

Overview of ionization cooling for a Muon Collider

Diktys Stratakis

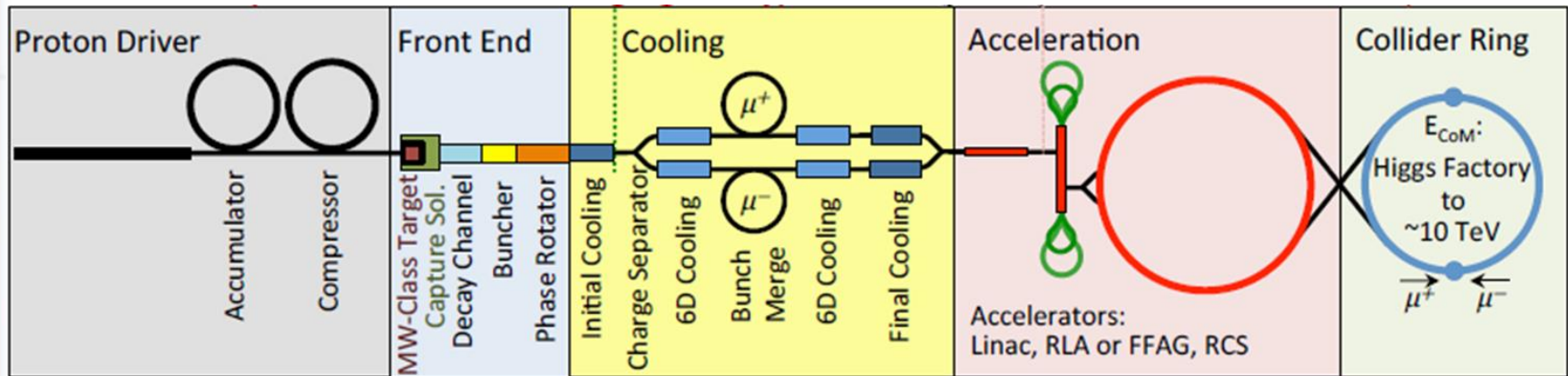
Fermi National Accelerator Laboratory

PITT PACC Workshop: Muon Collider Physics (online only)
December 01, 2020

Outline

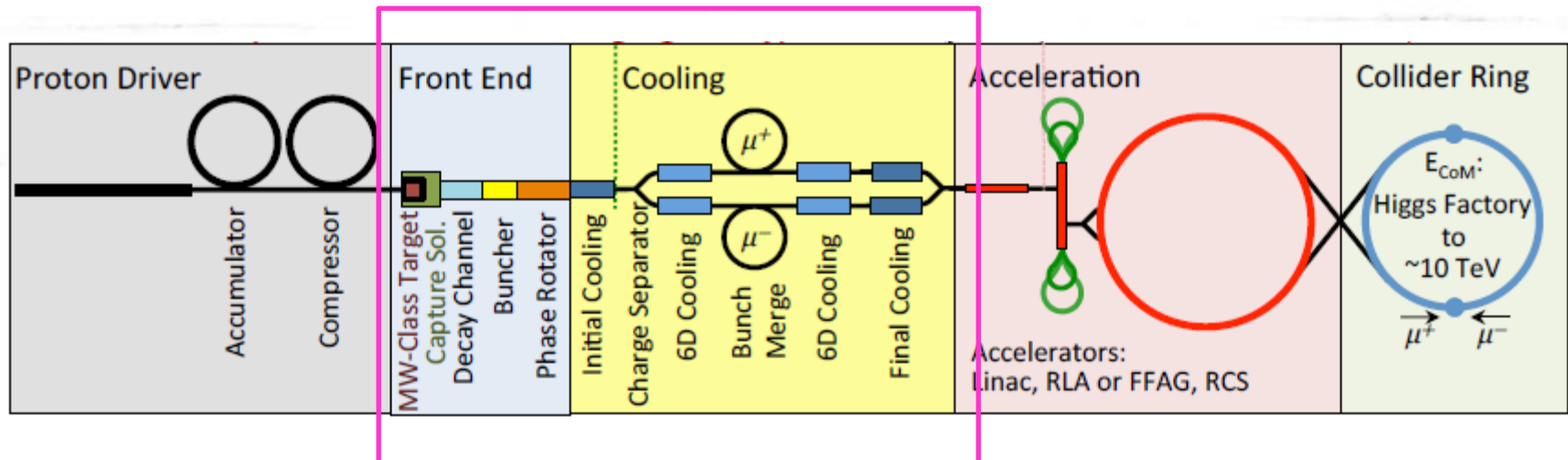
- Muon Collider overview
- Ionization cooling
 - Basic theory
 - Application to a Muon Collider
 - Simulation results
- Future work and possible improvements
- Summary

Muon Collider components



- The desired 6D emittance for a Muon Collider (MC) is 5-6 orders of magnitude less from the emittance of the muon beam at the production target
- As a result, significant “muon cooling” is required.

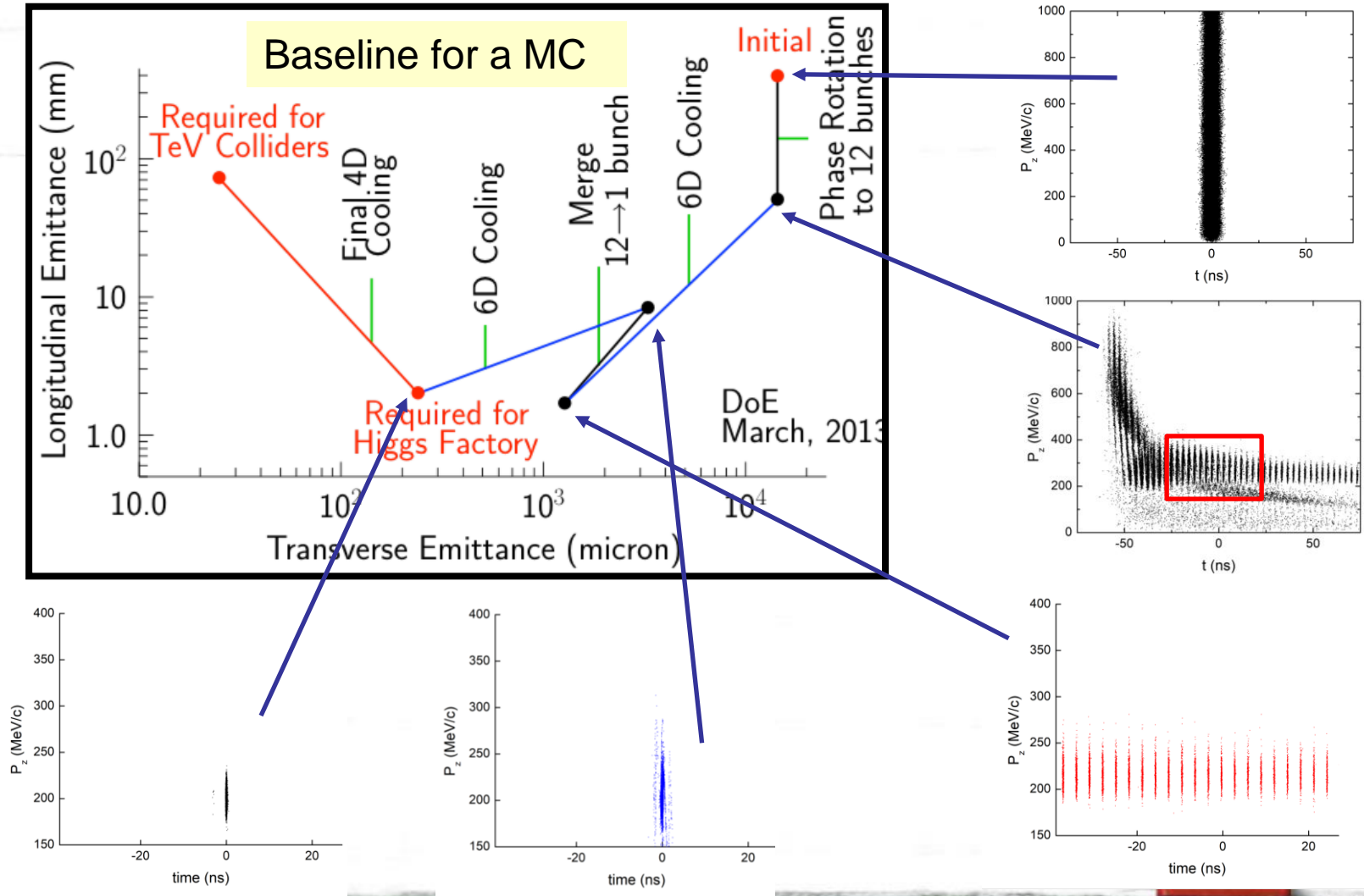
Cooling scheme for a Muon Collider



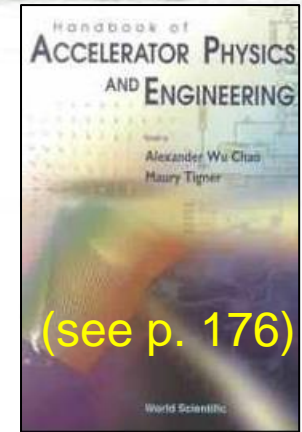
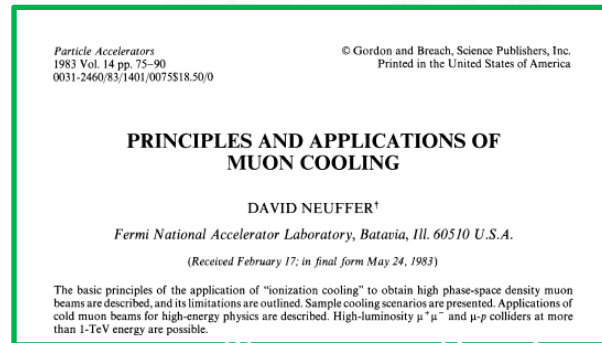
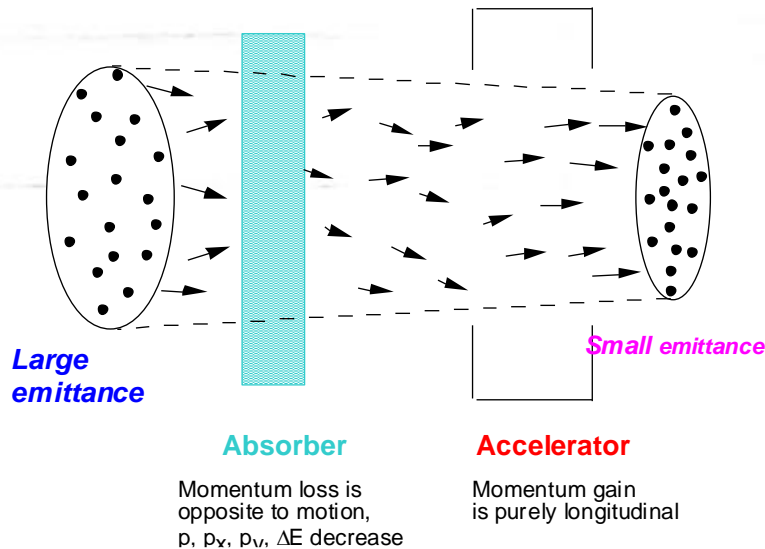
- Front-end produces 21 well aligned muon bunches
- Two sets of 6D cooling schemes
 - One before recombination (trans $\epsilon \approx 1.5$ mm)
 - One after recombination (trans $\epsilon \approx 0.30$ mm or less)
- Final cooling (if necessary)

