

News on Exclusive Diffractive Higgs studies



Marek Taševský

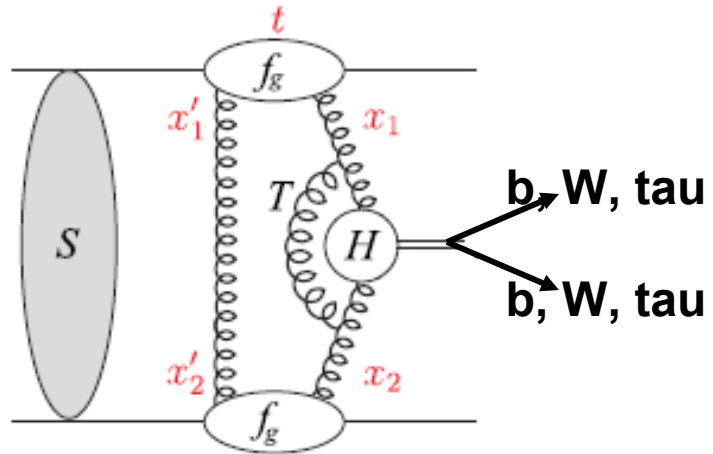
Institute of Physics, Academy of Sciences, Prague

Low-x workshop, Kavala, Greece - 25/06 2010

CED Higgs in SM and MSSM

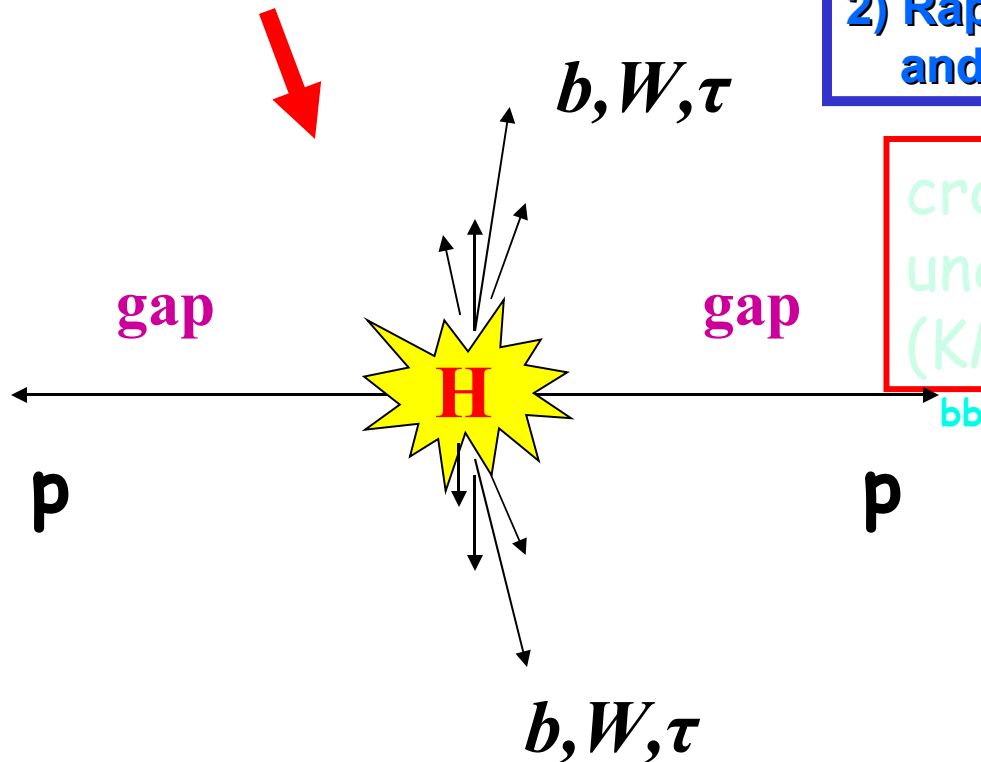
Updates to EPJC 53 (2008) 231

Central Exclusive Diffraction: Higgs production



This process is the core of the physics case of Forward detector upgrades (AFP in ATLAS, HPS in CMS)

- 1) Protons remain undestroyed and can be detected in forward detectors
- 2) Rapidity gaps between leading protons and Higgs decay products



cross-sections predicted with uncertainty of 3 !! (or more?) (KMR group, CHIDe model)

: at 120 GeV needs a special diffractive trigger

Central Exclusive Diffraction: Higgs production

Advantages:

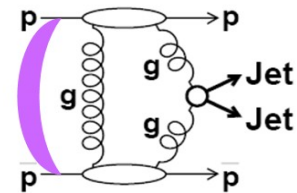
- I) Forward detectors give much better mass resolution than the central detector
- II) $JZ = 0$, CP-even selection rule:
 - strong suppression of CED $gg \rightarrow bb$ background (by (m_b/M_H))
 - produced central system is 0 → **just a few events are enough to determine Higgs quantum numbers.** Standard searches need high stat. (ϕ -angle correlation of jets in VBF of Higgs) and **coupling to Vector bosons.**
- III) Access to main Higgs decay modes in one (CED) process: bb , WW , $\tau\tau$
 - ↓
 - information about Yukawa coupling**
(Hbb difficult in standard searches due to huge bg.)
- IV) In MSSM, CED Higgs process give very important information on the Higgs sector, and in addition, for sufficiently high $\tan\beta$ values, it allows direct measurement of the Higgs width.

Disadvantages:

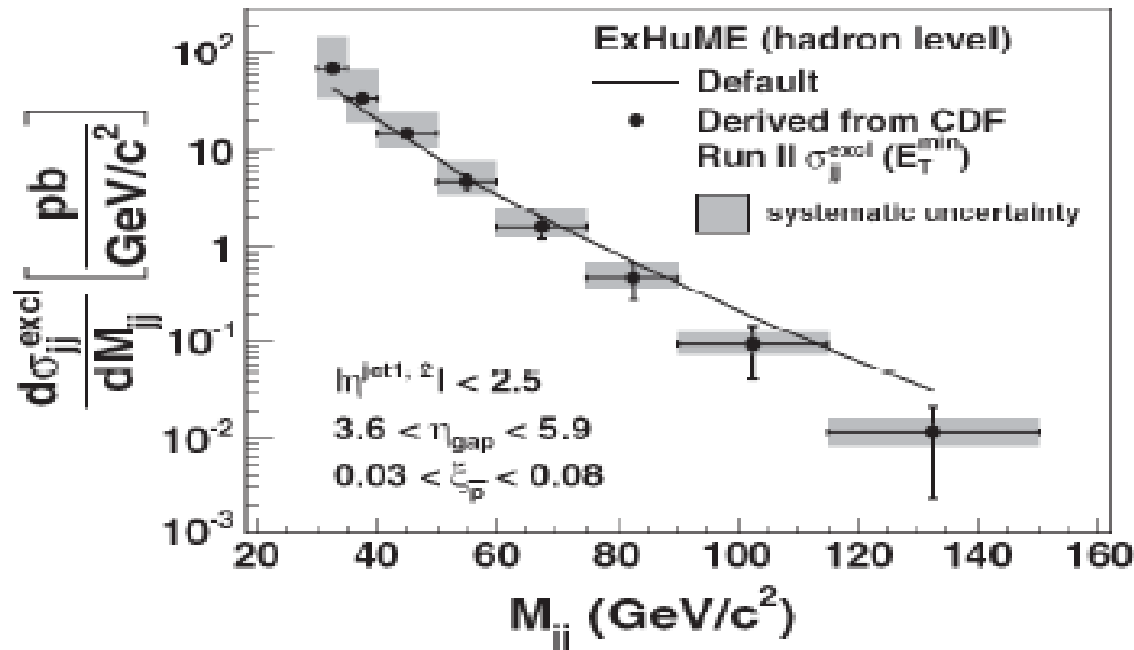
- Low signal x-section
- Large Pile-up



Exclusive dijet x-section vs. M_{jj}



arXiv:0712.0604 ,
PRD-2008



curve: ExHuME hadron-level exclusive dijet cross sections vs. dijet mass

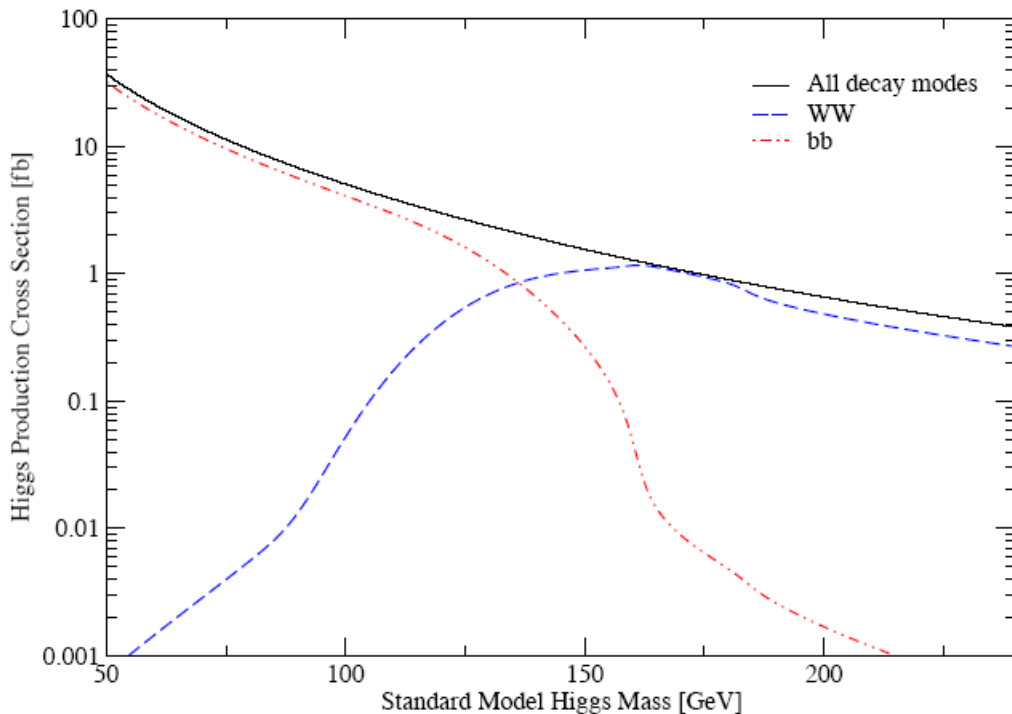
points: derived from CDF excl. dijet x-sections using ExHuME

Stat. and syst. errors are propagated from measured cross section uncertainties using M_{jj} distribution shapes of ExHuME generated data.

