

Particle-Flow Event Reconstruction from LEP to LHC

□ Outline

◆ Introduction : Basics of Particle-Flow Event Reconstruction

◆ First Lecture :

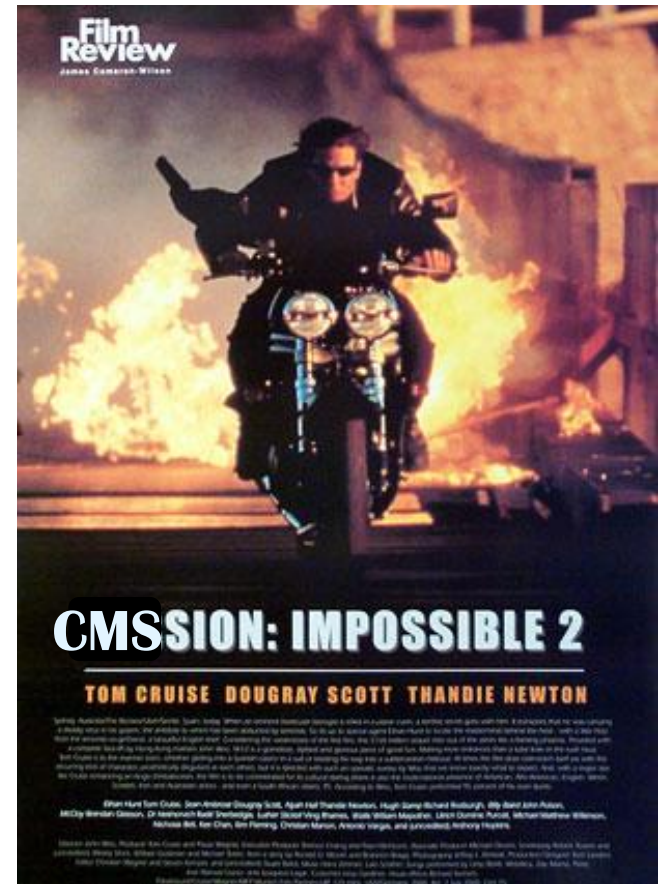
Second Lecture :



Intermezzo :

Lessons Learned ?

Exercises...

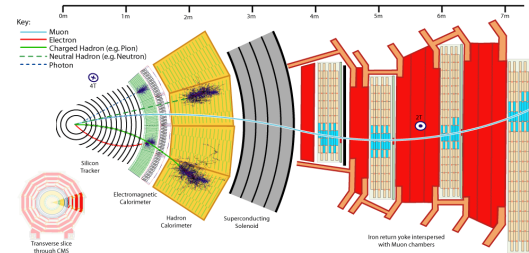


Basics of Particle-Flow Event Reconstruction (I)

□ What is Particle-Flow Event Reconstruction ?

- ◆ A better name could actually be

“PARTICLE FOLLOW RECONSTRUCTION”

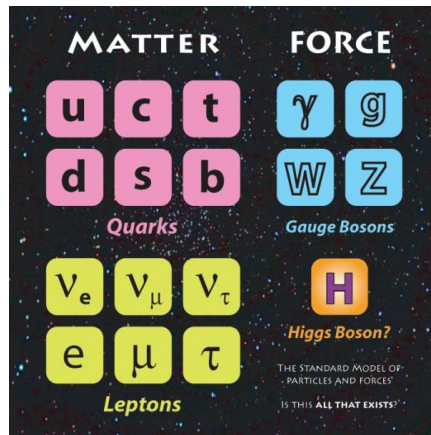


- ◆ Made easier by a smart detector design, carefully thought ahead of time
 - To be able to follow **individually** each particle arising from a collision
 - In their journey through the various sub-detectors
 - To be able to identify **each** particle
 - with their characteristic interaction in the various sub-detectors
 - To be able to measure the origin, direction, energy, charge of **each** particle
 - with an optimal combination of the measurements of all sub-detectors
 - with a decent accuracy (See later for the definition of “decent”)
 - Towards a global event description with a complete list of particles
 - As if it came directly from a Monte Carlo event generator

Basics of Particle-Flow Event Reconstruction (2)

□ What particles are to be followed ?

- ◆ Ideally : build a detector to follow all the standard model (SM) particles
 - By definition, any exotic particle decays to SM particles (+ dark matter)



+



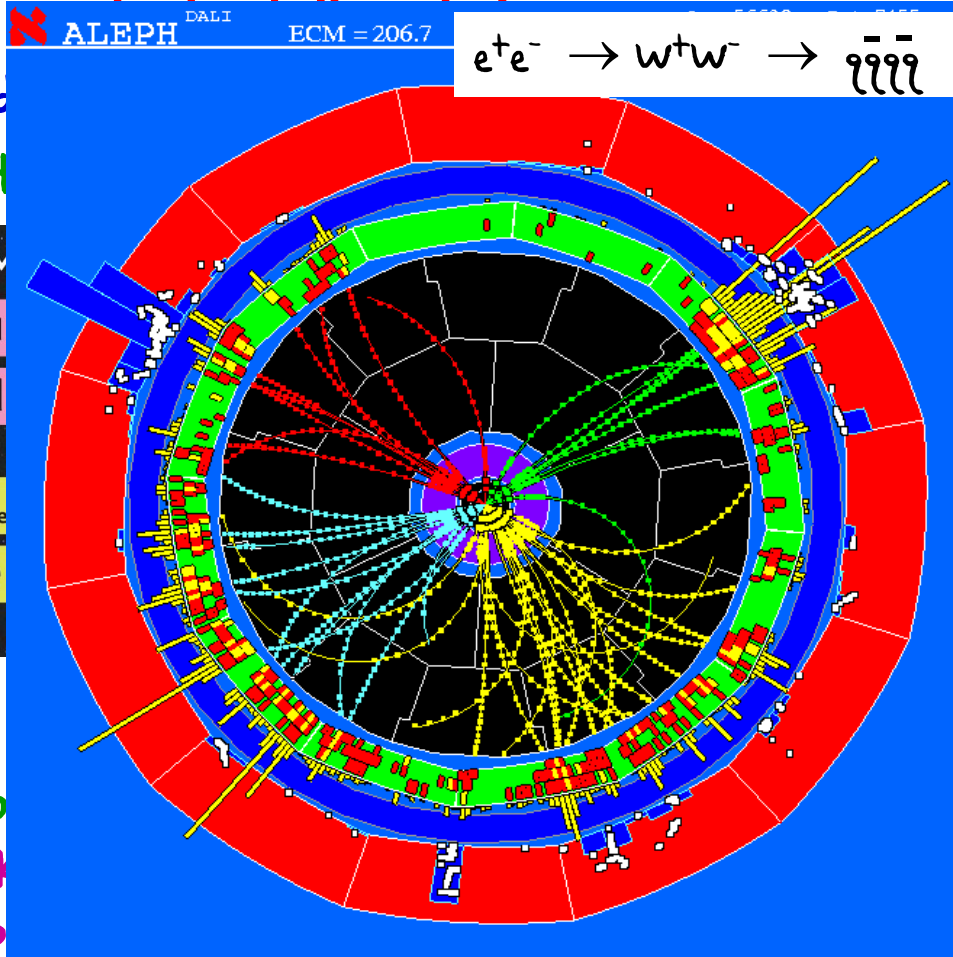
◆ Not quite ...

- Only e, γ, ν stable, and $\mu \rightarrow e \nu_e \nu_\mu$ but $\beta\gamma\tau \sim 6 \text{ Km/GeV}$, quasi stable
 - Quarks and gluons hadronize to give jets of hadrons, τ, W, Z decay to leptons and quarks, and H decay to pairs of every other particles

Basics of Particle-Flow Event Reconstruction (2)

□ what particles

- ◆ Ideally : build
- By definit



SM) particles
(+ dark matter)

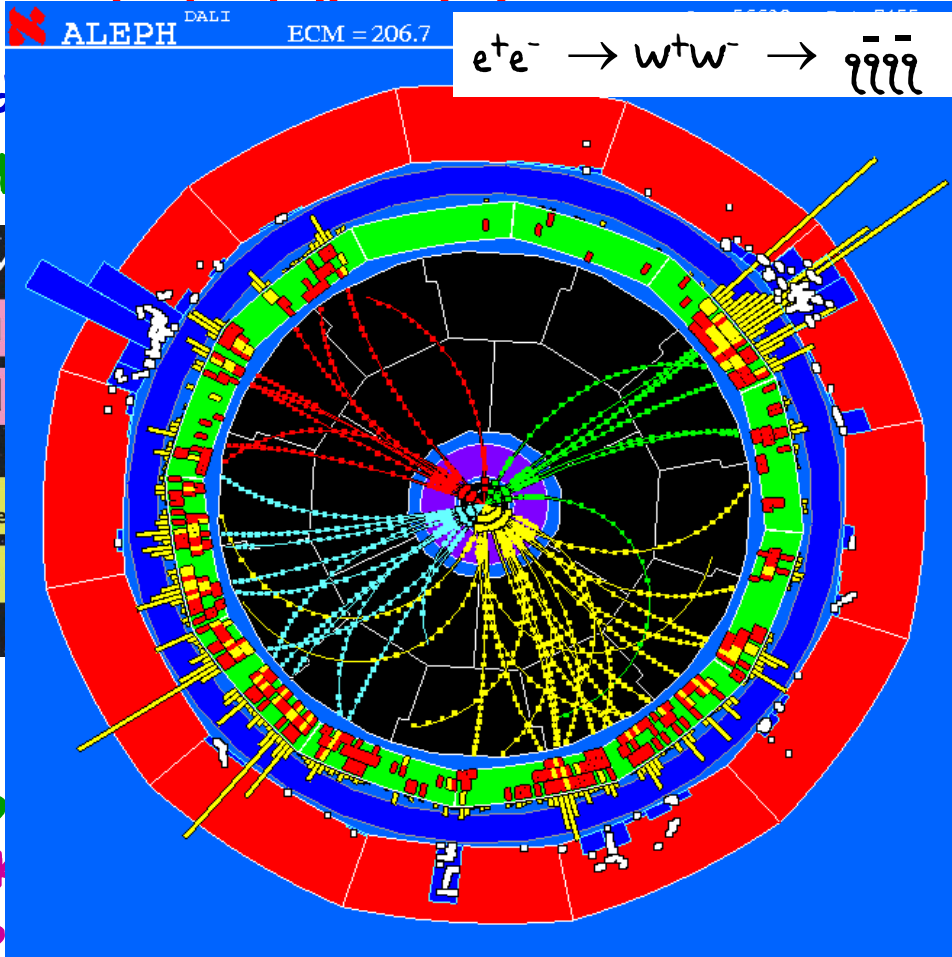
- ◆ Not quite ...
- Only e, γ, ...
- Quark
lepton

/GeV, quasi stable
τ, w, Z decay to
other particles

Basics of Particle-Flow Event Reconstruction (2)

□ What particles

- ◆ Ideally : build
- By definit



SM) particles
(+ dark matter)

◆ Not quite ...

- Only $e, \gamma, \mu, \tau, \nu_e, \nu_\mu, \nu_\tau$
- Quark
lepton

/GeV, quasi stable
 τ, w, Z decay to
ng

◆ Therefore : Follow electrons, muons, photons, stable hadrons, neutrinos and

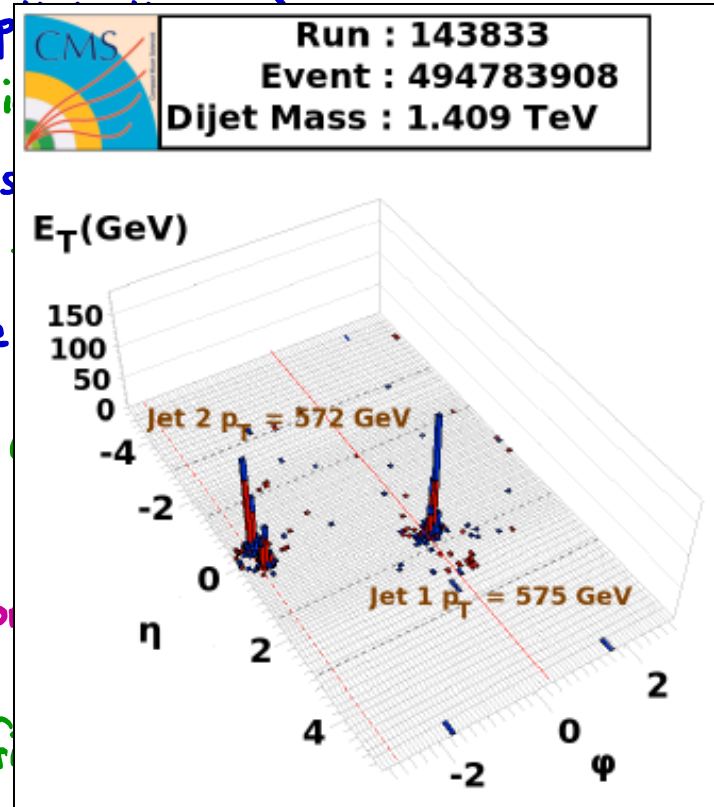
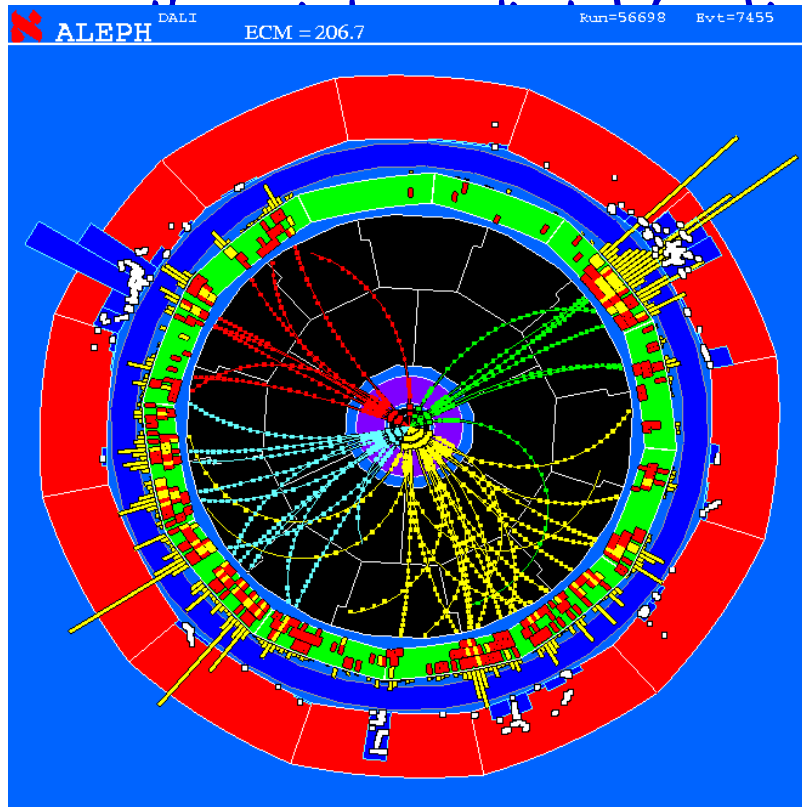
- And identify/reconstruct them individually, to go back to SM particles

Basics of Particle-Flow Event Reconstruction (3)

- why bother with individual particle reconstruction ?
 - ◆ It's seemingly complicated (combination, optimization, ...)
 - People usually like simple criteria for identification, reconstruction, ...
 - ◆ It requires an excellent knowledge of all sub-detectors of an experiment
 - People usually have the knowledge of the sub-detector they built
 - ◆ It requires to think ahead of time of the detector design, and of the interplay between sub-detectors
 - People usually optimize the design of "their" sub-detector, with at best one primary goal in mind
 - e.g., b tagging for vertex detector, tracks for tracker, isolated photons and isolated electrons for ECAL, jets for HCAL, ...
 - But with multiple technical/practical/financial constraints
 - which often take the precedence on the primary goal

Basics of Particle-Flow Event Reconstruction (3)

- why bother with individual particle reconstruction ?



