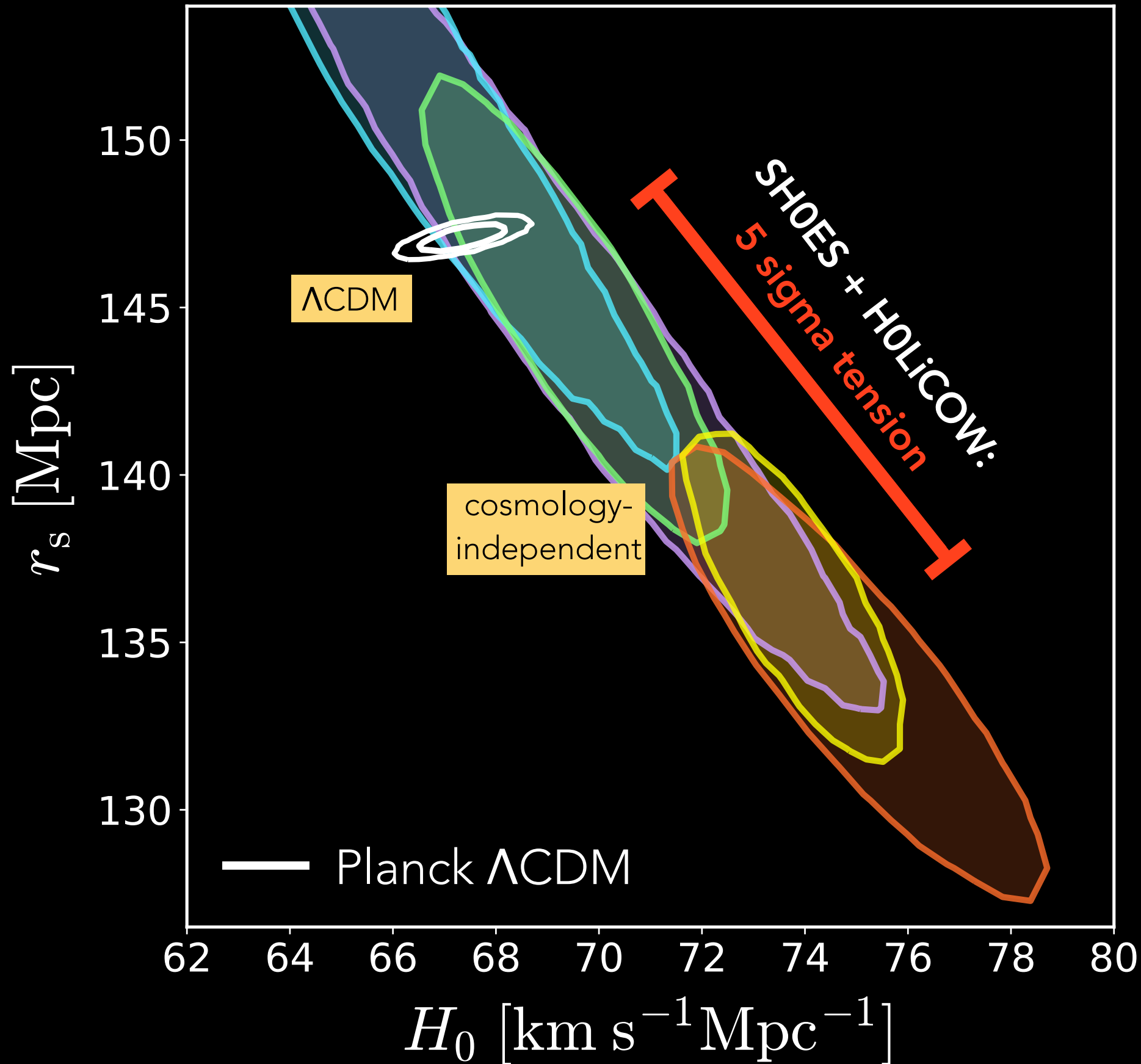
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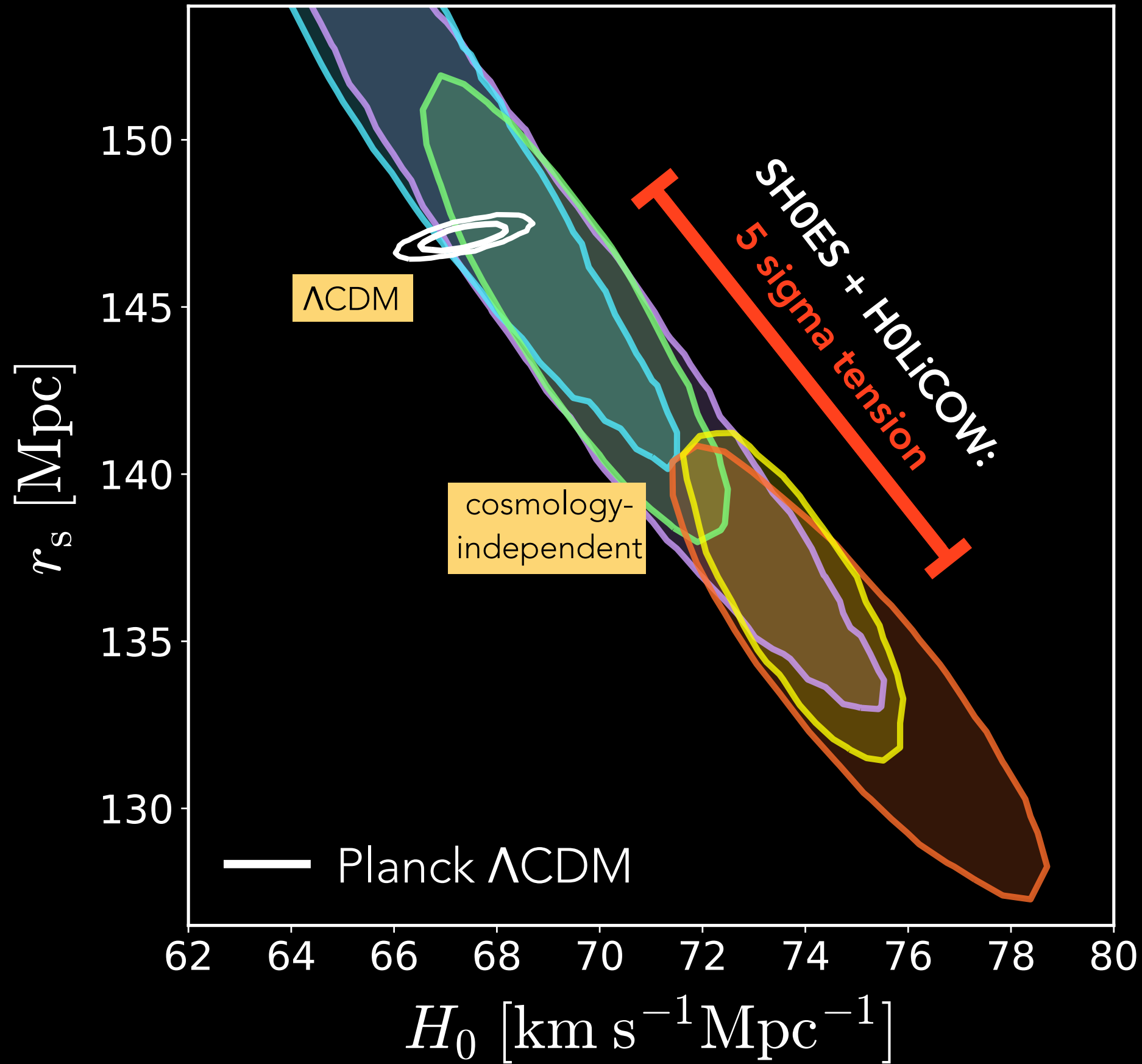
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Can we make changes to Λ CDM that reconcile the tension?

Consider these 4 extensions of Λ CDM:

Extra relativistic particle (N_{eff})	Redshift dependent DE (w CDM)
Early Dark Energy (EDE)	DE model 2 (PEDE)

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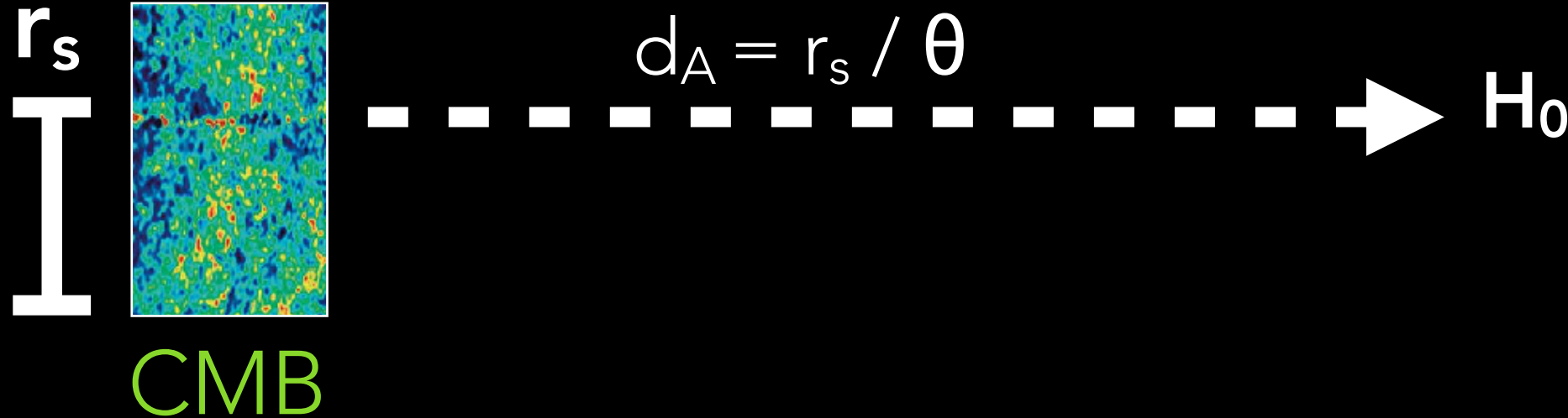
CHANGES TO EARLY PHYSICS

CHANGES TO LATE PHYSICS

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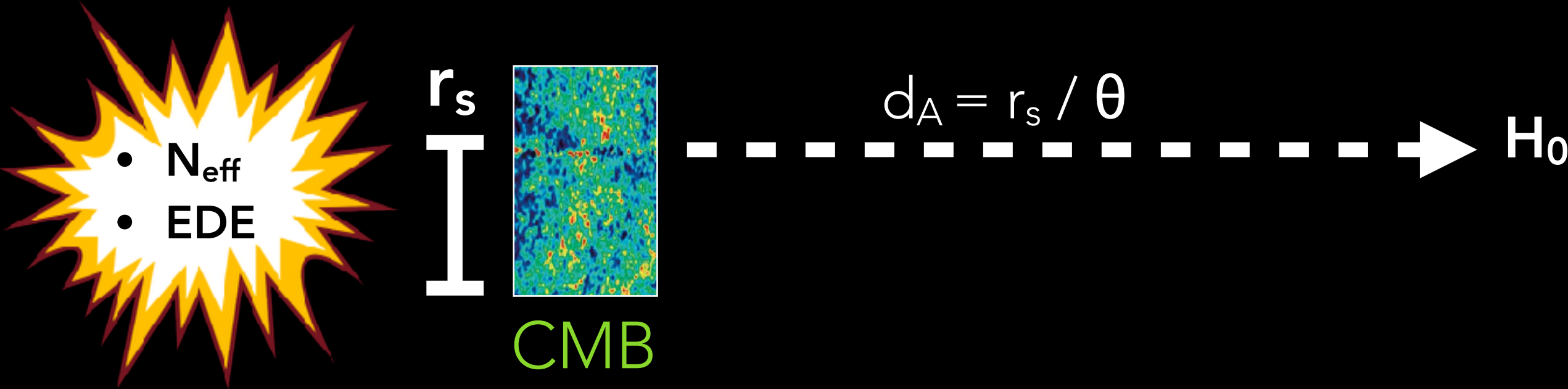


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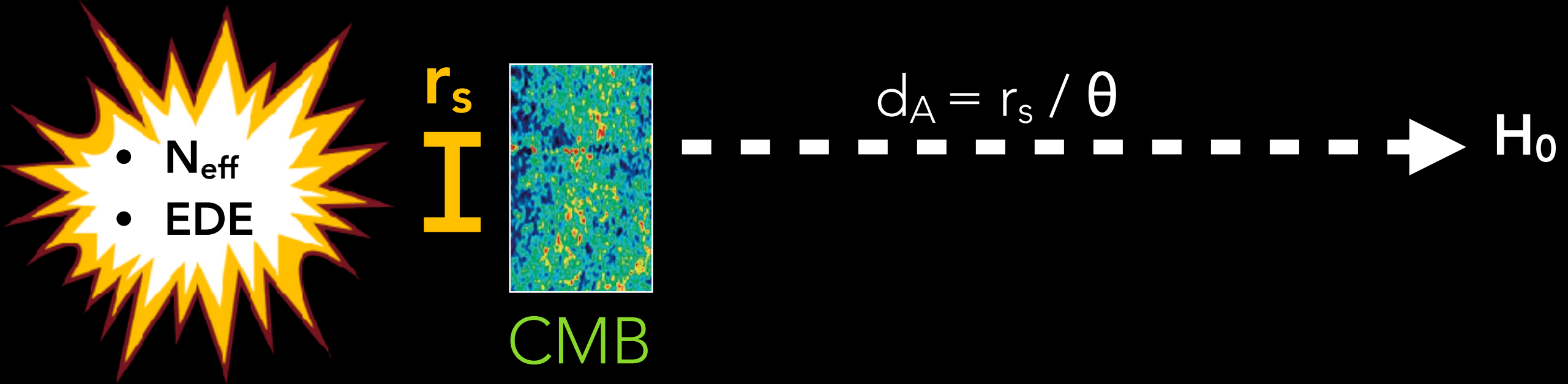