Focus of
this talk

Muon Collider cooling baseline



Ionization cooling formalism (1)



Energy loss
term

Multiple scattering term

- Transverse cooling:
- Minimum emittance:

$$\frac{d\varepsilon_T}{ds} = -\frac{1}{\beta^2 E} \frac{dE}{ds} \varepsilon_T + \frac{\beta \gamma \beta_T}{2} \frac{d\theta_0^2}{ds}$$

$$\varepsilon_T^{\text{eq}} = \left(\frac{dE}{ds} \right)^{-1} \frac{\beta_T (13.6 \text{ MeV})^2}{2 \beta m_\mu c^2 L_R}$$

L_R : Radiation length

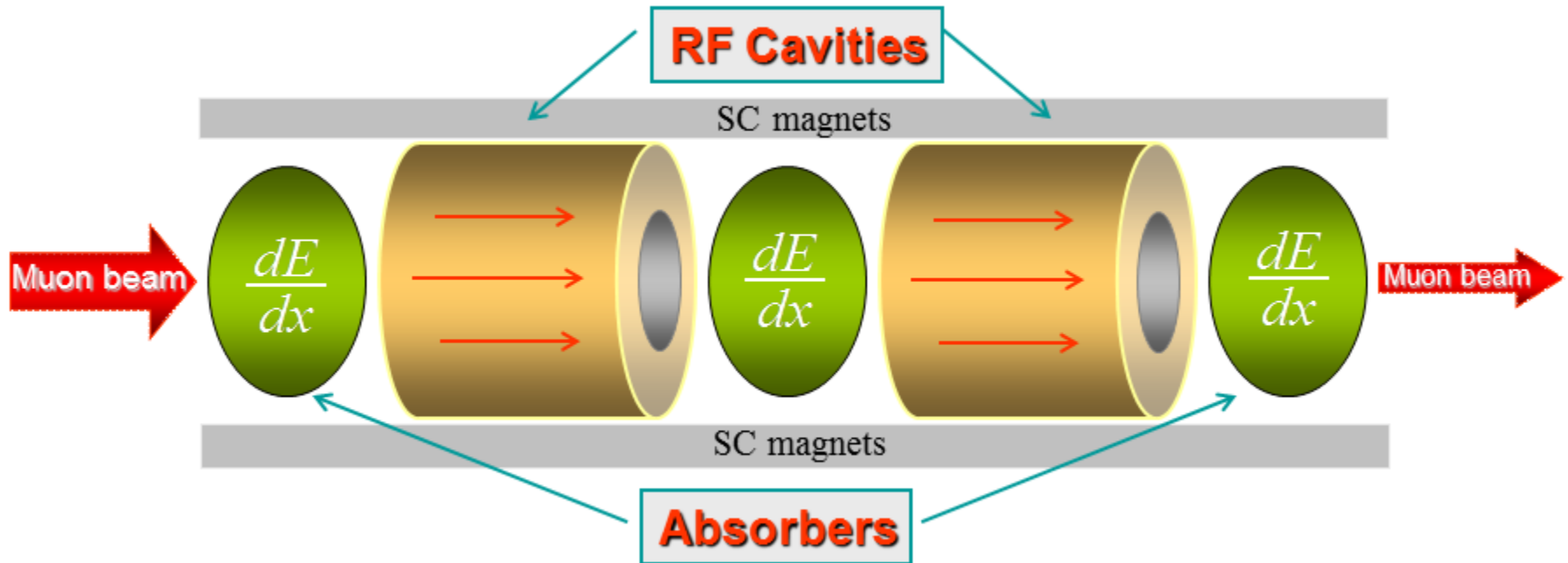
E : Muon energy

β_T : Transverse beta function

$\frac{dE}{ds}$: Energy loss

- Cooling can be controlled by material and magnetic focusing properties

Sample of a cooling channel

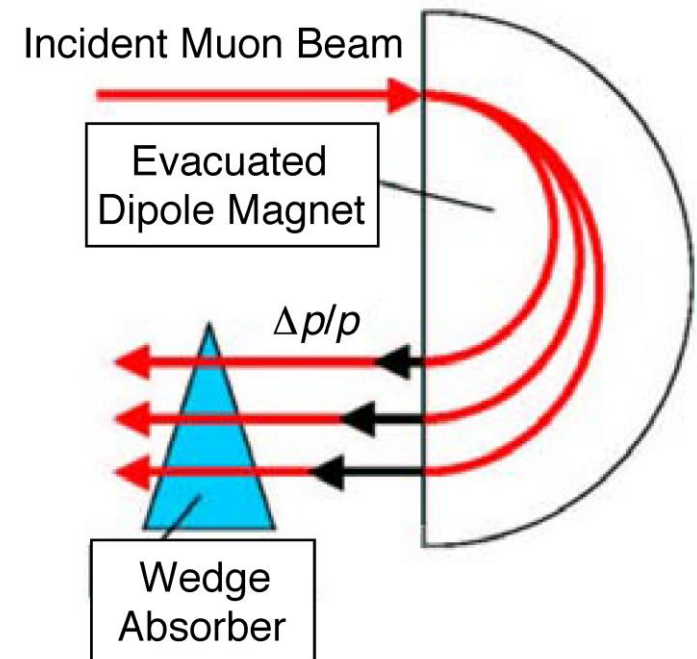


- Energy loss through ionization in the absorbers
- Restore the lost momentum in z with a longitudinal E-field
- A pillbox cavity is placed adjacent to the absorber
- This scheme cools the beam transversely only!

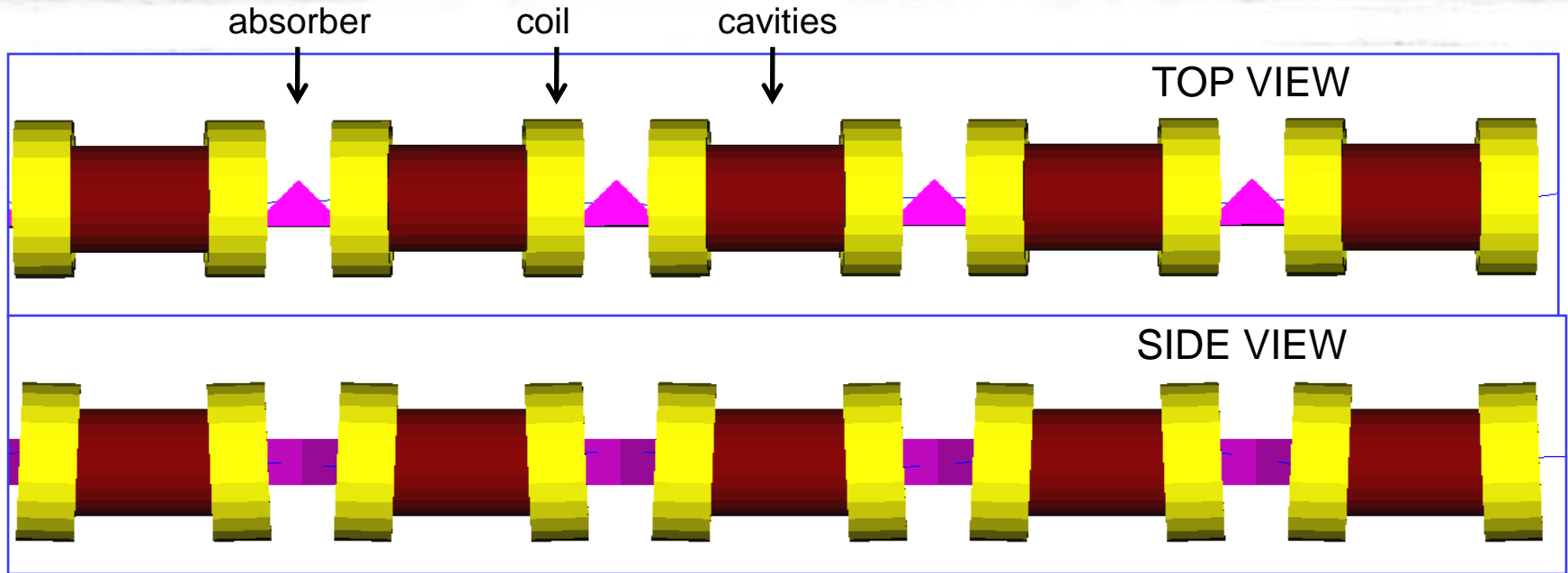
Ionization cooling formalism (2)

- Longitudinal cooling:
$$\frac{d\sigma_E^2}{ds} = -2 \frac{\partial \left(\frac{dE}{ds} \right)}{\partial E} \sigma_E^2 + \frac{d \langle \Delta E_{rms}^2 \rangle}{ds}$$

Cooling term
Straggling term
- Cooling occurs only if derivative:
$$\frac{\partial \left(\frac{dE}{ds} \right)}{\partial E} > 0$$
- Ionization loss does not naturally provide adequate longitudinal cooling
- Can be enhanced, if it is arranged that high energy muons lose more energy than low energy ones.



