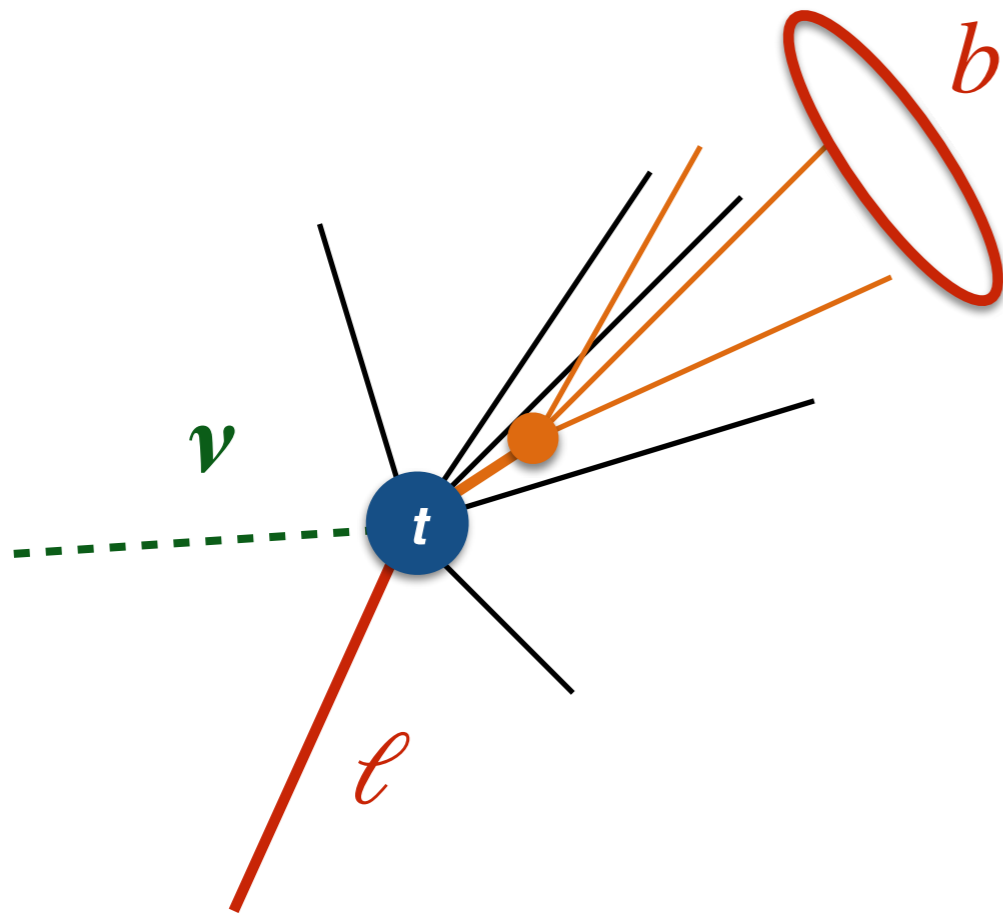


measuring b -fragmentation in $t\bar{t}$ events with ATLAS



Chris Pollard
on b -half of the ATLAS collaboration

**LHC Top WG meeting
2020 11 23**

<https://cds.cern.ch/record/2730444>

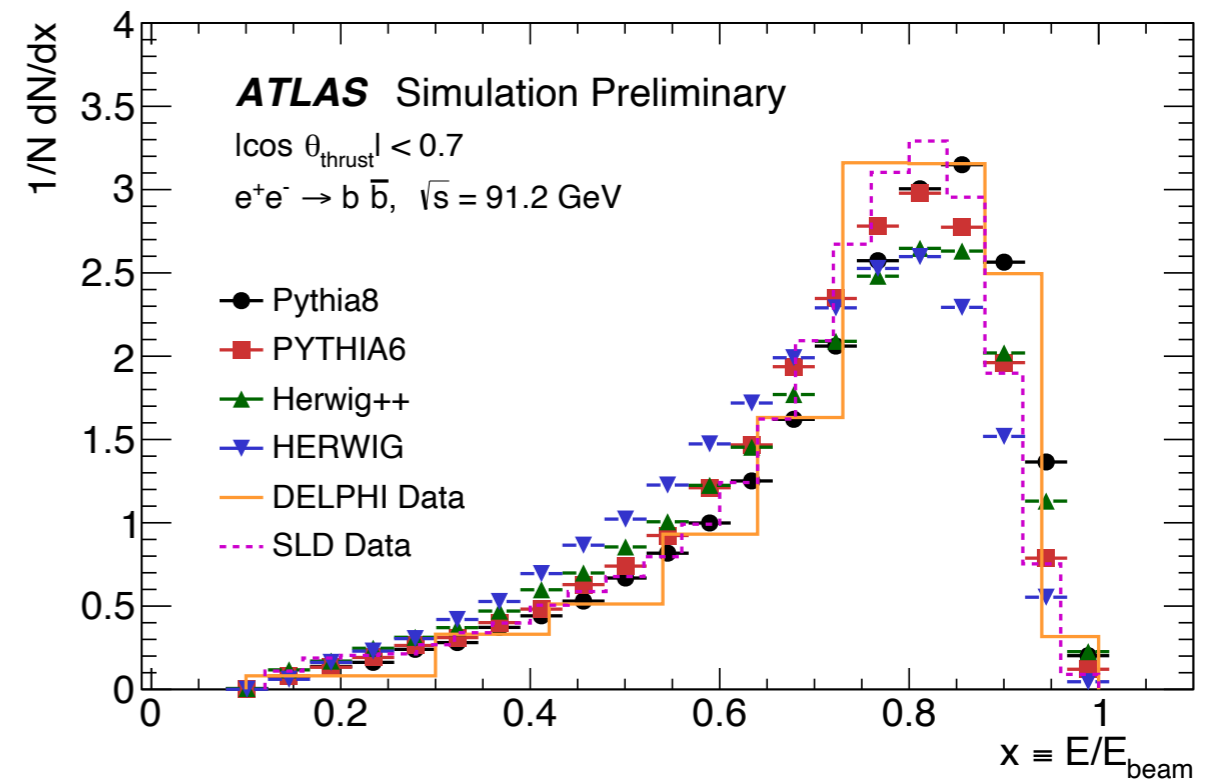
ATLAS-CONF-2020-050



motivation I

the fragmentation of b -quarks into hadrons is of interest for many reasons.

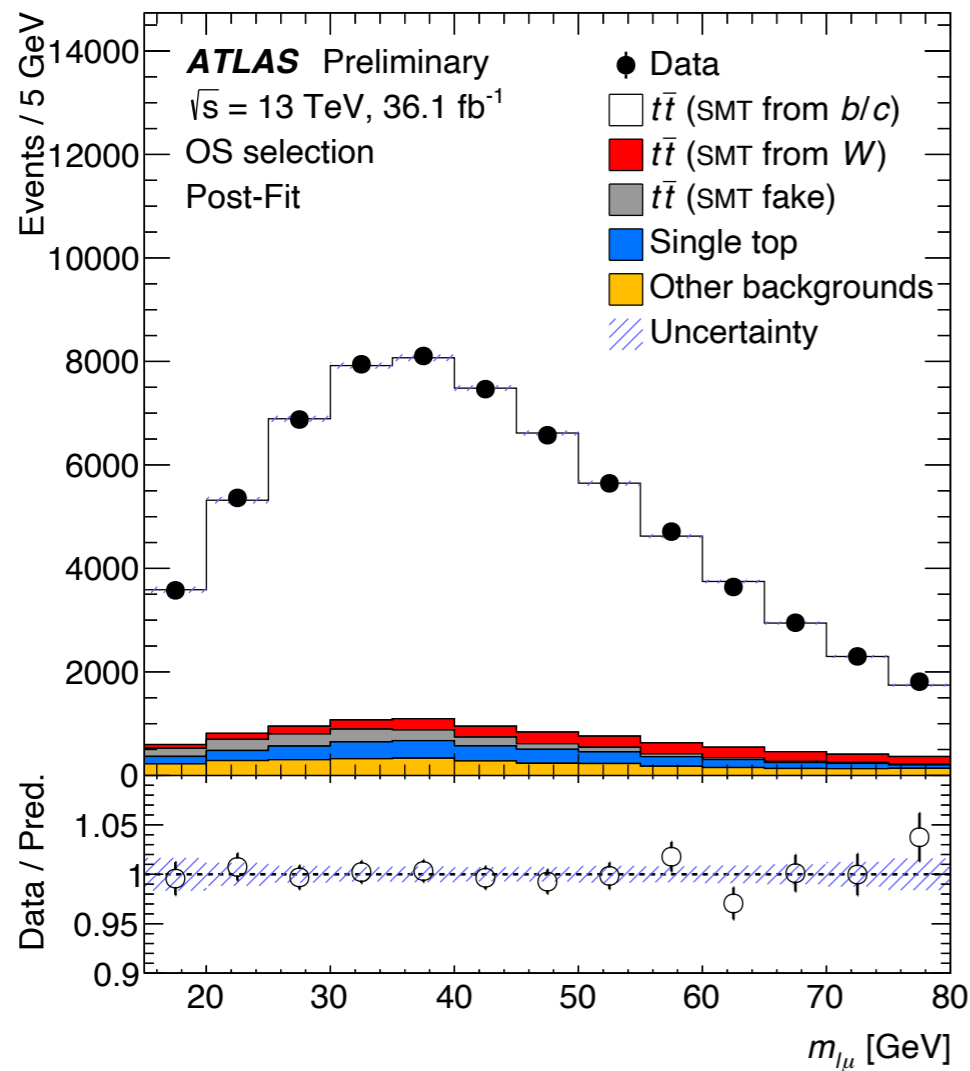
- b -hadrons leave a striking experimental signature and...
- a unique correspondence to the originating b -quarks
 - ergo a precise probe of QCD
- b -fragmentation currently tuned to e^+e^- data ...
- then extrapolated to the LHC environment
 - to what degree is this correct?



ATLAS-PHYS-PUB-2014-008

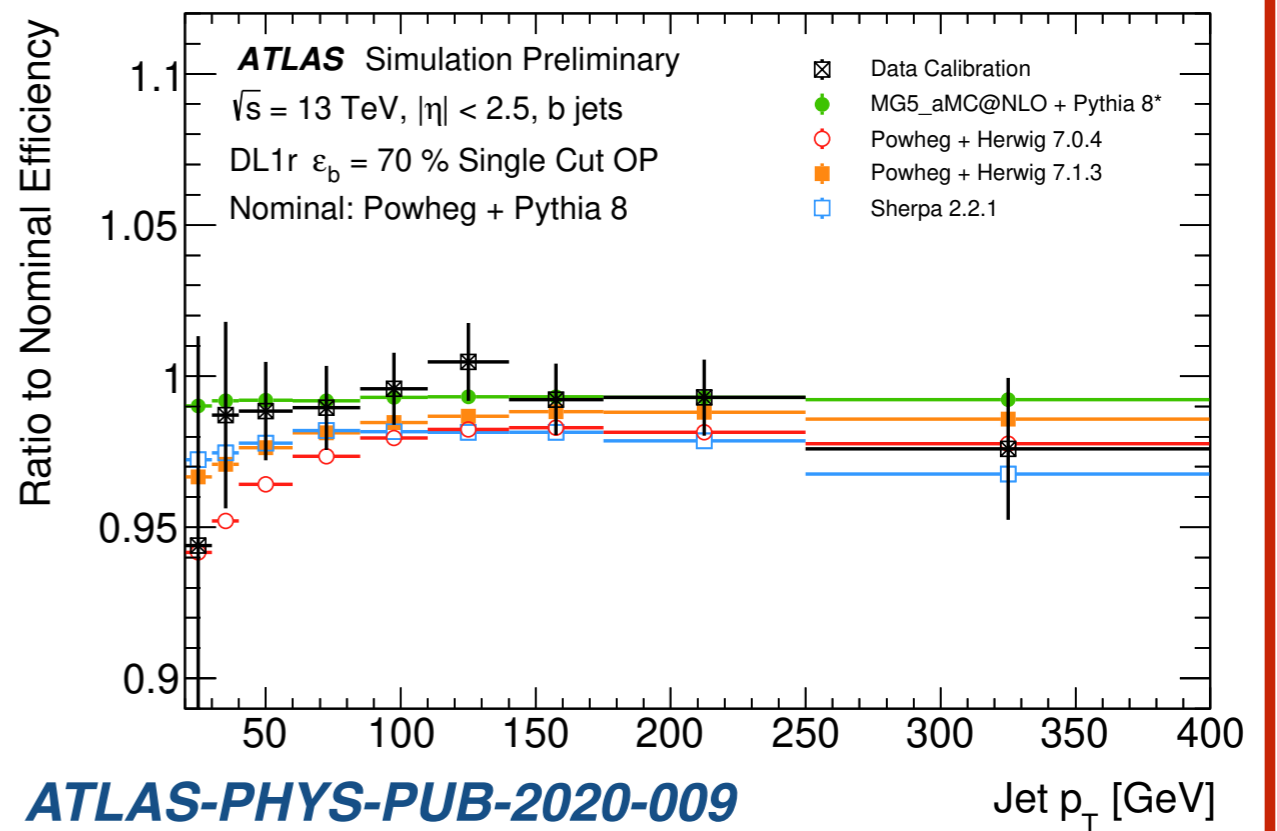
some tension between $e^+e^- \rightarrow b\bar{b}$ measurements of b -fragmentation
parton-shower generators are also not in good agreement

motivation II



$b \rightarrow \mu$ momentum transfer key for new top mass measurements

critical for delivering the best LHC physics results with b -jets.
 b -tagging efficiency and b -jet response are very sensitive to fragmentation.



