



Theoretical study of field emission from dielectric coated surfaces

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Introduction

- Electron field emission plays an essential role in a wide range of applications, including
 - Electron microscopes
 - X-ray sources
 - High current cathodes
 - High power electromagnetic sources and amplifiers
 - Vacuum microelectronics
 - Emerging quantum nanoelectronics
- Field emitters attract intensive attention, because of
 - High efficiency
 - High brightness
 - Low emittance
 - Miniaturized device size

Field emitters with coating

- **Artificial coating:** Ultra-thin coatings are fabricated onto metallic cathodes to provide chemical and mechanical protection, and longer current stability, smaller turn-on electric field and enhanced current emission due to the lowering of the effective potential barrier [1], [2].
- **Naturally formed coating:** Native oxides or foreign adsorbates can be easily formed on the surface of the emitter at low vacuum condition [3].
- Coating on the cathode surface forms a double-layer potential barrier
 - Strongly influences the field emission current.
 - Has its potential to change the electrons' mean transverse energy behavior that affects beam quality, important to photoinjectors for future x-ray free electron lasers (XFELs) [4]

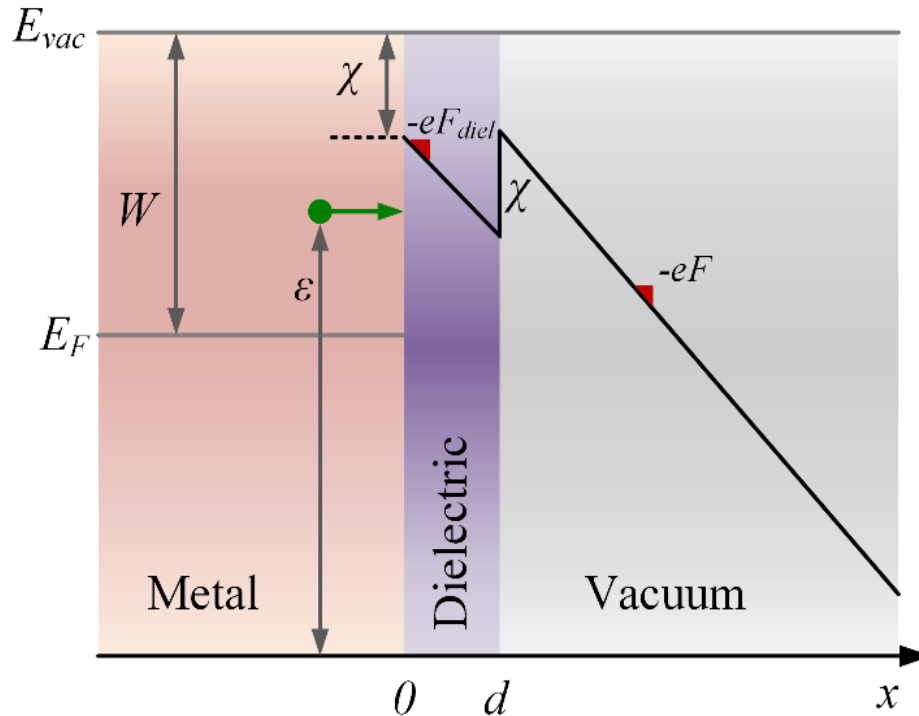
[1] L. W. Swanson and R. W. Strayer, J. Chem. Phys., 48(6), 2421–2442 (1968).

[2] X. Xiong et al., ACS Nano, 14(7), 8806–8815 (2020).

[3] M. Morita and T. Ohmi, Jpn. J. Appl. Phys. 33, 370 (1994).

[4] G. Wang, P. Yang, N. A. Moody, and E. R. Batista, npj 2D Mater Appl. 2, 17 (2018).

Exact Solution to Schrödinger Equation



Potential barrier

$$V(x) = \begin{cases} 0, & x < 0, \\ E_F + W - \chi - eF_{diel}x, & 0 \leq x < d, \\ E_F + W + e(F - F_{diel})d - eFx, & x \geq d. \end{cases}$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + [V(x) - \varepsilon] \psi(x) = 0$$

$\psi(x)$: complex electron wave function

Exact Solution to Schrödinger Equation

$x < 0$ (Inside metal)

$$\psi_I(x) = e^{ik_0x} + R_1 e^{-ik_0x}$$

↑ Incident plane wave ↑ Reflected waves

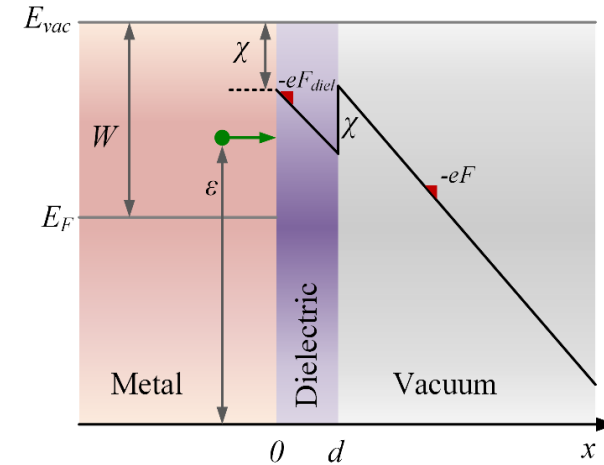
$$k_0 = \sqrt{2m\varepsilon/\hbar^2}$$

$0 \leq x < d$ (Inside dielectric)

$$\psi_{II}(x) = aAi(-\eta_1) + bBi(-\eta_1)$$

$x \geq d$ (Vacuum Region)

$$\psi_{III}(x) = T_3 [Ai(-\eta_2) - iBi(-\eta_2)]$$



$$\eta_1 = \left(\frac{2meF_{diel}}{\hbar^2} \right)^{1/3} \left(x + \frac{\varepsilon - V_1}{eF_{diel}} \right)$$

$$V_1 = W + E_F - \chi$$

$$\eta_2 = \left(\frac{2meF}{\hbar^2} \right)^{1/3} \left(x + \frac{\varepsilon - V_2}{eF} \right)$$

$$V_2 = W + E_F + e(F - F_{diel})d$$

Exact Solution to Schrödinger Equation

Boundary conditions at metal-dielectric interface & dielectric-vacuum interface

$$\left. \begin{aligned} \psi_I(x=0) &= \psi_{II}(x=0) \\ \left. \frac{\partial \psi_I}{\partial x} \right|_{x=0} &= \left. \frac{\partial \psi_{II}}{\partial x} \right|_{x=0} \\ \psi_{II}(x=d) &= \psi_{III}(x=d) \\ \left. \frac{\partial \psi_{II}}{\partial x} \right|_{x=d} &= \left. \frac{\partial \psi_{III}}{\partial x} \right|_{x=d} \end{aligned} \right\} \Rightarrow R_1, a, b, \text{ \& } T_3$$

Electron transmission probability

$$D(\varepsilon) = \frac{J_3(\varepsilon)}{J_i(\varepsilon)} \quad \text{with} \quad J = \frac{i\hbar}{2m} (\psi \nabla \psi^* - \psi^* \nabla \psi)$$