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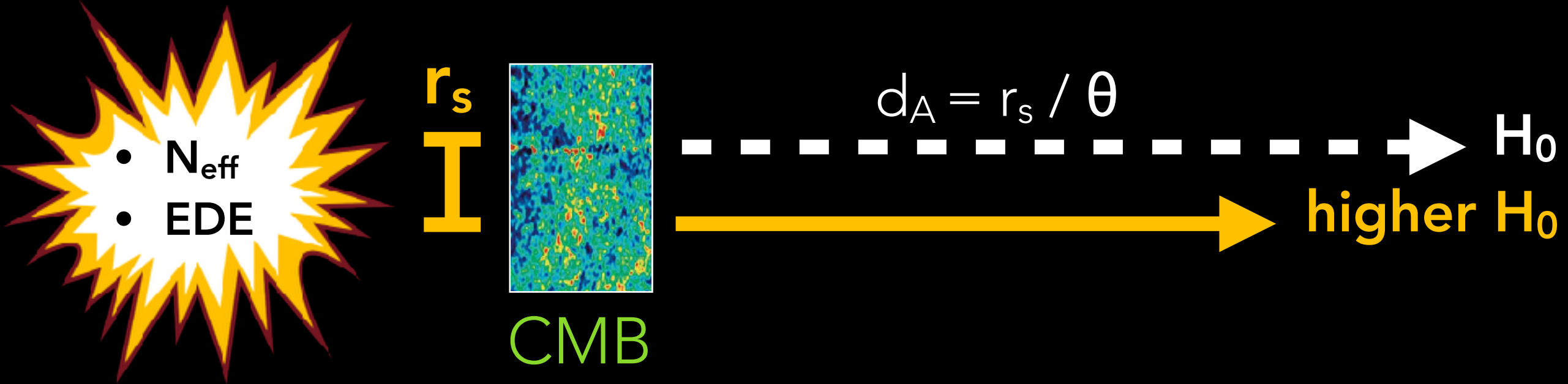


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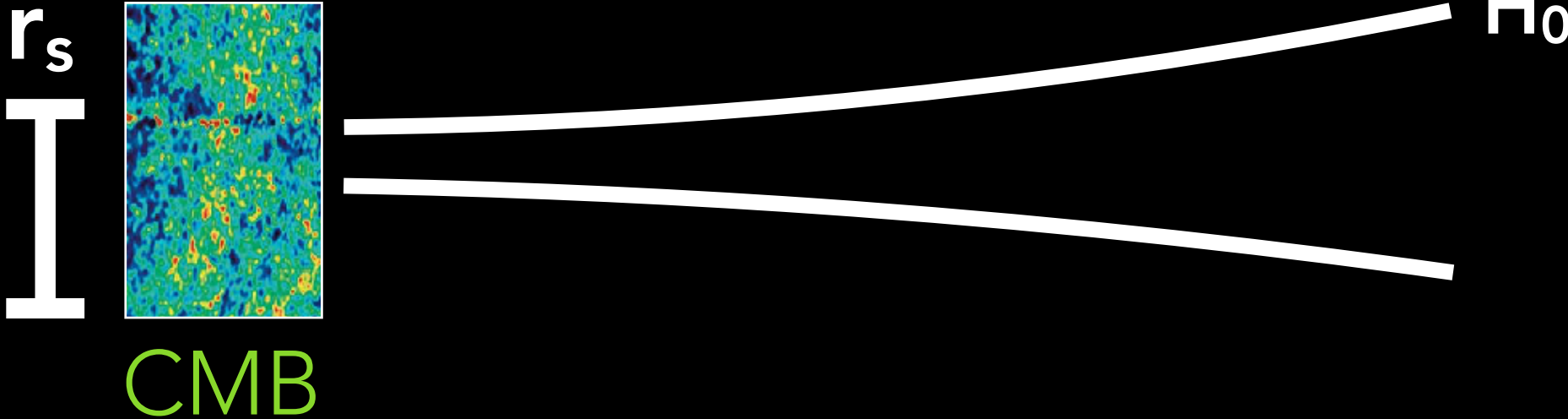
CHANGES TO LATE PHYSICS

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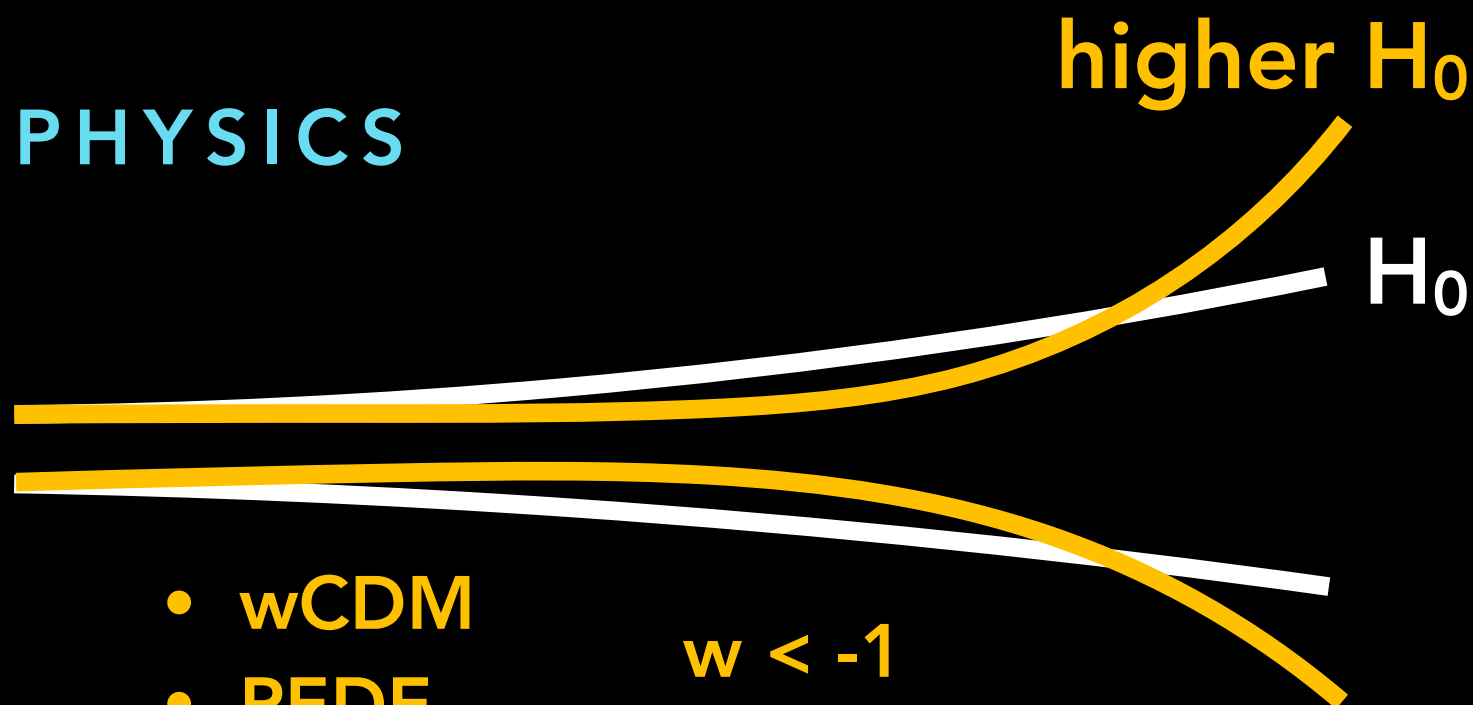
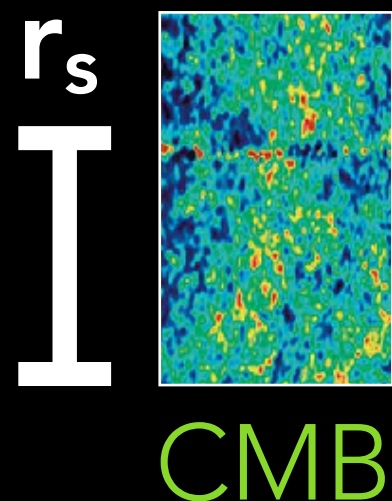


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CHANGES TO EARLY PHYSICS

CHANGES TO LATE PHYSICS



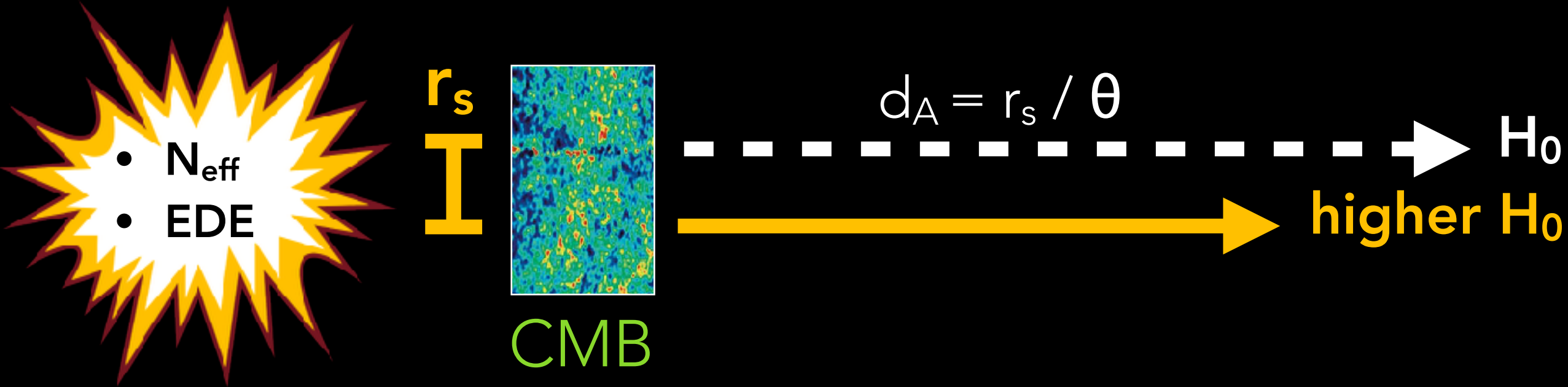
- wCDM
- PEDE

Li & Shafieloo (2019)

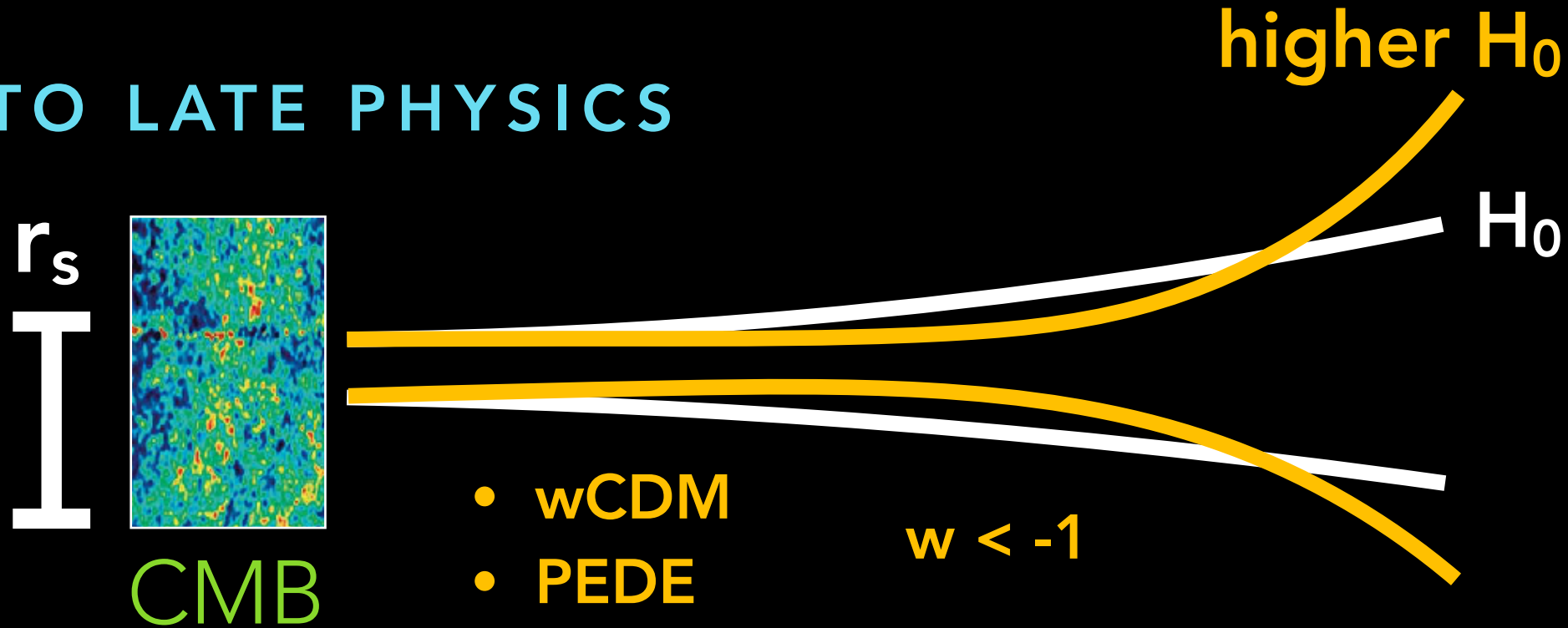
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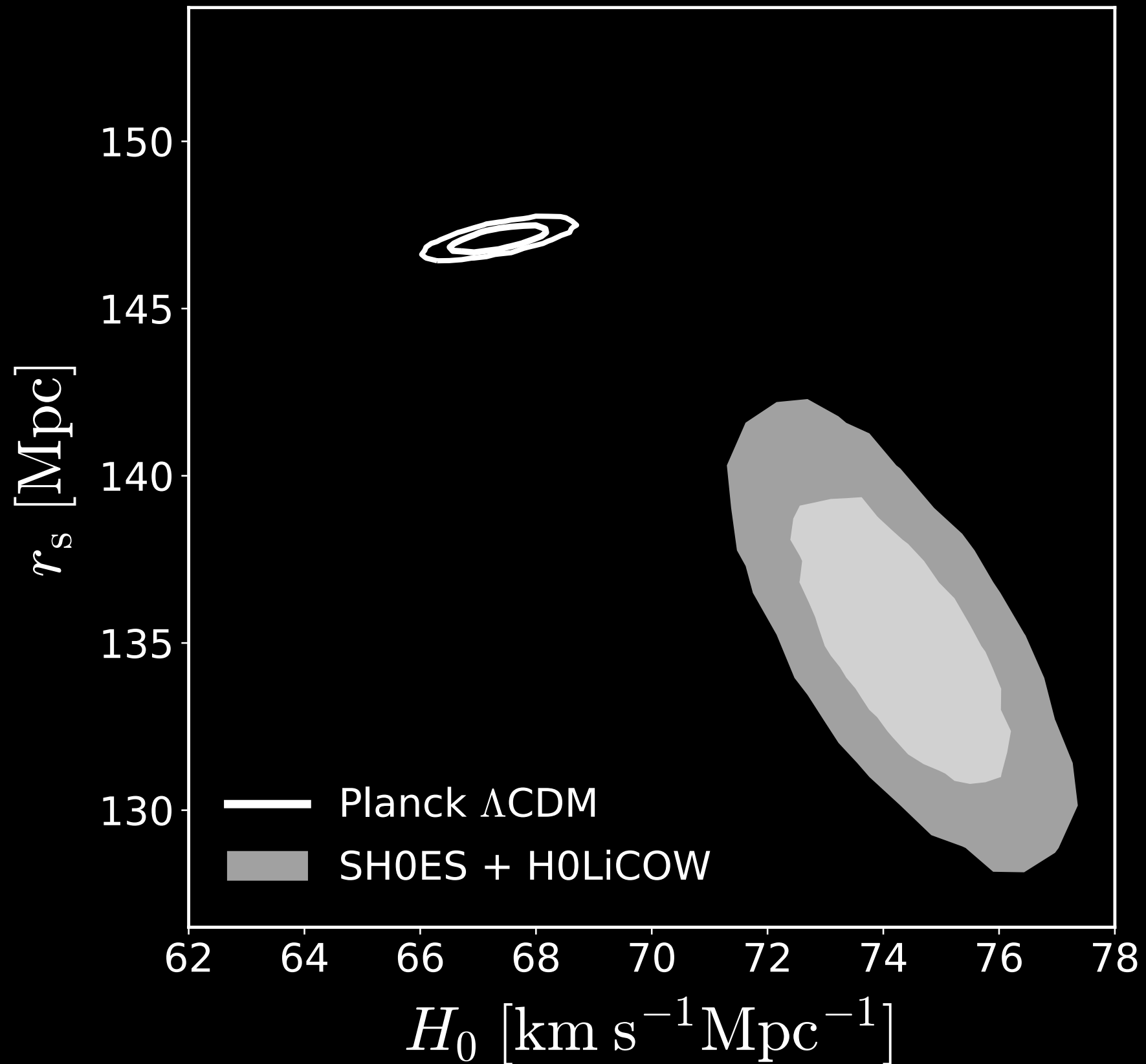