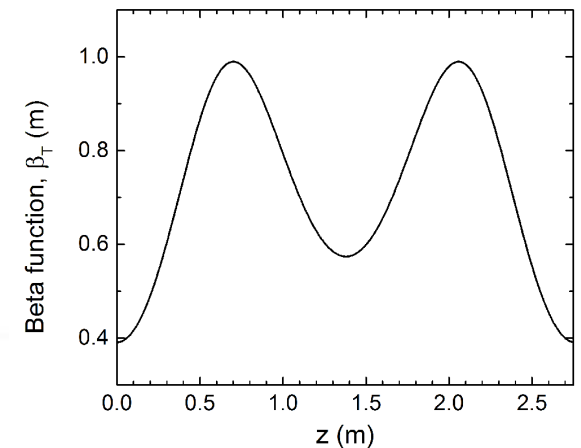
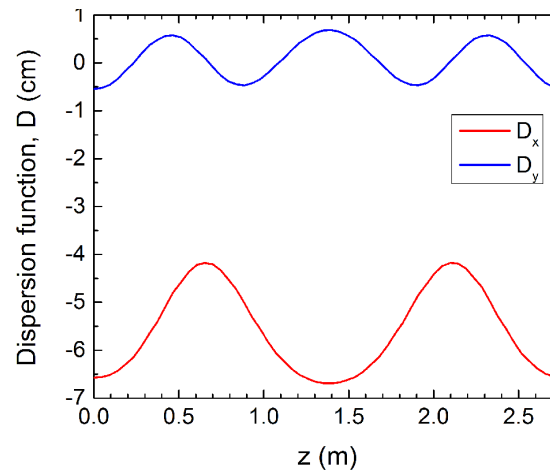
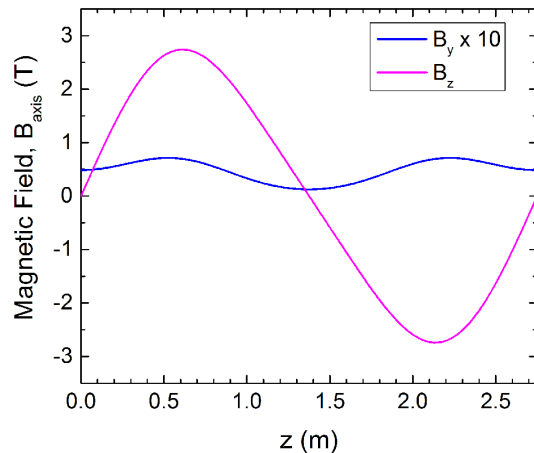
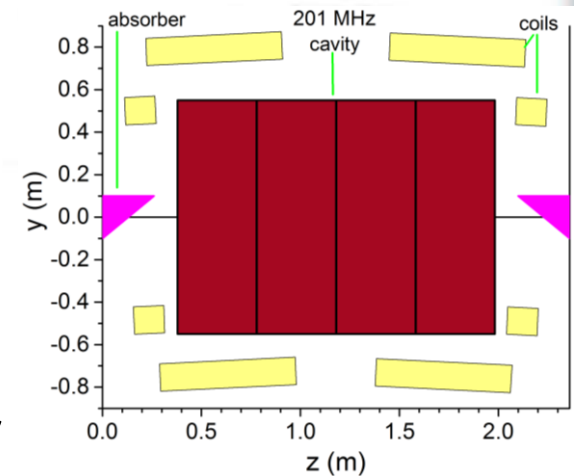
Cooling channel for a Muon Collider



- Straight geometry simplifies construction and relaxes several technological challenges
- Multiple stages with different cell lengths, focusing fields, rf frequencies to ensure fast cooling
-

Cooling channel: How it works

- Coils are slightly tilted to generate a B_y component
- This leads to dispersion, primarily in x .
- 6D cooling on wedge absorber
- Better, if beta is minimum at the absorber



MC cooling channel (8 stages)

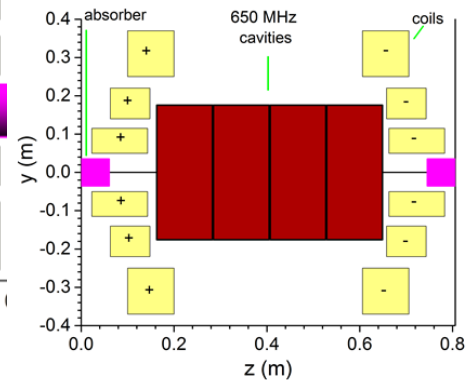
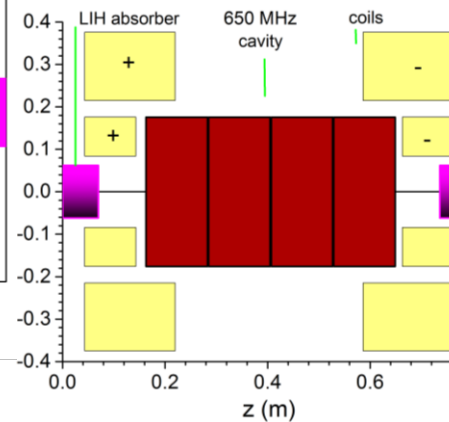
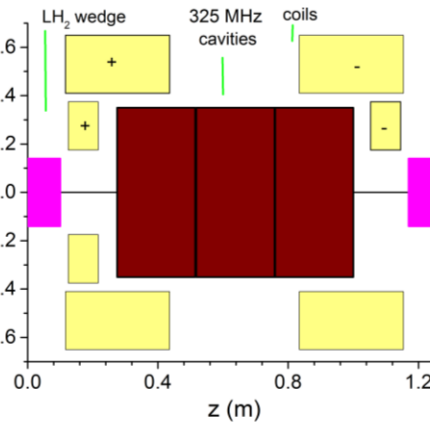
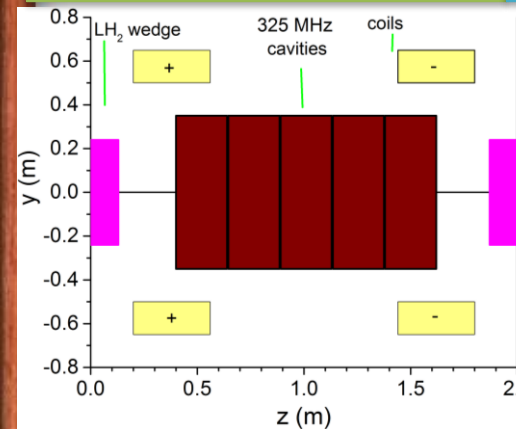


STAGE 2
64 m (32 cells)

STAGE 4
62.5 m (50 cells)

STAGE 6
62 m (77 cells)

STAGE 8
41.1 m (51 cells)



Absorber
TOP VIEW

LH or LiH

3.7 T (8.4 T)

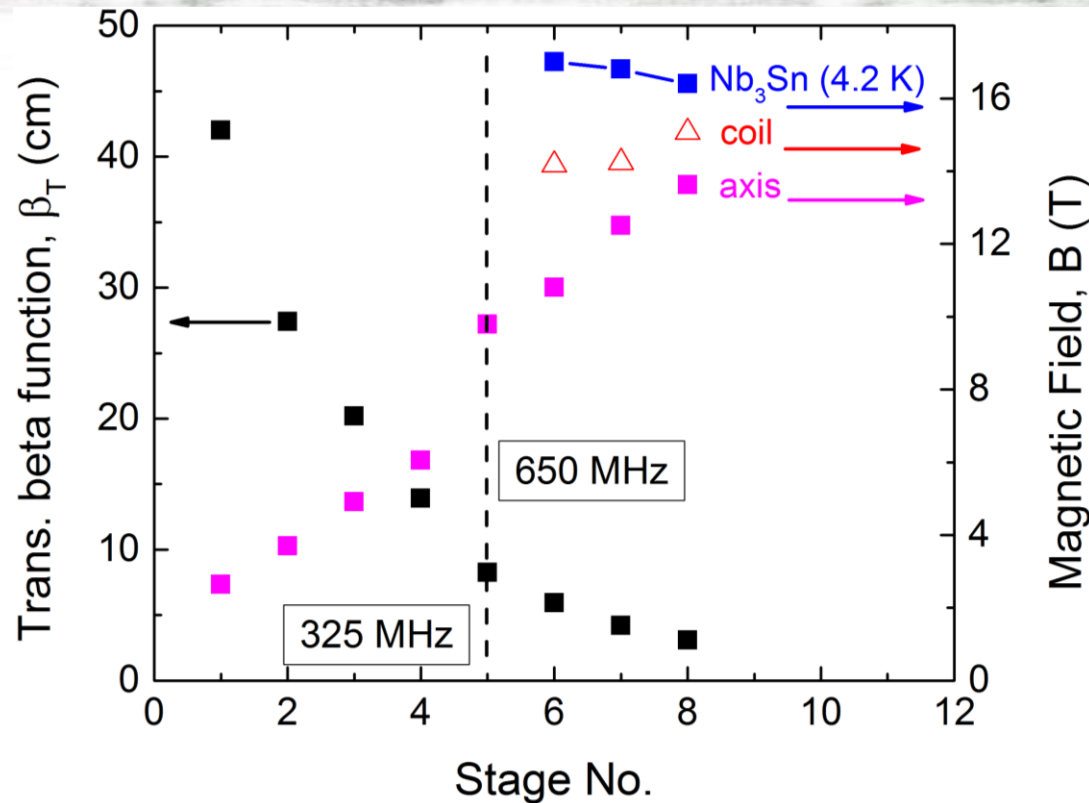
6.0 T (9.2 T)

10.8 T (14.2 T)

13.6 T (15.0 T)

Peak B-field on axis (coil)

Tapered lattice design: 8 stages

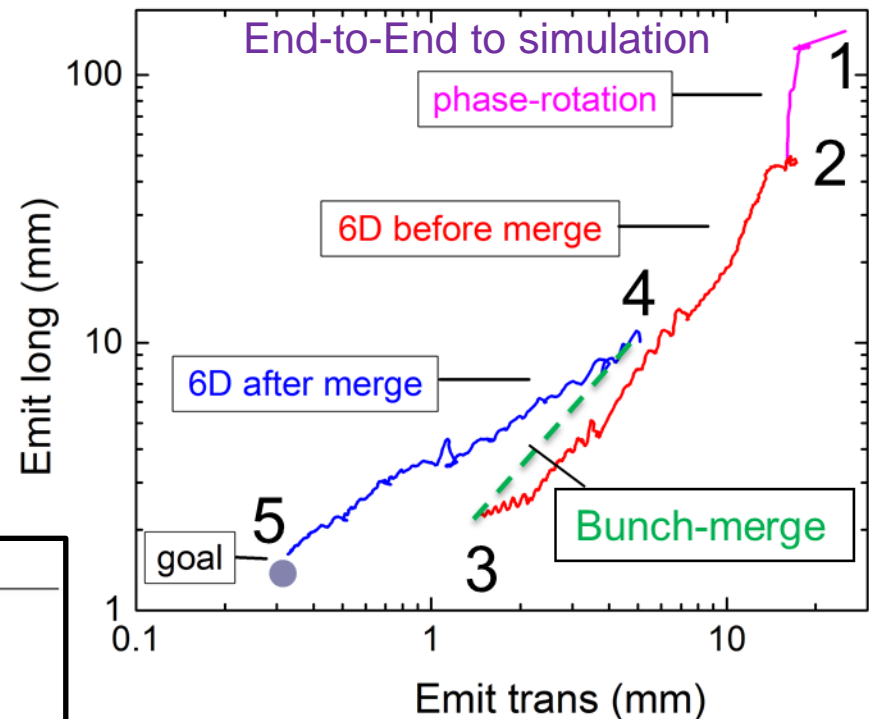


- We set two constraints in our (initial) design:
 - Peak fields on coils don't exceed Niobium Tin limits
 - Cavities within > 1 T operate at 50% of the achievable gradient at 0 T

