

# Top modelling and its uncertainties in Sherpa

Frank Siegert

LHC Top WG Meeting, 24 November 2020

Emmy  
Noether-  
Programm

Deutsche  
Forschungsgemeinschaft

DFG




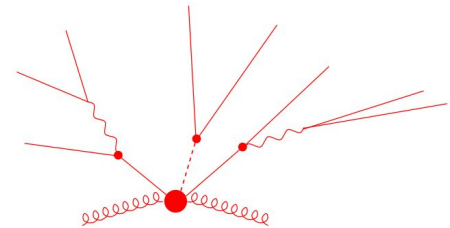
- A parton-shower Monte Carlo is not a fixed-order prediction
  - **It is much more powerful!**
  - **And at the same time much more ambiguous!**

**Ambiguities = Uncertainties**

(and in addition there can be bugs of course)

Let's review how they are addressed in Sherpa ...

- A parton-shower Monte Carlo is not a fixed-order prediction
  - **It is much more powerful!**
  - **And at the same time much more ambiguous!**
- Typical sources of ~~trouble~~ **ambiguities**:
  - Hard scattering
    - » Limited perturbative accuracy
      - ⇒ **ambiguity in scale and PDF choices** 
    - » Narrow-width approximation instead of full offshell WbWb
      - Diagram overlap between tt and tWb
        - ⇒ **ambiguity in overlap removal**
      - Spin correlations between production and decay MEs in the chain
        - ⇒ **ambiguity in polarisation treatment**
          - particularly tricky for tau decays, as they can be hadronic!
    - » NLO EW Combination of NLO QCD and NLO EW corrections
      - ⇒ **ambiguity in combination between NLO QCD and NLO EW**
    - » Multi-leg merging of ME & PS
      - ⇒ **ambiguity in transition**



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On-the-fly variations of scales and PDFs [[1606.08753](#)]

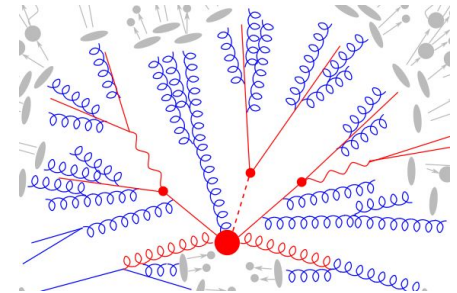
Spin correlations (→ later)



Additive and multiplicative schemes [[2005.12128](#)]

Variation of merging scale  $Q_{\text{cut}}$

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- ▶ Typical sources of ~~trouble~~ **ambiguities**:
  - Parton shower: QCD corrections with three major ambiguities
    - » Functional form of splitting kernels  
(approximation of real-emission MEs)  
⇒ **ambiguity which (finite) pieces to keep**
    - » Kinematics recoil  
(how to construct  $1 \rightarrow 2$  splittings with  $m=0$  away from collinear limit)  
⇒ **ambiguity where to distribute recoil for momentum conservation**
    - » Evolution variable  
(direction in which logs are resummed)  
⇒ **ambiguity what “from hard to soft” means exactly**
    - » Additionally many ambiguities for treatment of quark masses in the above!



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Switch between 2 different recoil schemes [[0912.3501](#)]

- » Kinematics recoil

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- » Evolution variable

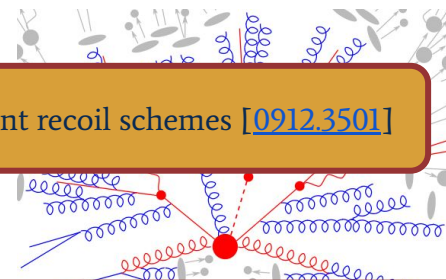
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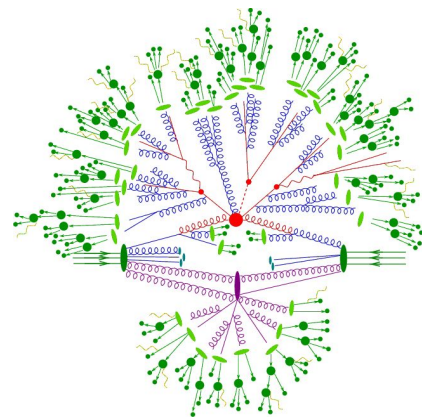
Switch between different evolution schemes ( $\rightarrow$  later).  
Choice of shower starting scale  $\mu_Q$ .

