



UNIVERSITÀ DEGLI STUDI DI TORINO

**XI NExT PhD Workshop: Probing fundamental physics at colliders and beyond 2021**

**A NEW SUBTRACTION SCHEME  
IN PERTURBATIVE QCD**

Gloria Bertolotti

In collaboration with:

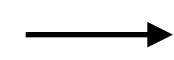
*L. Magnea, E. Maina, G. Pelliccioli, C. Signorile-Signorile, P. Torrielli, S. Uccirati*

*JHEP12(2018)107, JHEP02(2021)037*

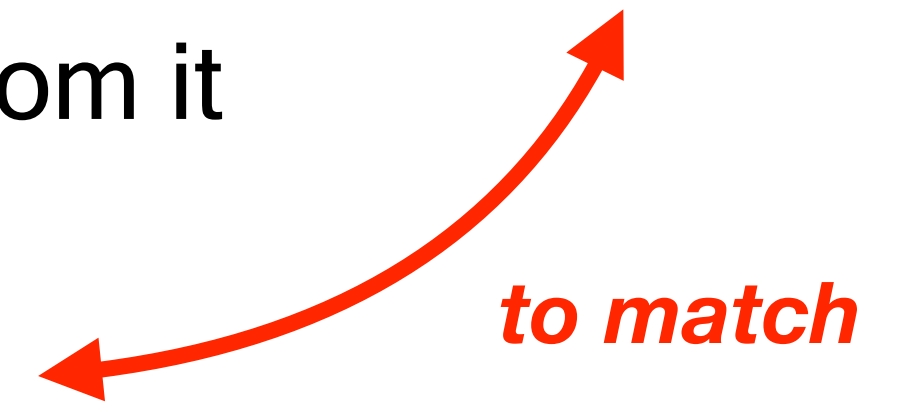
# Motivation

- Future high-precision phase of the LHC = growing experimental data accuracy
- Better control on SM expectation to the identify any deviations from it

↓  
**How?**



Improve theoretical prediction  
of the main multi-particle  
processes in colliders

