

On the difference between pp and $ppbar$ cross sections

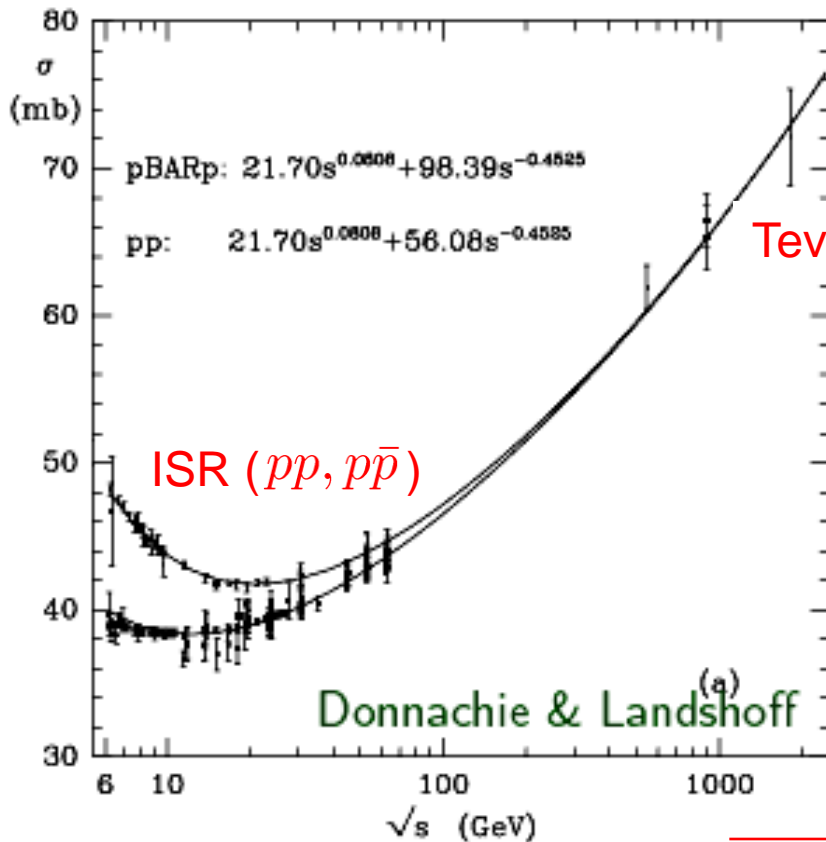
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JHEP 1003:037 (2010); 0912.3806 [hep-th]

with Emil Avsar and Toshihiro Matsuo

The total cross section difference

Consider the difference $\Delta\sigma \equiv \sigma_{tot}^{p\bar{p}} - \sigma_{tot}^{pp}$



What's the fate of $\Delta\sigma$
as $s \rightarrow \infty$?

Up to the ISR energy, ($\sqrt{s} = 53\text{GeV}$)

$$\Delta\sigma \sim s^{-0.5} > 0$$

→
LHC(pp), cosmic rays

First data from LHC



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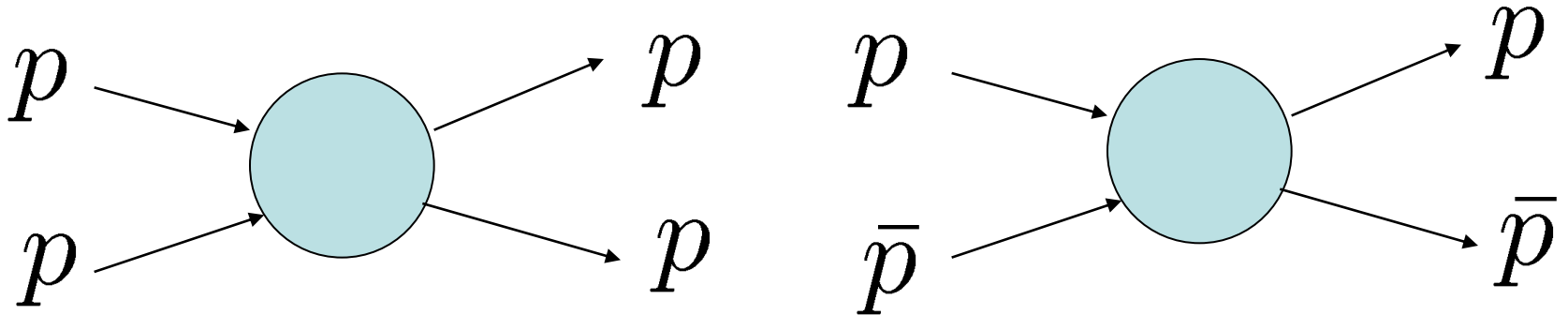
Charged-particle multiplicities in pp interactions at $\sqrt{s} = 900$ GeV measured with the ATLAS detector at the LHC $\star, \star\star$

