

Searching for new physics through neutrino non-standard interactions

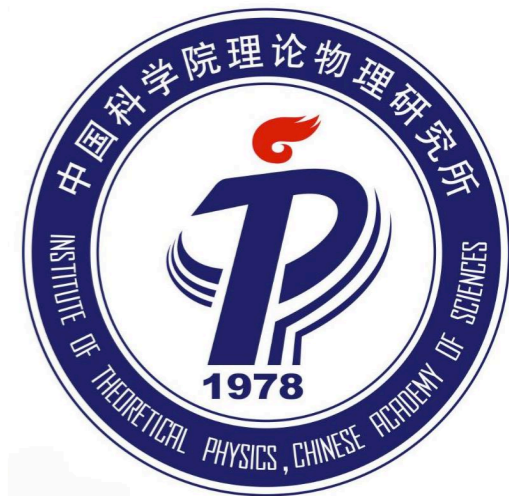
Yong Du

email: yongdu@itp.ac.cn

BSM-2021, March 31 2021

YD, H-L. Li, J. Tang, S. Vihonen, J-H. Yu, JHEP 03(2021) 019

YD, J-H. Yu, arXiv: 2101.10475 (To appear in JHEP)



Overview

* Not to be complete

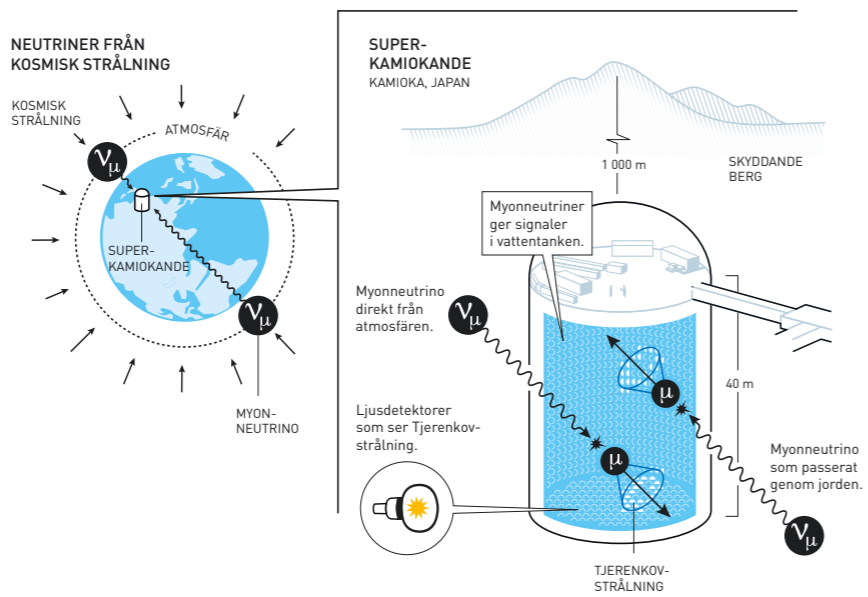


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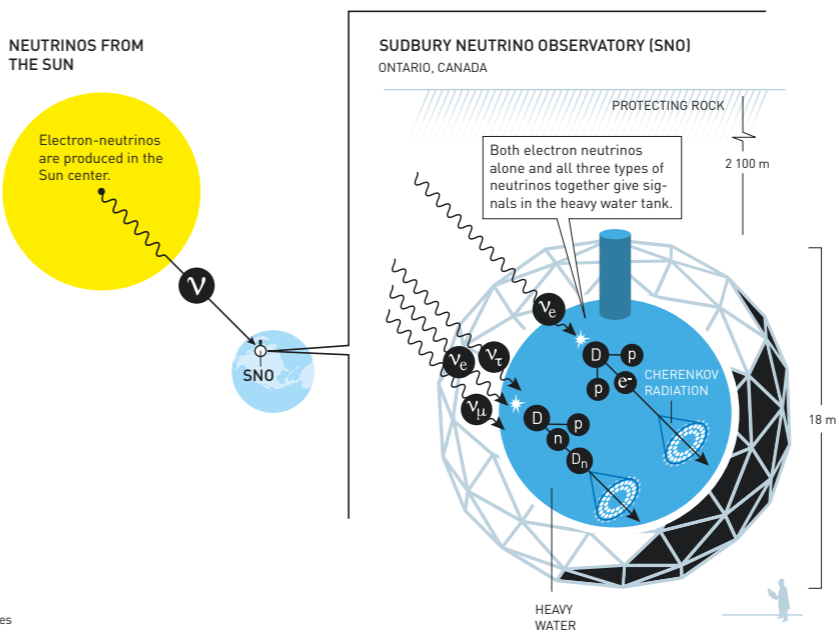
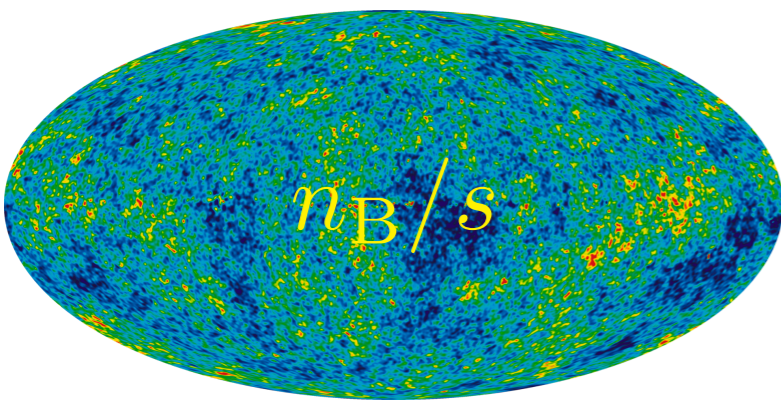


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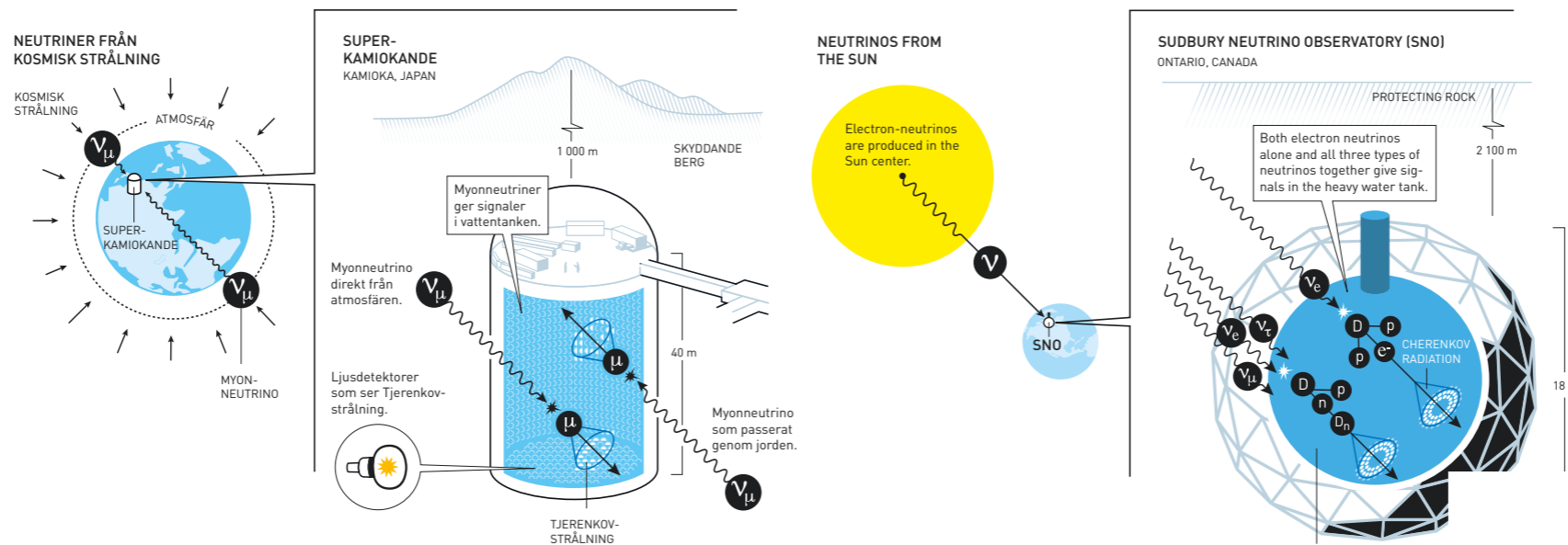
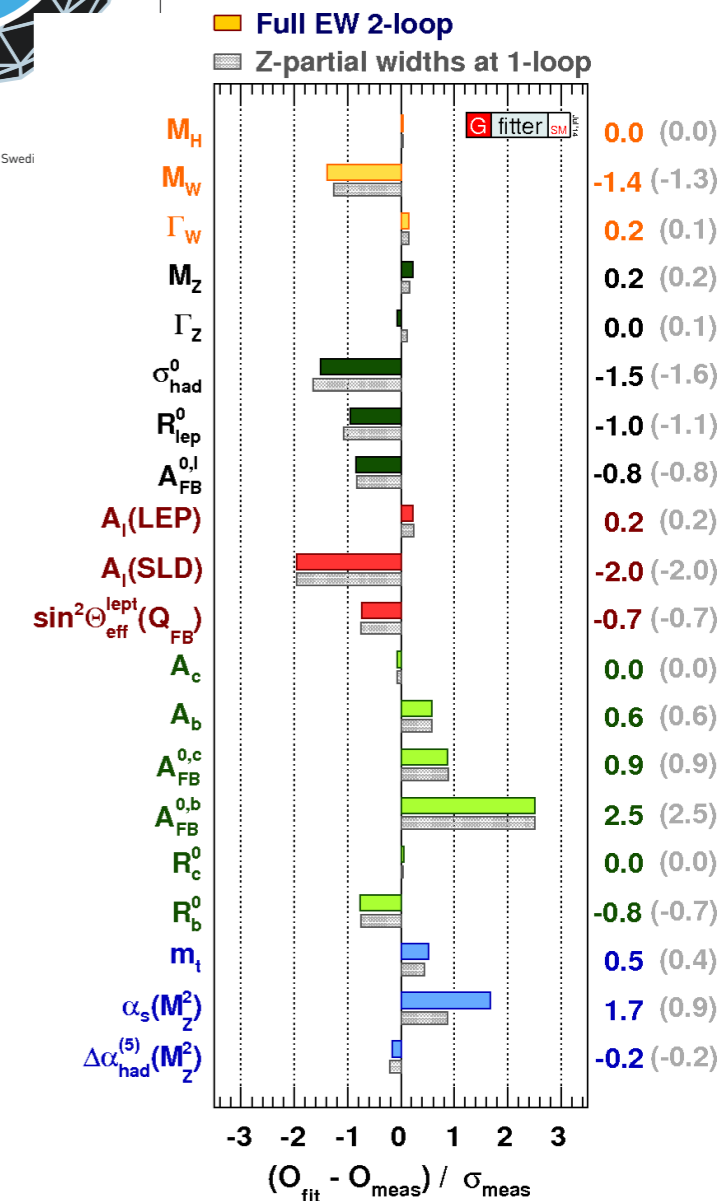
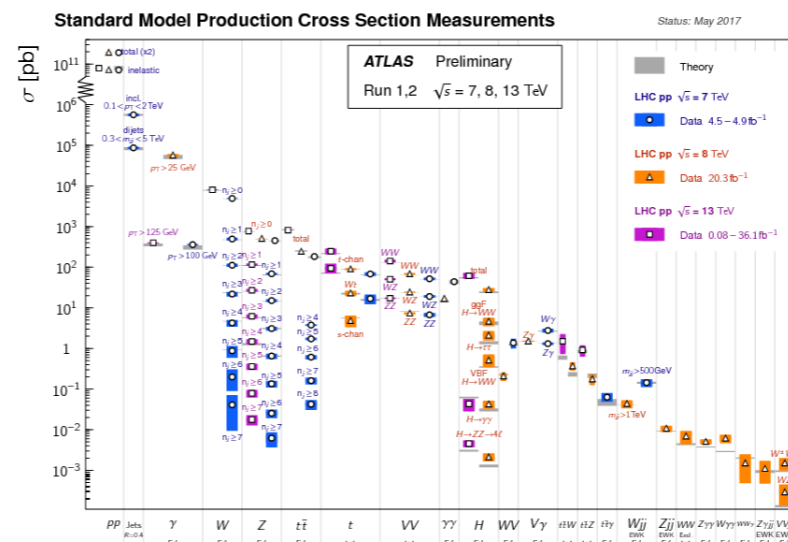
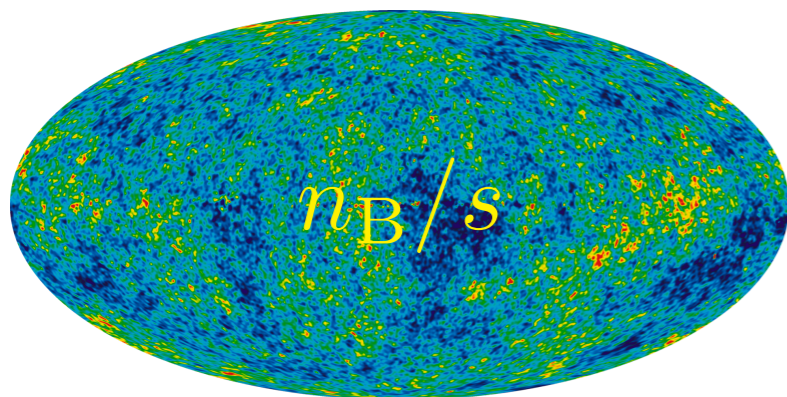
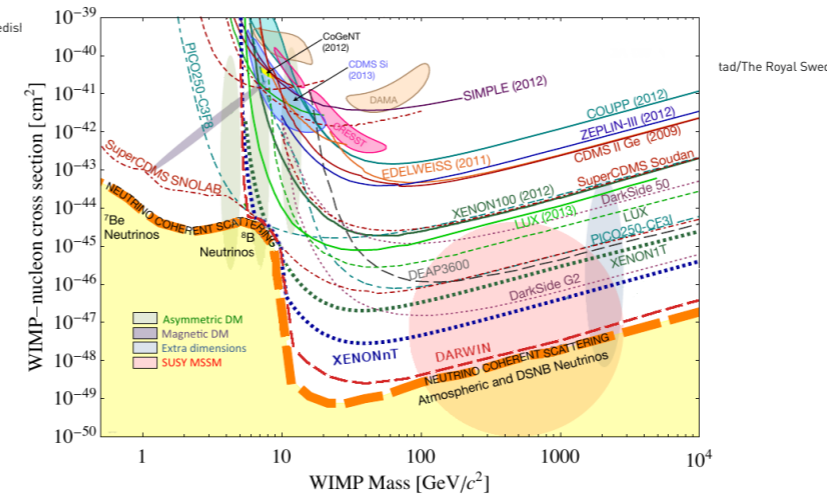
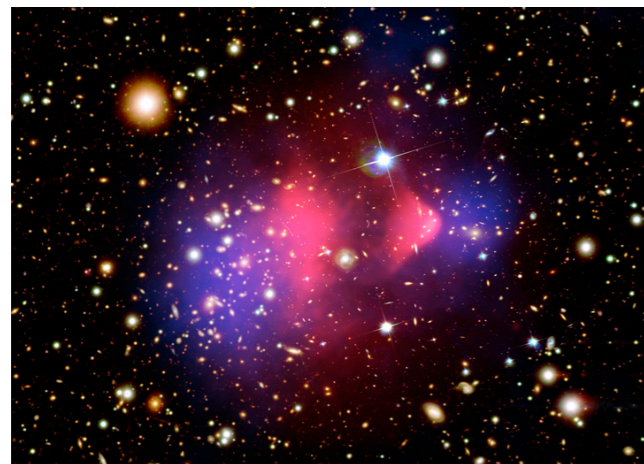