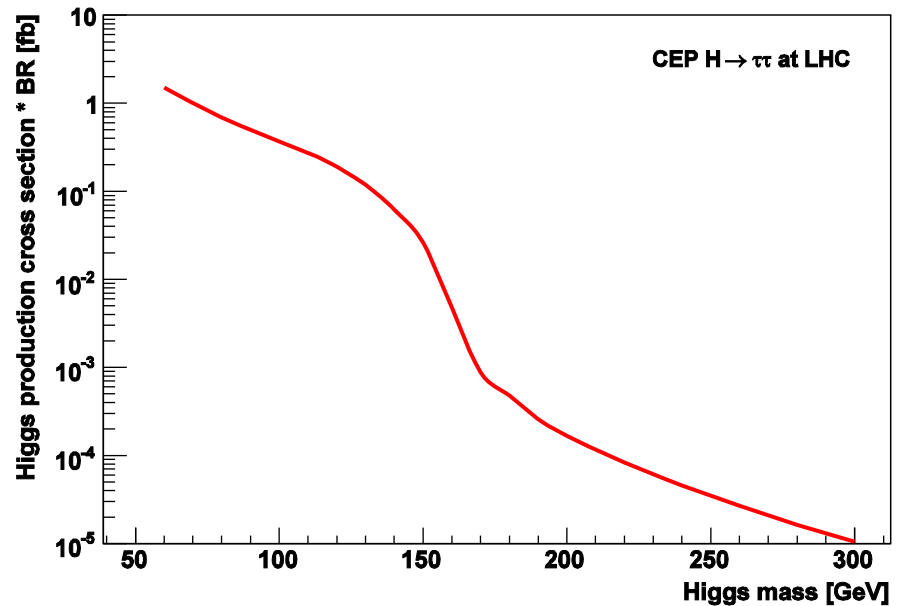
Use ExHuMe as the standard event generator

Cross sections (KMR) and FD Acceptances

M_H [GeV]	σ (bb) [fb]	σ (WW*) [fb]	Acc (420+420)	Acc(420+220)
120	1.9	0.37	0.20	0.17
130		0.70	0.15	0.24
140	0.6	0.87	0.11	0.31
160	0.045	1.10	0.04	0.43
180	0.0042	0.76	0.01	0.53



Acceptances for AFP detectors at (220,420) m and 2.5mm,4mm from the beam (1mm dead space)



Analyses on CED Higgs production

CMS:

$H \rightarrow bb$: fast simulation, $100 < M_H < 300$ GeV, $d_{220} \sim 1.5$ mm, $d_{420} \sim 4.5$ mm, $\text{Acc} = \text{Acc}(\xi, t, \varphi)$

- track variables (N_c, N_c^\perp) not used \rightarrow PU bg overestimated
- L1 trigger: single-sided FD220 .AND. Etjet1>40.AND. Etjet2>40. No efficiencies applied
- published in CMS-Totem document CERN/LHCC 2006-039/G-124
- signal selection efficiencies used in MSSM study ([EPJC 53 \(2008\) 231](#))

ATLAS:

$H \rightarrow bb$: 1) gen.level + smearing of basic quantities, $M_H = 120$ GeV

- L1 trigger: fixed rates of dijet triggers with prescales
- one MSSM point ($\tan\beta = 40$): [JHEP 0710 \(2007\)090](#)

2) fast simulation, $M_H = 120$ GeV

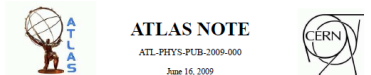
- L1 trigger: a dedicated $H \rightarrow bb$ trigger (ATLAS-COM-DAQ-2009-062).
Efficiencies and prescales easily applicable offline
- common ATLAS note just released

$H \rightarrow WW$: fast + full simulation, $M_H = 160$ GeV

- common ATLAS note just released

$H \rightarrow \tau\tau$: full simulation, $100 < M_H < 300$ GeV (designed for a MSSM study)

All ATLAS analyses use $d_{220} \sim 2-3$ mm, $d_{420} \sim 4-5$ mm, $\text{Acc} = \text{Acc}(\xi, t)$



Central Exclusive Diffractive production of the Standard Model Higgs boson decaying into bb and WW as seen in the ATLAS detector

A. Brandt, V. Jurasek, A. Pal and M. Talevsky

Abstract

We present a feasibility study of the Standard Model Higgs boson production in the Central Exclusive Diffractive processes with Higgs boson decaying into bb and WW. Results are based on events generated by Monte Carlo event generators only. Dedicated event generators such as ExHicMe and Pomwig are used to generate central exclusive and inclusive diffractive processes. The effect of pile-up is studied using PYTHIA and HERWIG event generators. Acceptances and resolutions of proposed forward detectors to be placed at 220 and 420 m from the interaction point are used together with the response of the ATLAS detector obtained from the ATLAS fast simulation. The predicted cross sections for the signal process is low and therefore its detection requires collecting data at the highest instantaneous luminosities where the pile-up is an issue. After applying a number of severe cuts to reduce the background coming from the pile-up events, the signal event yield turns out to be rather moderate, however keeping the signal-to-background ratio close to unity.

Analysis strategy for $H \rightarrow bb$

- 1) Proton detection:** in Forward proton taggers at 220m and 420m
- 2) jets:** two b-tagged jets: $ET1 > 45 \text{ GeV}$, $ET2 > 30 \text{ GeV}$, $|\eta_{1,2}| < 2.5$, $3.0 < |\phi_1 - \phi_2| < 3.3$
- 3) Kinematics constraints - matching criteria:** $0.75 < R_j < 1.2$, $|\Delta y| < 0.1$
- 4) L1 triggers:**
 - 420+220:** $J20J40 + FD220 + \bar{\eta} < 0.5 + |\Delta\eta| < 2 + fT > 0.45 \rightarrow$ special diffractive trigger
 - 420+420:** $J20J40 + \bar{\eta} < 0.5 + |\Delta\eta| < 2 + fT > 0.45 \rightarrow$ FD420 cannot be included in L1
- 5) Mass windows:** $117.6 < M_{420} < 122.4$,
 $114.2 < M_{420+220} < 125.8$ (3σ - window)

6) PU bg suppressors: Few tracks outside the dijet reduction factor ~ 20 from fast timing detector

Take the experimental efficiencies ϵ and calculate

- 7) Cross sections:** $\sigma_{SM} \times \epsilon \times \Gamma(gg \rightarrow H)_{NP} / \Gamma(gg \rightarrow H)_S$
- 8) Decays:** $BR(H \rightarrow bb)_{NP}$ from FeynHiggs
(MSSM: incl. Δb dep.)

Efficiencies for SM $H \rightarrow bb$ (CMS+Totem)

M_H [GeV]	Acc_{420}	Acc_{comb}	Acc_{220}	ϵ_{420}	ϵ_{comb}	ϵ_{220}
100	0.37	0.13	0.0	0.012	0.008	0.0
120	0.31	0.25	0.0			
140	0.25	0.37	0.0			
160	0.19	0.49	0.0			
180	0.14	0.60	0.0			
200	0.09	0.69	0.0			
300	0.0	0.76	0.13			



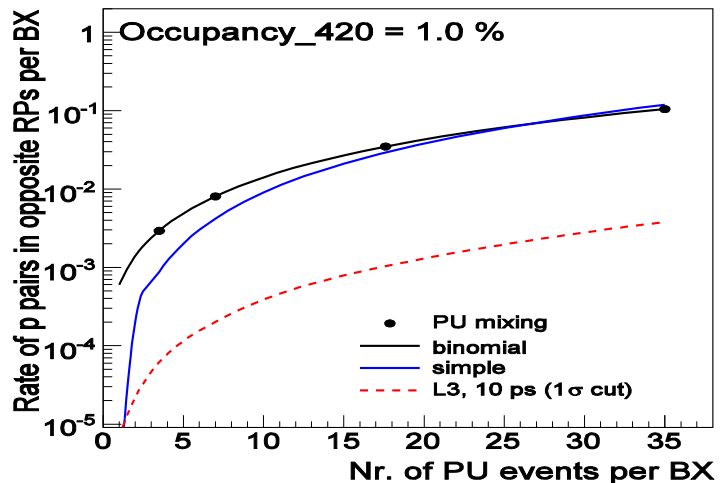
CERN LHC 2006-039-G-124
CMS Note-2007-002
TOTEM Note-06-5
21 December 2006

Prospects for
Diffractive
and Forward Physics
at the LHC

The CMS and TOTEM
diffractive and forward physics
working group

CED experimental challenges: Pile-up

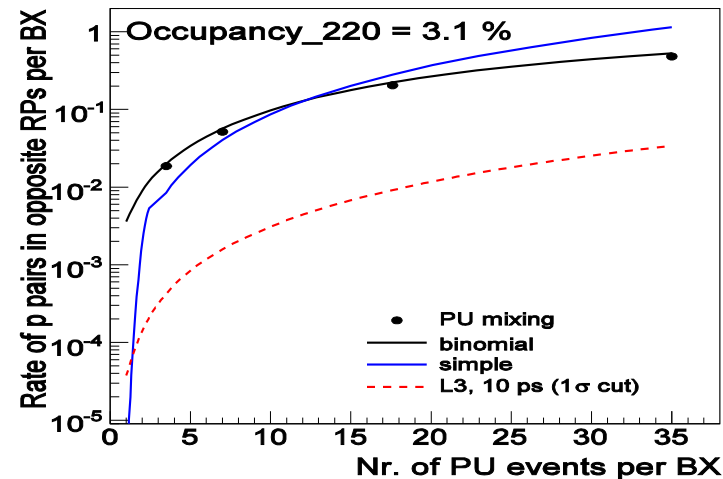
Huge Pile-up bg for diffractive processes: overlap of three events ($2 \times SD + \text{non-diffr. Dijets}$). Can be reduced by Fast Timing detectors: t-resol. required: few ps for high lumi!



Huge rates

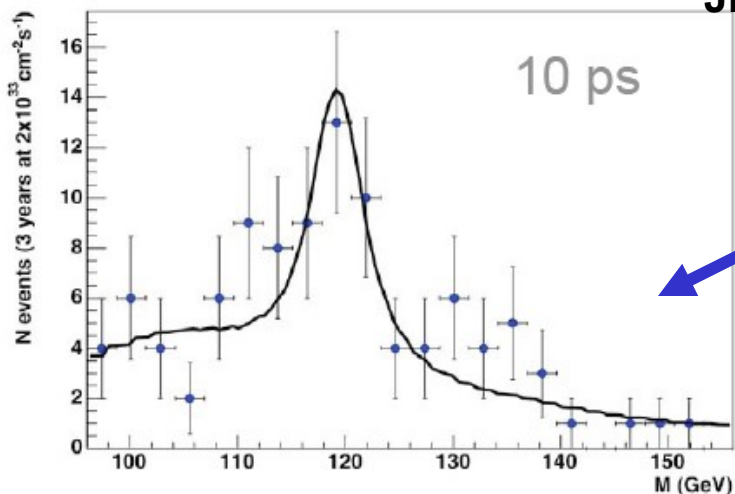
Reduced by

Fast timing Detectors (factor ~20)