A B S T R A C T

The first measurements from proton-proton collisions recorded with the ATLAS detector at the LHC are presented. Data were collected in December 2009 using a minimum-bias trigger during collisions at a centre-of-mass energy of 900 GeV. The charged-particle multiplicity, its dependence on transverse momentum and pseudorapidity, and the relationship between mean transverse momentum and charged-particle multiplicity are measured for events with at least one charged particle in the kinematic range $|\eta| < 2.5$ and $p_T > 500$ MeV. The measurements are compared to Monte Carlo models of proton-proton collisions and to results from other experiments at the same centre-of-mass energy. The charged-particle multiplicity per event and unit of pseudorapidity at $\eta = 0$ is measured to be $1.333 \pm 0.003(\text{stat.}) \pm 0.040(\text{syst.})$, which is 5–15% higher than the Monte Carlo models predict.

Discrepancy between the data and MC tuned to Tevatron and CERN $p\bar{p}$ data

Odderon and Reggeon



The difference $\Delta\mathcal{A} \equiv \mathcal{A}_{pp \rightarrow pp}(s, t) - \mathcal{A}_{p\bar{p} \rightarrow p\bar{p}}(s, t)$

is odd under crossing, generated by the exchange of **C-odd** objects in QCD

\longrightarrow
 $\left\{ \begin{array}{l} \text{Reggeon (vector mesons)} \\ \text{Odderon (C-odd glueballs)} \end{array} \right.$
Lukaszuk, Nicolescu (1973)

$$\Delta\sigma(s) = \frac{1}{s} \text{Im}\Delta\mathcal{A}(s, t=0) \sim s^{-0.5}$$

Attributed to the Reggeon exchange. Any room for the Odderon?

The AdS/CFT correspondence

Maldacena (1997)

Conjecture: N=4 SYM at **strong** coupling is dual to type IIB at **weak** coupling on $AdS_5 \times S^5$

$$ds^2 = R^2 \overbrace{\frac{-dx^\mu dx_\mu + dz^2}{z^2}}^{\text{our universe}} + R^2 \overbrace{d\Omega_5^2}^{\text{extra dimensions}}$$

N=4 SYM

string

(anomalous) dimension



mass

\lambda



curvature radius R^4/α'^2

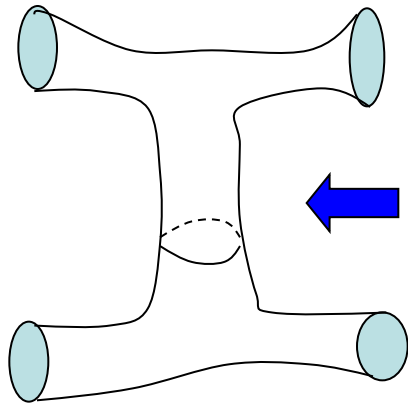
number of colors $1/N_c$



string coupling constant g_s

High energy scattering in AdS

Regge trajectory in string theory $m^2 = \frac{4n}{\alpha'} = \frac{2(j-2)}{\alpha'}$



$g_{\mu\nu}, \phi, A_\mu, B_{\mu\nu}, \dots$

graviton, dilaton, SO(6) gauge boson, **B-field**....

