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Search for Higgs boson decays to beyond-the-Standard-Model light bosons in four-lepton final states with the ATLAS detector at the LHC

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1 Motivation-Theoritical context-

**2** Overview

**3** HM analysis selection

**4** HM analysis backgrounds estimation

**5** Results

**1** Motivation-Theoritical context-

**2** Overview

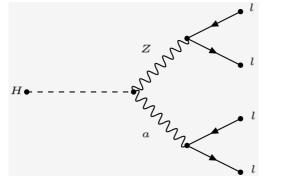
**3** HM analysis selection

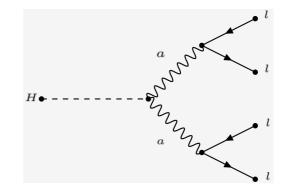
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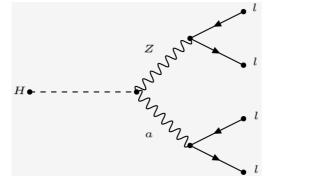
Two BSM benchmark models considered:

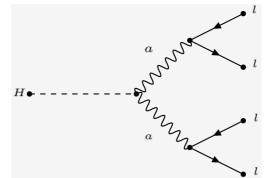
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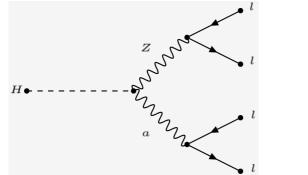
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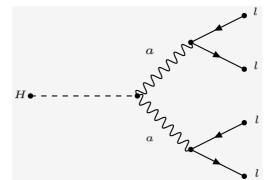




- $\rightarrow\,$  It is in this model where the prediction goes for the Higgs boson decays to 1 or 2 pseudoscalar a.
- $\rightarrow$  The decays of  $a \rightarrow 2I$  are determined by the Yukawa couplings of a to fermions.

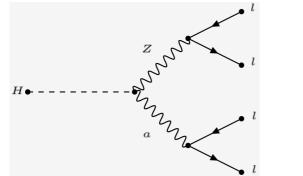
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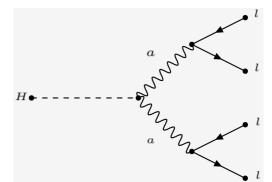




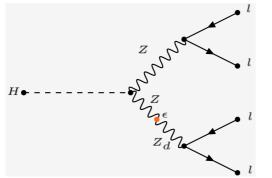
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- HAHM model: Hidden Abelian Higgs Model

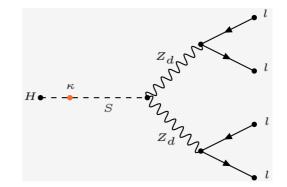
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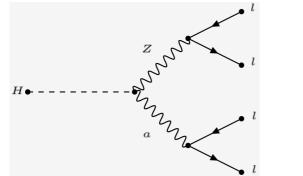
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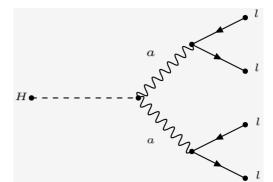




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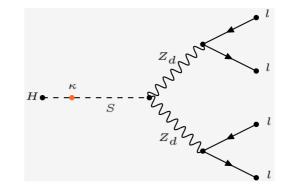
• 2HDM+S model: 2 Higgs doublet model with an additional singlet field





- $\rightarrow\,$  It is in this model where the prediction goes for the Higgs boson decays to 1 or 2 pseudoscalar a.
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- HAHM model: Hidden Abelian Higgs Model

 $H \bullet \cdots \bullet \sum_{\substack{z \\ z_{d}}}^{z} \int_{z_{d}}^{z} \int_{z_{d}}^{z}$ 



- $\rightarrow$  Introduce an additional U(1) dark gauge symmetry mediated by a dark gauge boson  $Z_d$ .
- $\rightarrow$  The Z<sub>d</sub> boson interacts with a SM gauge particle and the strength of this coupling is defined by the Kinetic mixing parameter  $\epsilon$ .
- $\rightarrow$  When the  $U(1)_d$  is broken by a dark Higgs boson, the SM Higgs boson is then mixing with a dark Higgs boson and their coupling is controled by the strenght parameter  $\kappa$ .