Simulation results

- Complete end-to-end simulation from the target (point 1)
- 6D emittance reduction by five orders of magnitude (point 5)
- Achieved emittances and transmissions specified by MAP
- Overall distance ~ 900 m

Parameters end of cooling channel	MAP Goal	Channel
Emit., Trans. (mm)	0.30	0.28
Emit., Long. (mm)	1.50	1.57
Particles #	4.7×10^{12}	5.9×10^{12}



PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS **18**, 031003 (2015)

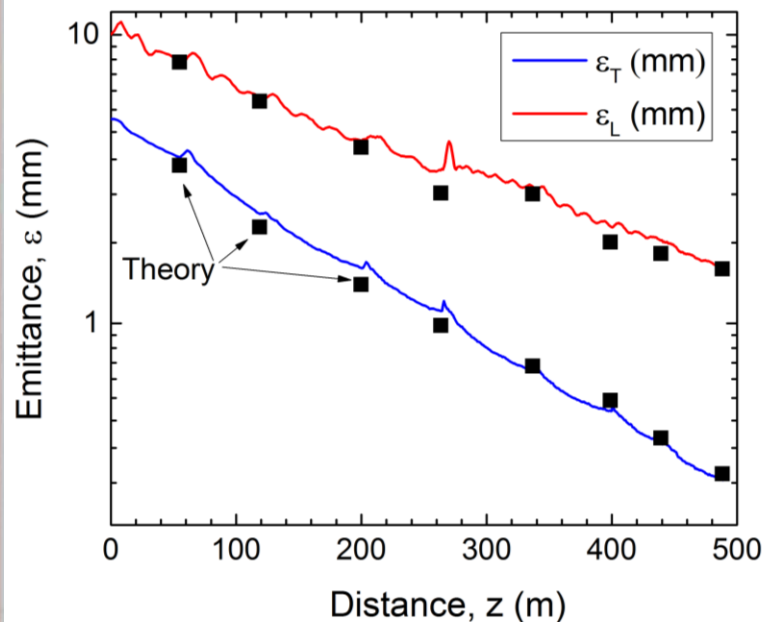


Rectilinear six-dimensional ionization cooling channel for a muon collider:
A theoretical and numerical study

Diktys Stratakis and Robert B. Palmer
Brookhaven National Laboratory, Upton, New York 11973, USA
(Received 25 September 2014; published 6 March 2015)

Theory and simulation

- Cooling channel performance can be predicted theoretically
- We found good agreement between simulation & theory



Transverse Cooling:

$$\frac{d\varepsilon_T}{ds} = -\frac{g_T}{\beta^2 E} \frac{dE}{ds} \varepsilon_T + \frac{\beta_T E_s^2}{2\beta^3 m_\mu c^2 L_R E}$$

$$\varepsilon_T^{\text{eq}} = \left(\frac{dE}{ds}\right)^{-1} \frac{\beta_T E_s^2}{2\beta g_T m_\mu c^2 L_R}$$

$$g_T = 1 - D/w$$



Longitudinal Cooling:

$$\frac{d\varepsilon_L}{ds} = -\frac{g_L}{\beta^2 E} \frac{dE}{ds} \varepsilon_L + \frac{\beta_L d\langle\Delta E^2\rangle}{2 ds}$$

$$\varepsilon_L^{\text{eq}} = \left(\frac{dE}{ds}\right)^{-1} \frac{\beta^2 E \beta_L d\langle\Delta E^2\rangle}{2 g_L ds}$$

$$g_L = \frac{2\gamma^2 - 2\ln[K(\gamma^2 - 1)]}{\gamma^2 \ln[K(\gamma^2 - 1)] - (\gamma^2 - 1)} + \frac{D}{w}$$

Emittance evolution:

$$\varepsilon^{\text{calc}}(s) = \varepsilon^{\text{eq}} + (\varepsilon^0 - \varepsilon^{\text{eq}}) \exp\left(-\frac{s}{s^{\text{calc}}}\right)$$

$$s_T^{\text{calc}} = \frac{\beta^2 E}{g_T} \left\langle \frac{dE}{ds} \right\rangle^{-1}$$

$$s_L^{\text{calc}} = \frac{\beta^2 E}{g_L} \left\langle \frac{dE}{ds} \right\rangle^{-1}$$

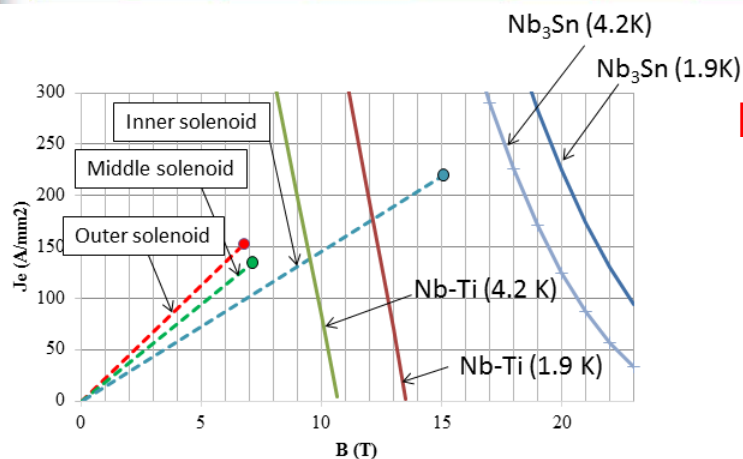
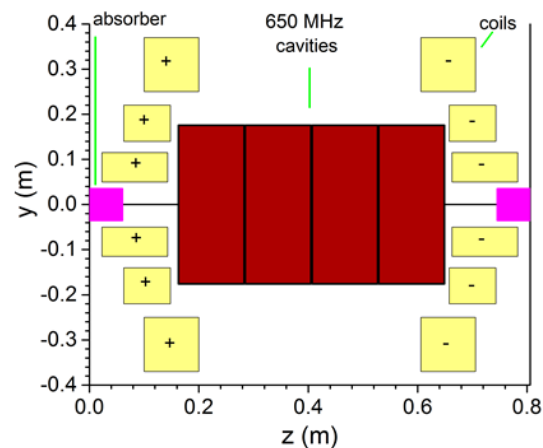
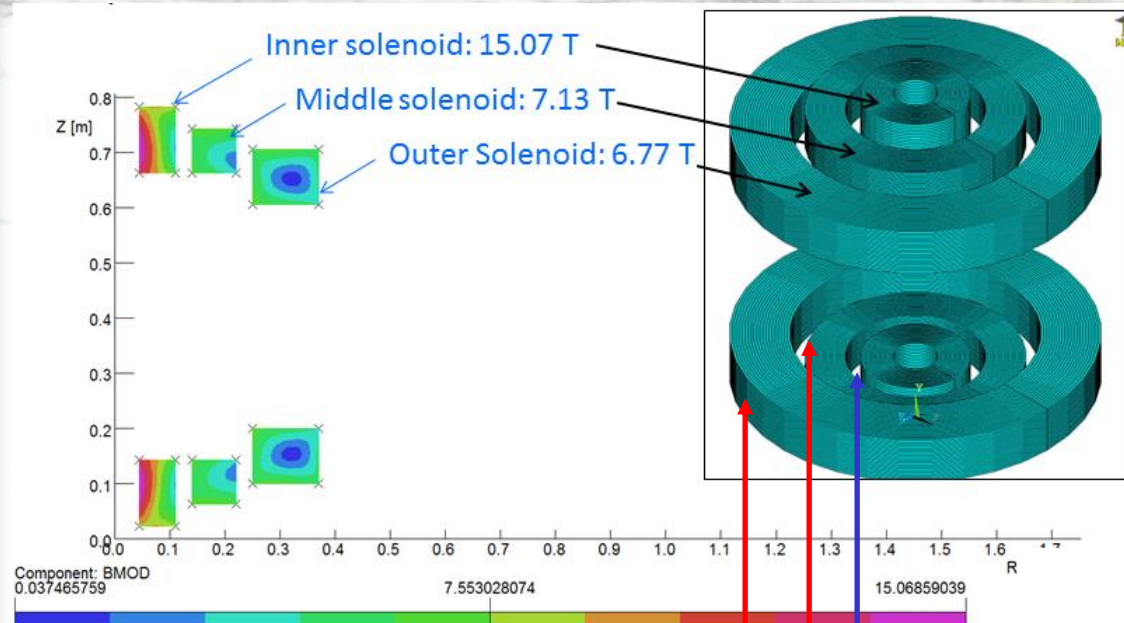
TUPME021

Proceedings of IPAC2014, Dresden, Germany

THEORETICAL FRAMEWORK TO PREDICT EFFICIENCY OF IONIZATION COOLING LATTICES*

Diktys Stratakis[#], Brookhaven National Laboratory, Upton, NY, USA
David Neuffer, Fermi National Accelerator Laboratory, Batavia, IL, USA

Magnet technology



NbTi
Nb₃Sn

	% of the load line at operational current		
	Inner solenoid	Middle solenoid	Outer solenoid
Nb-Ti @ 4.2 K	-	76%	74%
Nb-Ti @ 1.9 K	-	59%	58%
Nb3Sn @ 4.2 K	88%	-	-
Nb3Sn @ 1.9 K	81%	-	-