Exact Solution to Schrödinger Equation

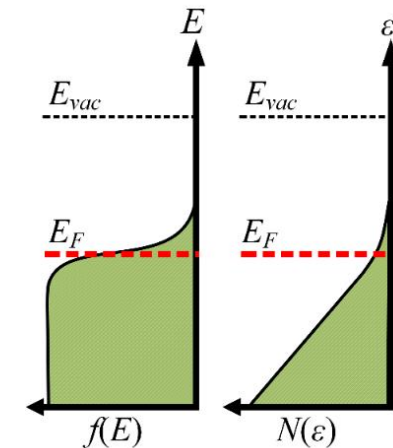
Electron transmission probability

$$D(\varepsilon) = \frac{1}{\pi} \frac{1}{k_0} \left(\frac{2meF}{\hbar^2} \right)^{1/3} |T_3|^2$$

Emission current density

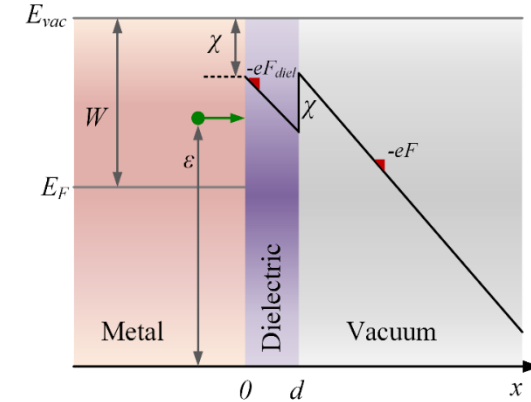
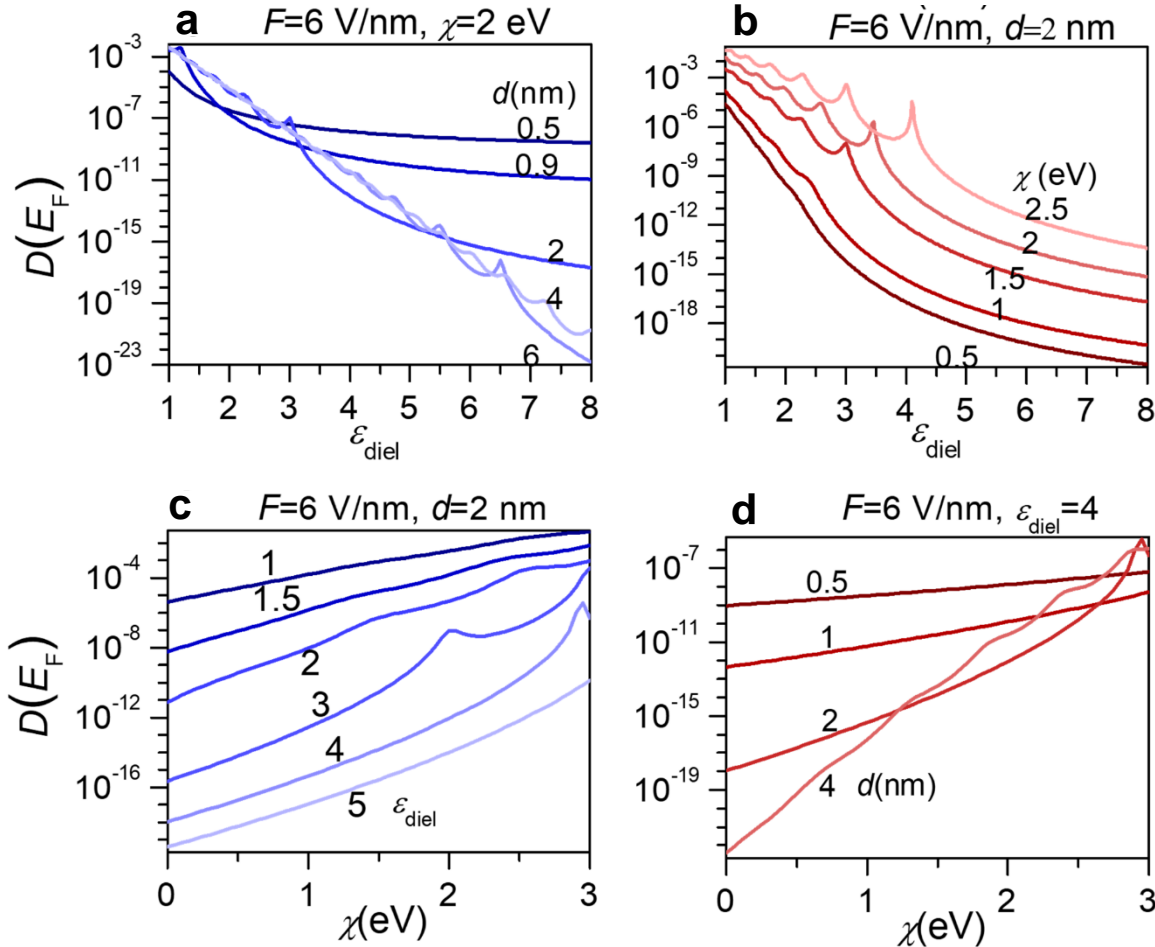
$$J = e \int_0^{\infty} D(\varepsilon) N(\varepsilon) d\varepsilon$$

$$N(\varepsilon) = \frac{mk_B T}{2\pi^2 \hbar^3} \ln \left[1 + \exp \left(\frac{E_F - \varepsilon}{k_B T} \right) \right]$$



where $N(\varepsilon)d\varepsilon$ is the number density of electrons inside metal with longitudinal energy between ε and $\varepsilon+d\varepsilon$ impinging on the surface of metal per unit time [a].

The effects of dielectric properties on electron transmission probability

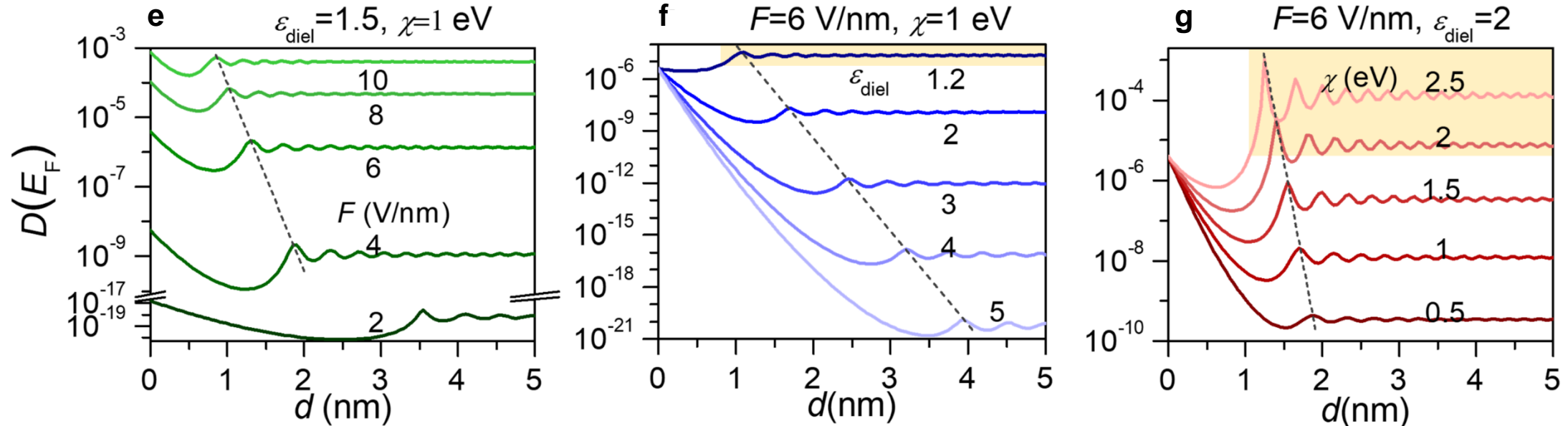


For 1D surface: $F_{\text{diel}} = F / \epsilon_{\text{diel}}$

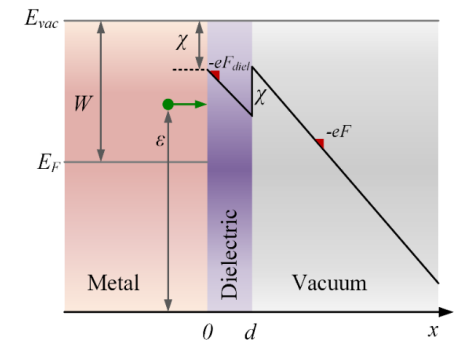
- Electron transmission probability **decreases** with ϵ_{diel}
- **More resonances** for emitters with **thicker** coating
- **Stronger resonances** for **larger** electron affinity
- Electron transmission probability **increases** with χ