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- The results using 2015-6 data  $(36 fb^{-1})$  are published (using release 20) [arXiv:1802.03388]
- These results conver three channels with 4e,  $2e2\mu$  and  $4\mu$  in the final state:

High-Mass (HM):  $H \rightarrow Z_d Z_d(aa) \rightarrow 4/$ , 15 GeV  $< m_{Z_d}(m_a) < 60$  GeV.

Low Mass(LM):  $H \rightarrow Z_d Z_d(aa) \rightarrow 4\mu$ , 1 GeV  $< m_{Z_d}(m_a) < 15$  GeV.

 $\mid$  ZZ $_d$ :  $H \rightarrow$  ZZ $_d \rightarrow$  4/ , 15 GeV  $< m_{Z_d} <$  55 GeV.

# **Today**, we are going to focus on the **HM** channel.

1 Motivation-Theoritical context-

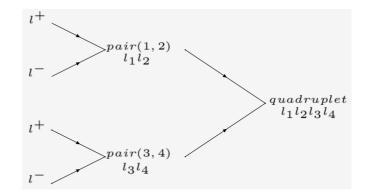
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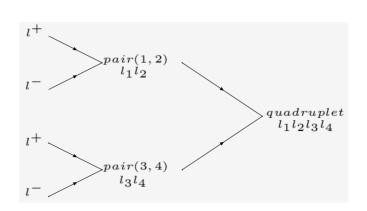
- Quadruplet formation and selection
  - in each event, a quadrulpet is formed from two lepton pairs each with same flavour opposite sign leptons: "1,2" and "3,4"
  - each lepton should fire at least 1 trigger.
  - Three leading-*pt* leptons must have: *pt* > 20, 15 and 10 GeV.
  - $\Delta R(I, I') > 0.10(0.20)$  for same-flavour (different-flavour) leptons in the quadruplet



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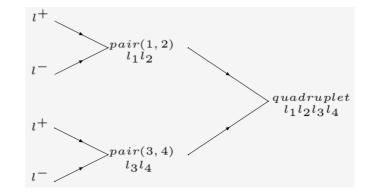
# • Quadruplet ranking



• The selected quadruplet should have the smallest difference in mass between lepton pairs:  $\Delta m_{ll} = |m_{12} - m_{34}|$ 

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# Event selection

- Higgs boson mass window:115 GeV  $< m_{4l} < 130$  GeV
- Z veto: 10 GeV  $< m_{12,34} <$  64 GeV and 5 GeV  $< m_{14,32} <$  75 GeV
- Quarkonia veto: event is rejected if either (or both) condition are fulfilled

 $(m_{J/\Psi} - 0.25 \text{ GeV}) < m_{12,34,14,32} < (m_{\Psi(2S)} + 0.30 \text{ GeV}) \text{ or } (m_{\Upsilon(1S)} - 0.70 \text{ GeV}) < m_{12,34,14,32} < (m_{\Upsilon(3S)} + 0.75 \text{ GeV})$ 

• Medium Signal Region (SR):  $m_{34}/m_{12} > 0.85$ 

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The background processes considered in this analysis are as followed:

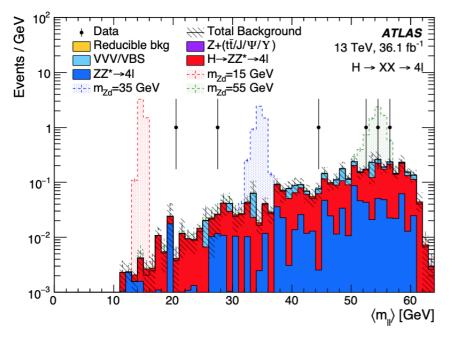
# **Dominant backgrounds** • $H \rightarrow ZZ^* \rightarrow 4/$ represents 63% of the total • non-resonant $ZZ^* \rightarrow 4/$ represents 19% of the total