- » Additionally many ambiguities for treatment of quark masses in the above!

Only probed to small extent by on-the-fly variations of scales and PDFs! [[1606.08753](#)]



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  - Hadronisation: Soft QCD modelling without “first principles”
    - » B-hadron production from partons  
⇒ **ambiguity of flavours formed** (e.g. meson or baryon,  $B^*$  or  $B$ , ...)
  - Hadron decays: Effective field theories for heavy-flavour decays
    - » B-hadron decays  
⇒ **ambiguity of decay matrix elements** (form factor models)



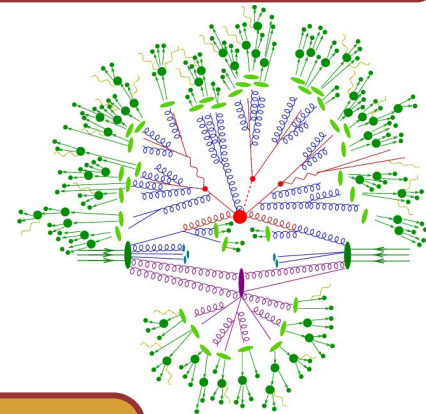
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Interface to alternative hadronisation from Pythia6



Several form factor models for important B and D decays.  
Branching ratio defaults from PDG, but can be varied.



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Let's look at some typical Sherpa top setups  
and their modelling & uncertainties ...

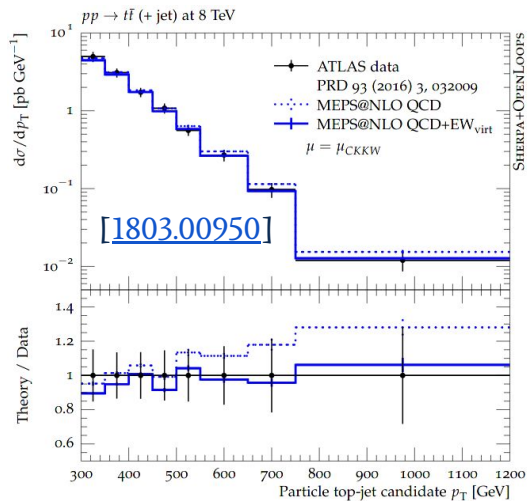
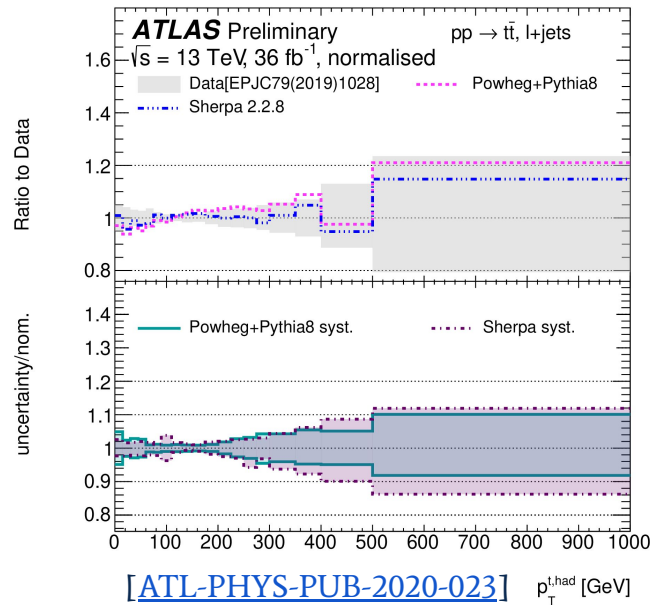
▶ Multileg NLO QCD merging Sherpa+OpenLoops [[1402.6293](#)]

- $tt + ttj$  (+  $ttjj$ ) @ NLO
- $(ttjj + )$   $ttjjj + ttjjjj$  @ LO

▶ Recent ATLAS & CMS studies (→ just shown by [Marino](#)):

