## Isospin extrapolation as a method to study inclusive b->sll decays

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Isospin extrapolation as a method to study inclusive $\bar{B} \rightarrow X_{s} \ell^{+} \ell^{-}$decays
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 approach. Several experimental advantages are seen which have the potential to make measurements of inclusive $\bar{B} \rightarrow X_{s} \ell^{+} \ell^{-}$decays tractable at a hadron collider.

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## b->sll decays

- b->sll decays are loop suppressed semileptonic decays. Their loop suppression allows for NP sensitivity up to $\sim 50 \mathrm{TeV}$.
- They have been part of LHCb's core program for years.
$b$

- Focus has been on exclusive decays, whereby the strange quark hadronises into a specific final state.



- Exclusive decays are fully reconstructed $->$ signal peaks at the $B$ mass.


## Where are we

- Two sets of deviations with the interpretation limited either by theory or statistics.

Limited by statistical uncertainties



Limited by theoretical interpretation


- Inclusive $\mathrm{b}->$ sll measurements offer a way forward for both these limitations.


## Inclusive b—>sll decays

- Instead of reconstructing a specific hadronic final state, allow the strange quark to hadronise whatever it likes
b

- Inclusive decays have complementary (and generally more precise) theoretical uncertainties compared to exclusive ones.


- For the branching fraction, uncertainty saturated by experimental uncertainties rather than theoretical ones.


## Methods to study inclusive b->sll decays

- Inclusive $b->$ sll decays have been the domain of the B-factories.
- They employ a sum-of-exclusives approach:
- Reconstruct as many exclusive final states as possible (typically $50 \%$ coverage).
- Extrapolate missing modes using a hadronisation model (e.g. with JETSET).

| Belle, Phys. Rev. D 93, 032008 (2016) |  |  |  |
| :--- | :--- | :--- | :---: |
| $\bar{B}^{0}$ decays |  | $B^{-}$decays |  |
|  | $\left(K_{S}^{0}\right)$ | $K^{-}$ |  |
| $K^{-} \pi^{+}$ | $\left(K_{S}^{0} \pi^{0}\right)$ | $K^{-} \pi^{0}$ |  |

- For Belle-Il a fully inclusive approach, whereby only the two leptons are reconstructed, is also foreseen.
- This has no systematic uncertainty associated with the extrapolation, but suffers from larger background.


## Our approach

- Our approach is to reconstruct an additional charged kaon in addition to the two leptons.

- This can be seen as a hybrid of the fully inclusive and sum-of-exclusives modes.
- Still needs an extrapolation, but hopefully cleaner (or at least complementary) than from a sum-ofexclusives method.
- We are not claiming to have invented isospin extrapolation here. This has been used to fill in some gaps in the sum-of-exclusives method. We instead promote this to the main extrapolation of the analysis.


## Fast simulation

- To explore some experimental advantages, generate some fast simulation with RapidSim. axdiv. 112.07899
- B-hadrons produced with kinematics expected within the LHCb acceptance.
- Smearing to account for reconstruction.
- We generate two exclusive channels as a proxy for inclusive decays.
- $B^{+}->K^{+} \pi^{-} \pi^{+} \mu^{+} \mu^{-}$
- $\mathrm{B}^{+}->\mathrm{K}^{+} \pi^{-} \mu^{+} \mu^{-}$
- In both cases the pions are missing from the visible signature.
- Apply $\mu p_{T}>300 \mathrm{MeV} / \mathrm{c}$ to account for trigger effects in run III.


## Background to these decays

- There are two main backgrounds to an inclusive analysis:
- Combinatorial, whereby accidental combinations of different B/D decays are made.
- Double-semileptonic: $B \rightarrow\left(D \rightarrow K^{-} \ell^{+} \nu_{\ell} X\right) \ell^{-} \nu_{\ell} X$

- Combinatorial is easier to distinguish but less well understood.


## The sideband

- Combinatorial background is extrapolation into signal region using a sideband (above the B mass).

$K \mu \mu$ visible mass

- Signals are substantially closer to the sideband than in a fully inclusive approach.


## The mass of the strange hadron, $\mathrm{m}_{\mathrm{x}}$

- An important discriminating variable is the mass of the strange hadron.
- Also selected to reduce background in sum-of-exclusives analyses.
- If we use the rest frame approximation [1] to calculate $m x_{s}$, see an improvement

$$
\left(p_{B}\right)_{z}=\frac{m_{B}}{m_{\text {reco }}}\left(p_{\text {reco }}\right)_{z}
$$



- KIl signature has better resolution on the mass c.f. fully inclusive approach.


## Other advantages

- Other advantages include:
- A better defined vertex (three tracks instead of two).
P. Alvarez-Cartelle, W. Altmannshofer, Beyond the flavour anomalies II workshop
- Flavour tagged for $A_{C P} / A_{\text {fB }}$ measurements.
- Access to opposite sign $\mathrm{m}_{\mathrm{KI}}$ combination.

- In order to fully understand the advantages and remaining level of background, a detailed study with full simulation would be required (beyond the scope of the paper).


## Comments on the extrapolation

- The extrapolation boils down to calculating the fraction of inclusive $\mathrm{b}->$ sll that produce a charged kaon.