Illustration: © Johan Jarnestad/The Royal Swedisl



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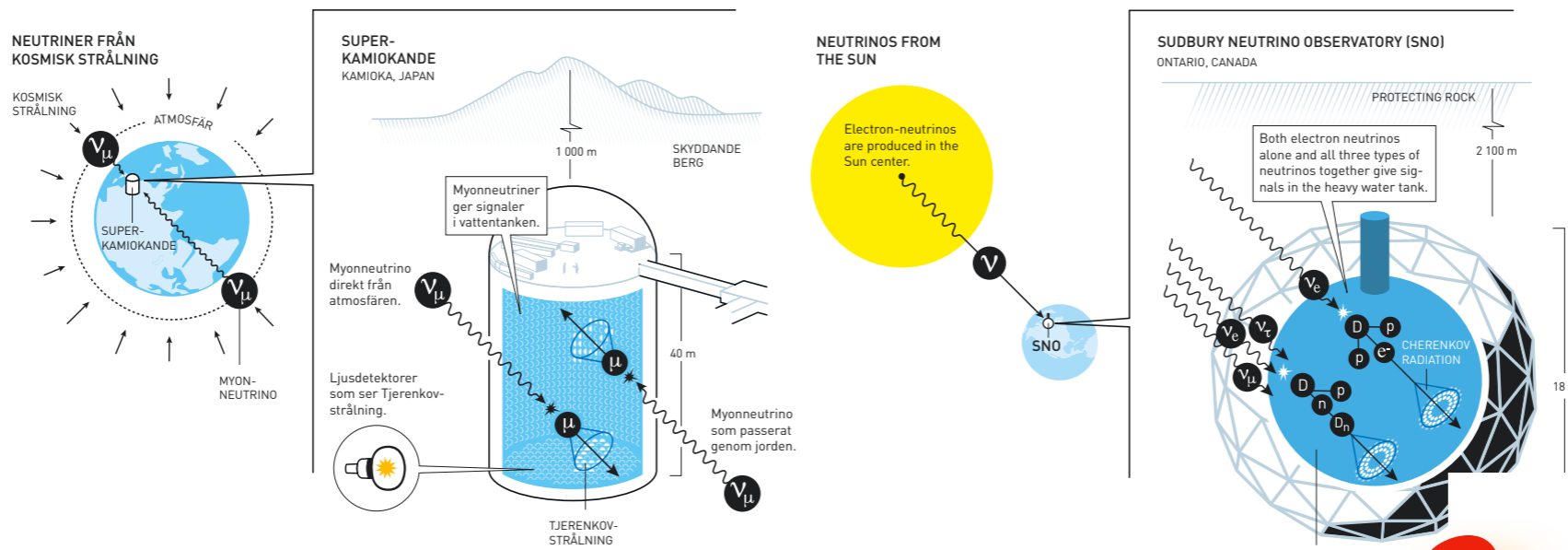
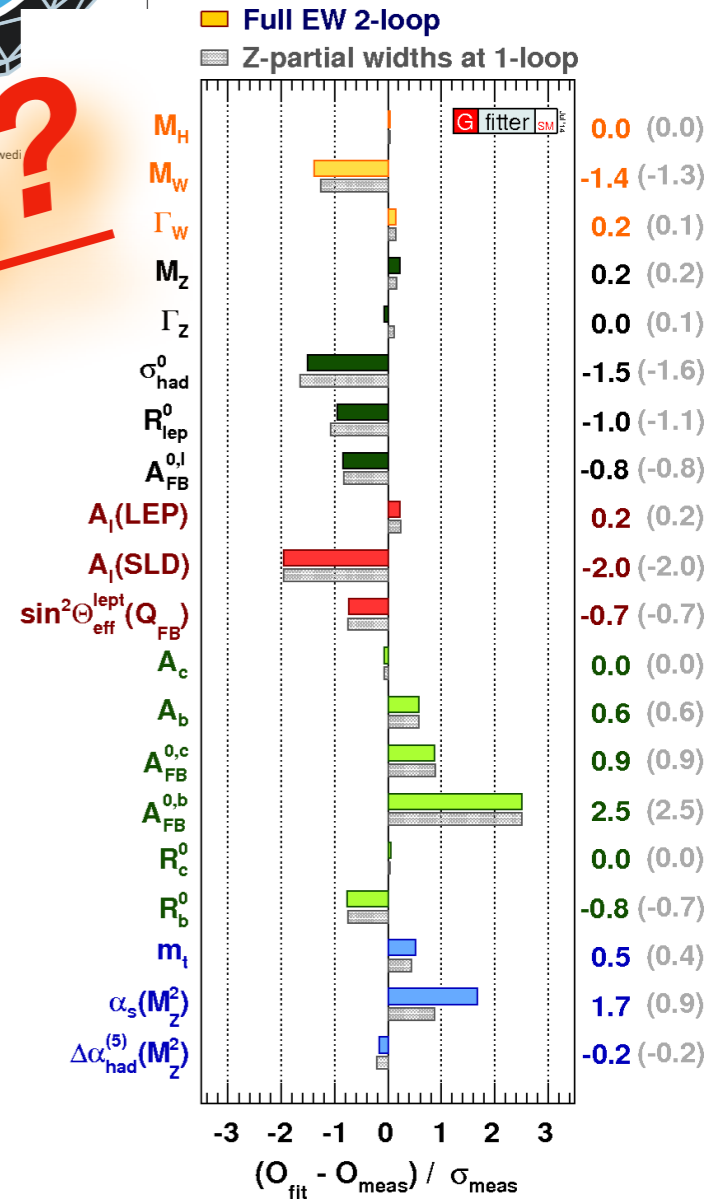
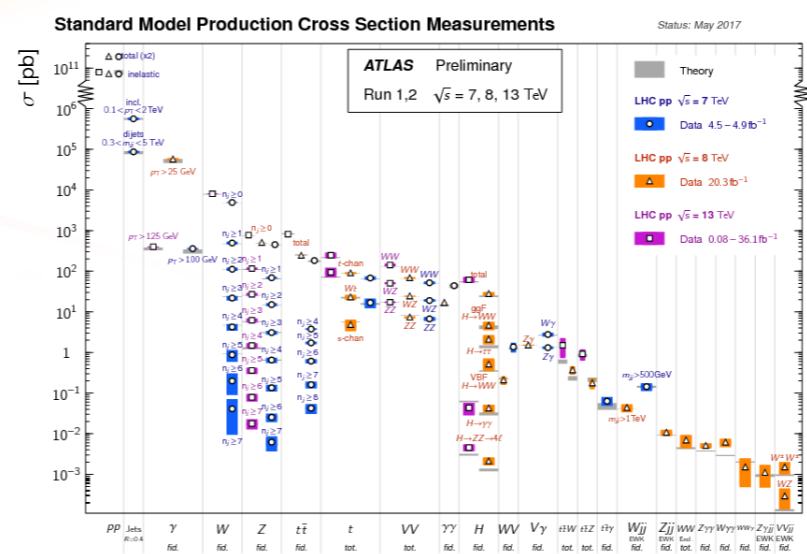
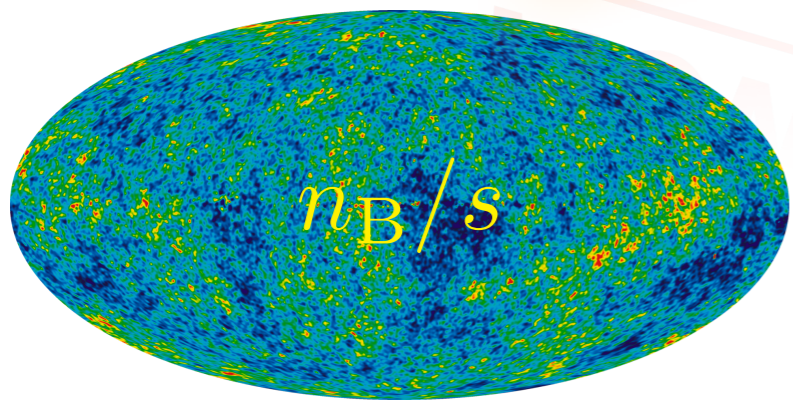
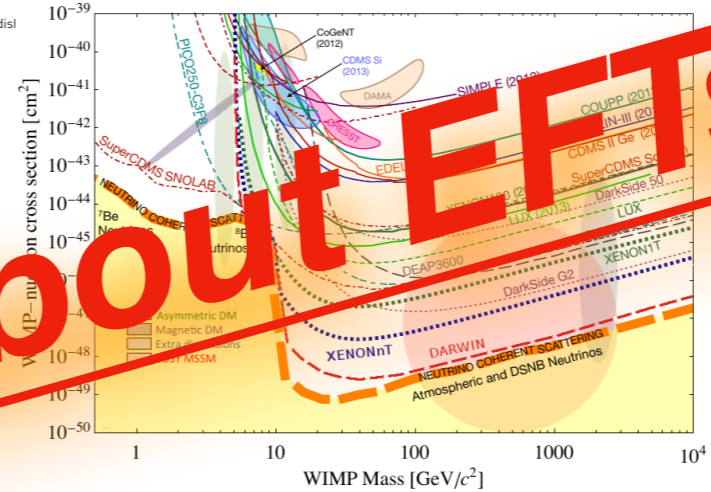


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How about EFTs?



Overview

In this talk, I will only focus on **neutrino NSIs** from an EFT approach

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❖ **Charge-Current (CC) NSIs**: from terrestrial neutrino oscillation experiments (dim-6 SMEFT operators only)

YD, H-L. Li, J. Tang, S. Vihonen, J-H. Yu, JHEP 03(2021) 019

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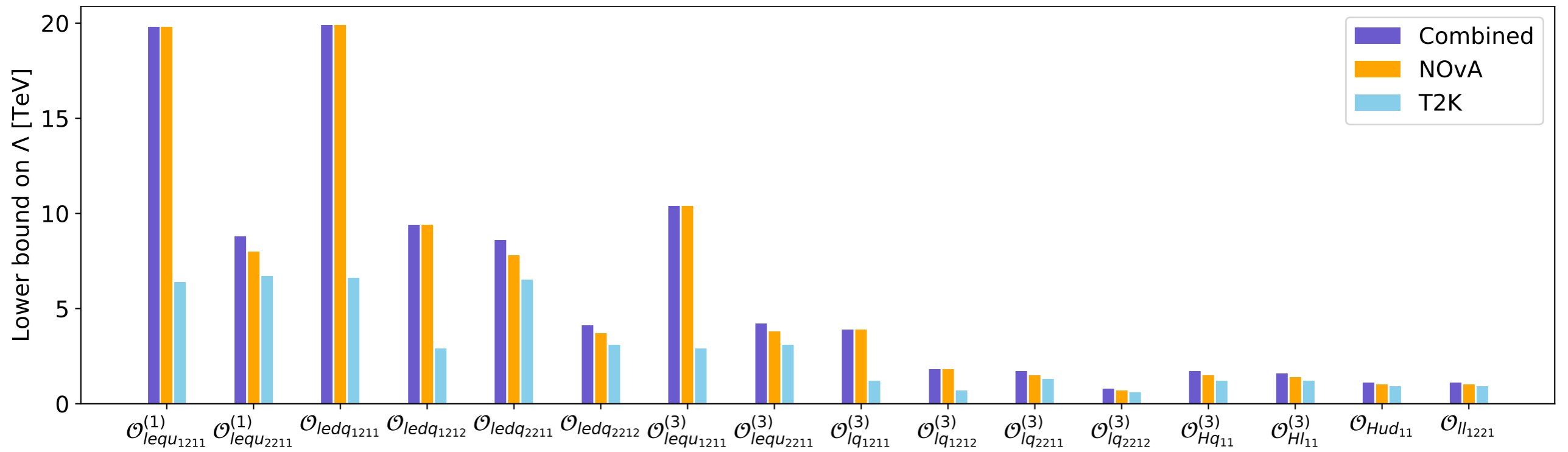
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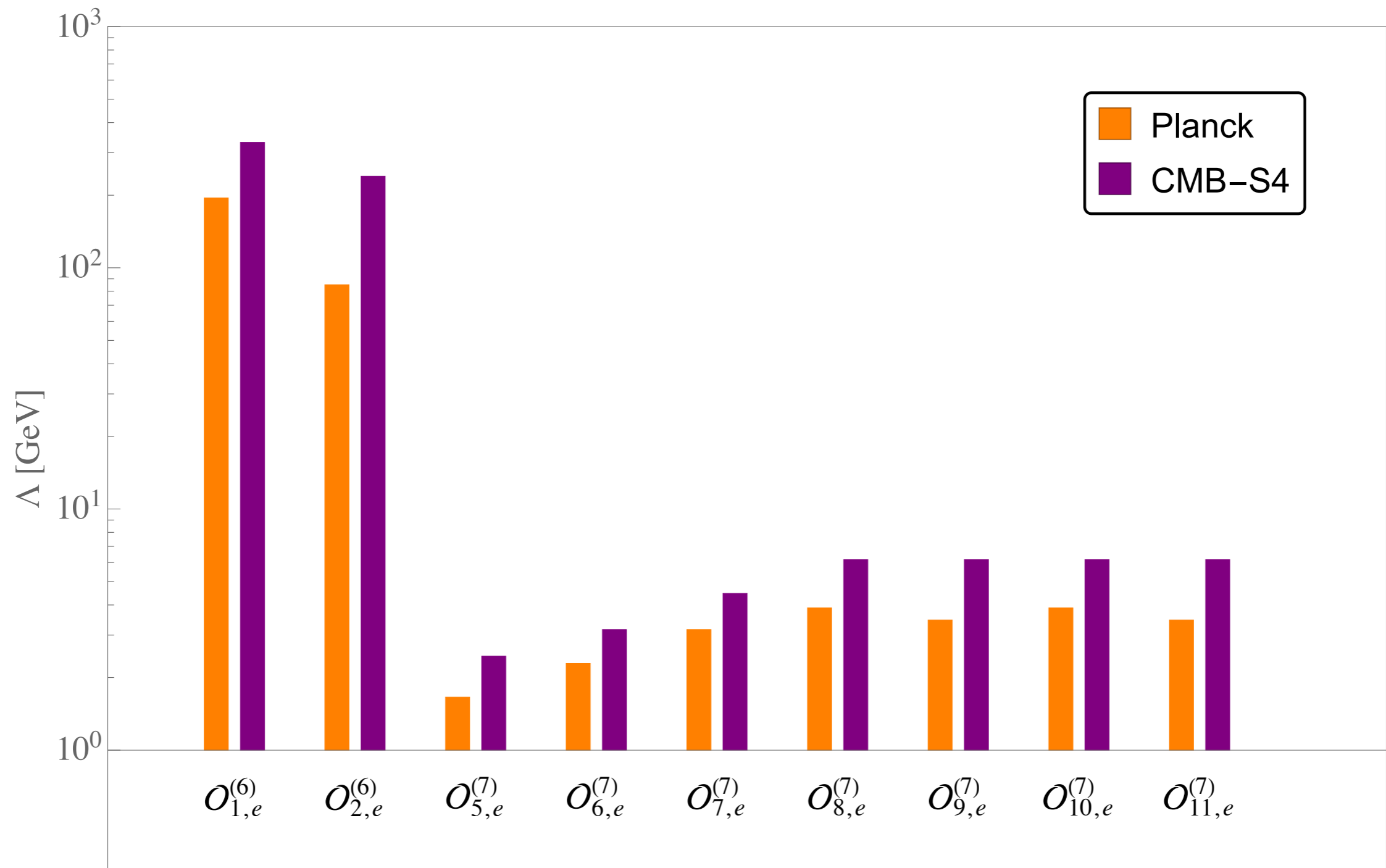
❖ **Neutral-Current (NC) NSIs**: Neff from Planck and CMB-S4 (ν - ν , ν - e , ν - γ operators up to dim-7)