Good modelling of  $tt$ +jets & interesting uncertainty comparison:

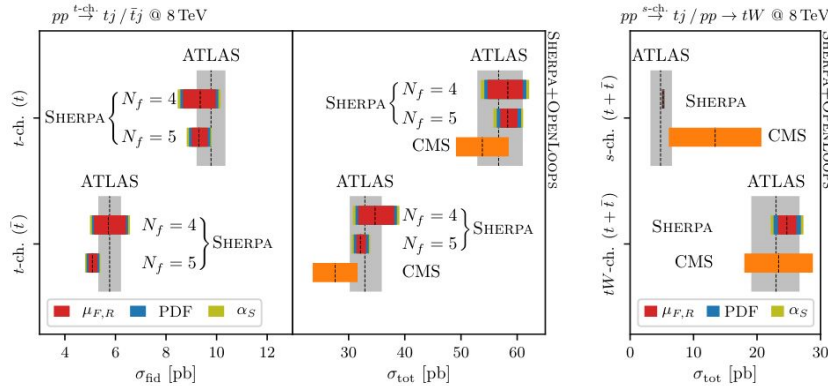
- PH+Py8:  $\mu^{\text{ME}}$ ,  $\text{PDF}^{\text{ME}}$ ,  $\text{hdamp}$ ,  $\mu^{\text{PS,ISR}}$ ,  $\mu^{\text{PS,FSR}}$ , PS recoil
  - Sherpa  $tt$  MEPS@NLO:  $\mu^{\text{ME+PS}}$ ,  $\text{PDF}^{\text{ME+PS}}$ ,  $Q_{\text{cut}}$ ,  $\alpha_S$ ,  $\mu_Q$ , PS recoil
- Similar order of magnitude, WIP to understand differences.



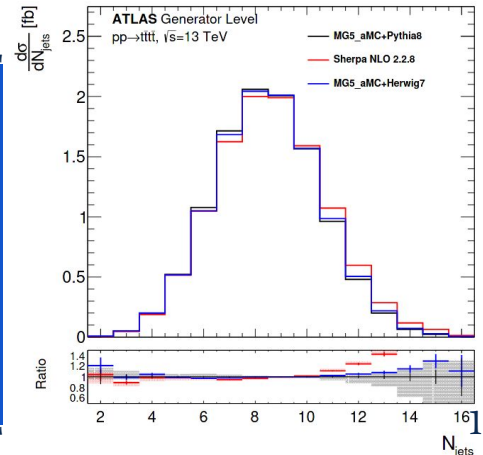
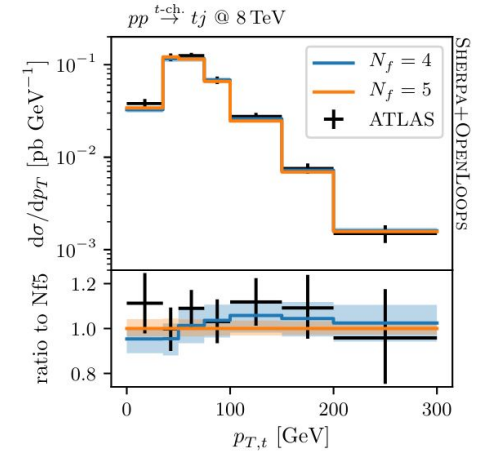
▶ New: NLO EW corrections directly in MC sample

- using EWvirt approximation in MEPS@NLO
- typical EW Sudakov effects relevant in high  $p_T$  regions
- now available “for free” as on-the-fly weights in samples
  - » since 2.2.9 even in additive and multiplicative scheme
  - useful to define uncertainty!