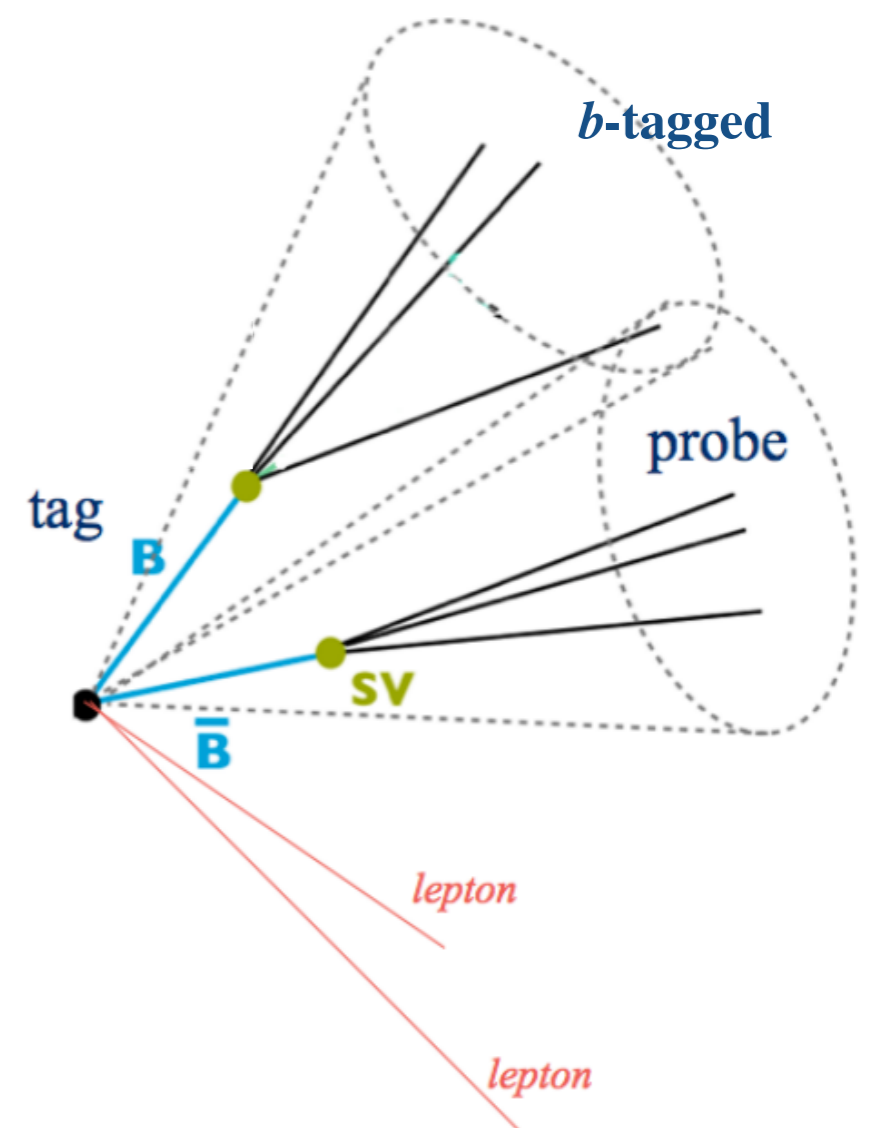
analysis strategy

1. obtain sample of b -jets from $t\bar{t} \rightarrow e\mu\nu\nu b\bar{b}$ events
2. identify b -hadron candidates using ATLAS single secondary vertex finder (SSVF)
3. measure properties of the b -hadron sensitive to the modeling of b -quark fragmentation

overview: detector-level event selection

obtain sample of b -jets from $t\bar{t} \rightarrow e\mu\nu\nu bb$ events

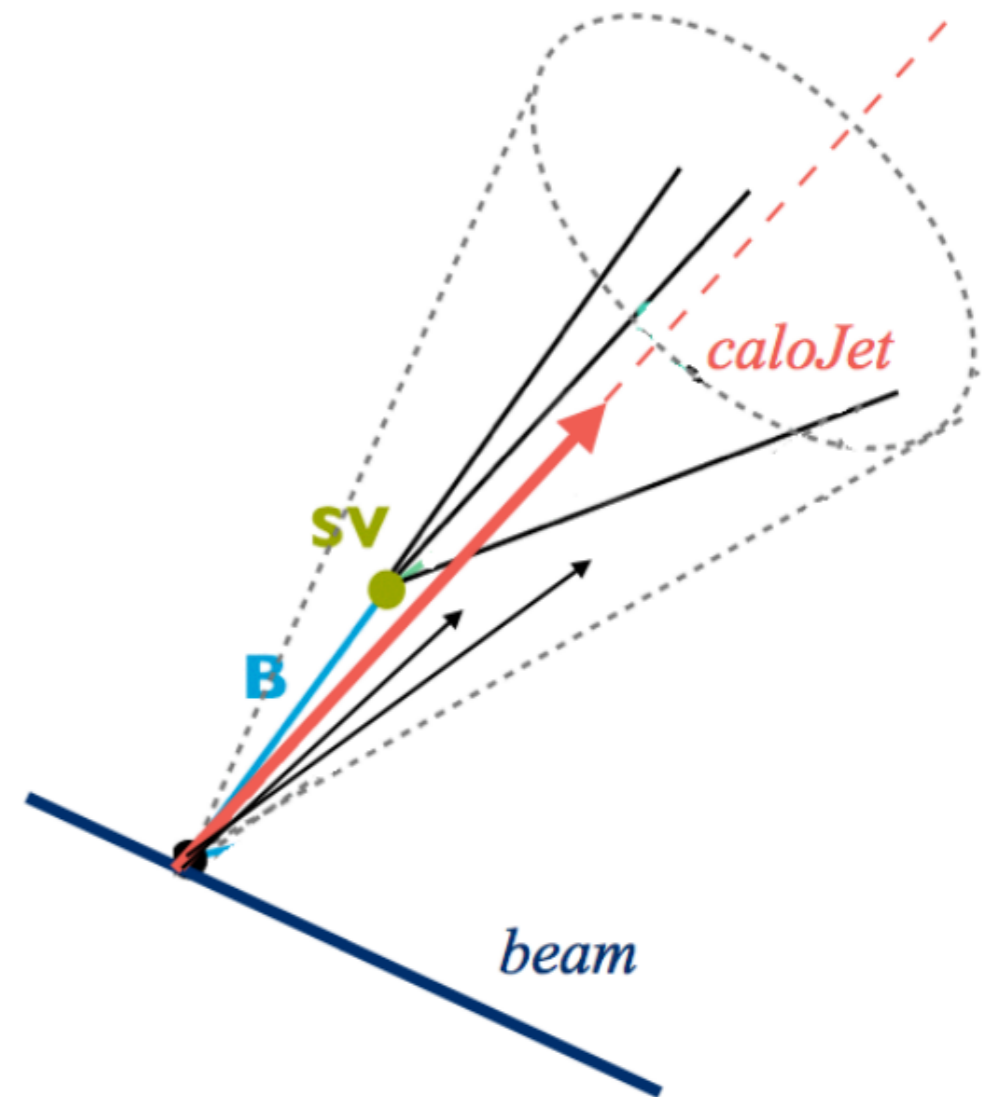
- require $\sum 1 e$ and $\sum 1 \mu$ with opposite charge
- $\sum 2$ anti- k_t $R = 0.4$ calorimeter jets
 - $|\eta| < 2.5, p_T > 30$ GeV
 - at least **one jet must be b -tagged** (“the tag”)
 - we measure the other jet (“the probe”)
 - if both jets are b -tagged, **both jets are measured**



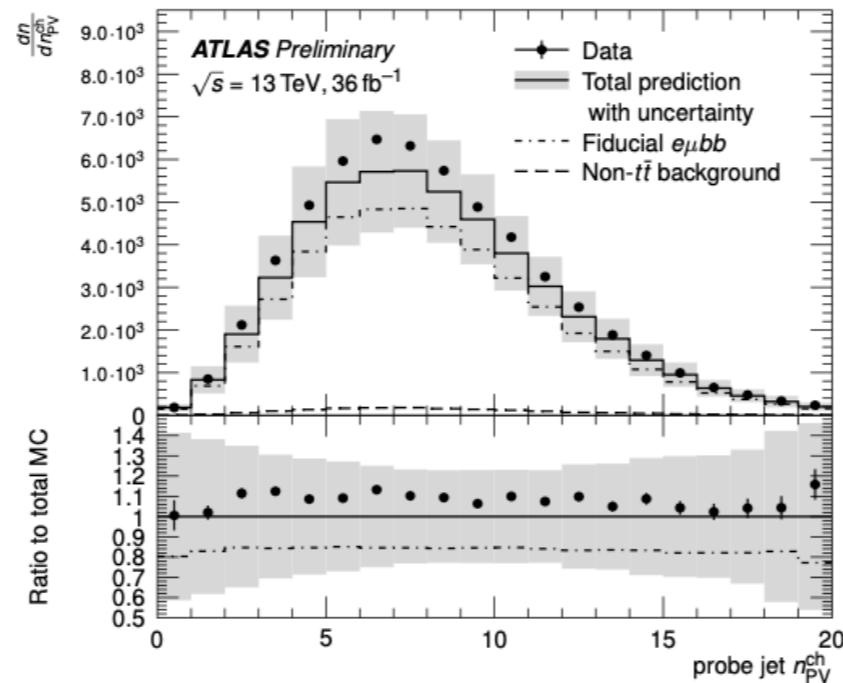
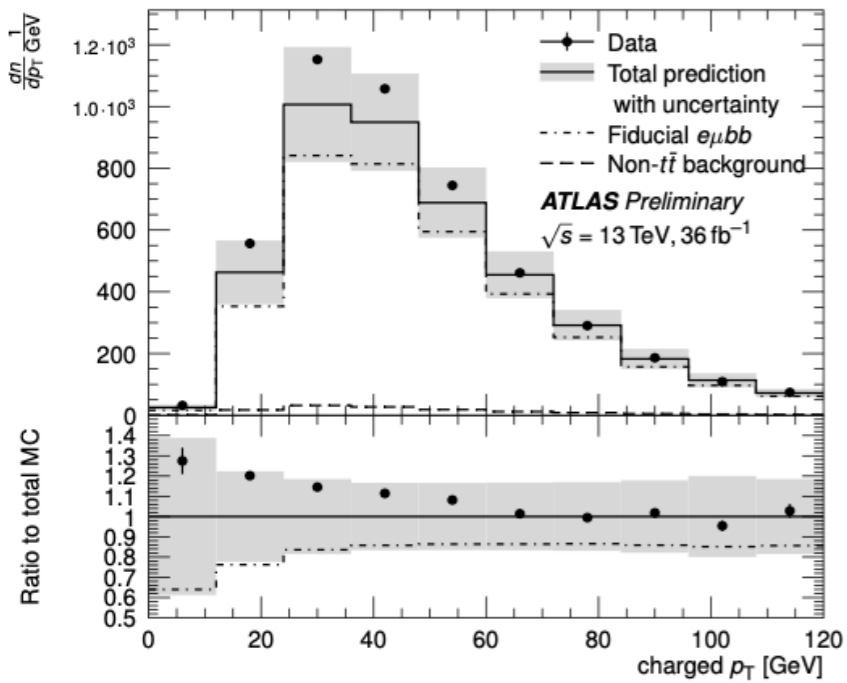
b-hadron reconstruction

identify *b*-hadron candidates using ATLAS's single secondary vertex finder (SSVF)

- we require a reconstructed SSVF vertex with ≥ 3 tracks for “probe” jets
- this is *nearly* a “second *b*-tag” and significantly reduces non-*b* backgrounds.
 - **95+% *b*-jet purity**
- the SSVF vertex **only has tracking information**
- track constituents of the SV make up the *b*-hadron candidate.
- tracks ghost-associated to the jet make up the jet charged component.

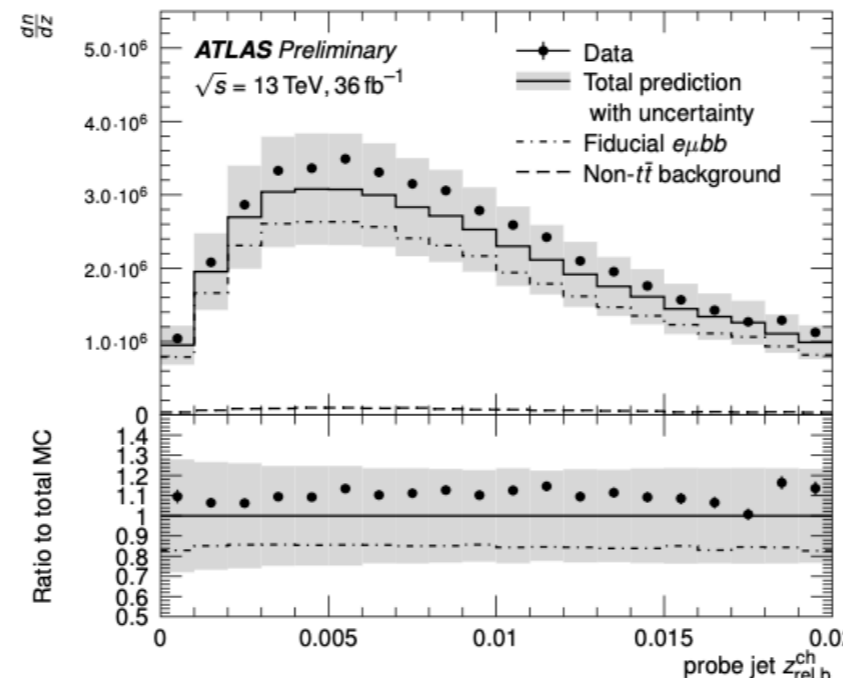
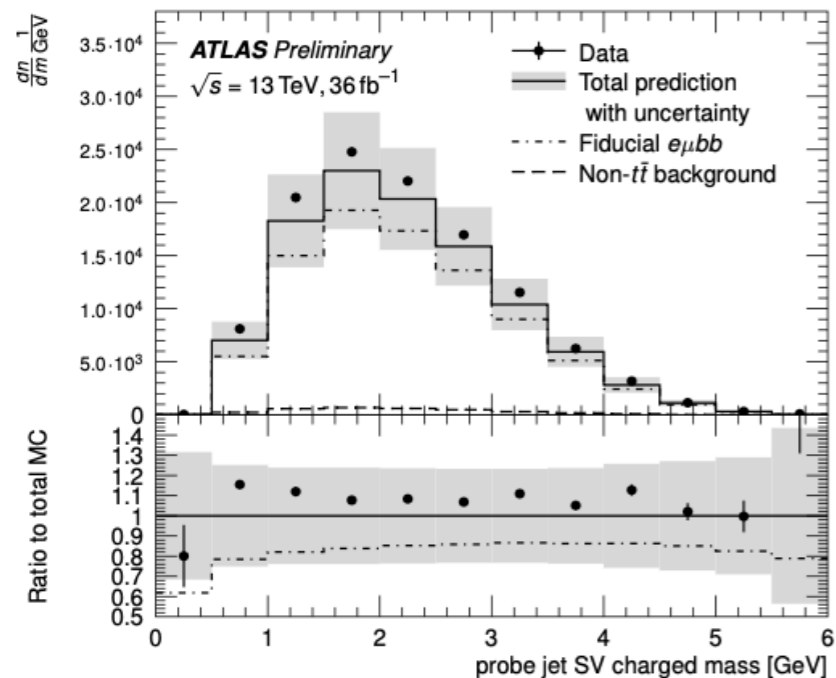


detector-level control plots



charged component of jet p_T
has similar slope to jet p_T
(data vs MC taken as uncertainty)