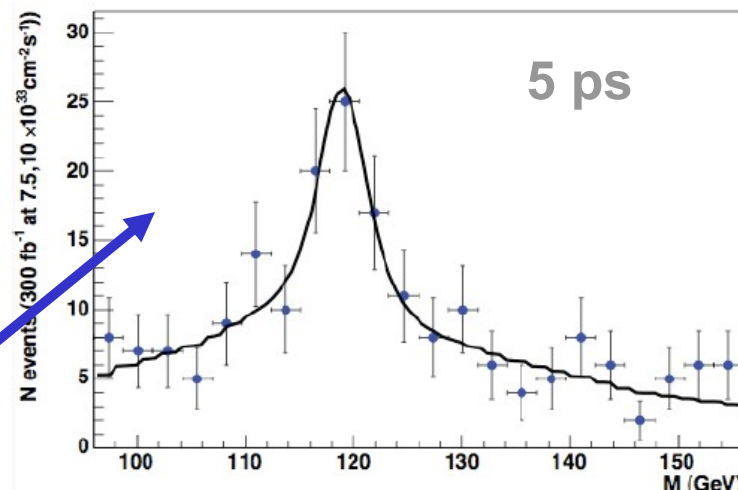
$h \rightarrow bb$, mhmax scenario, standard ATLAS L1 triggers, 420m only, 5 mm from beam

JHEP 0710:090,2007



60 fb collected at $2 \times 10 \text{ cm s}$

150 fb at $7.5 \times 10 \text{ cm s}$
plus 150 fb at 10 cm s



MSSM and CED go quite well together

The **intense coupling regime** is where the masses of the 3 neutral Higgs bosons are close to each other and $\tan\beta$ is large

Extended Higgs sectors: “typical” features

Search for heavy MSSM Higgs bosons ($M_A, M_H \gg M_Z$):

Decouple from gauge bosons

⇒ no HVV coupling

⇒ no Higgs production in weak boson fusion

⇒ no decay $H \rightarrow ZZ \rightarrow 4\mu$

Large enhancement of coupling to $b\bar{b}$, $\tau^+\tau^-$ for high $\tan\beta$

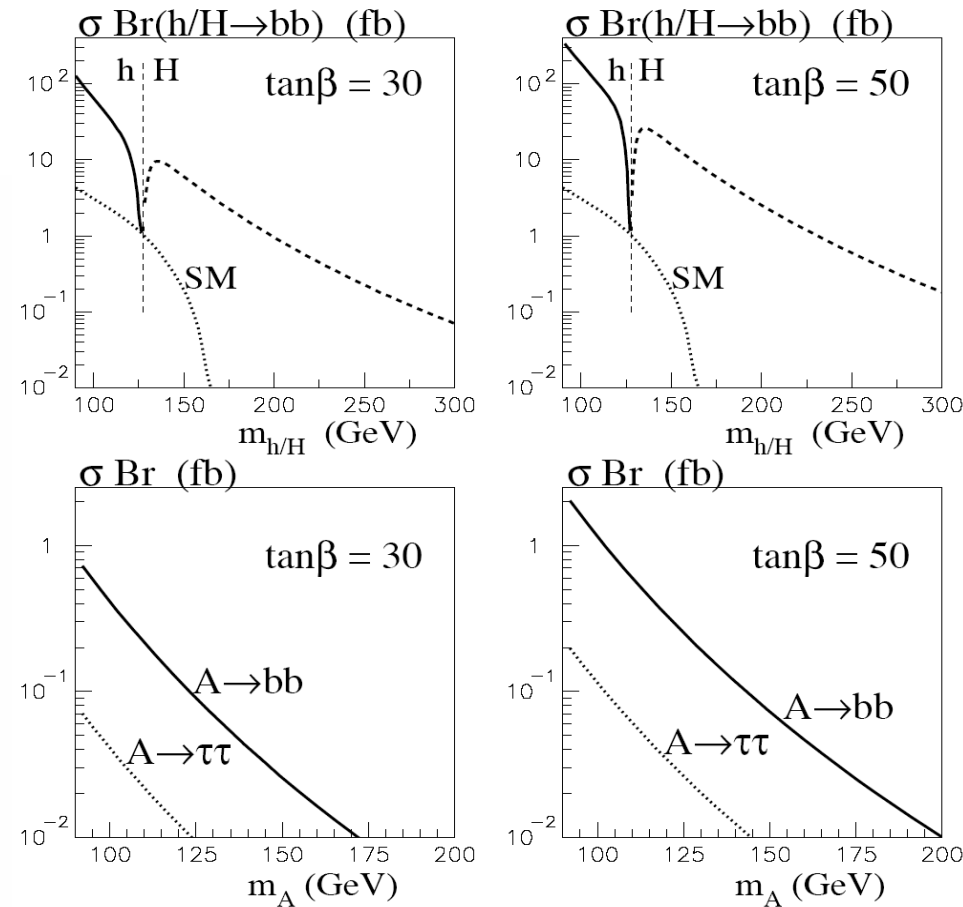
⇒ Decays into $b\bar{b}$ and $\tau^+\tau^-$ play a crucial role

“Typical” features of models with an extended Higgs sector:

- A light Higgs with SM-like properties, couples with about SM-strength to gauge bosons
- Heavy Higgs states that decouple from the gauge bosons

Studying the MSSM Higgs Sector by Forward Proton Tagging at the LHC, Georg Weiglein, EPS07, Manchester, 07/2007 – p.3

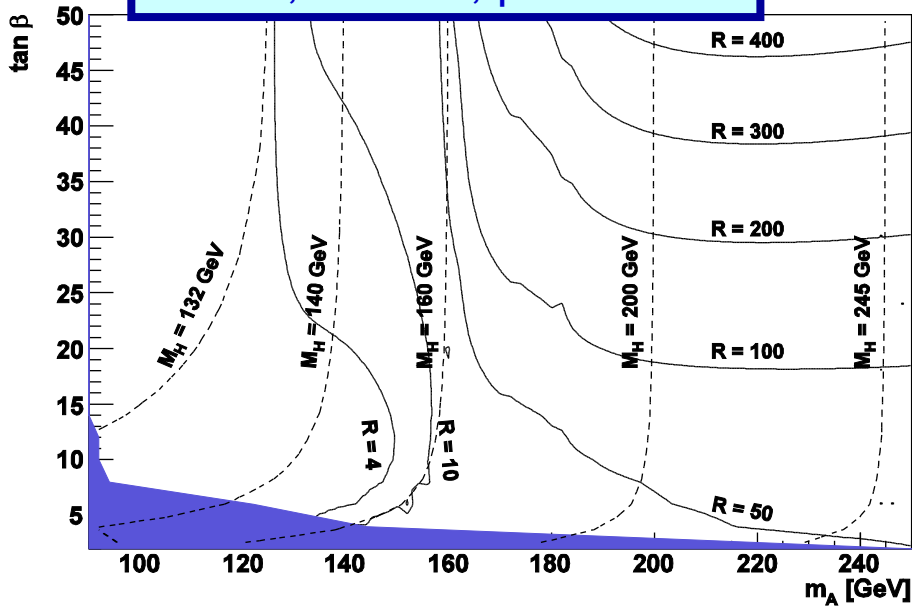
Central exclusive diffractive production



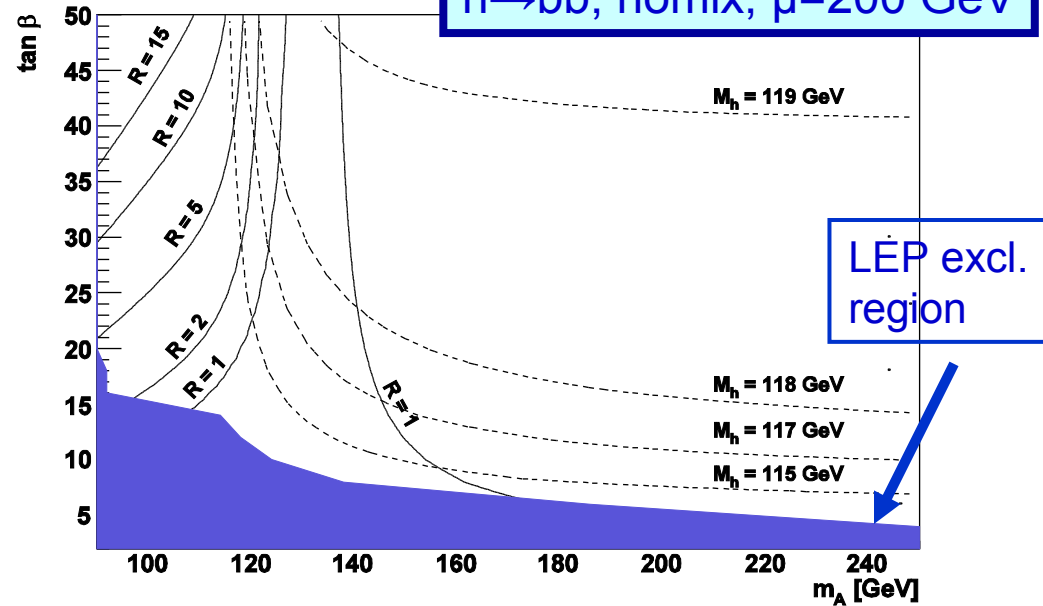
Well known difficult region for conventional channels, tagged proton channel may well be the **discovery channel** and is certainly a powerful **spin/parity filter**

R=MSSM[M, tanβ] / SM[M]