- Single-top production [\[1711.02568\]](#)
  - NLO+PS setup with OpenLoops virtual MEs
  - For t-channel, s-channel, tW using diagram removal approach

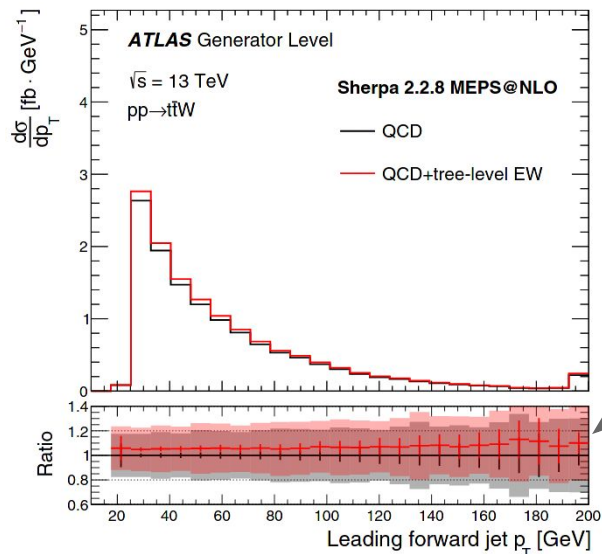


- 4-top production
  - NLO+PS samples with OpenLoops virtual MEs
  - Recent ATLAS comparisons with aMC@NLO
  - WIP: include NLO EW (and subleading trees) through EWvirt mechanism

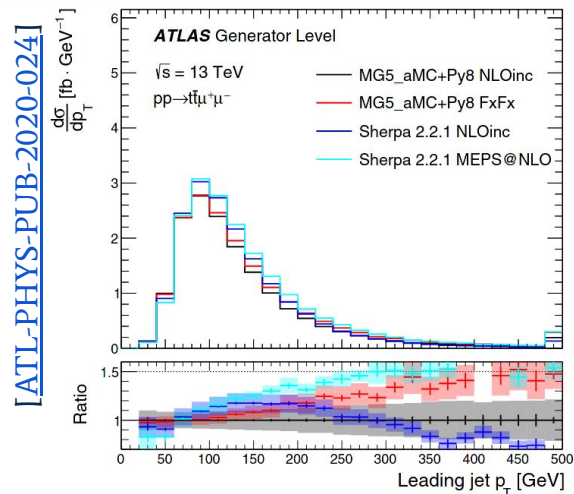


[\[ATL-PHYS-PUB-2020-024\]](#)

- ▶ ttV+jets and ttll+jets with Sherpa+OpenLoops
  - MEPS@NLO tt{V,ll} + 0,1j@NLO + 2,3j@LO
  - Merged setups predict higher rates and harder spectra
  - ttV+1j@NLO helps capture large parts of NNLO effects relevant for global normalisation

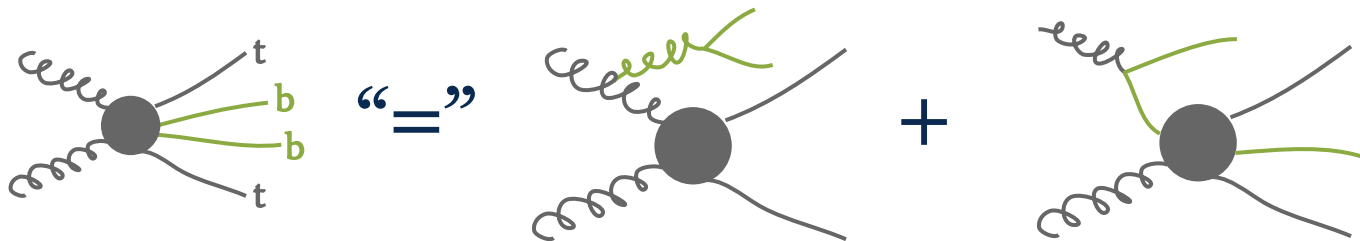


- ▶ NLO EW weights available through EWvirt approximation
  - particularly interesting because of subleading EW tree contributions
    - » detailed studies ongoing to compare to fixed-order literature
    - » EWvirt approach makes this available differentially in MC sample!