other relevant PV and SV
quantities are well-modeled



SV mass and relative momentum
transverse to jet axis,

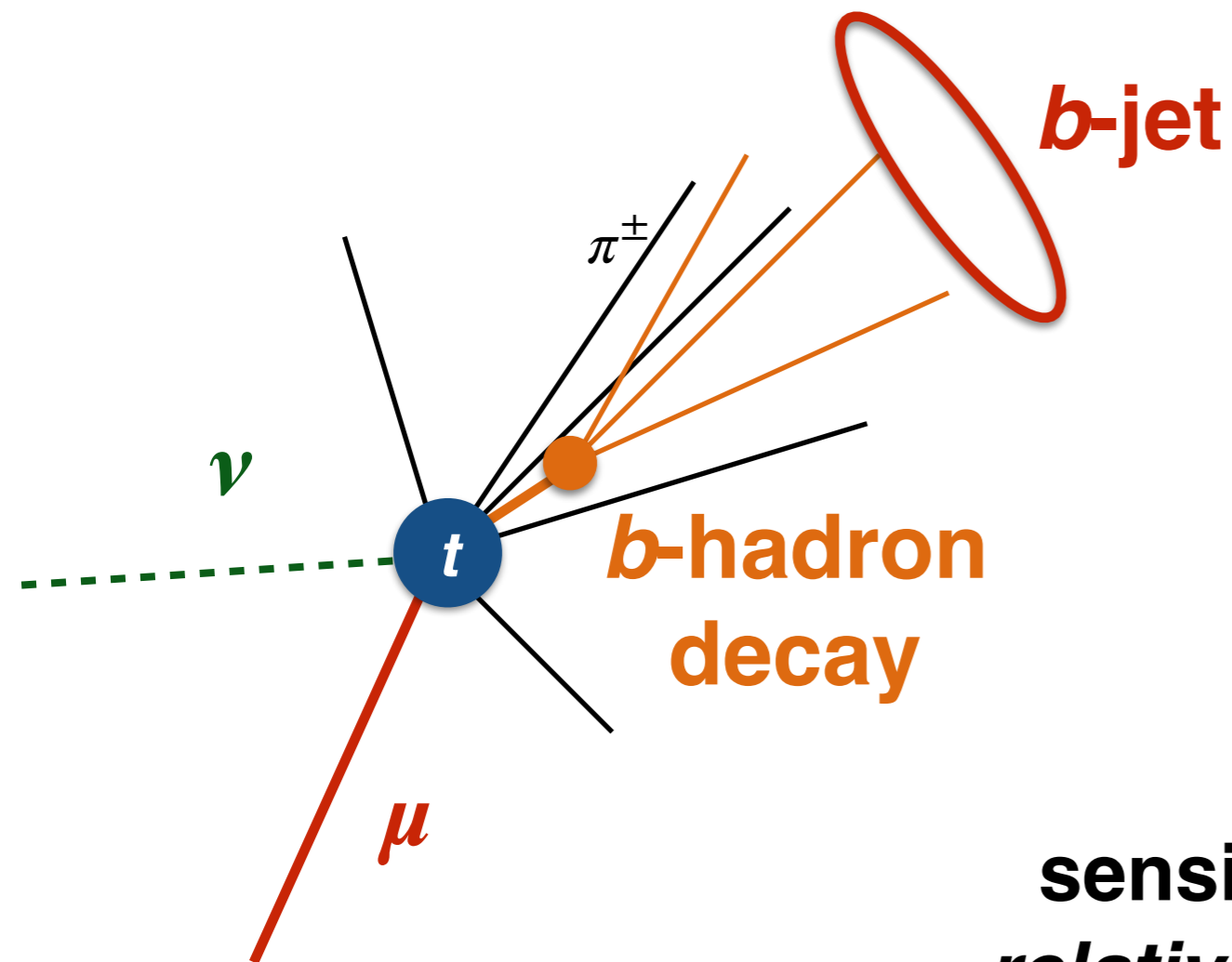
$$z_{\text{rel},b}^{\text{ch}} = (\vec{p}_b^{\text{ch}} \times \vec{p}_{\text{jet}}^{\text{ch}}) / |\vec{p}_{\text{jet}}^{\text{ch}}|^2$$

are also shown.

fiducial definition

- $t\bar{t}$ events
 - single top is treated as background for ease of interpretation
- == 1 μ , == 1 e with dressed $p_T > 25$ GeV, $|\eta| < 2.5$
- == 2 b -jets with $|\eta| < 2.1$, $p_T > 30$ GeV, $\Delta R(\text{jet}, \text{jet}) > 0.5$
- b -jet defined as
 - anti- k_t $R = 0.4$ jet of stable particles
 - prompt charged leptons and all neutrinos are excluded
 - == 1 weakly-decaying b -hadron ($p_T > 5$ GeV) within $\Delta R(\text{jet}, \text{hadron}) < 0.3$
- **stable charged particles with $p_T > 500$ MeV used to calculate observables**
 - stable charged b -hadron decay products make up \vec{p}_b^{ch}
 - stable charged particles ghost-associated to jet make up $\vec{p}_{\text{jet}}^{\text{ch}}$

observables I



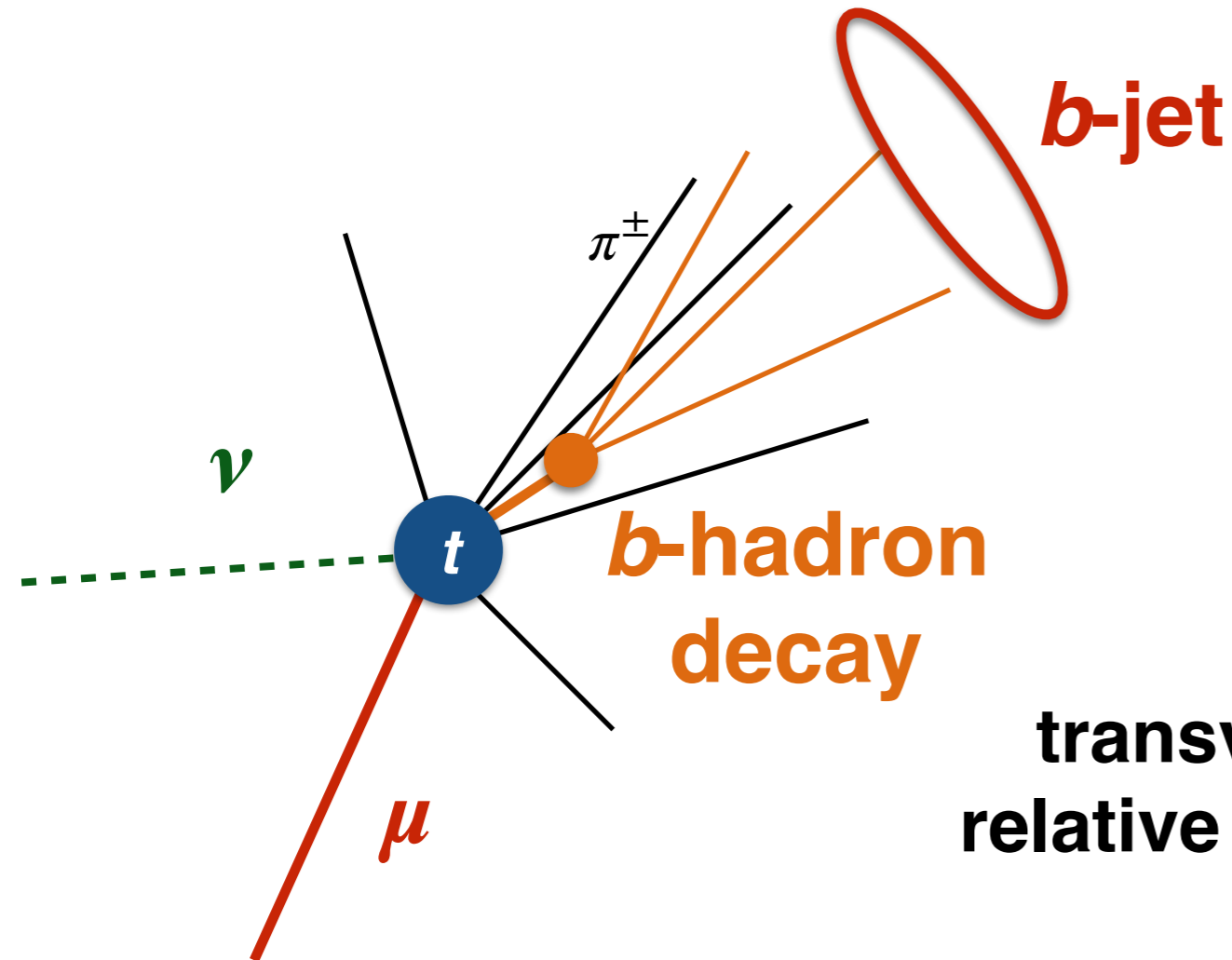
$$z_{L,b}^{\text{ch}} = \frac{\vec{p}_b^{\text{ch}} \cdot \vec{p}_{\text{jet}}^{\text{ch}}}{|p_{\text{jet}}^{\text{ch}}|^2},$$

$$z_{T,b}^{\text{ch}} = \frac{p_{T,b}^{\text{ch}}}{p_{T,\text{jet}}^{\text{ch}}},$$

sensitive to momentum of b -hadron
relative to the nearby hadronic activity

seeded by the particle jet direction, but **only charged particles are used to calculate fiducial observables**

observables II



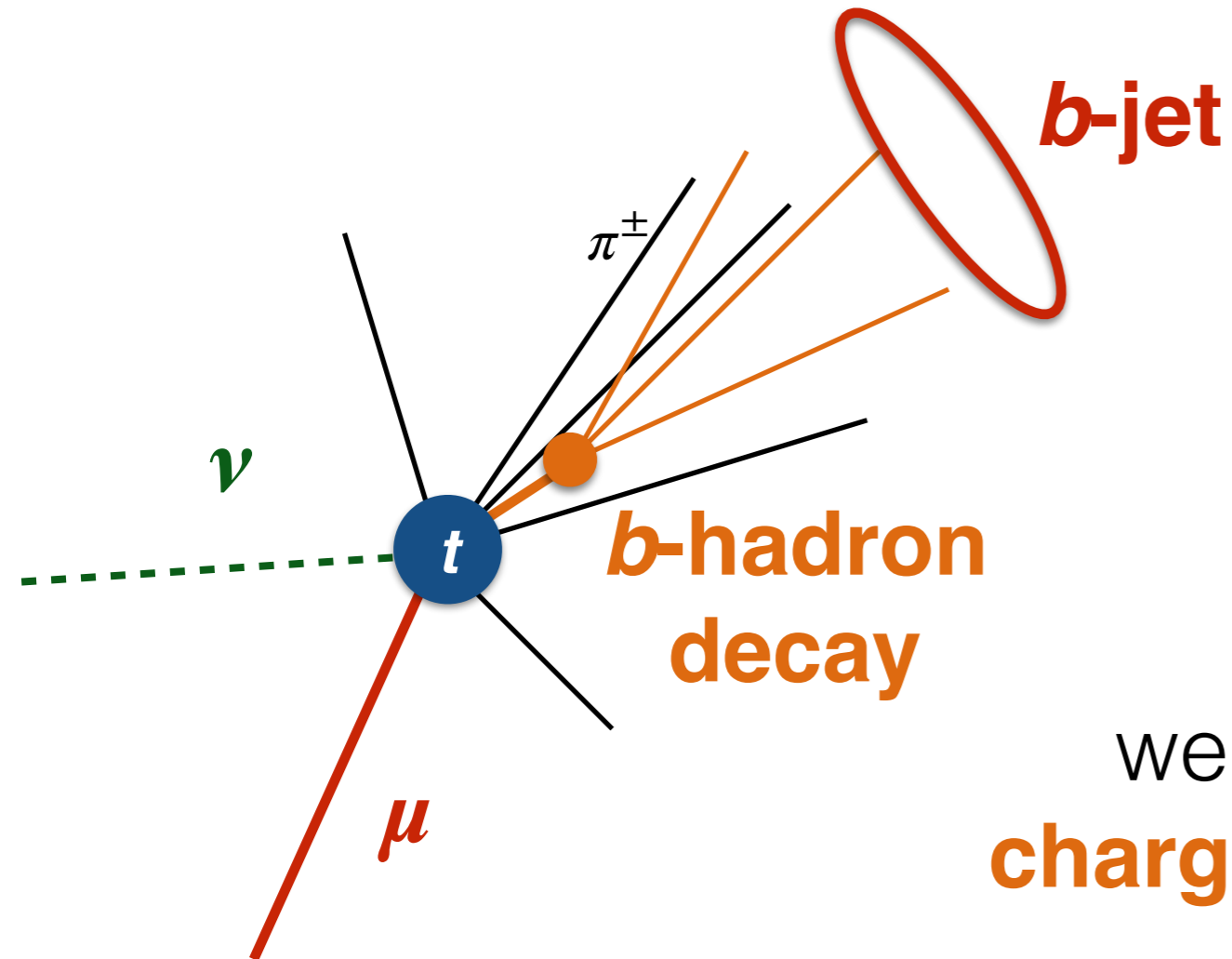
$$\rho = \frac{2p_{T,b}^{\text{ch}}}{p_T^e + p_T^\mu}$$

transverse momentum of b -hadron
relative to the average pt of the leptons
in the event

sensitive to choice of e.g. $\alpha_S(\text{FSR})$ but **not** $\alpha_S(\text{ISR})$

sensitivity also to wide-angle radiation
in top-quark decays

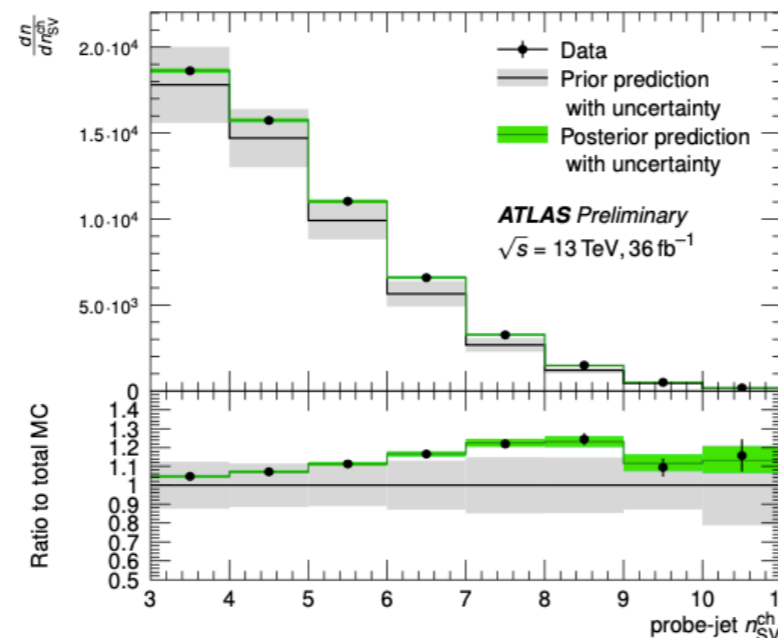
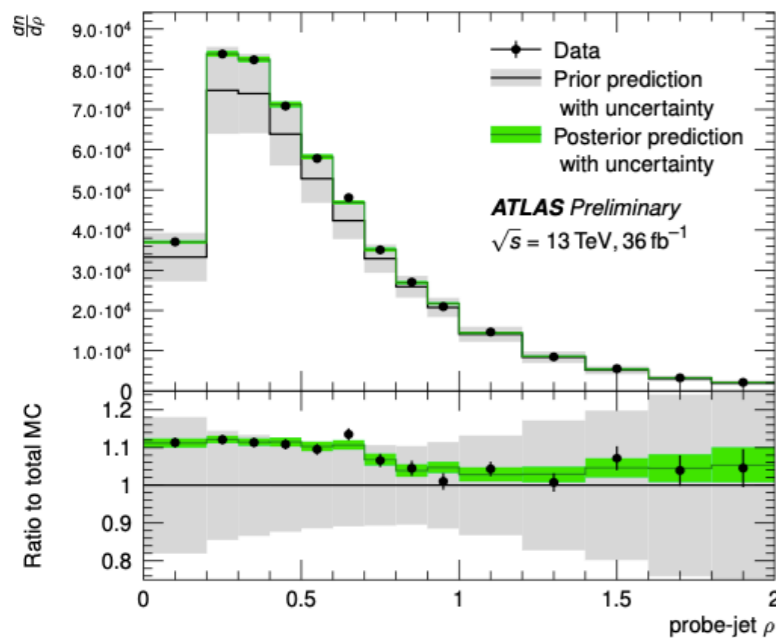
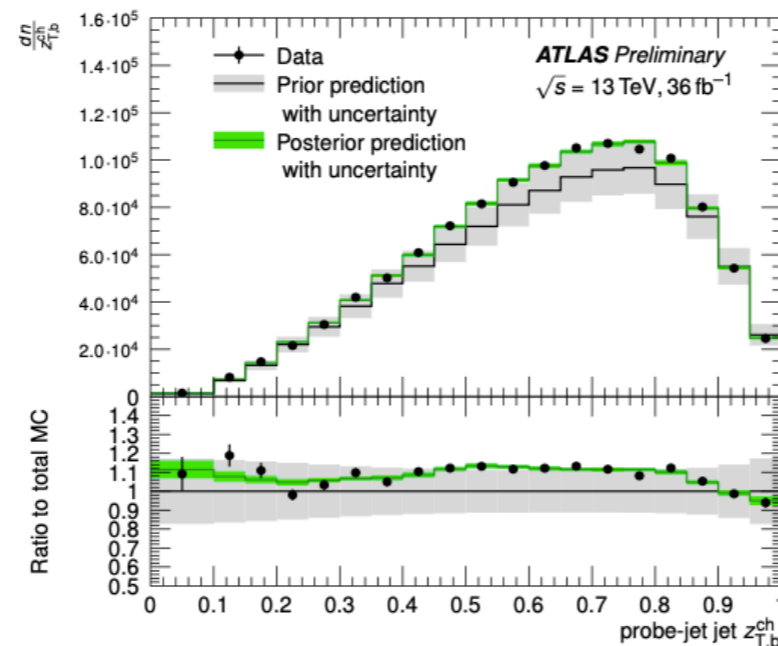
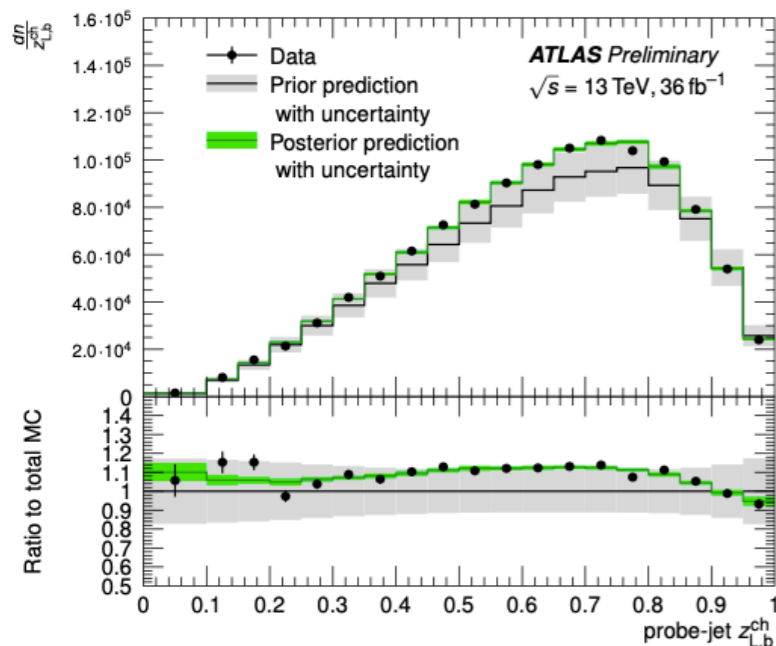
observables III