YD, J-H. Yu, arXiv: 2101.10475 (To appear in JHEP)

Spoiler: CC NSIs



Spoiler: NC NSIs



CC NSIs

What neutrino experimentalists measure: Mismatch between production and detection

QM: Production/detection parameters

$$|\nu_\alpha^s\rangle = \frac{(1 + \epsilon^s)_{\alpha\gamma}}{N_\alpha^s} |\nu_\gamma\rangle, \quad \langle\nu_\beta^d| = \langle\nu_\gamma| \frac{(1 + \epsilon^d)_{\gamma\beta}}{N_\beta^d}$$

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NSI parameter	Upper bound	Experiments
$ \epsilon_{\mu e}^s $	0.004	
$ \epsilon_{\mu\mu}^s $	0.021	T2K [21, 72, 73], NO ν A [24]
$ \epsilon_{\mu\tau}^s $	0.080	
$ \epsilon_{ee}^d $	0.007	
$ \epsilon_{\mu e}^d $	0.018	
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CC NSIs

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YD, H-L. Li, J. Tang, S. Vihonen, J-H. Yu, JHEP 03(2021) 019

Q: What is the implication on the UV physics?

CC NSIs

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What particle physicists care about: UV physics that induces these interactions

QFT: NSI parameters

$$\mathcal{L} \supset -\frac{2V_{ud}}{v^2} \left\{ [1 + \epsilon_L]_{\alpha\beta} (\bar{u}\gamma^\mu P_L d) (\bar{\ell}_\alpha \gamma_\mu P_L \nu_\beta) + [\epsilon_R]_{\alpha\beta} (\bar{u}\gamma^\mu P_R d) (\bar{\ell}_\alpha \gamma_\mu P_L \nu_\beta) \right. \\ \left. + \frac{1}{2} [\epsilon_S]_{\alpha\beta} (\bar{u}d) (\bar{\ell}_\alpha P_L \nu_\beta) - \frac{1}{2} [\epsilon_P]_{\alpha\beta} (\bar{u}\gamma_5 d) (\bar{\ell}_\alpha P_L \nu_\beta) + \frac{1}{4} [\epsilon_T]_{\alpha\beta} (\bar{u}\sigma^{\mu\nu} P_L d) (\bar{\ell}_\alpha \sigma_{\mu\nu} P_L \nu_\beta) + \text{h.c.} \right.$$

CC NSIs

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Connection between the two:

$$\epsilon_{e\beta}^s = \left[\epsilon_L - \epsilon_R - \frac{g_T}{g_A} \frac{m_e}{f_T(E_\nu)} \epsilon_T \right]_{e\beta}^*, \quad (\beta \text{ decay})$$

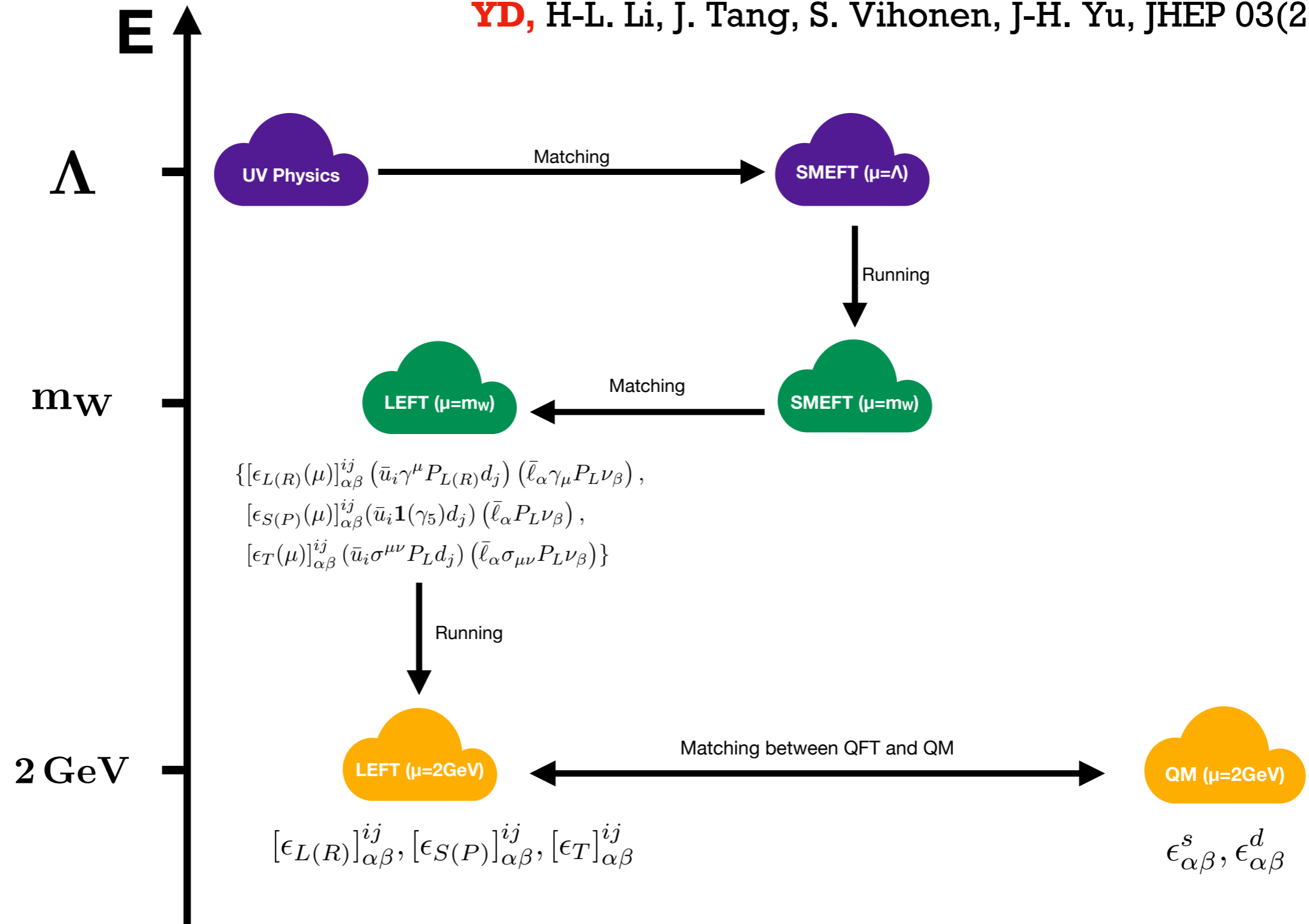
$$\epsilon_{\beta e}^d = \left[\epsilon_L + \frac{1 - 3g_A^2}{1 + 3g_A^2} \epsilon_R - \frac{m_e}{E_\nu - \Delta} \left(\frac{g_S}{1 + 3g_A^2} \epsilon_S - \frac{3g_A g_T}{1 + 3g_A^2} \epsilon_T \right) \right]_{e\beta},$$

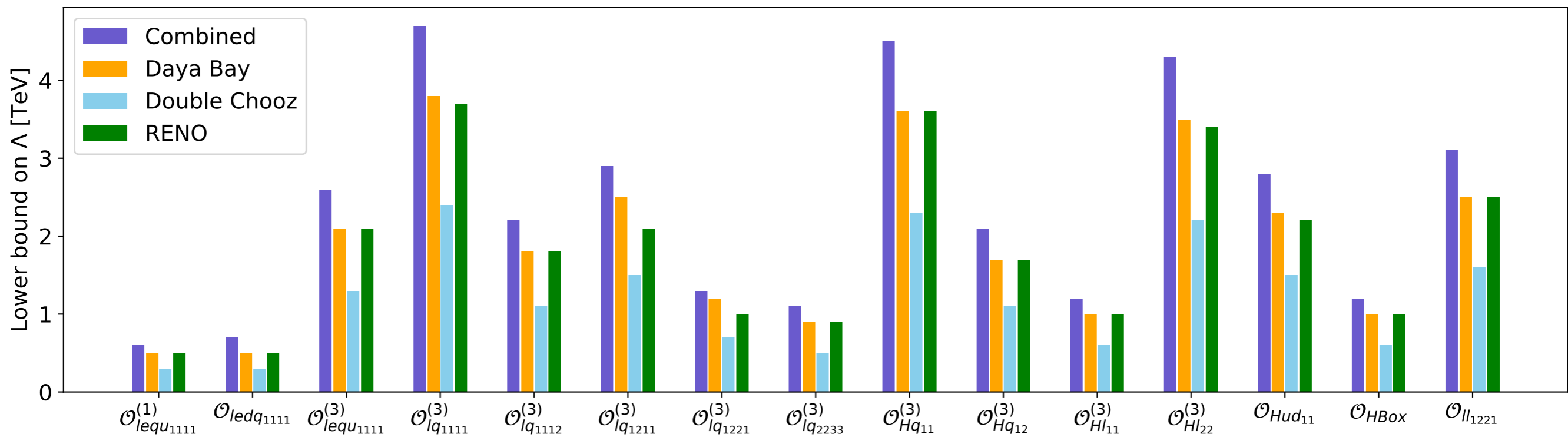
Falkowski, Gonzalez-Alonso, Tabrizi, JHEP11(2020)048

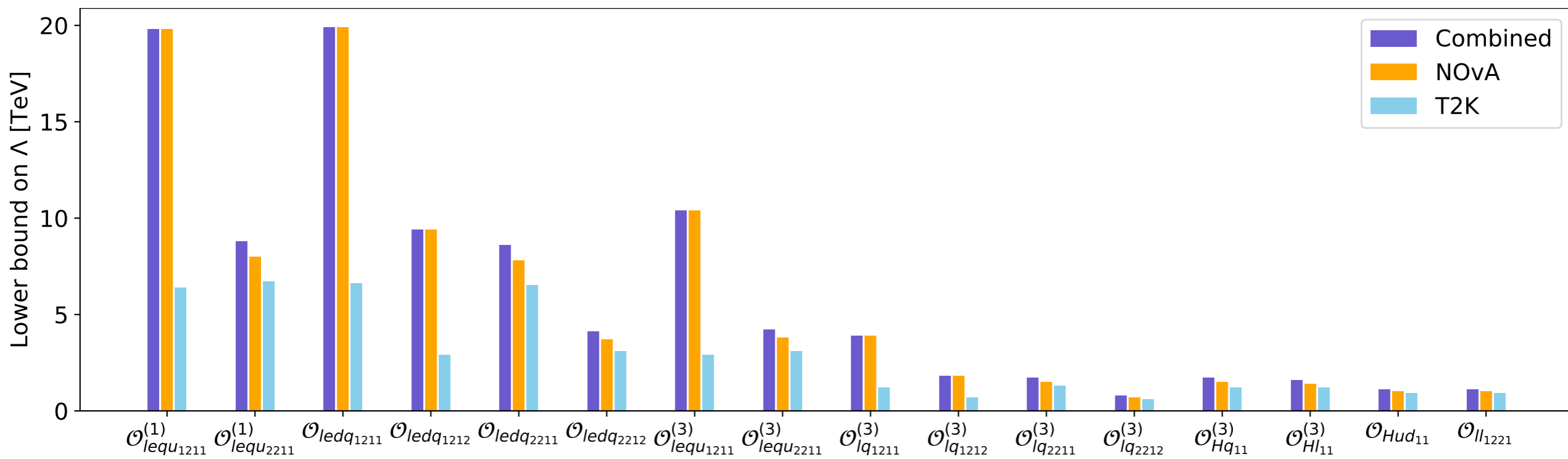
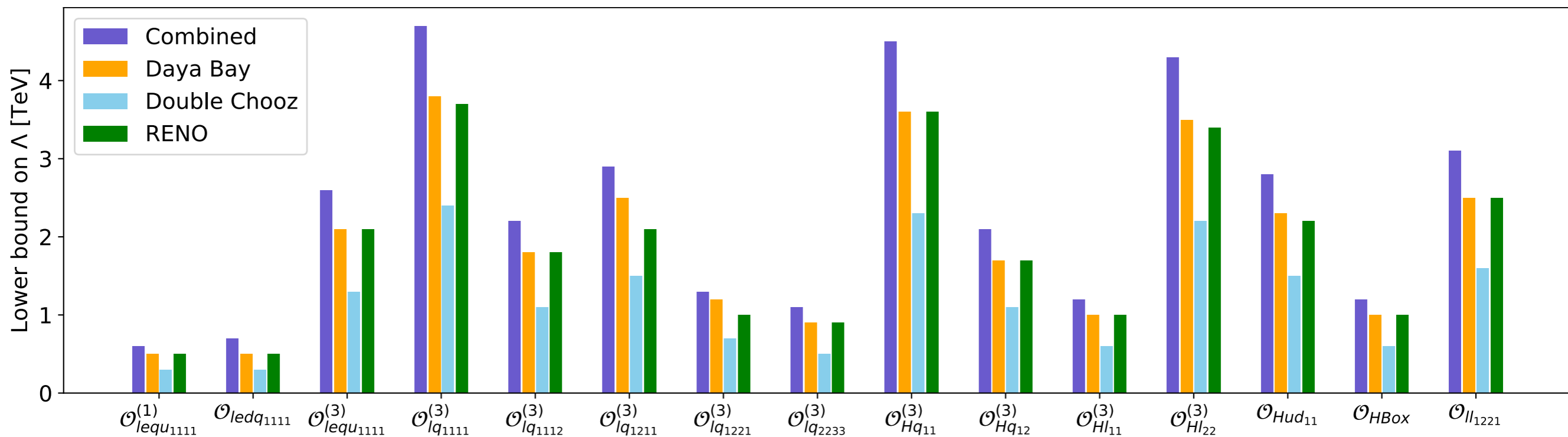
$$\epsilon_{\mu\beta}^s = \left[\epsilon_L - \epsilon_R - \frac{m_\pi^2}{m_\mu (m_u + m_d)} \epsilon_P \right]_{\mu\beta}^*, \quad (\text{pion decay})$$