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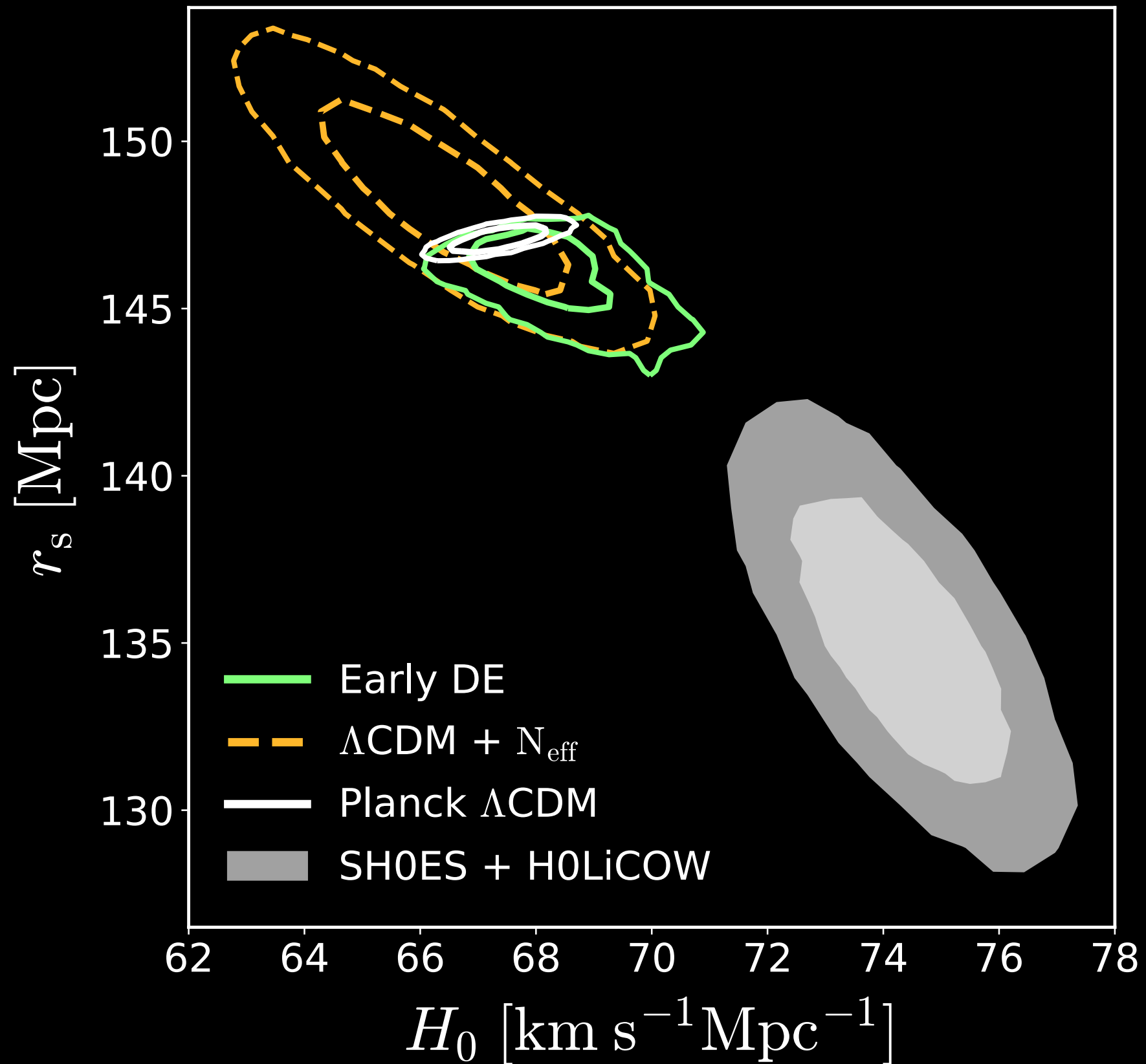


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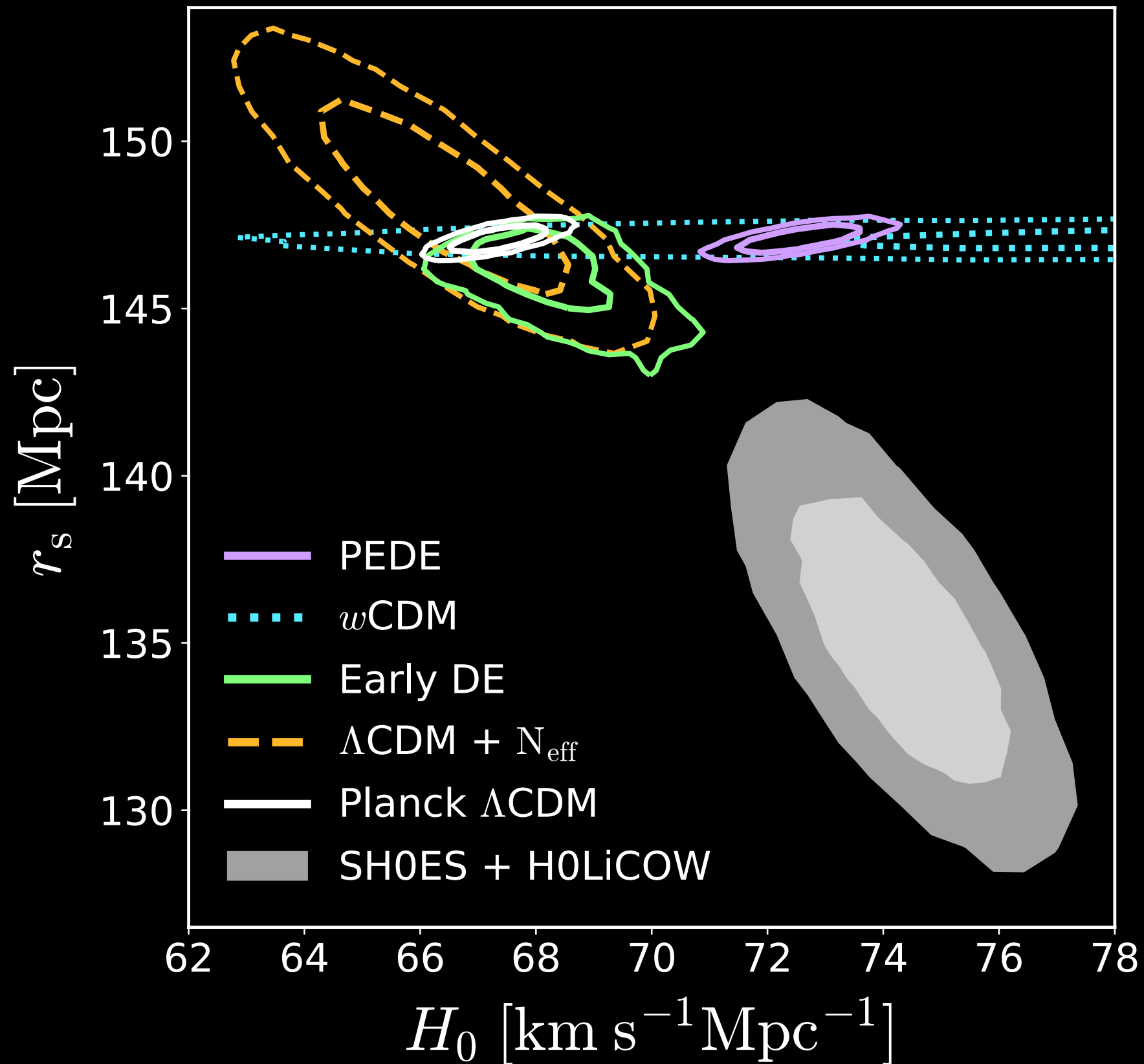
RESULTS



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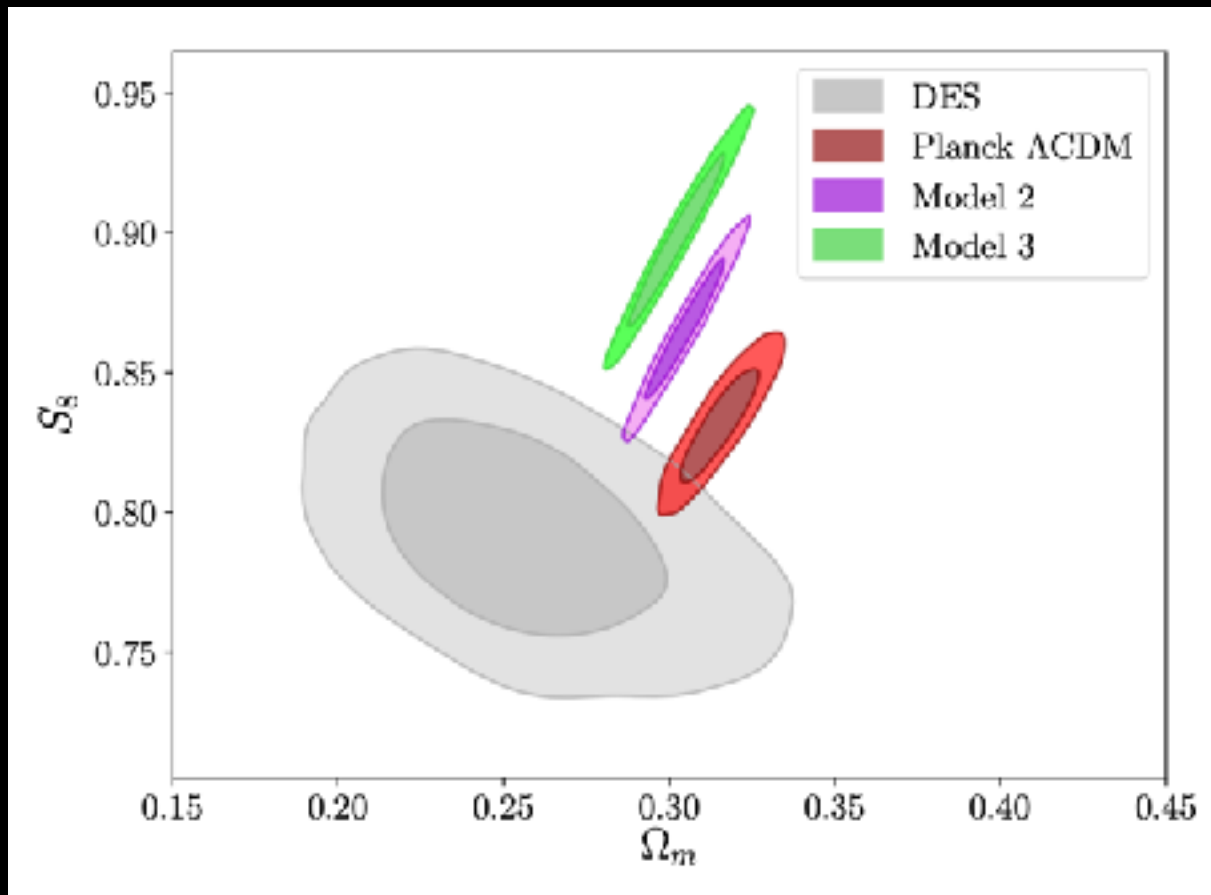
RESULTS