Total cross section \rightarrow

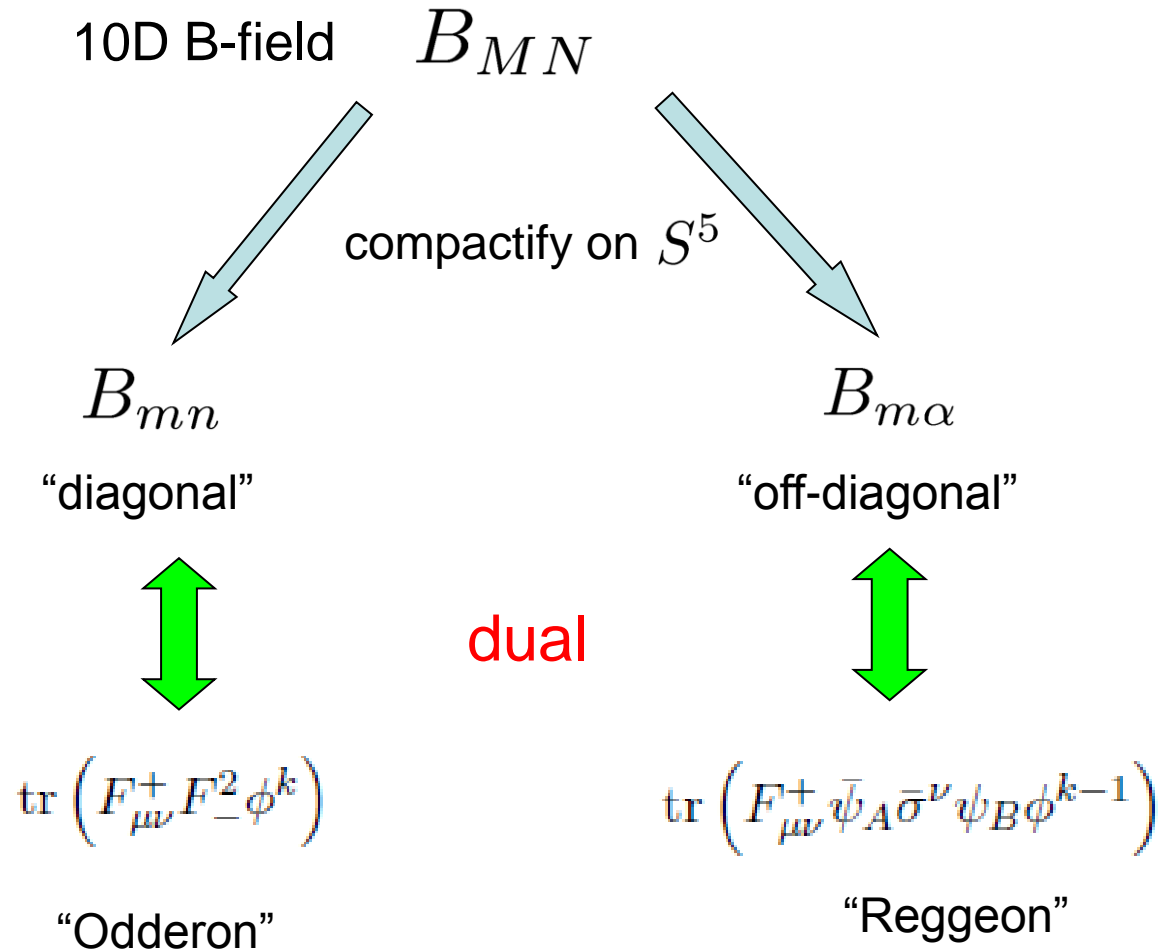
Reggeized graviton = **Pomeron**

Total cross section difference \rightarrow

Reggeized B-field = **Odderon**

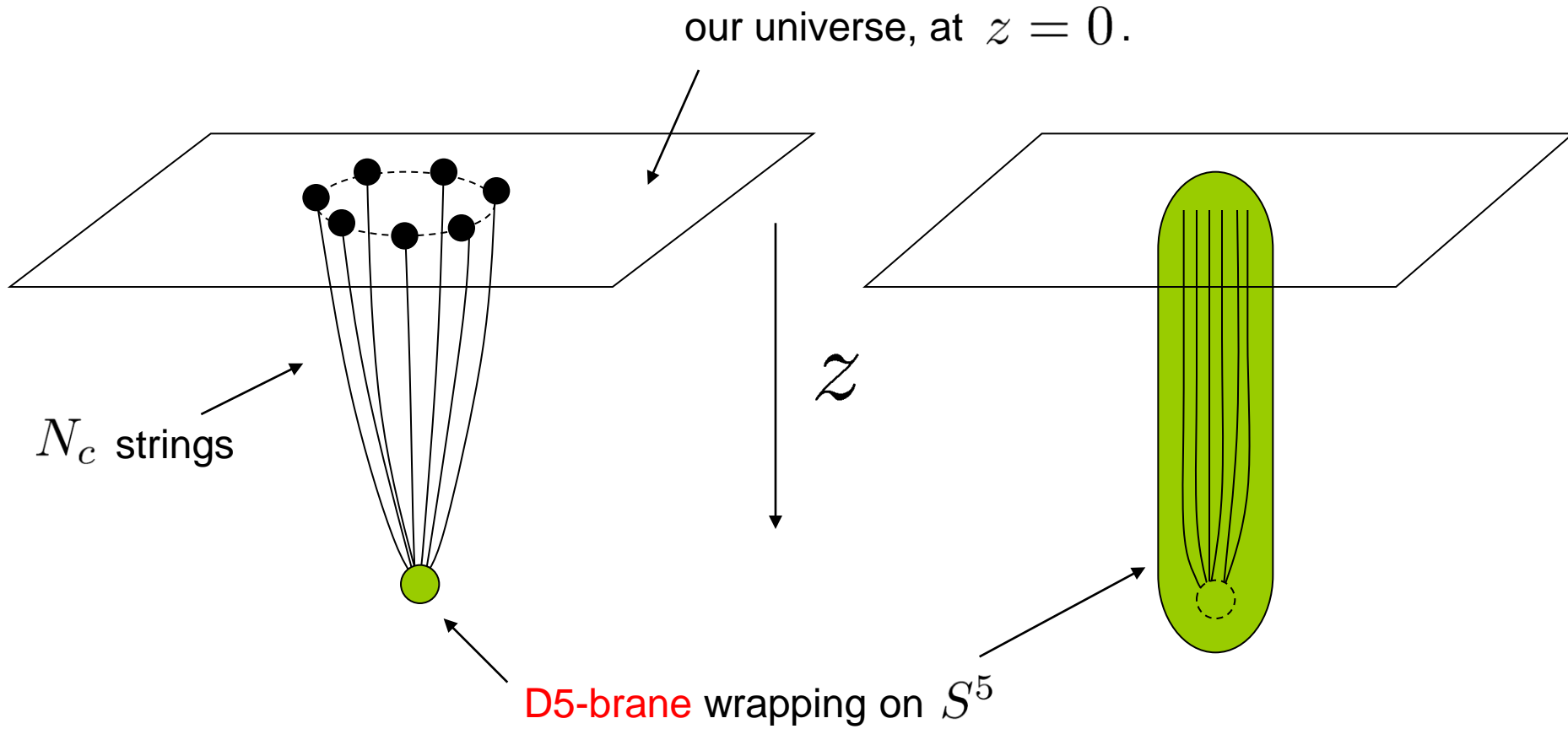
Brower, Djuric, Tan (2009)

The antisymmetric B-field



Odderon and Reggeon unified in 10 dimensions

Baryons in AdS/CFT



Witten (1998)

Imamura (1999)
Callan, Guijosa, Savvidy (1999)