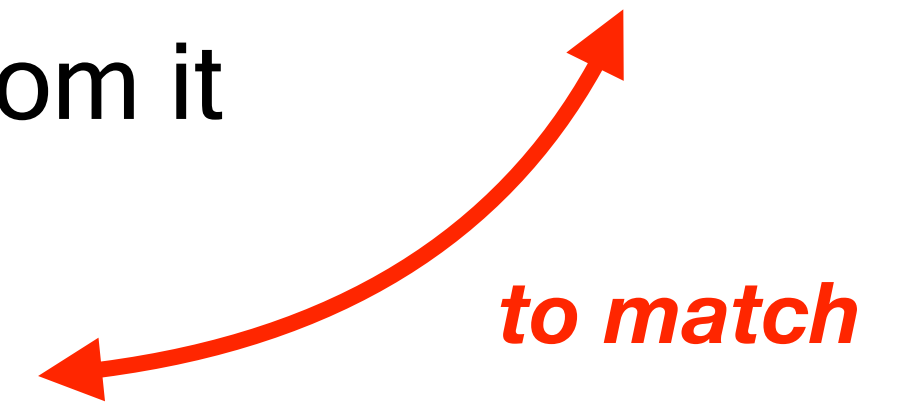


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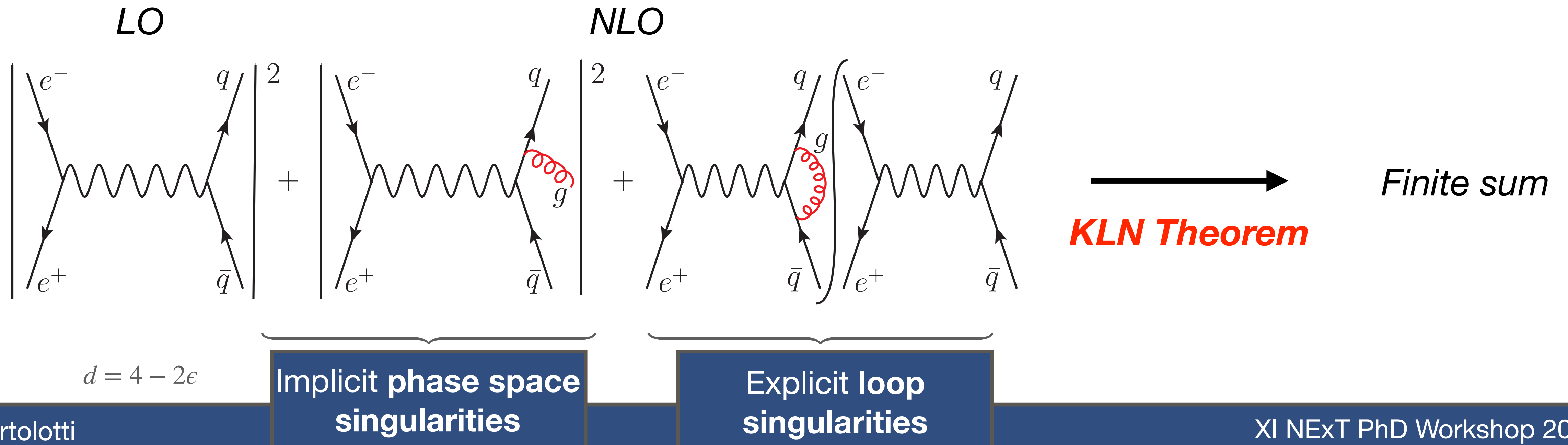
Improve theoretical prediction of the main multi-particle processes in colliders



*to match*

## Problem

Higher orders corrections in perturbative QCD manifest **IR singularities**



## General solution?

### Subtraction procedure

Well established schemes at NLO:

- Catani-Seymour (CS) [\[9602277\]](#)
- Frixione-Kunst-Signer (FKS) [\[9512328\]](#)
- Nagy-Soper [\[0308127\]](#)

### ➔ What about NNLO?

Many proposed schemes but a general and efficient method is still missing due to the increasing complexity of the emerging singularities

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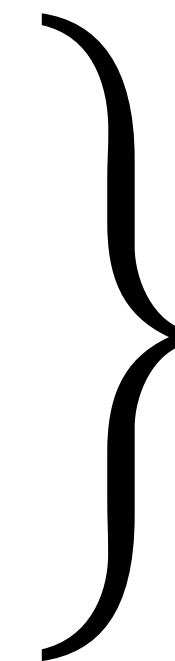
### ➔ What about NNLO?

Many proposed schemes but a general and efficient method is still missing due to the increasing complexity of the emerging singularities

## Local Analytic Sector Subtraction

Why a new scheme?

- Generality
- Locality
- Simplicity of analytical calculation
- Structure minimality



**Ideal features  
for a fully general scheme**

# Local Analytic Sector Subtraction

(massless partons,  $d = 4 - 2\epsilon$ )

- NLO**

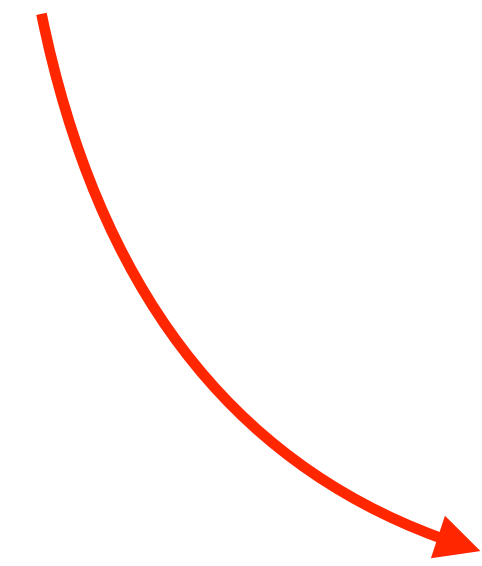
$$\frac{d\sigma^{\text{NLO}}}{dX} = \lim_{d \rightarrow 4} \left\{ \int d\Phi_n V_n \delta_n + \int d\Phi_{n+1} R_{n+1} \delta_{n+1} \right\}$$

Counterterm

$$\frac{d\sigma_{ct}^{\text{NLO}}}{dX} = \int d\Phi_{n+1} \bar{K}_{n+1}$$

Integrated counterterm

$$I_n = \int d\Phi_{\text{rad}} \bar{K}_{n+1}$$



$$\frac{d\sigma^{\text{NLO}}}{dX} = \int d\Phi_n \underbrace{\left( V_n + I_n \right)}_{\text{finite in } \epsilon} \delta_n + \int d\Phi_{n+1} \underbrace{\left( R_{n+1} \delta_{n+1} - \bar{K}_{n+1} \delta_n \right)}_{\text{integrable in } d = 4}$$