→ which often take the precedence on the primary goal

- ◆ There are easier/faster ways to reconstruct “physics objects” (e.g., jets...)
 - e.g., purely tracker-based, or purely calorimeter-based

Basics of Particle-Flow Event Reconstruction (4)

□ So, indeed ... why bother ?

◆ Philosophical answer

- A list of particles is the closest one can get from the actual collision
 - Giving a complete and fully consistent view of the event
 - Making reconstructed events very similar to generated events
 - Greatly simplifies the analysis design process, for any final state

◆ Practical answer

- Each sub-detector response depends on the particle type. After identification of a particle, their combination

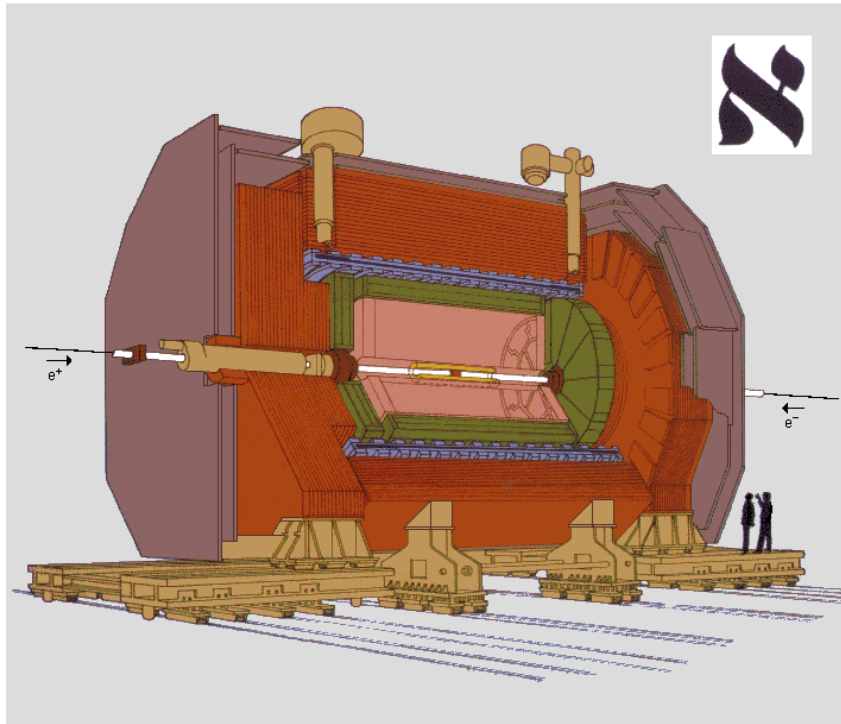
- Returns the best energy, direction, (mass) determination for each type
- Gives in turn the optimal response for jets, photons and leptons
- Is expected to improve the performance of any data analysis

◆ Financial answer : detector are expensive, thus make optimal use of them !

◆ Question : is the improvement with respect to easier approaches worth it ?

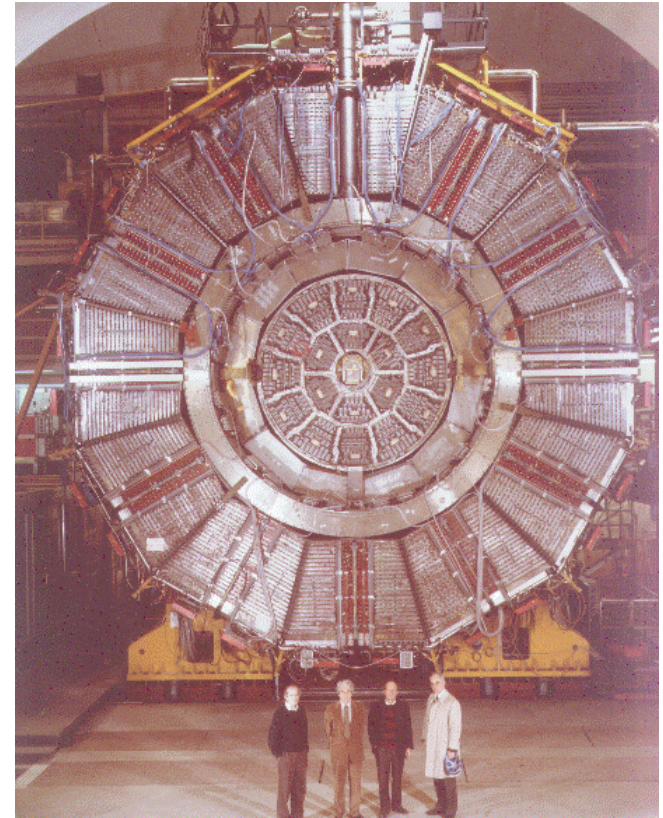
A detector thought for Particle Flow : ALEPH

- ALEPH is not very different from the standard HEP detector



The ALEPH Detector

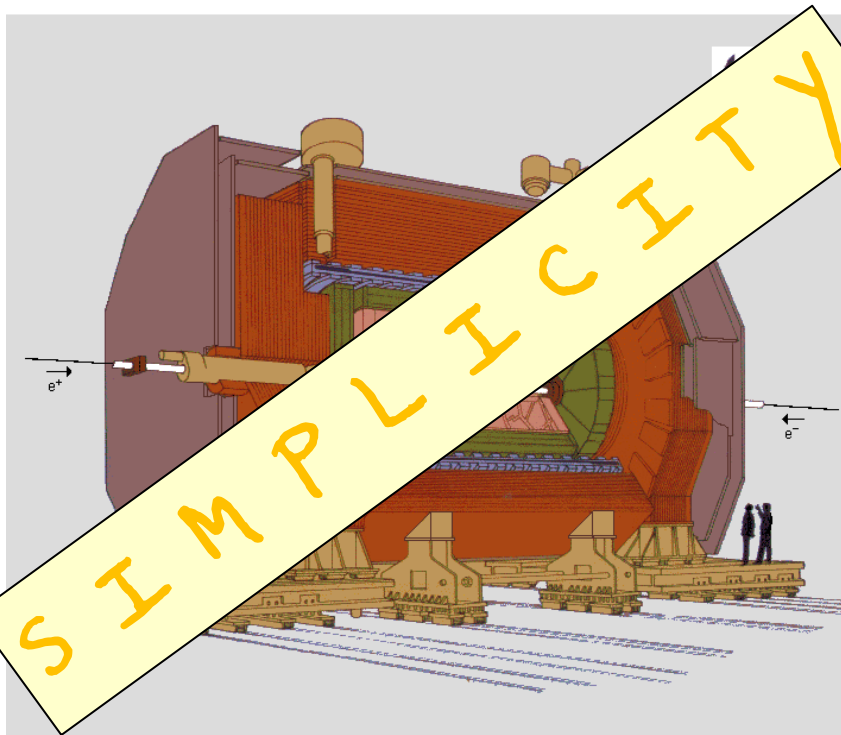
- Vertex Detector
- Inner Tracking Chamber
- Time Projection Chamber
- Electromagnetic Calorimeter
- Superconducting Magnet Coil
- Hadron Calorimeter
- Muon Chambers
- Luminosity Monitors



- ◆ Tracking, hermetic EM and HAD calorimetry, muon chambers, large axial B field
 - So, what is so special about the design of this detector ?

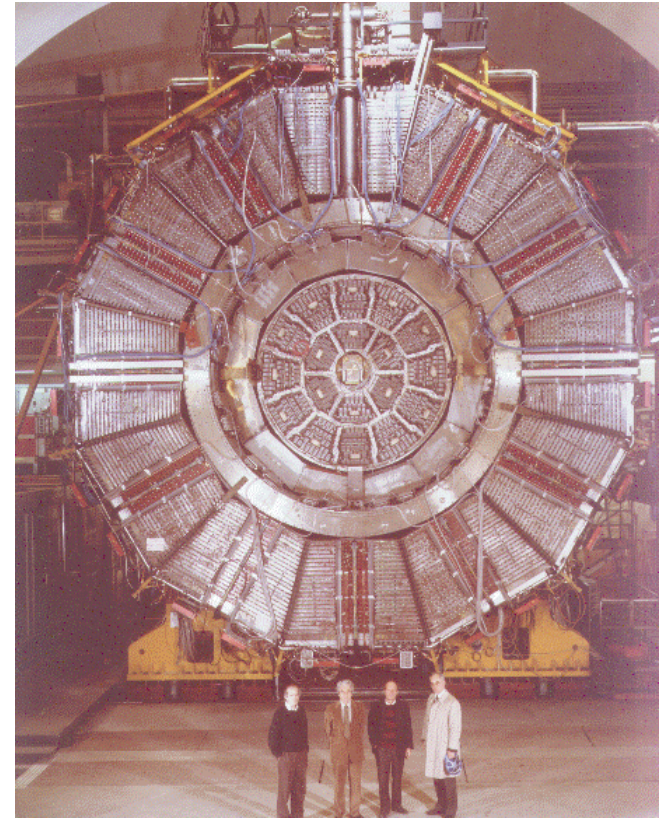
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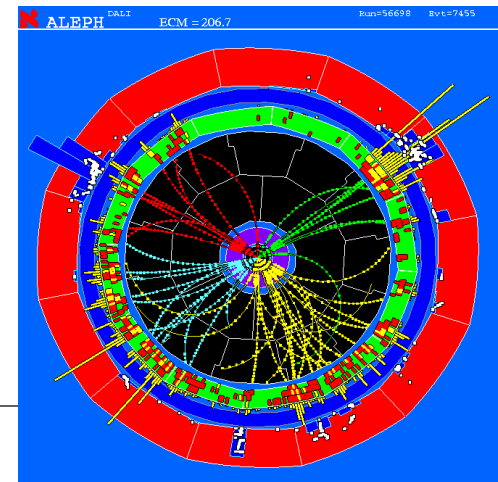
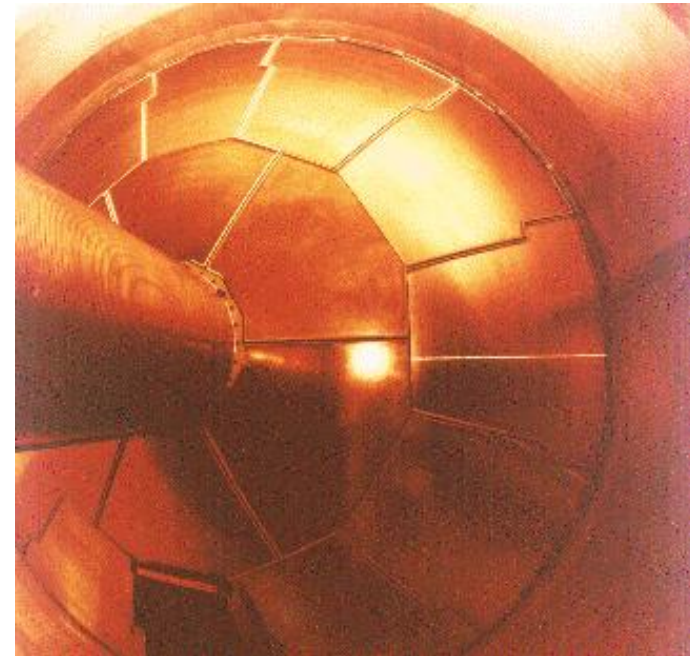
- ◆ Tracking, hermetic EM and HAD calorimetry, muon chambers, large axial B field
 - So, what is so special about the design of this detector ?

A tracker thought for Particle Flow (1)

□ ALEPH choice : a Time Projection Chamber

- ◆ Large volume mostly empty (filled with gas)
 - 1% X_0 , non destructive
- ◆ 21 three-dimensional measurements up to $R=1.80\text{m}$
 - No track-to-track ambiguity
 - 100% tracking efficiency, even in jets
 - Even if originating far from main IP
 - $\sigma(1/p_T) = 6.10^{-4} \text{ GeV}^{-1}$, no charge flip
 - Almost perfect in the LEP p_T range
 - No fake track, hence no fake energy reconstructed
 - All hits are displayed in this figure !
- ◆ Event charged-particle energy reconstructed perfectly
 - Origin, energy, direction, charge, but yet no ID.
 - (some e/π separation below 5 GeV, dE/dx)

B Field 1.5T



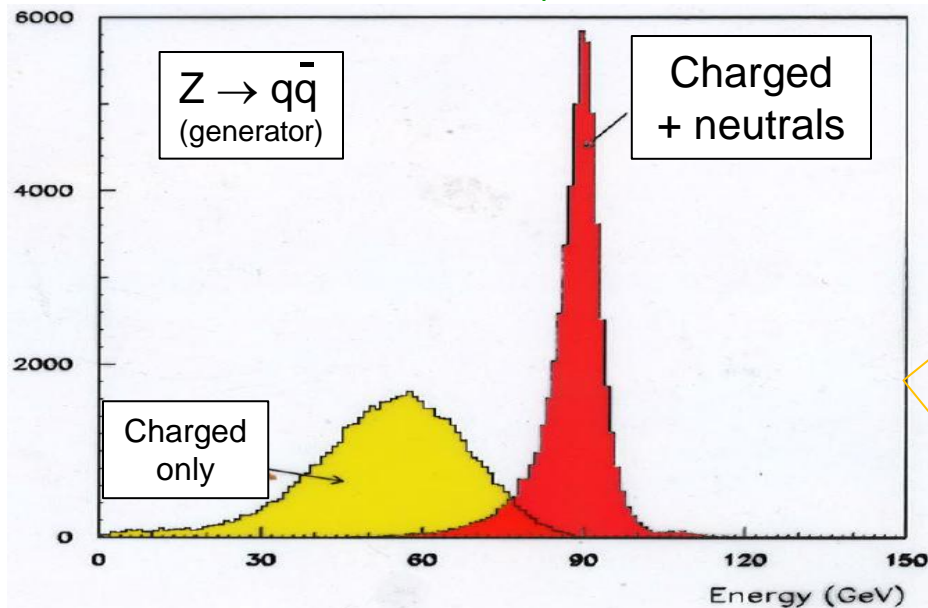
A tracker thought for Particle Flow (2)

□ Is that perfection enough ?

◆ Not really ...