[ATL-PHYS-PUB-2020-024]

- 4FS NLO+PS setup for ttbb with massive b-quarks in  $2 \rightarrow 4$  matrix element



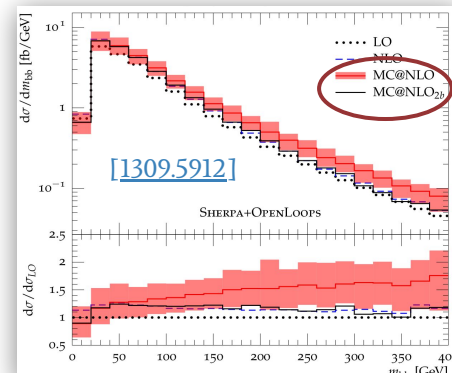
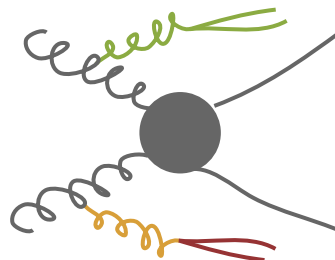
Final state  $g \rightarrow bb$  **dominant**

- massive b's  $\rightarrow$  no (jet) cuts!
- collinear  $g \rightarrow bb$  produced in ME

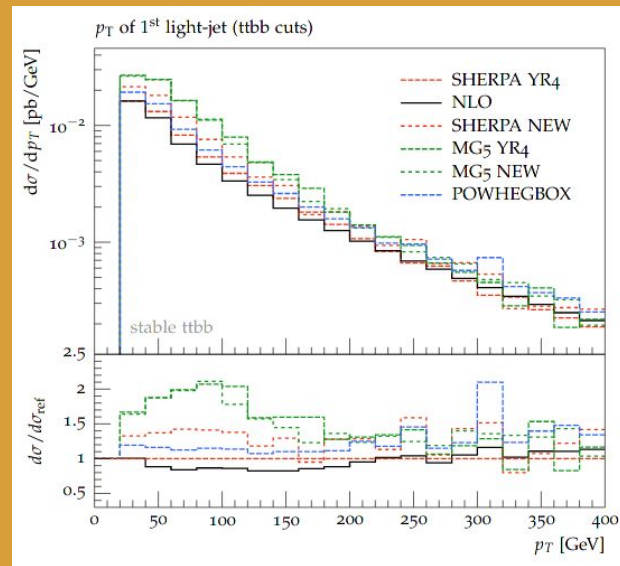
No initial state b in MEs

- 4FS PDFs
- IS  $g \rightarrow bb$  in ME

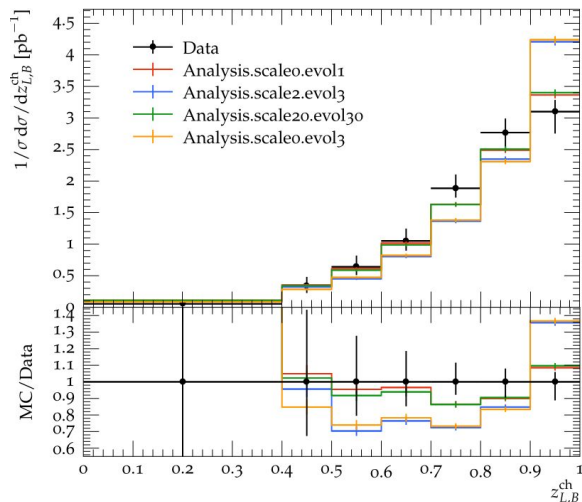
- Matched to parton shower for additional emissions
  - “**double-splitting**” contribution becomes relevant!



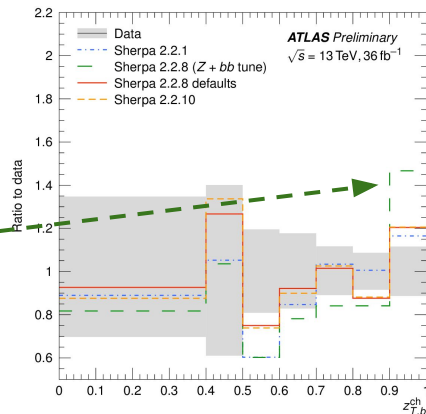
Many ambiguities for ttbb  
 discussed by [Andrea](#) yesterday  
 and [studied recently](#) in detail  
 in LHC Higgs XS WG  
 → watch for a summary soon!