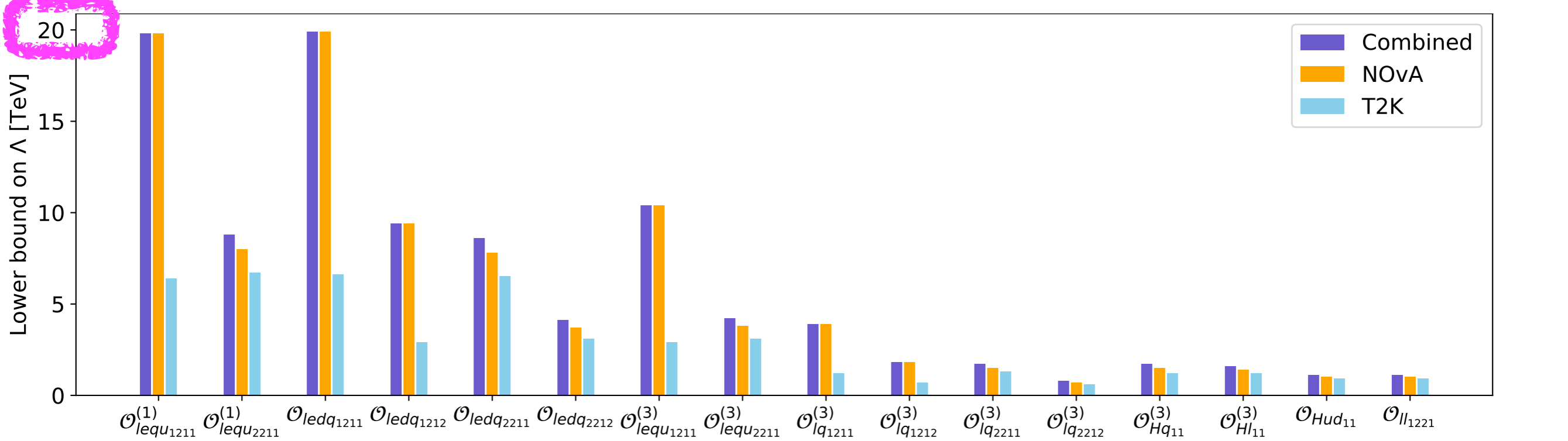
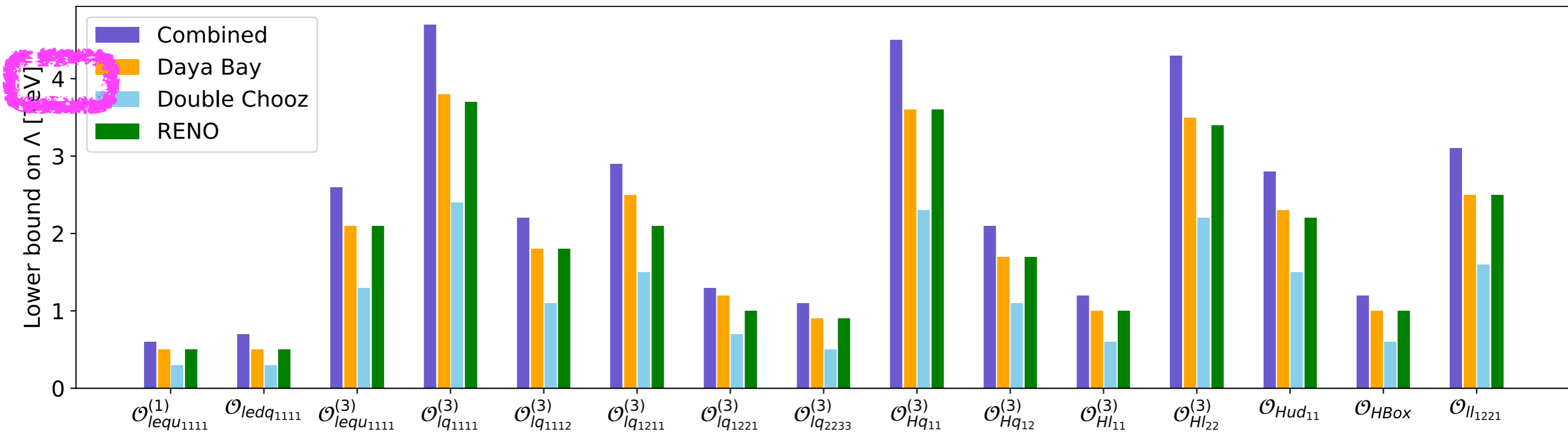
CC NSIs

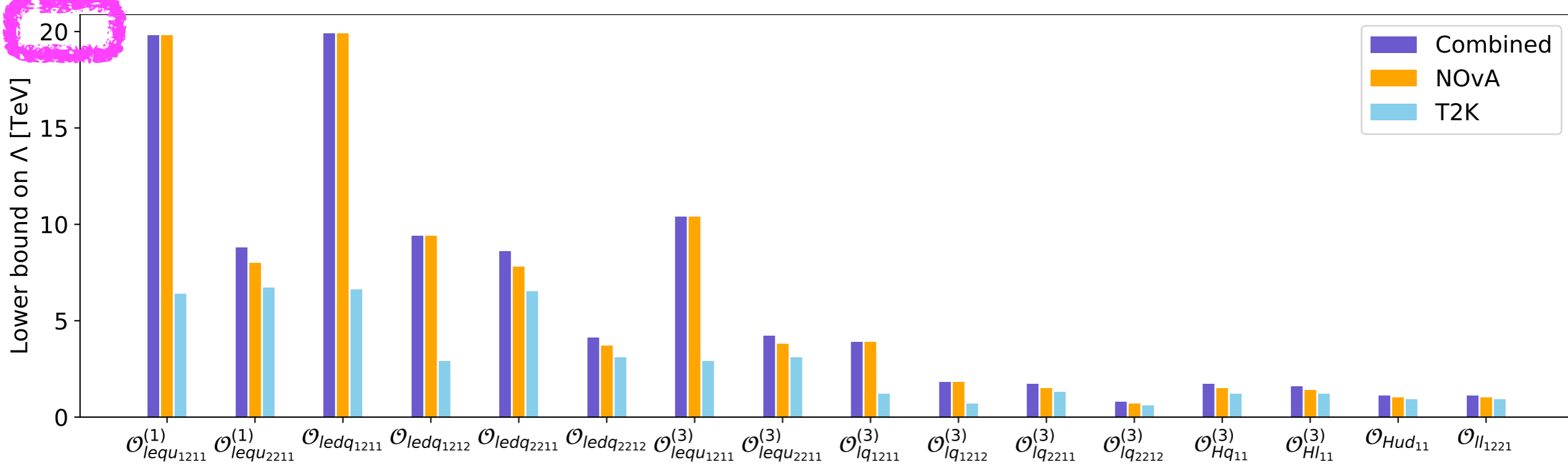
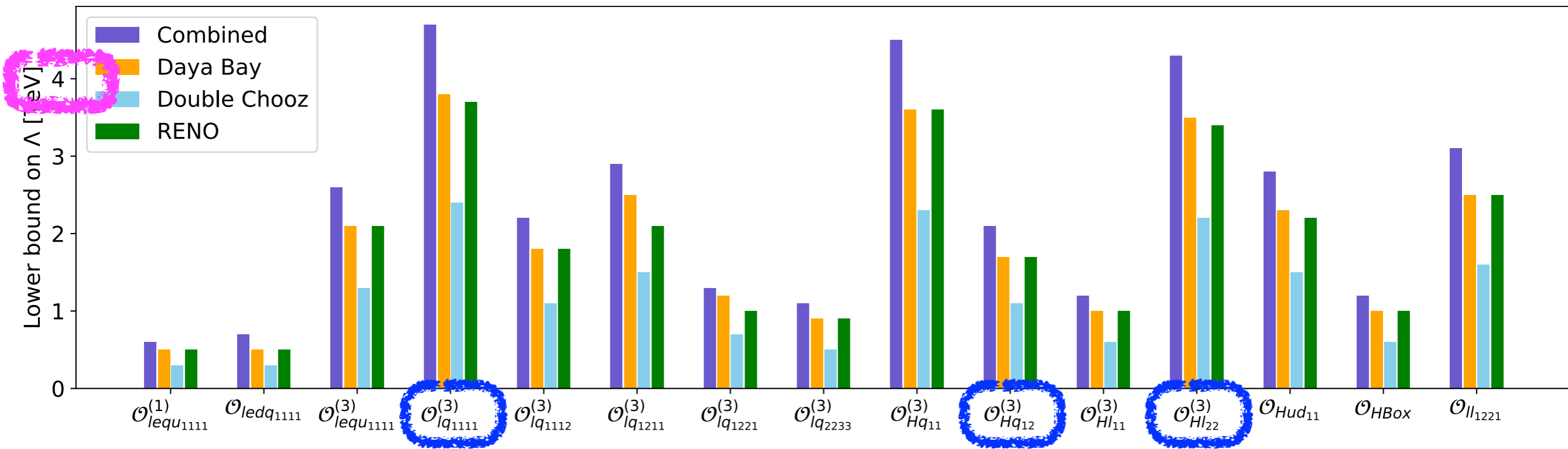
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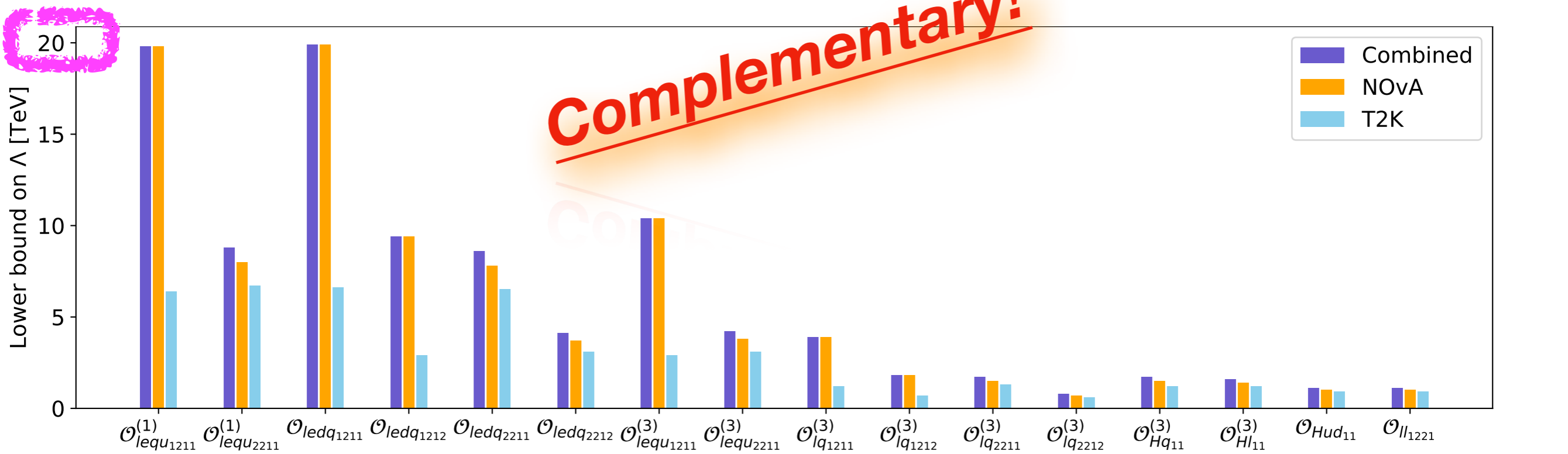
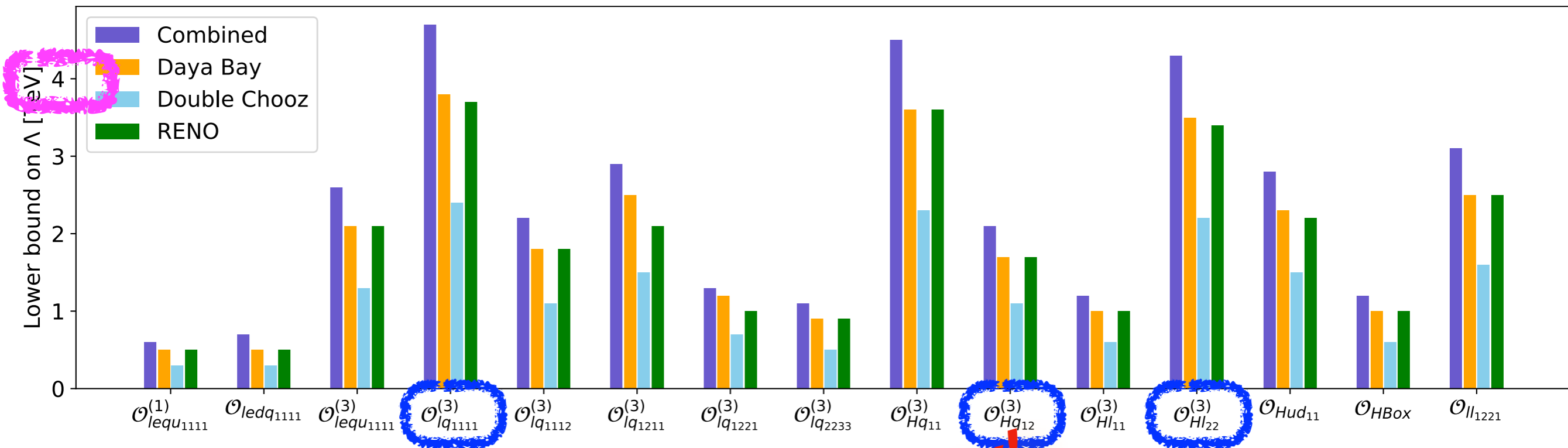