(π^\pm , K^\pm , p , \bar{p} , e^\pm , μ^\pm)

- In a hadronic Z decay, charged particles carry 65% of the total energy
→ with a broad distribution (from 0 to 100%)
- Back to almost 100% if neutral particles are identified perfectly



A proof of principle
for particle flow reco

→ Calorimeter resolution and acceptance taken into account

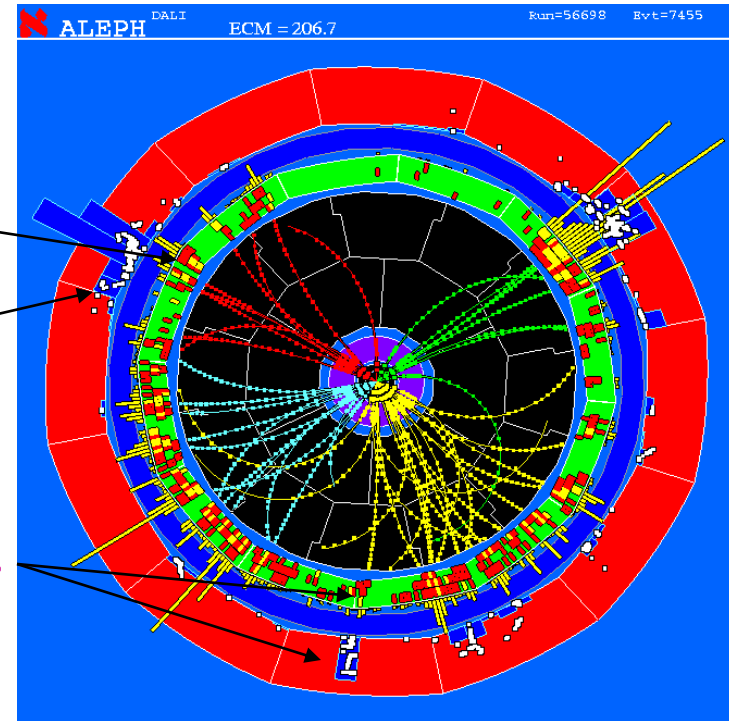
A calorimetry thought for Particle Flow ? (1)

□ Remember our famous $e^+e^- \rightarrow WW$ event

- ◆ Neutrals as energy deposits in the calorimeters
 - 25% are photons (mostly from π^0 decay)
 - Detected in ECAL
 - 10% are neutral hadrons (K_L^0 , n , \bar{n} , ...)
 - Detected in ECAL and HCAL
 - But 65% are charged particles
 - Electrons in ECAL, Muons in HCAL
 - Charged hadrons in ECAL and HCAL

□ Weapons to identify what is what

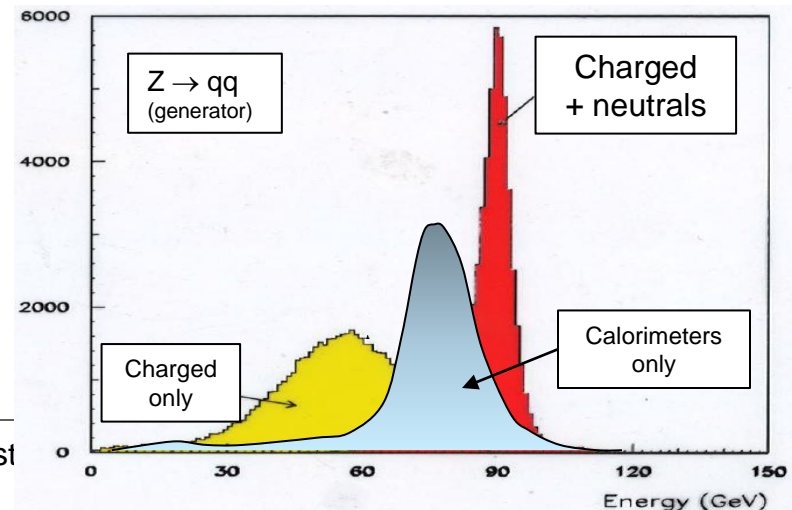
- ◆ Large B field + large tracking volume
 - separate charged from neutrals, and charged from charged
- ◆ Calorimetry fine (3D) segmentation + small Moliere radius ("isolate" all deposits)
- ◆ No/little material in front of calorimetry (one particle = one deposit)



A calorimetry thought for Particle Flow ? (2)

□ Why not using calorimeters only ?

- ◆ worse energy resolution for charged hadrons ($100\%/\sqrt{E}$ vs perfect) and worse direction determination for charged hadrons (magnetic field)
 - Affects 65% of the event energy in hadronic final states
- ◆ Muons are just minimum ionizing particles
 - Energy information is lost, so need anyway to use tracking
- ◆ ECAL response for charged/neutral hadrons is smaller than for photons/electrons
 - A large fraction of the energy is not underestimated
- ◆ Low momentum charged particles do not reach calorimeters
 - Their energy is lost
- ◆ No particle list
- ◆ For Z hadronic decays
 - Reconstruct 72 ± 10 GeV
 - Need a posteriori corrections



An EM calorimetry thought for Particle Flow (1)

□ The ECAL, for electrons and photons

◆ 36 modules (12 in barrel and 12/end-cap)

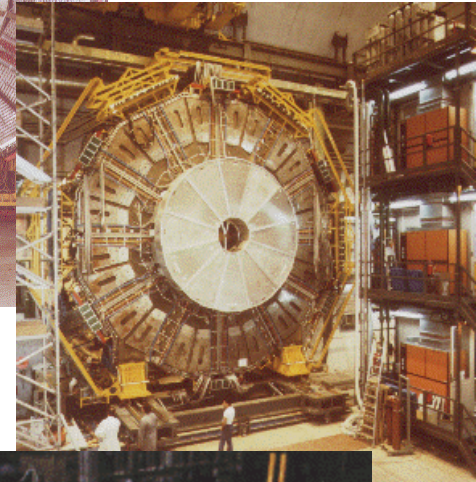
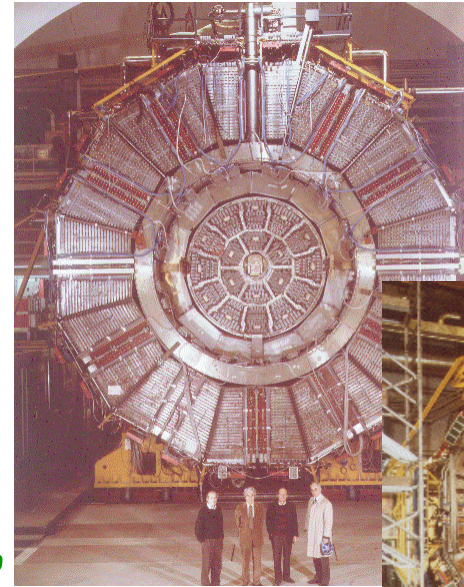
◆ For each module:

$$R_M = 1.5 \text{ cm}$$

- 45 planes of lead / wire chamber
- Total thickness $22 X_0$
- Transverse segmentation of $3 \times 3 \text{ cm}$
- Longitudinal segmentation $4/9/9 X_0$
- 75000 x 3 cathode readout "towers"
- Each of the 45 wire planes is readout too

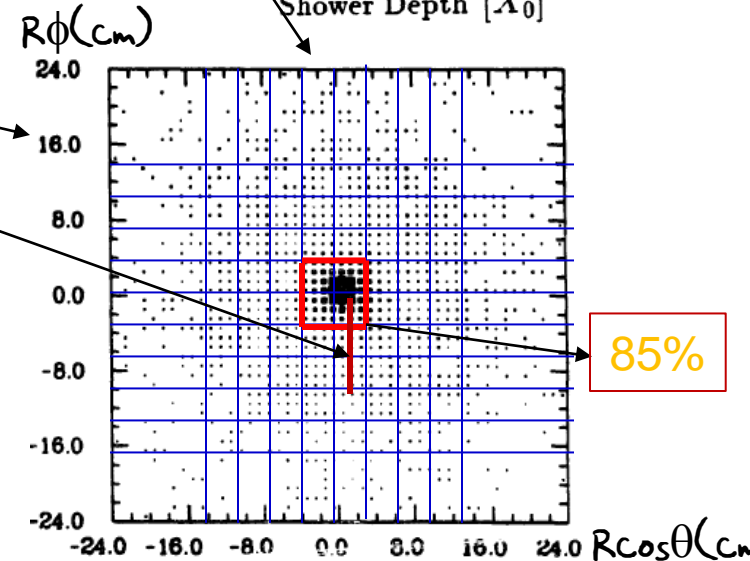
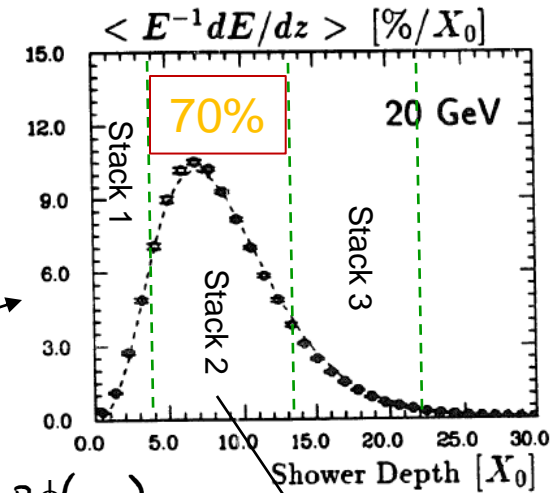
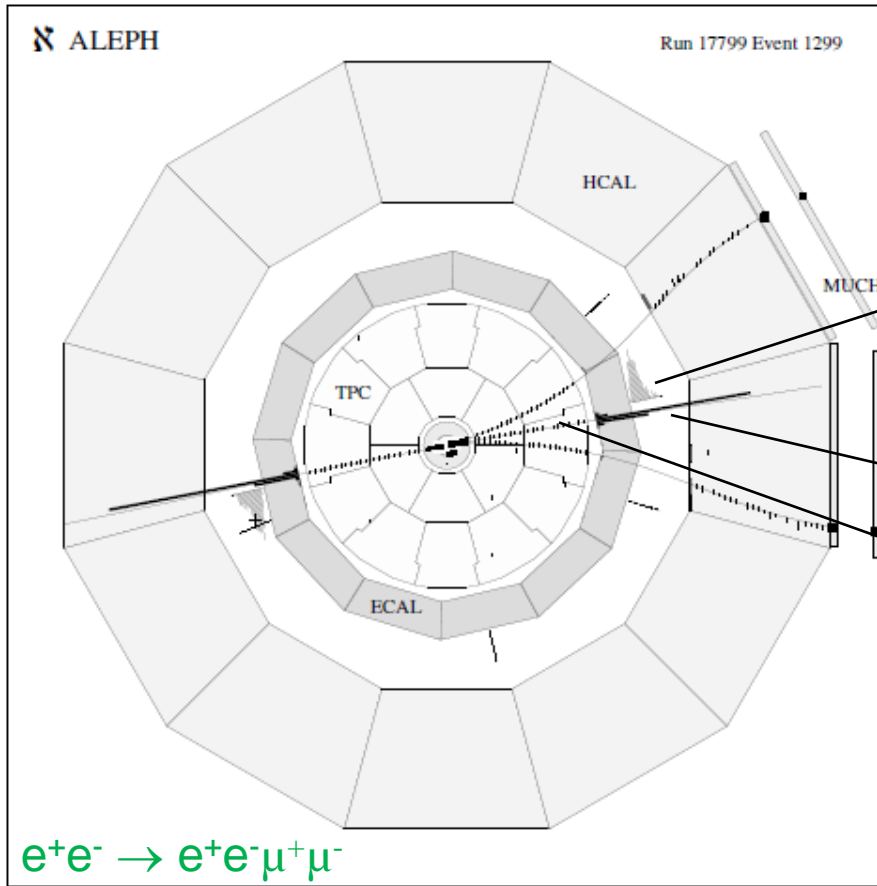
◆ Main characteristics

- Hermeticity
- 3D fine granularity
- Redundancy
- Simplicity



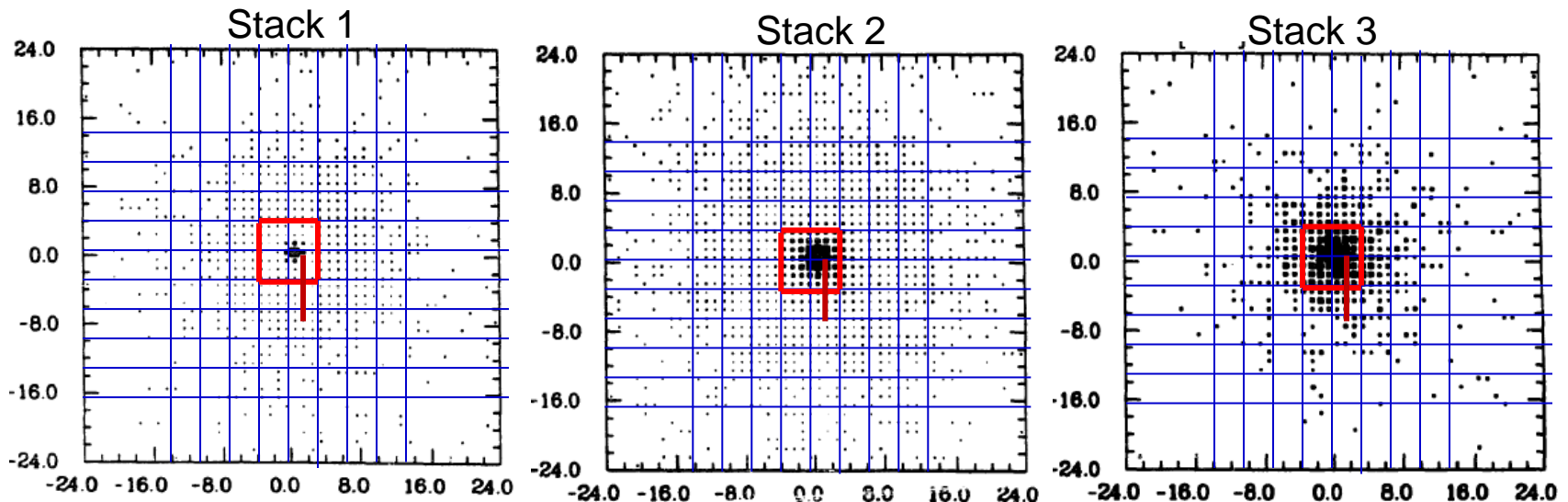
An EM calorimetry thought for Particle Flow (2)

- How do electromagnetic showers look like in this ECAL ?



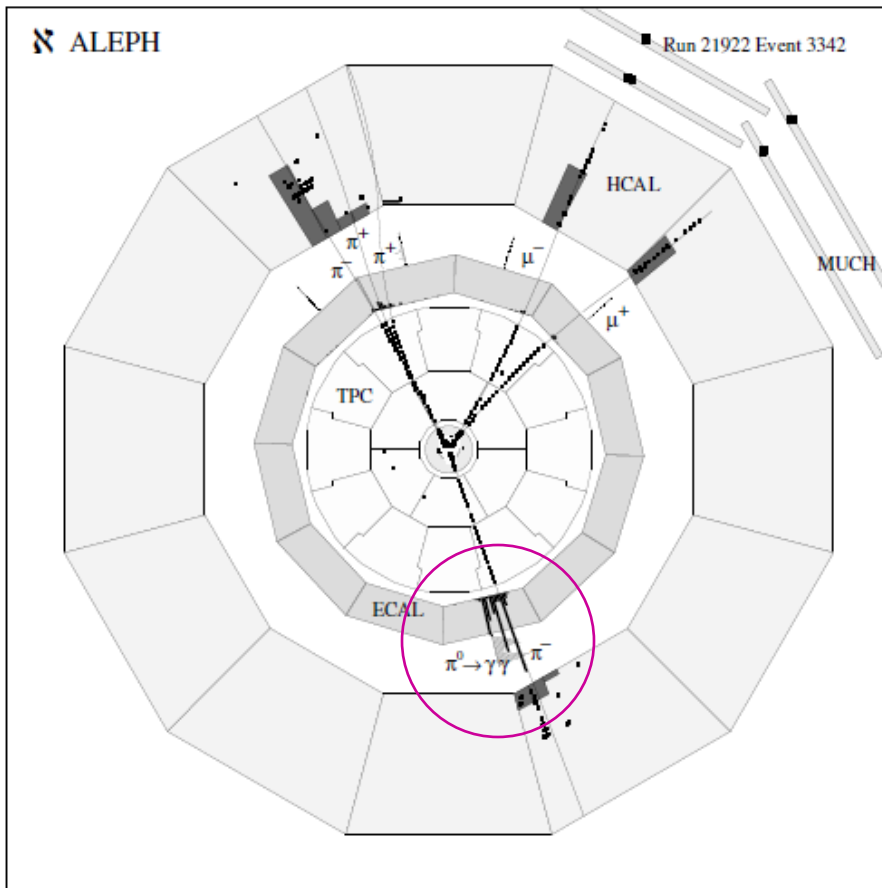
An EM calorimetry thought for Particle Flow (3)

- EM energy deposits (γ , e^\pm) are reconstructed as follows
 - ◆ Seed with a cell in stack 2 with $E_2 > 30$ MeV, and larger than its 8 neighbours
 - ◆ Correlate with a cell in stack 1 and/or in stack 3 with $E_{1,3} > 15$ MeV
 - ◆ Estimate the energy from the four central towers E_4 (and divide by 0.85)
 - ◆ Call it photon if the closest track extrapolation is > 2 cm away from barycentre
 - ◆ Otherwise, test E_4/p , $\Sigma E_{x_i}/E_4$, dE/dx against the electron hypothesis
 - and call it either electron or charged hadron