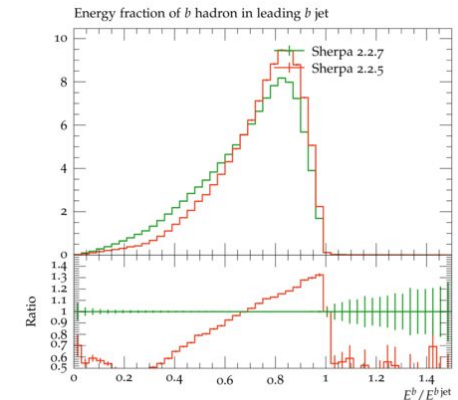
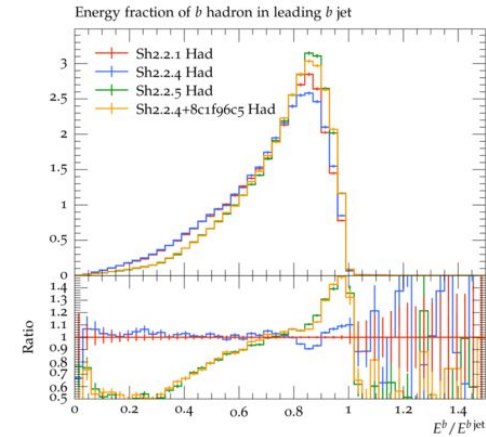
- ▶ New schemes for evolution variable in HF splittings
  - CSS\_EVOLUTION\_SCHEME=3 improves  $g \rightarrow bb$  splitting by appropriate b-mass term in evolution variable
  - also applied to  $b \rightarrow bg$  splitting, although not necessary there
  - worse data/MC agreement in [b-fragmentation analysis](#)
    - » dominated by  $b \rightarrow bg$  splittings



- ▶ Future in Sherpa  $\geq 2.2.9$ : **hybrid scheme (“30”)**
  - $b \rightarrow bg$  retains original evolution
    - » b-fragmentation in tt improved
  - relevant mass term included only in  $g \rightarrow bb$ 
    - » relevant e.g. for Zbb



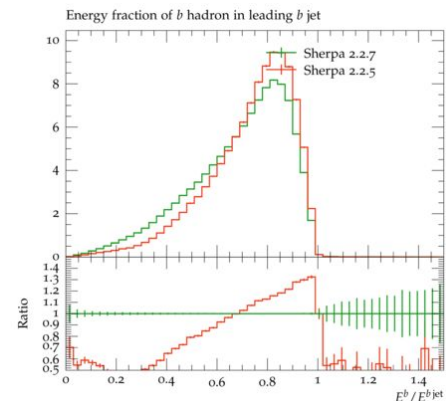
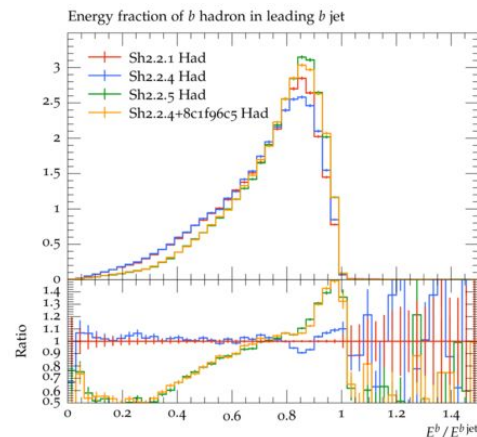
- ▶ Example of impact from shower splitting functions
  - Sherpa 2.2.5 contained bugfix for showering in coloured decays based on developments in [[1709.08615](#)]
  - bugfix contained wrong definition of mass term in shower splitting function
  - Large impact on fragmentation function
  - Details in <https://gitlab.com/sherpa-team/sherpa/-/issues/176>
  
- ▶ First found in Sherpa 2.2.6 ATLAS validation, fixed in Sherpa 2.2.7:





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- ▶ First found in Sherpa 2.2.6 ATLAS validation, fixed in Sherpa 2.2.7:

A bug, not an ambiguity.  
 But difficult to disentangle!