# Local Analytic Sector Subtraction (massless partons, $d = 4 - 2\epsilon$ )

$$\frac{d\sigma^{\text{NLO}}}{dX} = \int d\Phi_n \left( V_n + I_n \right) \delta_n + \int d\Phi_{n+1} \left( R_{n+1} \delta_{n+1} - \bar{K}_{n+1} \delta_n \right)$$

## Key features

- **Phase-space partition (FKS)** : identify each singular contribution exactly

Sector functions  $\mathcal{W}_{ij}$  :

$$R = \sum_{i,j} R \mathcal{W}_{ij} = R \mathcal{W}_{31} + R \mathcal{W}_{32} + \dots$$

Sums rules

$$\sum_{i \in \mathcal{F}} \sum_{j \neq i} W_{ij} = 1$$

$$\sum_{j \neq i} S_i W_{ij} = 1$$

$$C_{ij} (W_{ij} + W_{ji}) = 1$$

$$S_i C_{ij} W_{ij} = C_{ij} S_i W_{ij} = 1$$

$$W_{ij} = \frac{\sigma_{ij}}{\sigma} \quad \sigma_{ij} = \frac{1}{\epsilon_i \omega_{ij}} = \frac{S_{qj}}{S_{ij}}$$

$$\sigma = \sum_{i \in \mathcal{F}} \sum_{j \neq i} \sigma_{ij}$$

# Local Analytic Sector Subtraction (massless partons, $d = 4 - 2\epsilon$ )

$$\frac{d\sigma^{\text{NLO}}}{dX} = \int d\Phi_n \left( V_n + I_n \right) \delta_n + \int d\Phi_{n+1} \left( R_{n+1} \delta_{n+1} - \bar{K}_{n+1} \delta_n \right)$$