An EM calorimetry thought for Particle Flow (4)

- Efficiency for photons ($E > 250 \text{ MeV}$) in jets and tau decays $> 95\%$
 - ◆ And 100% in isolation



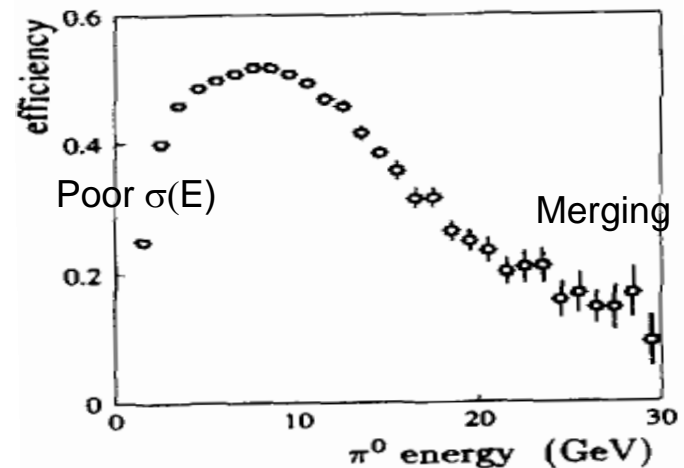
Example : $e^+e^- \rightarrow \tau^+\tau^-\mu^+\mu^-$

$\tau^- \rightarrow \rho^- \nu_\tau \rightarrow \pi^- \pi^0 \nu_\tau \rightarrow \pi^- \gamma \gamma \nu_\tau$

$\tau^+ \rightarrow \pi^+ \pi^+ \pi^- \nu_\tau$

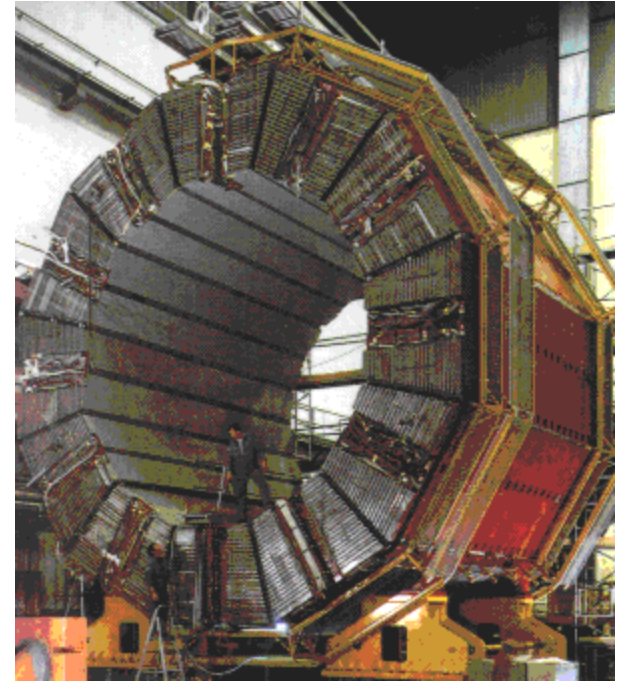
ECAL granularity performance:

π^0 's can even be resolved in jets
(mostly irrelevant for PF purposes)



A HAD calorimetry thought for Particle Flow ? (1)

- The HCAL for charged/neutral hadrons, ... and muons
 - ◆ 36 modules (24 in barrel, 6/end-cap)
 - ◆ Placed behind the magnet coil
 - Hadrons lose some energy in the coil
 - ◆ For each module
 - 23 planes of 5cm iron / drift streamer tubes
 - Total thickness $7.2 \lambda_0$'s
 - Cathode readout, transverse size 15cm x 15cm
 - But no longitudinal segmentation
 - Digital individual drift tube readout (Yes/No)
 - But only 2D granularity
 - ◆ Seems not particularly optimal for hadron/hadron separation
 - We'll see that in a couple slides
 - ◆ What about muons ?



A HAD calorimetry thought for Particle Flow ? (2)

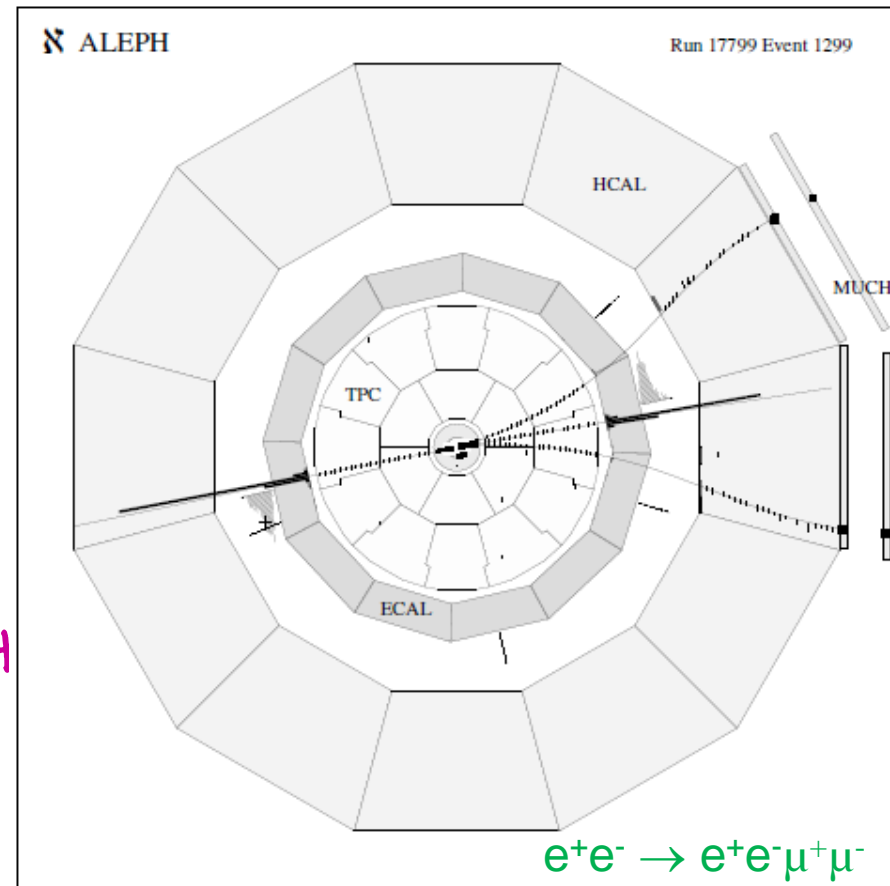
□ Muon tracking in HCAL

◆ Remember our $e^+e^- \mu^+\mu^-$ event

- Drift tubes used to track muons
 - All the way to mu-chambers
- “Energy” measured : 3-5 GeV
 - About 300 MeV/layer (mip)
 - Easily link-able to the track
- Unambiguous identification
 - width, penetration depth, μ CH
 - 95% identification efficiency
- Almost no difference in jets
 - 0.5% fake rate (decays)

□ Summary so far : γ 's, e 's, μ 's (tracks + calo energies) well identified

- ◆ Thanks to 3D (2D) segmentation of tracker, ECAL, HCAL (+ B field)

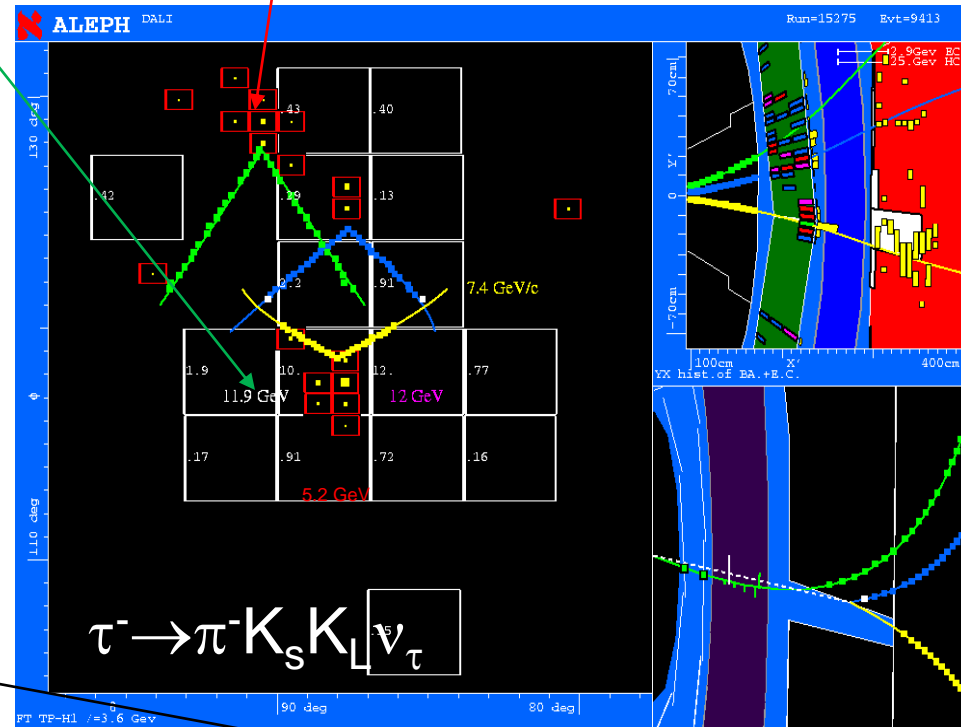


A HAD calorimetry thought for Particle Flow ? (3)

- We are left with charged and neutral hadrons
 - ◆ i.e., with several sets of tracks remaining after e and μ identification
 - Possibly linked to one HCAL cluster and ECAL cluster(s)
 - Unlinked HCAL clusters give rise to a neutral hadron

□ In this particular event:

- ◆ Three charged particles
- ◆ Each linked to an ECAL cluster
- ◆ Two linked to one HCAL cluster
- ◆ All identified as charged hadron
 - with $E = p_{Track}$
- ◆ But clear excess of Calo energy
 - 5, 25 GeV in ECAL, HCAL
 - 7.4 and 4.8 GeV in Tracker
- ◆ Signs the presence of a neutral hadron



→ Note the $K_S \rightarrow \pi^+ \pi^-$ decay

A HAD calorimetry thought for Particle Flow ? (4)

- Determination of the neutral hadron energy
 - ◆ HCAL cluster without a track pointing to it
 - $E_{\text{Neutral Hadron}} = E_{\text{HCAL cluster}}$
 - Keep if $E_{\text{Neutral Hadron}} > 500 \text{ MeV}$
 - ◆ HCAL cluster with one or more tracks pointing to it
 - $E_{\text{Neutral Hadron}} = E_{\text{HCAL cluster}} + c_{e/\pi} \sum E_{\text{ECAL cluster}} - \sum p_{\text{Track}}$
 - with $c_{e/\pi} = 1.30$ (ratio of electron to pion response in ECAL)
 - Keep if $E_{\text{Neutral Hadron}} > 100\% \sqrt{\sum p_{\text{Track}}}$
 - i.e., if larger than 1σ of the expected hadron shower fluctuations
- The HCAL of ALEPH could have been more carefully designed
 - ◆ only 50% efficiency (and 50% purity) for neutral hadrons
 - Better granularity, better material, would have helped (more expensive)
 - Could the 2D granularity of the streamer tubes have been exploited ?
 - The HCAL could have been put inside the magnet (more expensive)

Does it work as initially designed for ? (1)

Remember the ideal projection

In hadronic Z decays, the energy is shared

- 65% are charged particles.

→ with perfect energy resolution

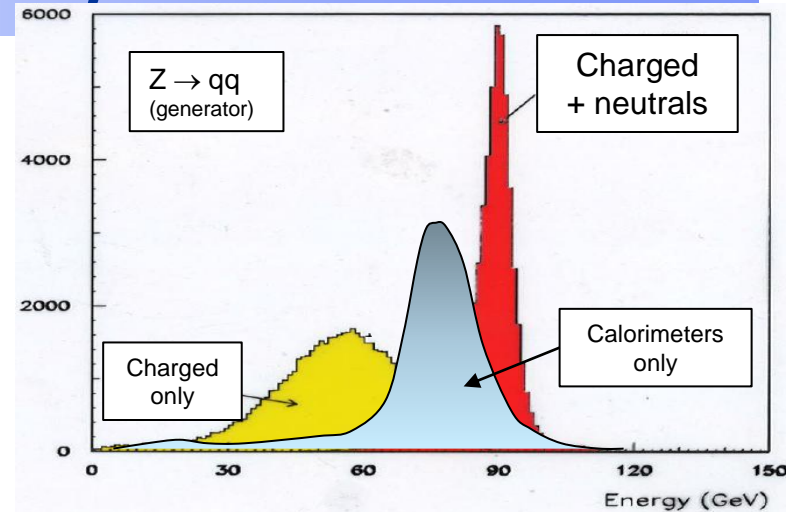
- 25% are photons. For single photons:

→ $\sigma(E_\gamma) = 20\% \sqrt{E_\gamma}$

- 10% are neutral hadrons. For single hadrons

→ $\sigma(E_{\text{had}}) = 100\% \sqrt{E_{\text{had}}}$

The total energy resolution, if all particles were ideally identified, would be



$$\sigma(E_{\text{tot}}) = \left[\sum_{\text{tracks}} \sigma^2(p_{\text{track}}) + \sum_{\text{photons}} (0.20\sqrt{E_\gamma})^2 + \sum_{\text{neutral hadrons}} (1.00\sqrt{E_{\text{had}}})^2 \right]^{1/2}$$

$$= [0 + 0.04 \times 0.25 E_{\text{tot}} + 1.00 \times 0.10 E_{\text{tot}}]^{1/2}$$

$$= [0 + 0.01 + 0.10]^{1/2} \sqrt{E_{\text{tot}}} = 33\% \sqrt{E_{\text{tot}}} = 3.1 \text{ GeV} \quad (\text{For } E_{\text{tot}} = m_Z)$$

Charged hadrons

Photons

Neutral hadrons

Does it work as initially designed for ? (2)

□ Not quite for the energy resolution ...

◆ Measured resolution twice as large:

- $E_{\text{tot}} = 90.5 \pm 6.2 \text{ GeV}$

◆ But remember with calorimeters only:

- $E_{\text{tot}} = 72 \pm 13 \text{ GeV}$

□ Reasons for the difference ?

