

# Nucleon tomography with GPDs at the EIC

Hyon-Suk Jo

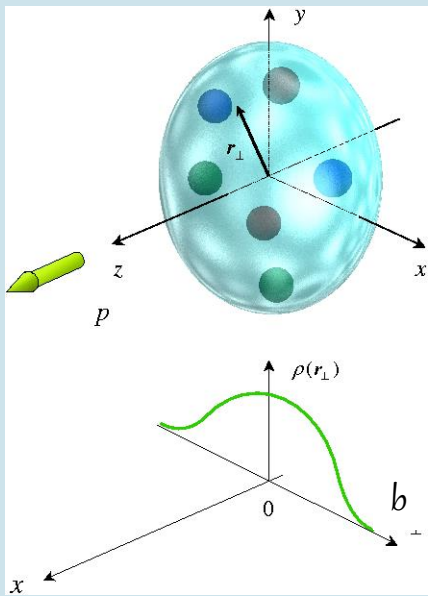
Kyungpook National University

2020 KPS Fall Meeting

Focus Session : Cornerstone for future collider projects

November 5, 2020

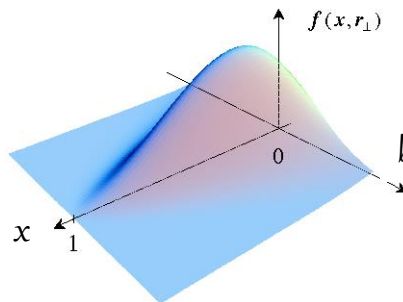
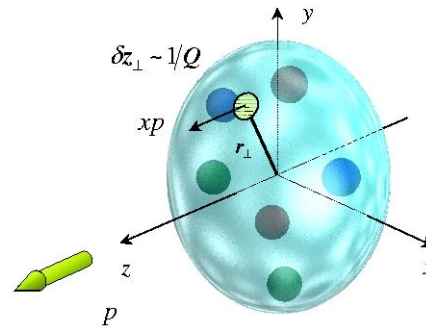
# Generalized parton distributions (GPDs)



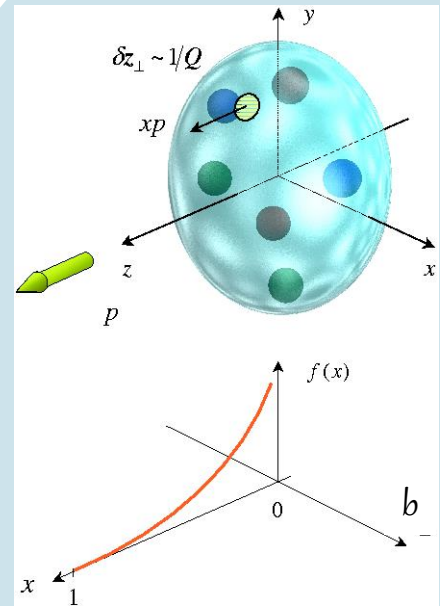
Form Factors (elastic scattering) :  
**Transverse position** of the quarks in the nucleon

GPDs, accessible via **exclusive reactions**, provide a **correlation** between the **transverse position** and the **longitudinal momentum** of the quarks in the nucleon

## Nucleon tomography



*Transverse position ( $b$ ) as a function of longitudinal momentum fraction ( $x$ )*

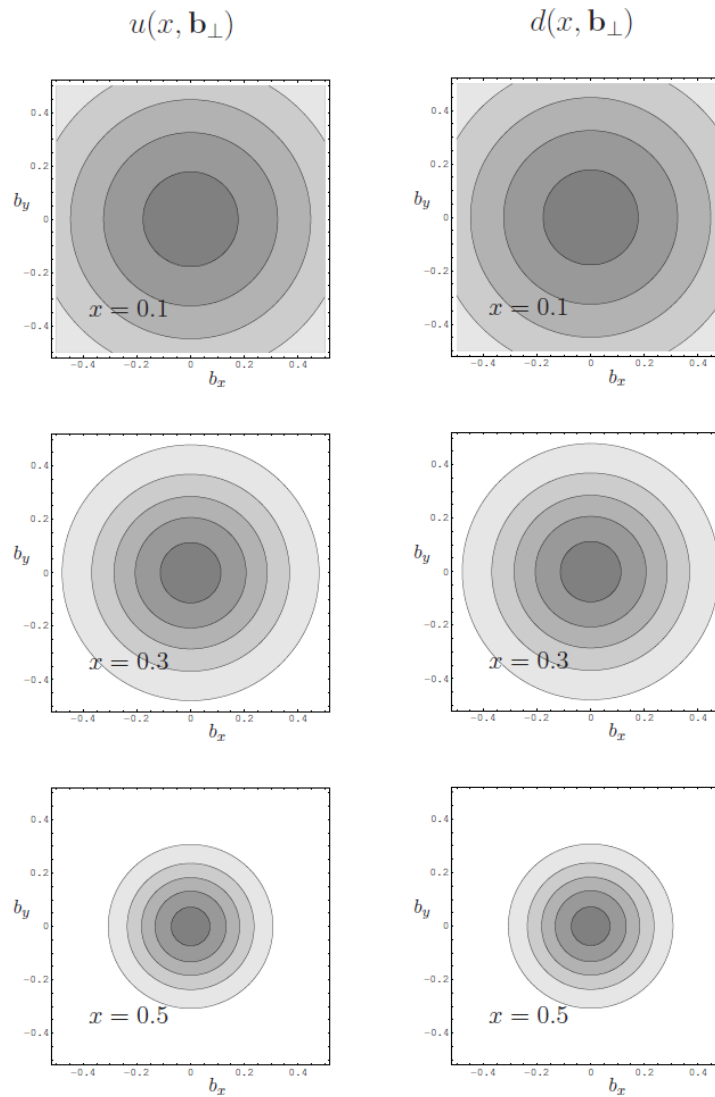


Parton Distribution Functions (deep inelastic scattering) :  
**Longitudinal momentum** of the quarks in the nucleon

*In this model, valence quarks (high  $x$ ) are at the heart of the nucleon and sea quarks (low  $x$ ) extend to its periphery*

# Interpretation of GPDs : impact parameter $b_{\perp}$ as a function of $x$

*Transverse position  $b_{\perp}$  of the partons inside the nucleon for different values of longitudinal momentum fraction  $x$*

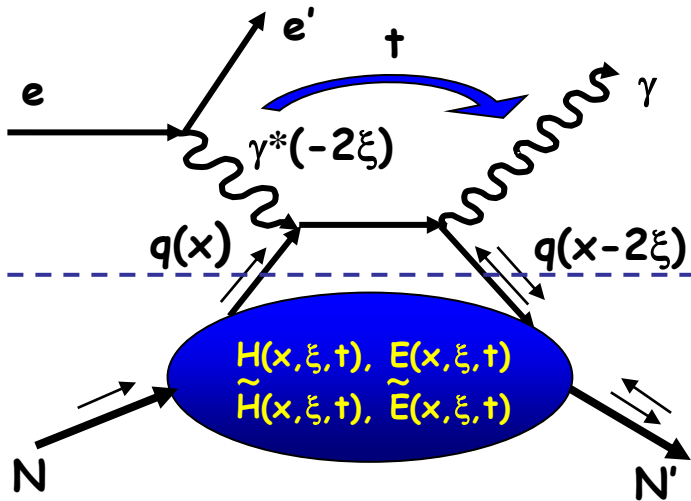


# Deeply Virtual Compton Scattering (DVCS) and GPDs

high  $Q^2$ , small  $t$ , fixed  $x_B$

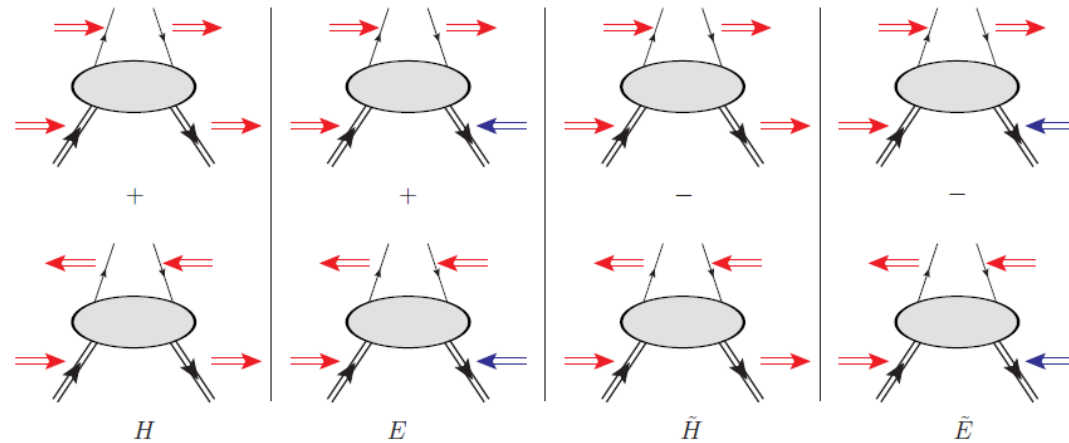
## DVCS

$$x_B = \frac{2\xi}{1+\xi}$$



$$e N \longrightarrow e' N' \gamma$$

At leading-order QCD, leading twist, there are 4 chiral-even (parton helicity is conserved) GPDs for each parton

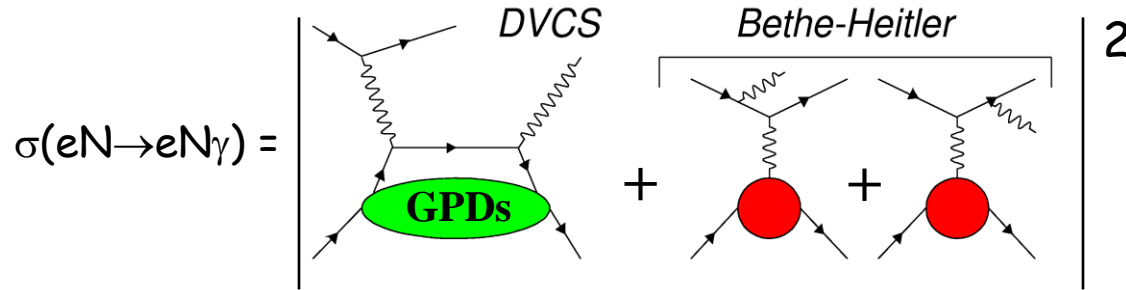


DVCS is the key reaction to access the GPDs as it offers the simplest interpretation in terms of GPDs

$H^{q,g}(x, \xi, t)$	$E^{q,g}(x, \xi, t)$	for sum over parton helicities
$\tilde{H}^{q,g}(x, \xi, t)$	$\tilde{E}^{q,g}(x, \xi, t)$	for difference over parton helicities
nucleon helicity conserved	nucleon helicity changed	

# DVCS and Bethe-Heitler processes

BH fully calculable in QED



DVCS and Bethe-Heitler (BH) **experimentally undistinguishable**  
interference between the 2 processes

$$T^{DVCS} \sim \int_{-1}^{+1} \frac{H(x, \xi, t)}{x \pm \xi + i\epsilon} dx + \dots \sim P \int_{-1}^{+1} \frac{H(x, \xi, t)}{x \pm \xi} dx - i\pi H(\pm\xi, \xi, t) + \dots$$

Unpolarized Cross Section

$$\frac{d^4 \sigma}{dQ^2 dx_B dt d\phi} \approx |T^{DVCS} + T^{BH}|^2 = |T^{DVCS}|^2 + |T^{BH}|^2 + I$$