WEPR1103

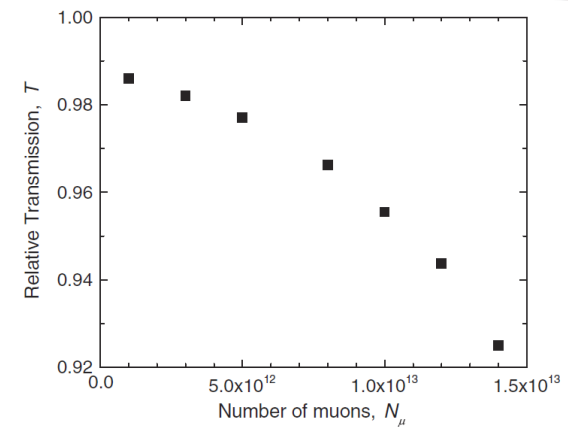
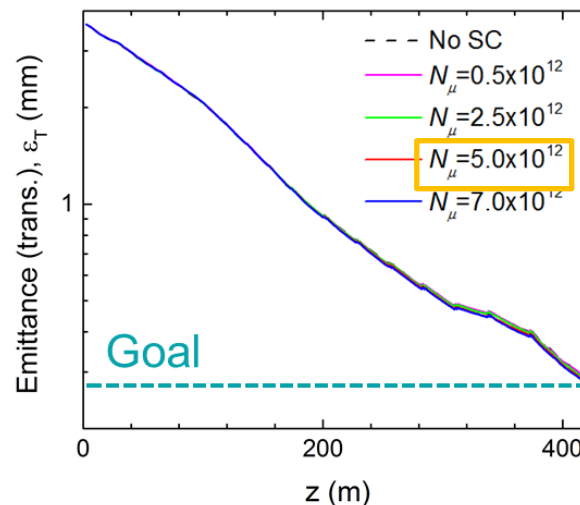
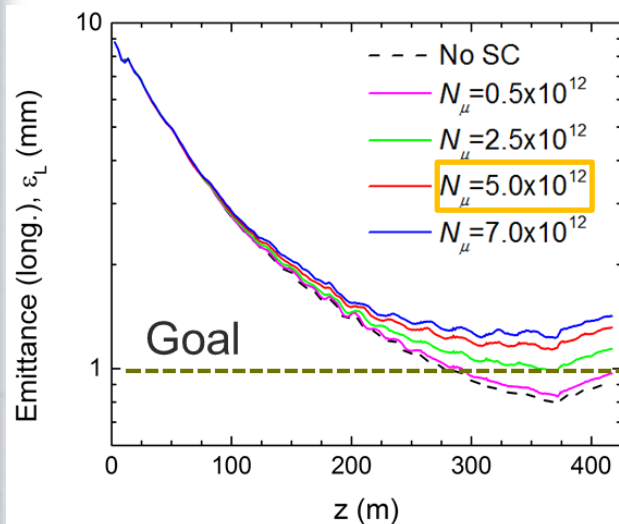
Proceedings of IPAC2014, Dresden, Germany

MAGNET DESIGN FOR A SIX-DIMENSIONAL RECTILINEAR COOLING CHANNEL - FEASIBILITY STUDY*

H. Witte[†], D. Stratakis, J. S. Berg, R. B. Palmer, Brookhaven National Laboratory, Upton, NY, USA
F. Borgnolutti, Lawrence Berkeley National Laboratory, Berkeley, CA, USA

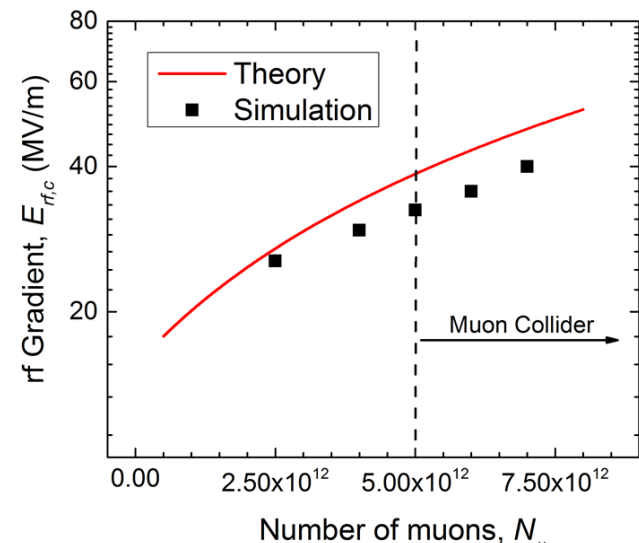
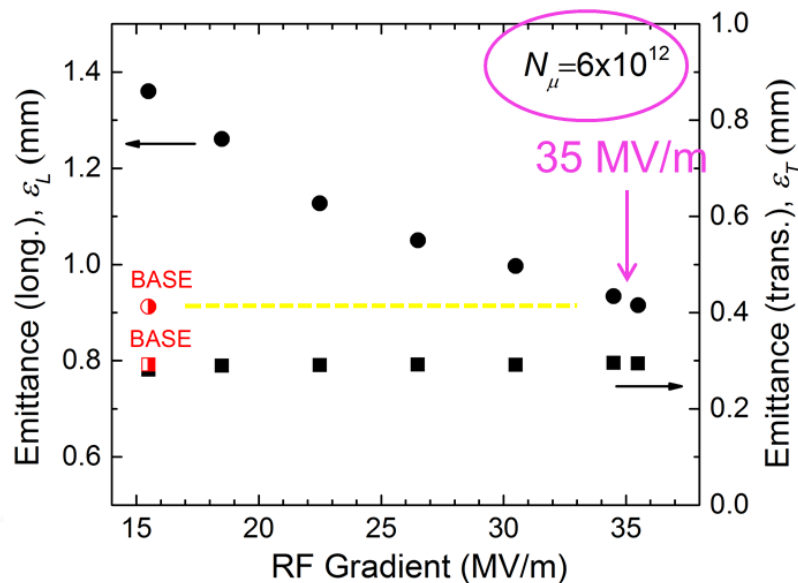
Influence of space-charge

- At the end of cooling, 5×10^{12} muons are squeezed within a 2 cm rms bunch. There is a concern for space-charge (SC)
- Simulation revealed that SC causes particle loss & longitudinal emittance growth



Space-charge compensation

- Increasing the rf gradient can compensate SC. Compensation gradient is coupled to the beam intensity
- For a MC to obtain a $\varepsilon_L \sim 1.0$ mm the rf gradient of a 805 MHz cavity needs to surpass 32 MV/m



PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS 18, 044201 (2015)

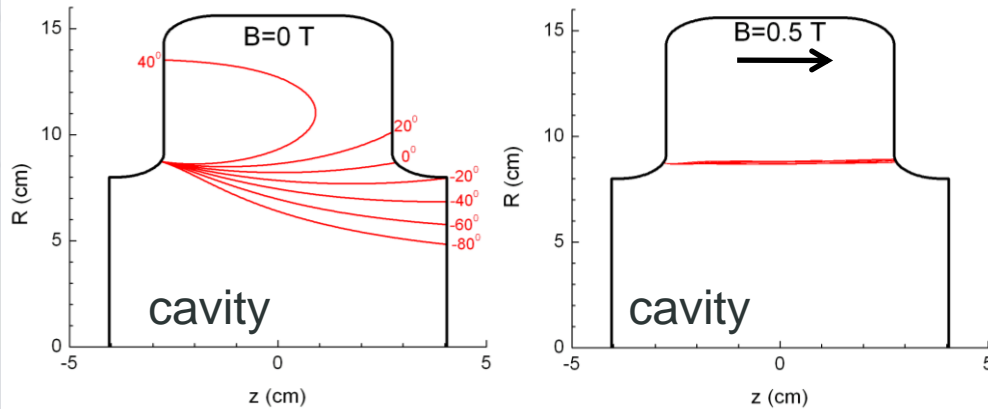


Influence of space-charge fields on the cooling process of muon beams

Diktys Stratakis and Robert B. Palmer
Brookhaven National Laboratory, Upton, New York 11973, USA

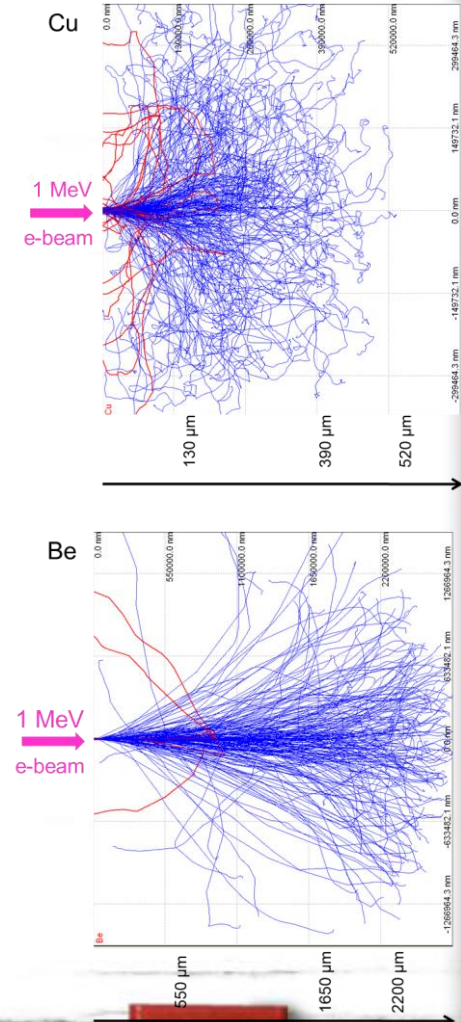
David P. Grote
Lawrence Livermore National Laboratory, Livermore, California 94550, USA
(Received 15 November 2014; published 7 April 2015)

Operation of rf cavities in B-fields

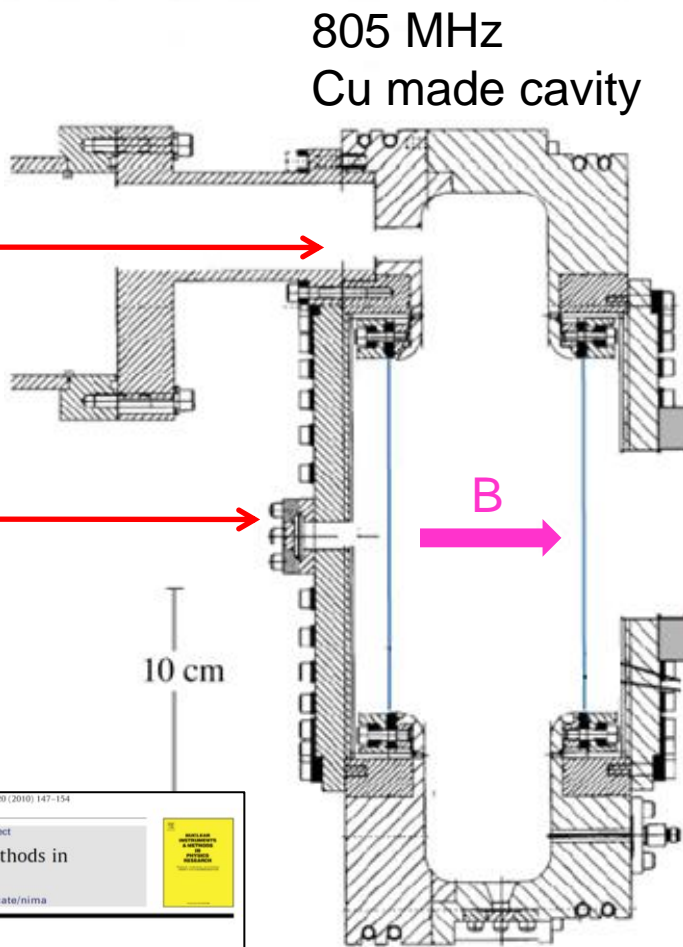


Damage on a 805 MHz rf cavity immersed in a multi-T magnetic field.