we also unfold the **number of charged, stable b -hadron children with $p_T > 500$ MeV.**

sensitive to the of b -hadron species production rates

detector-level data/MC comparison of observables



for now please ignore posterior prediction and green uncertainty bands (we'll return to this)

reasonable agreement observed, with some small systematic differences

unfolding (FBU)

- based on the "Fully Bayesian Unfolding" technique (arxiv: 1201.4612)
- we use the following likelihood to unfold the detector effects:

$$\mathcal{L}(d|\sigma, \Lambda) = \prod_{i \in \text{recobins}} \text{Pois}(d_i|x_i(\sigma, \Lambda))$$

$$x_i(\sigma, \Lambda) = L(\Lambda) \times (b_i(\Lambda) + M_{ij}(\Lambda) \sigma_j)$$

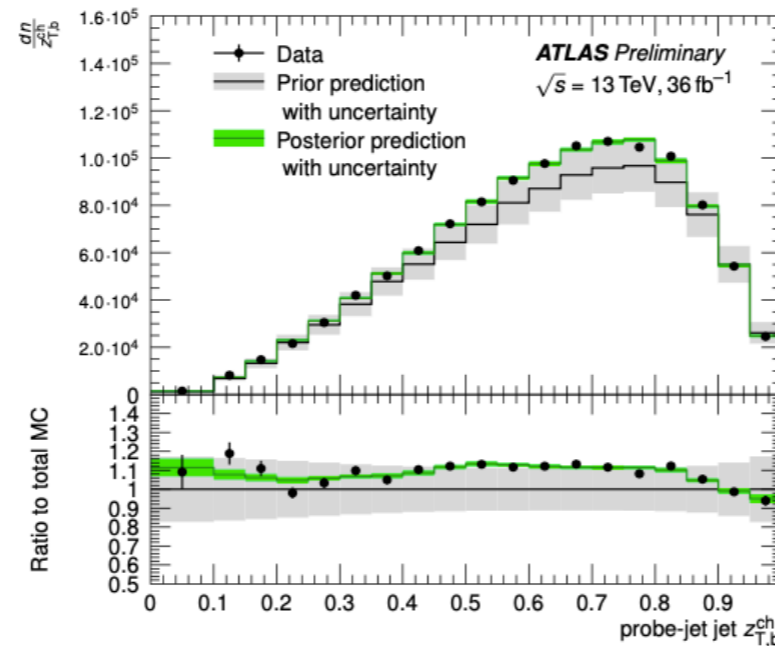
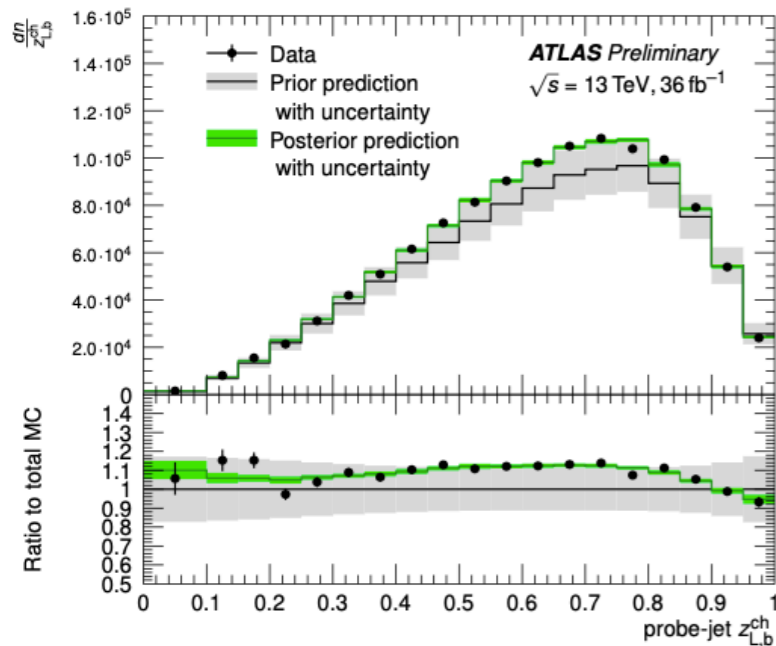
- where \vec{x} is the total number of predicted events in detector-level bins, L is the luminosity, \vec{b} are predicted background events, M_{ij} is probability that a particle-level jet in bin i is reconstructed in detector-level bin j , and Λ are the systematic nuisance parameters.
- we marginalize over nuisance parameters through MCMC sampling to build the marginal posteriors on $\vec{\sigma}$, the parameters-of-interest.

systematic uncertainties

the dominant uncertainties in the unfolding model come from

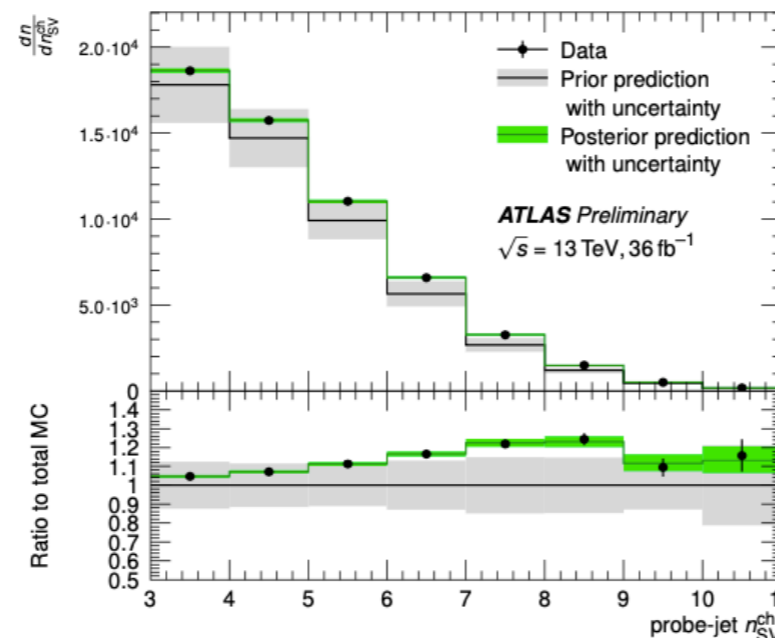
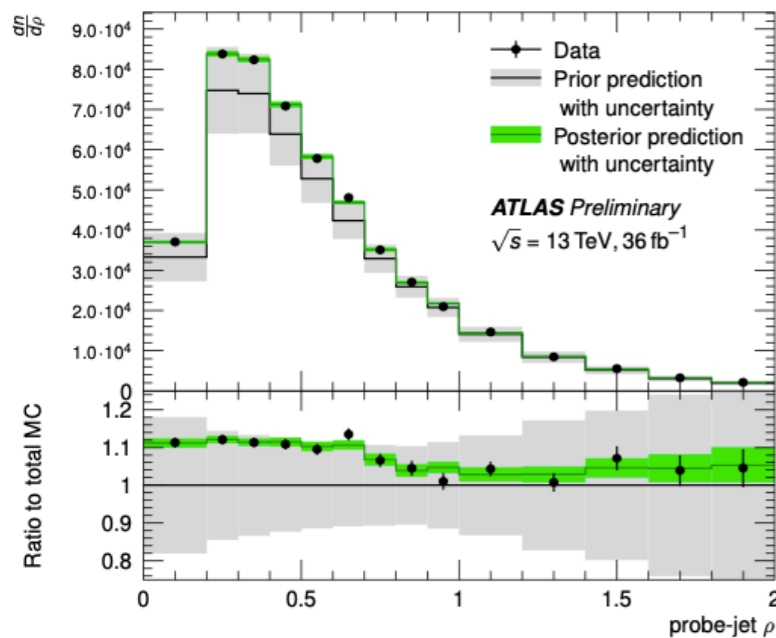
- track efficiency and impact-parameter resolutions
- $t\bar{t}$ modeling: non-fiducial background and detector response
 - Pythia vs Herwig
 - A14 α_S^{FSR} tune variations
 - jet charged p_T and n_{SV}^{trk} mismodeling
- pileup $\langle\mu\rangle$ distribution

detector-level data/MC comparison before/after unfolding



detector-level observables
before (prior) and
after (posterior) unfolding

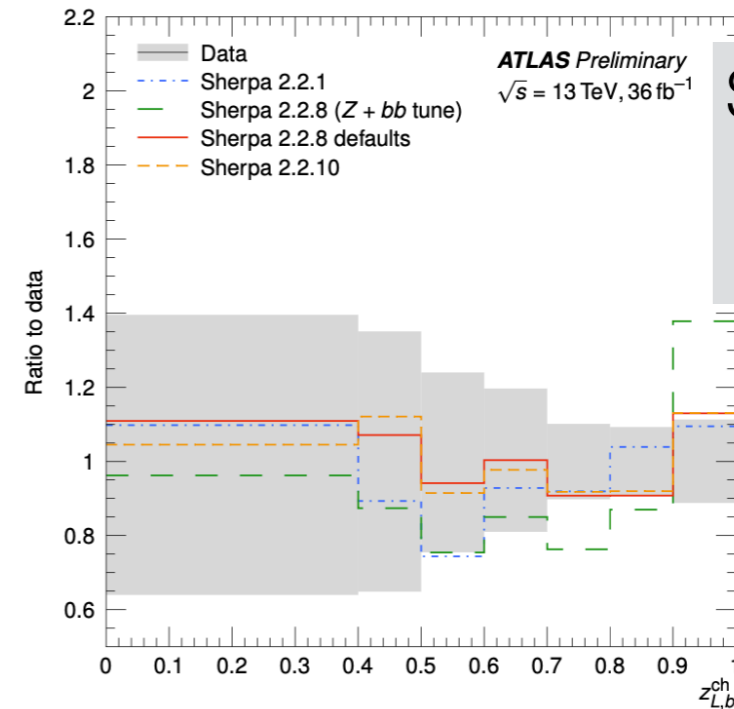
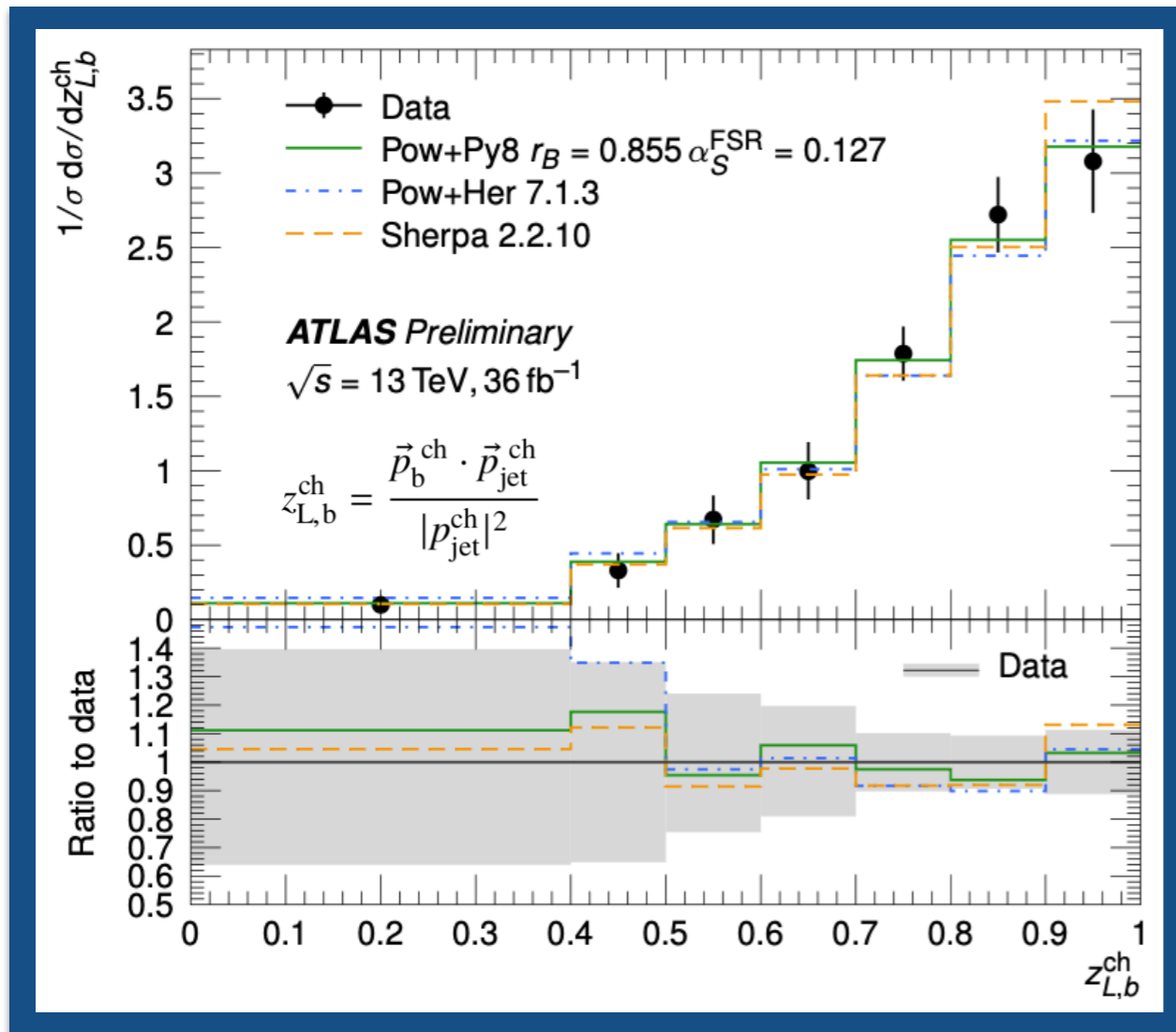
unfolding model is capable of describing observed data