Beam-polarized Cross-  
Section difference

$$\frac{d^4 \vec{\sigma}}{dQ^2 dx_B dt d\phi} - \frac{d^4 \leftarrow{\sigma}}{dQ^2 dx_B dt d\phi} \propto \text{Im}(T_{DVCS}) \times T_{BH}$$

# GPDs and proton spin puzzle

Proton spin puzzle :  
The origin of the proton spin is still unknown

$$\frac{1}{2} = J^q + J^g = \frac{1}{2}\Delta\Sigma + \Delta G + L_q + L_g$$

Orbital angular momentum

GPDs H and E provide access to the total angular momentum of the partons in the nucleon

Ji's angular momentum sum rule:

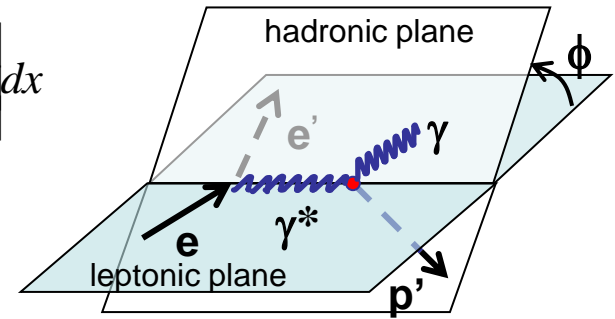
$$J^{q,g} = \frac{1}{2} \int_{-1}^1 x dx (H^{q,g}(x, \xi, t=0) + E^{q,g}(x, \xi, t=0))$$

# Compton Form Factors (CFFs) and DVCS observables

Compton Form Factors (CFFs)

$$\begin{cases} \text{Re}\mathcal{H}_q = e_q^2 P \int_0^{+1} \left( H^q(x, \xi, t) - H^q(-x, \xi, t) \right) \left[ \frac{1}{\xi - x} + \frac{1}{\xi + x} \right] dx \\ \text{Im}\mathcal{H}_q = \pi e_q^2 \left[ H^q(\xi, \xi, t) - H^q(-\xi, \xi, t) \right] \end{cases}$$

$$\xi = x_B / (2 - x_B) \quad k = t / 4M^2$$

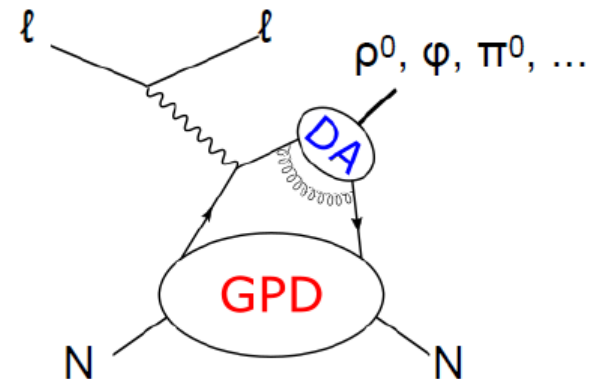
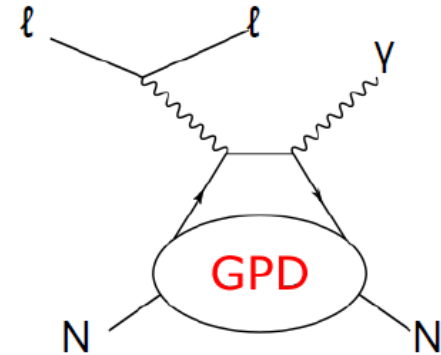


Each DVCS observable is sensitive to a different combination of GPDs

	Proton Neutron
Polarized beam, unpolarized target: $\Delta\sigma_{LU} \sim \sin\phi \text{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E} + \dots\}$	$\Rightarrow \begin{matrix} \text{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p, \mathcal{E}_p\} \\ \text{Im}\{\mathcal{H}_n, \tilde{\mathcal{H}}_n, \mathcal{E}_n\} \end{matrix}$
Unpolarized beam, longitudinal target: $\Delta\sigma_{UL} \sim \sin\phi \text{Im}\{F_1\tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E}) - \xi kF_2\tilde{\mathcal{E}}\}$	$\Rightarrow \begin{matrix} \text{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ \text{Im}\{\mathcal{H}_n, \mathcal{E}_n\} \end{matrix}$
Polarized beam, longitudinal target: $\Delta\sigma_{LL} \sim (A+B\cos\phi) \text{Re}\{F_1\tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E}) + \dots\}$	$\Rightarrow \begin{matrix} \text{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ \text{Re}\{\mathcal{H}_n, \mathcal{E}_n\} \end{matrix}$
Unpolarized beam, transverse target: $\Delta\sigma_{UT} \sim \cos\phi \sin(\phi_s - \phi) \text{Im}\{k(F_2\mathcal{H} - F_1\mathcal{E}) + \dots\}$	$\Rightarrow \begin{matrix} \text{Im}\{\mathcal{H}_p, \mathcal{E}_p\} \\ \text{Im}\{\mathcal{H}_n\} \end{matrix}$

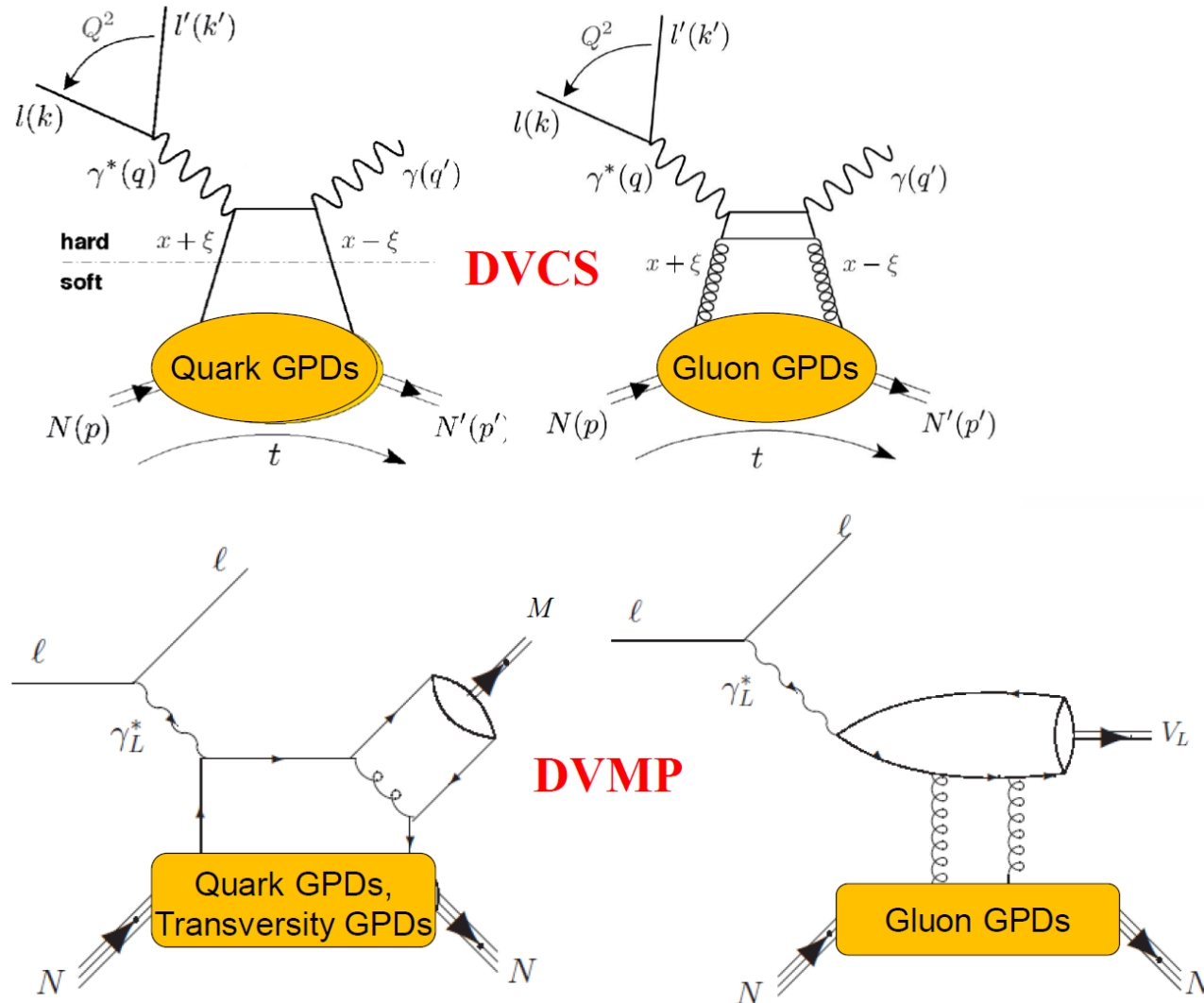
# Quark-flavor separation of GPDs

- 1<sup>st</sup> method : Neutron DVCS
  - A combined analysis of DVCS observables for proton and neutron (deuterium) targets is needed to perform a quark-flavor separation of the GPDs
  - High sensitivity to GPD E with a polarized beam
- 2<sup>nd</sup> method : Deeply virtual meson production (DVMP)
  - DVMP cross-section measurements are another way to access quark-flavor separation of GPDs
  - Gives also access to transversity GPDs (chiral-odd GPDs, where parton helicity is changed)