ADDITIONAL CHALLENGES

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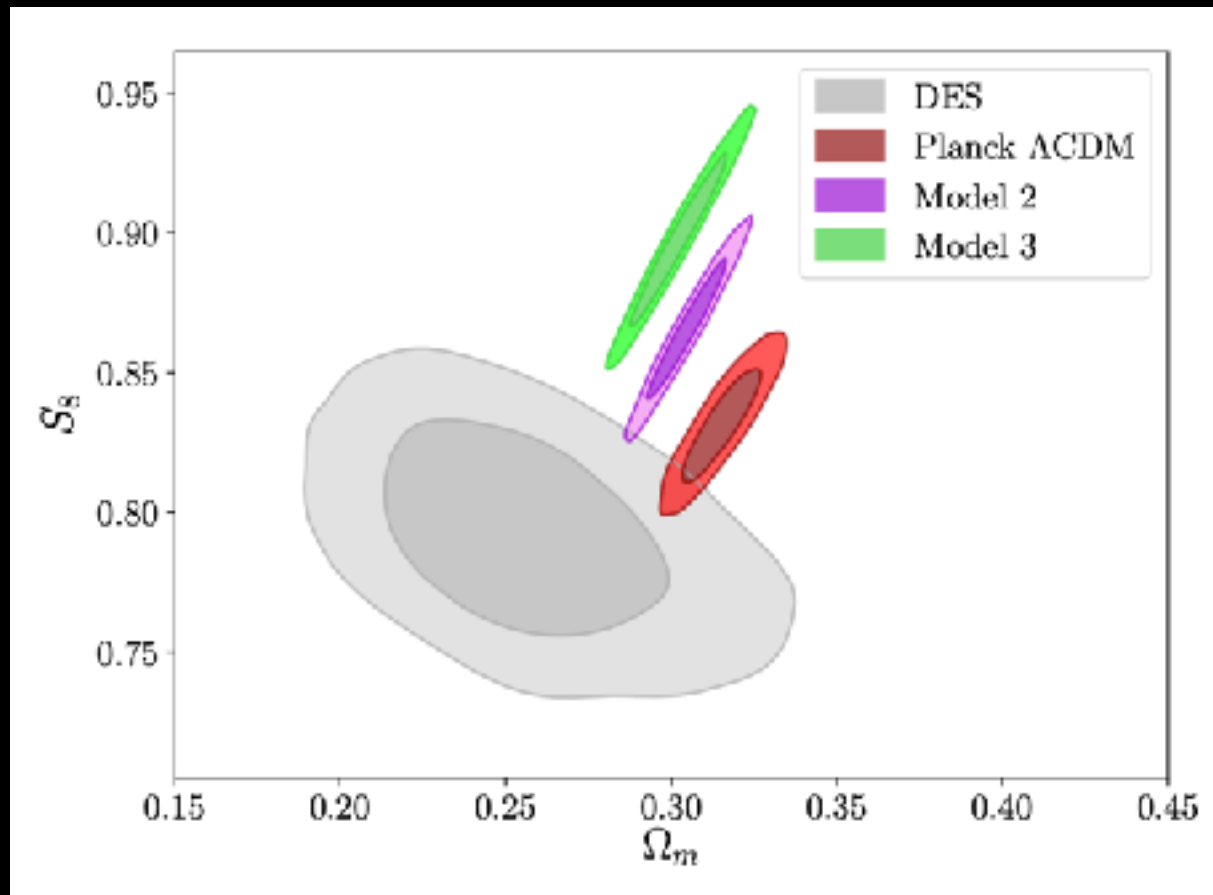
Early-time modifications worsen
the growth tension...



Jedamzik, Pogosian, Zhao (2020)

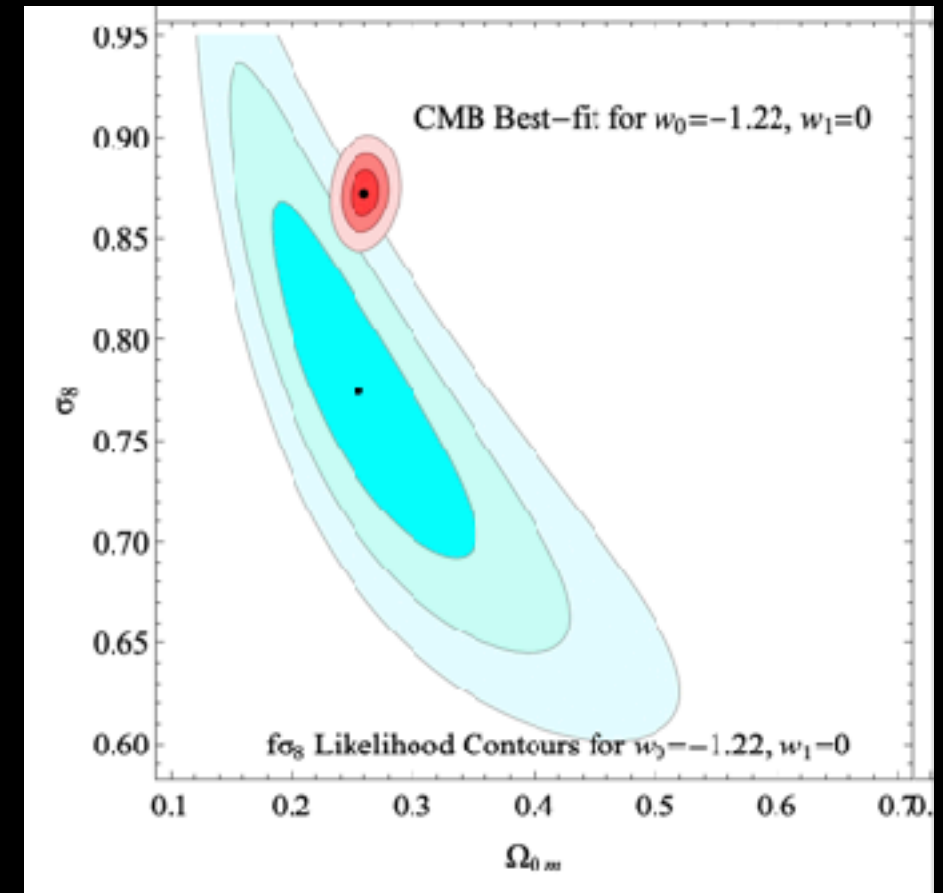
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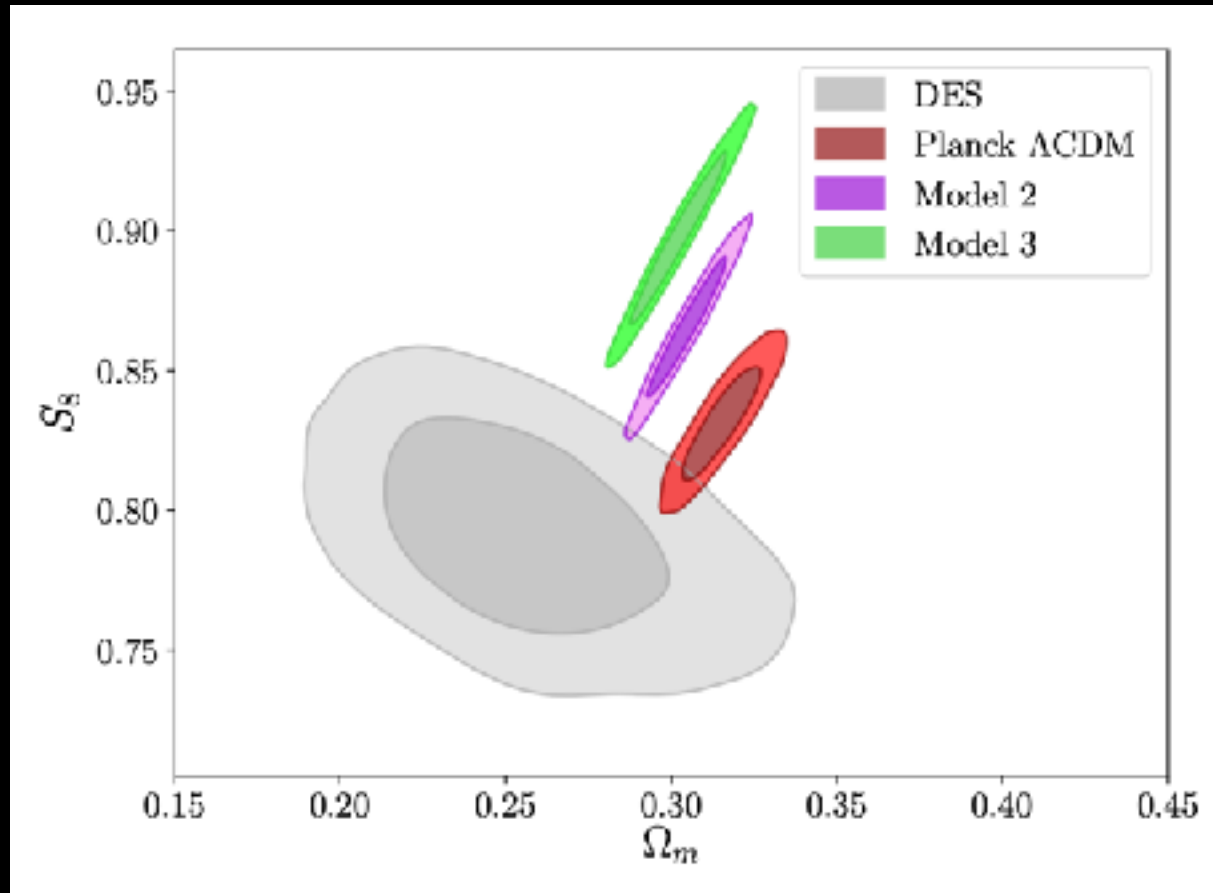
Aleatas & Perivolaropoulos (2021)



...and late-time modifications as well!

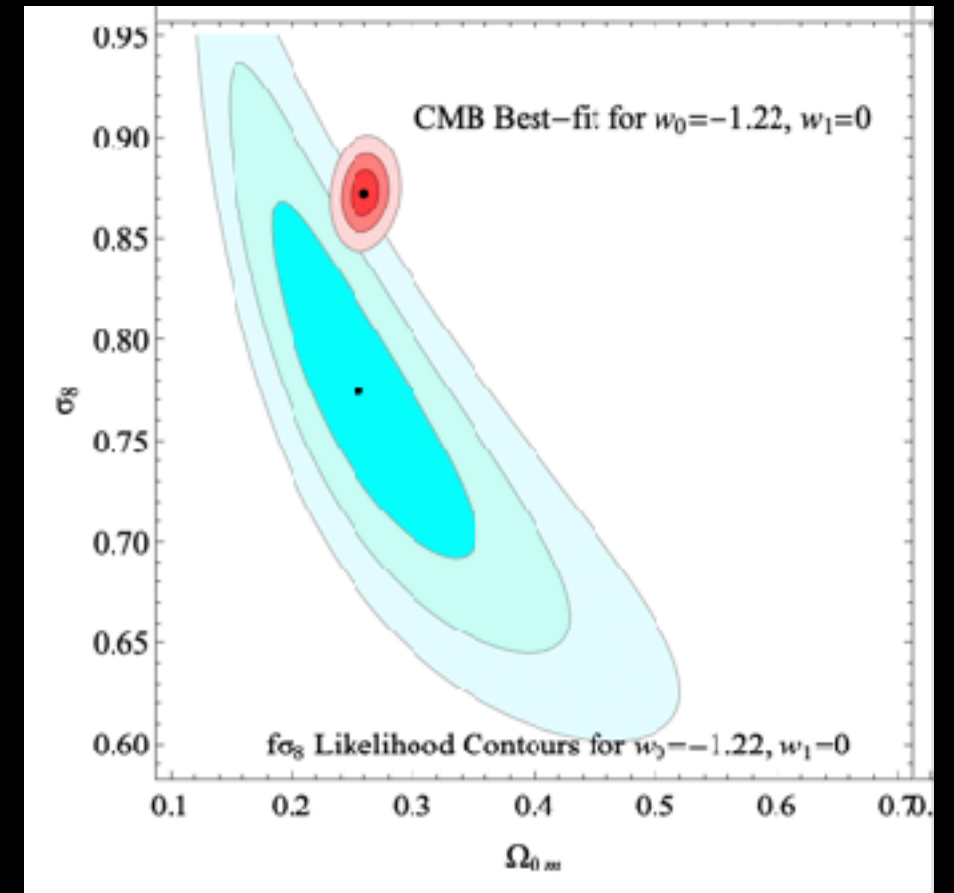
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We should adopt a holistic approach to resolve the Hubble tension

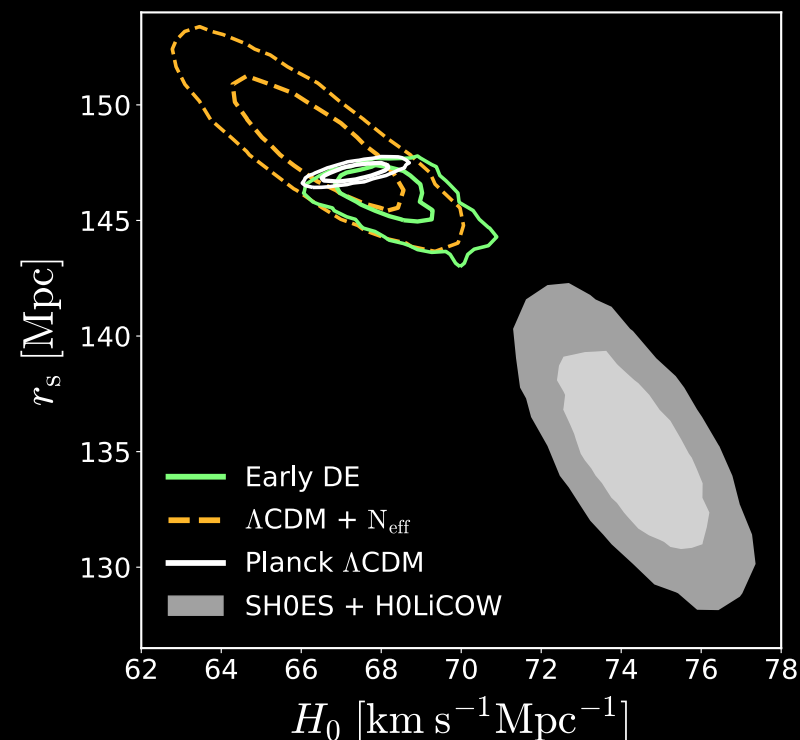
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NO