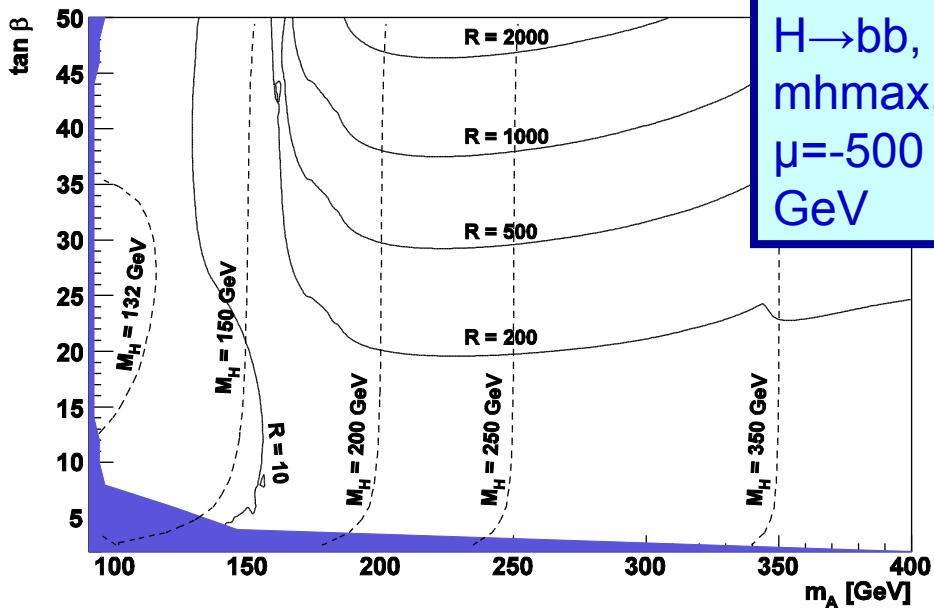
H→bb, mhmax, μ=200 GeV



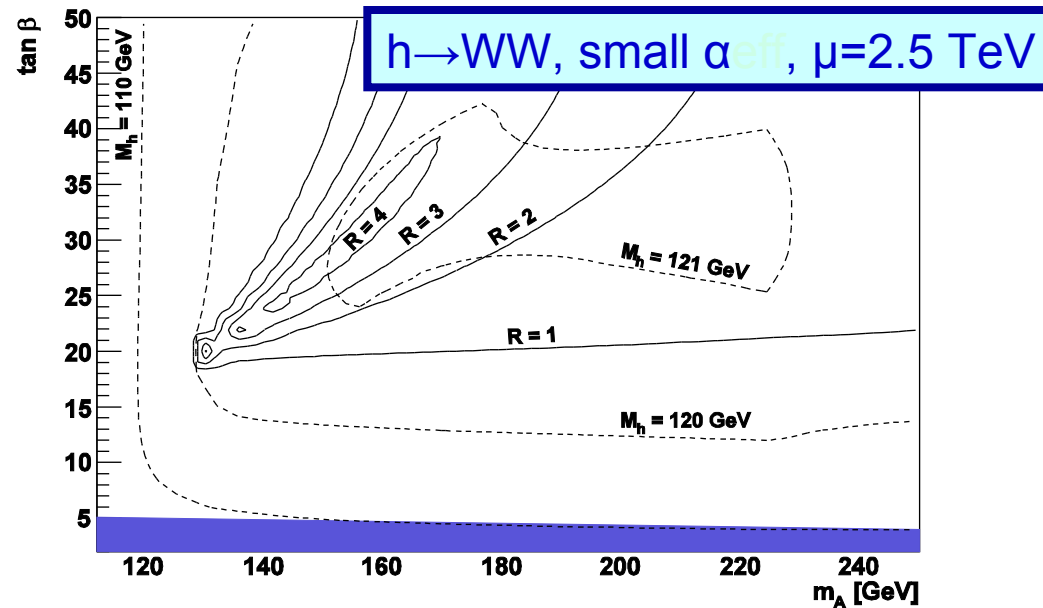
h→bb, nomix, μ=200 GeV



H→bb, mhmax, μ=-500 GeV



h→WW, small α_{eff}, μ=2.5 TeV



Update of EPJC 53 (2008) 231

(HKRSTW = Heinemeyer, Khoze, Ryskin, Stirling, Tasevsky, Weiglein)

1) FeynHiggs 2.3.0 → most recent FeynHiggs 2.6.2

2) NLO formula for CED $gg \rightarrow bb$ background

3) Tevatron exclusion areas added

4) A few Cold Dark Matter scenarios tried

5) 4 generation

A new paper is close to publication

Four luminosity scenarios (ATLAS+CMS):

1) 60 fb⁻¹ – low lumi (no pile-up)

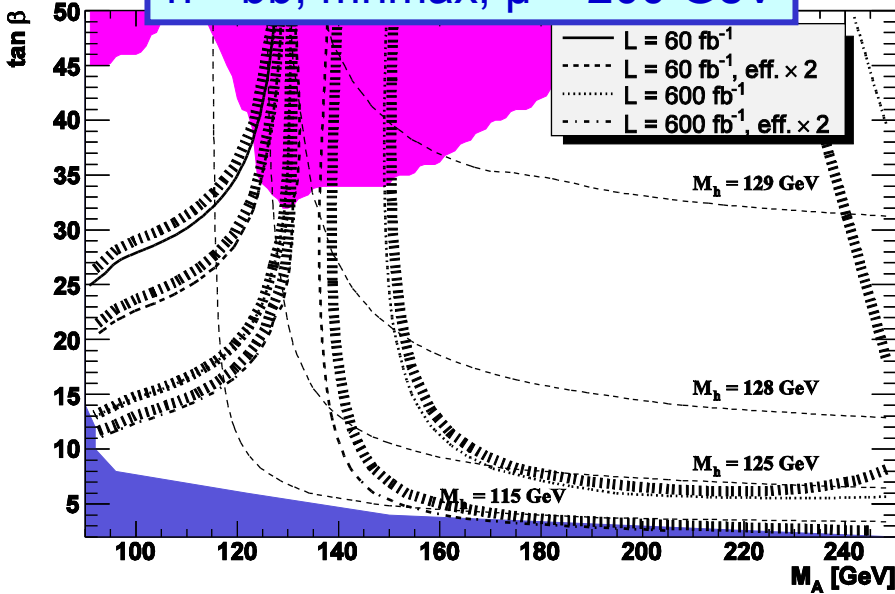
2) 60 fb⁻¹ x 2 – low lumi (no pile-up) but improved signal efficiency

3) 600 fb⁻¹ – high lumi (pile-up suppressed)

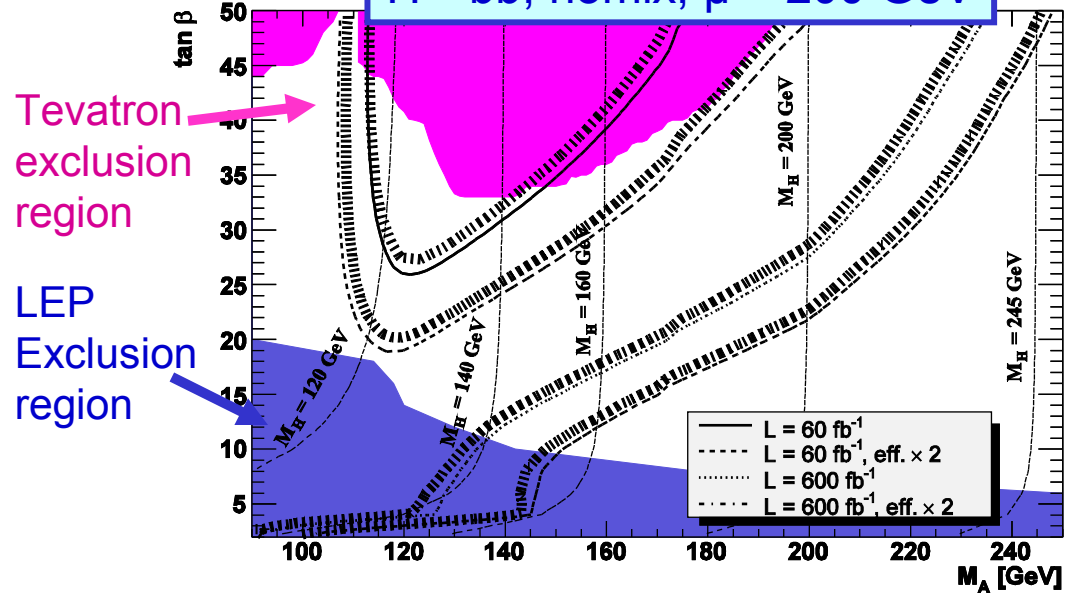
4) 600 fb⁻¹ x2 – high lumi (pile-up suppressed) but improved signal efficiency

3 σ -contours

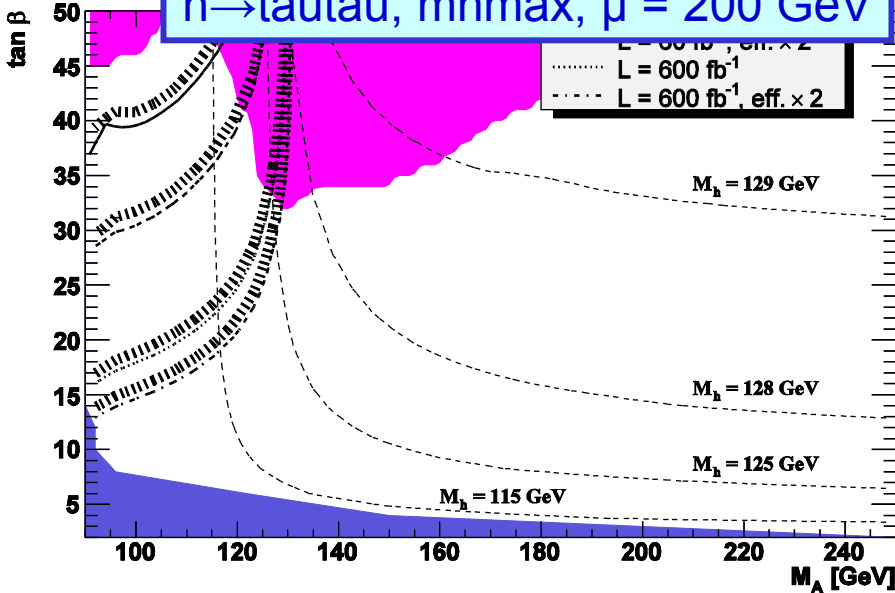
$h \rightarrow bb$, mhmax, $\mu = 200$ GeV



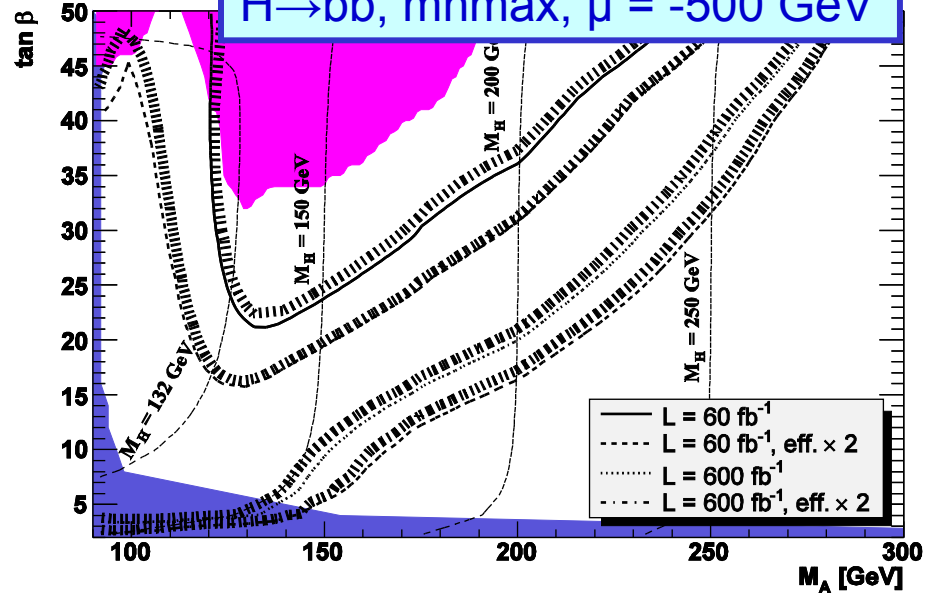
$H \rightarrow bb$, nomix, $\mu = 200$ GeV