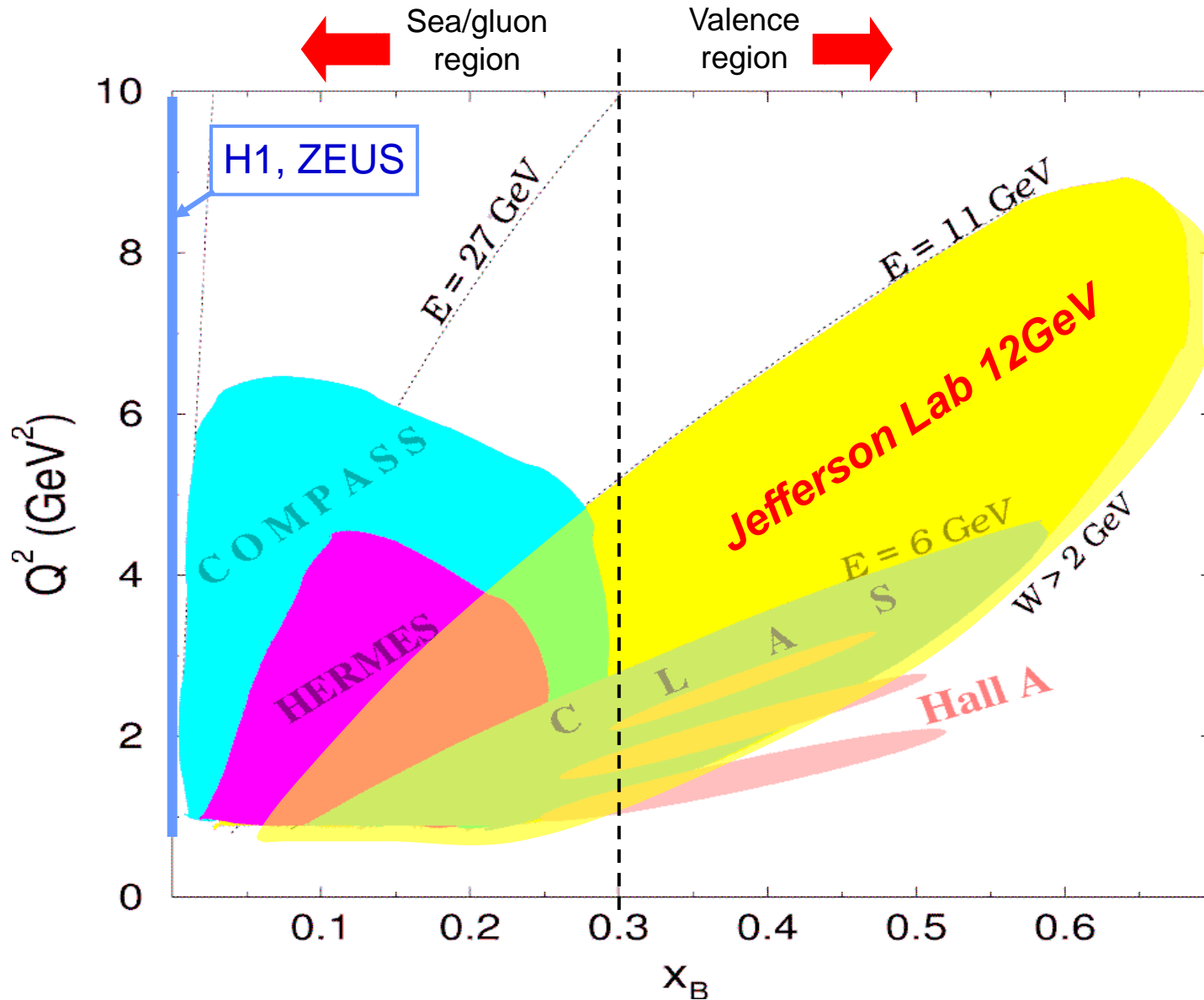




# DVCS and DVMP diagrams for gluons GPDs

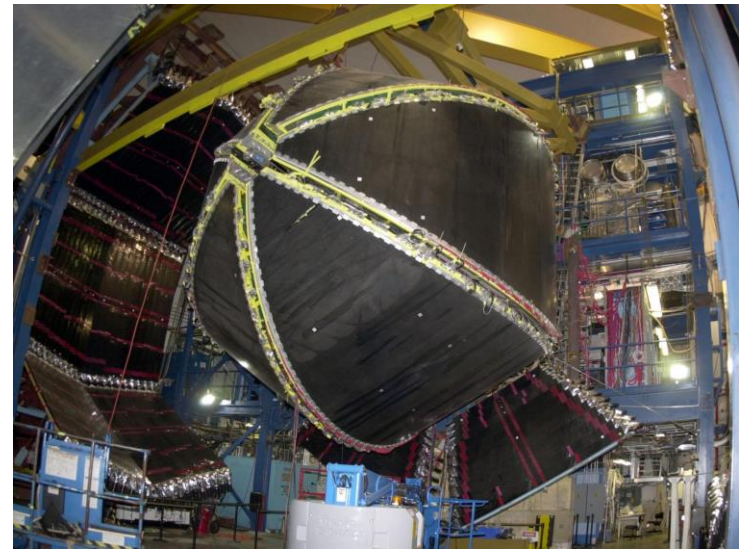
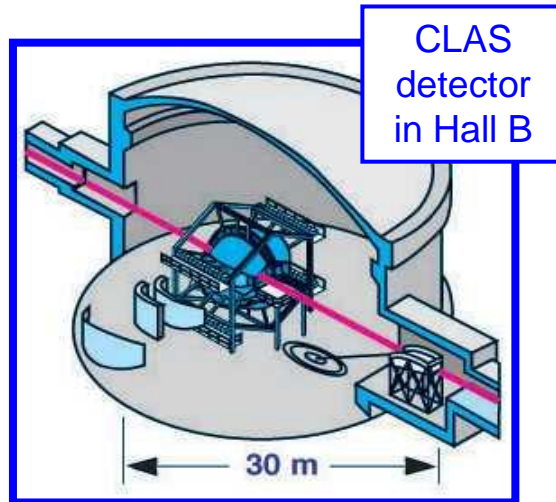
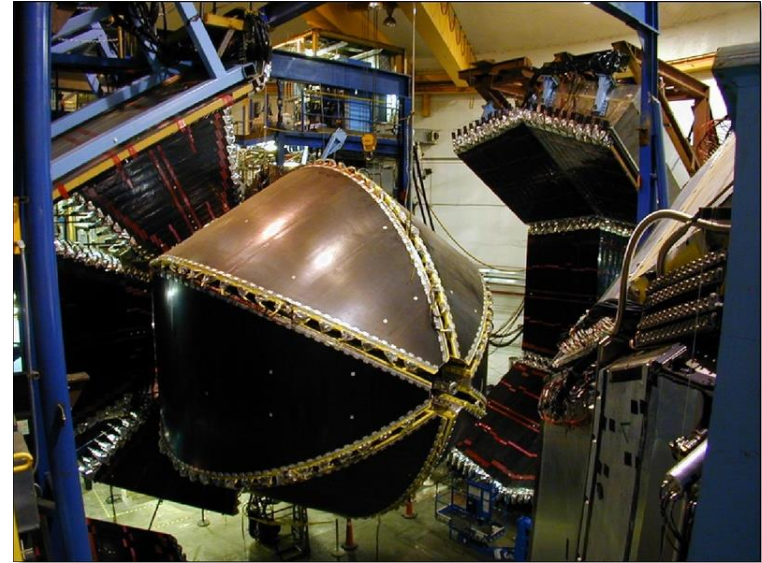
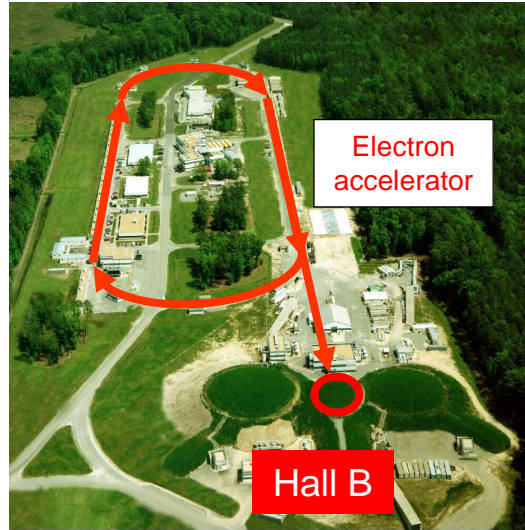


# Kinematic coverage of past and current DVCS experiments



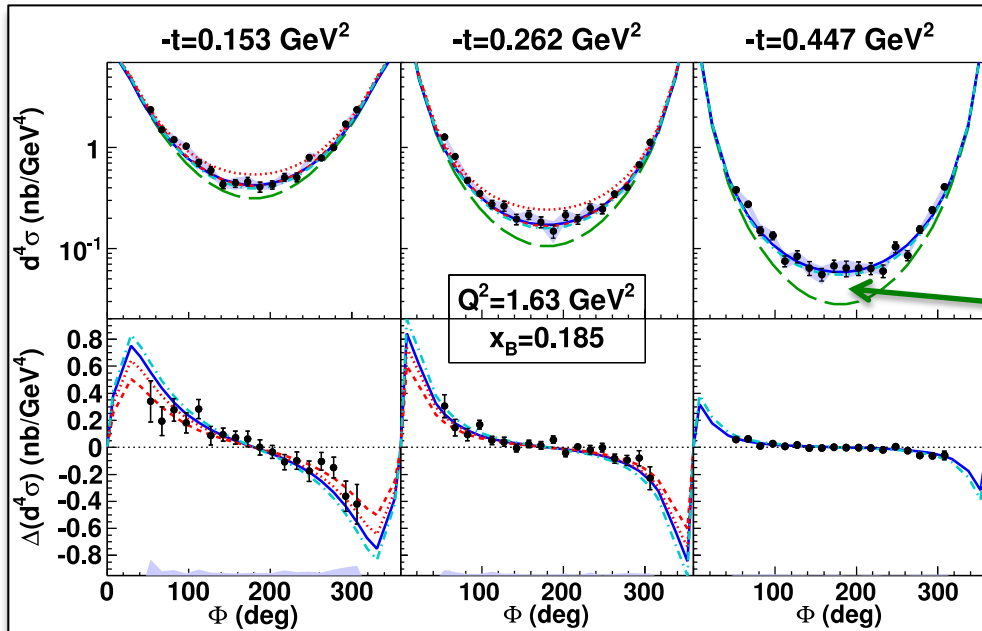
# Jefferson Lab 6 GeV and the CLAS detector

Jefferson Lab



# DVCS unpolarized and beam-polarized cross sections from CLAS data

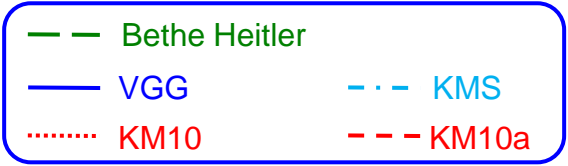
H.S. Jo *et al.* (CLAS Collaboration),  
Phys. Rev. Lett. 115, 212003 (2015)



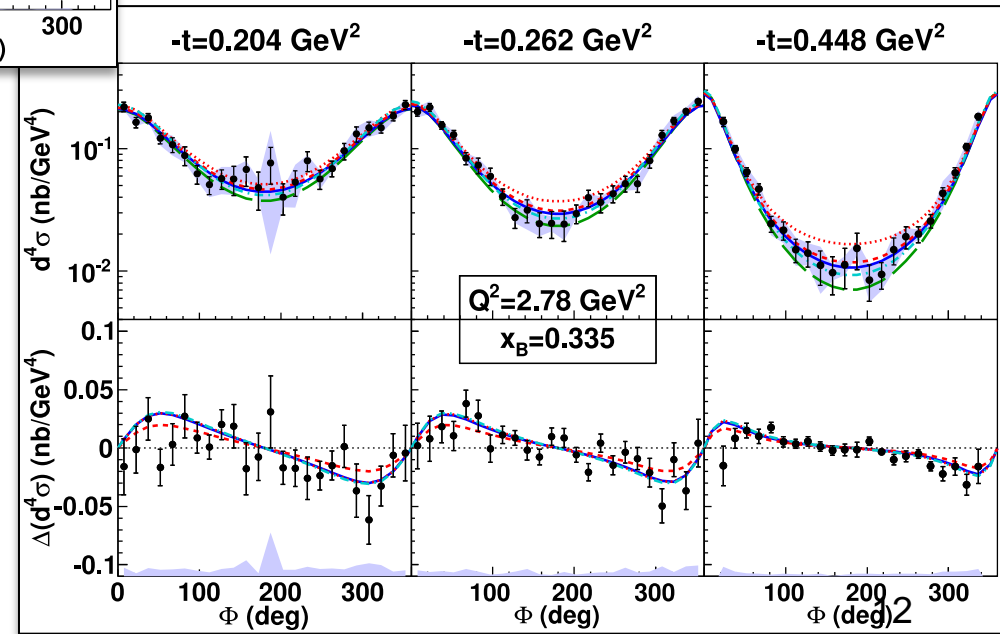
$$\frac{d^4 \sigma_{ep \rightarrow ep\gamma}}{dQ^2 dx_B dt d\Phi} \text{ (nb/GeV}^4\text{)}$$

DVCS + interference

$$\frac{1}{2} \left( \frac{d^4 \bar{\sigma}_{ep \rightarrow ep\gamma}}{dQ^2 dx_B dt d\Phi} - \frac{d^4 \bar{\sigma}_{ep \rightarrow ep\gamma}}{dQ^2 dx_B dt d\Phi} \right) \text{ (nb/GeV}^4\text{)}$$