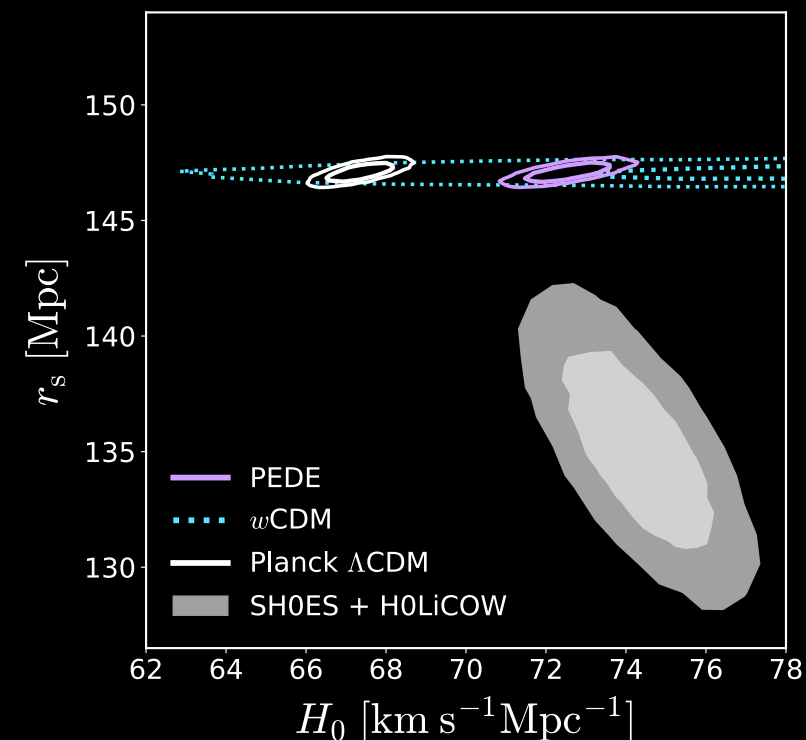
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Early-time modifications

- Small shift of H_0 and r_s to local values
- Worsen growth tension



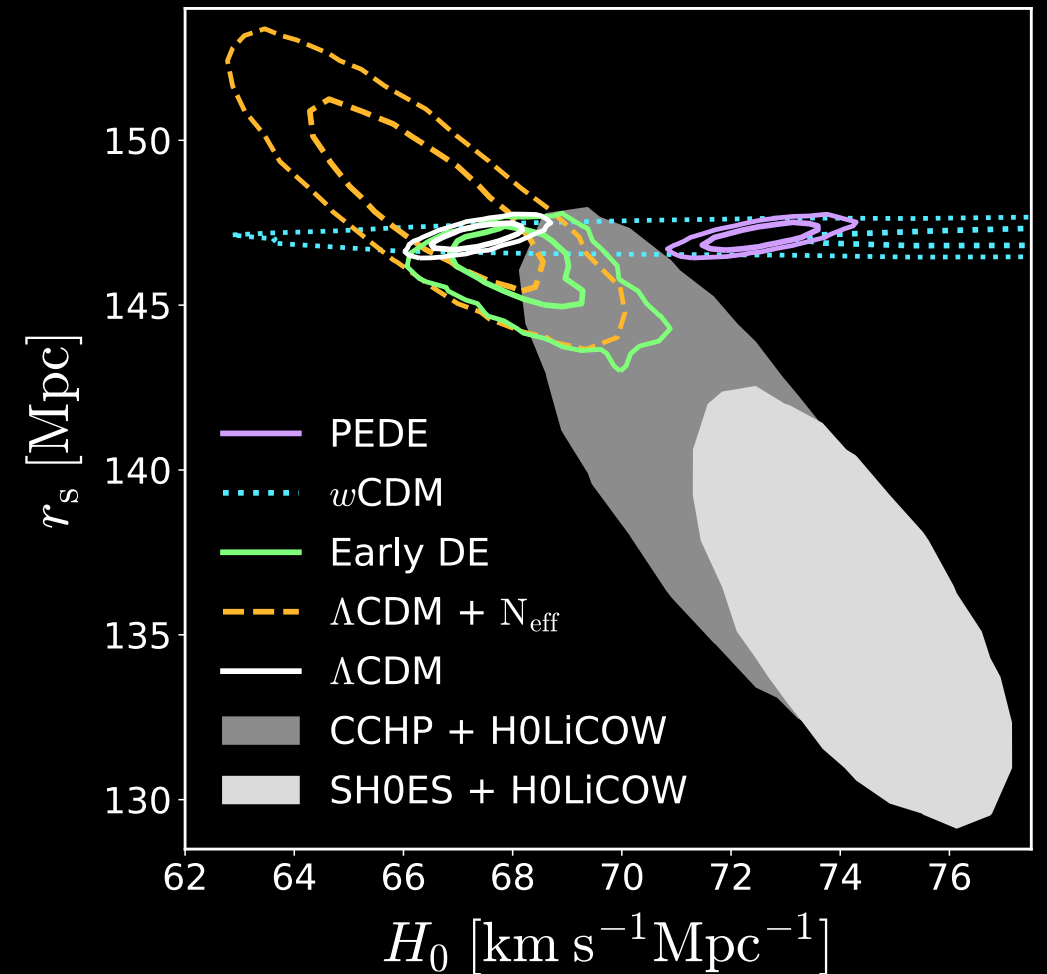
Late-time modifications

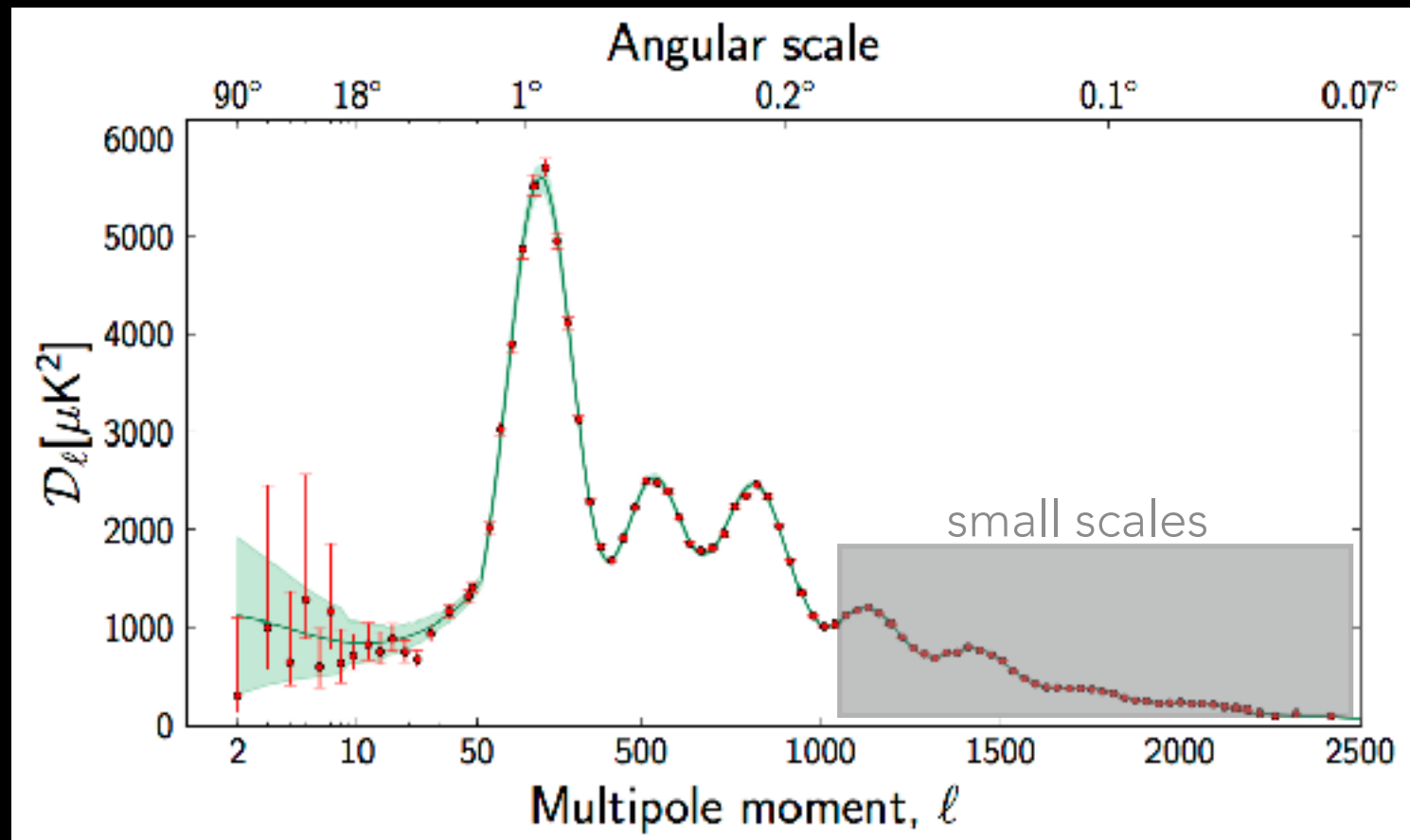
- Only solves tension in H_0 , not in r_s
- Worsen growth tension



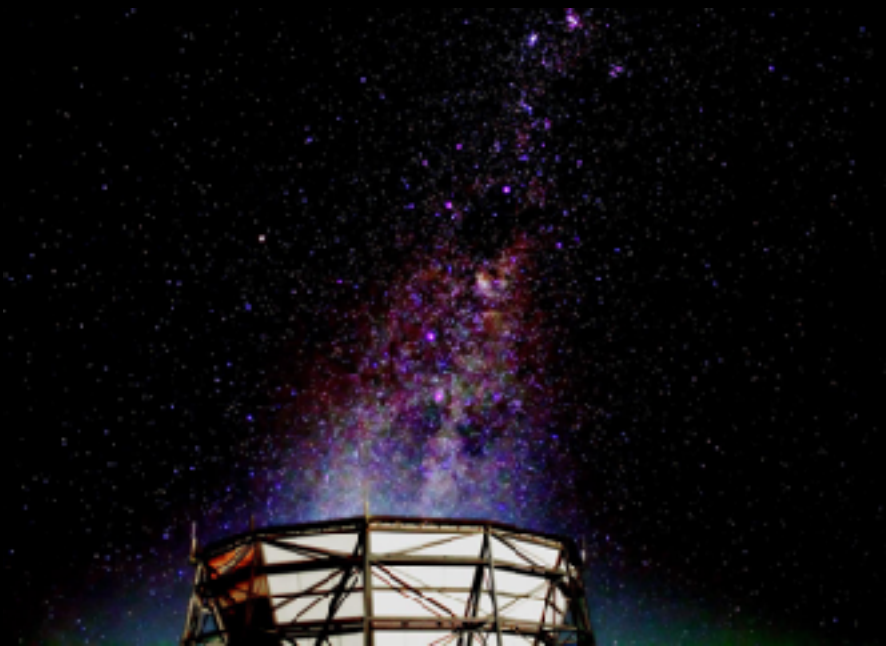
CONCLUSIONS

- There is a tension between local and CMB-based measurements of H_0 and r_s .
- The measurements using local data can be carried out without assuming a cosmological model.
- CMB-based measurements require a cosmological model. Extensions to Λ CDM can follow two strategies:
 - Changing the early (pre-CMB) physics slightly decreases the tension (but not enough).
 - Changing the late physics by adding some additional form of dark energy can only change H_0 and not r_s , and does therefore not solve the problem.





FUTURE CMB MEASUREMENTS

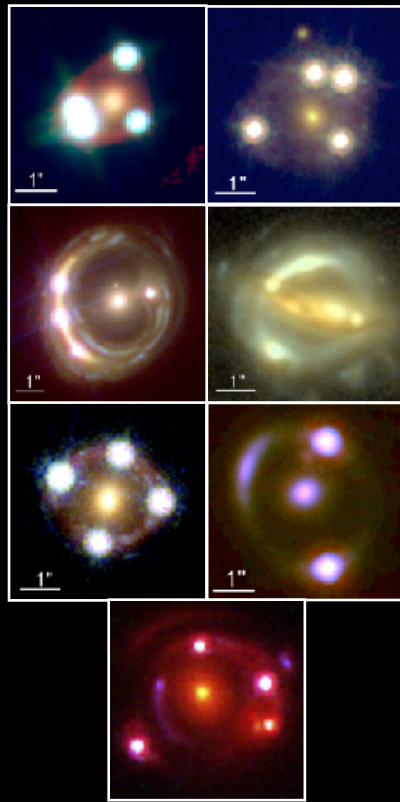


Atacama Cosmology Telescope (ACT)
Aiola et al. (2020)

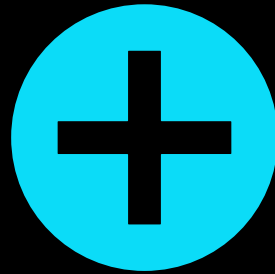


South Pole Telescope (SPT)

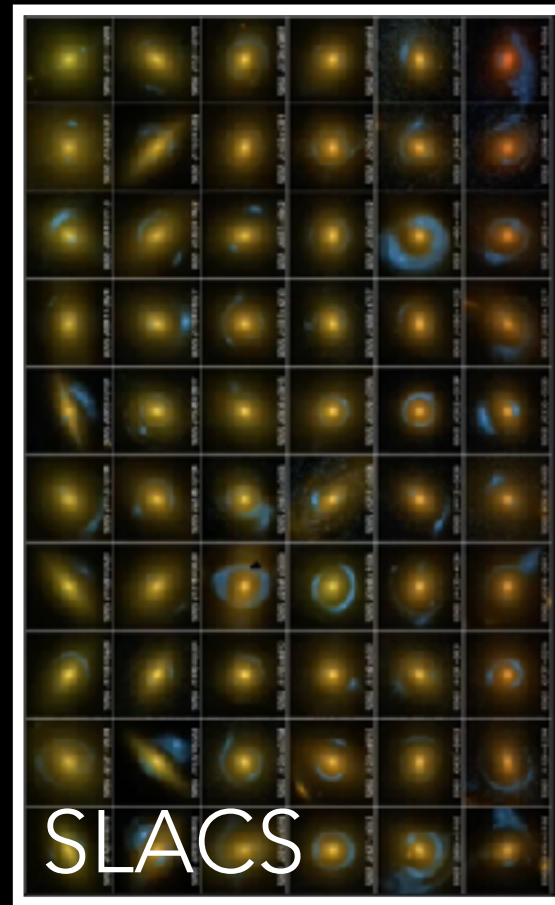
TD COSMO



HOLICOW + STRIDES



Millon et al. (2020)
Birrer et al. (2020)



Hubble constant:

67.4

± 4 km/s/Mpc

Precision forecast:

Currently **5%**

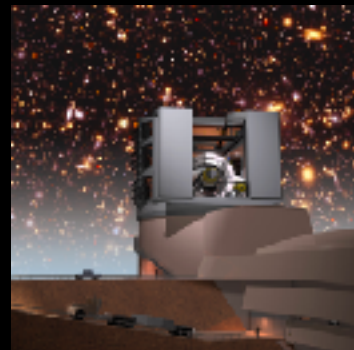
Future **1.5%**

Birrer & Treu (2020)

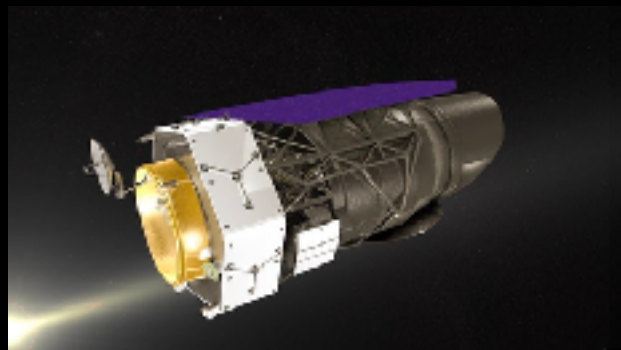
NEXT-GENERATION MISSIONS

10^5 galaxy-galaxy lenses

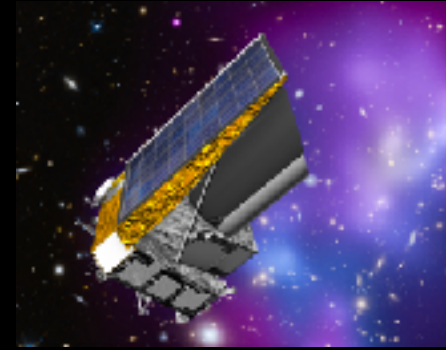
10^3 quasar-galaxy lenses



Rubin Observatory



Roman Telescope

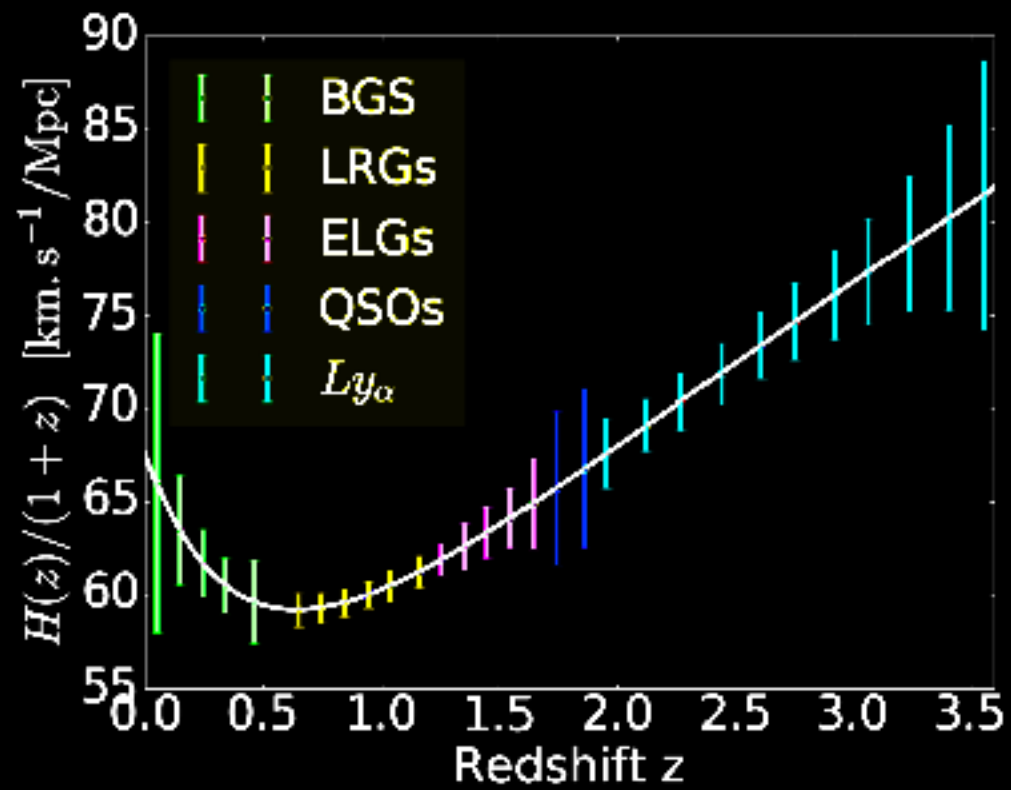


EUCLID (2022)



DESI

DESI

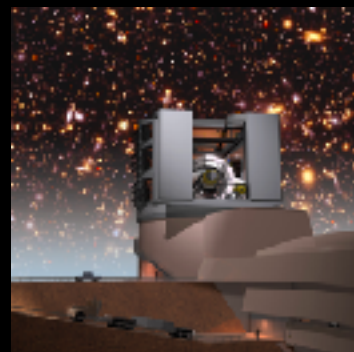


Vargas-Magaña, DESI collaboration, et al. (2019)

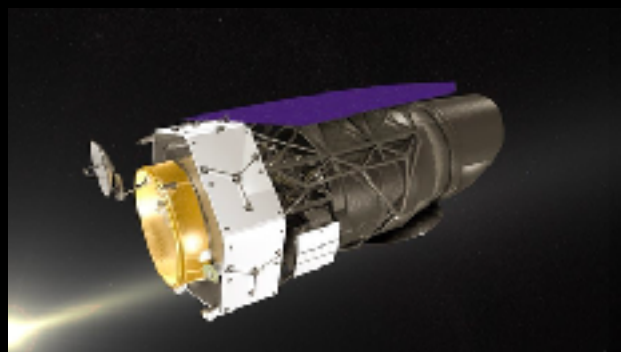
Order of magnitude
improvement on $H_0 r_s$

Pogosian et al. (2020)

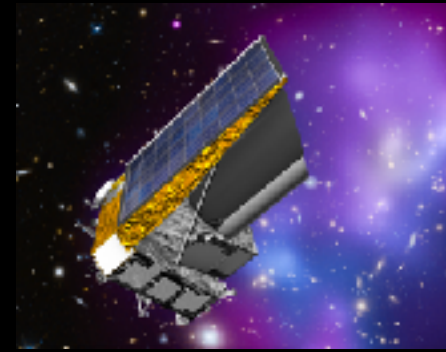
NEXT-GENERATION MISSIONS



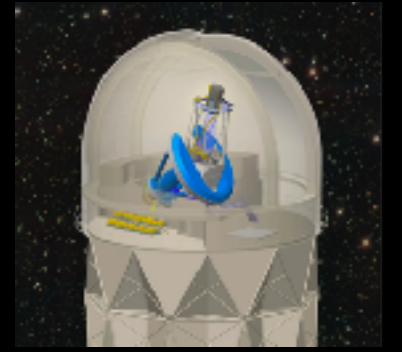
Rubin
Observatory



Roman Telescope



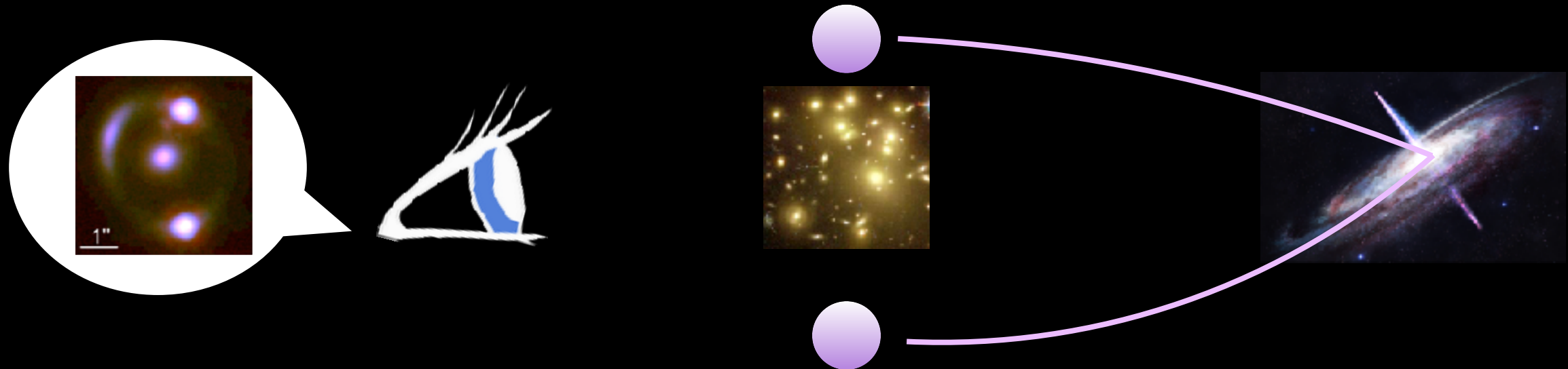
EUCLID (2022)



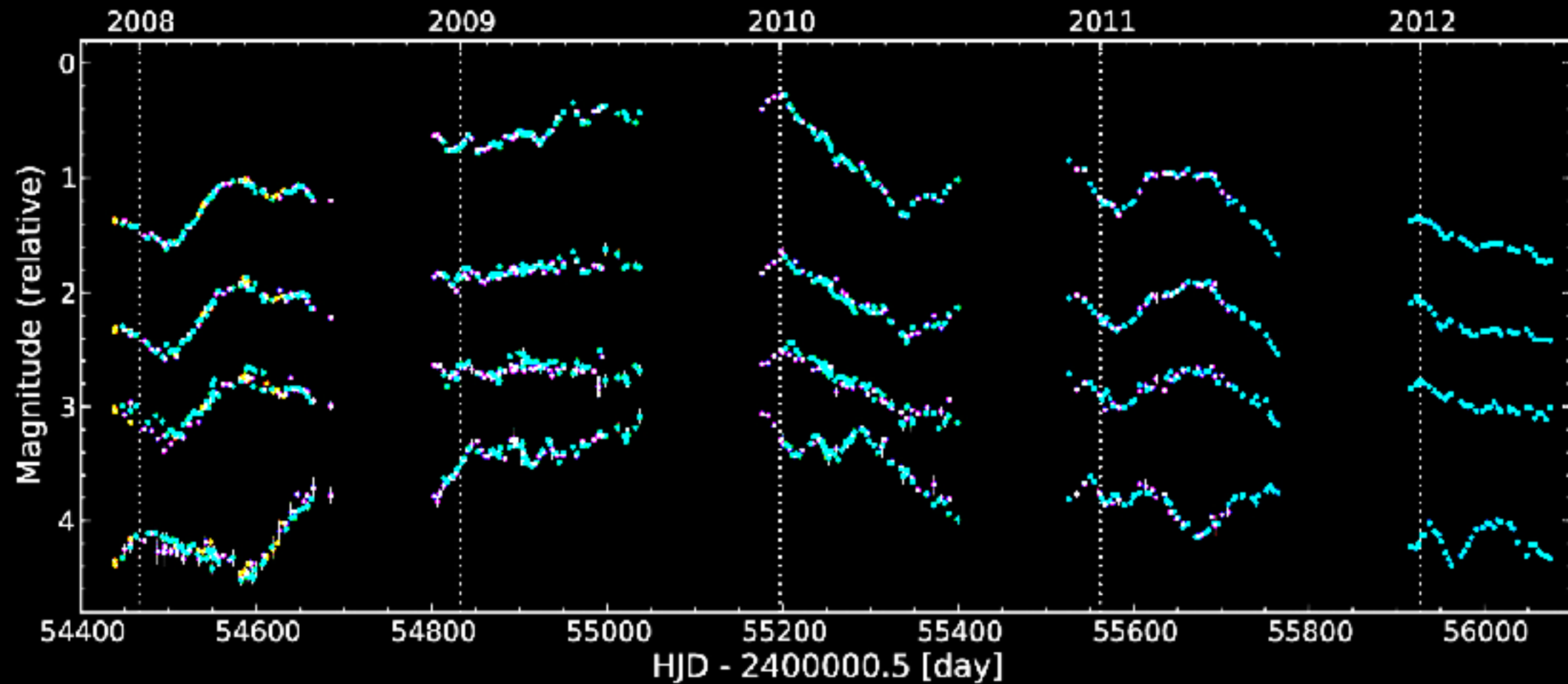
DESI

HOLICOW

H0 Lenses in COSMOSGRAB Wellspring



RXJ1131



HOLICOW

H0 Lenses in COSMOGRAIL Wellspring

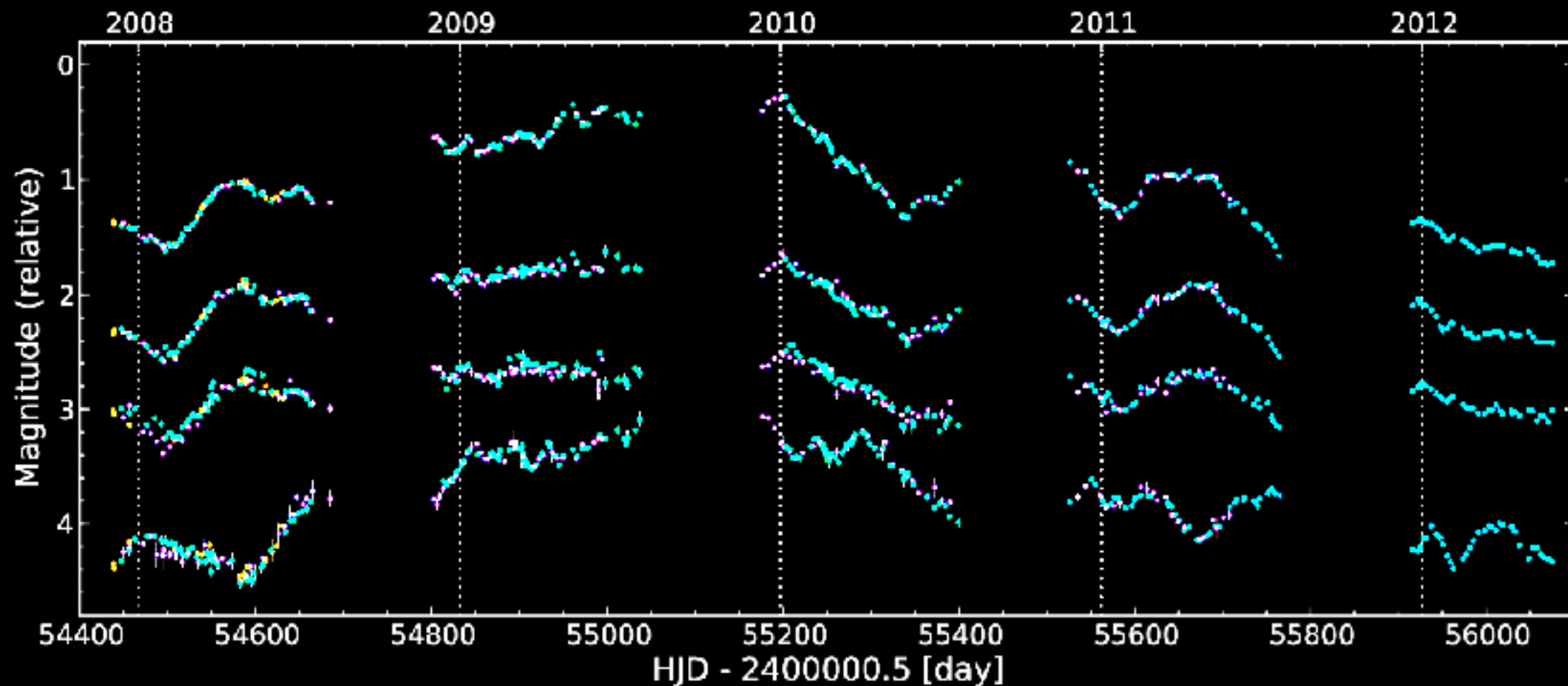

Time delay relative to the case of no lensing:

$$t = \frac{D_{\Delta t}}{c} \phi_{\text{lens}}$$

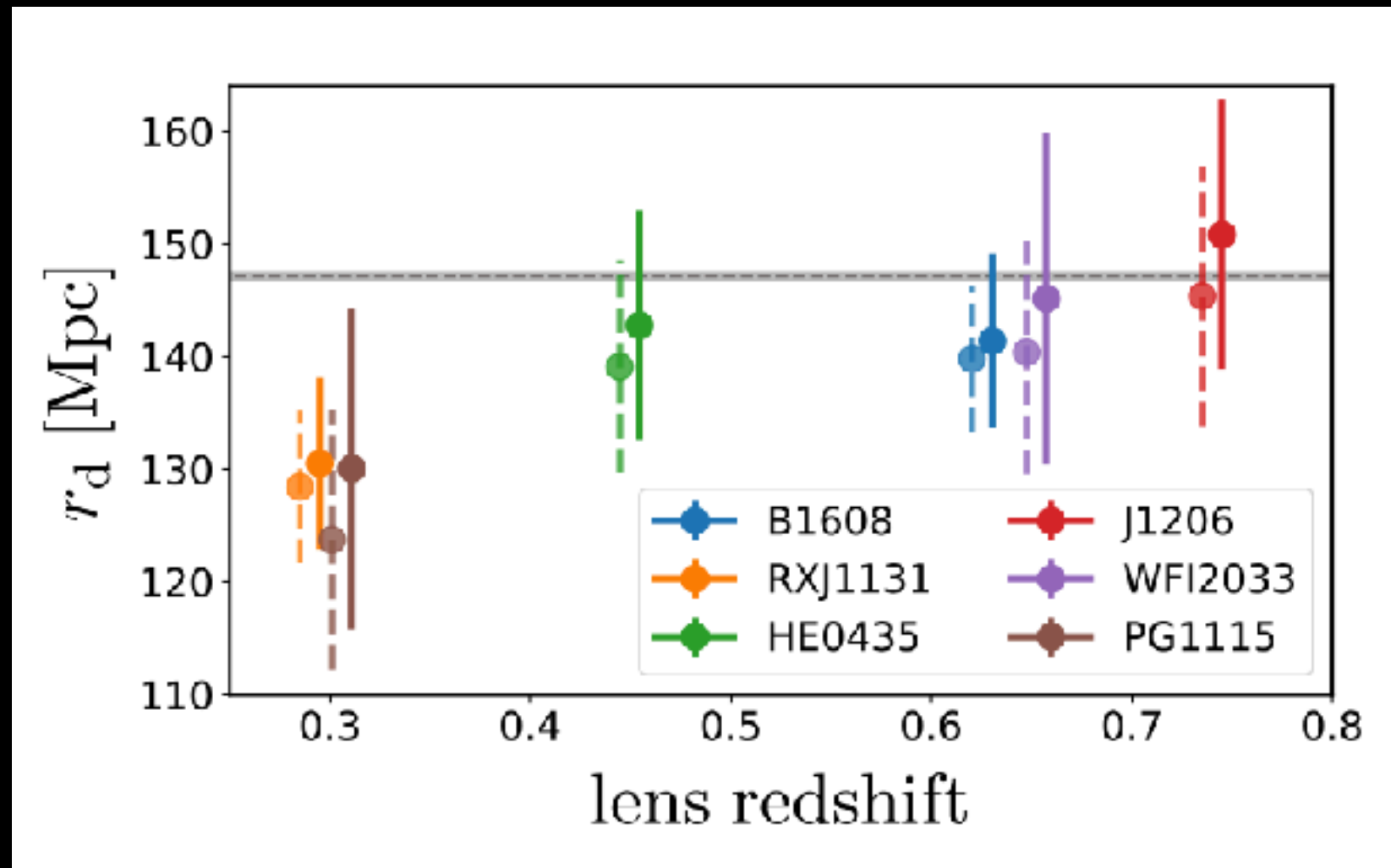
Time delay distance:

$$D_{\Delta t} = (1 + z_l) \frac{D_l D_s}{D_{ls}}$$

Absolute distances!



Redshift trend in the H0LiCOW lens sample

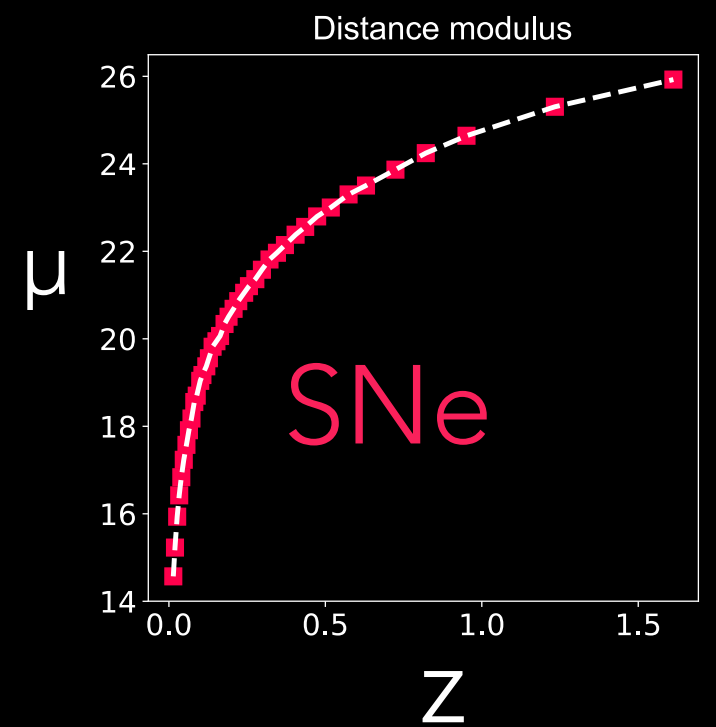
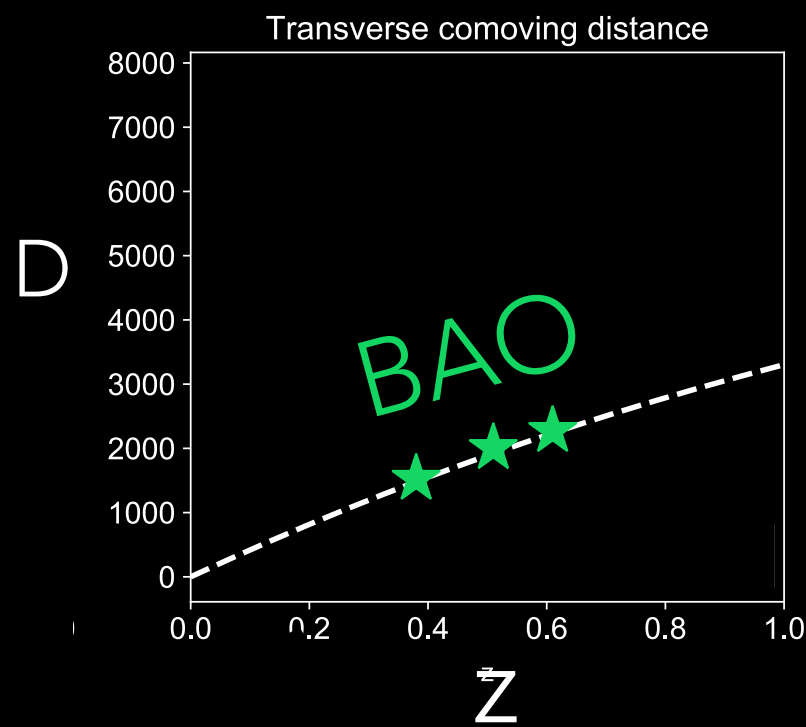
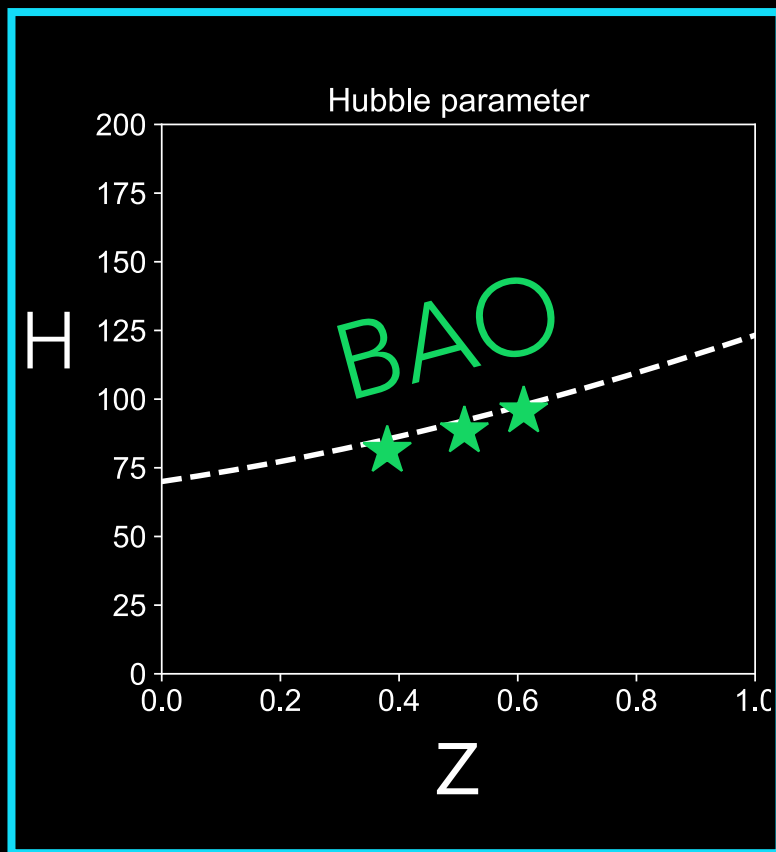


Friedman equations:

$$H(z) = H_0 \sqrt{\Omega_m (1+z)^3 + \Omega_\Lambda}$$

Taylor expansion:

$$H(z) = H_0 + c_1 z + c_2 z^2 + \mathcal{O}(z^3) \quad c_1 = \left. \frac{dH}{dz} \right|_0 \quad c_2 = \frac{1}{2} \left. \frac{d^2 H}{dz^2} \right|_0$$