Sub-Dominant backgrounds
WZ, VVV/VBS processes
tt

tt

tt

- All backgrounds estimates for this search rely basically on using MC simulations.
- Dominant backgrounds are cross-checked using validation regions.
- The data-driven ABCD method is used to estimate the reducible backgrounds.

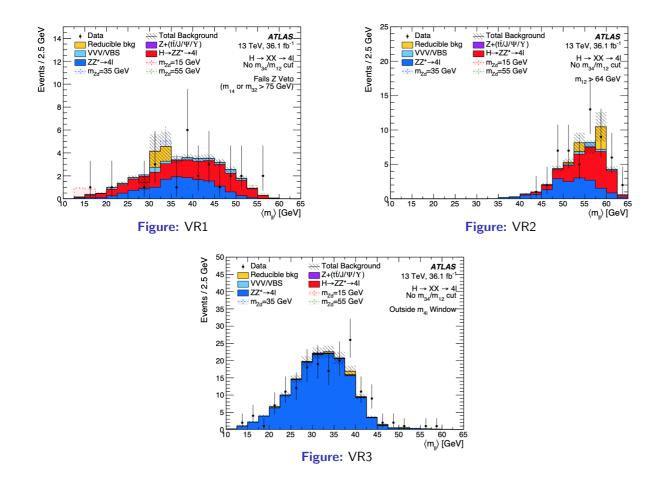


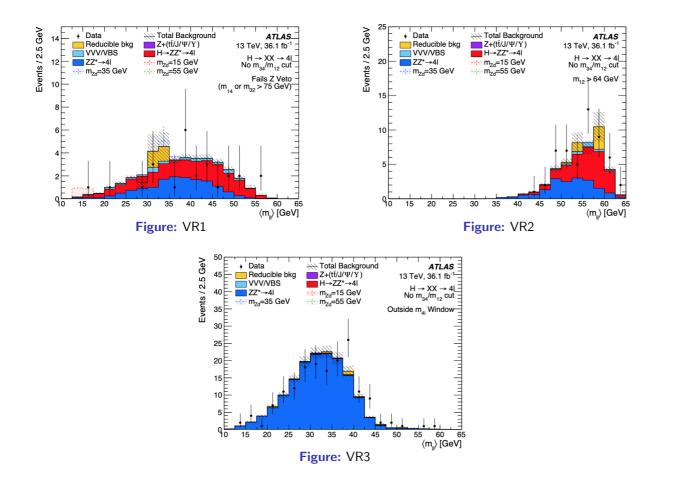
**Figure:**  $\langle m_{ll} \rangle$  distribution in Signal region with all channels combined in 2015-6 data.

The signal distributions correspond to the expected yield normalized with:  $\sigma(pp \rightarrow H \rightarrow Z_d Z_d \rightarrow 4I) = \frac{1}{10}\sigma_{SM}(pp \rightarrow H \rightarrow ZZ^* \rightarrow 4I)$ 

Process	Yield
$ZZ^* \rightarrow 4\ell$	$0.8 \pm 0.1$
$H \to ZZ^* \to 4\ell$	$2.6 \pm 0.3$
VVV/VBS	$0.51 \pm 0.18$
$Z + (t\bar{t}/J/\Psi) \rightarrow 4\ell$	$0.004 \pm 0.004$
Other Reducible Background	Negligible
Total	$3.9 \pm 0.3$
Data	6