◆ Neutral hadron energy losses

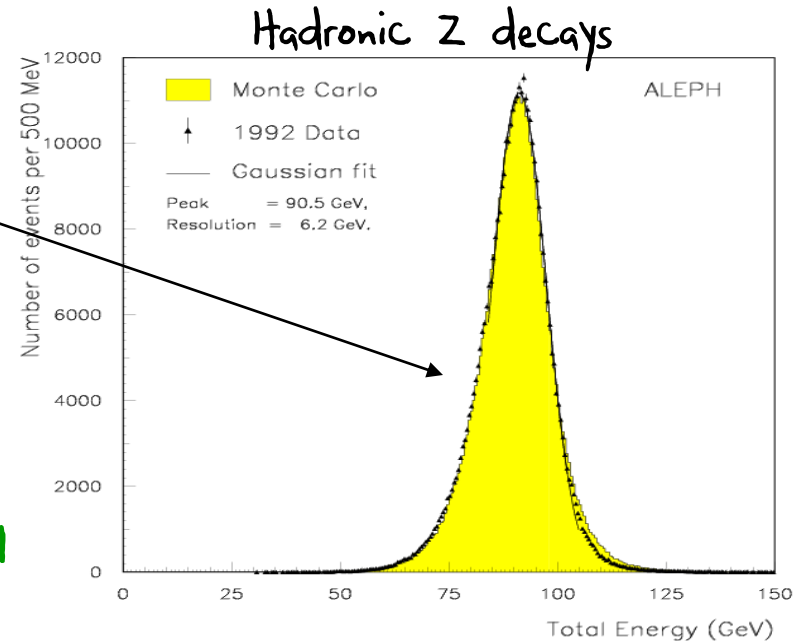
- e.g., neutral hadron shower in the coil
- e.g., unresolved neutral hadron shower

◆ Neutral hadron energy double counting

- e.g., charged hadron shower upwards fluctuation, interaction in the coil
→ giving rise to a fake neutral hadron

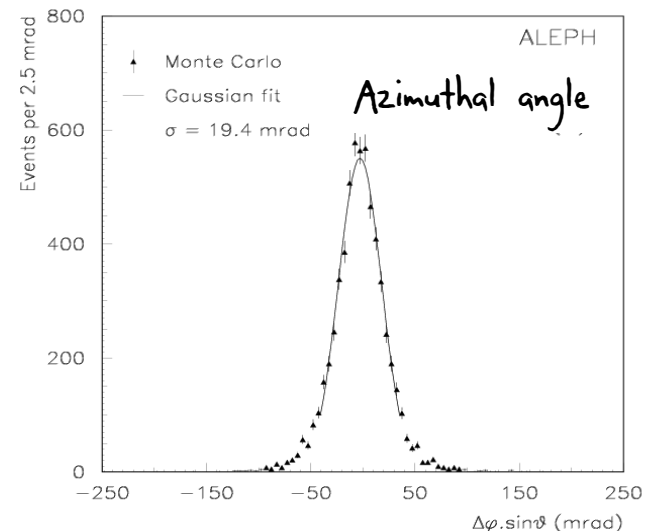
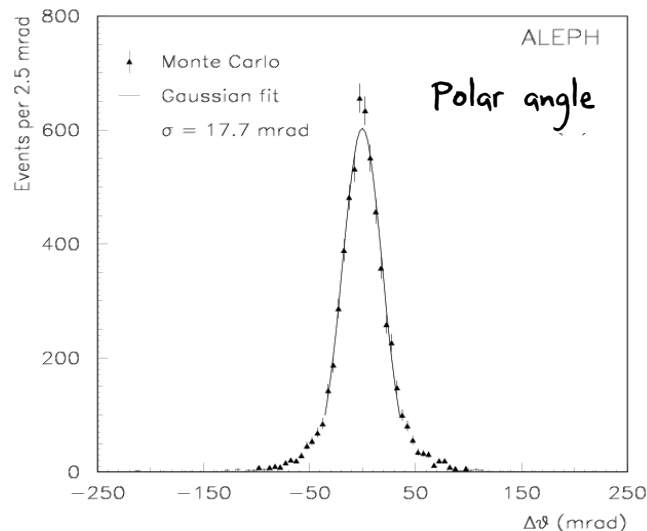
□ Could have been alleviated with a better HCAL design

- ◆ And maybe with an optimal use of streamer tubes (was tried, not hard enough)?



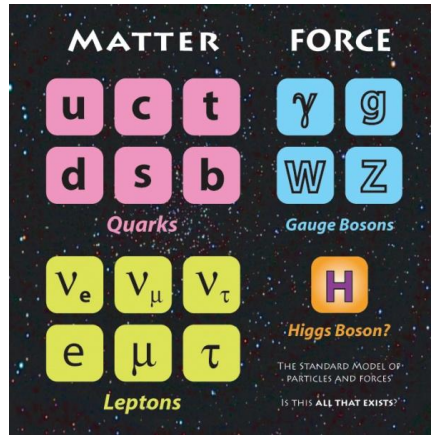
Does it work as initially designed for ? (3)

- But very much so for the quark angular resolution
 - ◆ Performance on hadronic Z decays (cont'd)
 - Jet angular resolutions (with respect to parton) of about 18 mrad
 - would have been 13 mrad with an ideal detector
 - was 64 mrad with tracks only, and 100 mrad with calos only
 - Made a breakthrough in b-tagging efficiency, in particular
 - By a better determination of track impact parameter wrt the jet axis



Does it work as initially designed for ? (4)

- And what about the very initial and very ambitious aim ?
 - ◆ i.e., identify all standard model particles and possibly dark matter particle ?



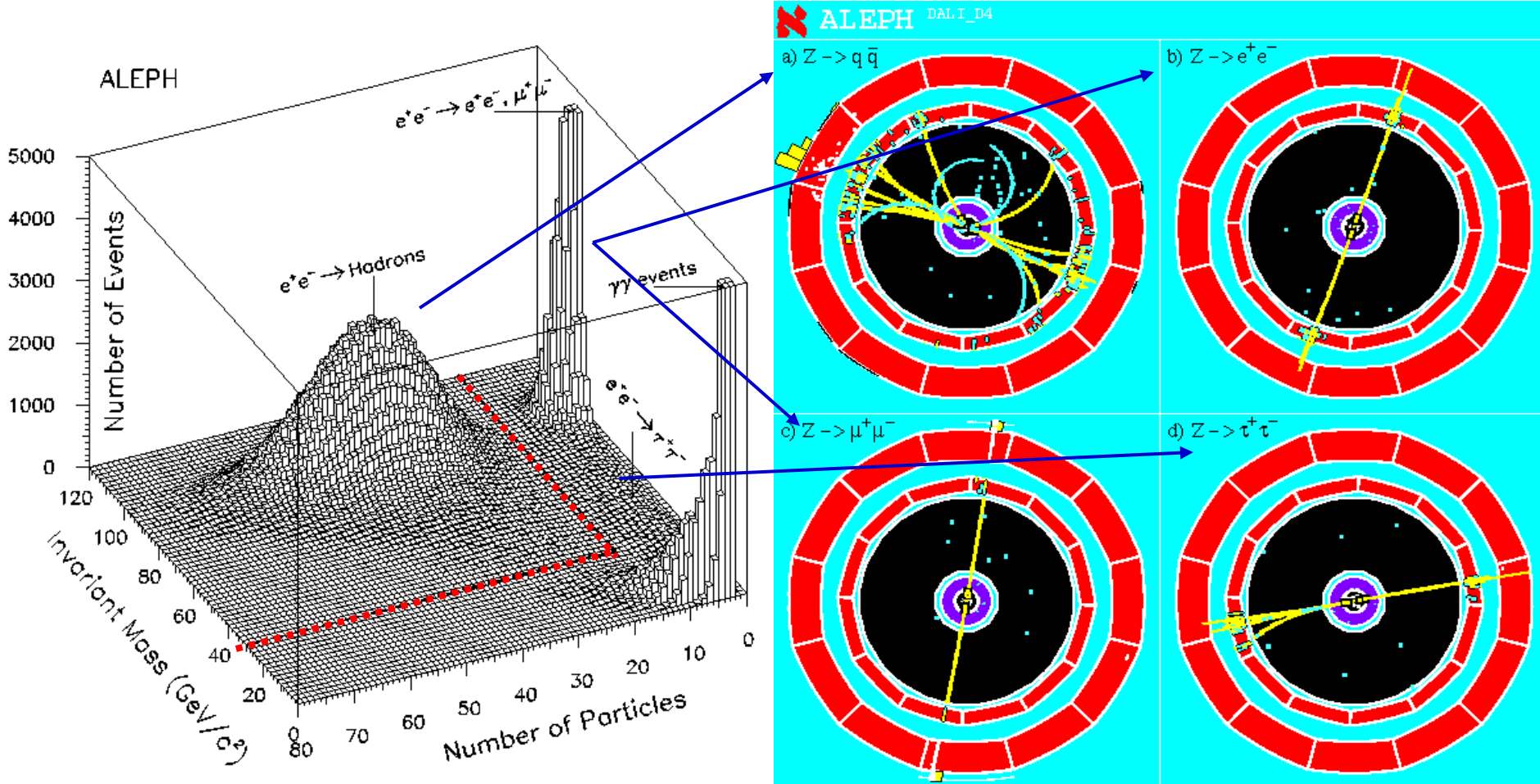
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- ◆ We already went through a lot here
 - Electrons
 - Muons
 - Photons
 - Taus
 - Quark and gluons (jets)

Does it work as initially designed for ? (5)

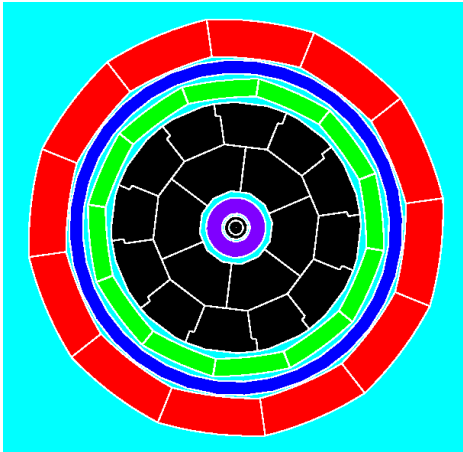
- Identification of Z (or W, H) through decay products
 - ◆ $M > 30 \text{ GeV}/c^2$, $N > 15$ particles - 99.6% efficient for hadronic Z decays



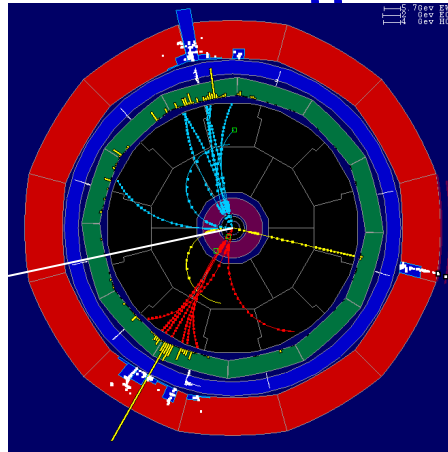
Does it work as initially designed for ? (b)

□ And what about neutrinos ?

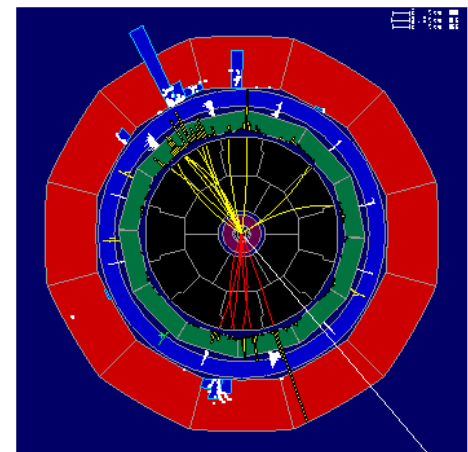
◆ $Z \rightarrow \nu\bar{\nu}$



$W^+W^- \rightarrow q\bar{q}\mu\nu$



$HZ \rightarrow b\bar{b}\nu\bar{\nu}$



◆ First event not detectable (no trigger), but neutrinos in the other two are !

• From energy-momentum conservation: $\vec{p}_\nu = - \sum_{\text{particles}} \vec{p}_i$ and $E_\nu = \sqrt{s} - \sum_{\text{particles}} E_i$

→ Hence the importance of measuring well the energy/direction of all particles in the event. Particle-Flow reconstruction is the tool for ν 's

→ The same statement holds true for

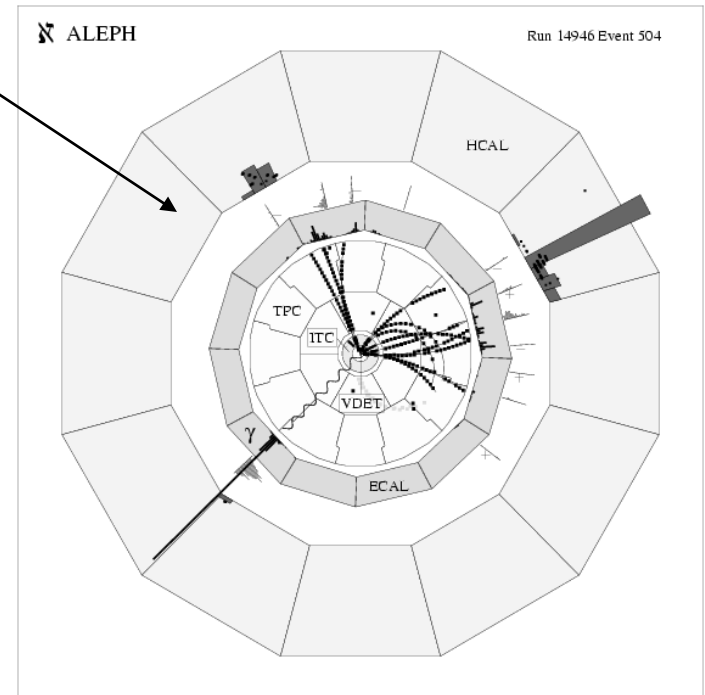


Does it work as initially designed for ? (1)

- Are neutrinos (or any kind of missing energy) well measured ?
 - ◆ Difficult to know, as neutrinos are not directly detected, but ...
 - Can use $e^+e^- \rightarrow q\bar{q}\gamma$ events, and fake a neutrino by “removing” the photon
 - And the energy of the neutrino is known in that case !

□ Exercise

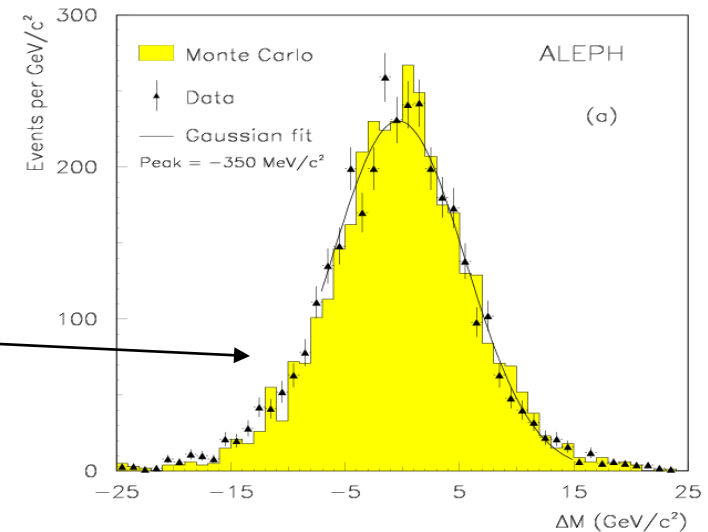
- ◆ Establish that
$$E_\gamma^{\text{recoil}} = \frac{s - m_{\text{had}}^2}{2\sqrt{s}}$$
- ◆ One can now compare
 - $E_\gamma(\text{measured})$ and $E_\gamma(\text{recoil})$
- ◆ Or equivalently
 - $m_{\text{had}}(\text{measured})$ and $m_{\text{had}}(\text{recoil})$
 - The latter determined from E_γ
- ◆ And establish the detector calibration for ν



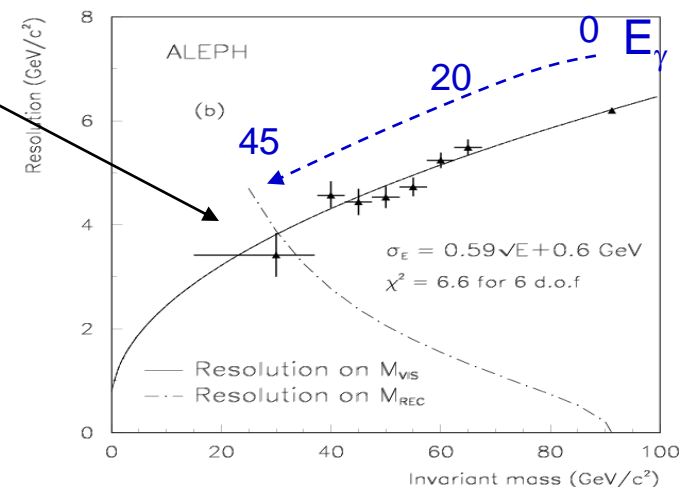
Does it work as initially designed for ? (8)