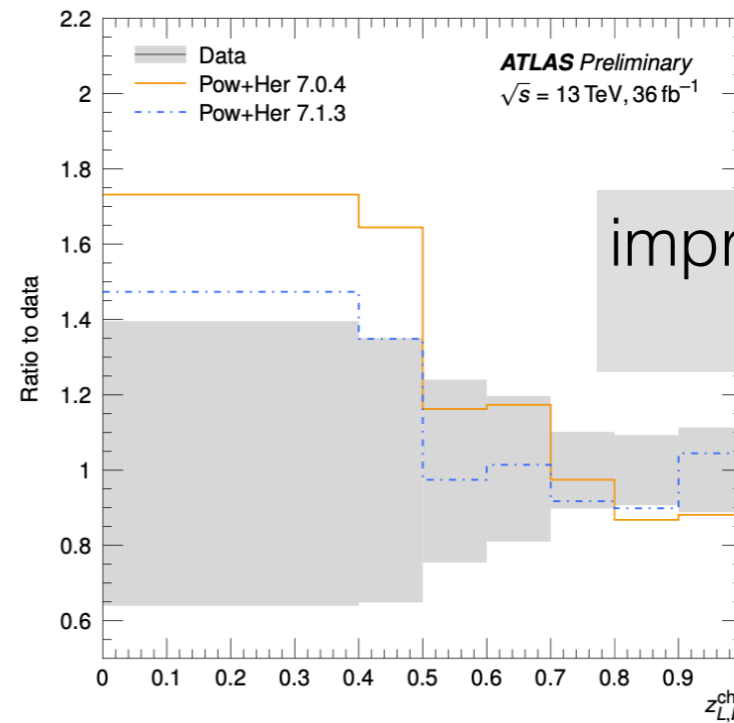
parameter anti-correlations cause reduction in total uncertainty

posterior uncertainty closely follows data stat uncertainty

fiducial data/MC comparisons I



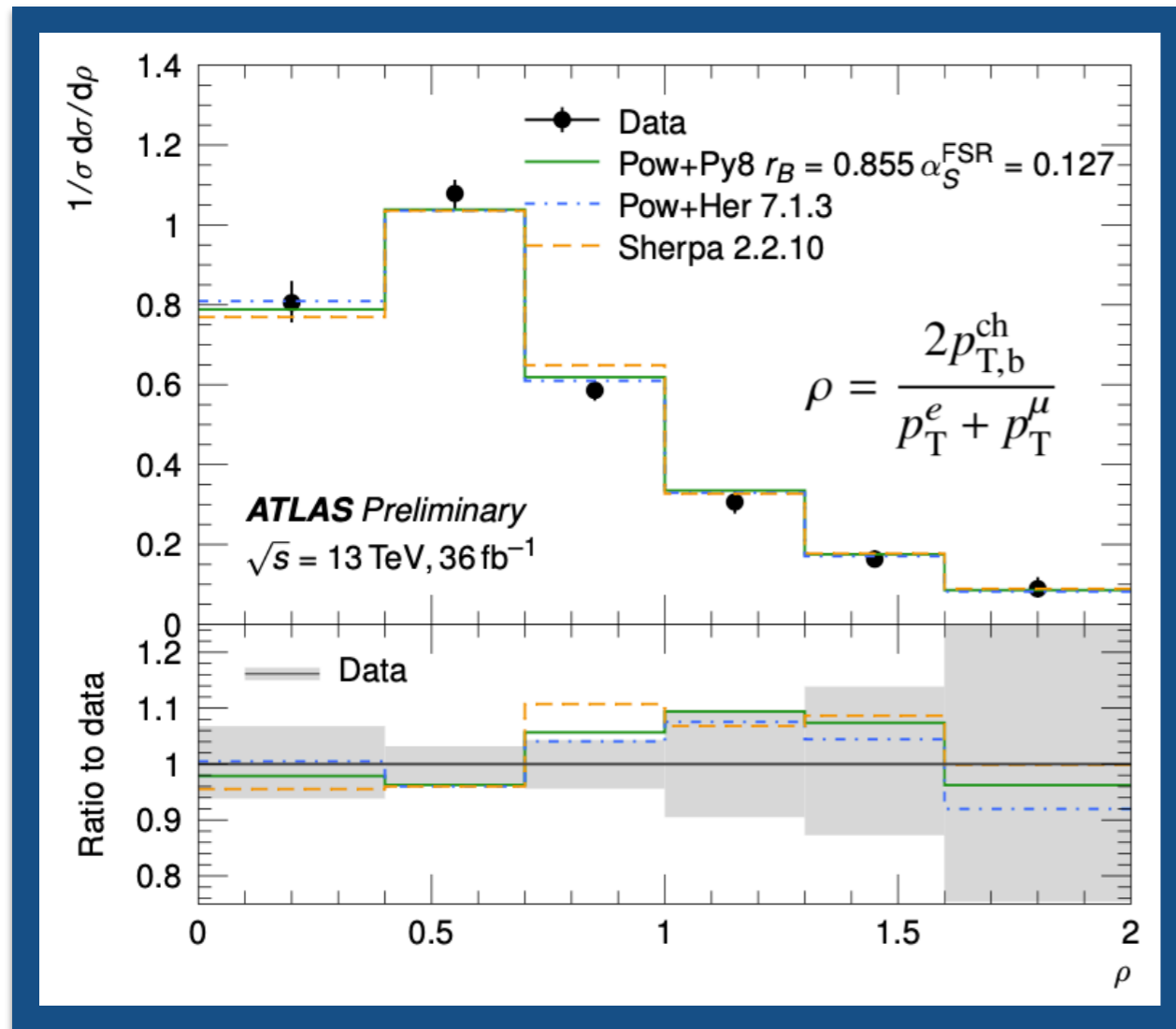
Sherpa: 2.2.10 with default tune looks good.



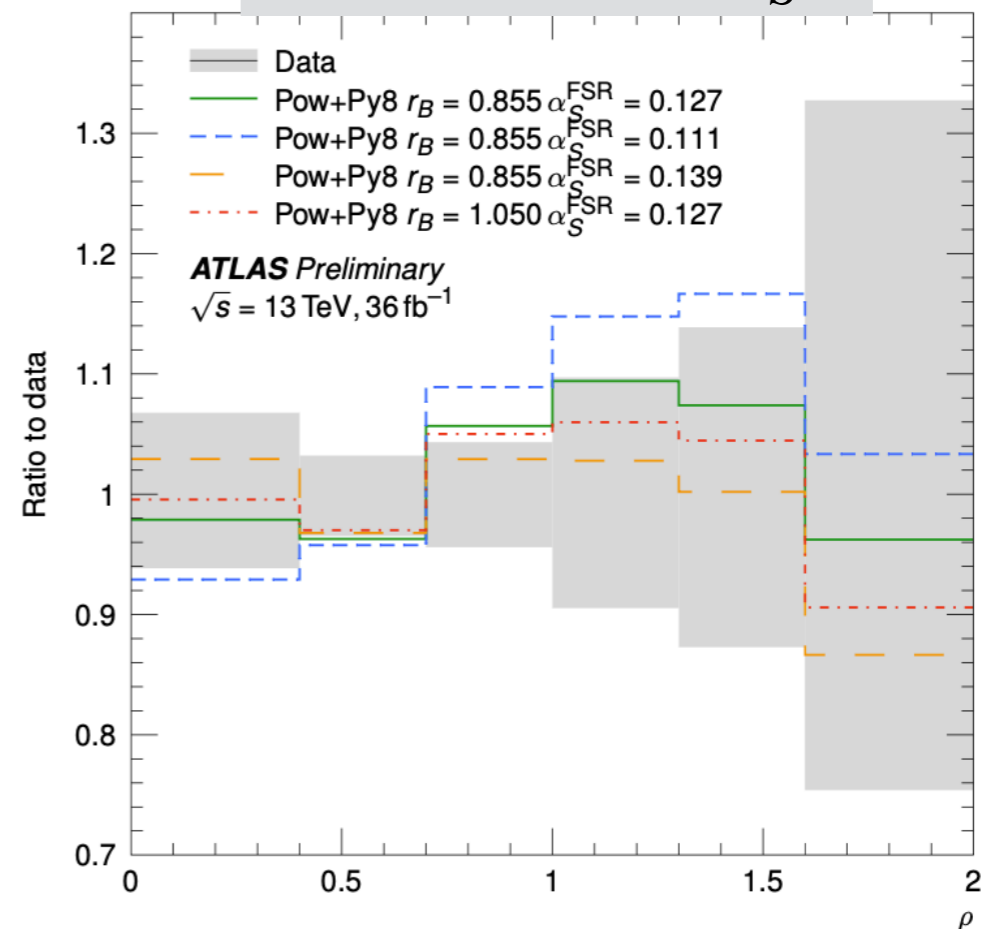
improvements in newer versions of H7

recent generator tunes to $e^+e^- \rightarrow Z \rightarrow bb$ do a reasonable job.

fiducial data/MC comparisons II

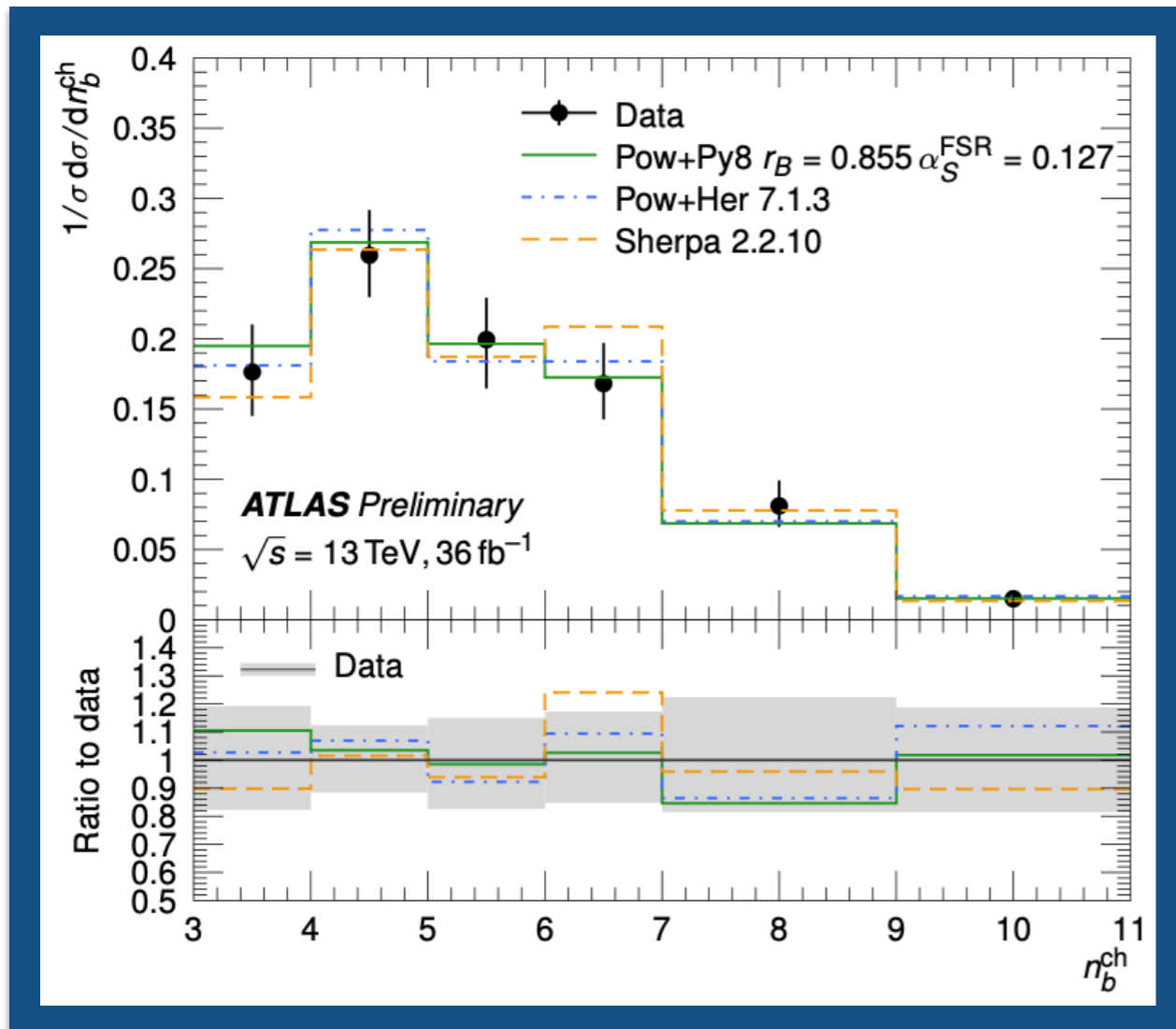


large impact from
 A14 α_S^{FSR} variations
 no sensitivity to α_S^{ISR}

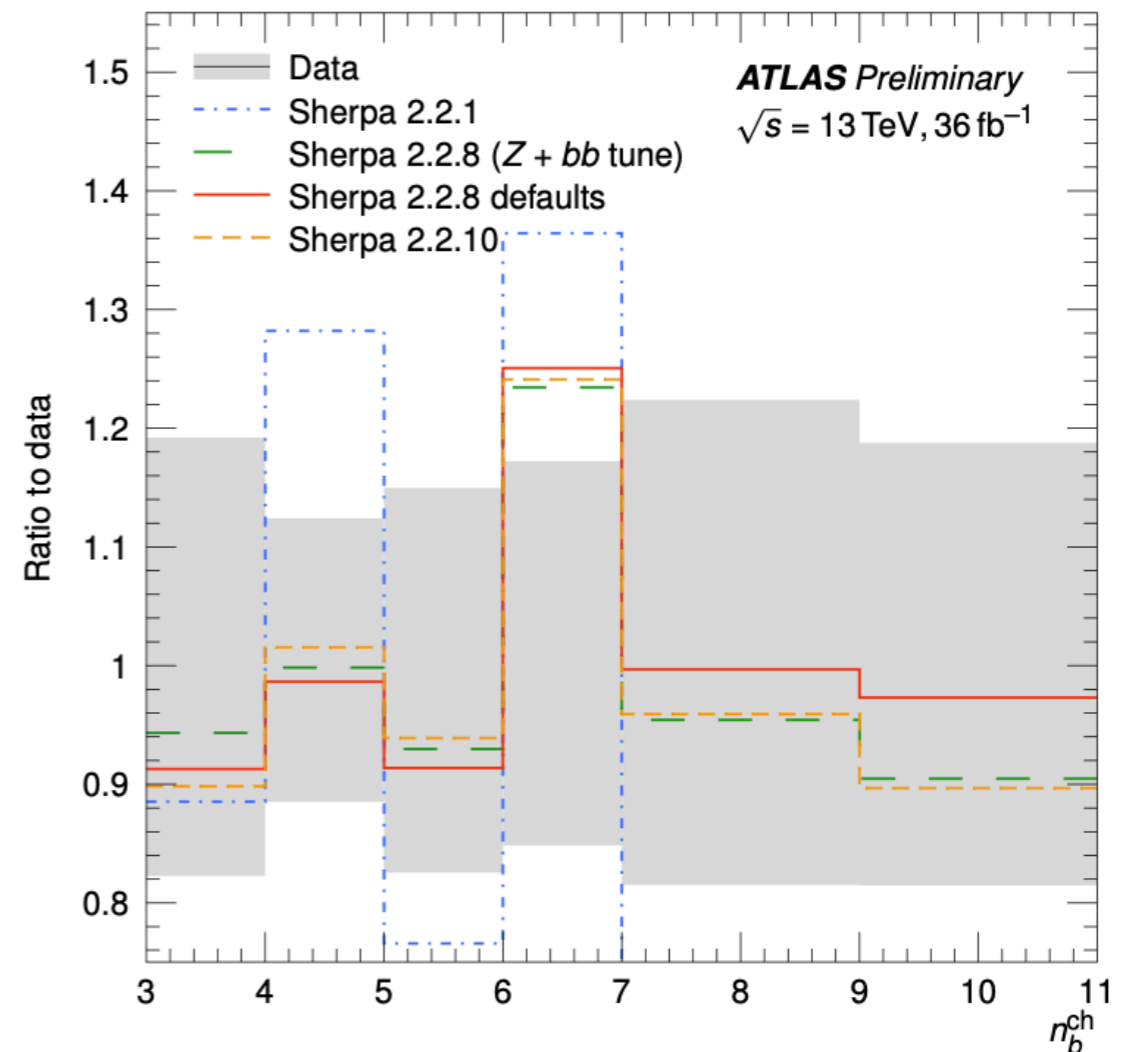


$\sim 1\sigma$ tensions
 with recent generators
 → further investigation.

fiducial data/MC comparisons III



poor choice of heavy baryon enhancement in ATLAS's Sherpa 2.2.1 setup has now been fixed.