Baryon-Odderon coupling

Born-Infeld and Chern-Simons action

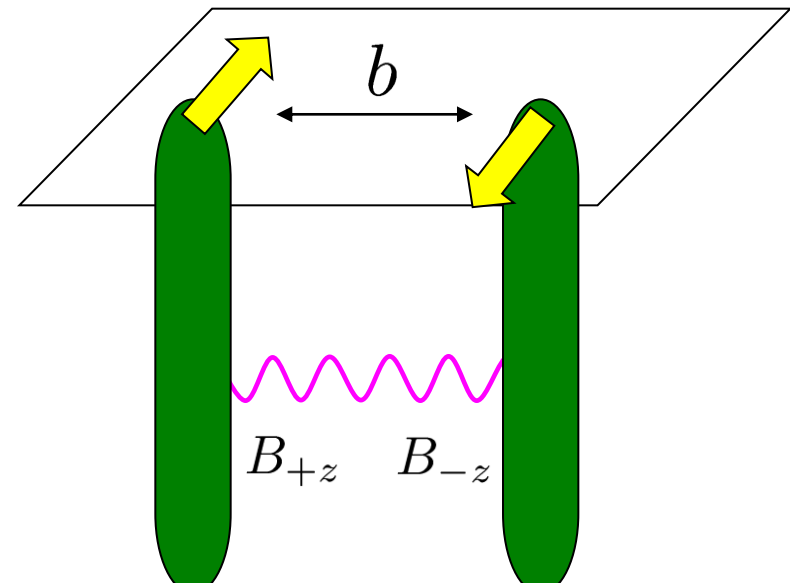
$$S = T_5 \int d^6\xi \left\{ -\sqrt{-\det(\tilde{G} + \tilde{B} + 2\pi\alpha'F)} + 2\pi\alpha'F \wedge c_{(4)} \right\}$$

$$= \frac{1}{\text{Vol}_{S^4}} \int dx^+ d\theta d\Omega_4 \frac{\partial \mathcal{L}}{\partial F_{+\theta}} \frac{\tilde{B}_{+\theta}}{2\pi\alpha'}$$

$$= \frac{\pm N_c}{2\pi\alpha' \text{Vol}_{S^4}} \int dx^+ d\theta d\Omega_4 \left(B_{+\theta} + \frac{\partial z}{\partial \theta} B_{+z} \right)$$

Reggeon
(decouple)

Odderon



Odderon amplitude, fixed impact parameter

$$i\mathcal{A}^\pm(s, b) = \pm i^2 \left(\frac{N_c}{2\pi\alpha' \text{Vol}_{S^4}} \right)^2 \sum_k \int dx^+ dz d\Omega_4 Y^{(k)}(\Omega)$$

$$\times \int dx'^- dz' d\Omega'_4 Y^{(k)}(\Omega') \langle B_{+z}^{(k)}(x^+, 0, b, z) B_{-z'}^{(k)}(0, x'^-, 0, z') \rangle$$

B-field propagator

Bena, Nastase, Vaman (2001)

$$\langle B_{mn} B_{m'n'} \rangle = T_{mnm'n'}^1 D(u) - \partial_m V_{n,m'n'} + \partial_n V_{m,m'n'} + T_{mnm'n'}^3 K(u)$$

$$\left\{ \begin{array}{l} T_{mnm'n'}^1 = R^4 (\partial_m \partial_{m'} u \partial_n \partial_{n'} u - \partial_m \partial_{n'} u \partial_n \partial_{m'} u) \\ u = \frac{(z - z')^2 + (x_\perp - x'_\perp)^2 - 2(x^+ - x'^+)(x^- - x'^-)}{2zz'} \\ D(u) \sim \frac{1}{z^2 \partial_z^2 - z \partial_z + 1 - M^2} \end{array} \right. \quad \text{“chordal distance”}$$

Up to this point, the amplitude is purely **real**.

Analytic continuation in spin

Include higher spin string excited states.