Y. Zhou and P. Zhang, "Theory of field emission from dielectric coated surfaces", Phys. Rev. Res., 2, 043439 (2020)

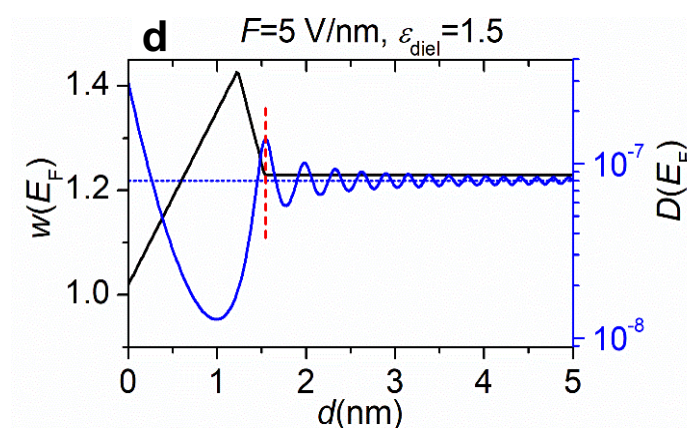
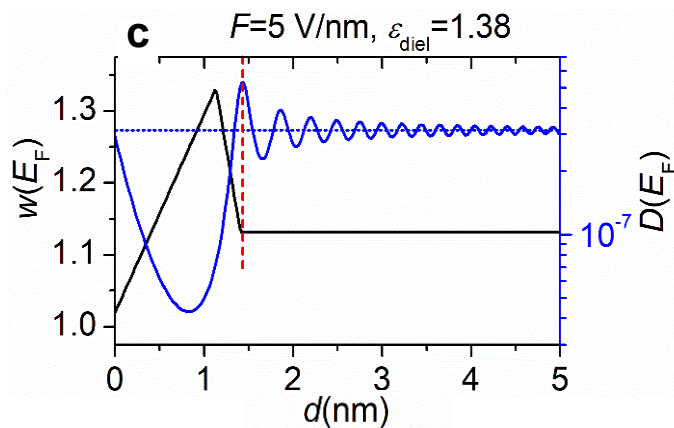
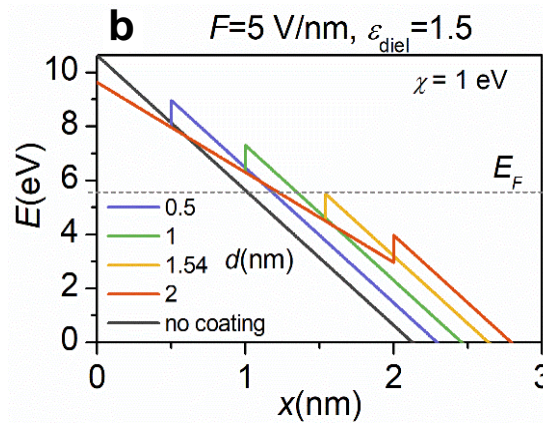
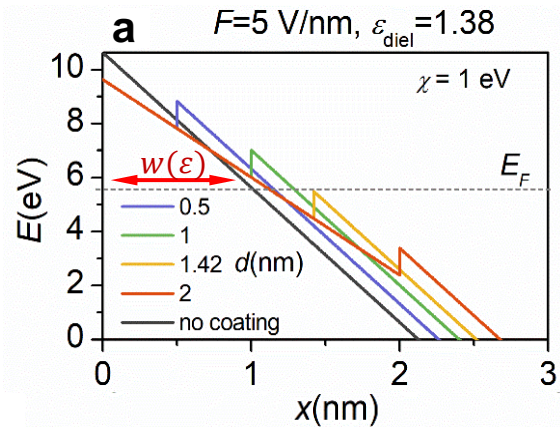
The effects of dielectric properties on electron transmission probability



- **Resonant peaks** for $d > d_0 \approx \epsilon_{diel}W/eF$
- Increasing F or χ shifts the **maximum** peak to **smaller** thickness
- Increasing ϵ_{diel} shifts the **maximum** peak to **larger** thickness
- **Coated case** can induce a transmission probability **larger** than the bare case



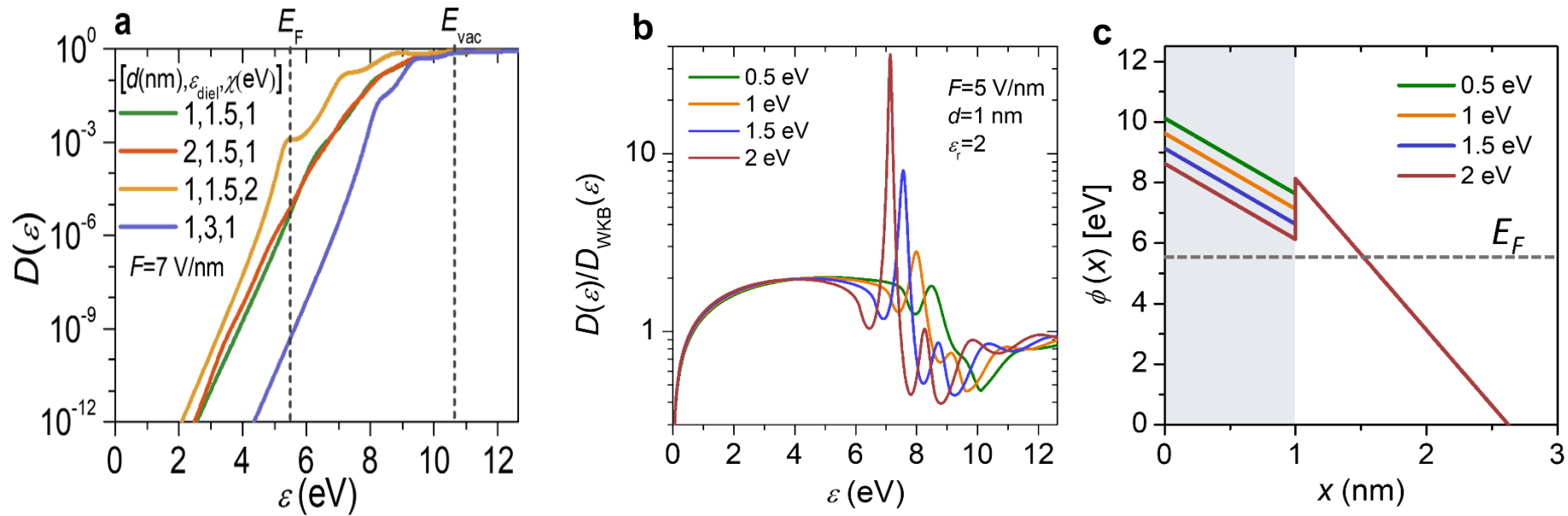
Potential profile and transmission probability



- The first resonance peak is at $d_0 \approx \varepsilon_{diel} W / eF$
- $D(E_F)$ oscillates around a **constant**, which is the tunneling probability through a single triangular barrier **in thick dielectric**.