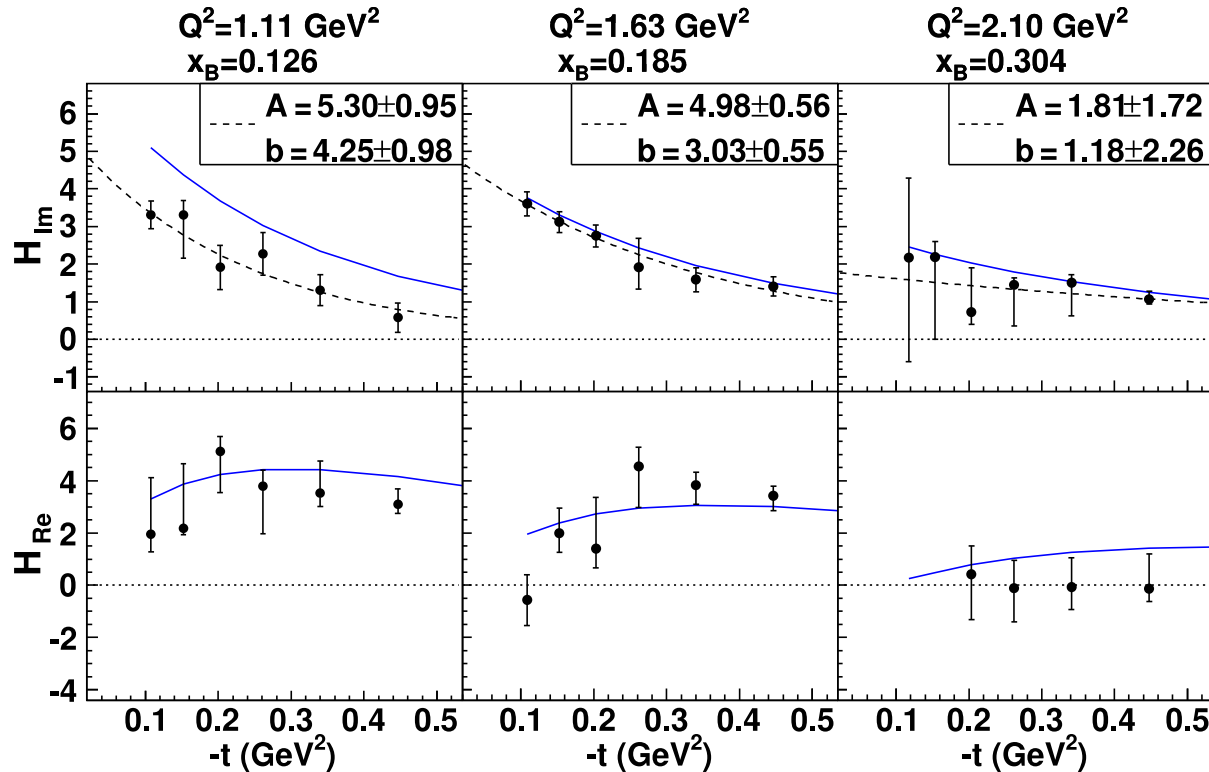


- Jefferson Lab's polarized electron beam (energy ~ 6 GeV, polarization ~ 80%) + LH<sub>2</sub> target
- Luminosity  $L = 2.10^{34} \text{ cm}^{-2}\text{s}^{-1}$

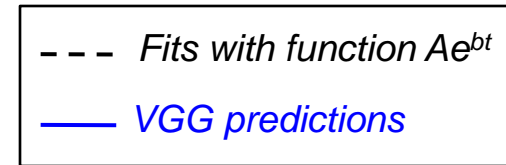


# Interpretation of fit results obtained from the cross sections

H.S. Jo *et al.* (CLAS Collaboration),  
Phys. Rev. Lett. 115, 212003 (2015)



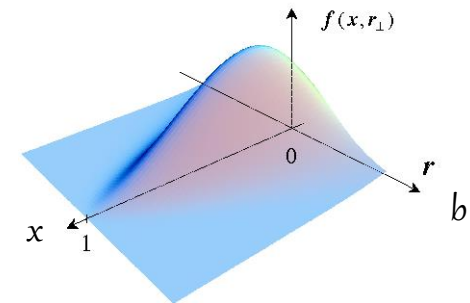
←  $b$  : transverse position of the quarks in the nucleon



The results tend to show that valence quarks (high  $x$ ) are at the heart of the nucleon and sea quarks (low  $x$ ) extend to its periphery

The transverse position  $b$  decreases with increasing  $x_B$

The results suggest that the nucleon size decreases at higher parton-momentum values, thus revealing from the experiment a **first tomographic image of the nucleon**



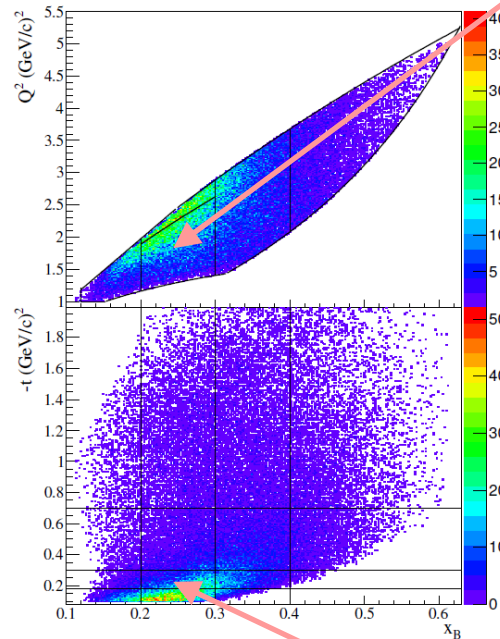
# DVCS on longitudinally polarized target from CLAS data

- **eg1-dvcs experiment**
- Beam energy  $\sim 6$  GeV
- CLAS + IC to detect forward photons
- Target: **longitudinally polarized NH<sub>3</sub>** (P $\sim 80\%$ )
- **3 DVCS observables**

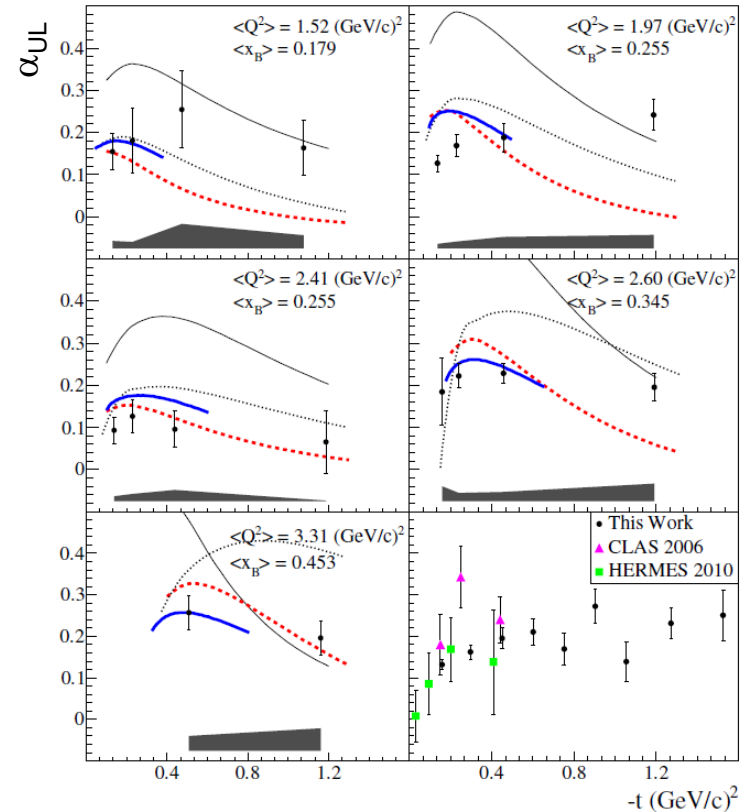
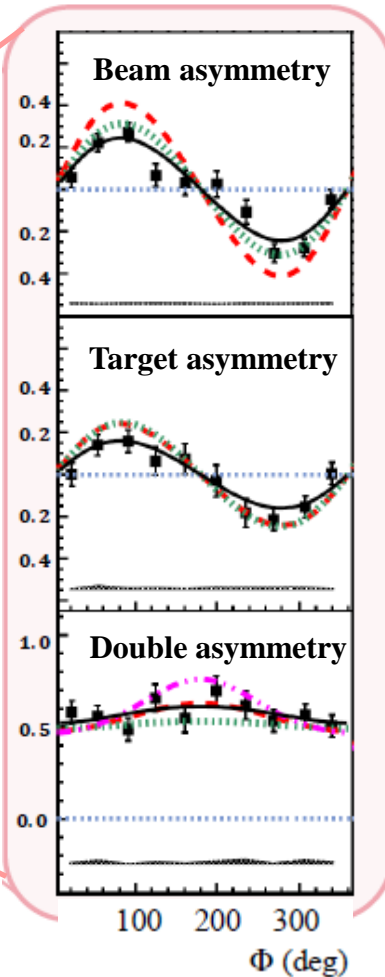
$$\vec{e}\vec{p} \rightarrow e\vec{p}\gamma$$

$$A_{UL} \sim \text{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$

$$A_{LU} = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-}$$



- **5 ( $Q^2-x_B$ ) bins**
- **4  $t$  bins**
- **10  $\phi$  bins**



- Improved statistics  $\times 10$  at low  $-t$
- Extended kinematic coverage

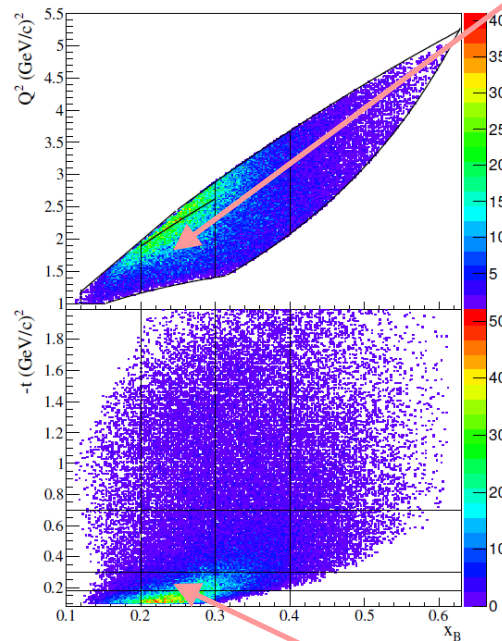
E. Seder *et al.* (CLAS Collaboration),  
Phys. Rev. Lett. 114, 032001 (2015)