$h \rightarrow \tau\tau$, mhmax, $\mu = 200$ GeV

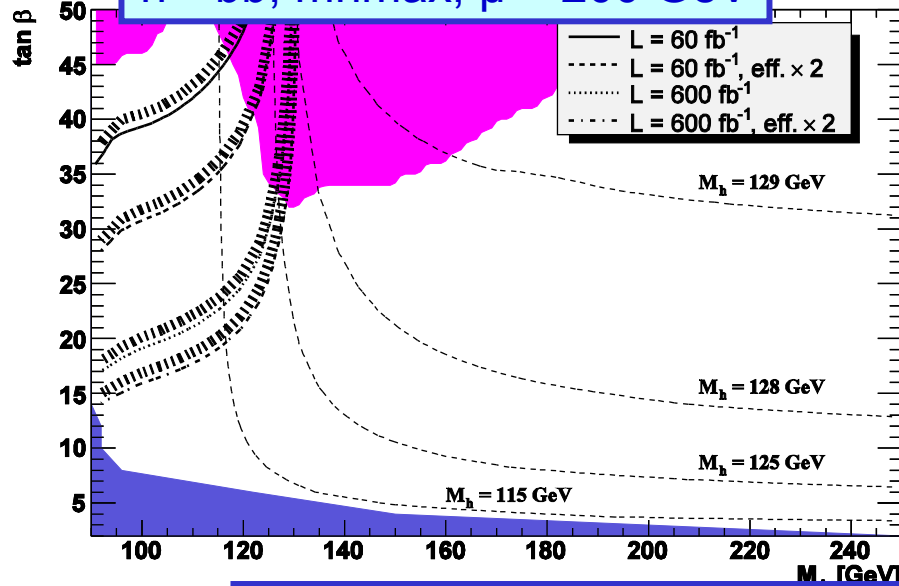


$H \rightarrow bb$, mhmax, $\mu = -500$ GeV

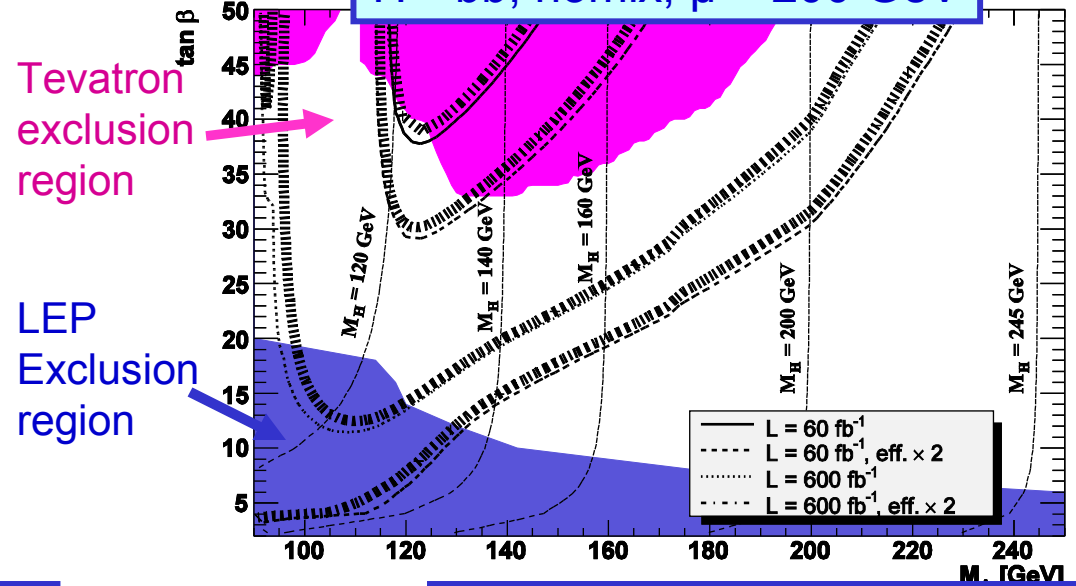


5 σ -contours

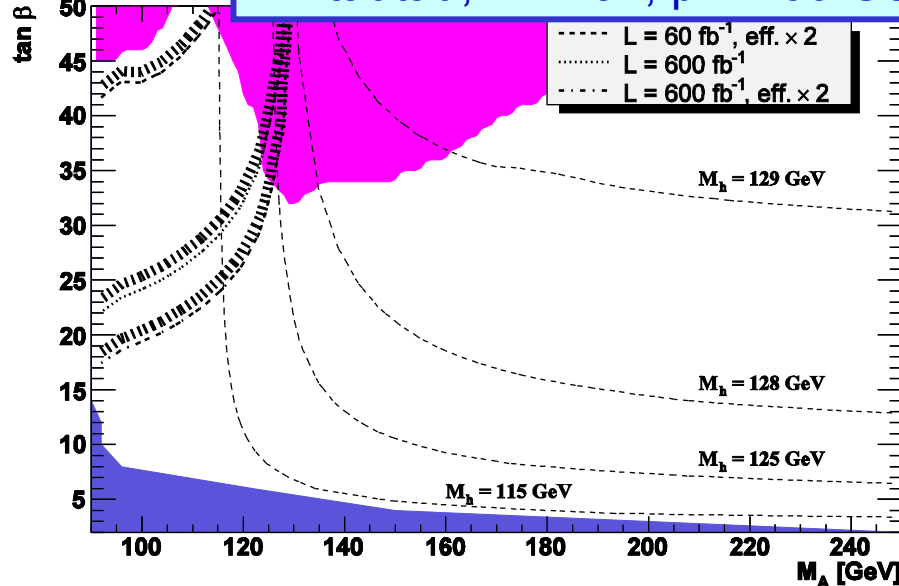
$h \rightarrow bb$, mhmax, $\mu = 200$ GeV



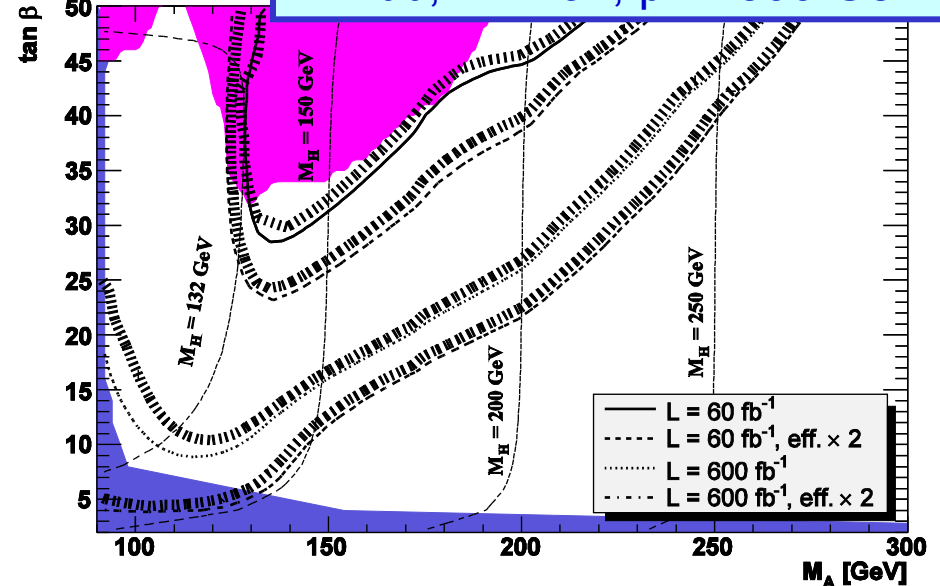
$H \rightarrow bb$, nomix, $\mu = 200$ GeV



$h \rightarrow \tau\tau$, mhmax, $\mu = 200$ GeV



$H \rightarrow bb$, mhmax, $\mu = -500$ GeV



Cold Dark Matter scenarios

- Follow the ideas of [arXiv: 0709.0098v2](#) (J.Ellis, T.Hahn, S.Heinemeyer, K.Olive and G.Weiglein):
- Explore new benchmark surfaces for MSSM Higgs phenomenology so that the supersymmetric relic density is compatible with the cosmological density of cold dark matter inferred from WMAP and other observations.

The allowed range of cold dark matter density: $0.0882 < \Omega_{\text{CDM}} h^2 < 0.1204$

- The benchmark surfaces may be presented as $(M_A, \tan\beta)$ planes with fixed or systematically varying values of NUHM parameters, such as scalar mass m_0 , gaugino mass $m_{1/2}$, trilinear parameter A_0 and the Higgs mixing parameter μ .

Plane P1: $m_0=800 \text{ GeV}$, $\mu=1000 \text{ GeV}$, $A_0=0$, varying $9/8M_A - 12.5 \text{ GeV} < m_{1/2} < 9/8M_A + 37.5 \text{ GeV}$

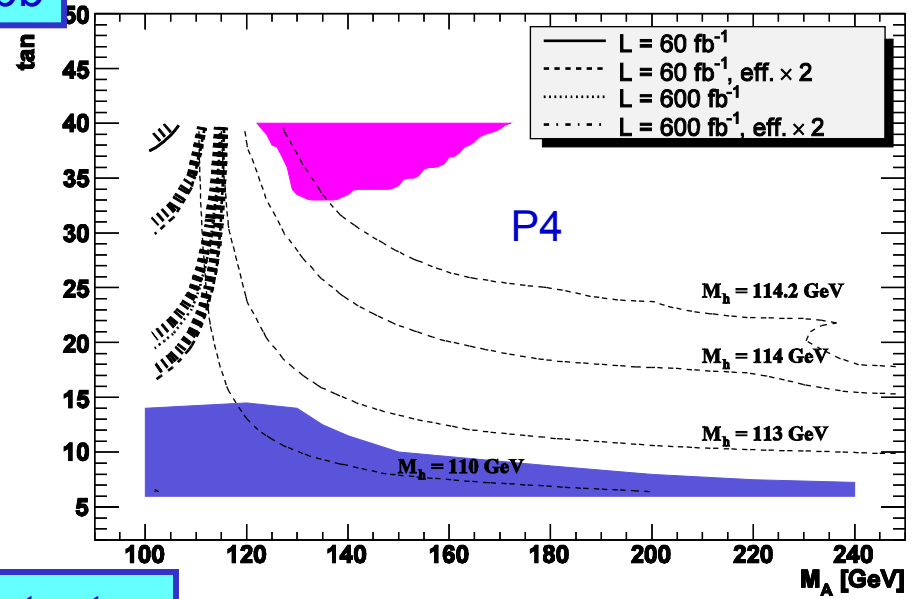
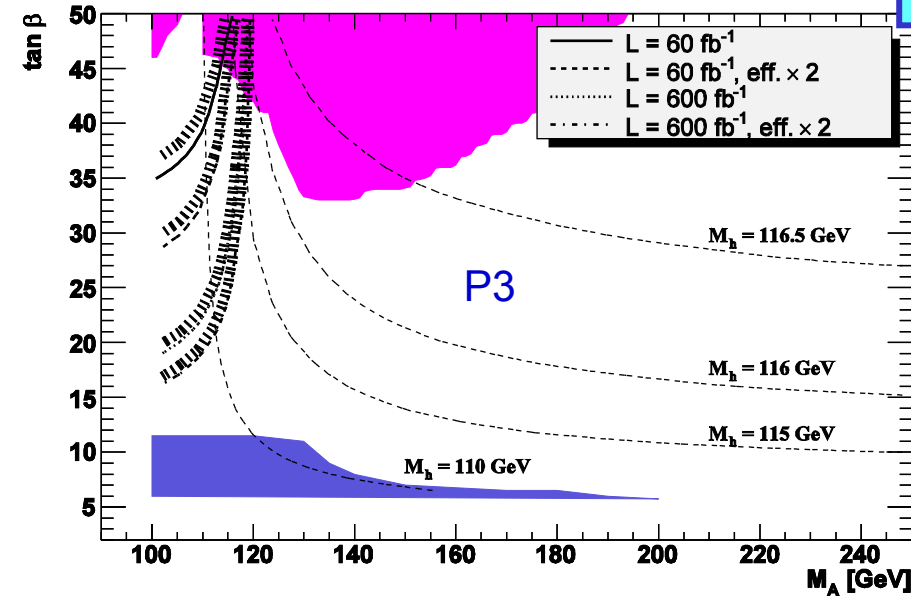
Plane P2: $m_0=300 \text{ GeV}$, $\mu=800 \text{ GeV}$, $A_0=0$, varying $1.2M_A - 40 \text{ GeV} < m_{1/2} < 1.2M_A + 40 \text{ GeV}$