Y. Zhou and P. Zhang, "Theory of field emission from dielectric coated surfaces", Phys. Rev. Res., 2, 043439 (2020)

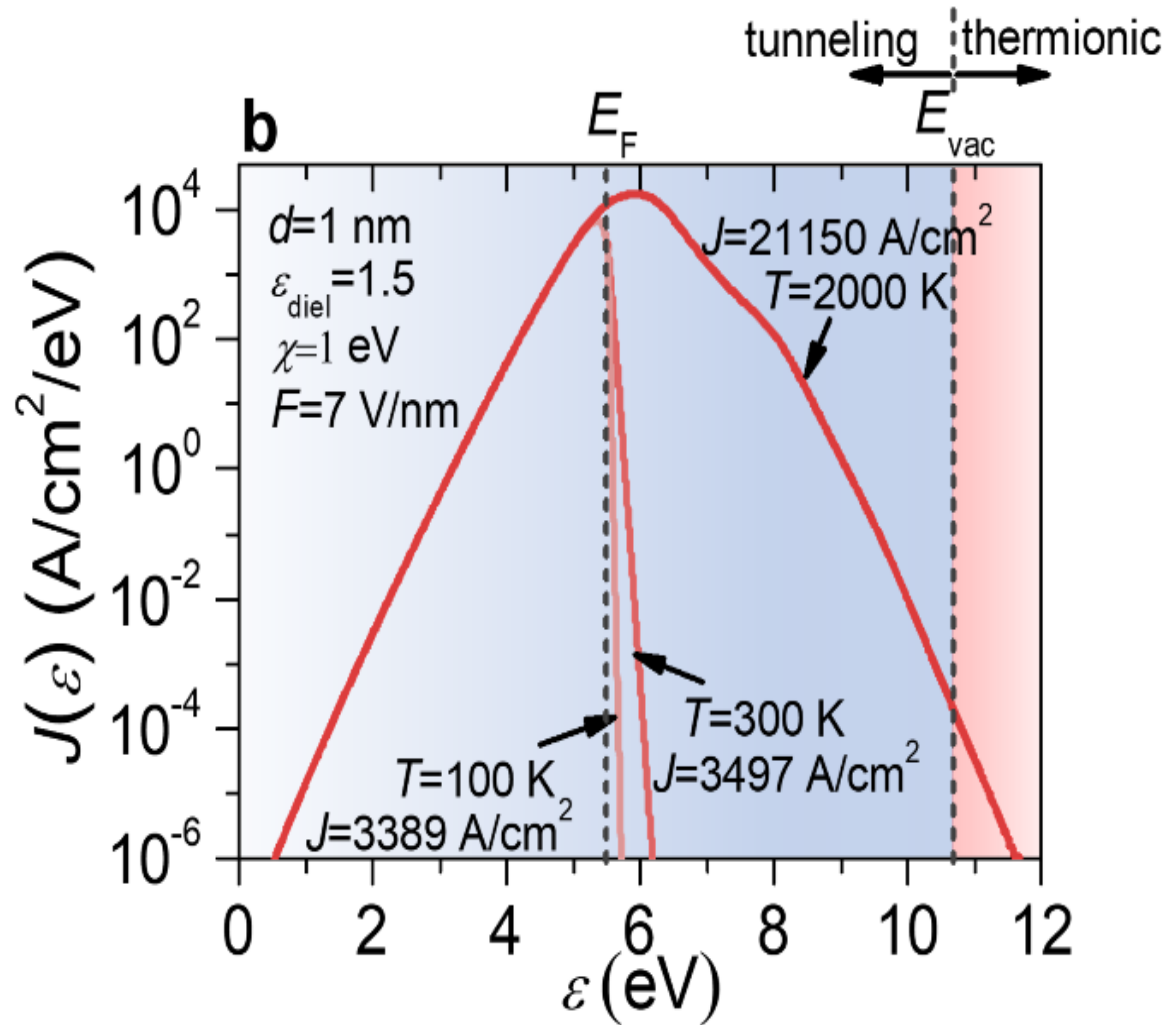
Potential profile and transmission probability



- **Stronger resonance** is observed for **larger electron affinity**

Y. Zhou and P. Zhang, "Theory of field emission from dielectric coated surfaces", Phys. Rev. Res., 2, 043439 (2020)

Field emission and thermionic emission



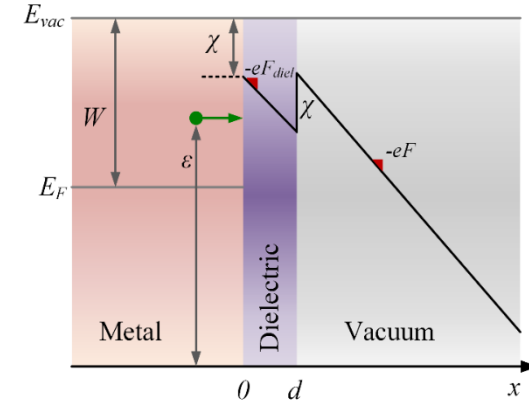
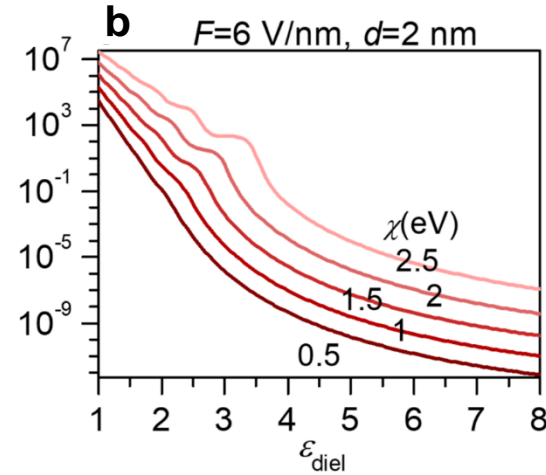
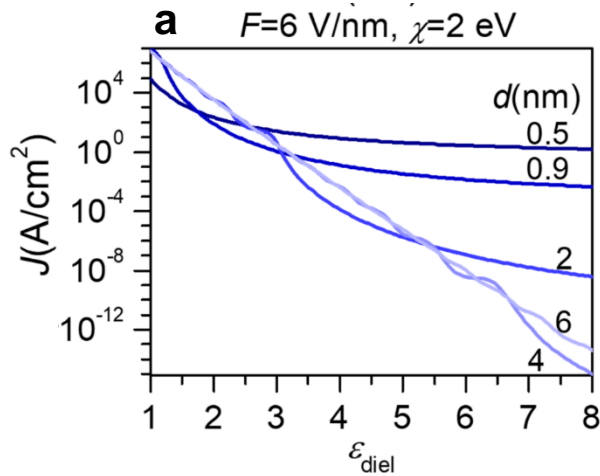
$$J(\epsilon) = D(\epsilon)N(\epsilon)$$

$$N(\epsilon) = \frac{mk_B T}{2\pi^2 \hbar^3} \ln \left[1 + \exp \left(\frac{E_F - \epsilon}{k_B T} \right) \right]$$