# DVCS on longitudinally polarized target from CLAS data

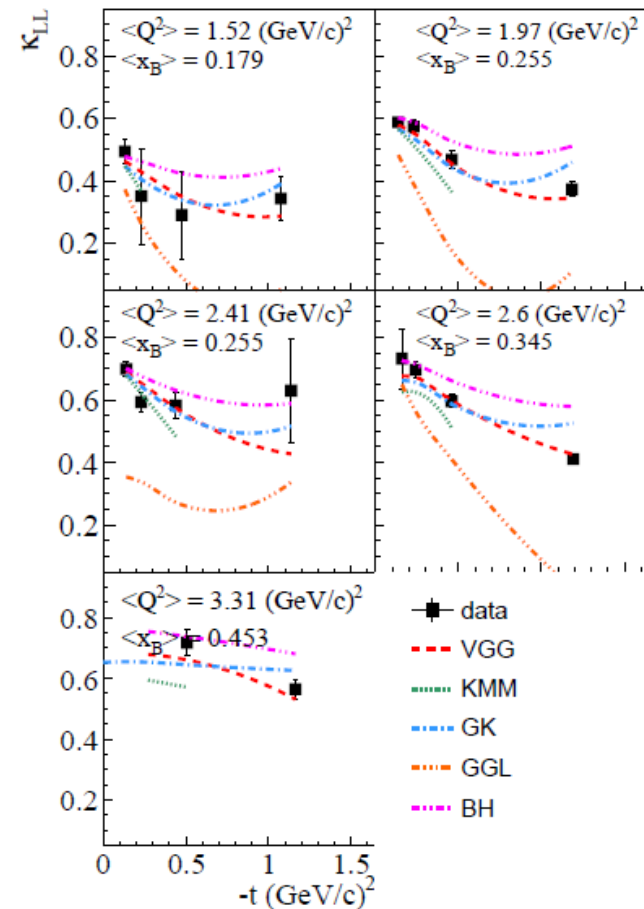
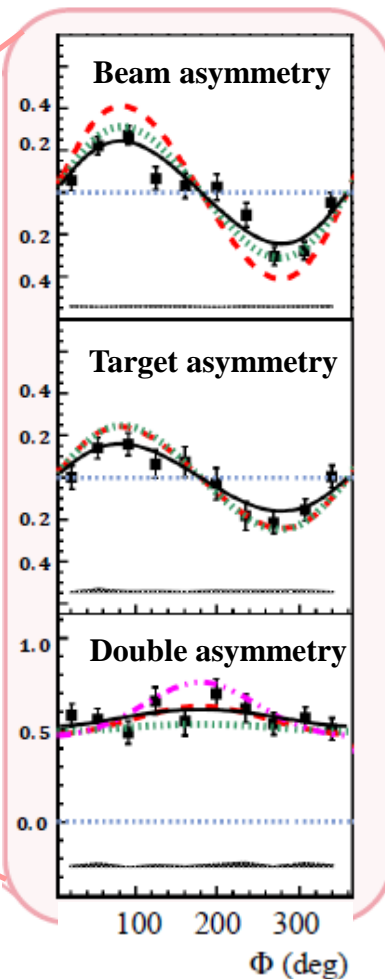
- **eg1-dvcs** experiment
- Beam energy  $\sim 6$  GeV
- CLAS + IC to detect forward photons
- Target: **longitudinally polarized NH<sub>3</sub>** (P $\sim 80\%$ )
- **3 DVCS observables**

$$\vec{e}\vec{p} \rightarrow e\vec{p}\gamma$$

$$A_{LL} \sim \text{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$

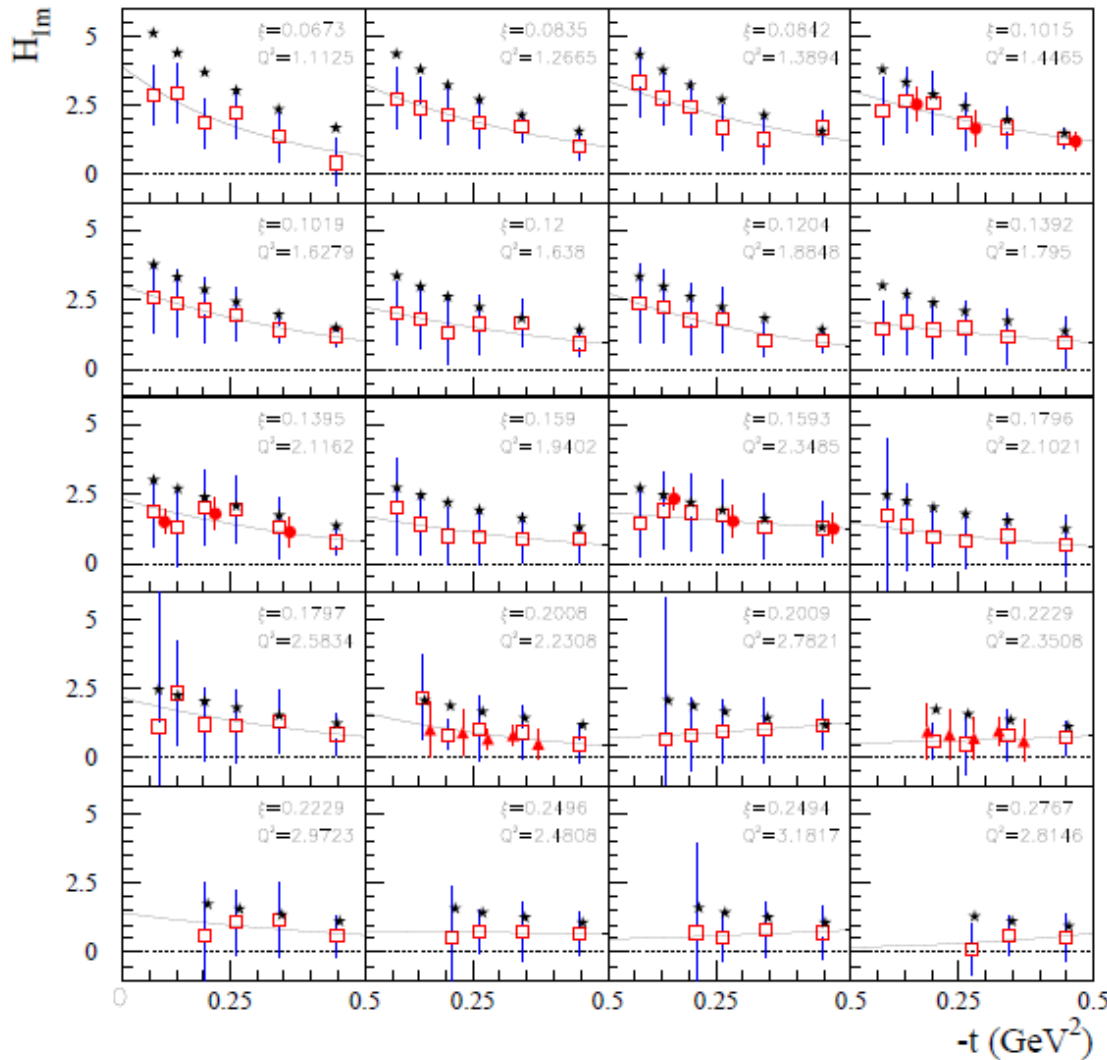


- **5 ( $Q^2$ - $x_B$ ) bins**
- **4  $t$  bins**
- **10  $\phi$  bins**



S. Pisano *et al.* (CLAS Collaboration),  
Phys. Rev. D 91, 052014 (2015)

# Extraction of $H_{Im}$ from the fits of Jefferson Lab 6 GeV data



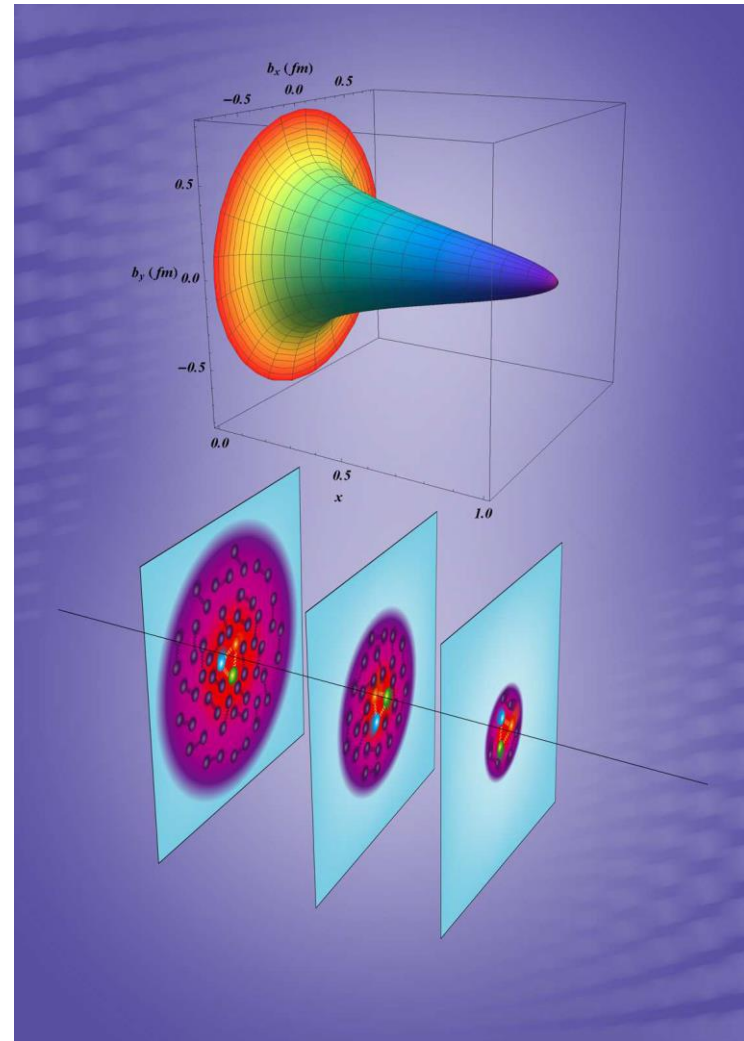
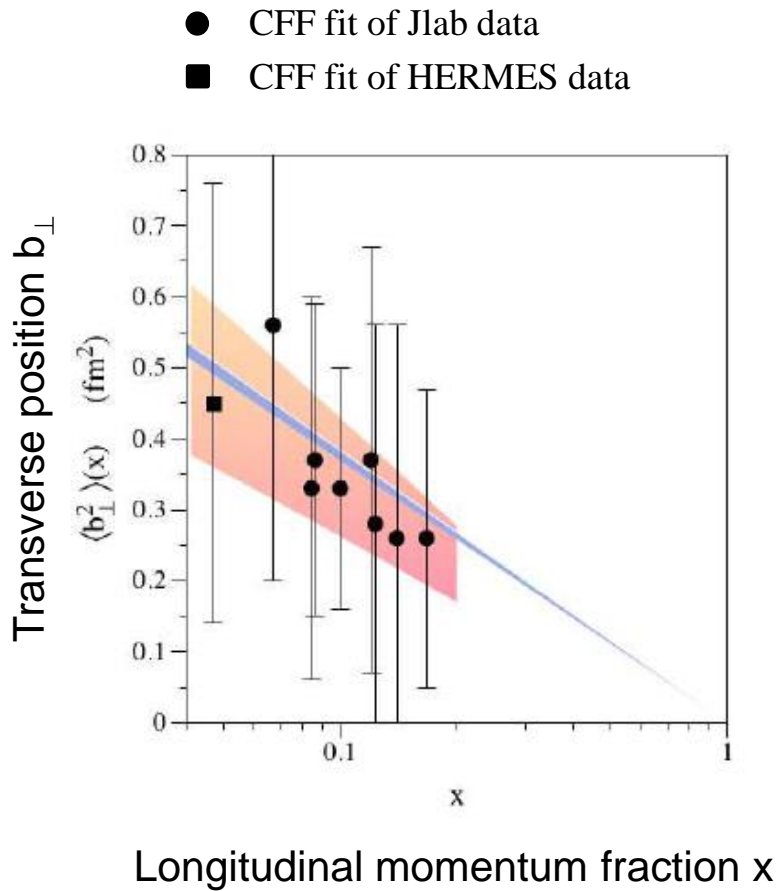
Fits in each  $(Q^2, x_B, t)$  bin of data