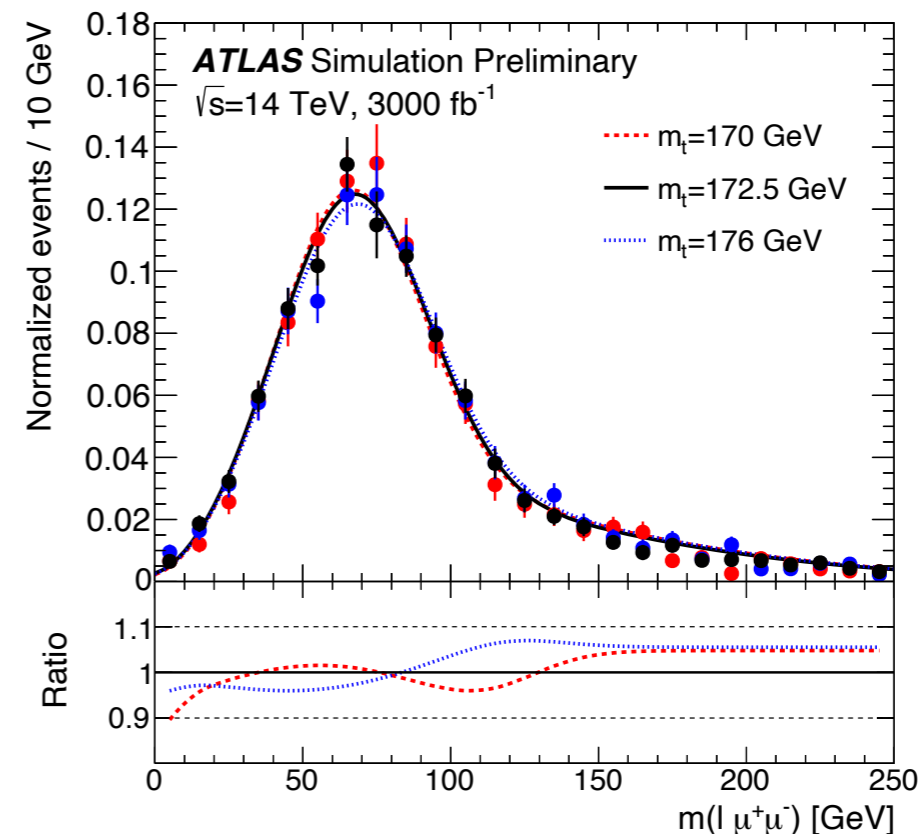
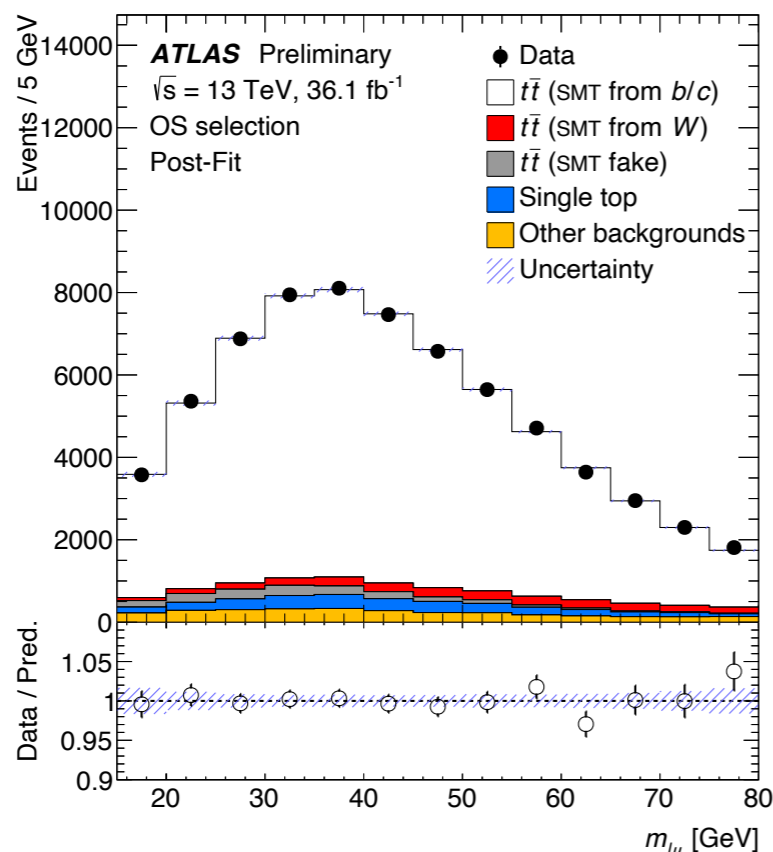


Py8/Her7 + EvtGen and Sherpa 2.2.10 all perform admirably.

impact: top mass

at our current precision, **generators well-tuned to $e^+e^- \rightarrow bb$ data yield good predictions in $t\bar{t}$ events and top-quark decays.**

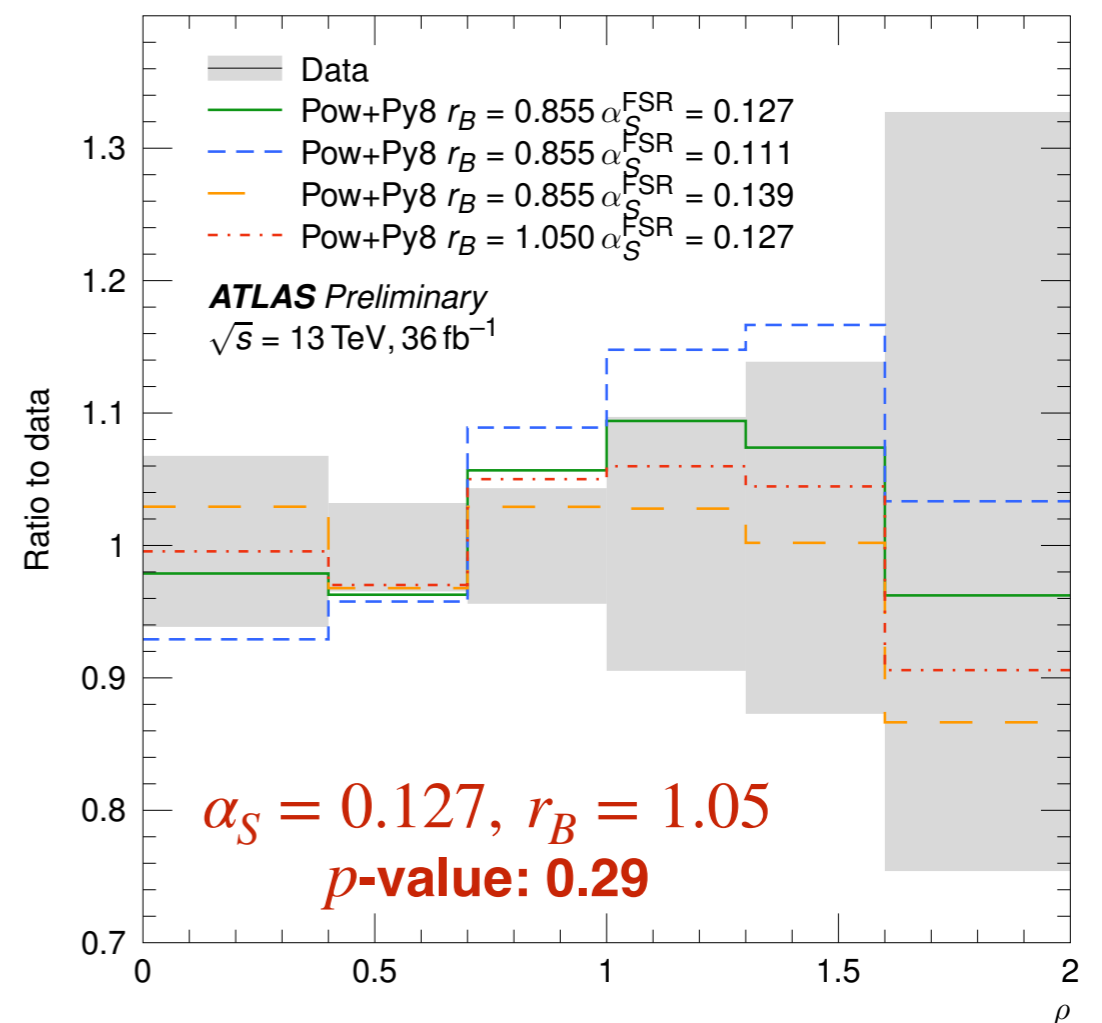
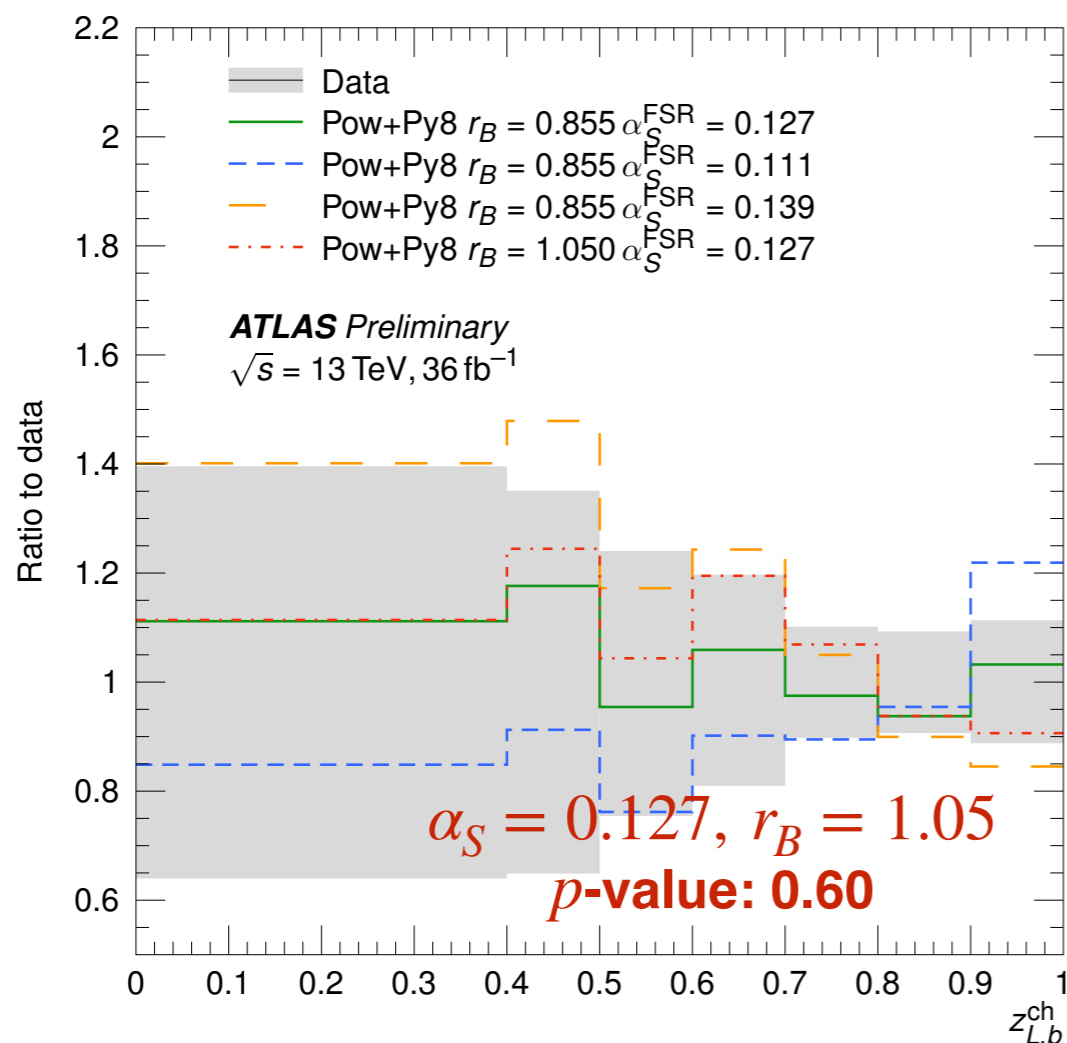
- we now have unfolded data to guide our choice of related uncertainties.
- precise b -fragmentation \rightarrow smaller uncertainties on the b -quark to b -jet transfer
- a better understanding of $t \rightarrow b^{\text{quark}} \rightarrow b^{\text{hadron}}$ crucial for fully-leptonic template mass extractions ($m_{\ell\mu}$ or $m_{\ell,J/\psi}$)
- joint fragmentation+mass measurements in leptonic modes **should be investigated**