- The model includes not only **field emission** but also **thermionic emission**

Y. Zhou and P. Zhang, "Theory of field emission from dielectric coated surfaces", Phys. Rev. Res., 2, 043439 (2020)

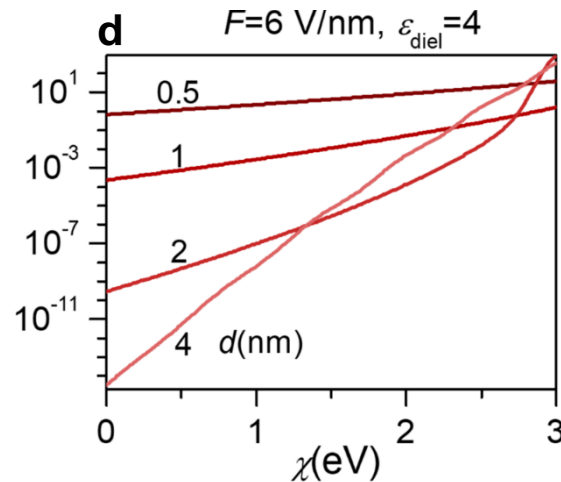
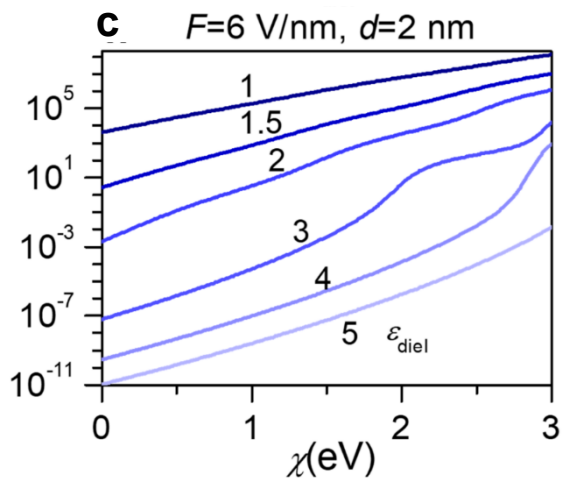
The effects of dielectric properties on field emission current density



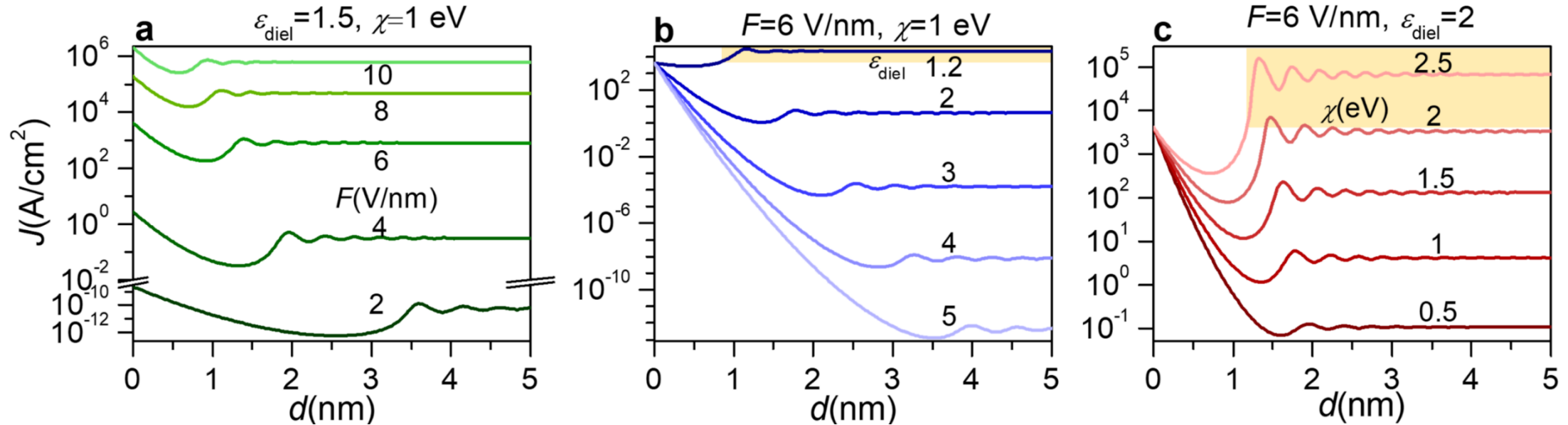
For 1D surface: $F_{\text{diel}} = F / \epsilon_{\text{diel}}$

- J decreases with ϵ_{diel}
- Resonance peaks are not as sharp as in $D(E_F)$ vs ϵ_{diel}
- J increases with χ

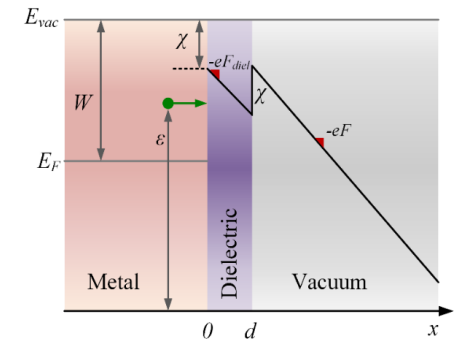
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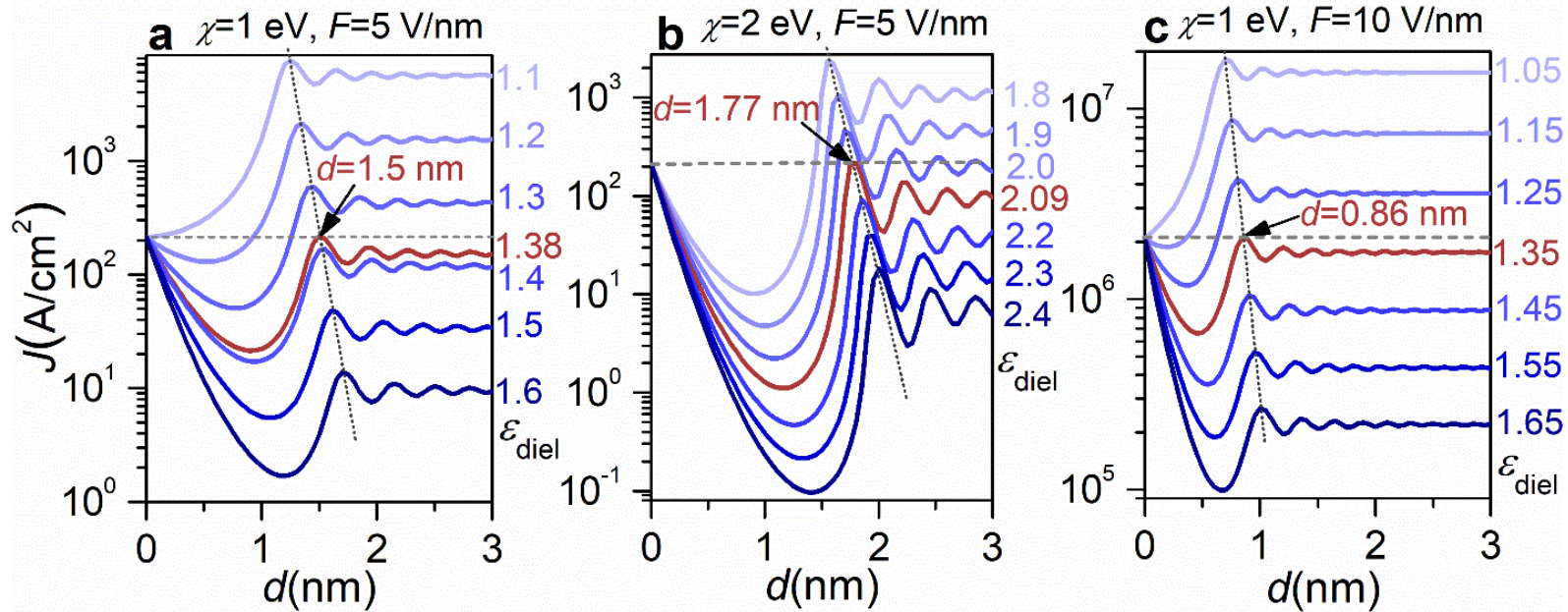
The effects of dielectric properties on field emission current density