$$H_{Im}(\xi, t) = A(\xi)e^{b(\xi)t}$$

$$\xi \approx \frac{x_B}{2 - x_B}$$



# From CFFs to proton tomography

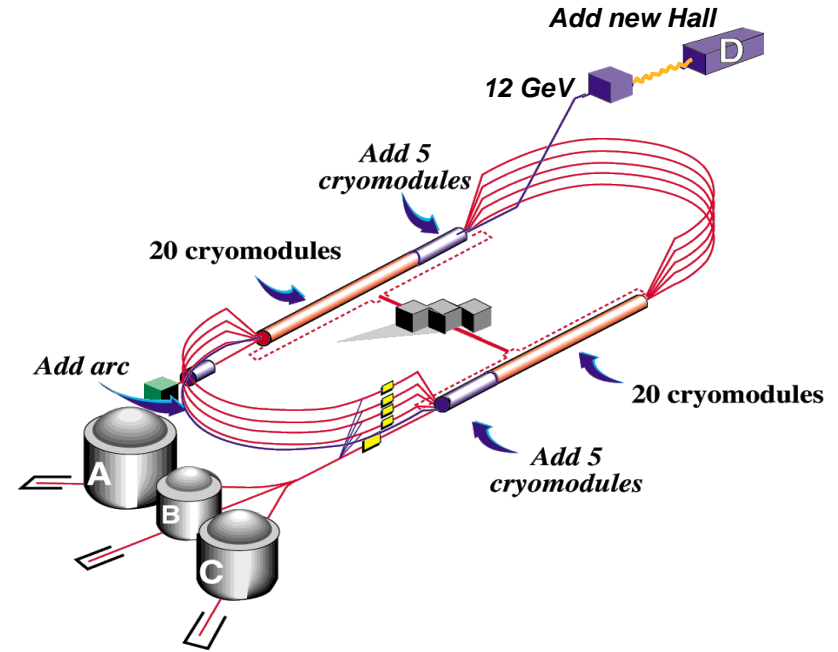
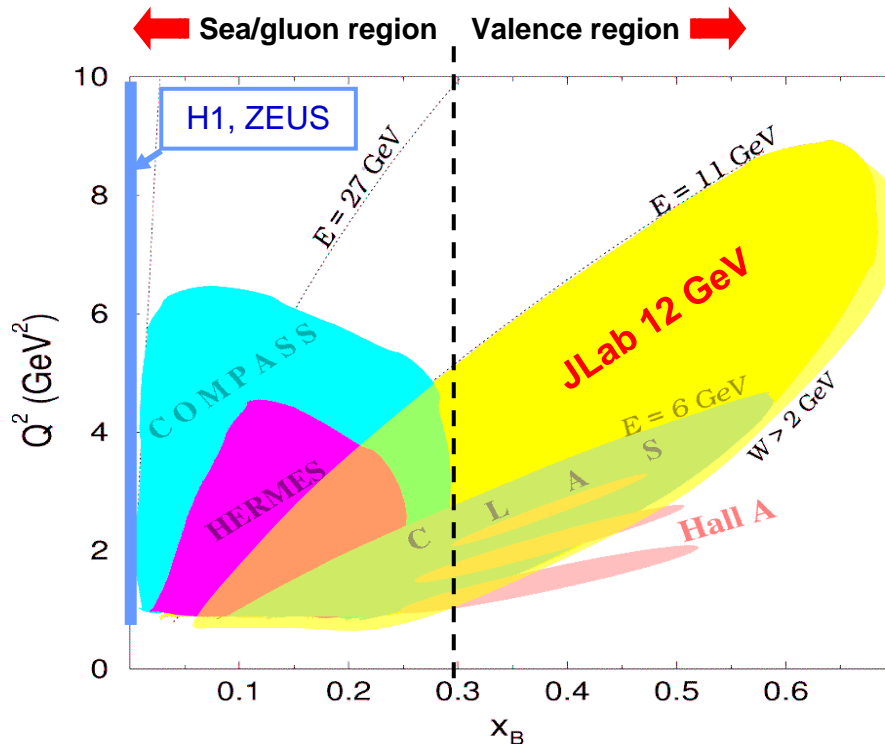


# Jefferson Lab upgrade to 12 GeV

$E = 2.2, 4.4, 6.6, 8.8, 11$  GeV  
for the Halls A, B, C

Beam polarization  $> 80\%$

Accelerator 12 GeV upgrade

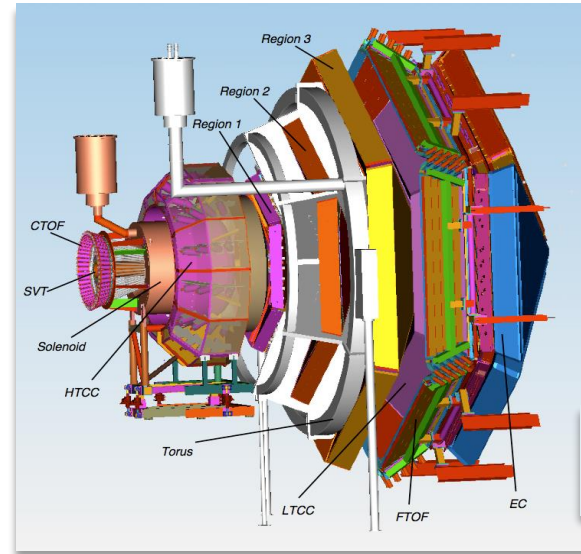
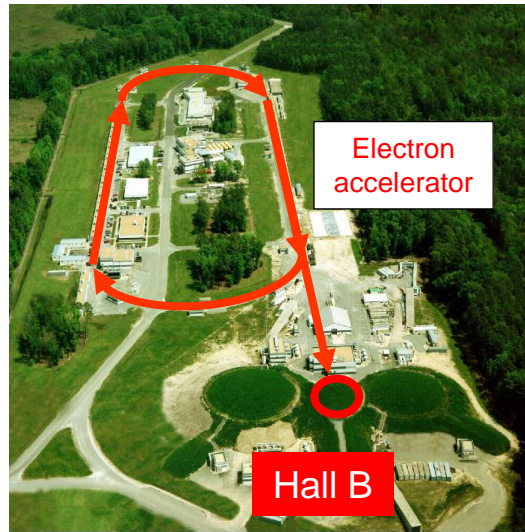


Study of high  $x_B$  domain  
requires high luminosity

The 12-GeV upgrade is  
well matched to studies in  
the valence-quark regime

# Jefferson Lab 12 GeV and the new CLAS12 detector

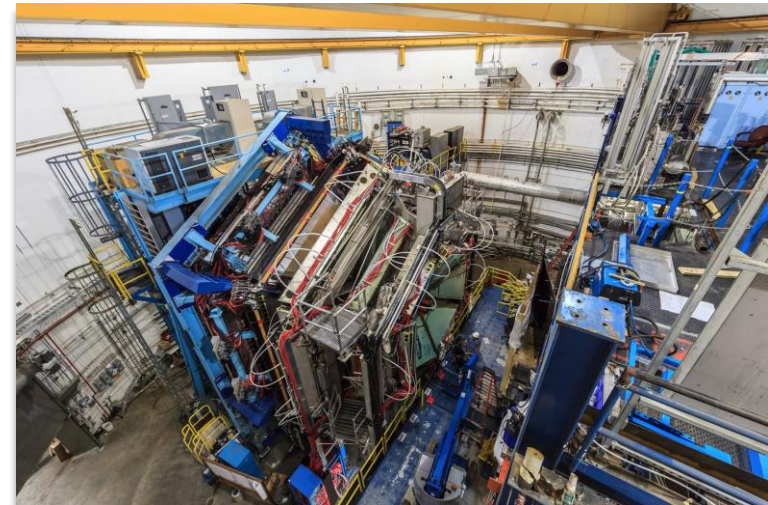
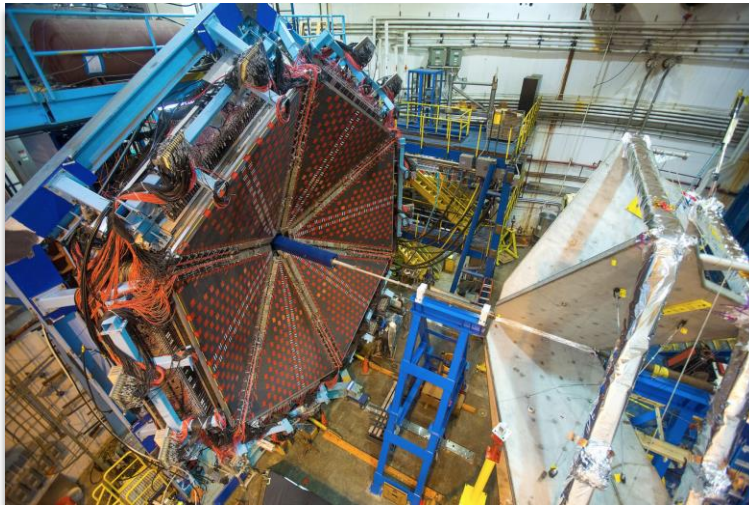
Jefferson Lab



CLAS12  
detector  
in Hall B

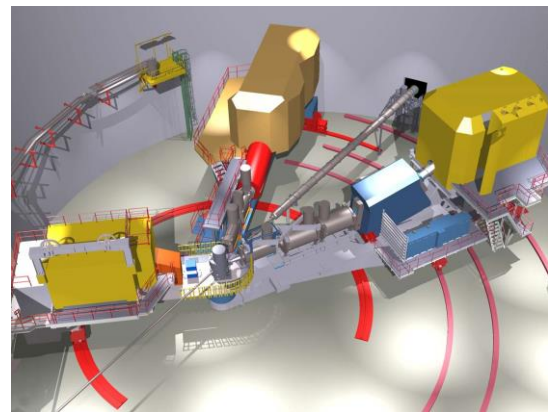
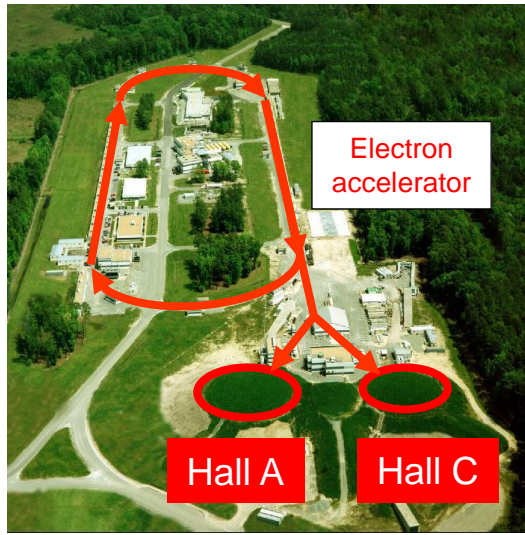
Design luminosity  
 $L \sim 10^{35} \text{ cm}^{-2}\text{s}^{-1}$