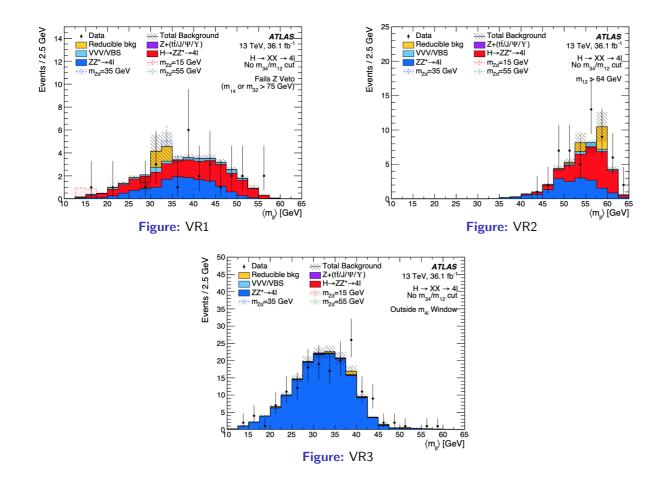
A total of 6 events are observed, with a total predicted background of 3.9  $\pm$  0.3 events.





• VR1

 $\begin{array}{l|l} & {\rm SR \ cut:} \ m_{12,34} < 64 \ {\rm GeV} \ {\rm and} \\ & m_{14,32} < 75 \ {\rm GeV}. \\ & | \ {\rm VR1 \ cut:} \ {\rm No} \ m_{34}/m_{12} \ {\rm cut} \\ & m_{14} \ {\rm or} \ m_{32} > 75 \ {\rm GeV}. \end{array}$ 

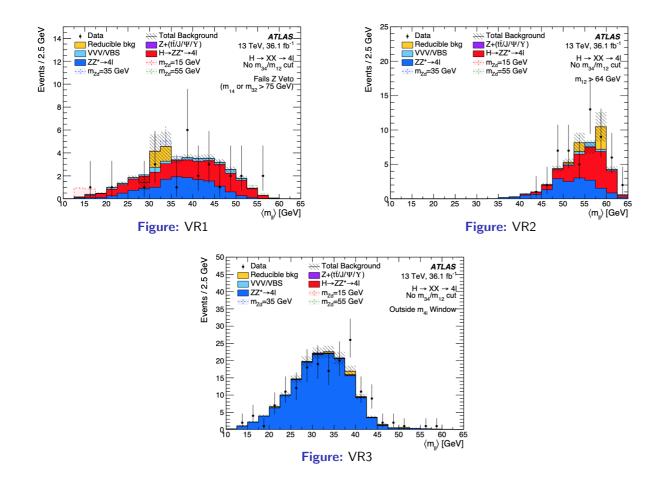




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• VR2

SR cut:  $m_{12,34} < 64$  GeV and  $m_{14,32} < 75$  GeV. VR2 cut: No  $m_{34}/m_{12}$  cut  $m_{12} > 64$  GeV.



• VR1 | SR cut:  $m_{12,34} < 64$  GeV and  $m_{14,32} < 75$  GeV. | VR1 cut: No  $m_{34}/m_{12}$  cut  $m_{14}$  or  $m_{32} > 75$  GeV.

• VR2 | SR cut:  $m_{12,34} < 64$  GeV and  $m_{14,32} < 75$  GeV. | VR2 cut: No  $m_{34}/m_{12}$  cut  $m_{12} > 64$  GeV.

• VR3

 $\begin{array}{l} | \ \, {\rm SR} \ {\rm cut}: m_{12,34} < 64 \ \, {\rm GeV} \ \, {\rm and} \\ m_{14,32} < 75 \ \, {\rm GeV}. \\ | \ \, {\rm VR3} \ \, {\rm cut}: \ \, {\rm No} \ \, m_{34}/m_{12} \ \, {\rm cut} \\ m_{4I} > 130 \ \, {\rm GeV} \ \, {\rm or} \ \, m_{4I} < 115 \ \, {\rm GeV}. \end{array}$ 

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Different statistical analysis procedure are constructed to interpret the results according to our benchmark model:

Model-independent limits

Set limits on the cross section in a fiducial volume such that the limit is suitably model-independent.