- Electrons impact rf surface and deposit heat in a small volume
- Surface damage via pulsed heating
- Effect is amplified with a B-field: A Cu made cavity may be damaged when $B > 1\text{ T}$
- Solution: Use dense materials like Be

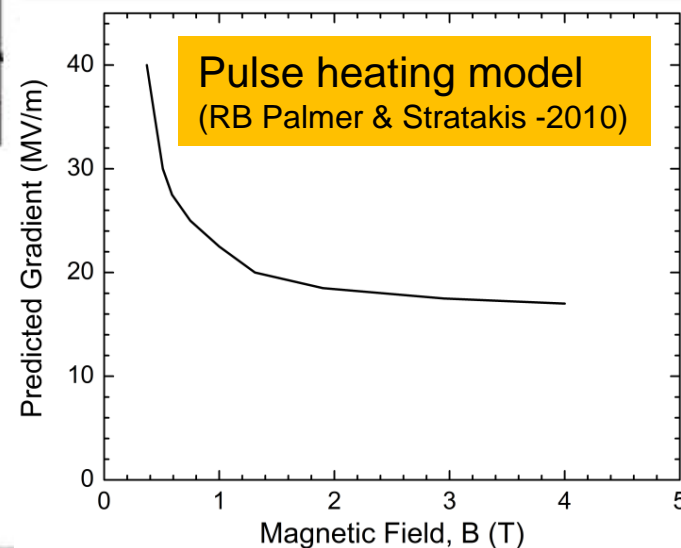
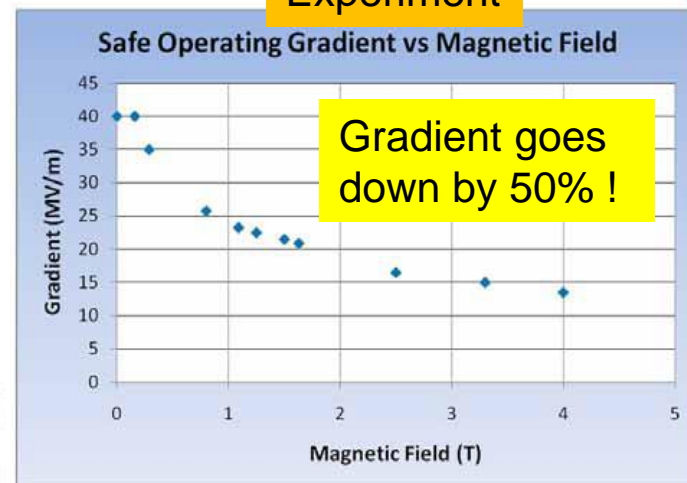


Data vs. model predictions

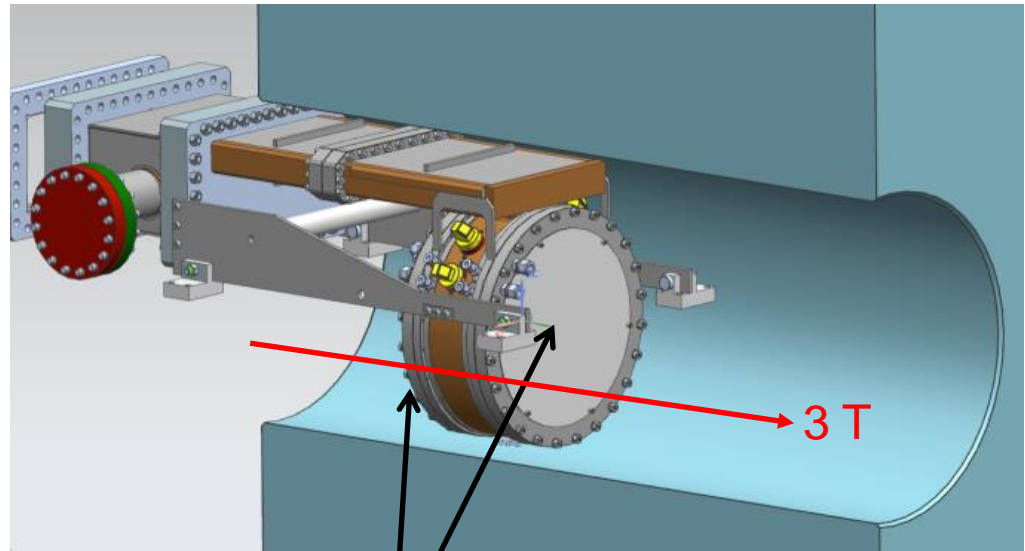
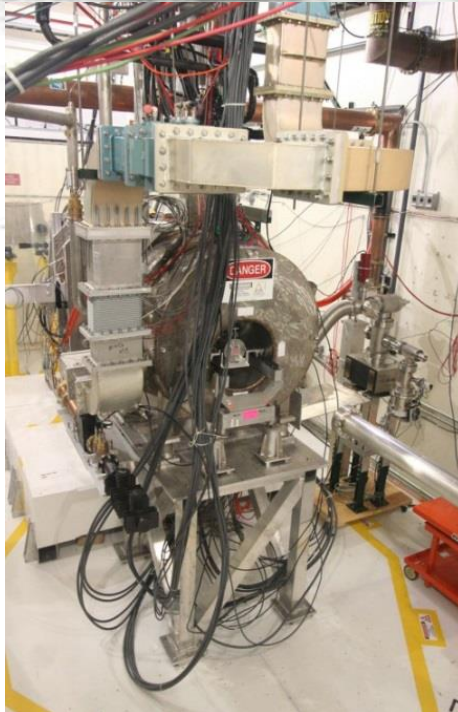


805 MHz
Cu made cavity

Experiment



Modular cavity test: A game changer



removable plates (Cu, Al, Be)

Material	B -field (T)	SOG (MV/m)	BDP ($\times 10^{-5}$)
Cu	0	24.4 ± 0.7	1.8 ± 0.4
Cu	3	12.9 ± 0.4	0.8 ± 0.2
Be	0	41.1 ± 2.1	1.1 ± 0.3
Be	3	$> 49.8 \pm 2.5$	0.2 ± 0.07
Be/Cu	0	43.9 ± 0.5	1.18 ± 1.18
Be/Cu	3	10.1 ± 0.1	0.48 ± 0.14

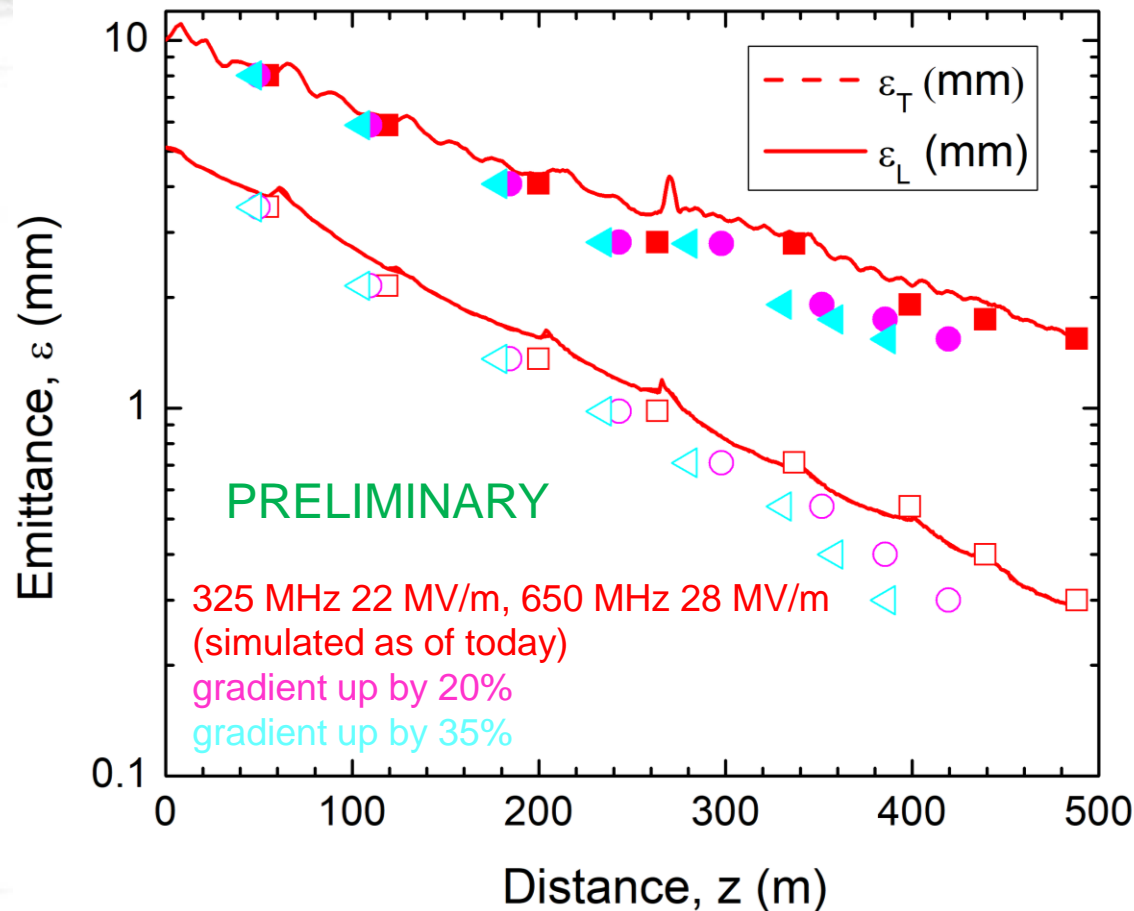
PHYSICAL REVIEW ACCELERATORS AND BEAMS **23**, 072001 (2020)

Operation of normal-conducting rf cavities in multi-Tesla magnetic fields for muon ionization cooling: A feasibility demonstration

D. Bowring[✉], A. Bross, P. Lane[✉], M. Leonova, A. Moretti, D. Neuffer[✉], R. Pasquinelli, D. Peterson[✉], M. Popovic, D. Stratakis, and K. Yonehara
Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA

- A Beryllium based cavity sustained a high gradient in the presence of multi-tesla B-fields!