- Recent efforts to find physics compromises which reduce negative weight fraction
- Thorough validation against full options to ensure physics is unchanged

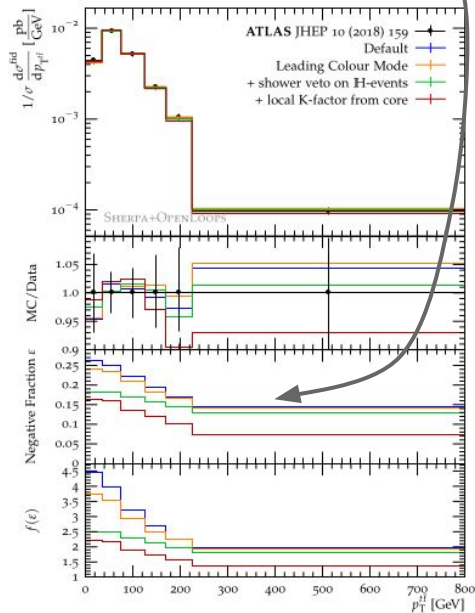
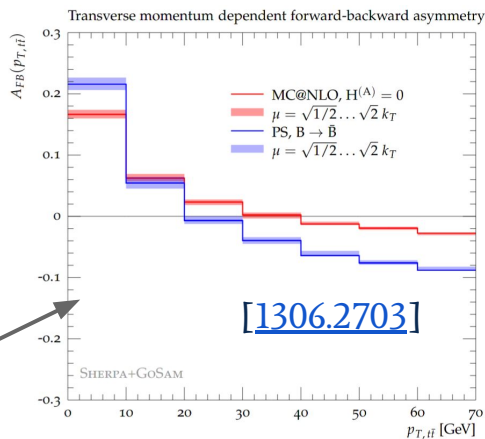
- Largest effect disabling full-colour mode in matched shower evolution

- use simple shower kernels instead of exact subtraction dipoles
- generally no significant impact, but careful with sensitive observables, e.g. forward-backward-asymmetry!

- Subleading effect: shower veto on H-events

- basic idea similar to part of [MC@NLO  \$\Delta\$ -matching](#)
- natural within merged sample  $\rightarrow$  no counter terms/folding necessary

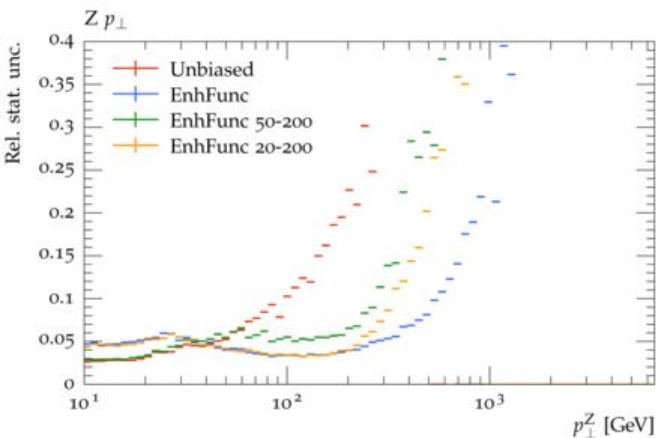
tt+0,lj@NLO+2,3j@LO	Negative Weight Fraction
Default	24.8%
Leading Colour Mode	18.7%
+ shower veto on H-events	14.5%
+ local K-factor from core	12.6%



- ▶ Experiments often slice samples (e.g. in  $p_{T\perp}$ ) to populate non-bulk regions
- ▶ Alternative: phase-space biasing of (otherwise) unweighted evts
  - (Re-)Implemented in Sherpa 2.2.8 in two variants (demonstrated here in Z+jets):