Particle-Flow calibration with $q\bar{q}\gamma$ events

- ◆ Determine the hadronic mass
 - with particle flow: m_{had} (measured)
 - As recoiling to the photon: m_{had} (recoil)
- ◆ Plot the difference ΔM
 - Compatible with 0, establishes calibration



- ◆ Determine the mass resolution
 - As a function of E_γ (i.e., of m_{had})
 - $E_\gamma = 0$ means $Z \rightarrow q\bar{q}$
 - Result
 - $\sigma(m) = 59\% \sqrt{m}$ or $\sigma(E) = 59\% \sqrt{E}$
 - To be compared to 33% ideally



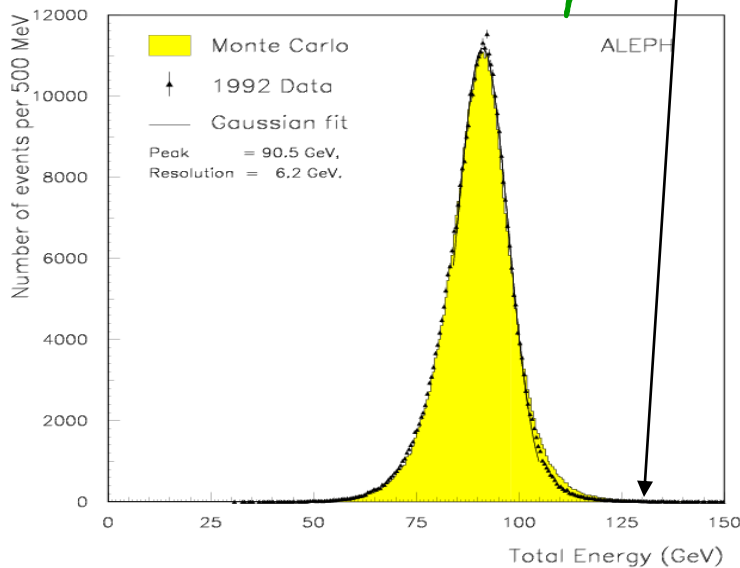
Is it ready for physics analysis ? (1)

- Core resolution vs tails : The devil is in the details
 - ◆ Particle-Flow reconstruction is a beautiful intellectual construction
 - With a superior core resolution and larger efficiencies for almost everything
 - Jet energy and direction, jet-jet invariant mass, missing energy, b tagging, tau selection, etc...
 - which a priori helps to separate different processes
 - ◆ But it might be subject to more reconstruction failures than simpler methods
 - Because it uses all sub-detectors of the experiments and needs a refined algorithm to put all pieces together
 - Again, the (apparent) simplicity of the detector design is essential here.
The simpler the detector, the simpler the algorithm ...
 - Pure tracking methods are not sensitive to calorimeter noise
 - Pure calorimetric methods are not sensitive to track-cluster link efficiency
 - ◆ Tails may therefore develop in various distributions (e.g., jet energy or direction)
 - which may become a showstopper when looking for rare events

Is it ready for physics analysis ? (2)

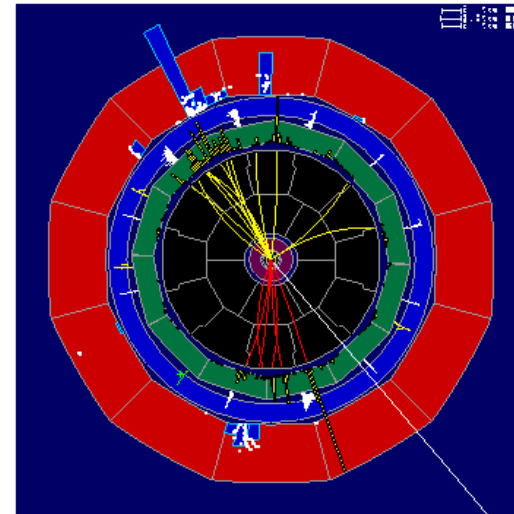
- Strategy, in the simulation and in the data (the best MC is the data)
 - ◆ 1) study events in the far tails, with a specific physics analysis
 - ◆ 2) fix the problem (if any - after all, in the data, it might be New Physics!)
 - ◆ 3) Go to 1) until no obvious problem remain
- More specifically, look here

Hadronic Z decays



... and here

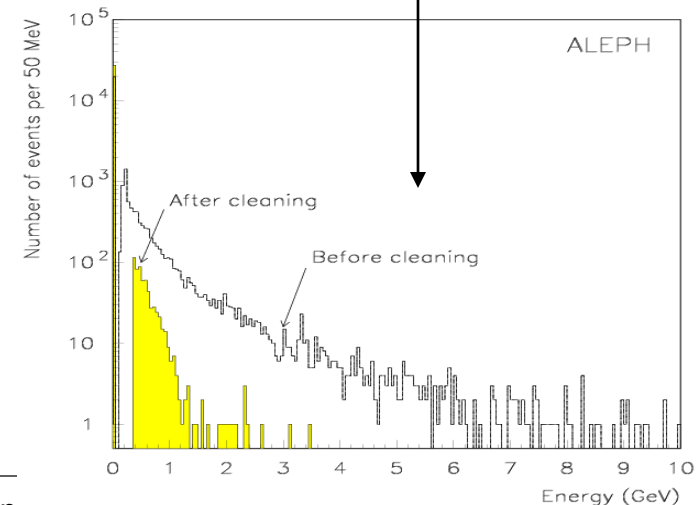
$HZ \rightarrow b\bar{b}v\bar{v}$



- ◆ Iterative process that improves the reliability, and ultimately the core resolution

Is it ready for physics analysis ? (3)

- Example : calorimeter noise (electronics, sparks, radioactive decays)
 - ◆ Typically creates high energy tails, fake particles in the list
 - Easily identified thanks to the redundancy of the calorimeter measurements
 - Pads vs wires in ECAL, towers vs drift tubes in HCAL
 - Noise cleaning requires compatibility between the measurements
 - Check effectiveness of the criterion on randomly triggered events and on Z hadronic decays (large tail must disappear)
 - Check over-cleaning in the $H\nu\bar{\nu}$ search (in MC, or in the first data)
- Repeat the exercise with $H\nu\bar{\nu}$ candidates
 - ◆ Track-cluster link failure
 - ◆ Particle identification failure
 - ◆ PF Algorithm bugs
 - ◆ Unforeseen configurations
 - with stubbornness and pragmatism



Is it ready for physics analysis ? (4)

- First analysis done in ALEPH with particle flow : $H\nu\bar{\nu}$ search
 - ◆ Two acoplanar jets, accompanied with missing energy (from the neutrinos)
 - LEP started in July 1989, delivered 25K Z in 1989
 - Used to develop and “tune” the algorithm
 - Particle-Flow algorithm written between January and April 1990
 - Algorithm and analysis first presented in May 1990, with 100K Z
 - Outperformed other analyses in ALEPH (tracks only) by a factor 2
 - Took the leadership in Higgs boson search until the end of LEP
- At that point, particle-flow reconstruction
 - ◆ Became the standard way of doing physics analysis in ALEPH
 - For precision measurements
 - For searches for Higgs boson and new particles
 - Almost 400 publications that use PF reconstruction

Lessons learned [?] (1)

- Particle Flow is potentially a very powerful event reconstruction
 - ◆ with possibly benefits in any physics analysis, if several conditions are met
- The detector better be designed with particle flow in mind
 - ◆ Overall simplicity of the design (over the full 4π solid angle)
 - Combination of all sub-detector information in PF algorithm
 - ◆ 3D granularity for all sub-detectors
 - ALL : Tracker, calorimeters, muon detectors, towards efficiency and purity
 - Good energy resolution is a bonus, but not crucial. Granularity comes first.
 - ◆ Large magnetic field (and large tracking volume)
 - To measure charged particle momenta and separate their energy deposits
 - ◆ Little material in front of the calorimeters
 - Must think of tracker, coil, services, ...
 - ◆ Redundancy of the measurements
 - To fight against fake, noise, ...

Lessons learned [?] (2)

- The algorithm development is also quite demanding
 - ◆ It requires an excellent knowledge of all sub-detectors
 - Towards the optimal use of the information
 - for particle identification, reconstruction, and for “cleaning”
 - ◆ It requires an excellent knowledge of particle/detector interaction theory
 - Towards the understanding and handling of “special” cases
 - ◆ It requires a specific developer profile
 - Stubbornness
 - There are no problems, there are only solutions
 - Pragmatism
 - Each special case is indeed special and requires specific treatment
 - ◆ It requires to be developed for/with a physics analysis, possibly with data
 - Towards performance improvement with superior motivation
 - And natural finding/solving of “special” cases

Lessons learned [?] (3)

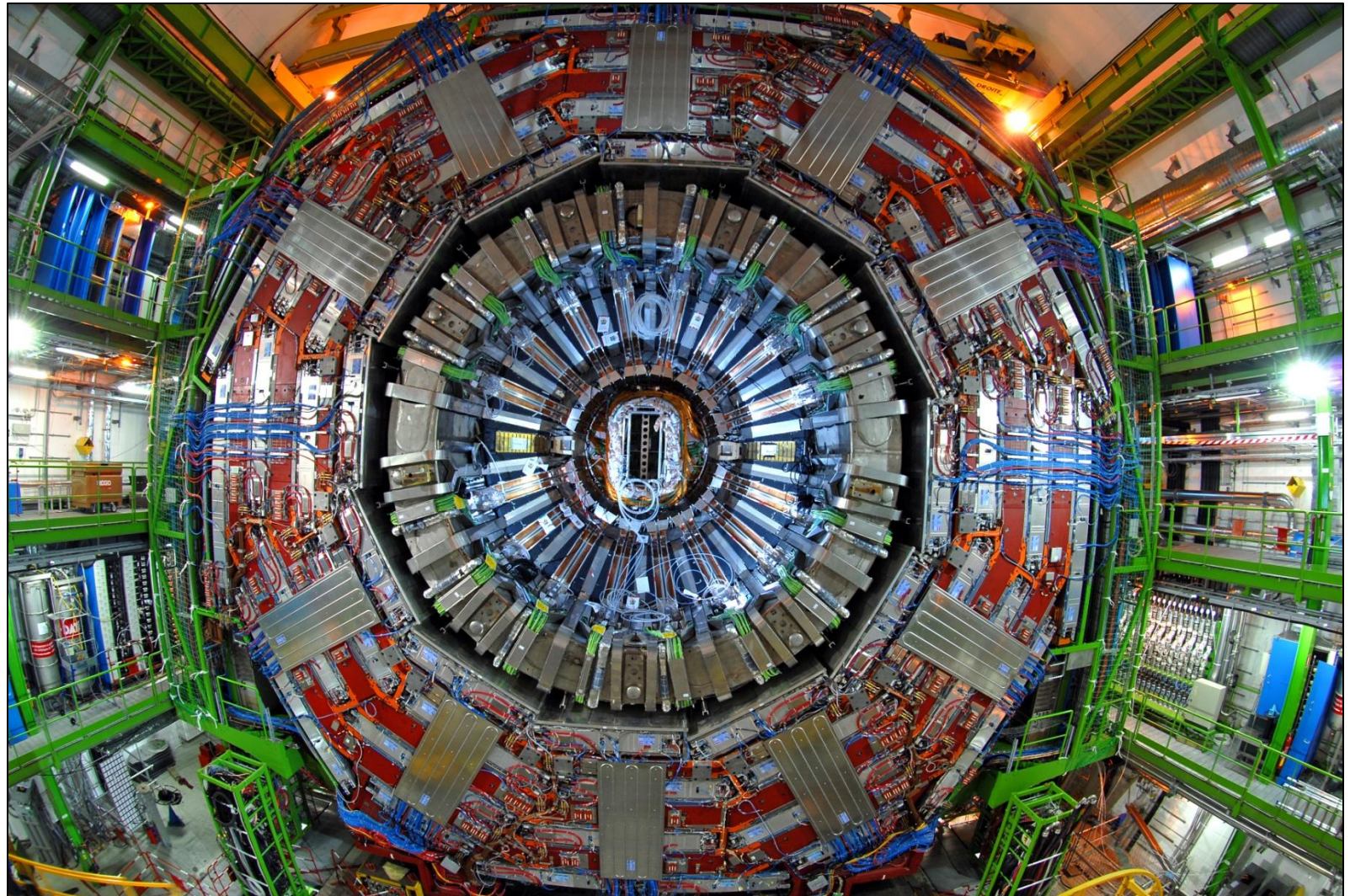
□ Exercise (1)

- ◆ Study the design of other experiments : DELPHI, OPAL, L3, and ATLAS
- ◆ Answer the following questions
 - Was/is a particle-flow event reconstruction developed ?
 - If yes, was/is its performance significantly better than simpler methods ?
 - Was/is it used in physics analyses ?
 - With the help of slides 35/36, find the reason(s) for the answers above.

□ Exercise (2) and hint for Exercise (1)

- ◆ Another method, now called "energy flow", combines tracking and calorimeters
 - It opens a road around each track (momentum p) extrapolation
 - It either masks the calorimeter energy measured in this road, replaced by p
 - Or it replaces the charged-hadron expected energy by p
- ◆ Did/do the above experiments use this method ? If yes, why ?
- ◆ Explain why, in ALEPH, the performance was poorer than PF ($\sigma_E > 10\text{GeV}$)

A detector not designed for Particle Flow: CMS (1)



A detector not designed for Particle Flow: CMS (2)

□ Original/official list of design requirements

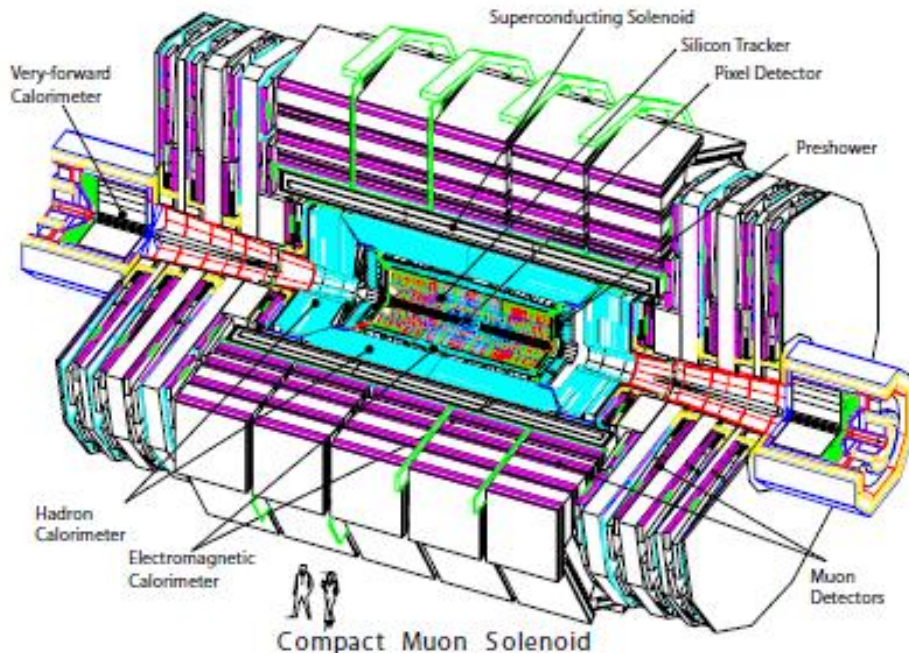
- Good muon identification and momentum resolution over a wide range of momenta and angles, good dimuon mass resolution ($\approx 1\%$ at 100 GeV), and the ability to determine unambiguously the charge of muons with $p < 1$ TeV;
- Good charged-particle momentum resolution and reconstruction efficiency in the inner tracker. Efficient triggering and offline tagging of τ 's and b -jets, requiring pixel detectors close to the interaction region;
- Good electromagnetic energy resolution, good diphoton and dielectron mass resolution ($\approx 1\%$ at 100 GeV), wide geometric coverage, π^0 rejection, and efficient photon and lepton isolation at high luminosities;
- Good missing-transverse-energy and dijet-mass resolution, requiring hadron calorimeters with a large hermetic geometric coverage and with fine lateral segmentation.