Analytically continue the B-field propagator to $j \neq 1$ and sum over j

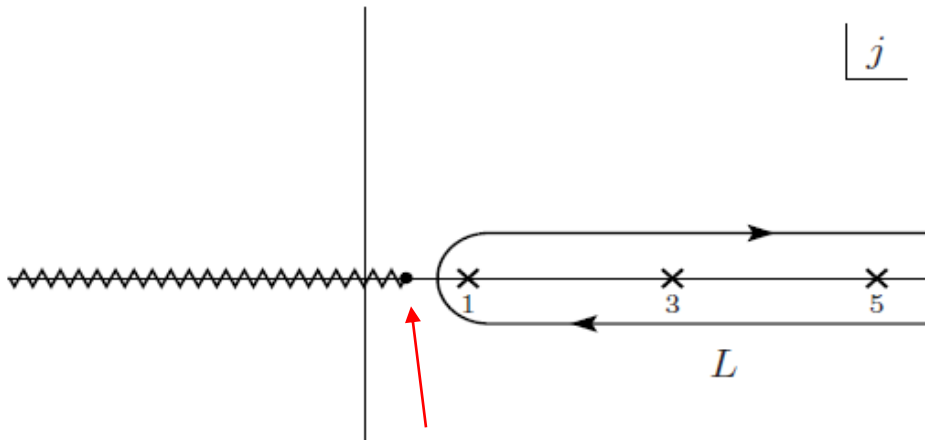
$$D(u) \sim \frac{1}{z^2 \partial_z^2 - z \partial_z + 1 - M^2}$$



j

$$\frac{1}{z^2 \partial_z^2 - z \partial_z + 1 - M_j^2}$$

$$M_j^2 = M^2 + 2\sqrt{\lambda}(j - 1)$$



A cut beginning at $j_0 = 1 - \frac{M^2}{2\sqrt{\lambda}}$

Total cross section difference

$$\Delta\sigma = \sigma^{BB} - \sigma^{\bar{B}\bar{B}} = 2 \int d^2b \operatorname{Im} \mathcal{A}^-(s, b) - 2 \int d^2b \operatorname{Im} \mathcal{A}^+(s, b)$$

$$= -\frac{\pi\sqrt{\lambda}}{4(\operatorname{Vol}_{S^4})^2} \sum_{I,k} \frac{M_I + \frac{1}{M_I}}{k+2} \int dz d\Omega_4 Y^{(k)}(\Omega_5) \int dz' d\Omega'_4 Y^{(k)}(\Omega'_5) \left(\frac{zz's}{4\sqrt{\lambda}} \right)^{\alpha_O(0)-1} .$$

Odderon intercept

$$\alpha_O = j_O + \frac{1}{2\sqrt{\lambda}}$$

Note that $\Delta\sigma$ is **negative** !

...in conflict with the ISR data...

How can $\Delta\sigma$ be negative?

sign of $\Delta\sigma \rightarrow$ sign of $\text{Im}\mathcal{A} \rightarrow$ sign of the interaction

exchange of the B-field \rightarrow **repulsion** $\rightarrow \Delta\sigma > 0??$

However, this may not be true in a **curved space** !

Look at the tensor part of the B-field propagator

$$T_{+z,-'z'}^1 = \underbrace{(-\partial_+ \partial_{-'})}_{\eta_{+-}} u \underbrace{(-\partial_z \partial_{z'})}_u - \dots$$

positive if the derivatives act on the numerator
negative if the derivatives act on the **denominator**

$$u = \frac{(z - z')^2 + b^2 + \eta_{+-}(x^+ - x'^+)(x^- - x'^-)}{zz'}$$

Sign flip from the warp factor

$$T_{+z, -'z'}^1 \sim \eta_{+-} \left(\underbrace{-1 - v}_{\text{negative}} + \underbrace{\frac{z}{z'} + \frac{z'}{z}}_{\text{positive (expected)}} \right) \quad v = \frac{(z - z')^2 + b^2}{2zz'}$$

The negative contribution is enhanced after integrating over b ,

dominates over the positive contribution when $\ln s \gg \sqrt{\lambda}$

Attraction between like charges and repulsion between opposite charges !

$\Delta\sigma < 0$ due to the curvature of AdS.

A scenario (“prediction”)

Reggeon gives a positive contribution $\Delta\sigma \sim s^{\alpha_R - 1} > 0$

$$\alpha_R(0) = 1 - \frac{9}{2\sqrt{\lambda}}, \quad 1 - \frac{16}{2\sqrt{\lambda}}, \dots,$$

Odderon gives a negative contribution $\Delta\sigma \sim s^{\alpha_O - 1} < 0$
when $\ln s \gg \sqrt{\lambda}$

$$\alpha_O(0) = 1, \quad 1 - \frac{3}{2\sqrt{\lambda}}, \quad 1 - \frac{8}{2\sqrt{\lambda}}, \quad 1 - \frac{15}{2\sqrt{\lambda}}, \dots$$

At the ISR energies, the Reggeon dominates. But the Odderon eventually flips the sign of $\Delta\sigma$, possibly at the LHC!