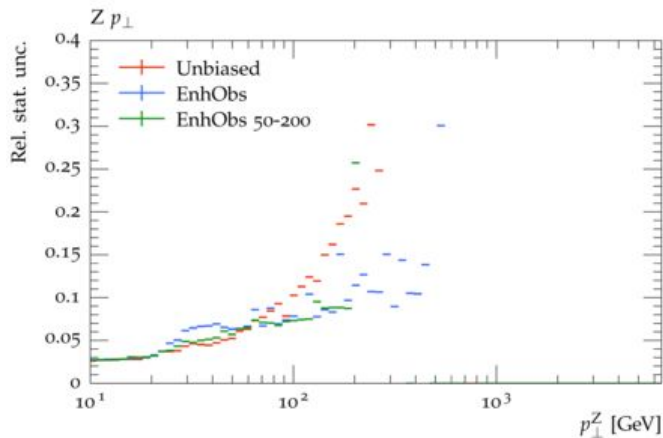
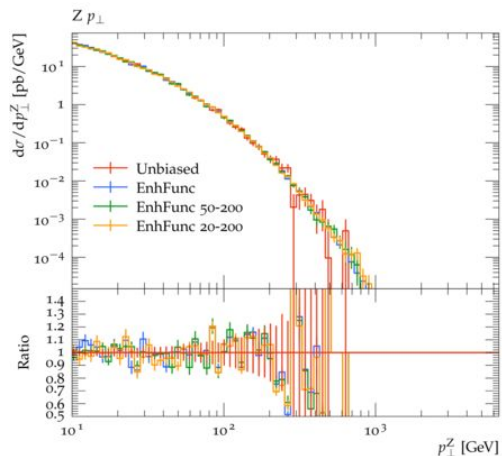
## Multiplicative function

“Enhance\_Function”



## Flattening in given observable

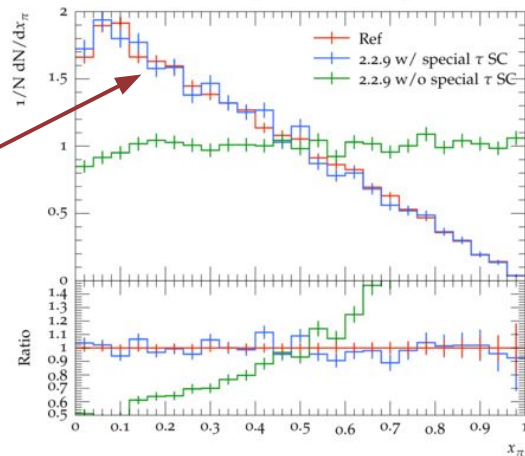
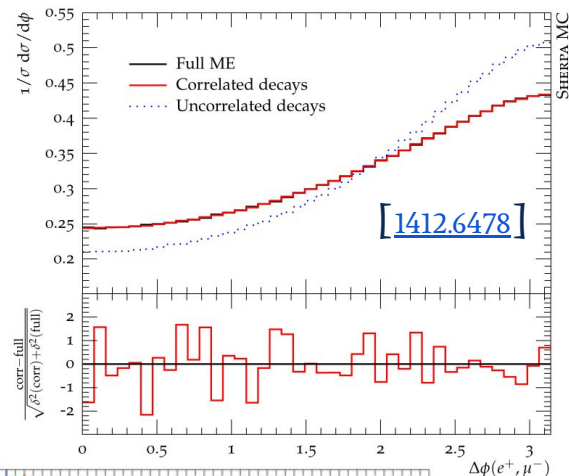
“Enhance\_Observable”



- What are spin correlations?

$$\mathcal{M} \sim j_1^\mu \left[ g_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right] j_2^\nu = \sum_{\lambda} \underbrace{j_1^\mu \varepsilon_\mu^*(\lambda)}_{\mathcal{M}_{\text{prod}}(\lambda)} \underbrace{\varepsilon_\nu(\lambda) j_2^\nu}_{\mathcal{M}_{\text{dec}}(\lambda)}$$

- Spin correlations for top and W decays were always fully taken into account in Sherpa
- For taus produced in W decays it was more difficult
  - separate decay chains, because taus decay hadronically
    - spin correlation algorithm only works recursively
  - dedicated [algorithm in Sherpa >=2.2.9](#) recovers polarisation of taus produced in hard decays
    - $\tau \rightarrow \pi \nu$  energy fraction correctly reproduced
    - important for analyses:
      - $\tau \rightarrow$  lepton angular correlations correct



- MC event generators are powerful but ambiguous
  - Even more so with complex features like NLO multi-leg merging or EW corrections→ Requires careful assessment of systematic uncertainties
- Sherpa is available for wide variety of top physics with state-of-the-art precision
  - Many recent features particularly for LHC top production processes.



Sherpas and tops go together quite naturally :)

**Thanks for your attention!**