Data taking with the new CLAS12 detector started in 2018

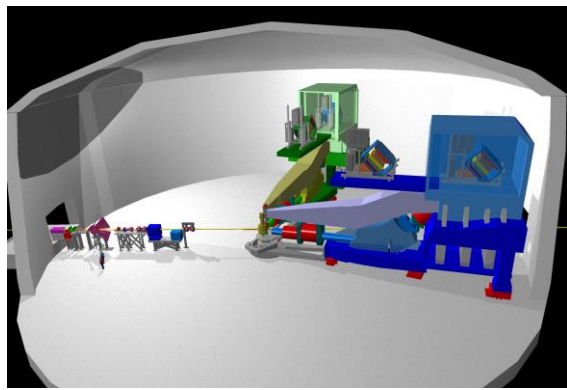


# DVCS experiments in Hall A and Hall C of Jefferson Lab

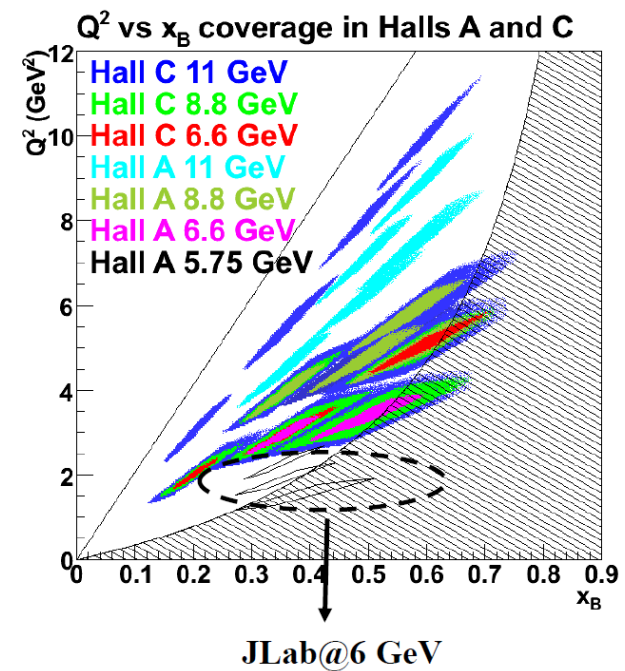
Jefferson Lab



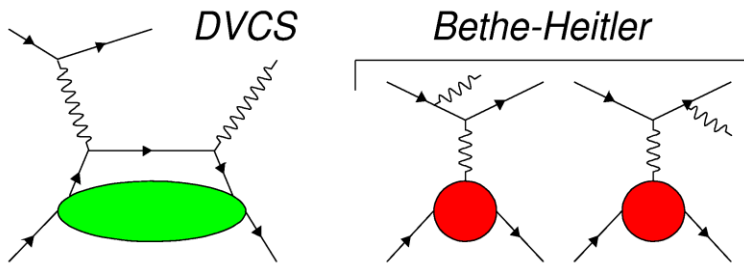
Hall C



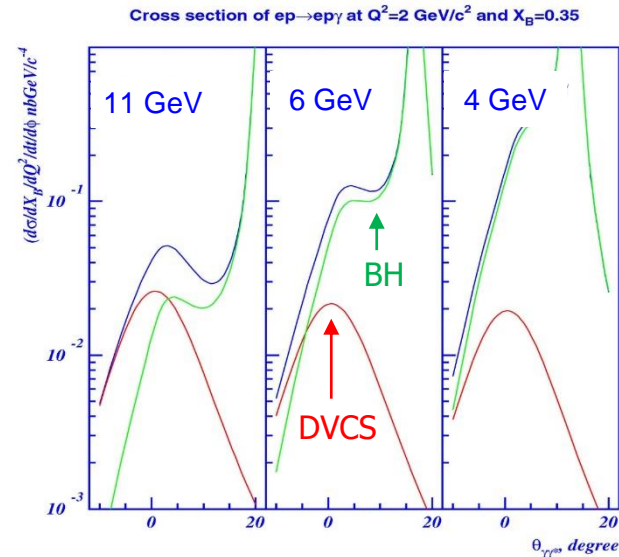
Hall A



# Projected results for CFFs with CLAS12



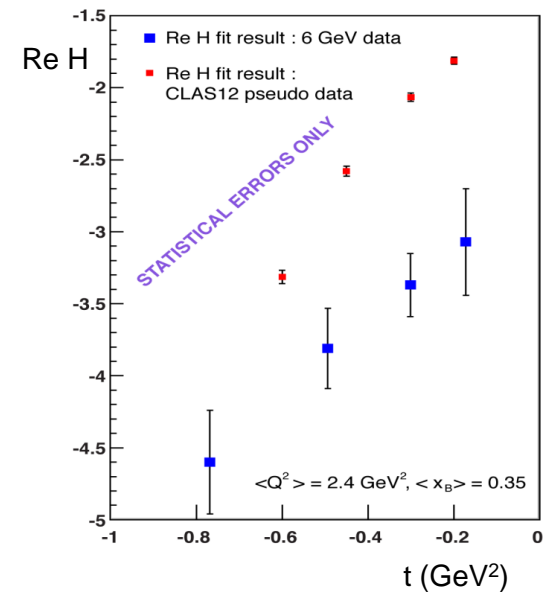
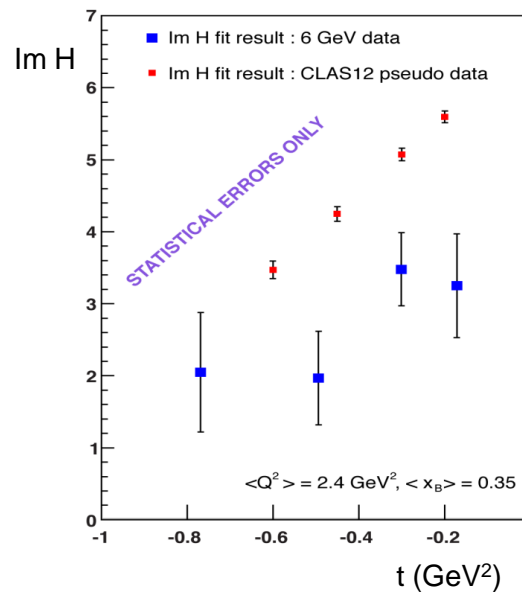
The contribution ratio from BH and DVCS to the  $ep \rightarrow e\gamma$  cross section changes with the electron beam energy



Impact of CLAS12 DVCS pseudo data on CFF fit

CLAS 6 GeV data

CLAS12 pseudo data



# First look at DVCS $A_{LU}$ with CLAS12 preliminary data

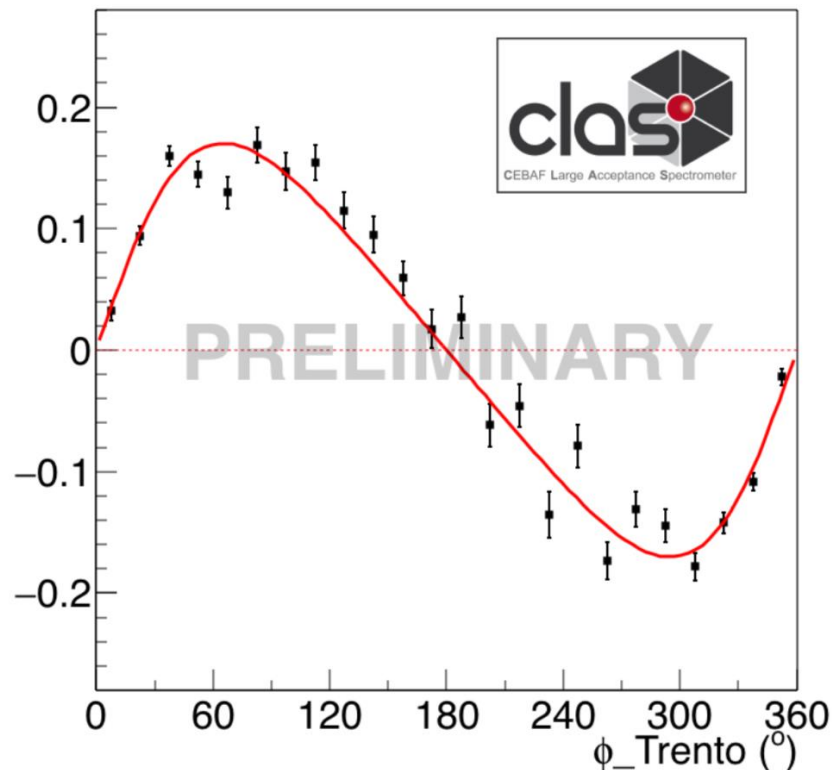
$$A_{LU} = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-}$$

## Beam spin asymmetry

$$A_{LU} = \frac{1}{P} \frac{N^+(\phi_{Trento}) - N^-(\phi_{Trento})}{N^+(\phi_{Trento}) + N^-(\phi_{Trento})}$$

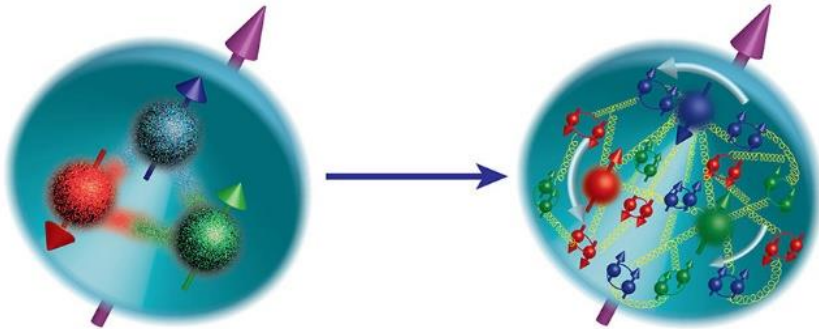
- $P$ : electron polarization
- $N^{+(-)}$ : number of photon electroproduction candidates with beam helicity  $+(-)$
- Analysis based on 2% of approved beam time

## Raw Beam-Spin Asymmetry $ep \rightarrow e\gamma$

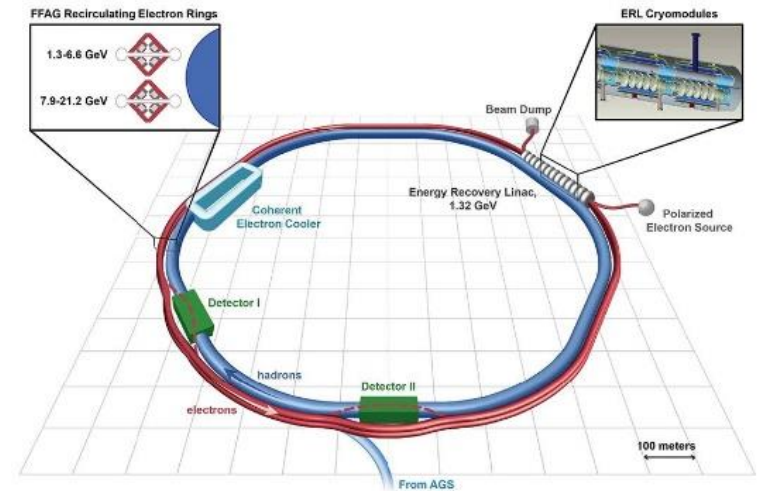


# DVCS at the Electron-Ion Collider (EIC)

Nucleon tomography of the gluons and sea quarks (low momentum fraction  $x$ )