- **Resonant peaks** for $d > d_0 \approx \epsilon_{\text{diel}}W/eF$
- Increasing F or χ shifts the **maximum** peak to **smaller** thickness
- Increasing ϵ_{diel} shifts the **maximum** peak to **larger** thickness
- **Coated case can induce a current density larger than the bare case**

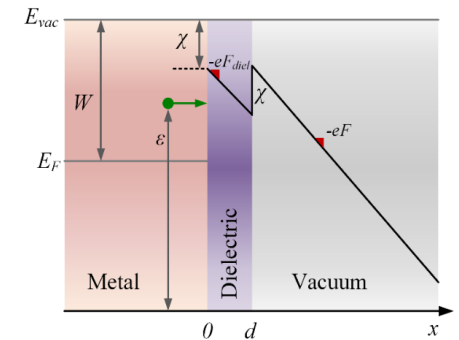


Determine the threshold thickness and threshold dielectric constant for current enhancement

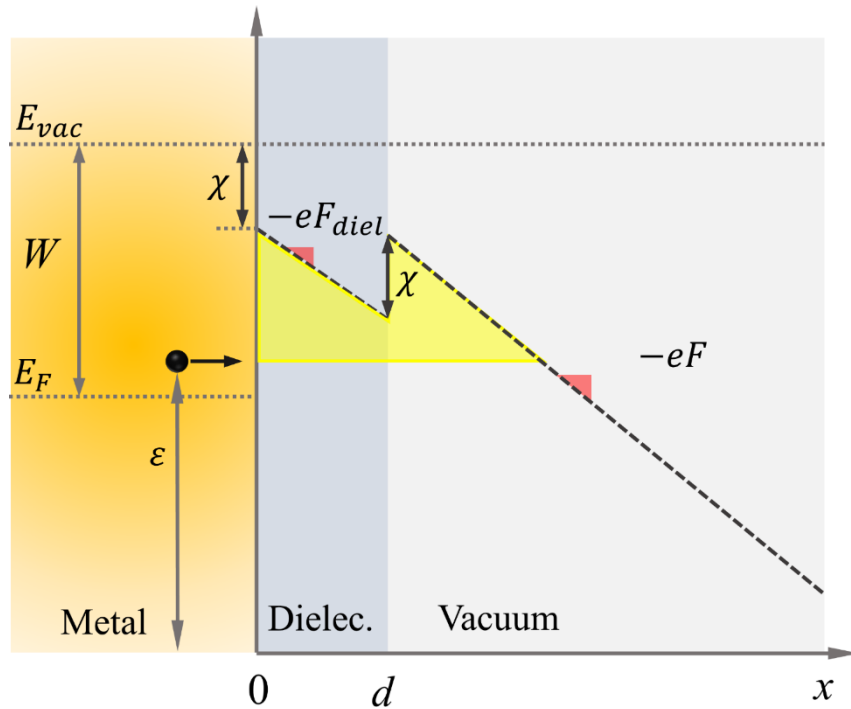


- An empirical relation between thickness threshold and dielectric constant threshold at room temperature:

$$d_{th} [\text{nm}] = \frac{\epsilon_{diel}^{th} W}{eF}$$



Comparison with modified double-barrier Fowler-Nordheim equation [1-3]



- According to WKB approximation,
- Electron transmission probability
- Electron emission current density

$$Q(\varepsilon) = 2 \int_0^{x_1} \sqrt{\frac{2m}{\hbar^2} [V(x) - \varepsilon]} dx$$

$$D(\varepsilon) = \exp[Q(\varepsilon)]$$

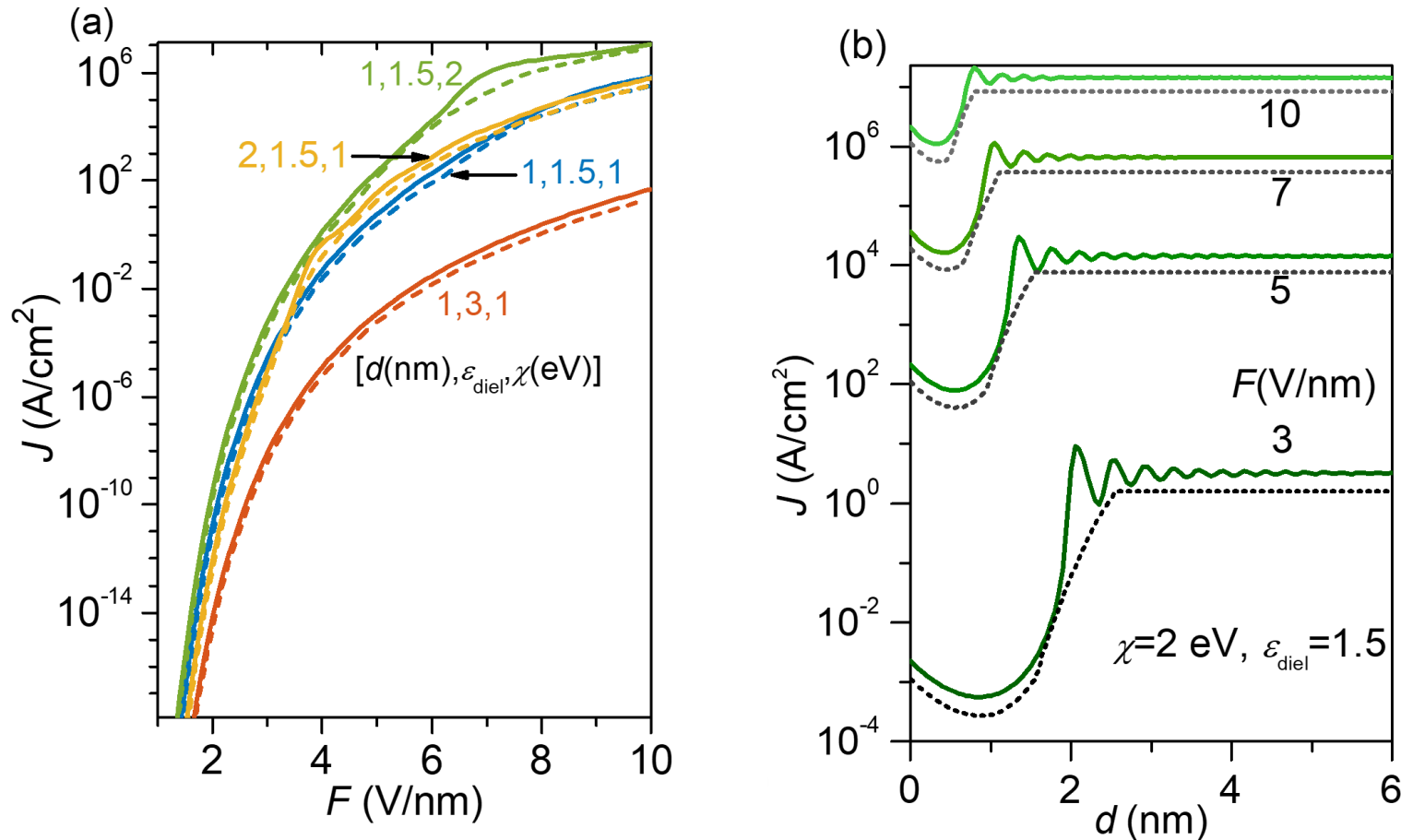
$$J = e \int_0^{\infty} D(\varepsilon) N(\varepsilon) d\varepsilon$$