impact: top mass

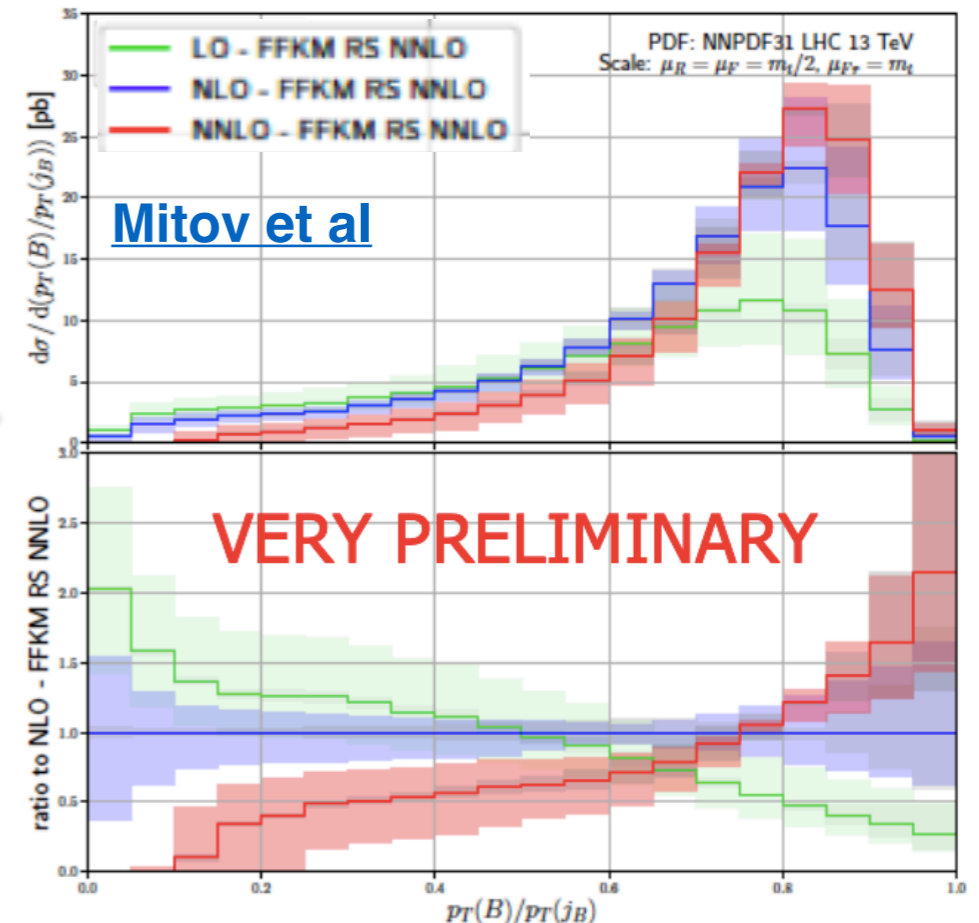
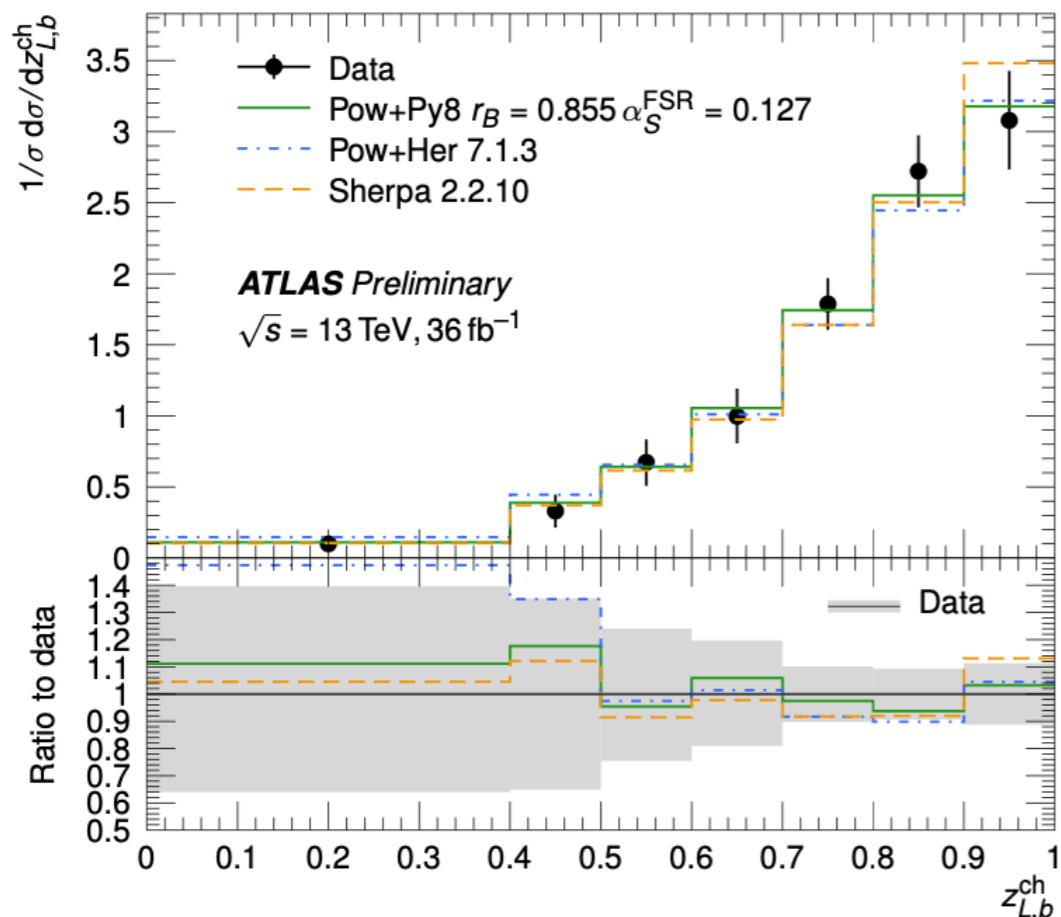
ATLAS retuned the Lund-Bowler r_B parameter to LEP data with the A14 α_S

- used in preliminary mass measurement ATLAS-CONF-2020-046
- this combination of $\alpha_S = 0.127$, $r_B = 1.05$ agrees with observed data to current precision. **(dash-dotted red curve)**



status of \sim analytic calculations

- Mitov et al. showed the first quasi-analytic calculations of similar distributions at TOP2020.
- will be very interesting to compare to unfolded data
- challenges for comparison:
 - calculations are for production of exclusive B-hadron species
 - jets are defined at parton-level in calculation
 - ATLAS measurement is to stable, charged particles



summary and outlook

- several differential observables sensitive to b -quark fragmentation have been measured in $t\bar{t} \rightarrow e\mu b\bar{b}$ events with the 2015+2016 ATLAS data
- comparisons to modern MC generator predictions are made
- we do not explicitly tune MC parameters in this paper
 - but these measurements will be used in future tunes.
- we will provide Rivet routine + HEPData entries for the unfolded data.
- we've identified several "sub-optimal" tunes that were used previously: hopefully we can prevent that in the future.
- many improvements possible: more precise measurements may be necessary for best "direct" top-quark mass extraction
- first ~analytic calculations of b -quark fragmentation in top-quark decays are on the horizon for comparison!

bonus

detector-level yields: predicted and observed

	events with $e\mu jj$ (≥ 1 b -tag)	probe-jets
process	predicted yields	
fiducial $t\bar{t}$	–	44000 \pm 12000
non-fiducial $t\bar{t}$	–	6700 \pm 1900
total $t\bar{t}$	76000 \pm 18000	51000 \pm 12000
single top	4400 \pm 1500	1580 \pm 600
Z+jets	125 \pm 45	13.0 \pm 5.1
diboson	90 \pm 34	9.7 \pm 3.9
total non- $t\bar{t}$	4600 \pm 1600	1600 \pm 600
b -jets	–	52200 \pm 12000
c -jets	–	180 \pm 60
other jets	–	250 \pm 70
total prediction	81000 \pm 18000	53000 \pm 12000
	observed yields	
data	88511	57476

$t\bar{t}$ dominant

high expected
 b -jet purity

7.5% data / MC
discrepancy

Table 1: Summary of expected and observed event and probe-jet yields.

unfolding procedure: background and response matrices

- to calculate the background predictions and response matrices at a given point in parameter space, we take linear combinations of the background predictions for nominal and the 1-sigma variation corresponding to a systematic uncertainty:

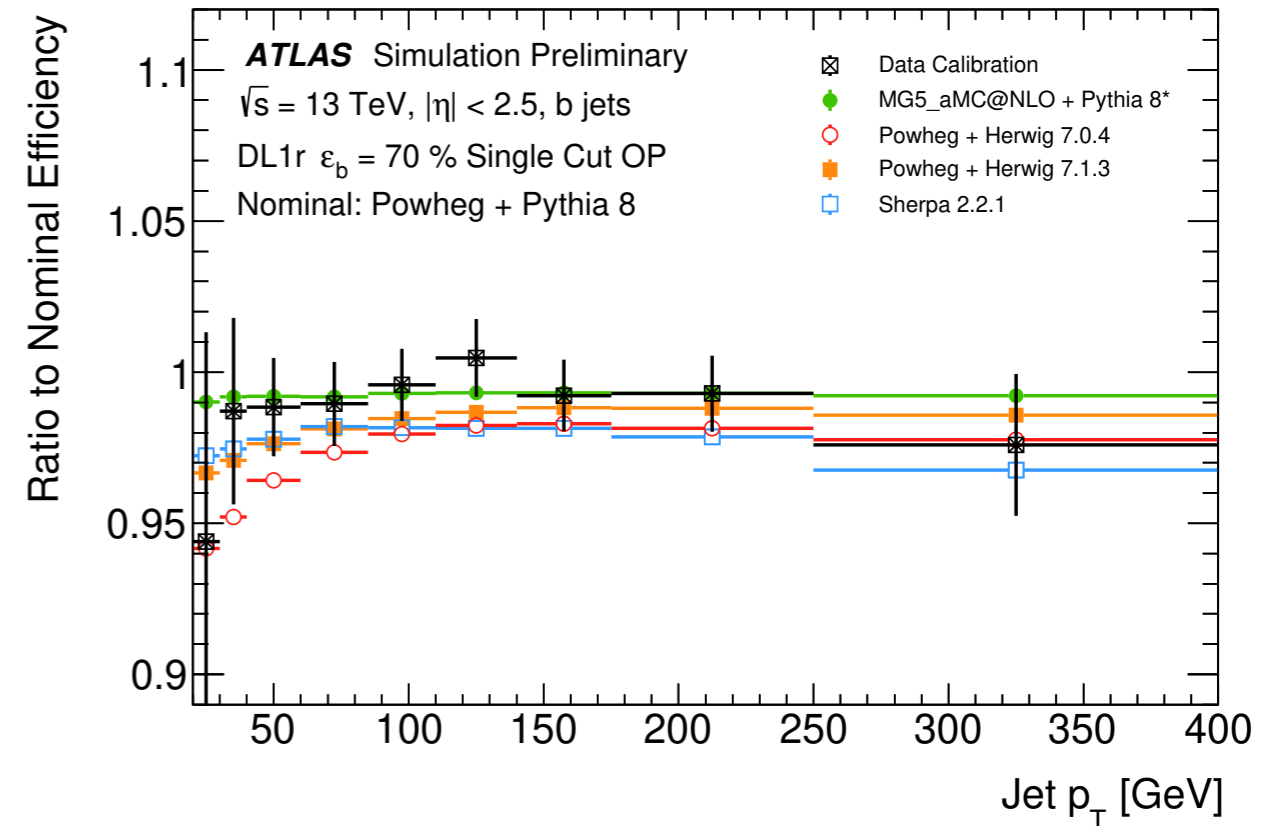
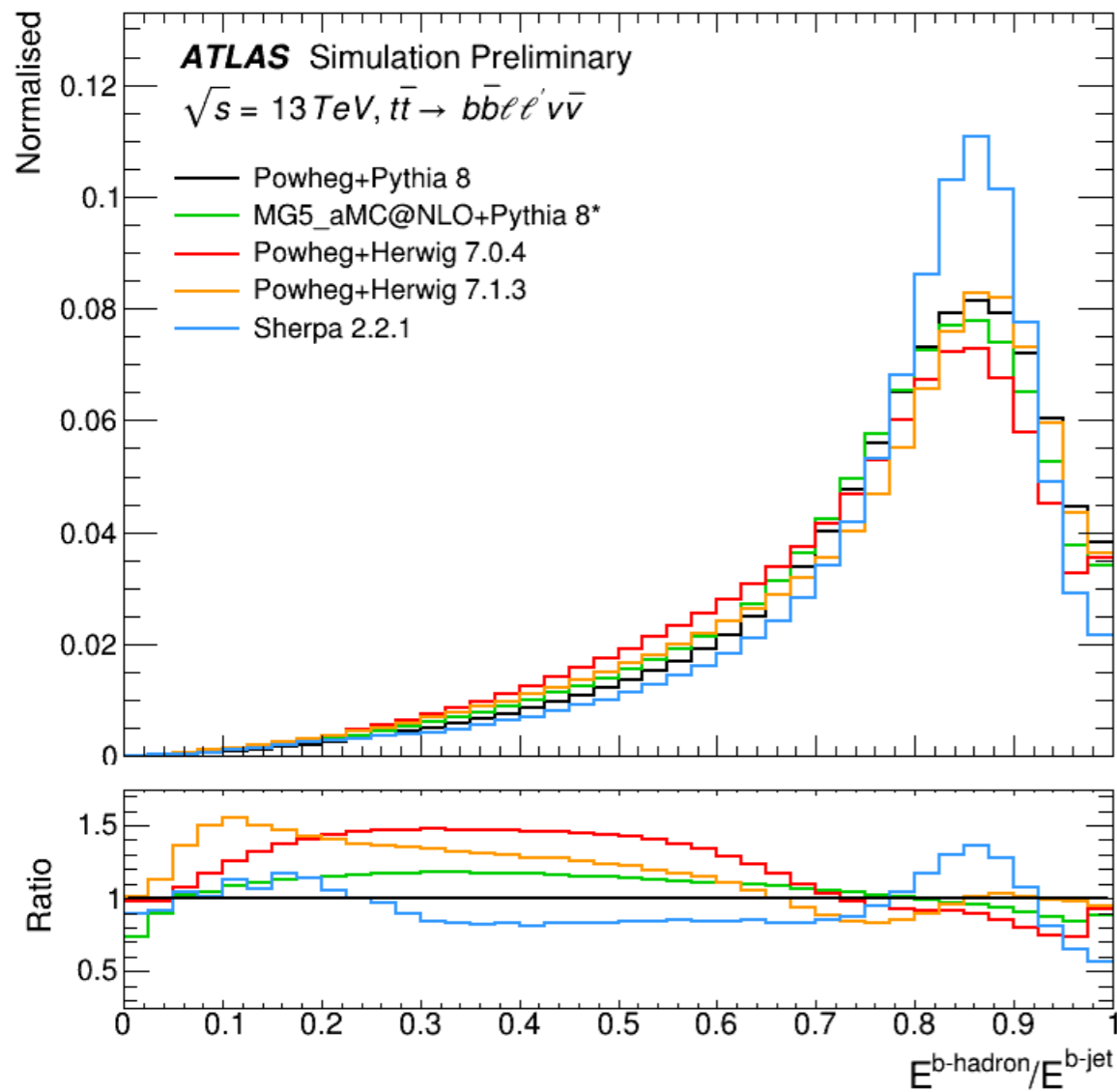
$$M_{ij}^{p \rightarrow d}(\vec{p}) = M_{ij,0}^{p \rightarrow d} + \sum_{k \in \text{systematics}} p_k * (M_{ij,k}^{p \rightarrow d} - M_{ij,0}^{p \rightarrow d})$$

$$b_i(\vec{p}) = b_{i,0} + \sum_{k \in \text{systematics}} p_k * (b_{i,k} - b_{i,0})$$

where the \vec{p} are the nuisance parameters at a particular point in parameter space.