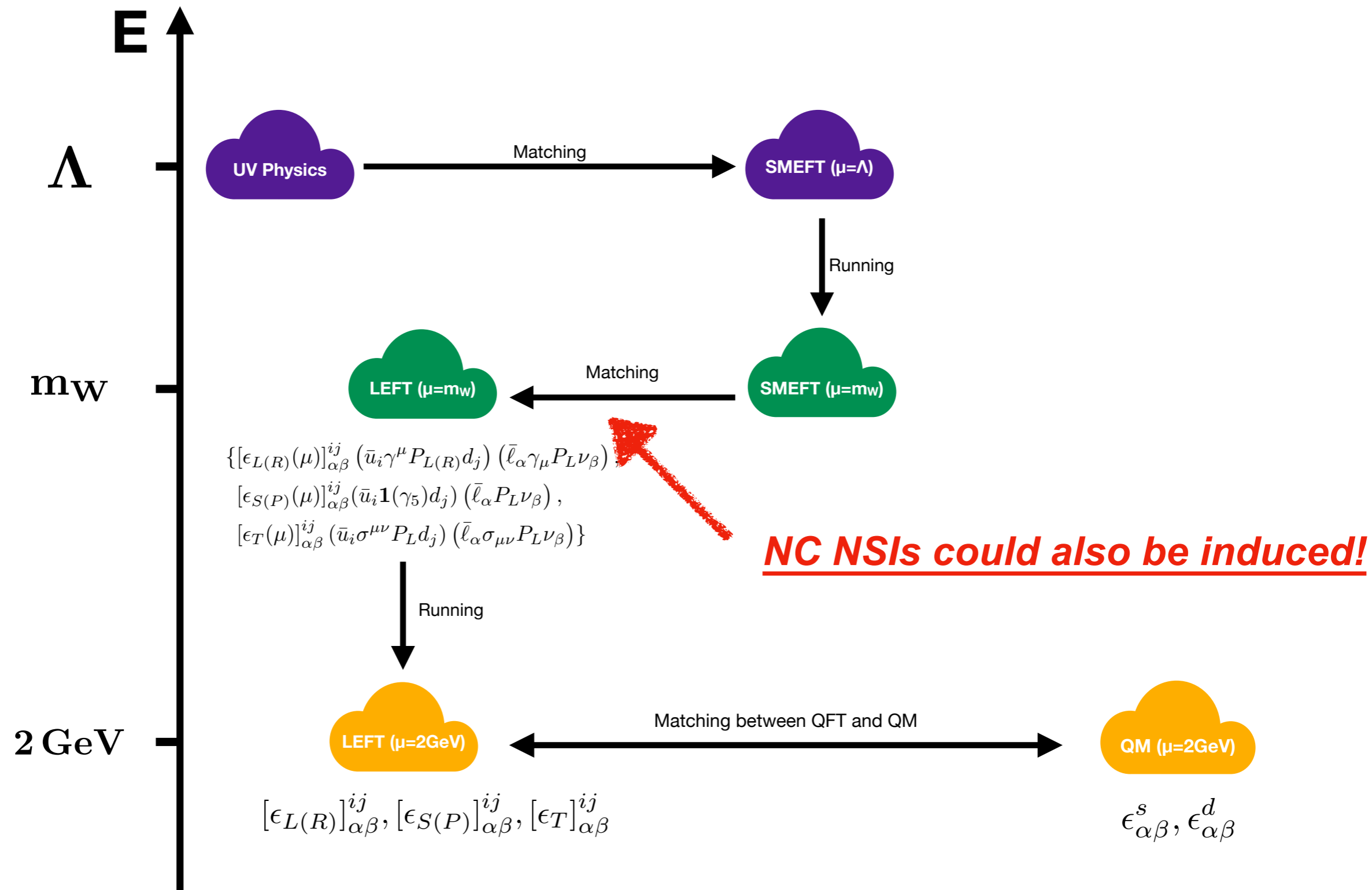








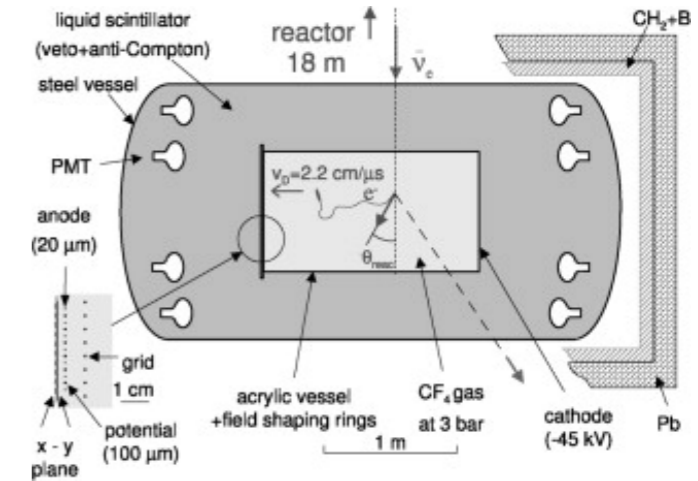
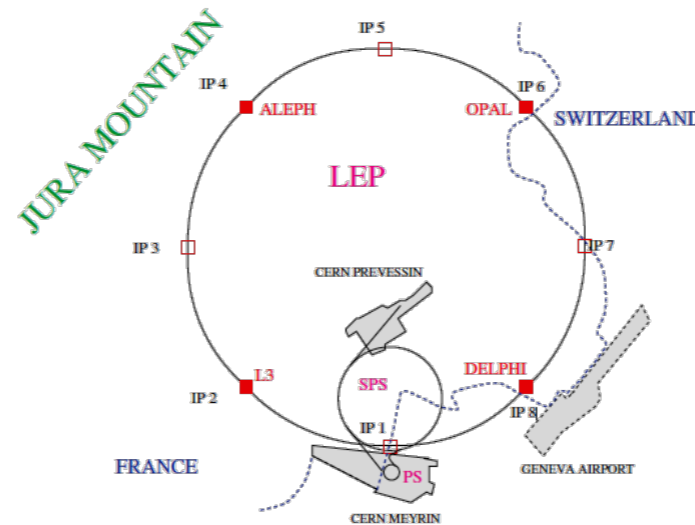
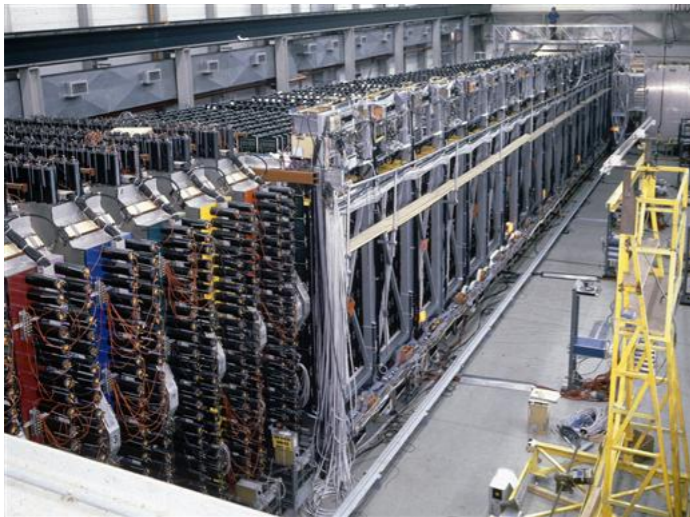
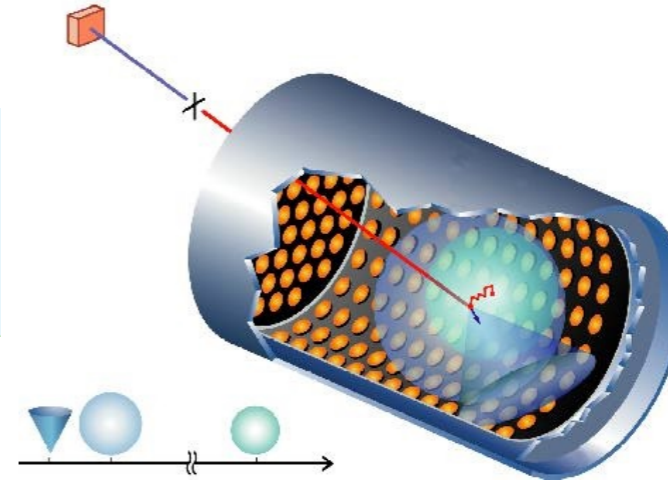
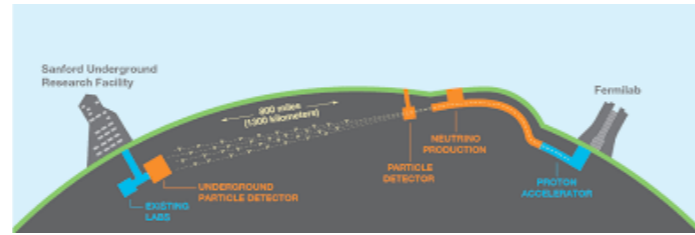
NC NSIs



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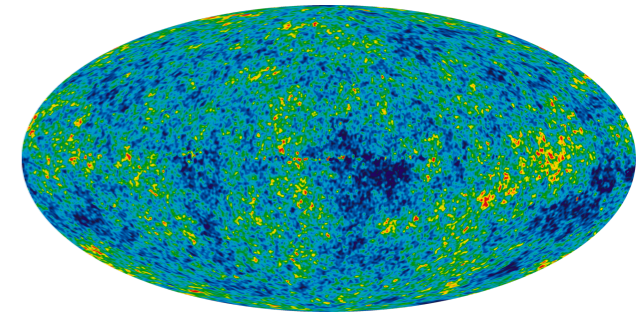
$$\mathcal{L}_{\text{NSI}}^{\text{NC}} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f,P} \epsilon_{\alpha\beta}^{f,P} (\bar{\nu}_\alpha \gamma_\mu P_L \nu_\beta) (\bar{f} \gamma^\mu P f)$$

Internal - Wiki



NC NSIs

$$\mathcal{L}_{\text{NSI}}^{\text{NC}} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f,P} \epsilon_{\alpha\beta}^{f,P} (\bar{\nu}_\alpha \gamma_\mu P_L \nu_\beta) (\bar{f} \gamma^\mu P f)$$

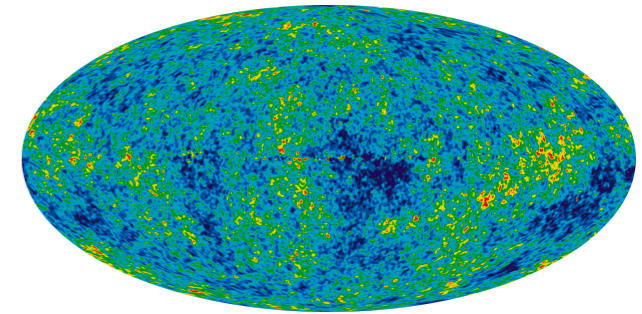


$$\rho_R = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{\frac{4}{3}} N_{\text{eff}} \right] \rho_\gamma$$

NC NSIs

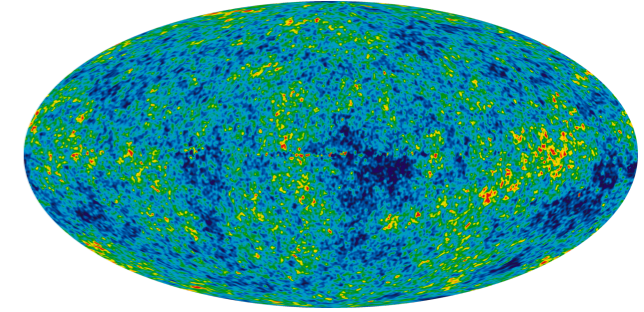
Q: How NC NSIs affect neutrino decoupling?

$$\mathcal{L}_{\text{NSI}}^{\text{NC}} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f,P} \epsilon_{\alpha\beta}^{f,P} (\bar{\nu}_\alpha \gamma_\mu P_L \nu_\beta) (\bar{f} \gamma^\mu P f)$$



$$\rho_R = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{\frac{4}{3}} N_{\text{eff}} \right] \rho_\gamma$$