[1] Q.-A. Huang, J. Appl. Phys. 79(7), 3703-3707, (1996).

[2] P. D. Keathley et al., Ann. Phys., 525(1-2), 144-150 (2013).

[3] K. L. Jensen, et al., J. Appl. Phys. 127(23), 235301, (2020).

Comparison with modified double-barrier Fowler-Nordheim equation [1-3]

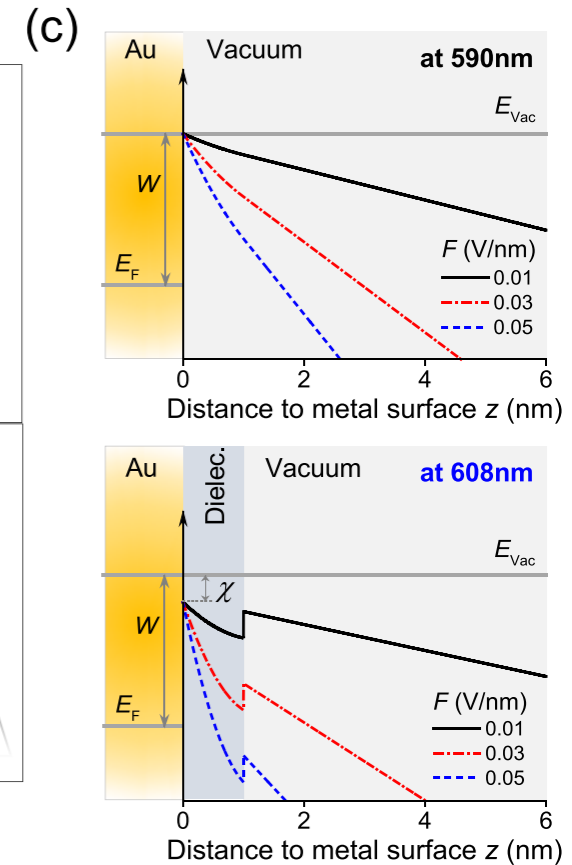
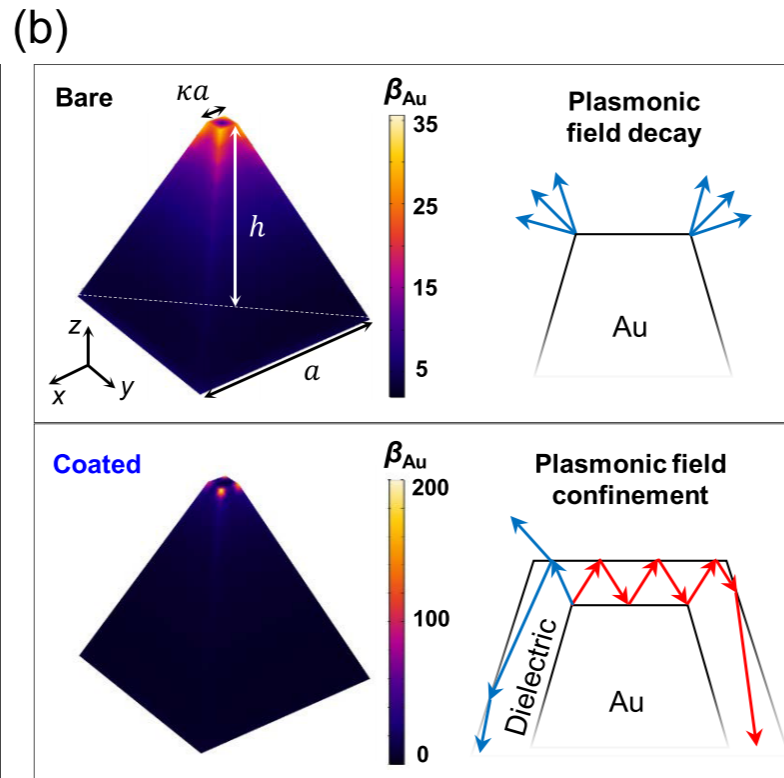
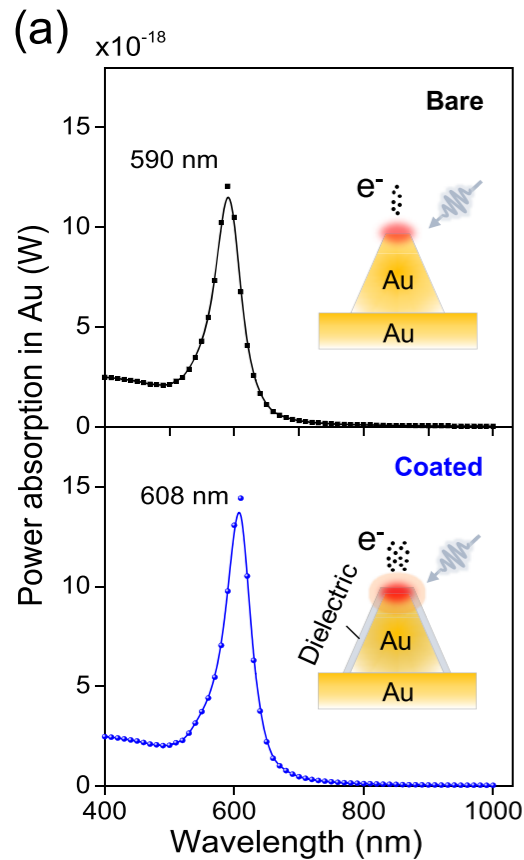


- Good **agreement** in the **scaling**
- Resonance behavior in J vs. d cannot be revealed by the modified FN equation