A recent prediction

Extrapolate the low-energy phenomenological fit to high energies.

Avila, Gauron, Nicolescu (2007)

LHC is also a good place to discover the Odderon. We predict

$$\begin{aligned}\sigma_T^{pp}(\sqrt{s} = 14 \text{ TeV}) &= 123.32 \text{ mb} , \\ \Delta\sigma(\sqrt{s} = 14 \text{ TeV}) &= -3.92 \text{ mb} ,\end{aligned}$$

Negative !

An old prediction

A Possible Interpretation of pp Rising Total Cross-Sections.

L. LUKASZUK (*) and B. NICOLESCU

*Division de Physique Théorique (**), Institut de Physique Nucléaire (***)
and Laboratoire de Physique Théorique et Hautes Energies (***) - Paris*

In the very first paper in
1973, they have predicted

$$\Delta\sigma < 0 \quad !!$$

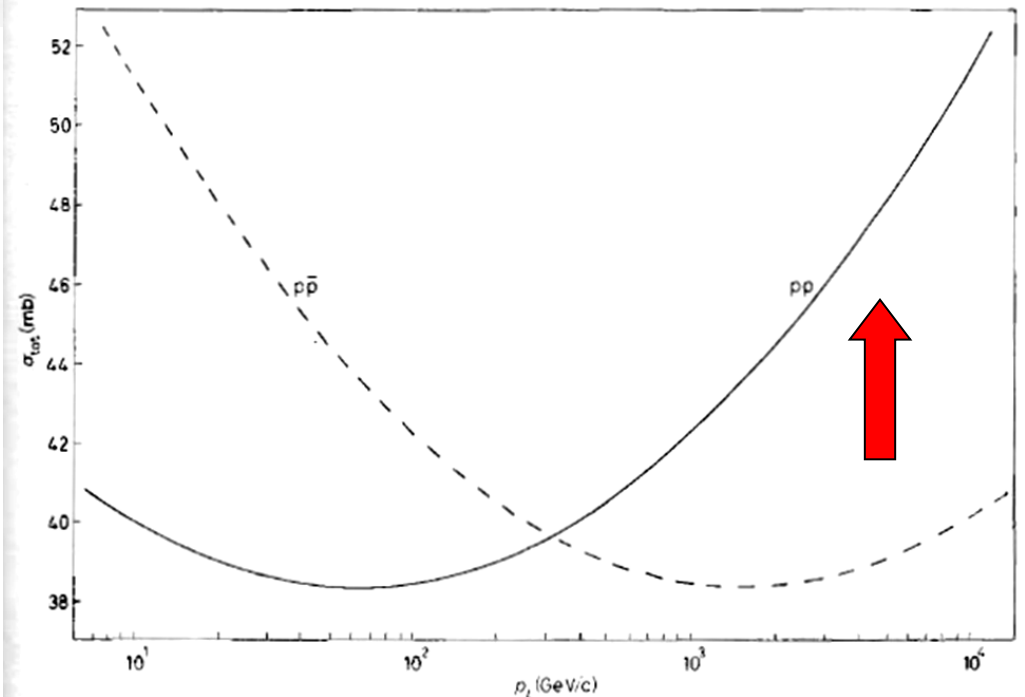


Fig. 3. - Comparison between the predictions of σ_{tot}^{pp} and $\sigma_{tot}^{p\bar{p}}$.

A hint from LHC?



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Need to produce more particles in pp !

Conclusion

- Single-odderon exchange between “baryons” calculated from AdS/CFT
- Analytical expression of the total cross section difference.
- The pp cross section becomes larger than the $ppbar$ cross section at very high energies, possibly at the LHC.