Dimensions	Operators	Wilson coefficients
dimension-5	$\mathcal{O}_1^{(5)} = \frac{e}{8\pi^2} (\bar{\nu}_\beta \sigma^{\mu\nu} P_L \nu_\alpha) F_{\mu\nu}$	$C_1^{(5)}$
dimension-6	$\mathcal{O}_{1,f}^{(6)} = (\bar{\nu}_\beta \gamma_\mu P_L \nu_\alpha) (\bar{f} \gamma^\mu f)$	$C_{1,f}^{(6)}$
	$\mathcal{O}_{2,f}^{(6)} = (\bar{\nu}_\beta \gamma_\mu P_L \nu_\alpha) (\bar{f} \gamma^\mu \gamma_5 f)$	$C_{2,f}^{(6)}$
	$\mathcal{O}_3^{(6)} = (\bar{\nu}_\beta P_L \nu_\alpha) (\bar{\nu}_{\beta'} P_L \nu_{\alpha'})^*$	$C_3^{(6)}$
	$\mathcal{O}_4^{(6)} = (\bar{\nu}_\beta \gamma_\mu P_L \nu_\alpha) (\bar{\nu}_{\beta'} \gamma_\mu P_L \nu_{\alpha'})^*$	$C_4^{(6)}$
	$\mathcal{O}_5^{(6)} = (\bar{\nu}_\beta \sigma^{\mu\nu} P_L \nu_\alpha) (\bar{\nu}_{\beta'} \sigma^{\mu\nu} P_L \nu_{\alpha'})^*$	$C_5^{(6)}$
dimension-7	$\mathcal{O}_1^{(7)} = \frac{\alpha}{12\pi} (\bar{\nu}_\beta P_L \nu_\alpha) F^{\mu\nu} F_{\mu\nu}$	$C_1^{(7)}$
	$\mathcal{O}_2^{(7)} = \frac{\alpha}{8\pi} (\bar{\nu}_\beta P_L \nu_\alpha) F^{\mu\nu} \tilde{F}_{\mu\nu}$	$C_2^{(7)}$
	$\mathcal{O}_{5,f}^{(7)} = m_f (\bar{\nu}_\beta P_L \nu_\alpha) (\bar{f} f)$	$C_{5,f}^{(7)}$
	$\mathcal{O}_{6,f}^{(7)} = m_f (\bar{\nu}_\beta P_L \nu_\alpha) (\bar{f} i \gamma_5 f)$	$C_{6,f}^{(7)}$
	$\mathcal{O}_{7,f}^{(7)} = m_f (\bar{\nu}_\beta \sigma^{\mu\nu} P_L \nu_\alpha) (\bar{f} \sigma_{\mu\nu} f)$	$C_{7,f}^{(7)}$
	$\mathcal{O}_{8,f}^{(7)} = (\bar{\nu}_\beta i \overleftrightarrow{\partial}_\mu P_L \nu_\alpha) (\bar{f} \gamma^\mu f)$	$C_{8,f}^{(7)}$
	$\mathcal{O}_{9,f}^{(7)} = (\bar{\nu}_\beta i \overleftrightarrow{\partial}_\mu P_L \nu_\alpha) (\bar{f} \gamma^\mu \gamma_5 f)$	$C_{9,f}^{(7)}$
	$\mathcal{O}_{10,f}^{(7)} = \partial_\mu (\bar{\nu}_\beta \sigma^{\mu\nu} P_L \nu_\alpha) (\bar{f} \gamma_\nu f)$	$C_{10,f}^{(7)}$
	$\mathcal{O}_{11,f}^{(7)} = \partial_\mu (\bar{\nu}_\beta \sigma^{\mu\nu} P_L \nu_\alpha) (\bar{f} \gamma_\nu \gamma_5 f)$	$C_{11,f}^{(7)}$

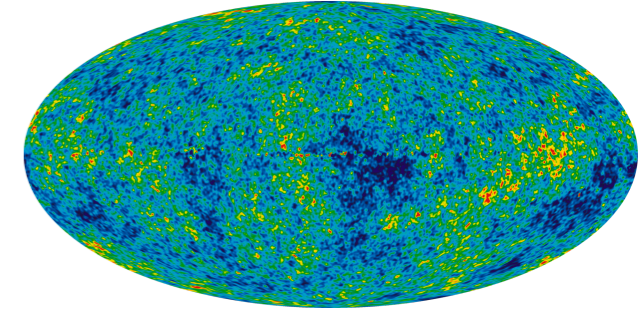


$$\rho_R = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{\frac{4}{3}} N_{\text{eff}} \right] \rho_\gamma$$

Majoron model

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	$\mathcal{O}_4^{(6)} = (\bar{\nu}_\beta \gamma_\mu P_L \nu_\alpha) (\bar{\nu}_{\beta'} \gamma_\mu P_L \nu_{\alpha'})^*$	$C_4^{(6)}$
	$\mathcal{O}_5^{(6)} = (\bar{\nu}_\beta \sigma^{\mu\nu} P_L \nu_\alpha) (\bar{\nu}_{\beta'} \sigma^{\mu\nu} P_L \nu_{\alpha'})^*$	$C_5^{(6)}$
dimension-7	$\mathcal{O}_1^{(7)} = \frac{\alpha}{12\pi} (\bar{\nu}_\beta P_L \nu_\alpha) F^{\mu\nu} F_{\mu\nu}$	$C_1^{(7)}$
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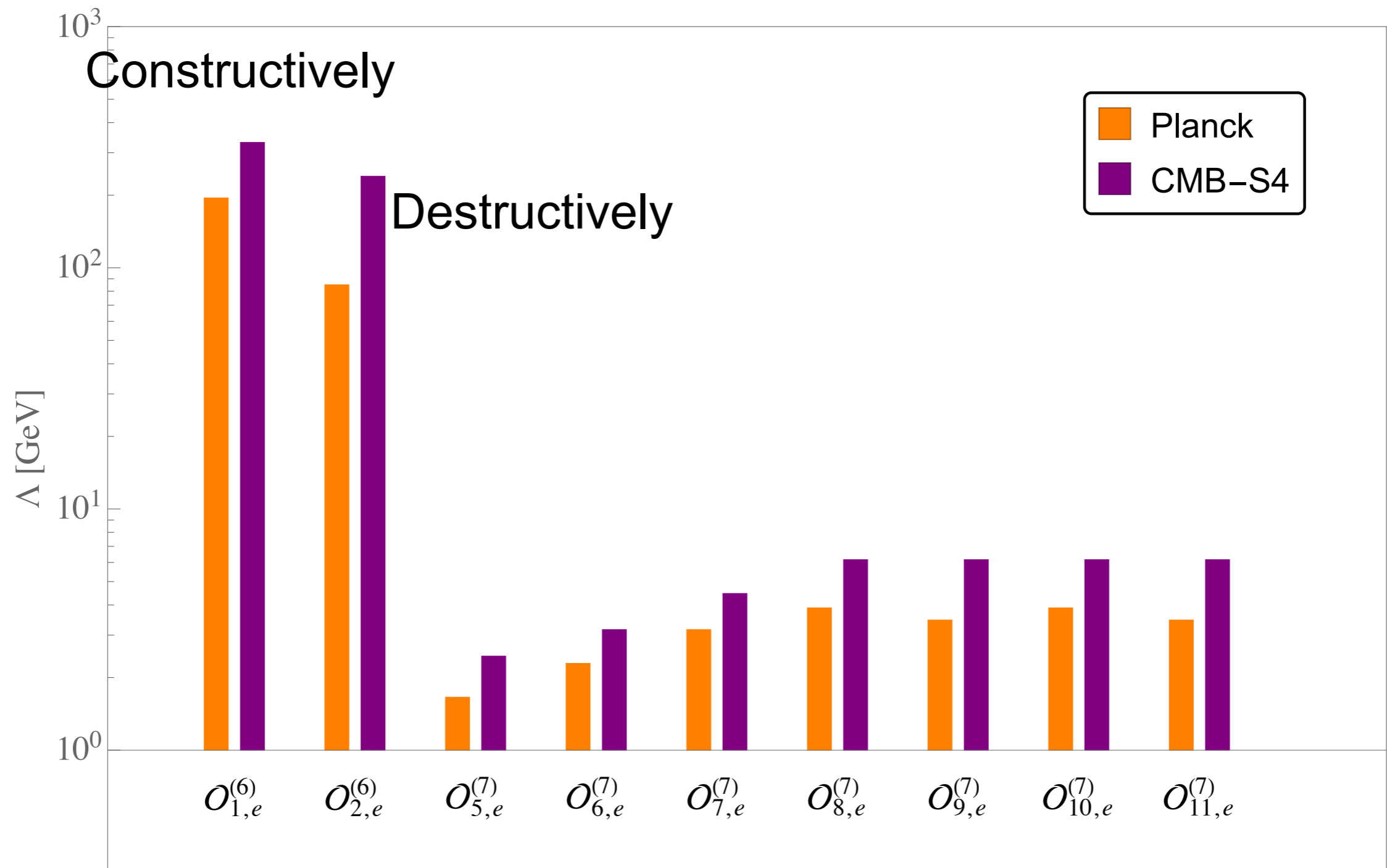
U(1)' model



$$\rho_R = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{\frac{4}{3}} N_{\text{eff}} \right] \rho_\gamma$$

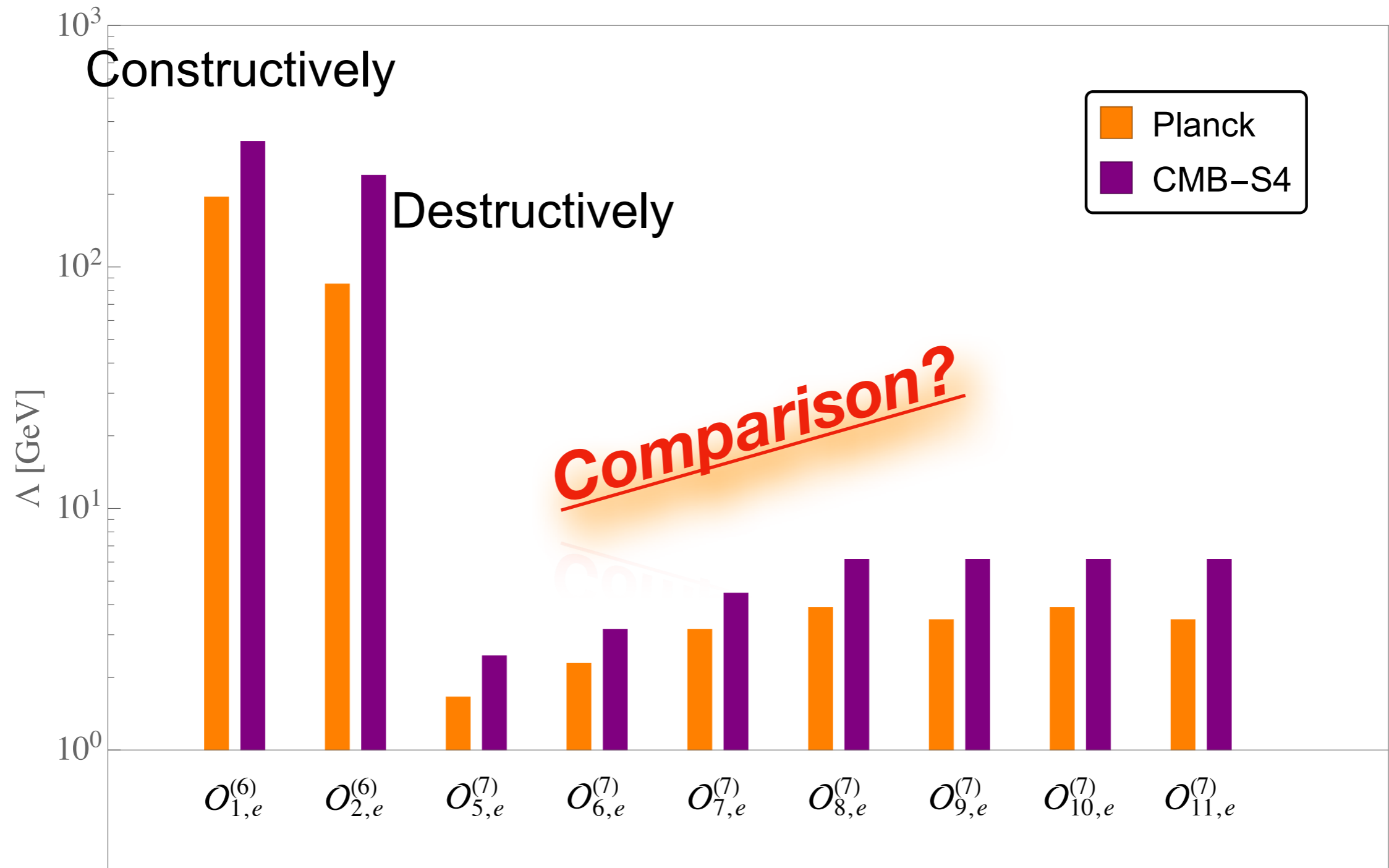
Results: NC NSIs

YD, J-H. Yu, arXiv: 2101.10475



Results: NC NSIs

YD, J-H. Yu, arXiv: 2101.10475



Results: NC NSIs comparison

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$$\mathcal{L}_{\text{NSI}}^{\text{NC}} = -2\sqrt{2}G_F \sum_{\alpha\beta, f, P} \epsilon_{\alpha\beta}^{f, P} (\bar{\nu}_\alpha \gamma_\mu P_L \nu_\beta) (\bar{f} \gamma^\mu P f)$$

ϵ 's	[103]	[97]	[82]	[83]	[84]	[85]	[90]	[98]	[35]	This work	
										Planck	CMB-S4
$\epsilon_{ee}^{e,L}$	[-0.010, 2.039]	[-1.53, 0.38]	[-0.07, 0.1]	[-0.05, 0.12]	[-0.03, 0.08]	[-0.036, 0.063]	[-0.017, 0.027] [-0.003, 0.003]	[-0.08, 0.08]	[-0.185, 0.380] [-0.130, 0.185]	[-1.6, 1.44]	[-0.61, 0.46]
$\epsilon_{e\mu}^{e,L}$	[-0.179, 0.146]	[-0.84, 0.84]	-	-	[-0.13, 0.13]	-	[-0.152, 0.152] [-0.055, 0.055]	[-0.33, 0.35]	[-0.025, 0.052] [-0.017, 0.040]	[-1.6, 1.44]	[-0.61, 0.46]
$\epsilon_{e\tau}^{e,L}$	[-0.860, 0.350]	[-0.84, 0.84]	[-0.4, 0.4]	[-0.44, 0.44]	[-0.33, 0.33]	-	[-0.152, 0.152] [-0.055, 0.055]	[-0.33, 0.35]	[-0.055, 0.023] [-0.042, 0.012]	[-1.6, 1.44]	[-0.61, 0.46]
$\epsilon_{\mu\mu}^{e,L}$	[-0.364, 1.387]	-	[-0.03, 0.03]	-	[-0.03, 0.03]	-	[-0.040, 0.04] [-0.010, 0.010]	-	[-0.290, 0.390] [-0.192, 0.240]	[-1.6, 1.44]	[-0.61, 0.46]
$\epsilon_{\mu\tau}^{e,L}$	[-0.035, 0.028]	-	[-0.1, 0.1]	-	[-0.1, 0.1]	-	-	-	[-0.015, 0.013] [-0.010, 0.010]	[-1.6, 1.44]	[-0.61, 0.46]
$\epsilon_{\tau\tau}^{e,L}$	[-0.350, 1.400]	-	[-0.5, 0.5]	-	[-0.46, 0.24]	[-0.16, 0.110] [0.41, 0.66]	[-0.040, 0.04] [-0.010, 0.010]	-	[-0.360, 0.145] [-0.120, 0.095]	[-1.6, 1.44]	[-0.61, 0.46]
$\epsilon_{ee}^{e,R}$	[-0.010, 2.039]	[-0.07, 0.08]	[-1, 0.5]	[-0.04, 0.14]	[0.004, 0.151]	[-0.27, 0.59]	[-0.33, 0.25] [-0.07, 0.07]	[-0.04, 0.06]	[-0.185, 0.380] [-0.130, 0.185]	[-1.6, 1.44]	[-0.39, 0.31]
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Results: NC NSIs comparison

YD, J-H. Yu, arXiv: 2101.10475

$$\mathcal{L}_{\text{NSI}}^{\text{NC}} = -2\sqrt{2}G_F \sum_{\alpha\beta, f, P} \epsilon_{\alpha\beta}^{f, P} (\bar{\nu}_\alpha \gamma_\mu P_L \nu_\beta) (\bar{f} \gamma^\mu P f)$$

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$\epsilon_{ee}^{e,R}$	[-0.010, 2.039]	[-0.07, 0.08]	[-1, 0.5]	[-0.04, 0.14]	[0.004, 0.151]	[-0.27, 0.59]	[-0.33, 0.25] [-0.07, 0.07]	[-0.04, 0.06]	[-0.185, 0.380] [-0.130, 0.185]	[-1.6, 1.44]	[-0.39, 0.31]
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Complementary!