Plane P3: $m_{1/2}=500 \text{ GeV}$, $m_0=1000 \text{ GeV}$, $A_0=0$, $200 < \mu < 400 \text{ GeV}$

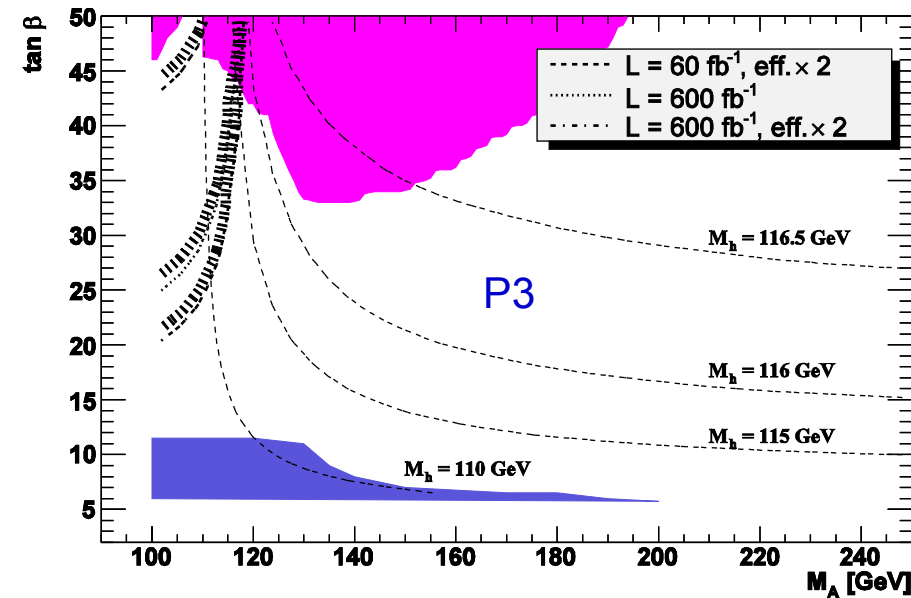
Plane P4: $m_{1/2}=300 \text{ GeV}$, $m_0=300 \text{ GeV}$, $A_0=0$, $200 < \mu < 350 \text{ GeV}$

Cold Dark Matter: $h \rightarrow bb, \tau\tau$, 5σ -contours

$h \rightarrow bb$

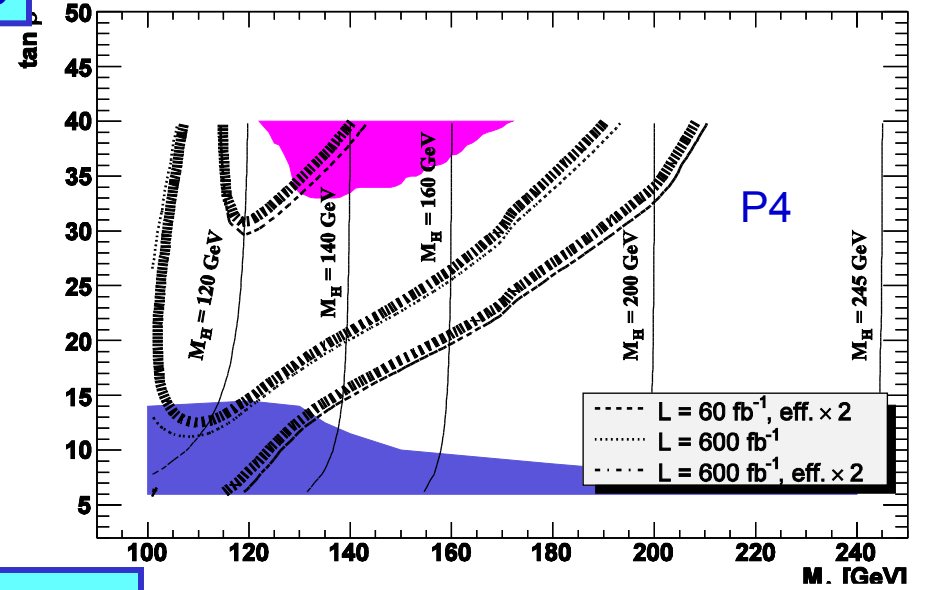
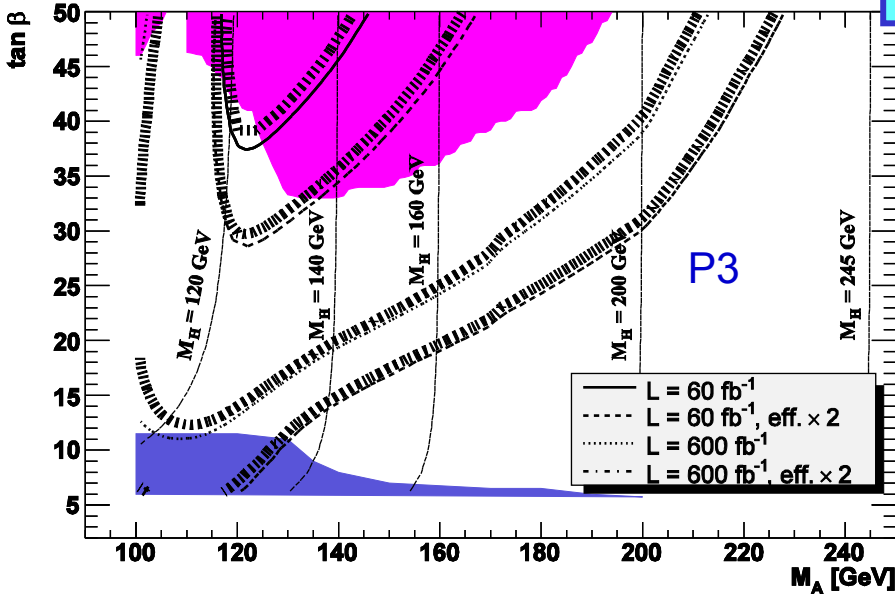


$h \rightarrow \tau\tau$

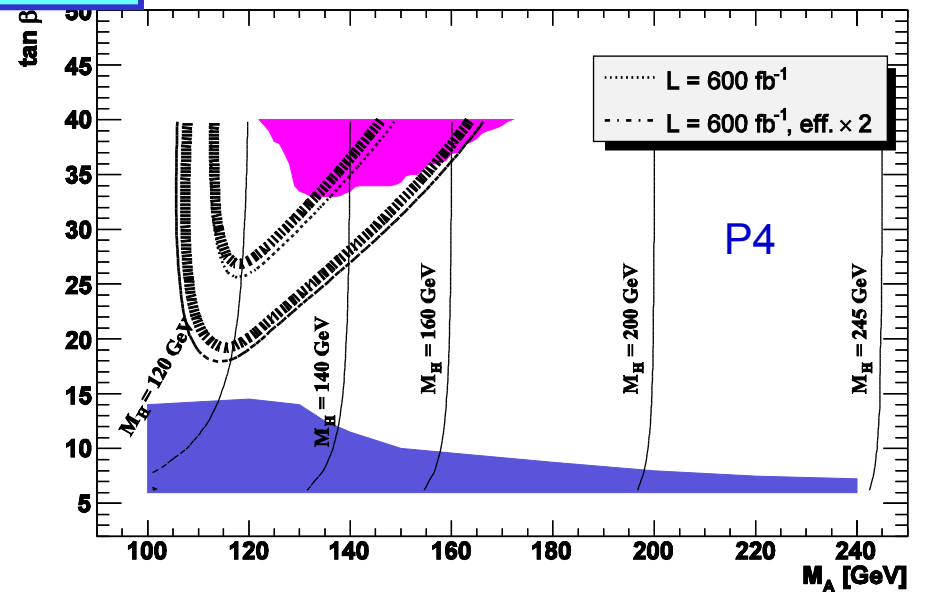


Cold Dark Matter: $H \rightarrow bb, \text{tautau}, 5\sigma\text{-contours}$

$H \rightarrow bb$



$H \rightarrow \text{tautau}$



Determination of Higgs CP properties

Existing SM analyses for LHC:

- rely largely on the coupling of Higgs boson to heavy gauge bosons:

$$\text{WBF } H \rightarrow W W \rightarrow ll\nu\nu$$

$$\text{WBF } H \rightarrow \tau \tau$$

$$\text{WBF } H \rightarrow ZZ \rightarrow 4l$$

Backgrounds: $gg \rightarrow H$, $WWjj(\text{EW})$, $WWjj(\text{QCD})$

tt , tW , $Zjj(\text{EW})$, $Zjj(\text{QCD})$

Prerequisites for these SM(-like) analyses: Higgs with

- sufficiently large HVV coupling
- sufficiently large $\text{BR}(H \rightarrow VV)$ – $M_H > 140$ GeV to suppress $H \rightarrow bb$;
- possibly large $\text{BR}(H \rightarrow \tau\tau)$

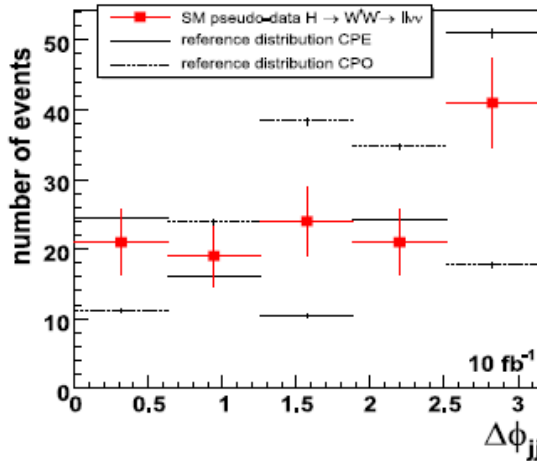
SM analyses of the structure of the HVV coupling: CP-even vs. CP-odd

[T. Plehn et al., 2001] (theory)

[V. Hankele et al., 2006] (theory)