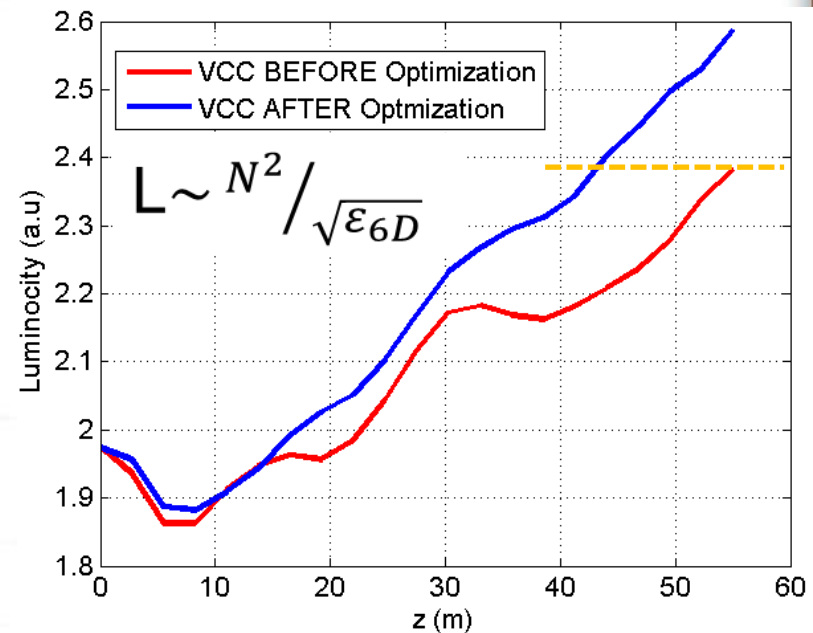
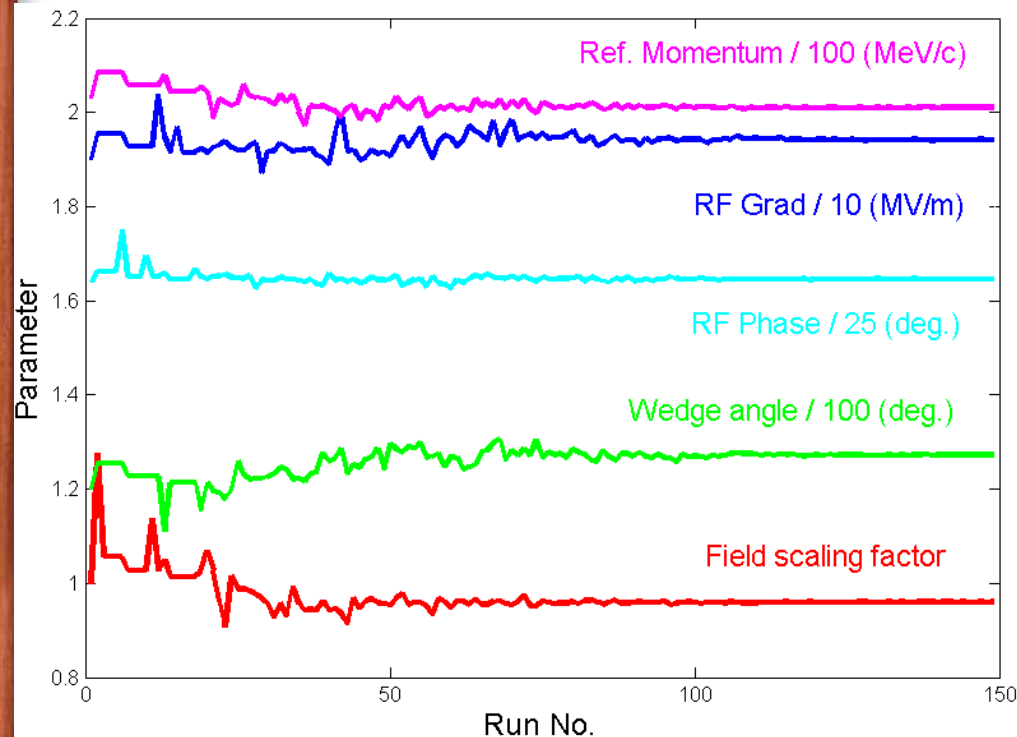
Future: Simulate with higher gradients



- Increasing the rf gradient can reduce the length of the cooling channel

Future: Multivariable optimization

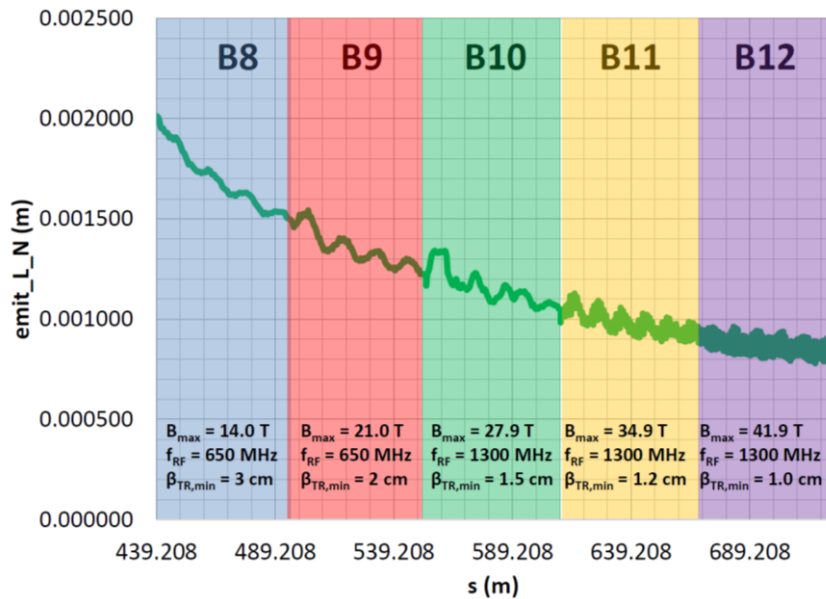
- Nelder-Mead algorithm: Objective is to maximize luminosity.
- Applied for VCC optimization: 8 parameters each time
- Promising results for first stage: 25% shorter channel!



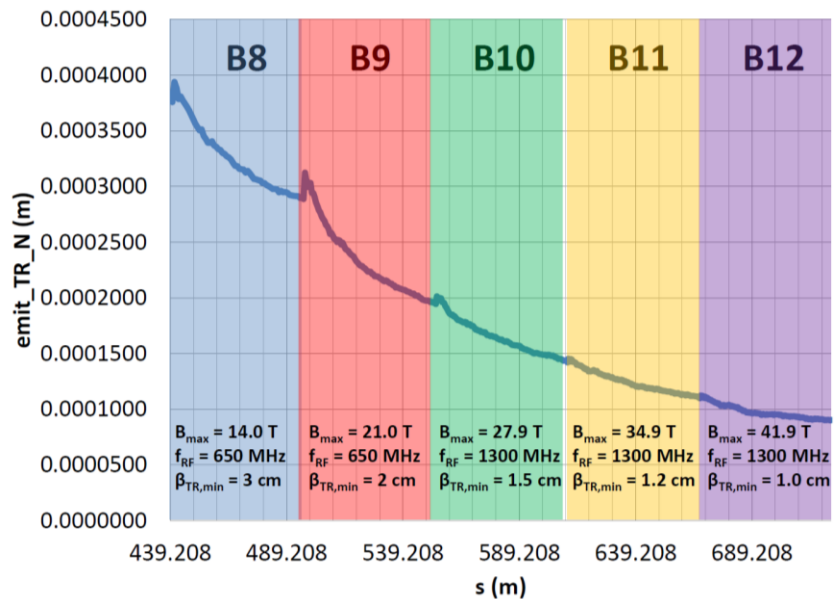
Future: Move towards higher fields

- The proposed channel can achieve substantial cooling if HTS magnet technology is considered

Longitudinal Cooling for Stages B8 - B12

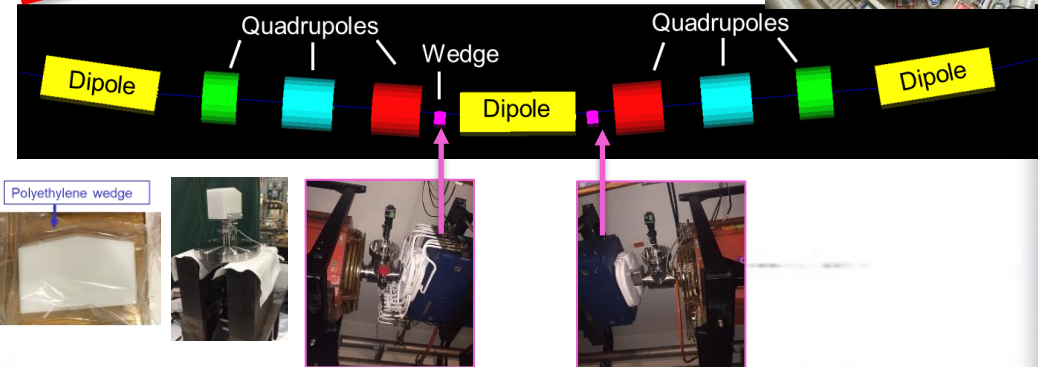
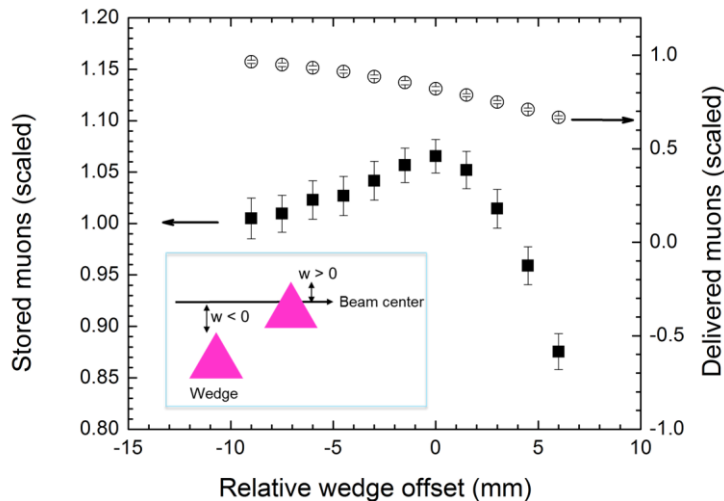
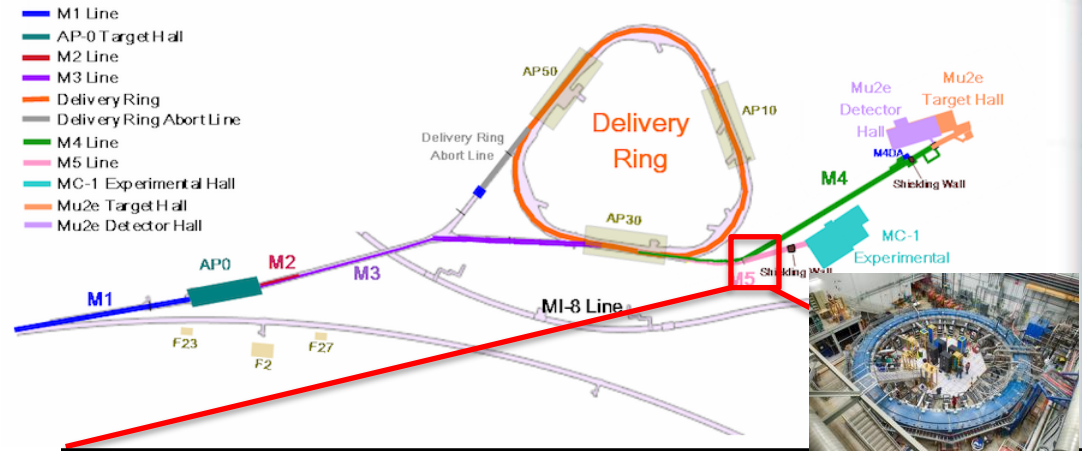
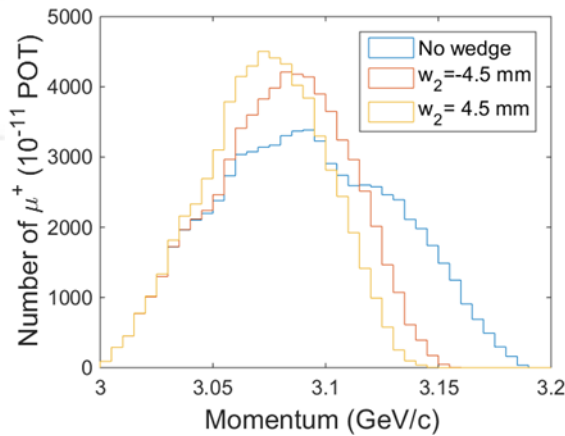