Summary

We investigate charge- and neutral-current neutrino NSIs in the EFT framework.

- ❖ For CC NSIs, we find reactor (Daya Bay, Double Chooze, RENO) and long baseline (T2K, NOvA) neutrino experiments are complementary, the latter are sensitive to new physics already at the $\sim 20\text{TeV}$ scale.
- ❖ For NC NSIs up to dim-7, constraints from precision measurements of N_{eff} (Planck, CMB-S4) are complementary to other type of neutrino experiments (COHERENT, collider, solar and reactor neutrino experiments, DUNE etc).

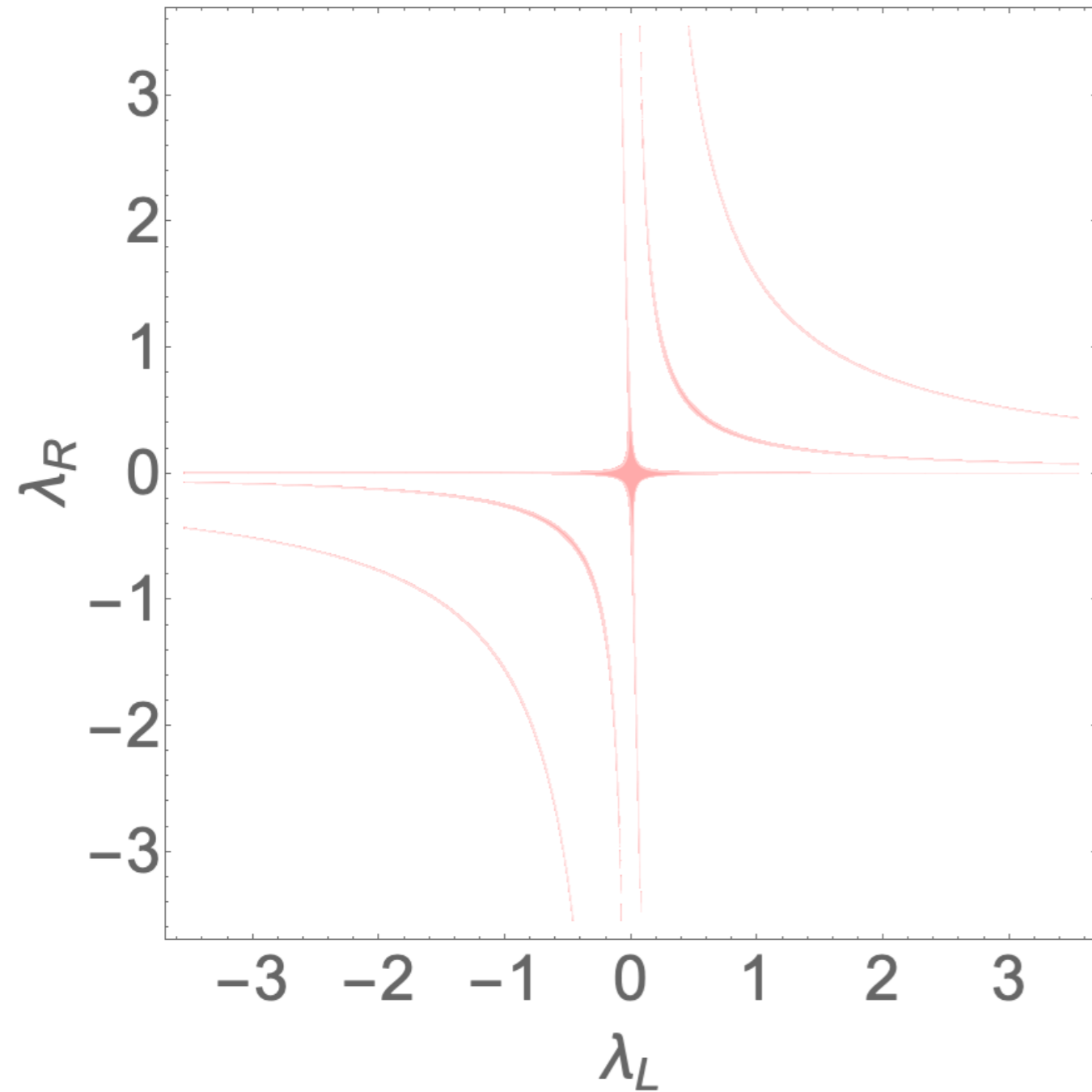
Back up

Simplified leptoquark model

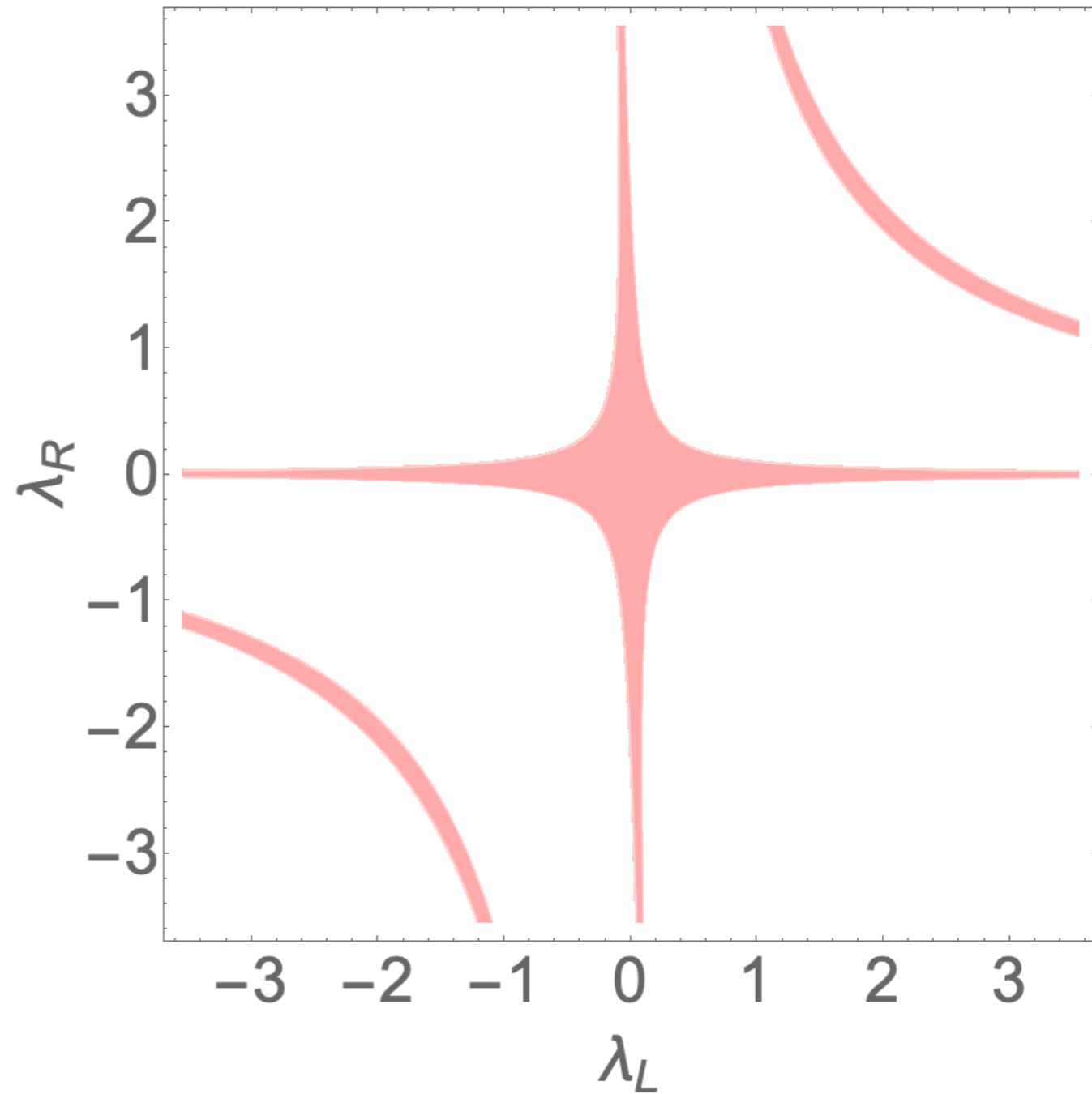
YD, H-L. Li, J. Tang, S. Vihonen, J-H. Yu, JHEP 03(2021) 019

$$\mathcal{L}_{\text{LQ}} = |D_\mu S|^2 - M_1^2 |S|^2 - \lambda_{H1} |H|^2 |S|^2 - \frac{c}{2} |S|^4 \\ + ((\lambda^L)_{i\alpha} \bar{q}_i^c \ell_\alpha + (\lambda^R)_{i\alpha} \bar{u}_i^c e_\alpha) S_1 + \text{h.c.}$$

$|\epsilon_{\mu e}^S|$ from π^\pm decay, $M=1$ TeV

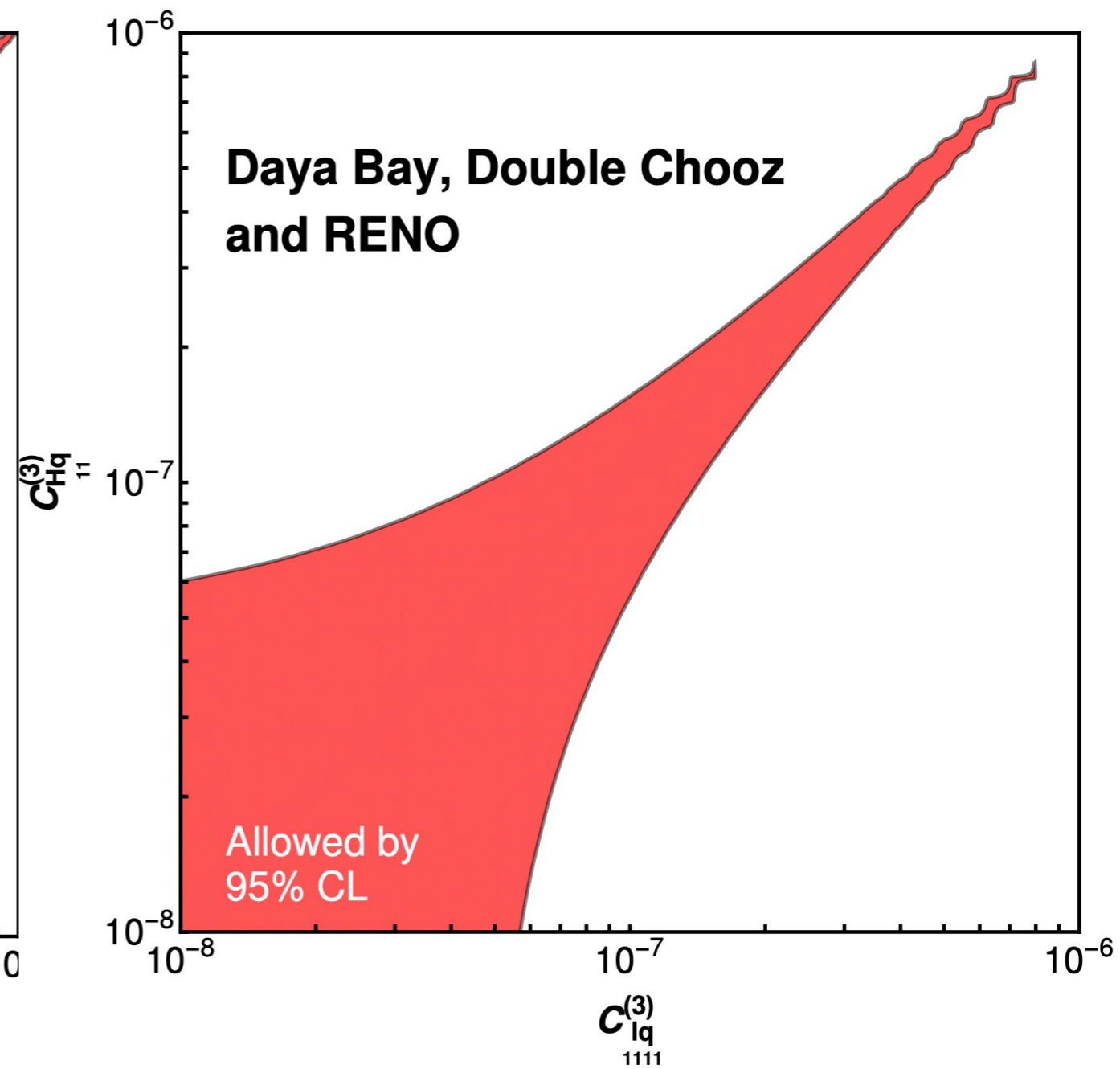
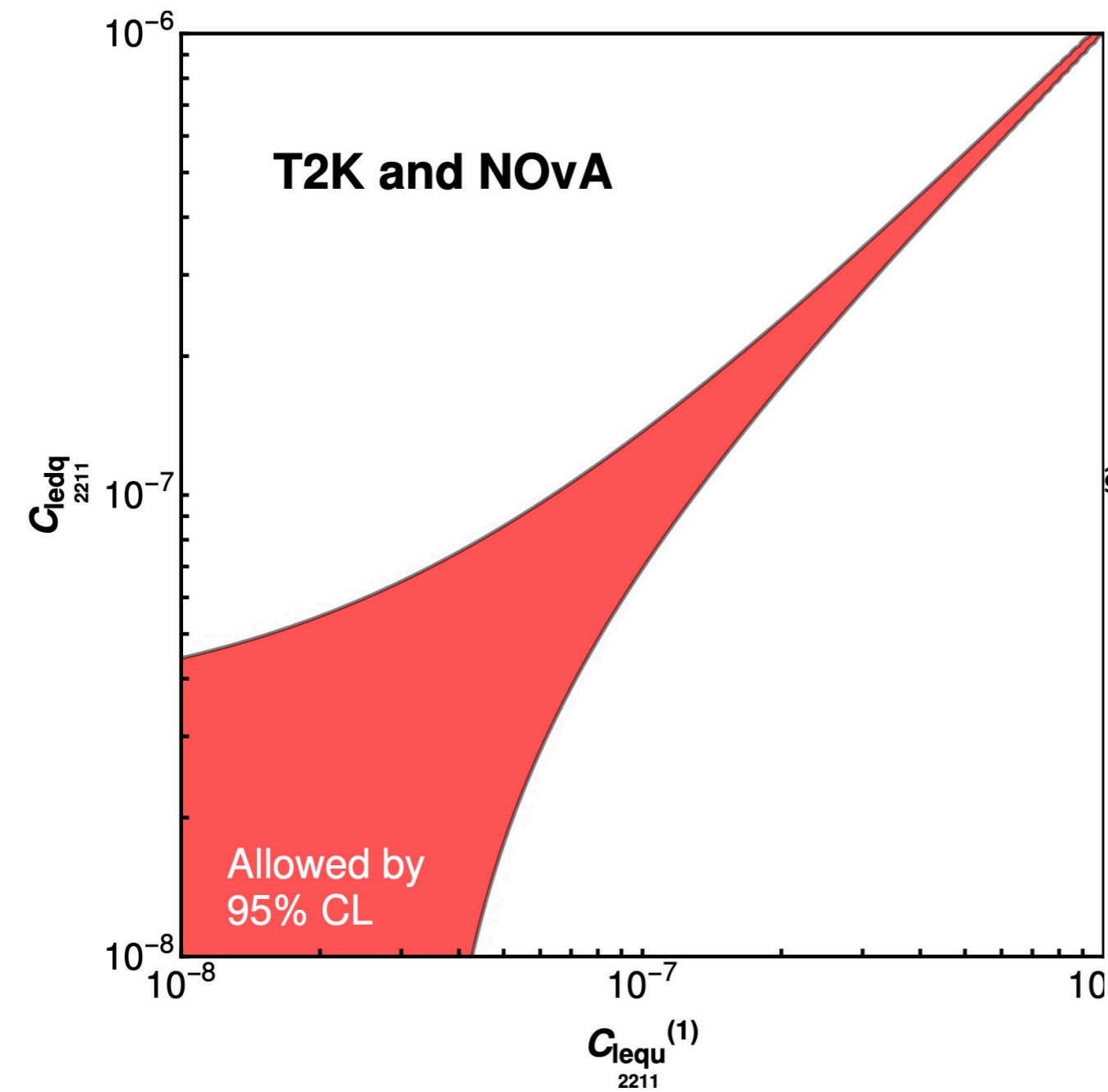