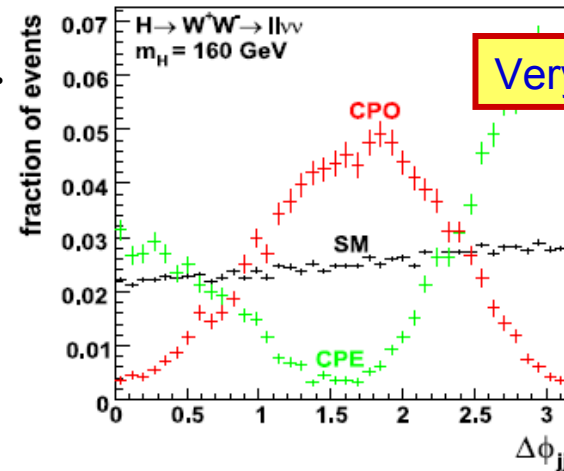
[C. Ruwiedel et al., 2007] (experiment)

Higgs coupling structure: HWW, HVV



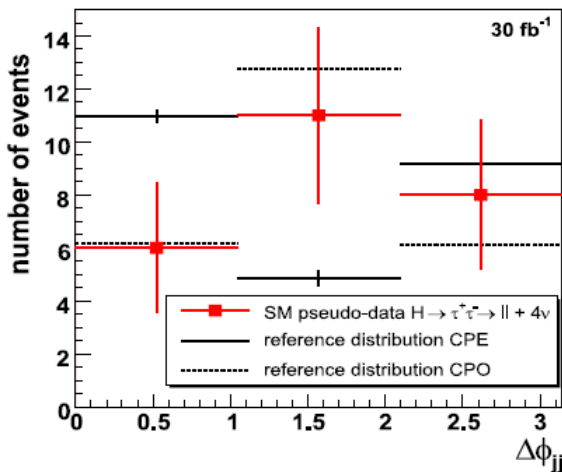
10 fb⁻¹ $H \rightarrow W^+W^- \rightarrow ll\nu\nu, M_H = 120$

$\Delta\phi_{jj} =$
Difference in azimuthal
angle of two tagging jets



5σ - discrimination of anomalous couplings already with 10 fb⁻¹

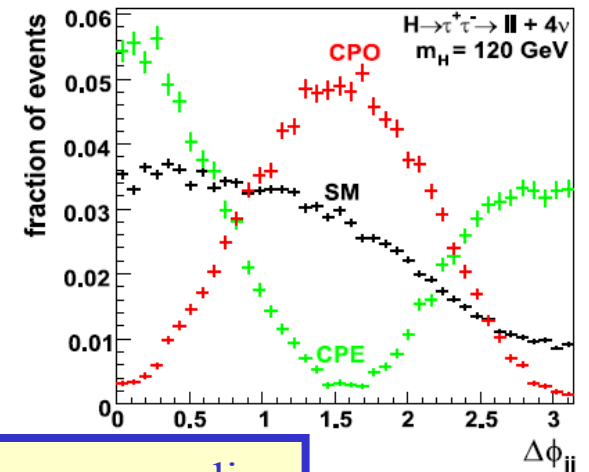
SM = signal with SM couplings + backgrounds, **CPE** = signal with CP-even couplings + backgrounds
CPO = signal with CP-odd couplings + backgrounds



30 fb⁻¹

$H \rightarrow \tau^+\tau^- \rightarrow ll + 4\nu, M_H = 120$ GeV

3σ - discrimination of anomalous CP-even coupling



Situation in MSSM

Light Higgs: $M_h < 135$ GeV: too small BR($h \rightarrow VV^{(*)}$)

Heavy Higgses:

$$\begin{aligned}g_{hVV} &= g_{HVV}^{\text{SM}} \times \sin(\beta - \alpha) \\g_{HVV} &= g_{HVV}^{\text{SM}} \times \cos(\beta - \alpha) \\g_{AVV} &= 0 \quad \text{at tree-level}\end{aligned}$$

H A

$M \approx M > 150$ GeV:

$\beta - \alpha \rightarrow \pi/2 \Rightarrow h$ has substantial VV coupling, but not sufficient BR($h \rightarrow \tau\tau$)

H and A have negligible VV coupling

H A

$M \approx M < 130$ GeV $\Rightarrow |\sin(\beta - \alpha)| \ll 1$ possible:

H has substantial VV coupling, but not sufficient BR($H \rightarrow \tau\tau$)

Heavy Higgses: method relying

Central Exclusive Diffraction!

Light

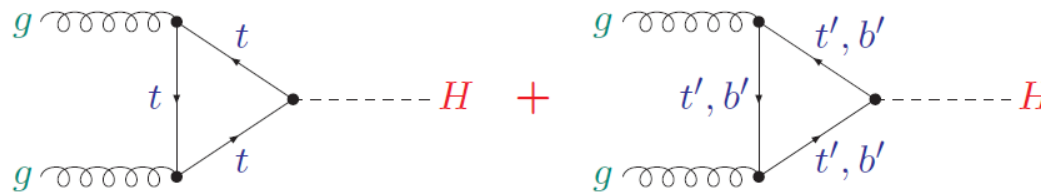
J = 0, C-even, P-even selection rule leads to a clear determination of quantum numbers of the centrally produced resonance. A few events are enough.

4th Generation Model

Assume the SM with a 4th generation of heavy fermions

Relevant changes:

1. additional contribution to $gg \rightarrow H$:



\Rightarrow factor of ~ 9 in Higgs production cross section

2. \Rightarrow factor of ~ 9 in $\Gamma(H \rightarrow gg)$

\Rightarrow reduced $\text{BR}(H \rightarrow b\bar{b})$, $\text{BR}(H \rightarrow \tau^+\tau^-)$

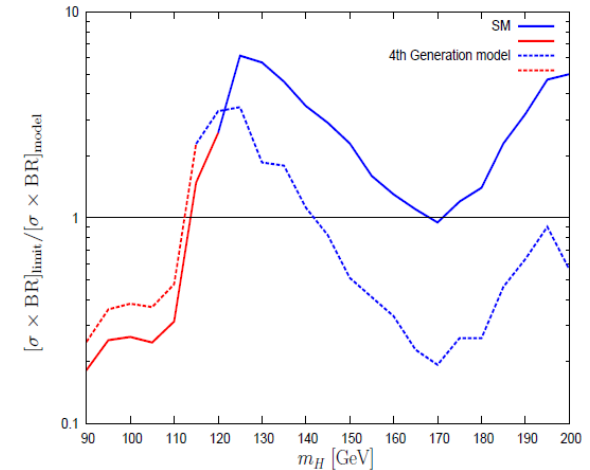
Evaluation of SM quantities with **FeynHiggs**

subsequent application of reduction and enhancement factors

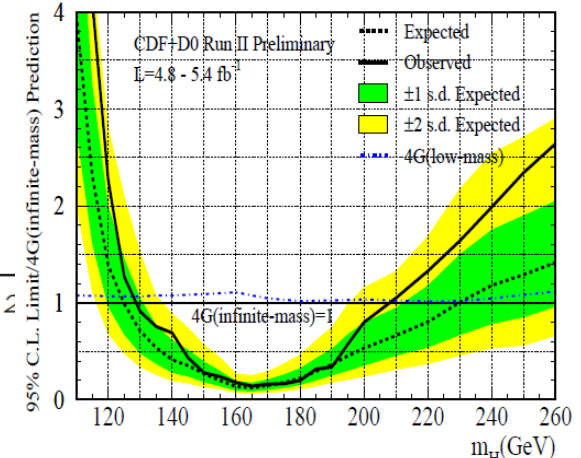
Sven Heinemeyer (for Marek Tasevsky) Physics at the LHC conference (DESY Hamburg), 08.06.2010

LEP and Tevatron limits

[P. Bechtle et al., 2008]



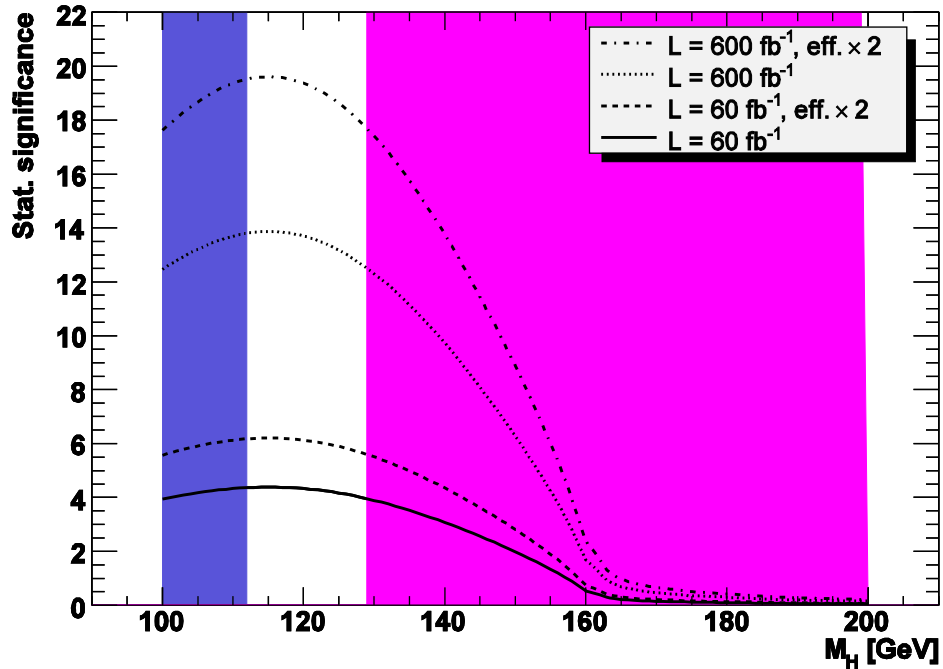
[CDF, D0, 2010]



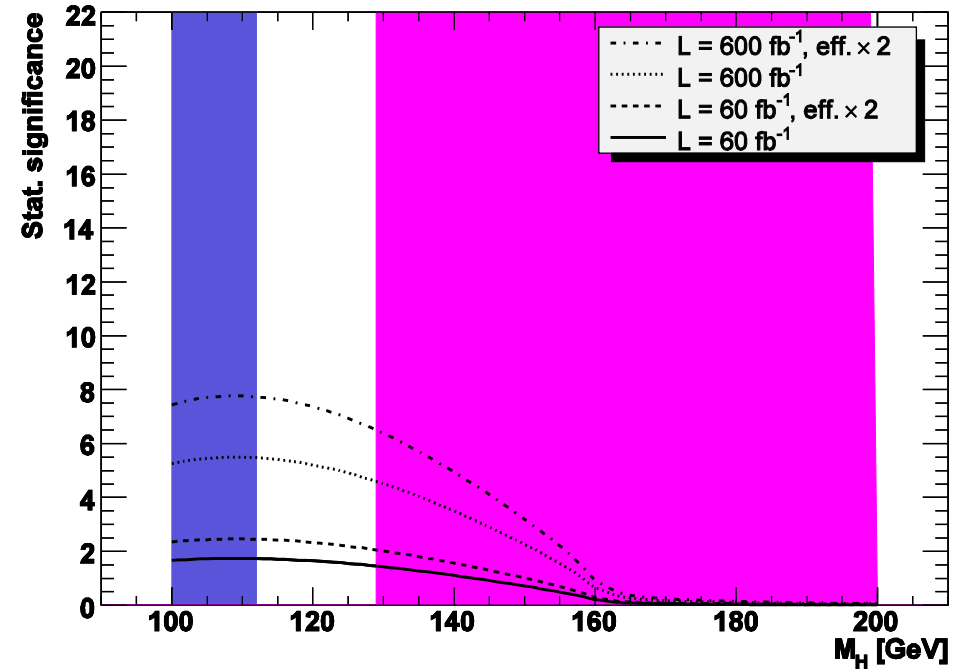
only $112 \text{ GeV} \lesssim M_H \lesssim 130 \text{ GeV}$, $M_H \gtrsim 210 \text{ GeV}$ still allowed

CED in 4th Generation Model

$H \rightarrow bb$



$H \rightarrow \tau\tau$



Good prospects for $H \rightarrow bb$ at low luminosities and for $H \rightarrow \tau\tau$ at high luminosities.