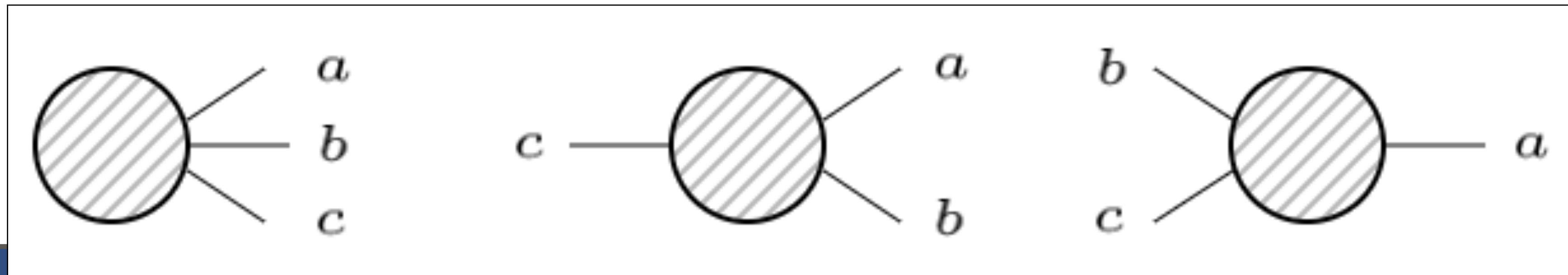
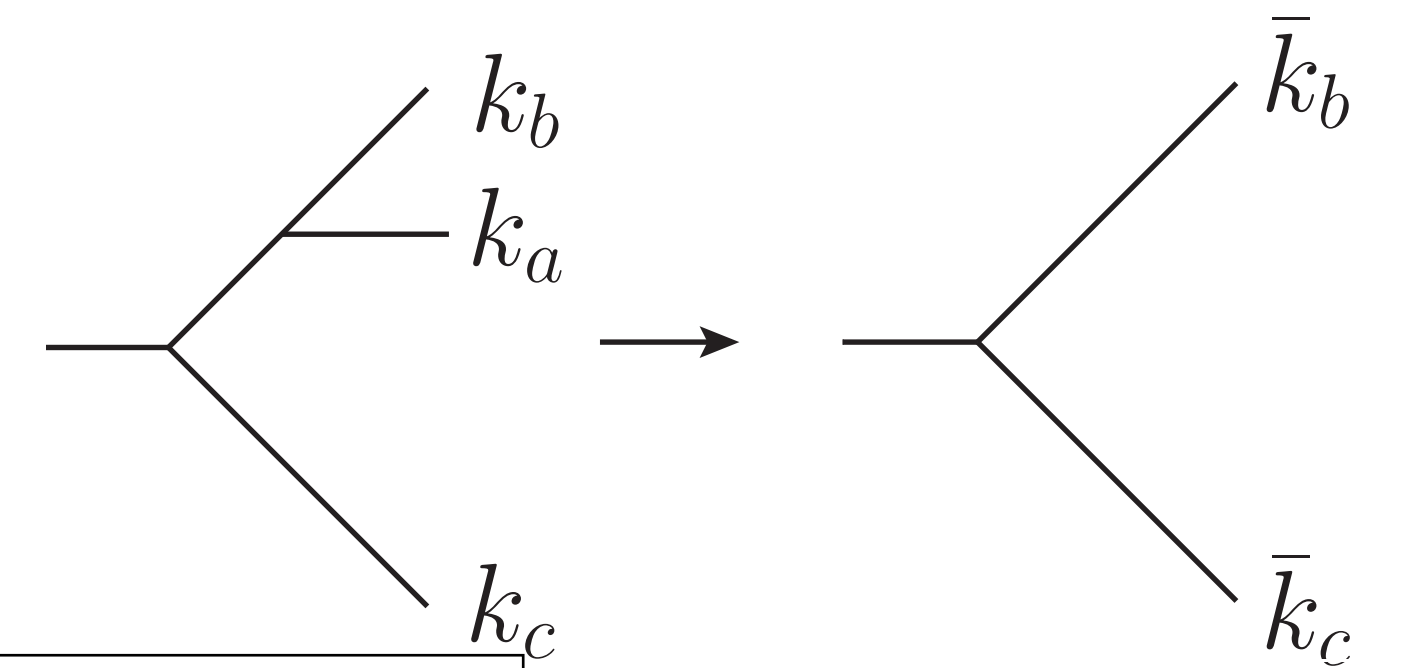
## Key features

- **Phase-space partition (FKS)** : identify each singular contribution exactly
- **Momenta mapping (CS)** : phase-space factorisation and on-shell momentum conserving Born-level kinematics

Phase space factorisation:  $d\Phi_{n+1} = d\bar{\Phi}_n d\bar{\Phi}_{\text{rad}}$

Mapped kinematics:  $\{\bar{k}\}^{(abc)} = \left\{ \{k\}_{a b \bar{c}}, \bar{k}_b^{(abc)}, \bar{k}_c^{(abc)} \right\}$

$$\bar{k}_b^{(abc)} + \bar{k}_c^{(abc)} = k_a + k_b + k_c$$





# Local Analytic Sector Subtraction (massless partons, $d = 4 - 2\epsilon$ )

$$\frac{d\sigma^{\text{NLO}}}{dX} = \int d\Phi_n \left( V_n + I_n \right) \delta_n + \int d\Phi_{n+1} \left( R_{n+1} \delta_{n+1} - \bar{K}_{n+1} \delta_n \right)$$

## Key features

- **Phase-space partition (FKS)** : identify each singular contribution exactly
- **Momenta mapping (CS)** : phase-space factorisation and on-shell momentum conserving Born-level kinematics
- **Candidate counterterm** with remapped kinematics

$$\bar{K}_{ij} = \left[ \bar{S}_i + \bar{C}_{ij} - \bar{S}_i \bar{C}_{ij} \right] R W_{ij}$$

Consistency relations

$$S_i \bar{S}_i R = S_i R$$

$$S_i \bar{S}_i \bar{C}_{ij} R = S_i \bar{C}_{ij} R$$

$$C_{ij} \bar{C}_{ij} R = C_{ij} R$$

$$C_{ij} \bar{S}_i \bar{C}_{ij} R = C_{ij} \bar{S}_i R$$

ensure cancellation  
in each sector

$$RW_{ij} - \bar{K}_{ij} \longrightarrow \textit{finite}$$

# Local Analytic Sector Subtraction (massless partons, $d = 4 - 2\epsilon$ )

$$\frac{d\sigma^{\text{NLO}}}{dX} = \int d\Phi_n \left( V_n + I_n \right) \delta_n + \int d\Phi_{n+1} \left( R_{n+1} \delta_{n+1} - \bar{K}_{n+1} \delta_n \right)$$