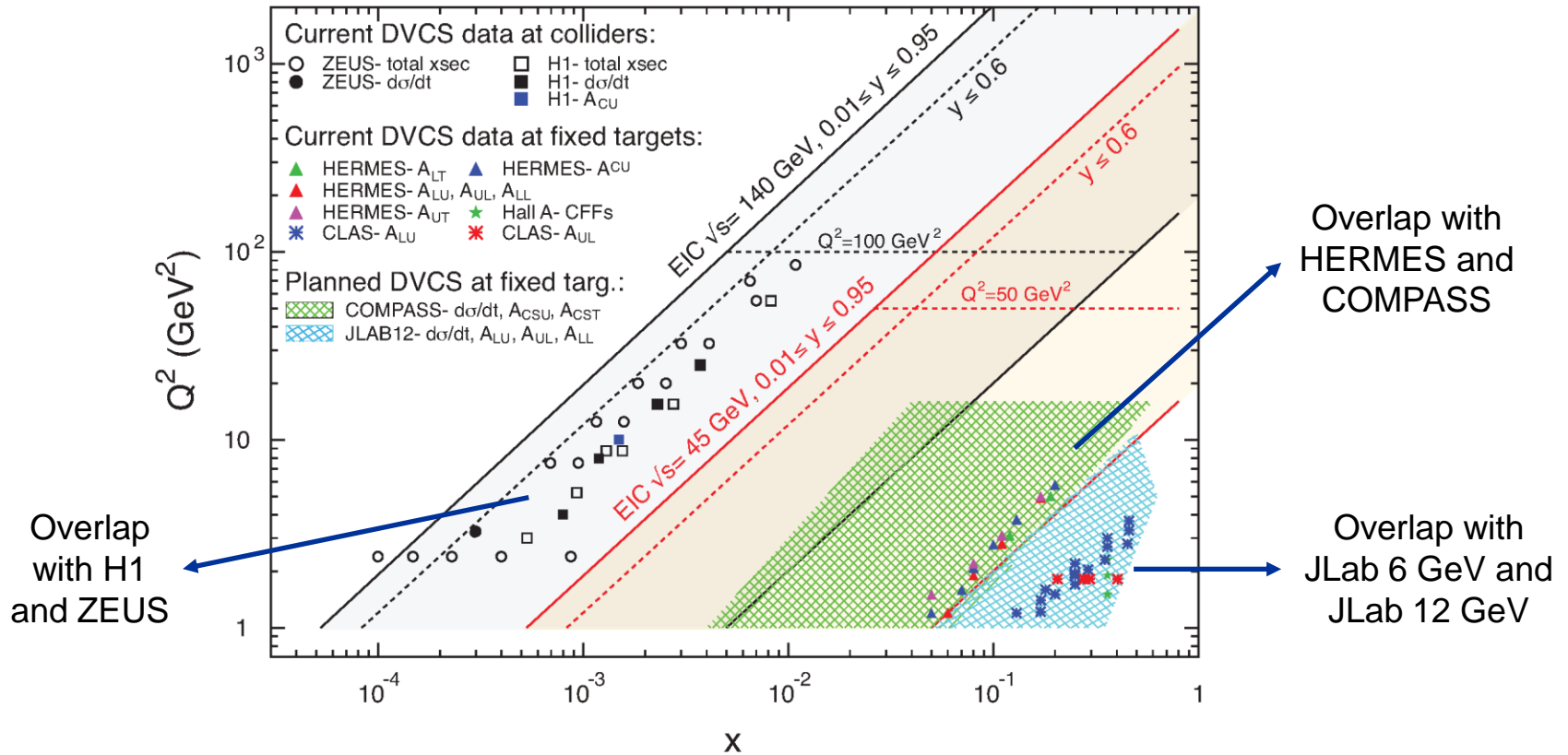


The Electron-Ion Collider (EIC) will be constructed at BNL



eRHIC design (BNL)

# DVCS at the EIC : gluons and sea quarks

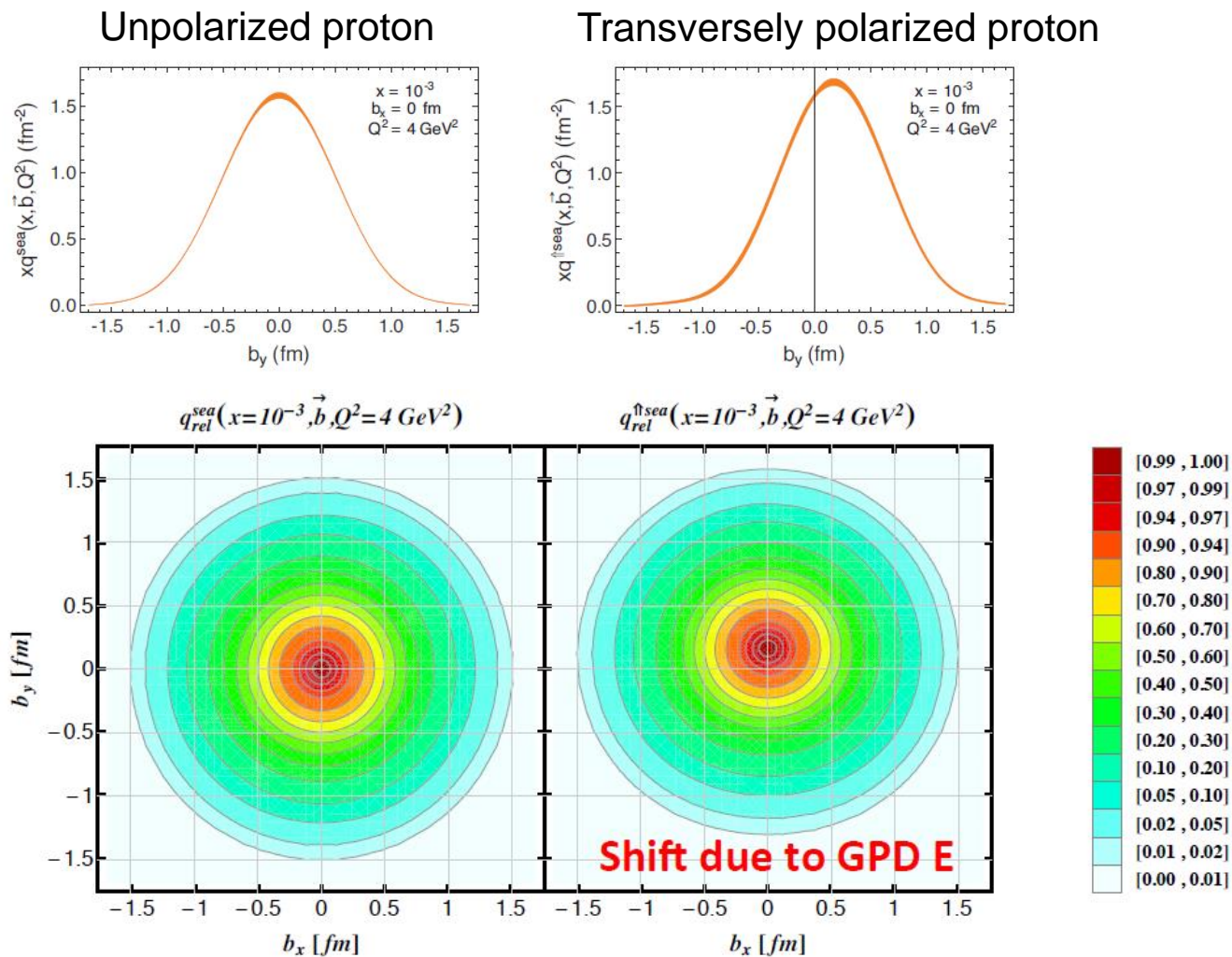


- Collision of polarized electrons with polarized protons, light and heavy nuclei
- High Luminosity :  $L_{ep} \geq 10^{33-34} \text{ cm}^{-2} \text{ s}^{-1}$  (100-1000 times HERA)



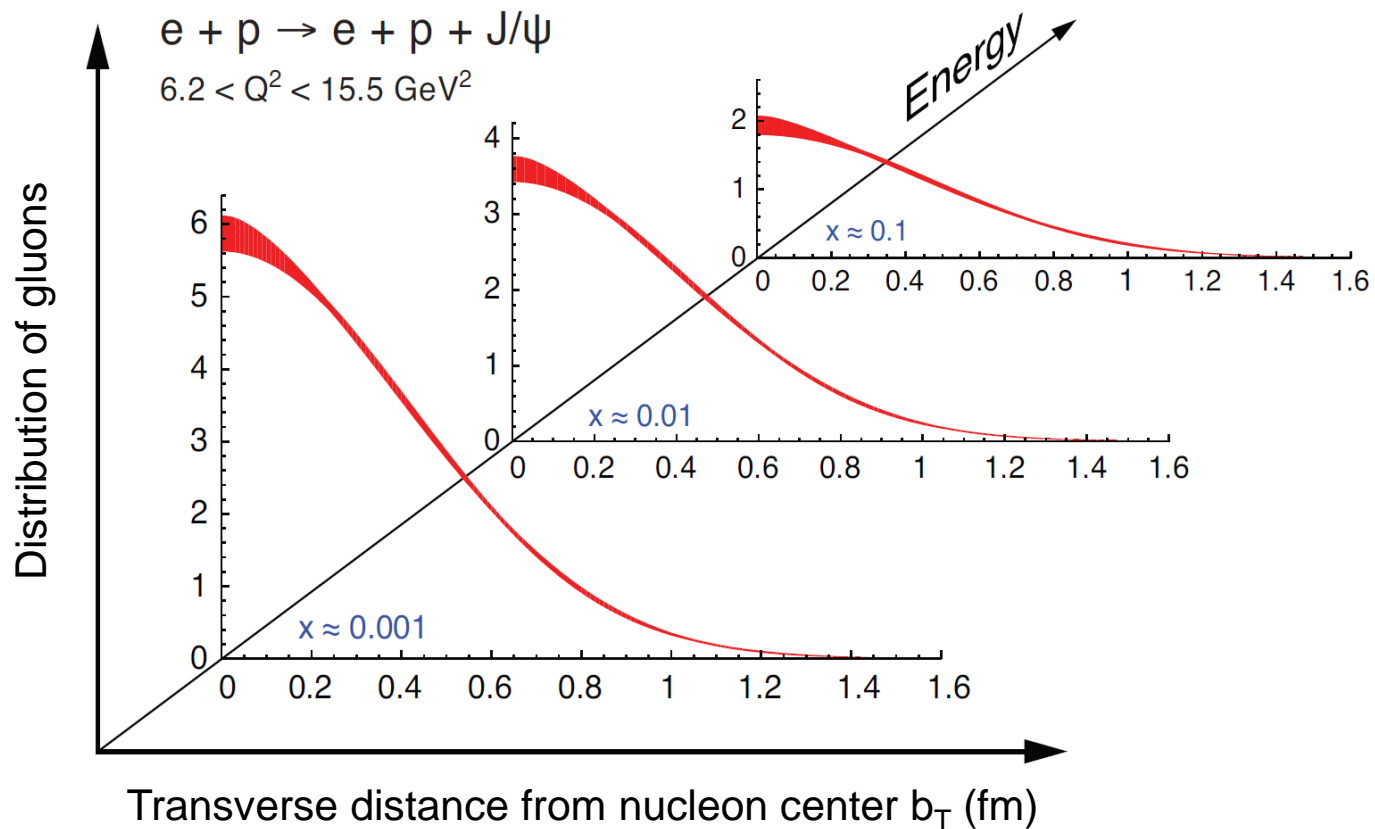
# Spatial distributions of sea quarks from EIC DVCS pseudo data

Spatial distributions obtained from GPD fits to EIC simulated DVCS data



# Gluon distributions from EIC $J/\psi$ pseudo-data

Projected gluon distributions vs transverse distance from the center of the nucleon from EIC pseudo-data of  $J/\psi$  production cross section



# Overview

---

- DVCS data from Jefferson Lab 6 GeV were used to extract a first experimental result of nucleon tomography.
- Jefferson Lab 12 GeV will provide high-precision data covering a large unexplored kinematic domain at high  $x$ .
- While Jefferson Lab is a unique facility to study the valence quarks, the future Electron-Ion Collider (EIC) will provide high-precision GPD measurements at low  $x$ , allowing us to perform nucleon tomography of the gluons and sea quarks.

Thank you