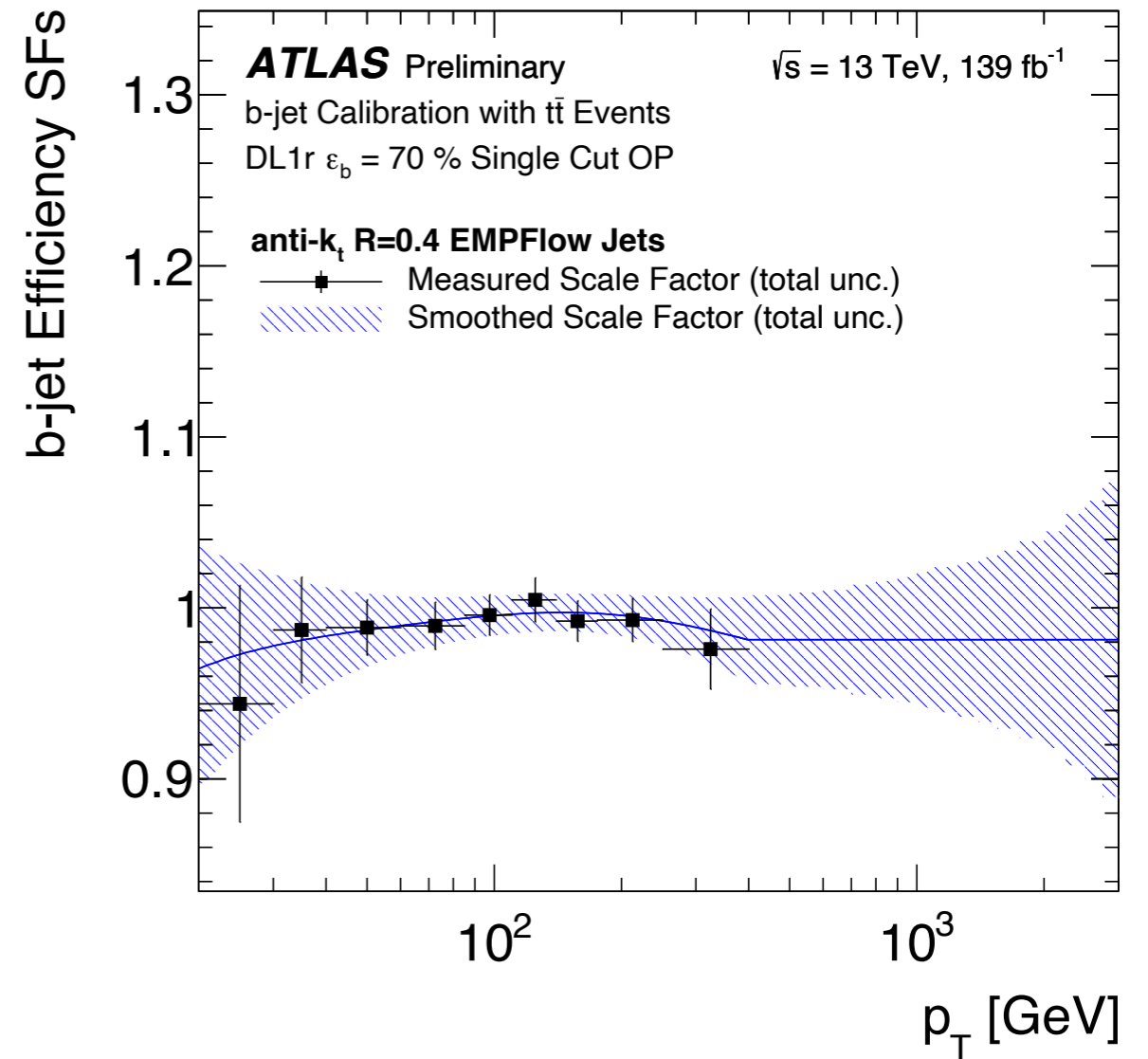
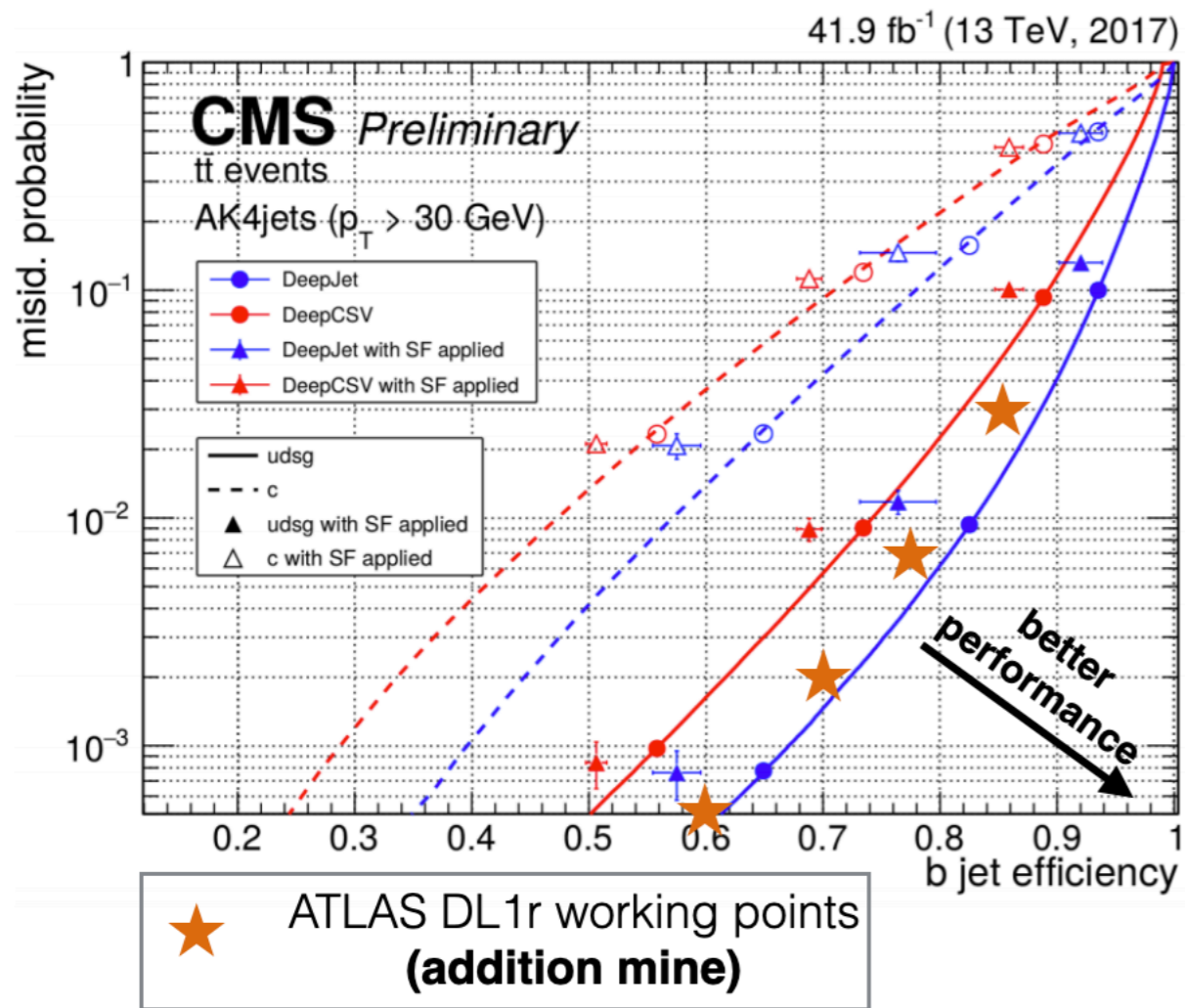
- in this way we can coherently and smoothly vary the response matrices and background predictions corresponding to systematic variations.
- we're wandering around the space of possible response matrices and background predictions available given our understanding of the predictions.

impact on b -tagging

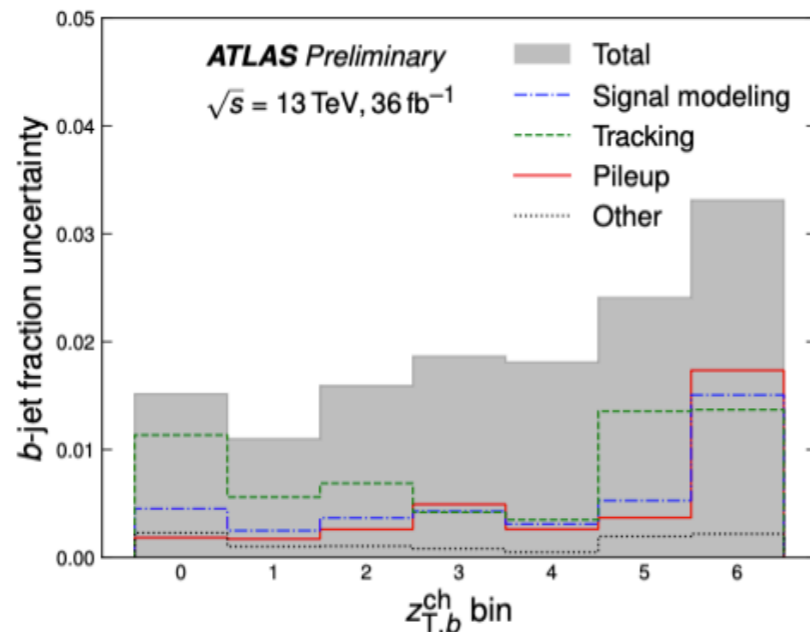


ATLAS-PHYS-PUB-2020-009

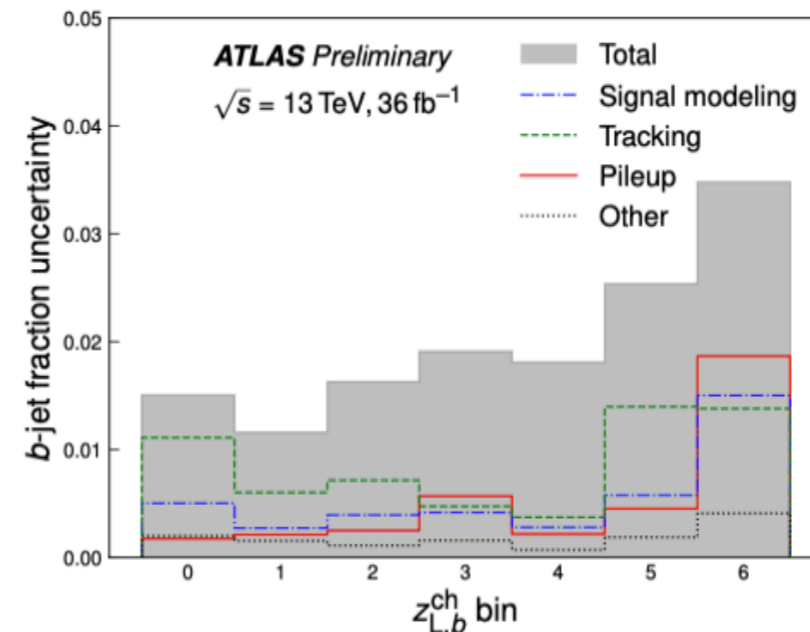
b -tagging SFs



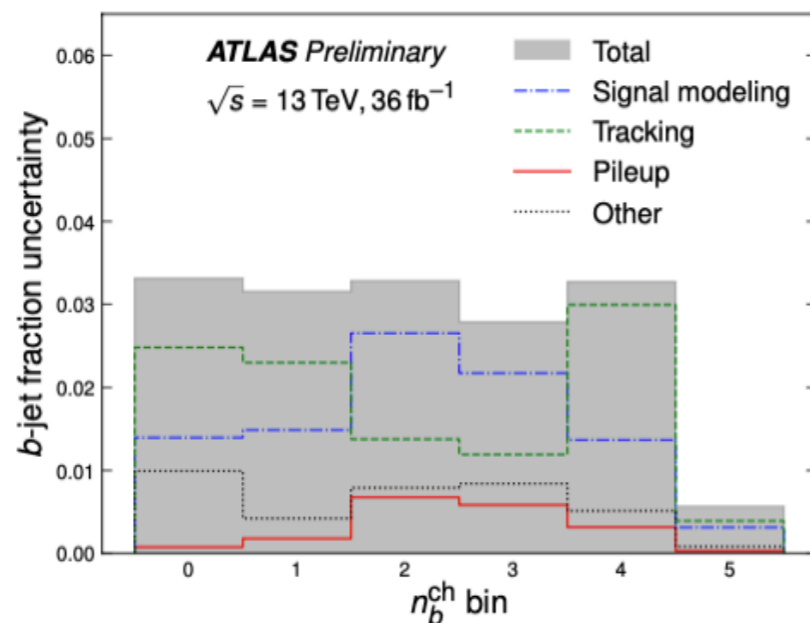
uncertainty breakdowns



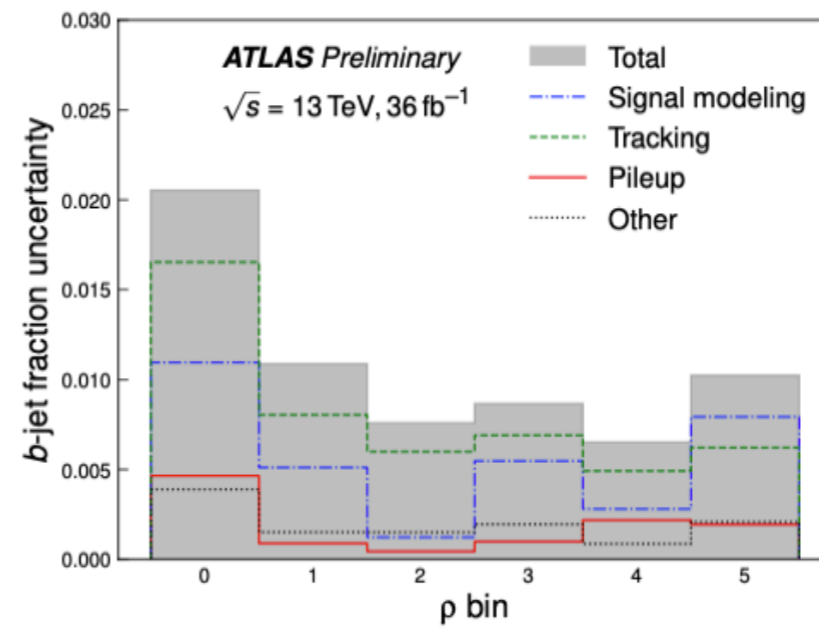
(a) $z_{T,b}^{\text{ch}}$



(b) $z_{L,b}^{\text{ch}}$



(c) n_b^{ch}



(d) ρ

p -values

generator configuration	$z_{T,b}^{\text{ch}}$	$z_{L,b}^{\text{ch}}$	n_b^{ch}	ρ
POWHEG+PYTHIA8 A14 $r_B = 0.855$ $\alpha_S^{\text{FSR}} = 0.127$	0.50	0.98	0.94	0.20
POWHEG+PYTHIA8 A14 $r_B = 0.855$ $\alpha_S^{\text{FSR}} = 0.139$	0.13	0.37	0.95	0.31
POWHEG+PYTHIA8 A14 $r_B = 0.855$ $\alpha_S^{\text{FSR}} = 0.111$	0.16	0.62	0.95	0.07
POWHEG+PYTHIA8 A14 $r_B = 1.050$ $\alpha_S^{\text{FSR}} = 0.127$	0.18	0.60	0.94	0.29
POWHEG+HERWIG7.0.4	0.22	0.50	0.95	0.12
POWHEG+HERWIG7.1.3	0.42	0.84	0.92	0.18
SHERPA2.2.1	0.14	0.45	0.01	0.52
SHERPA2.2.8 ($Z + bb$ tune)	0.004	0.02	0.45	0.24
SHERPA2.2.8	0.32	0.79	0.49	0.11
SHERPA2.2.10	0.27	0.84	0.32	0.08

Table 3: Summary of p -values for various MC generator configurations based on the observed data distributions