For $M_H > 200$ GeV: $\text{BR}(H \rightarrow bb, H \rightarrow \tau\tau)$ too small.

Summary

1) CED Higgs production in SM

- provides a moderate signal yields but it is attractive because

a few events are enough to establish the quantum numbers of a Higgs candidate. No need of coupling to vector bosons

- gives information about Hbb Yukawa coupling – which is difficult in standard searches

2) CED Higgs production in MSSM

- in MSSM the signal yields are greatly enhanced

- in MSSM it gives information about Higgs sector

- in MSSM the Higgs width may be directly measured (for large $\tan\beta$)

Update of the 2007 analysis:

- background NLO CED $gg \rightarrow bb$

- LEP/Tevatron exclusion regions (HiggsBounds)

- improved theory calculations (FeynHiggs)

- new CDM benchmark planes (similar results as for $m_{H\max}$ benchmark)

3) CED Higgs production in 4 Generation Model

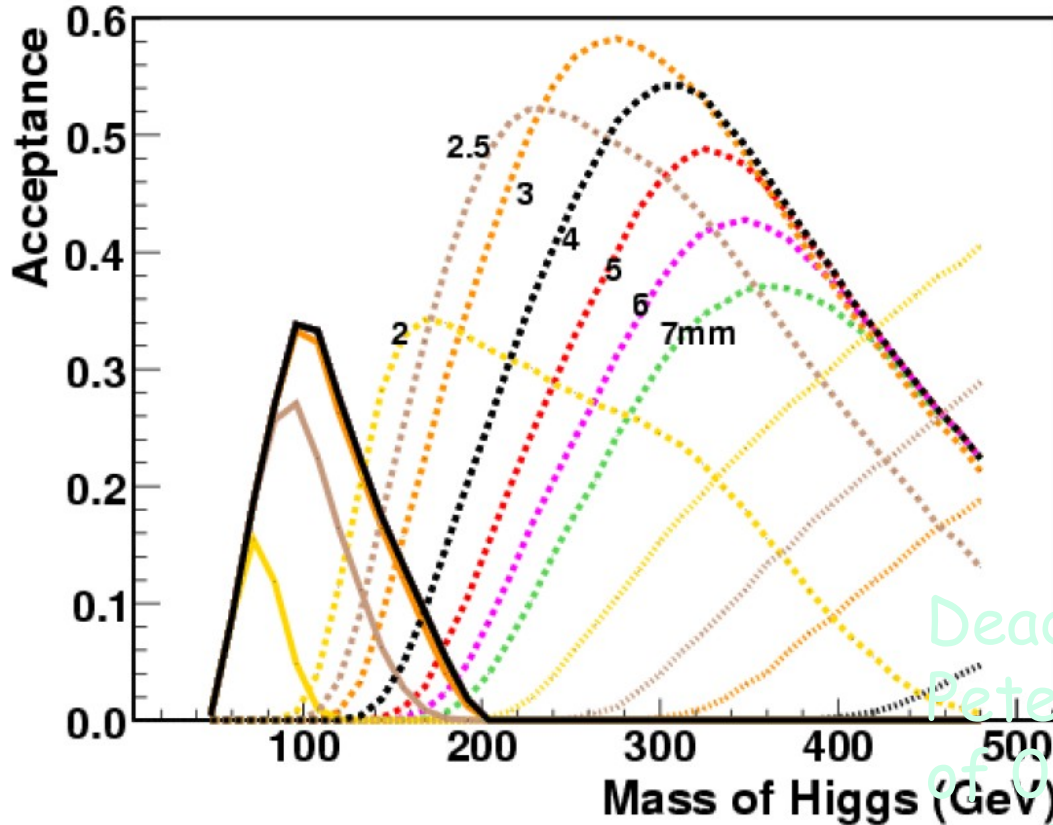
- LEP/Tevatron searches: $112 < M_H < 130$ GeV allowed

- good prospects for $H \rightarrow bb$ at low lumi, for $H \rightarrow \tau\tau$ at high lumi

B A C K U P S L I D E S

Acceptances

Acceptances depend heavily on the distance from the beam and dead space!
(if protons hit the dead space in 220 station, they are lost for 420 measurement)
Acceptance for 420+420, 420+220 and 220+220. Numbers mean total distances.
420 at 6 mm everywhere, 220 varying from 2mm to 7mm



Peter

Dead space = 1.1mm

220 at 2mm obstructs the tracking at 420 !

Dead space of 1.1 mm is too cautious.
Peter will make this plot for dead space of 0.5mm.
In the following analyses, dead space

15σ beam ~ 1.5 mm

(thin window (400 μ m) + safety offset (300 μ m) + edge (5 μ m) + alignment) ~ 0.7 mm

Basics of MSSM

Enlarged Higgs sector: Two Higgs doublets

$$H_1 = \begin{pmatrix} H_1^1 \\ H_1^2 \end{pmatrix} = \begin{pmatrix} v_1 + (\phi_1 + i\chi_1)/\sqrt{2} \\ \phi_1^- \end{pmatrix}$$

$$H_2 = \begin{pmatrix} H_2^1 \\ H_2^2 \end{pmatrix} = \begin{pmatrix} \phi_2^+ \\ v_2 + (\phi_2 + i\chi_2)/\sqrt{2} \end{pmatrix}$$

$$V = m_1^2 H_1 \bar{H}_1 + m_2^2 H_2 \bar{H}_2 - m_{12}^2 (\epsilon_{ab} H_1^a H_2^b + \text{h.c.}) \\ + \underbrace{\frac{g'^2 + g^2}{8}}_{\text{gauge couplings, in contrast to SM}} (H_1 \bar{H}_1 - H_2 \bar{H}_2)^2 + \underbrace{\frac{g^2}{2}}_{\text{gauge couplings, in contrast to SM}} |H_1 \bar{H}_2|^2$$

physical states: h^0, H^0, A^0, H^\pm Goldstone bosons: G^0, G^\pm

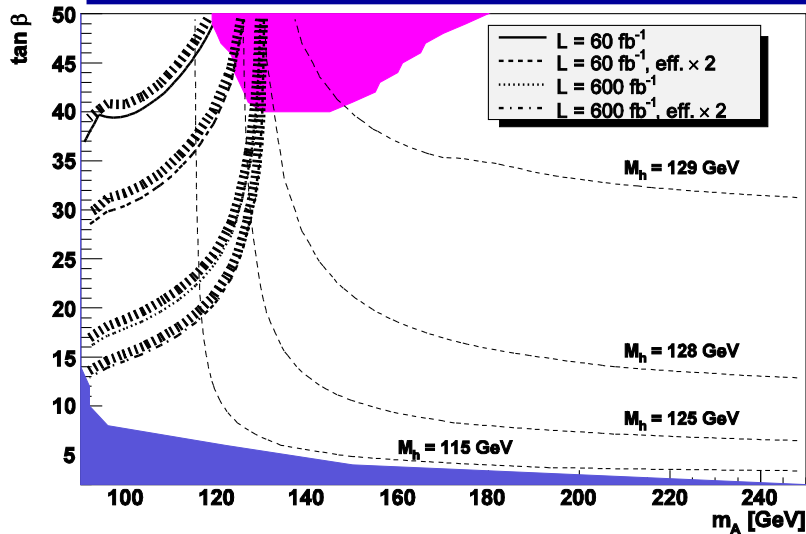
Input parameters: (to be determined experimentally)

$$\tan \beta = \frac{v_2}{v_1}, \quad M_A^2 = -m_{12}^2 (\tan \beta + \cot \beta)$$

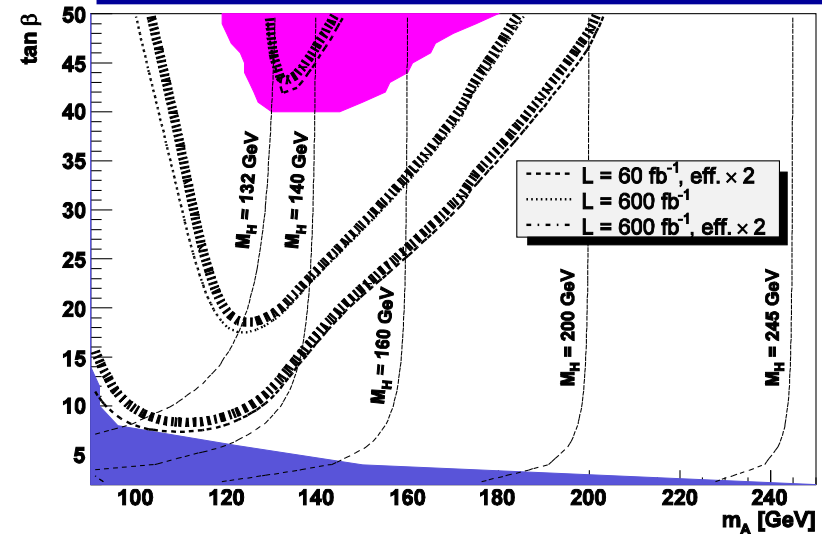
H → tau tau in MSSM

- In MSSM, $BR(h,H \rightarrow \tau\tau) \sim 10\%$, if decays to SUSY particles not allowed
- Studied in EPJC 53 (2008) 231: the same efficiencies assumed as in the bb case
- Background: QED $\gamma\gamma \rightarrow ll$ (suppressed by $p_T^{\text{prot}} > 0.2 \text{ GeV}$)
 CEP $gg \rightarrow gg$ (suppressed by $|\eta_{j1} - \eta_{j2}| < 1.1$ and $P(g/\tau) \sim 1/500$)

3σ: h → tautau, mhmax, μ=200 GeV



3σ: H → tautau, mhmax, μ=200 GeV



- In ATLAS the proper efficiencies now being estimated with full sim.
 - All tau-decays studied.
 - Backgrounds expected to be very low:
 - fully leptonic have high-pt leptons
 - fully hadronic have two tau-jets: very-few-particle jets going sharply back-to-back in φ
- No problems expected with triggering