Thin dielectric coated plasmonic photoemitter

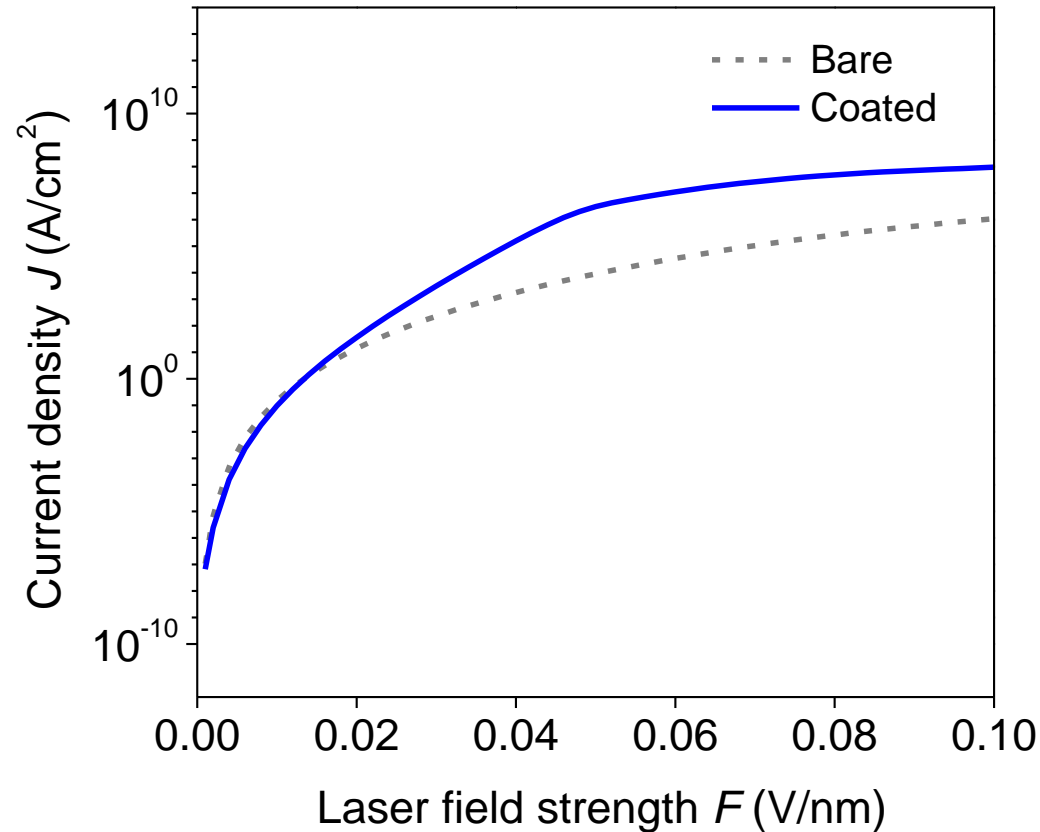
Plasmonic field confinement
Enhanced fields at Au surface

Electron affinity χ
Reduced tunneling barrier



X. Xiong, Y. Zhou, Y. Luo, X. Li, M. Bosman, L. K. Ang, P. Zhang, and L. Wu, "Plasmon-Enhanced Resonant Photoemission Using Atomically Thick Dielectric Coatings", ACS Nano, 14, 8806 – 8815 (2020).

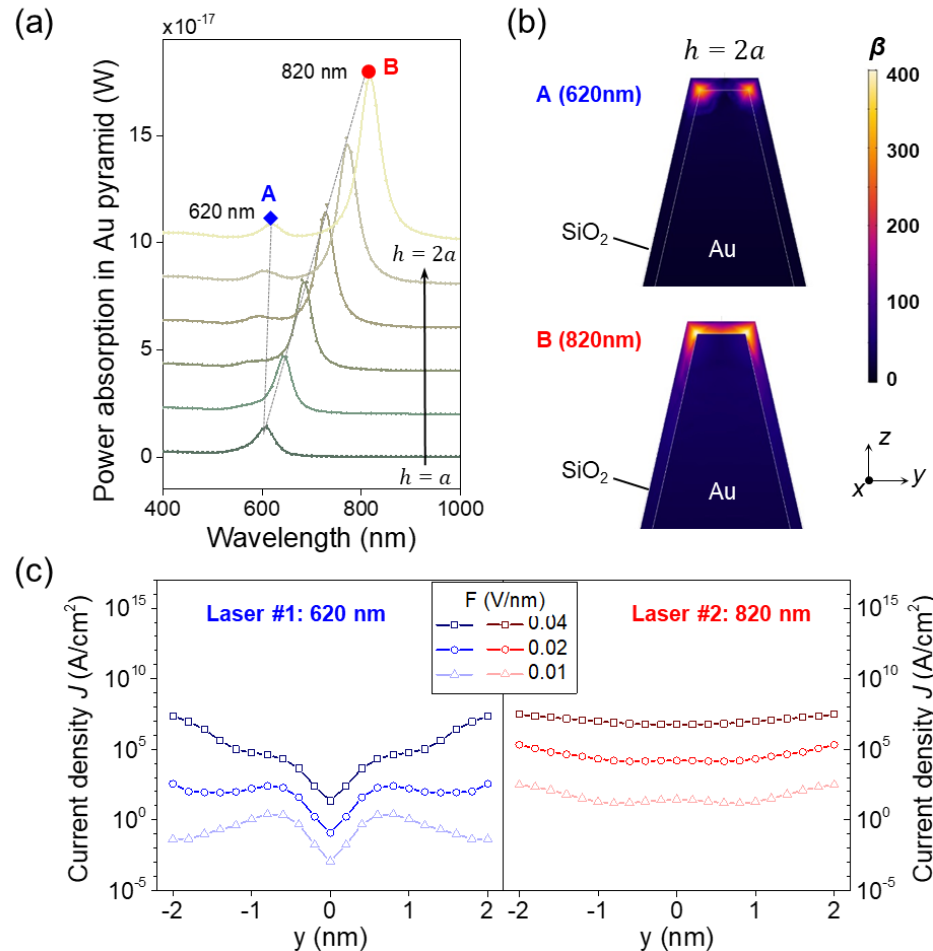
Thin dielectric coated plasmonic photoemitter



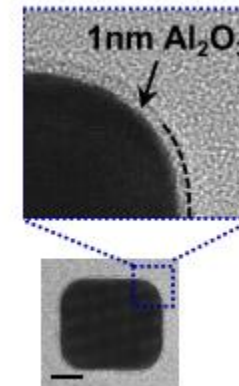
For $F = 0.014 - 1$ V/nm, J from the coated photoemitter is **enhanced by at least 2 orders of magnitude** as compared to the bare emitter.

X. Xiong, Y. Zhou, Y. Luo, X. Li, M. Bosman, L. K. Ang, P. Zhang, and L. Wu, "Plasmon-Enhanced Resonant Photoemission Using Atomically Thick Dielectric Coatings", ACS Nano, 14, 8806 – 8815 (2020).

Thin dielectric coated plasmonic photoemitter



Tunable photoemission
at different plasmonic
resonant wavelengths



X. Xiong, Y. Zhou, Y. Luo, X. Li, M. Bosman, L. K. Ang, P. Zhang, and L. Wu, "Plasmon-Enhanced Resonant Photoemission Using Atomically Thick Dielectric Coatings", ACS Nano, 14, 8806 – 8815 (2020).

Summary

- An **exact** analytical quantum model is constructed for electron emission from dielectric coated surfaces.
- Cathodes with coatings of **smaller dielectric constant** and **larger electron affinity** tend to emit a **larger** current density than the bare metal surfaces.
- Empirical relation of the threshold thickness and threshold dielectric constant is proposed.
- The model shows good agreement in scaling with modified double-barrier Fowler-Nordheim equation, but also reveals new **resonant** peaks.
- Plasmon resonant photoemission from dielectric coated photoemitter is also studied.



NanoPATH Group

Nanoelectronics, Plasmas, & Accelerator Technology

<http://www.egr.msu.edu/~pz>

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