- $\rightarrow$  The fiducial volume is defined in a way to mimic the selection aplied in this analysis and appropriate for a Higgs boson ( $m_H = 125 \text{ GeV}$ ) decaying to 2 intermidiate, on-shell, narrow X boson ( $Z_d$ , a).
- $\rightarrow$  Model-independent efficiency:  $\epsilon_c = \frac{N_{eco}^c}{N_{e,d}^c}$
- $\rightarrow \text{ Expectation: } N^{c}_{exp}(\langle m_{II} \rangle) = N^{c}_{bkg} + \sigma^{c}_{fid}.\mathcal{L}umi.\epsilon_{c}.\textit{Gaus}(\langle m_{II} \rangle, \overline{\langle m_{II} \rangle}, \sigma^{c}_{\langle m_{II} \rangle})$

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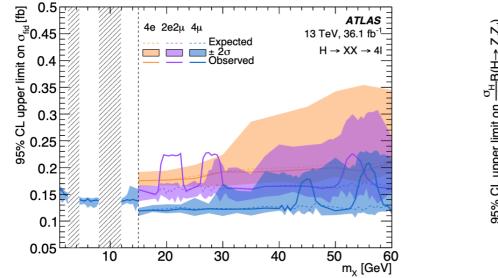
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- $\rightarrow$  Model-independent efficiency:  $\epsilon_c = \frac{N_{c_{eco}}^c}{N_{c_{eff}}^c}$
- $\rightarrow \text{ Expectation: } N^{c}_{exp}(\langle m_{II} \rangle) = N^{c}_{bkg} + \sigma^{c}_{fid}.\mathcal{L}umi.\epsilon_{c}.\textit{Gaus}(\langle m_{II} \rangle, \overline{\langle m_{II} \rangle}, \sigma^{c}_{\langle m_{II} \rangle})$

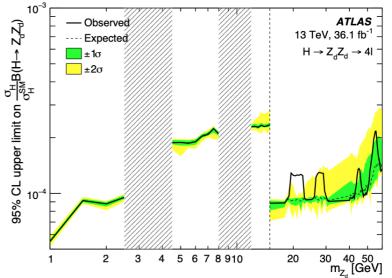
### Model-dependent limits

Set limits on total cross section for  $Z_d Z_d$  model.

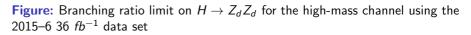
- $\rightarrow$  In that model, the total cross section for a specific model is in request:  $\sigma_H \mathcal{BR}(H \rightarrow Z_d Z_d \rightarrow 4I)$
- $\rightarrow$  Model-dependent acceptance in "total phase space" of a given model:  $\alpha_c = \frac{N_{dd}^{Fd}}{N^{Zd^c}}$

$$\rightarrow \text{ Expectation: } N_{Z_d,exp}^c(\langle m_{II} \rangle) = N_{Z_d,bkg}^c(\langle m_{II} \rangle) + \sigma_H \mathcal{BR}(H \rightarrow Z_d Z_d \rightarrow 4I). \mathcal{L}umi. \frac{\Gamma_{Z_d}^c}{\Gamma_{Z_d}^{4I}}. \alpha_c^{Z_d}. \epsilon_c. \textit{Gaus}(\langle m_{II} \rangle, \overline{\langle m_{II} \rangle}_c, \sigma_{\langle m_{II} \rangle}^c)$$