◆ Particle Flow reconstruction was not part of this list

- Although some of the requirements might be useful towards particle flow

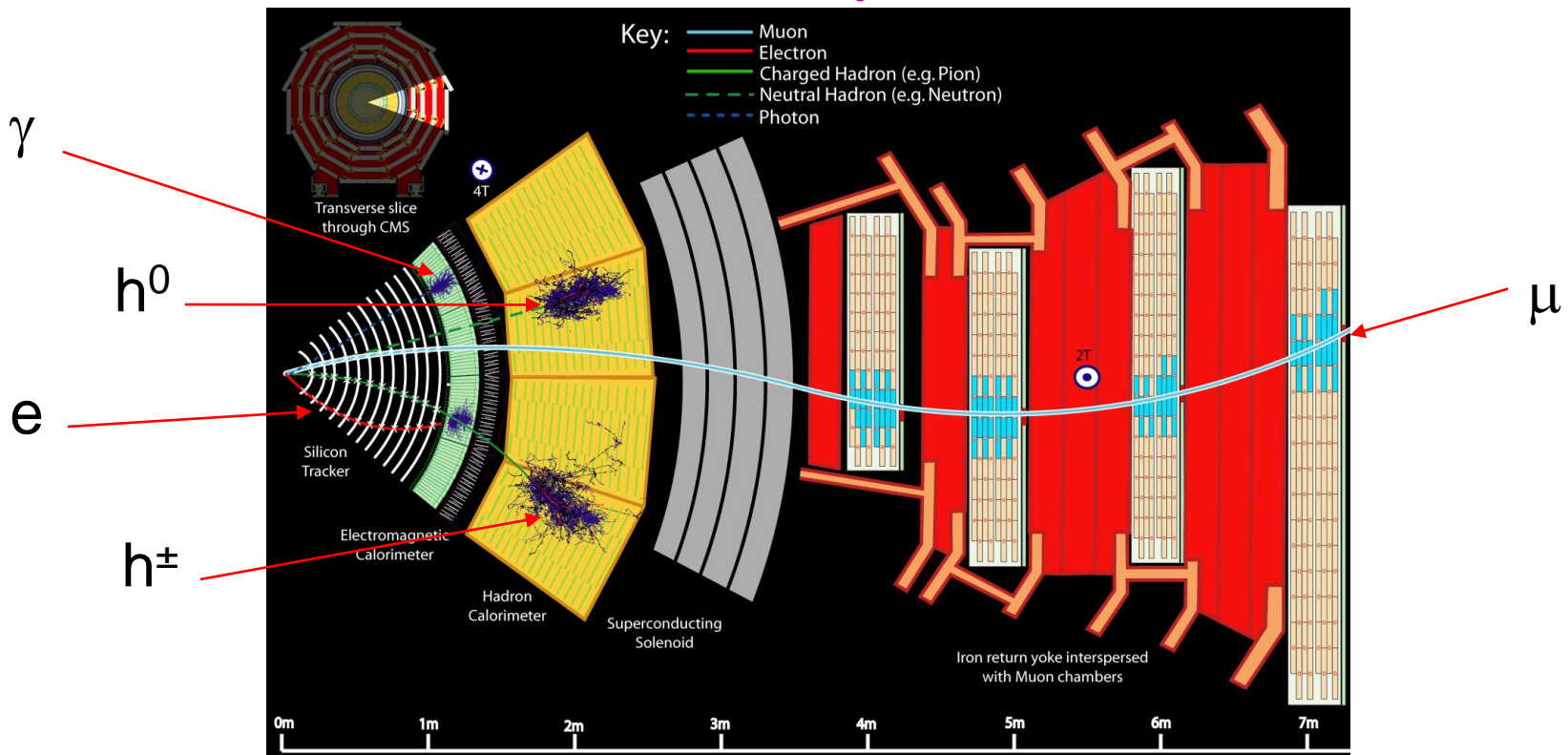
A detector not designed for Particle Flow: CMS (3)

- Let's check the main criteria for a PF-friendly detector
 - ◆ Simplicity
 - ◆ Hermeticity
 - ◆ Magnetic Field
 - ◆ Granularity
 - ◆ Material
 - ◆ Redundancy



Simplicity

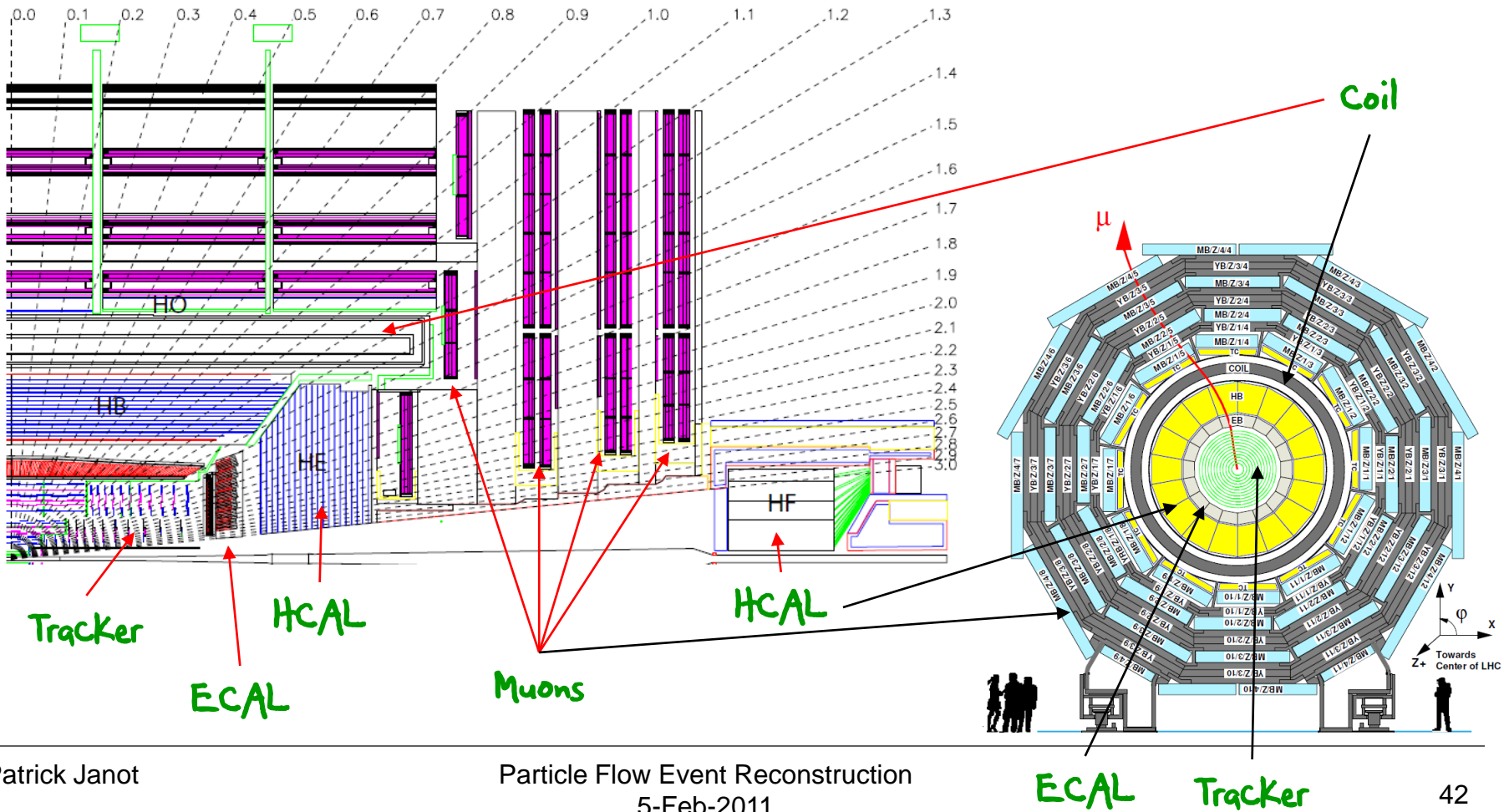
- The design looks “simple” : tracker, ECAL, HCAL, Muon chambers, B Field
 - ◆ Difficult to say at first glance that CMS was not designed for PF
 - It even seems able to follow each type of stable particles
 - Photon, electrons, muons, charged hadrons, neutral hadrons



Hermeticity

□ The detector is hermetic

◆ Calorimetry up to $|η| = 5.1$, tracking up to $|η| = 2.5$, muons up to $|η| = 2.4$



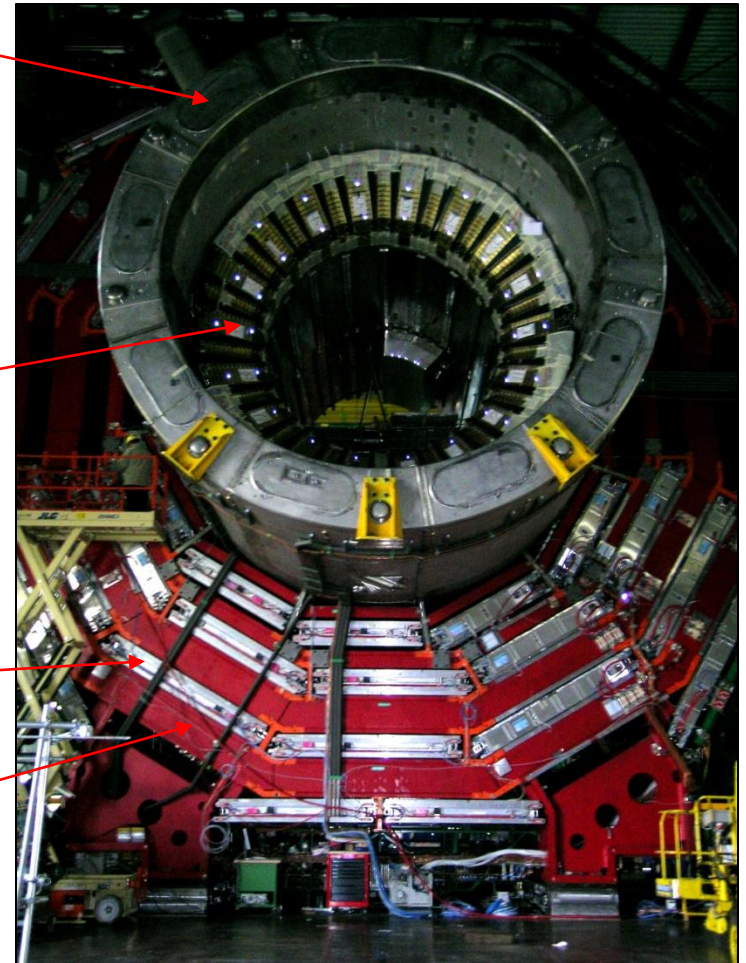
Large magnetic field

□ The magnetic field is large indeed

- ◆ Superconducting magnet
 - Length 12.50m
 - Diameter 6.30m
- ◆ Designed to deliver 4 T
 - Operated routinely at 3.8 T
 - A bonus over ALEPH
- ◆ Note : outside the HCAL !
 - Another bonus over ALEPH
 - See slides about Material

Muons

Iron return yoke



3D Granularity (I)