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$$\bar{K}_{ij} = \left[ \bar{S}_i + \bar{C}_{ij} - \bar{S}_i \bar{C}_{ij} \right] R W_{ij}$$

$$\bar{K} = \sum_{i \in \mathcal{F}} \bar{S}_i R + \sum_{i \in \mathcal{F}} \sum_{\substack{j \in \mathcal{F} \\ j > i}} \bar{C}_{ij} (1 - \bar{S}_i - \bar{S}_j) R + \sum_{i \in \mathcal{F}} \sum_{j \in \mathcal{I}} \bar{C}_{ij} (1 - \bar{S}_i) R$$

Soft

Hard-Collinear  
with final j

Hard-Collinear  
with initial j

# Local Analytic Sector Subtraction (massless partons, $d = 4 - 2\epsilon$ )

$$\frac{d\sigma^{\text{NLO}}}{dX} = \int d\Phi_n \left( V_n + I_n \right) \delta_n + \int d\Phi_{n+1} \left( R_{n+1} \delta_{n+1} - \bar{K}_{n+1} \delta_n \right)$$

## Key features

- **Phase-space partition (FKS)** : identify each singular contribution exactly
- **Momenta mapping (CS)** : phase-space factorisation and on-shell momentum conserving Born-level kinematics
- **Candidate counterterm** with remapped kinematics
- **Simple analytic calculation** of the counterterm

$$\bar{S}_i R(\{k\}) \propto \sum_{c,d \neq i} \frac{s_{cd}}{s_{ic} s_{id}} B_{cd}(\{k\}^{(icd)})$$



$$\bar{C}_{ij} R(\{k\}) \propto \frac{1}{s_{ij}} P_{ij}^{\mu\nu} B_{\mu\nu}(\{k\}^{(ijr)})$$

$$I^s \propto \sum_{c,d} \int d\Phi_{\text{rad}} \left( s_{cd}^{(icd)}; y, z, \phi \right) \frac{s_{cd}}{s_{ic} s_{id}} B_{cd}(\{\bar{k}\}^{(icd)})$$

$$\propto \sum_{c,d} B_{cd}(\{\bar{k}\}^{(icd)}) \frac{(4\pi)^{\epsilon-2} \Gamma(1-\epsilon) \Gamma(2-\epsilon)}{\epsilon^2 \Gamma(2-3\epsilon)}$$

## What about NNLO?

$$\begin{aligned}
 \frac{d\sigma^{NNLO}}{dX} = & \int d\Phi_n \left( \underbrace{VV + I^{(2)} + I^{(RV)}}_{\text{finite in } d=4, \text{ finite in PS}} \right) \delta_n(X) \\
 & + \int d\Phi_{n+1} \left[ \underbrace{\left( RV + I^{(1)} \right)}_{\text{finite in } d=4, \text{ div. in PS}} \delta_{n+1}(X) - \underbrace{\left( \overline{K}^{(RV)} + I^{(12)} \right)}_{\text{finite in } d=4, \text{ div. in PS}} \delta_n(X) \right] \\
 & + \int d\Phi_{n+2} \left[ \underbrace{RR \delta_{n+2}(X) - \overline{K}^{(1)} \delta_{n+1}(X) - \left( \overline{K}^{(2)} - \overline{K}^{(12)} \right) \delta_n(X)}_{\text{finite in } d=4, \text{ finite in PS}} \right]
 \end{aligned}$$

## Results

- ◆ NLO scheme recently extended to ISR for massless QCD  
+ ongoing code implementation
- ◆ NNLO scheme:  
main singular limits and analytic integrations are available

## Future plans

- ➔ Generalisation of the NNLO subtraction to initial-state massless radiation
- ➔ Numerical implementation of the NNLO Local Sector Subtraction
- ➔ Extension to the treatment of the massive case

*Thanks for  
your attention!*