Transverse Cooling for Stages B8 - B12



Longitudinal cooling demonstration for the Muon g-2 Experiment

- Proof-of-principle experiment: Demonstrated 8% gain



PHYSICAL REVIEW ACCELERATORS AND BEAMS 22, 053501 (2019)

Application of passive wedge absorbers for improving the performance of precision-science experiments

Diktys Stratakis
Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA

Summary

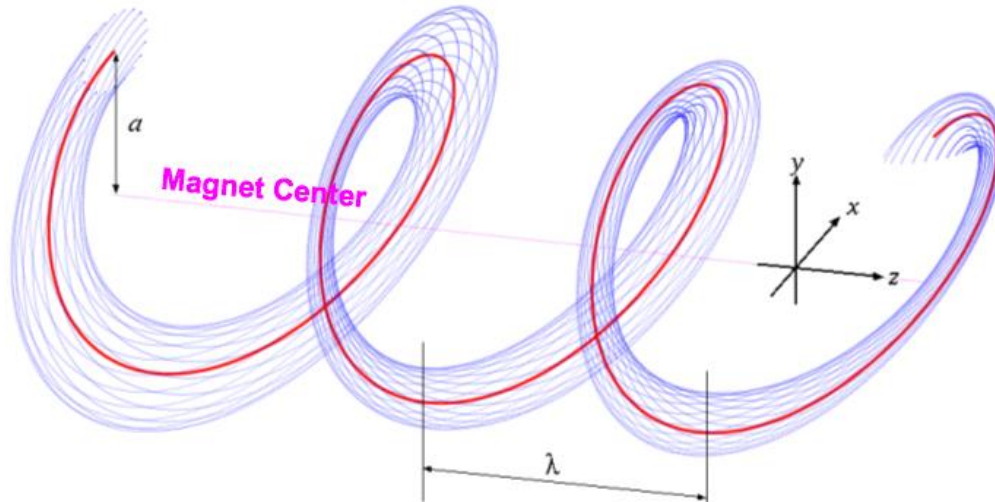
- We have a complete simulation model for a 6D cooling channel for a Muon Collider
- Beam at the end of the channel matches parameters specified by the Muon Accelerator Program within the limits of Nb-Sn magnet technology
- Further cooling is possible with HTS technology
- Operation of rf cavities in B-fields has been demonstrated:
 - Cooling channel can become shorter
 - Space-charge can be compensated -> higher transmission
- It is possible to reduce the channel length by >40% with a combination of higher rf gradients and multi-variable optimization

Further related work

- Neutrino factory cooling
 - D. Stratakis and D. Neuffer, Journal of Physics G: Nuclear and Particle Physics 41,
- Helical cooling channel
 - K. Yonehara, JINST 13, P09003 (2018)
- Final cooling
 - H. Sayed, Phys.Rev.ST Accel.Beams 18, 091001 (2015)
- Bunch merger
 - Y. Bao, Phys. Rev. Accel. Beams 19, 031001 (2016)
- Helical FOFO Snake
 - Y. Alexahin, JINST 13, P08013 (2018)

Beam dynamics in Helical Cooling Channel

Ya.S. Derbenev and R.P. Johnson, PRSTAB 8 041002 (2005)



Helical beam trajectories

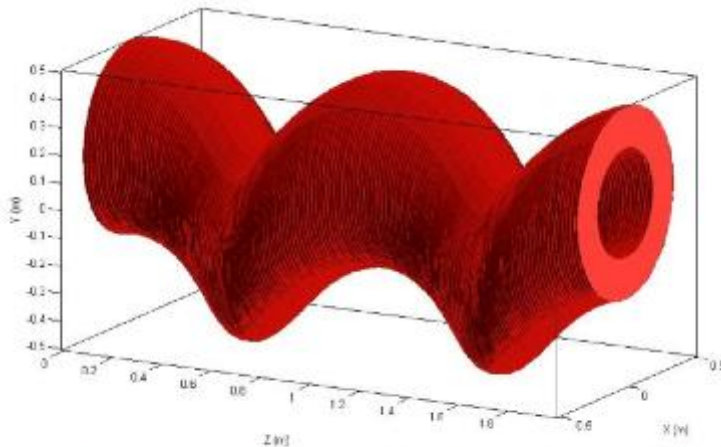


Figure 1: A helical solenoid generated with SolCalc.

Time- and z- independent Hamiltonian is a result from a helical coordinate system

- Helical coordinate system

$$\kappa = \frac{p_\phi}{p_z} = \frac{2\pi a}{\lambda} = k \cdot a$$

- Two Lorentz forces are applied on a muon and make a balance

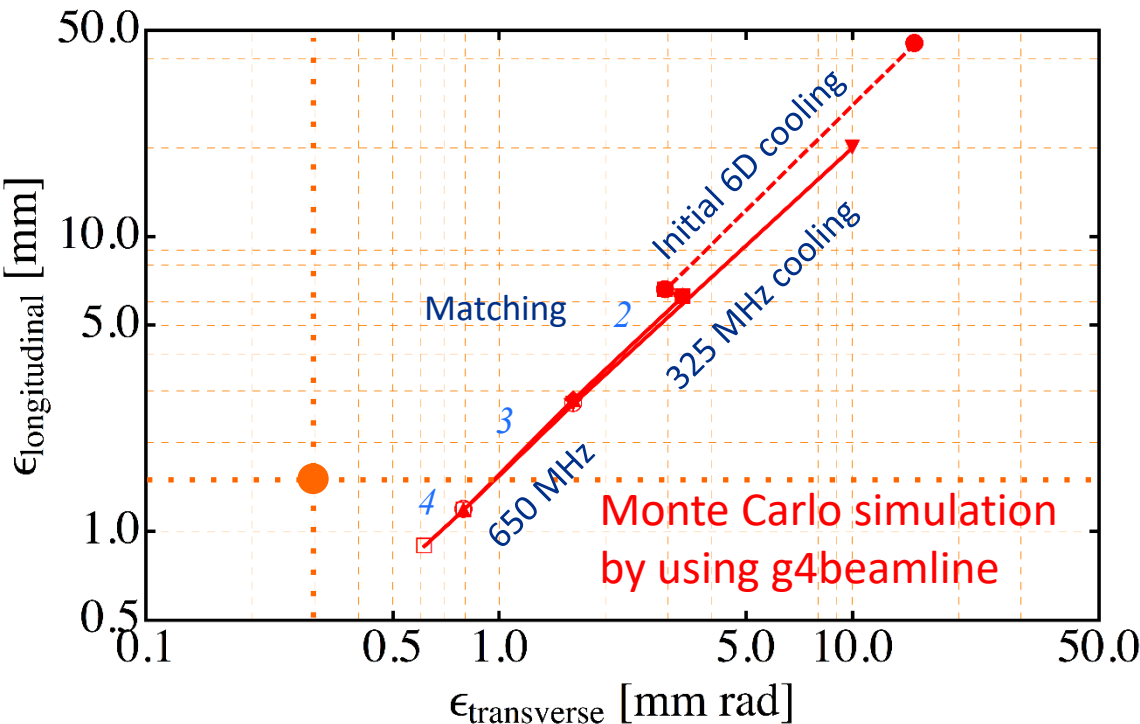
$$f_{balance} = \frac{e}{m_\mu} (p_z \cdot b_\phi - p_\phi \cdot b_z)$$

- Equation of motion of a reference (a red line in Top-Left Figure) is given

$$p(a) = \frac{\sqrt{1+\kappa^2}}{k} \cdot \left(b_z - \frac{1+\kappa^2}{\kappa^2} \cdot b_\phi \right)$$

- Solenoidal (b_z) and helical dipole (b_ϕ) fields are generated in a helical solenoid magnet
- Field gradient $b'_\phi = \frac{db_\phi}{dr}$ is naturally generated in the magnet and it induces focusing (a blue line in Top-Left Figure)

Simulated Cooling Performance



- RF parameter
 $E = 20 \text{ MV/m}$
 $\nu = 325 \text{ \& } 650 \text{ MHz}$
- Gas pressure
 $160 \text{ atm at } 300 \text{ K}$
 $43 \text{ atm at } 80 \text{ K}$
- Magnetic fields
 $B_z = 4 - 12 \text{ Tesla}$
- Equilibrium emittance
 $\epsilon_T = 0.6 \text{ mm (goal: } 0.3 \text{ mm)}$
 $\epsilon_L = 0.9 \text{ mm (goal: } 1.5 \text{ mm)}$
- Transmission (one cooling cycle)
 60%
- Channel length (one cooling cycle)
 280 meter

Helical cooling decrement can be explained by an old cooling theory

$$e_n(s) = (e_{n,0} - e_{n,equ}) e^{-L_x s} + e_{n,equ}$$

$$L_x = \frac{\frac{1}{2} g_x}{\frac{1}{2} b^2 E} \frac{dE}{ds} ; \text{cooling decrement}$$

$$e_{n,equ} = \frac{bg}{2} \left(\hat{b}_x S_x^2 \right) / \frac{1}{2} \frac{dE}{ds} ; \text{equilibrium emittance}$$

$$j = \int \frac{ds}{\hat{b}_x} \rightarrow \hat{b}_x = \frac{1}{w \hat{Q}_x} ; \text{beta function in HCC}$$

g_x ; partition of cooling

S_x^2 ; rms fraction due to stochastic process