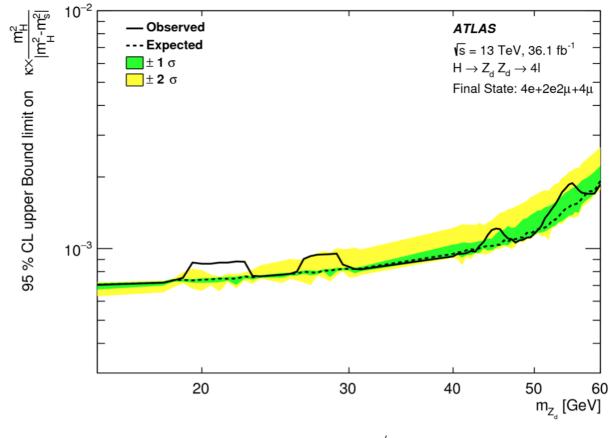




**Figure:** Upper limits at 95% CL on fiducial cross-sections for the  $H \rightarrow XX \rightarrow 4I$  process



- Model-independent fiducial cross-section limit (36  $fb^{-1}$ ) is calculated .
- The limit on  $\mathcal{BR}(H \to Z_d Z_d)$  is model dependent because one needs the  $\mathcal{BR}(Z_d \to II)$  from a model.
- The upper limit on the  $\mathcal{BR}(H \to Z_d Z_d)$  is approximately 0.01%



**Figure:** Upper limit on  $\kappa'$ 

The upper limit on the the higgs mixing parameter  $\kappa$  is around 0.1%

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- Searches are performed for exotic decays of SM Higgs boson to two  $Z_d$  givin 4 leptons in the final state.
- The data are found to be globally consistent with the SM backgroud expectations.
- Limits on fiducial cross-sections are computed in order to be used for testing other benchmark models than treated in this analysis.
- Upper limits; with dependence on the itermediate exotic boson's mass  $(m_{Z_d})$ ; are set on the branching ratio of Higgs boson to  $Z_d Z_d$  and on  $\kappa$ .



# BackUp

# Data-driven estimate for non-dominant processes with Fake leptons

Processes such as  $t\bar{t}$  and WZ may contribute to our selected events if non-lepton objects such as jets are incorrectly reconstructed as leptons.

- Events are selected in an inverted signal region (region B): a region defined by identical cuts to the normal signal region (region A) except with a few cuts inverted: the selected quadruplet in the event must contain one or two (but not more) leptons that are:
  - Electrons failing the LooseLH identification working point or failing the FixedCutLoose isolation working point, but not failing both of these requirements, or
  - Muons failing the FixedCutLoose isolation working point or  $d_0$  significance requirement
    - $\Rightarrow$  Those are bad leptons
- Events are selected in two regions that are rich in Z+jets events, where the event contains two leptons consistent with Z boson and exactly one other baseline reconstructed lepton. These third leptons, which are predominantly leptons faked by hadronic jets, either pass all requirements imposed in the standard analysis selection (the event then contributes to region C) or are leptons failing the cuts described above (the event contributes to region D).
- So Fake Factors are calculated as:  $f = \frac{N_C}{N_D}$
- Those factors are applied to the events in region **B**: events with exactly one bad lepton (**B1**) receive a weight given by the fake factor corresponding to the bad lepton, and events with exactly two bad leptons (**B2**) receive a weight given by the product of the fake factors of the two bad leptons and an additional factor of -1:  $N_{B_1}f N_{B_2}f_1f_2$
- The contribution to this estimate from processes producing four (or more) real leptons is estimated from the MC contribution to the inverted signal region  $(N_{B_1}^{real})$  and  $N_{B_2}^{real}$  with fake factors applied.

$$\mathsf{N}_{\mathcal{A}}^{\mathit{fake}} = (\mathsf{N}_{\mathcal{B}_1}\mathit{f} - \mathsf{N}_{\mathcal{B}_2}\mathit{f}_1\mathit{f}_2) - (\mathsf{N}_{\mathcal{B}_1}^{\mathit{real}}\mathit{f} - \mathsf{N}_{\mathcal{B}_2}^{\mathit{real}}\mathit{f}_1\mathit{f}_2)$$

For more details, see 2226555