- For $\mathrm{B}^{0}$ and $\mathrm{B}^{+}$decays, each $\mathrm{b}->$ sll decay is expected to either a charged or neutral kaon.
- Extrapolation is then done using isospin rules (naively expected to be around 50\%),
- Of course, we do not only produce $\mathrm{B}^{0}$ and $\mathrm{B}^{+}$mesons at the LHC.


## The complication from $B_{s}{ }^{0}$ and $\Lambda_{b}{ }^{0}$ hadrons

- Naively, isospin extrapolation should account for neutral kaons nicely for both $B_{s}{ }^{0}$ and $\Lambda_{b}{ }^{0}$ hadrons.

- Problem is that there is a fraction of inclusive $B_{s}{ }^{0}$ and $\Lambda_{b}{ }^{0}$ decays which do not produce kaons. This fraction is unknown and extrapolation appears difficult.

|  | $\boldsymbol{B}_{s}^{0}$ | $I\left(J^{P}\right)=0\left(0^{-}\right)$ |  |
| :---: | :---: | :---: | :---: |
| $\Gamma_{51}$ | $J / \psi(1 S) \phi$ |  | $(1.08 \pm 0.08) \times 10^{-3}$ |
| $\Gamma_{52}$ | $J / \psi(1 S) \phi \phi$ |  | ${ }^{(1.24-2.109)} \times 10^{-5}$ |
| $\Gamma_{53}$ | $J / \psi(1 S) \pi^{0}$ |  | $<1.2 \times 10^{-3}$ |
| $\Gamma_{54}$ | $J / \psi(15) \eta$ |  | $(4.0 \pm 0.7) \times 10^{-4}$ |
| $\Gamma_{55}$ | $J / \psi(1 S) K_{s}^{0}$ |  | $(1.92 \pm 0.14) \times 10^{-5}$ |
| $\Gamma_{56}$ | $J / \psi(1 S) \vec{K}^{*}(892)^{0}$ |  | $(4.1 \pm 0.4) \times 10^{-5}$ |
| $\Gamma_{57}$ | $J^{\prime \prime} \psi(1 S) \eta^{\prime}$ |  | $(3.3 \pm 0.4) \times 10^{-4}$ |
| $\Gamma_{58}$ | $J / \psi(11) \pi^{+} \pi^{-}$ |  | $(2.09 \pm 0.23) \times 10^{-4}$ |


| Decay Modes |  | $I\left(J^{P}\right)=0\left(3 / 2^{-}\right)$ |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |
| Mode |  | Fraction ( $\left.\Gamma_{i} / \Gamma\right)$ | Scale Factor/ Conf. Level | $\begin{aligned} & P \\ & (\mathrm{MeV} / \mathrm{c}) \end{aligned}$ |
| $\Gamma_{1}$ | $N \bar{K}$ | 20-30\% |  | 433 |
| $\Gamma_{2}$ | $\Sigma \pi$ | 20-40\% |  | 410 |
| $\Gamma_{3}$ | A $\sigma$ | $(5.0 \pm 2.0) \%$ |  |  |
| $\Gamma_{4}$ | Aлn | $\sim 25 \%$ |  | 419 |

## $B_{s}{ }^{0}$ and $\Lambda_{b}{ }^{0}$ hadron decays as background

- These decays have smaller production and branching fractions than the $\mathrm{B}^{+}$and $\mathrm{B}^{0}$ decays.

- We therefore propose to treat them as background and subtract them for the branching fraction.
- Dedicated auxiliary measurements can be useful:
- For $\mathrm{B}_{\mathrm{s}}: \mathcal{B}\left(B_{s} \rightarrow K^{+} K^{-} X \ell^{+} \ell^{-}\right)$
- For $\Lambda_{b}: \mathcal{B}\left(\Lambda_{b}^{0} \rightarrow p K^{-} X \ell^{+} \ell^{-}\right)$
- It is clear that a resulting systematic uncertainty will arise from this.


## Prospects for theoretically precise observables

- Of course none of these matters for observables which are either reliably zero in the SM (AcP) or hadronic uncertainties cancel (LFU ratios).
- In this case, missing an unknown fraction of the inclusive decay does not spoil the comparison with the SM.


- Here we note the fact that the inclusive BF is around order of magnitude higher than exclusive channels.
- By the end of run III, we expect around $1 \mathrm{M} X_{b} \rightarrow K^{+} X \ell^{+} \ell^{-}$candidates(!!).
- Due to the low reconstruction efficiency at LHCb, expect any LFU/Acp measurements to be statistically independent than exclusive ones (e.g. $\mathrm{R}_{\mathrm{k}}$ ).
- Can afford to be brutal with the selection and still have a large signal yield.


## Summary

- We propose to use the signature $X_{b} \rightarrow K^{+} X \ell^{+} \ell^{-}$as a proxy for inclusive b->sll decays.
- Several experimental advantages are expected with respect to a fully inclusive approach.
- Sideband closer to the signal - easier extrapolation for combinatorial background.
- Extrapolation complicated at LHC by presence of $\mathrm{B}_{\mathrm{s}}$ and $\Lambda_{\mathrm{b}}{ }^{0}$ hadrons - propose to treat them as background.
- Expect the largest sample of self-tagged $b->$ sll decays in the world with this method - could provide statistically independent measurements of clean observables.