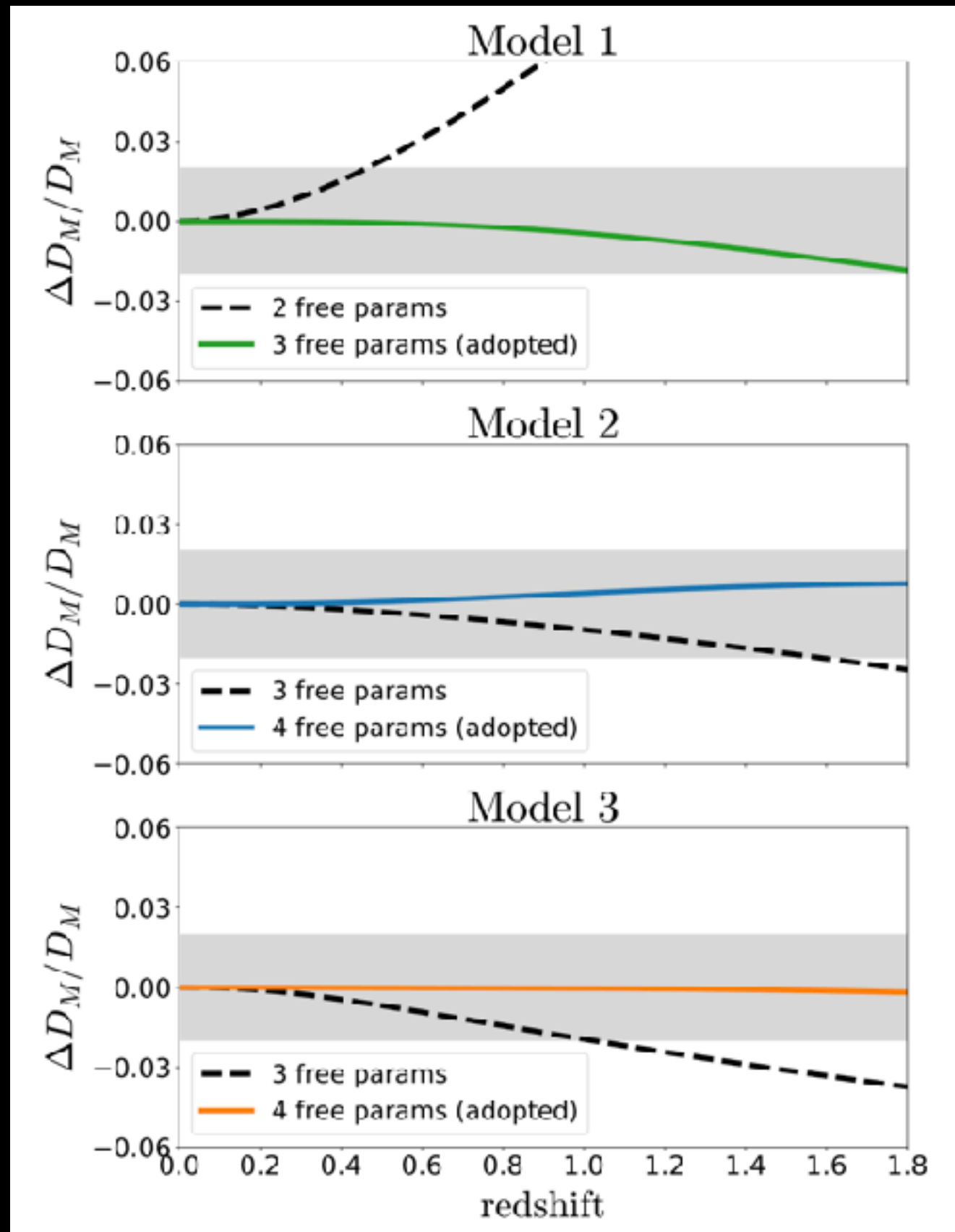
integrate

logarithm

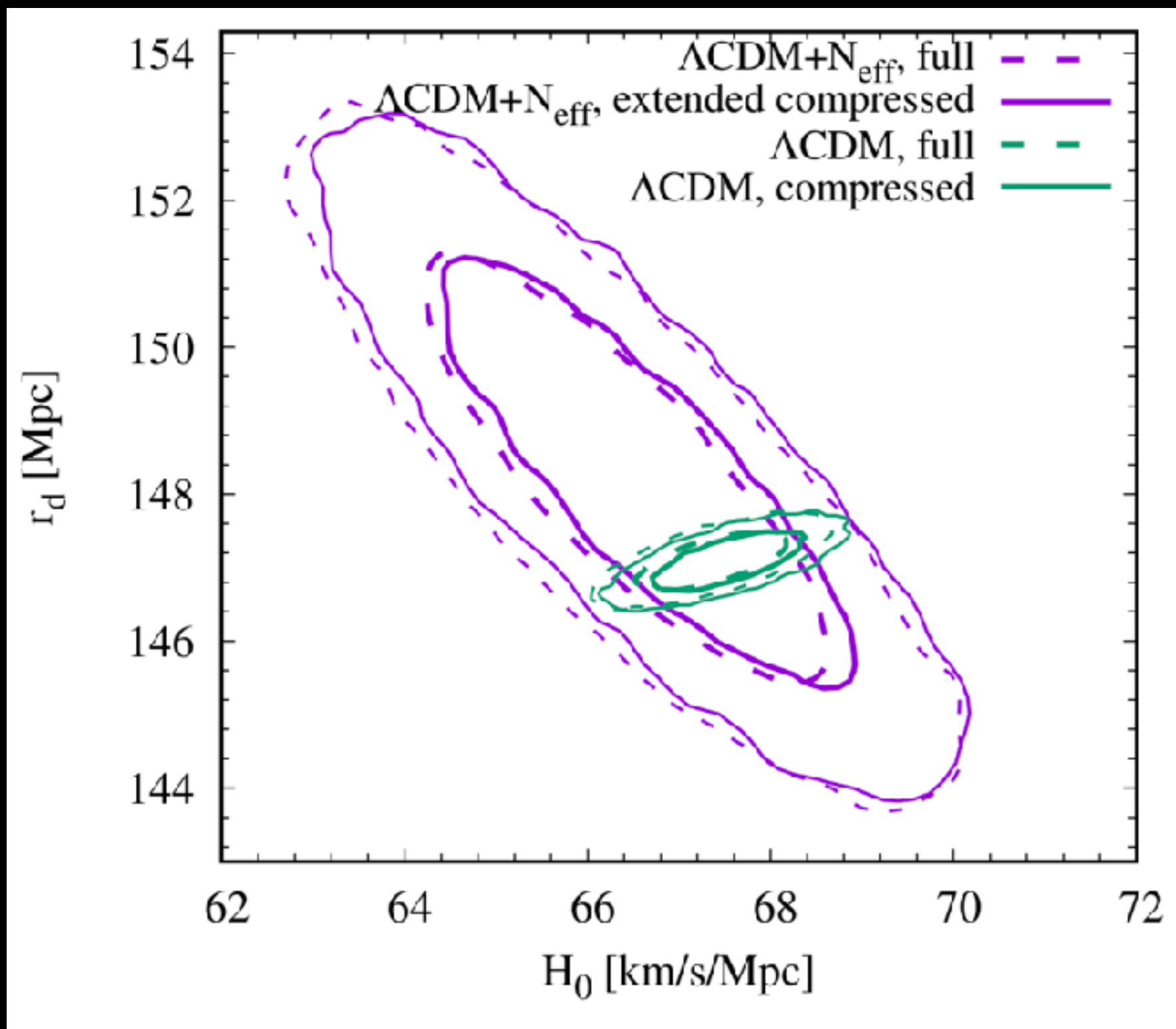
POLYNOMIAL PARAMETRIZATIONS

Model	Formula
1	$H(z) = H_0 (1 + b_1 z + b_2 z^2)$
2	$D_L(z) = \frac{c \ln(10)}{H_0} \left(\log(1 + z) + c_2 [\log(1 + z)]^2 + c_3 [\log(1 + z)]^3 + c_4 [\log(1 + z)]^4 \right)$
3	$D_M(z) = \frac{c}{H_0} \left(\frac{z}{1+z} + d_2 \left[\frac{z}{1+z} \right]^2 + d_3 \left[\frac{z}{1+z} \right]^3 + d_4 \left[\frac{z}{1+z} \right]^4 \right)$
4	$H(z) = H_0 \sqrt{\Omega_M (1 + z)^3 + \Omega_\Lambda + \Omega_k (1 + z)^2}$

...and how they perform in our redshift range

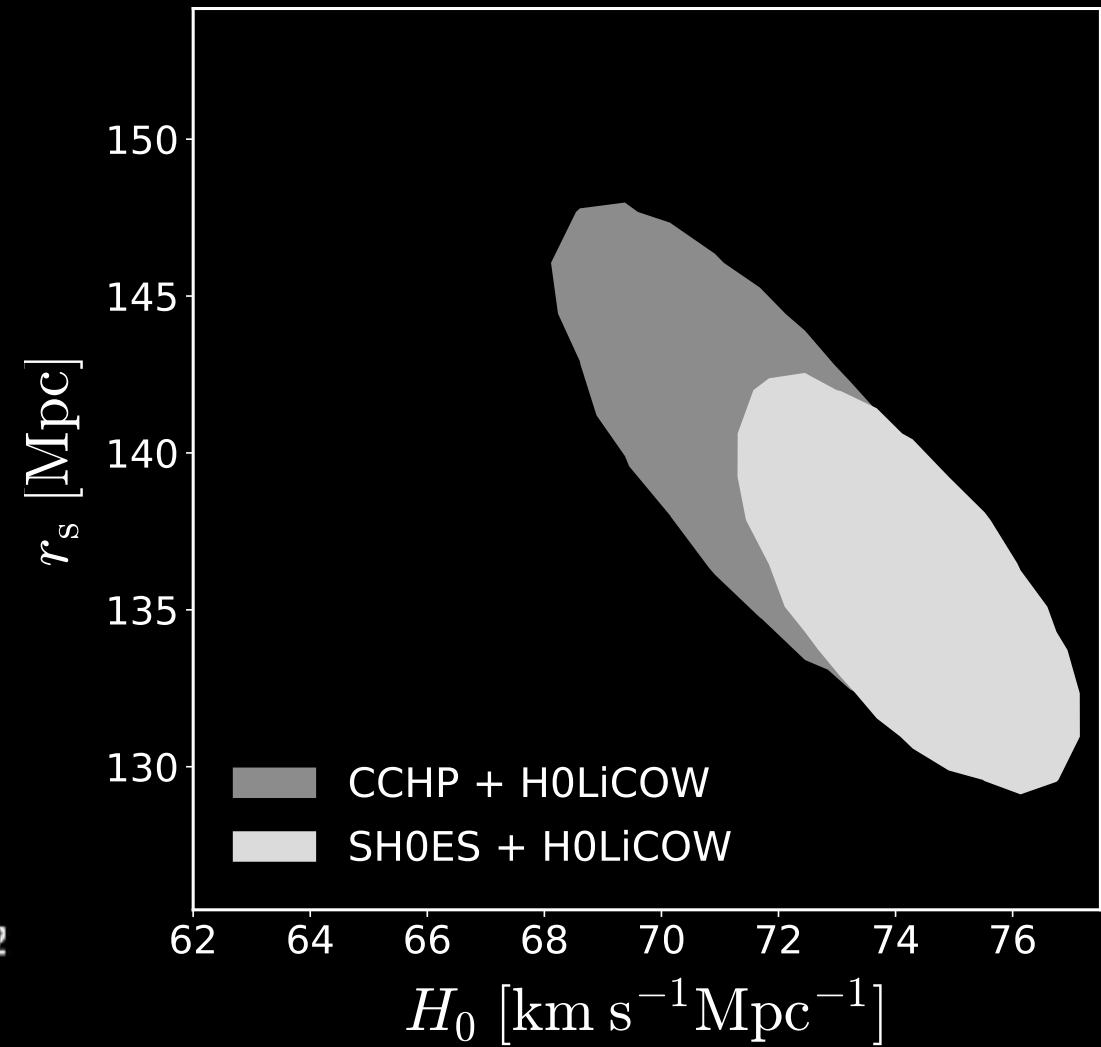
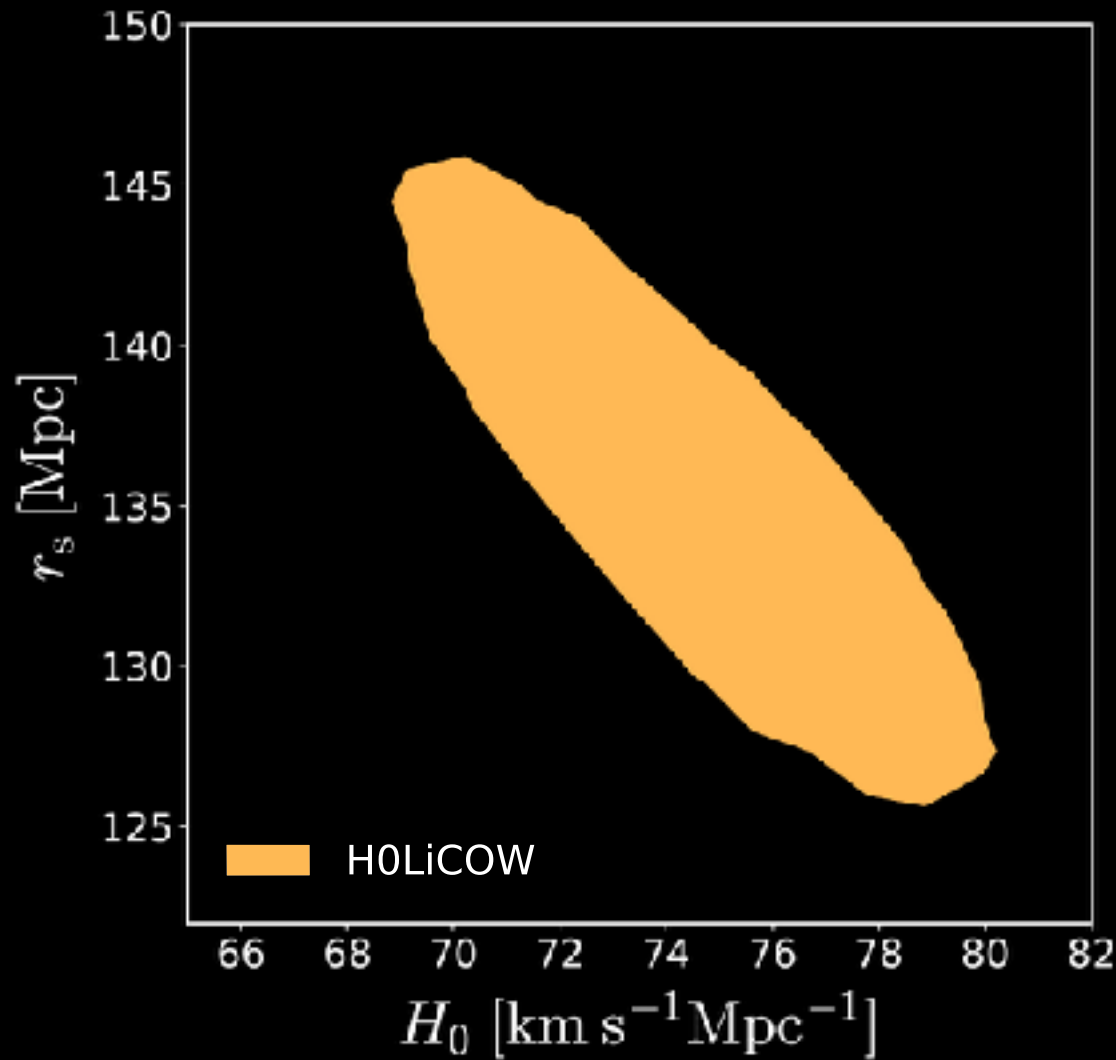


EXTENDED COMPRESSED PLANCK LIKELIHOOD



RESULTS

Arendse, Wojtak, Agnello et al. (2020)



H0LiCOW

$$H_0 = 74.5 \pm 2.2$$

$$r_s = 135.2 \pm 4.0$$

CCHP + H0LiCOW

$$H_0 = 71.9 \pm 1.5$$

$$r_s = 139.6 \pm 3.3$$

SHOES + H0LiCOW

$$H_0 = 74.2 \pm 1.2$$

$$r_s = 135.7 \pm 2.7$$

How to obtain $H(z)$ from D_M

From "General Cosmography Model with Spatial Curvature"

By En-Kun Li, Minghui Du, and Lixin Xu.

$$H(z, \Omega_K) = \frac{c}{\partial D_M(z)/\partial z} \sqrt{1 + \frac{H_0^2 \Omega_K}{c^2} D_M(z)^2},$$