$|\epsilon_{\mu e}^S|$ from π^\pm decay, $M=5$ TeV



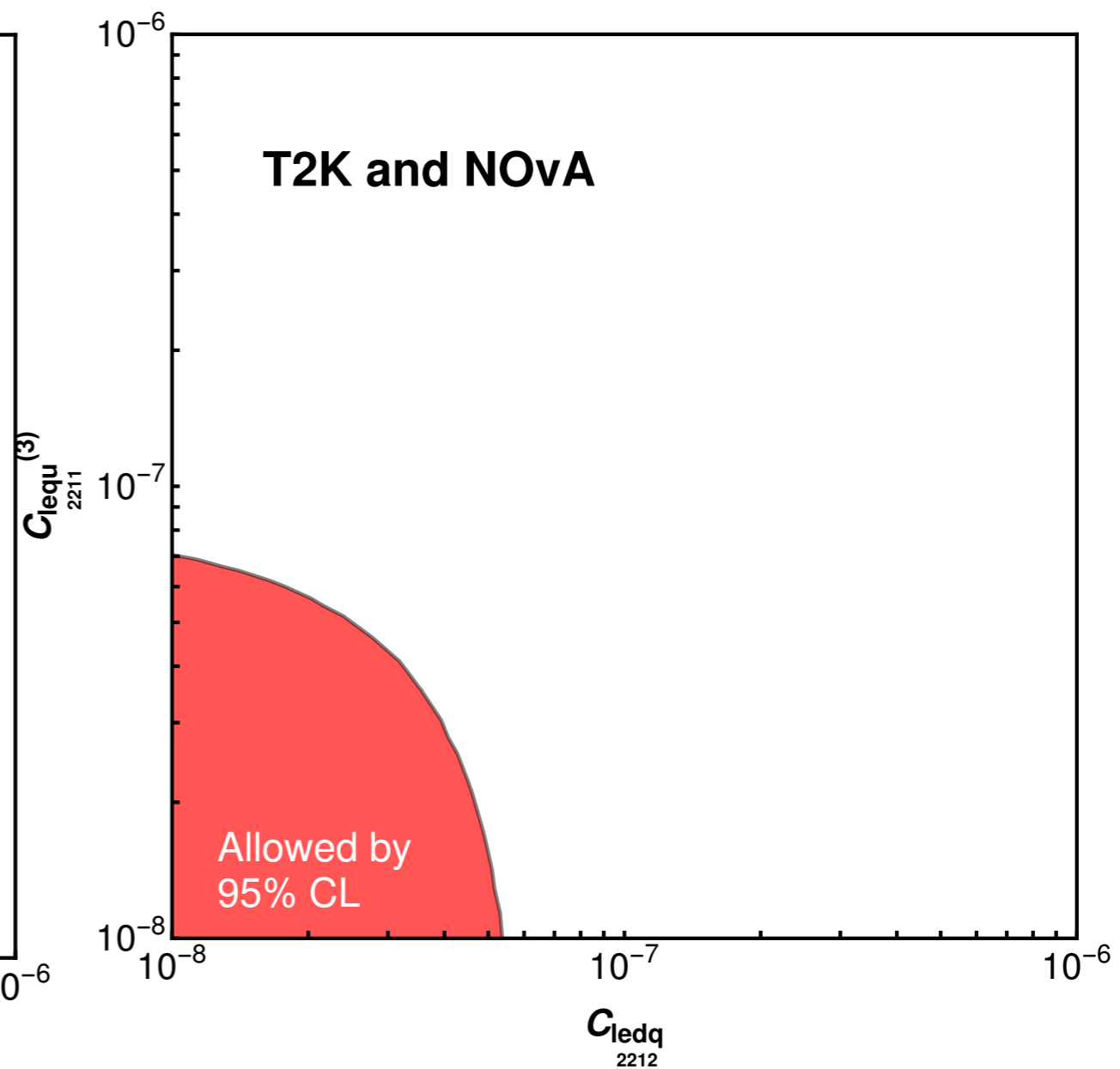
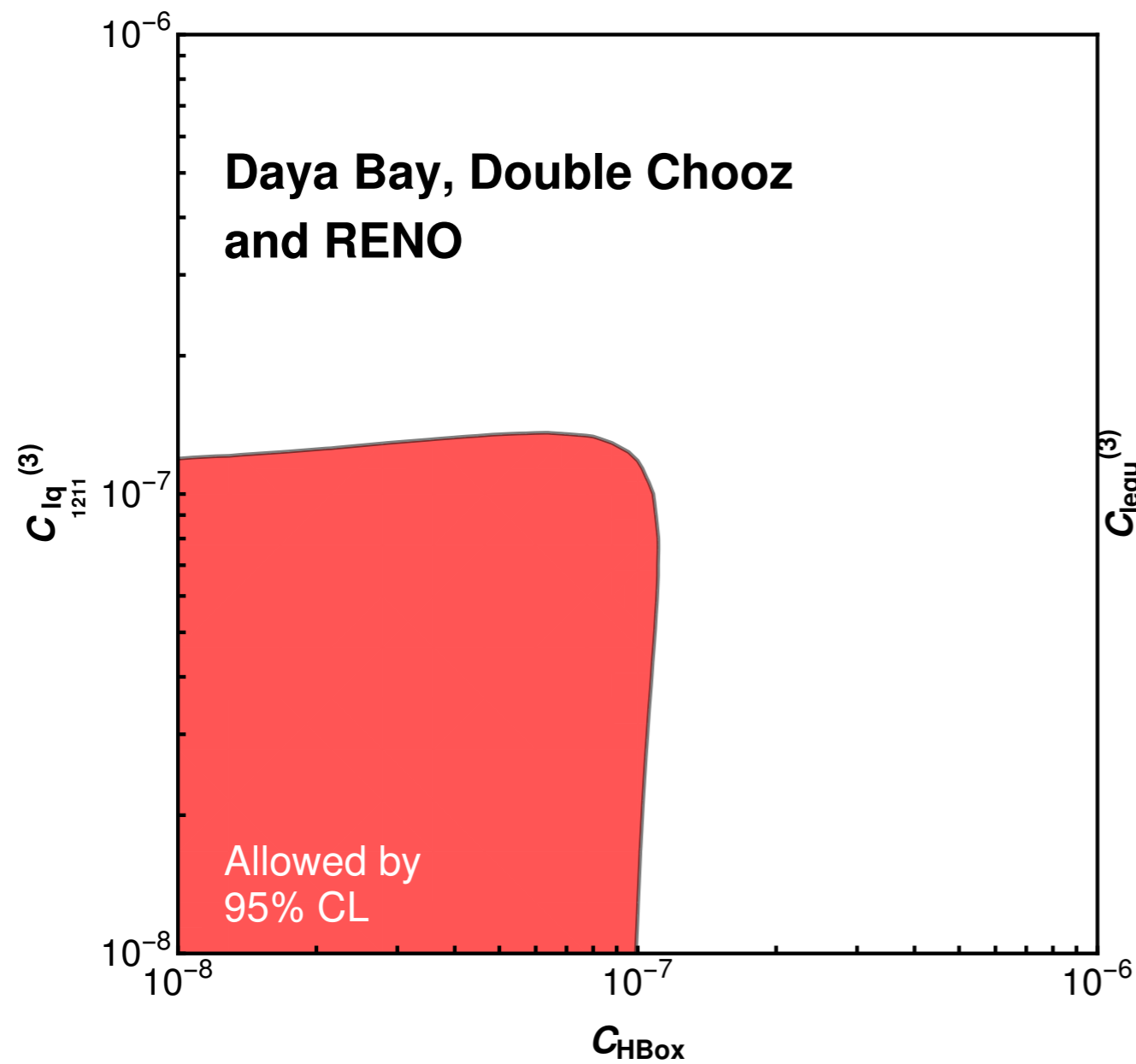
Multiple operators

YD, H-L. Li, J. Tang, S. Vihonen, J-H. Yu, JHEP 03(2021) 019



Multiple operators

YD, H-L. Li, J. Tang, S. Vihonen, J-H. Yu, JHEP 03(2021) 019



Reactor vs LBL neutrino experiments

Reactor

$$\epsilon_{e\beta}^s = \left[\epsilon_L - \epsilon_R - \frac{g_T}{g_A} \frac{m_e}{f_T(E_\nu)} \epsilon_T \right]_{e\beta}^*, \quad (\beta \text{ decay}) \quad (2.4)$$

$$\epsilon_{\beta e}^d = \left[\epsilon_L + \frac{1 - 3g_A^2}{1 + 3g_A^2} \epsilon_R - \frac{m_e}{E_\nu - \Delta} \left(\frac{g_S}{1 + 3g_A^2} \epsilon_S - \frac{3g_A g_T}{1 + 3g_A^2} \epsilon_T \right) \right]_{e\beta}, \quad (\text{inverse } \beta \text{ decay}) \quad (2.5)$$

LBL

$$\epsilon_{\mu\beta}^s = \left[\epsilon_L - \epsilon_R - \frac{m_\pi^2}{m_\mu (m_u + m_d)} \epsilon_P \right]_{\mu\beta}^*, \quad (\text{pion decay}) \quad (2.6)$$

NC NSIs: Comparison

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										Planck	CMB-S4
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Table 4. Summary of constraints on dimension-6 neutrino-electron NC NSIs from previous studies and this work. Constraints from a global fitting of all kinds of neutrino oscillation data plus the COHERENT result are obtained in Ref. [103], the TEXONO collaboration in Ref. [97], the LEP, LSND and CHARM-II experiments in Ref. [82], a global analysis of $\nu_e e$ and $\bar{\nu}_e e$ scattering data from LSND, Irvine, Rovno and MUNU experiments in Ref. [83], OPAL, ALEPH, L3, DELPHI, LSND, CHARM-II, Irvine, Rovno and MUNU experiments in Ref. [84], solar and reactor neutrino experiments in Ref. [85], low-energy solar neutrinos at source and detector from the Borexino experiment in Ref. [90], a global analysis of short baseline νe and $\bar{\nu} e$ data from LSND, LAMPF, Irvine, Rovno, MUNU, TEXONO and KRANOYARSK in Ref. [98], and DUNE in Ref. [35].

NC NSIs: Interference

$$\rho_{\nu\text{-total}}^{\text{interf.}}(\mathcal{O}_{1,e}^{(6)}) \simeq + \frac{256\sqrt{2}C_{1,e}^{(6)}G_F\sin^2\theta_W T_\gamma^9}{\pi^5\Lambda^2},$$

$$\rho_{\nu\text{-total}}^{\text{interf.}}(\mathcal{O}_{2,e}^{(6)}) \simeq - \frac{40\sqrt{2}C_{2,e}^{(6)}G_F T_\gamma^5 T_{\nu_e}^4}{\pi^5\Lambda^2} \times (1 + 4\sin^2\theta_W),$$

NC NSIs: Neff numbers

With the complete dictionary presented in section 4, one can readily solve the Boltzmann equations for T_γ and T_{ν_α} , and thus obtain corrections to N_{eff} . In what follows, we define these corrections as

$$\Delta N_{\text{eff}} = N_{\text{eff}}^{\text{SM+EFT}} - N_{\text{eff}}^{\text{SM}}, \quad (5.1)$$

where $N_{\text{eff}}^{\text{SM+EFT}}$ is the theoretical prediction of N_{eff} with the inclusion of the NC NSI operators, and $N_{\text{eff}}^{\text{SM}} = 3.044$ [123, 132] that from the pure SM. For Planck, we use the current result $N_{\text{eff}} = 2.99^{+0.34}_{-0.33}$ [114] at the 95% CL to obtain the constraints, and $\Delta N_{\text{eff}} < 0.06$ at 95% CL for CMB-S4 [117, 143, 144, 146].