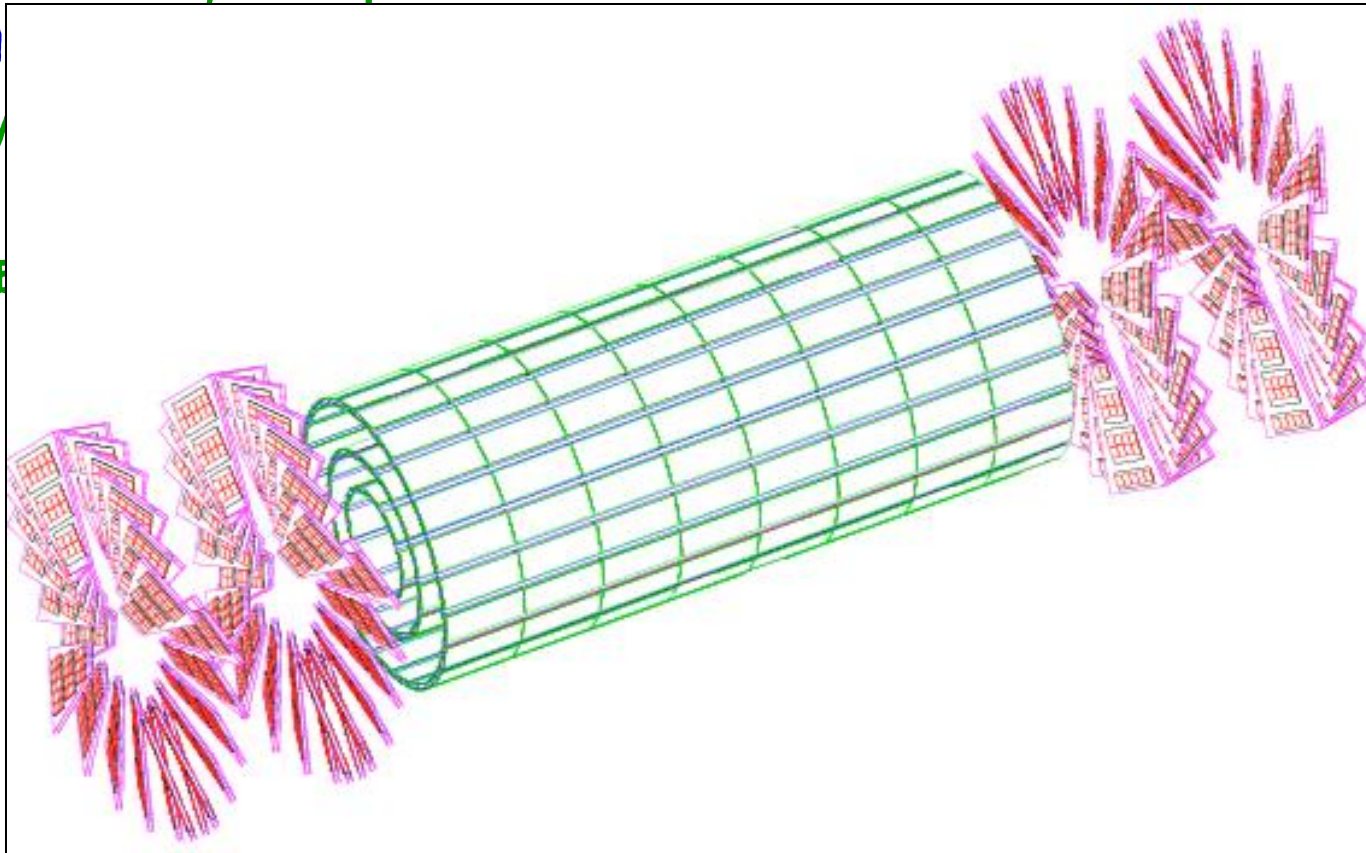
□ The tracker consists of

- ◆ A pixel detector with full 3D granularity, surrounding the beam pipe

- 3 barrel layers, 2 pixel disks

◆ A sil

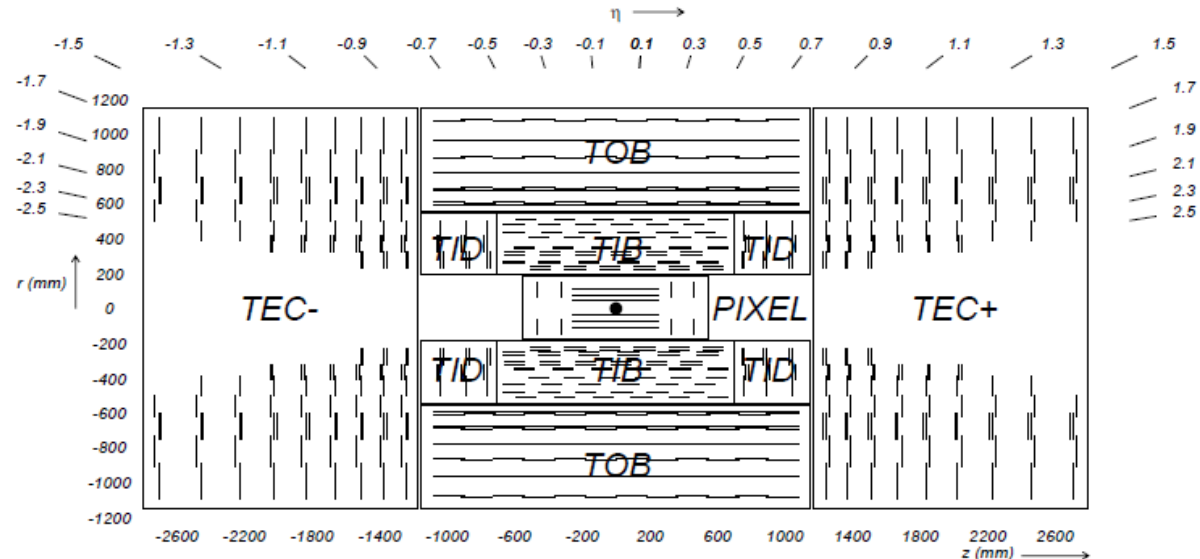
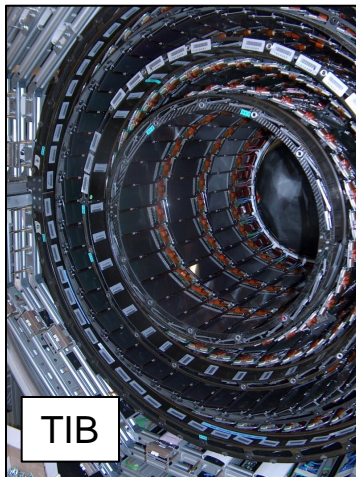
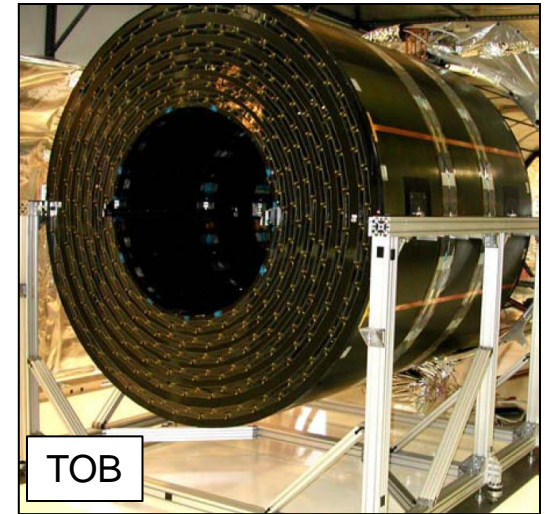
- A
- B



3D Granularity (2)

□ The tracker consists of

- ◆ A pixel detector with full 3D granularity
 - 3 barrel layers, 2 pixel disks
- ◆ A silicon-strip tracker with 3D-ish granularity
 - A few double sided layers
 - TIB 1,2; TOB 1,2; TID 1,2; TEC 1,2,5
 - Expect slightly more ambiguities than in ALEPH
 - Fake tracks ?

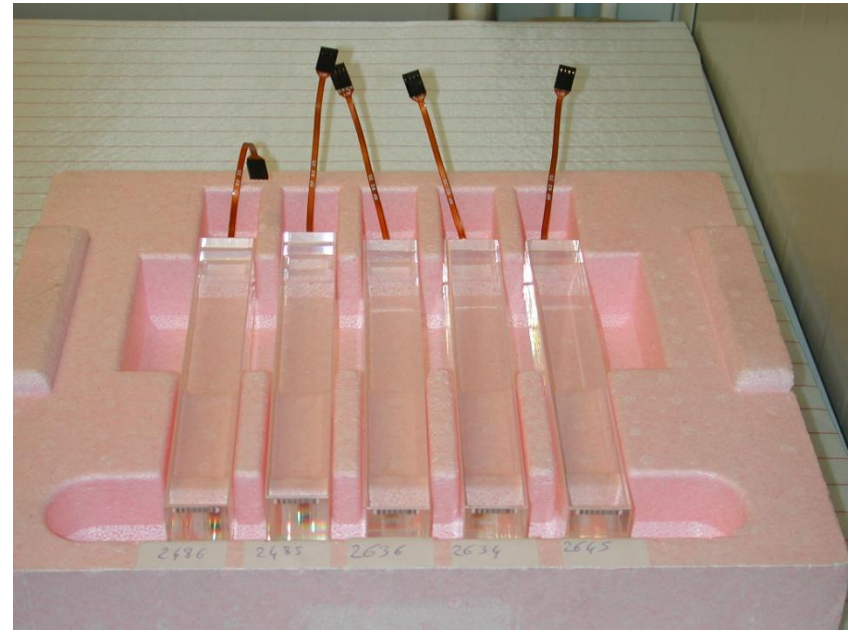


Note the large tracking volume as well

3D Granularity (3)

□ The ECAL consists of

- ◆ 75,848 crystals of PbWO_4
 - (75K towers in ALEPH)
- ◆ With transverse size 2.2 x 2.2 cm
 - (ALEPH : 3 x 3 cm)
- ◆ Moliere radius 2.2 cm
 - (ALEPH 1.6 cm)
- ◆ Radiation Lengths : $25.8 X_0$
 - (ALEPH : $22X_0$)



□ Transverse granularity very similar to that of ALEPH

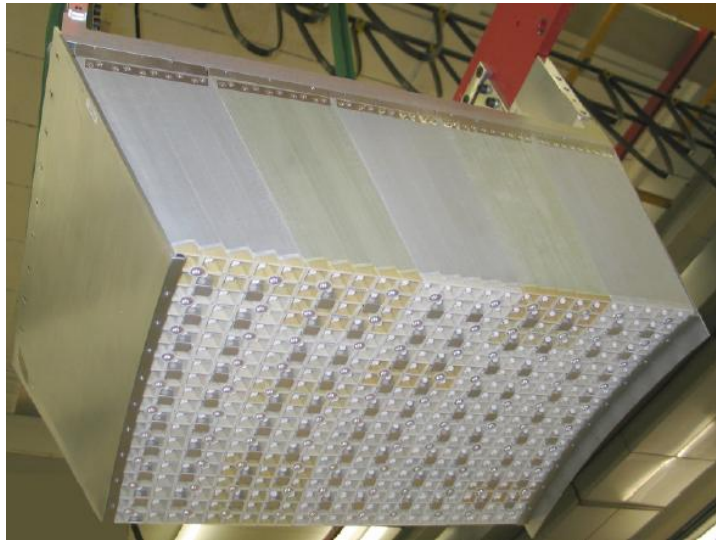
- ◆ But no longitudinal segmentation : photon/electron ID slightly less pure/efficient
 - (3 segmentation in depth in ALEPH + 45 wire planes/module)

□ Excellent energy resolution due to the homogenous material

$$\left(\frac{\sigma}{E}\right)^2 = \left(\frac{2.8\%}{\sqrt{E}}\right)^2 + \left(\frac{0.12}{E}\right)^2 + (0.30\%)^2, \quad \text{(ALEPH: } 20\%/ \sqrt{E}, \text{ but non essential)}$$

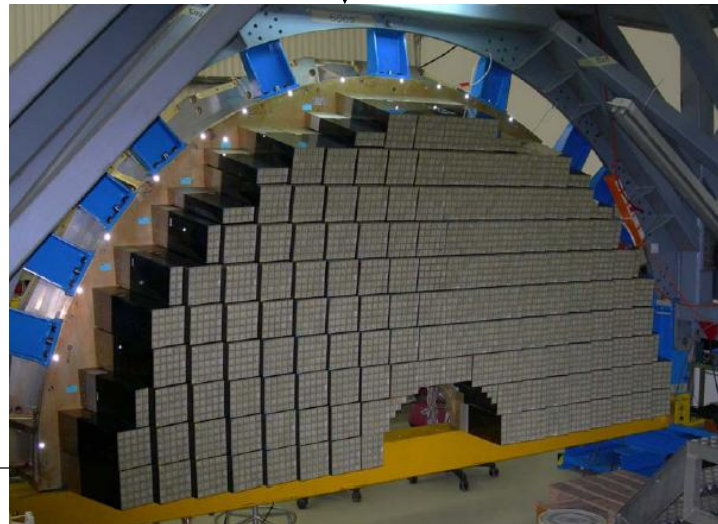
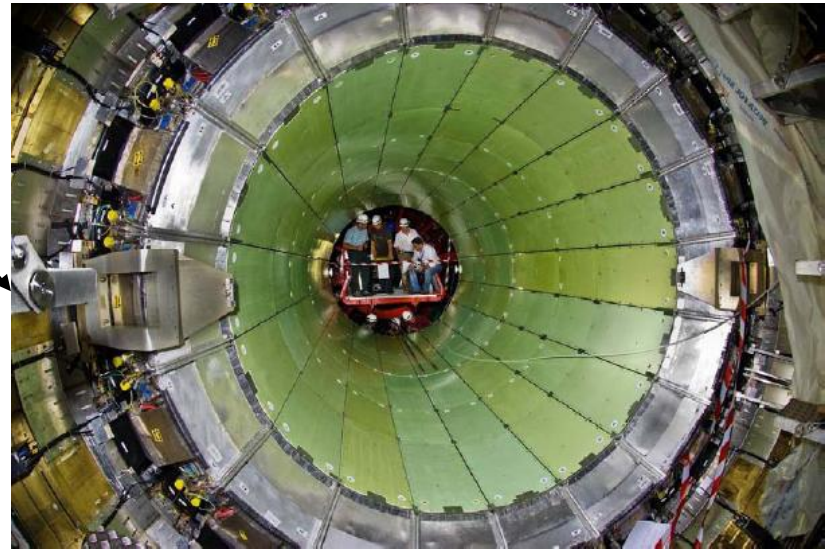
3D Granularity (4)

- ECAL crystals arranged towards quasi hermetic coverage



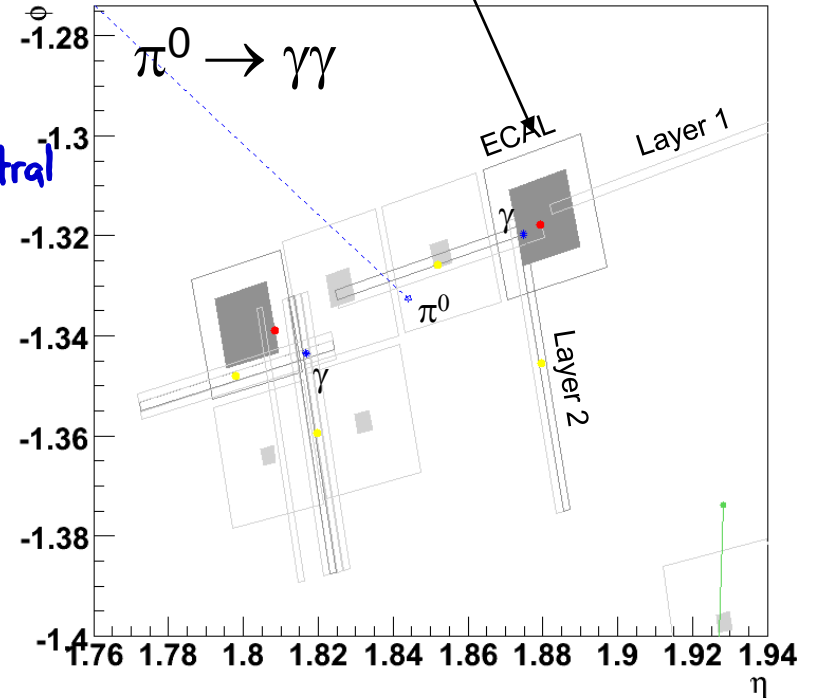
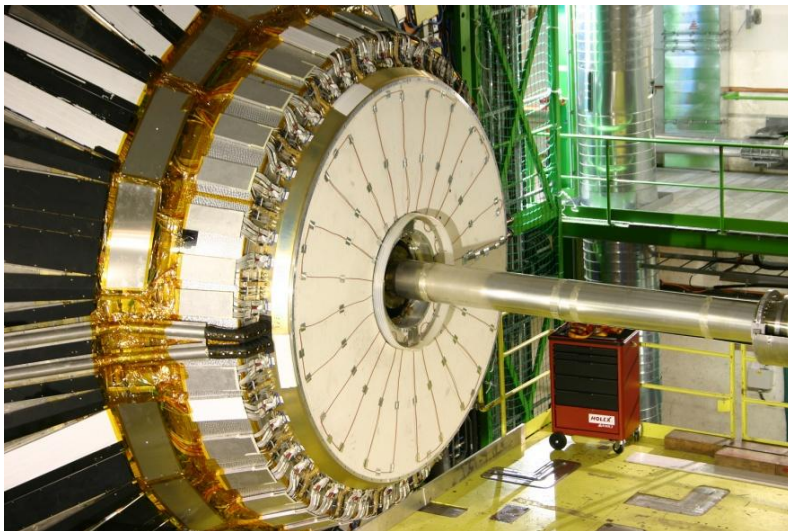
Barrel

End-caps



3D Granularity (5)

- End-caps covered by a preshower detector
 - ◆ with two layers of Lead + Silicon Strips (6 cm long, 2mm wide, $3X_0$)
 - Increase transverse granularity in the end-cap region
 - Add a modest level of longitudinal segmentation
 - Adds complexity
 - Worsens energy resolution
 - ◆ Overall outcome turned out to be neutral



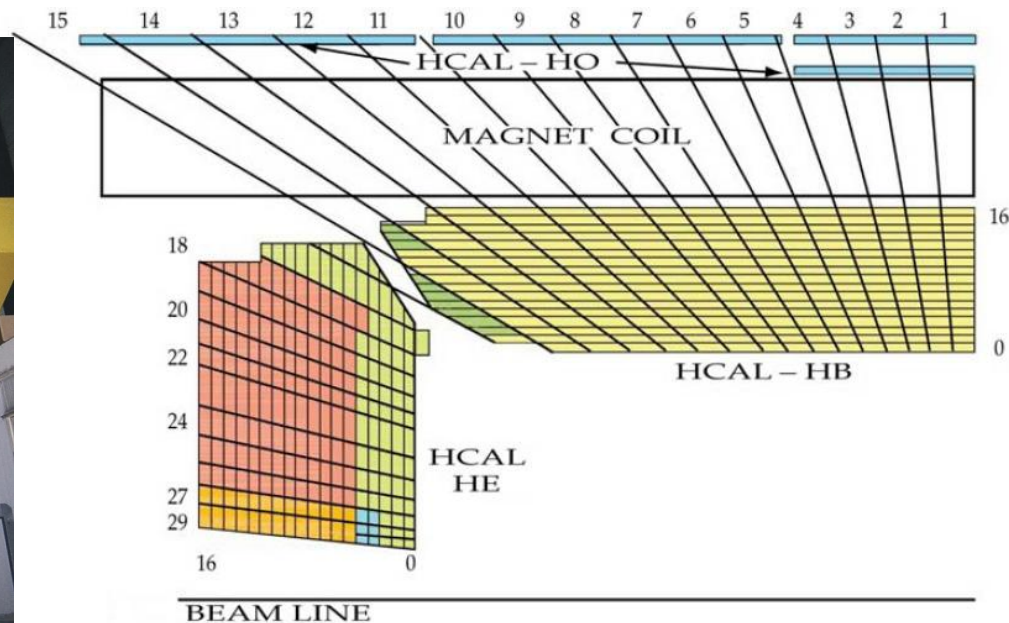
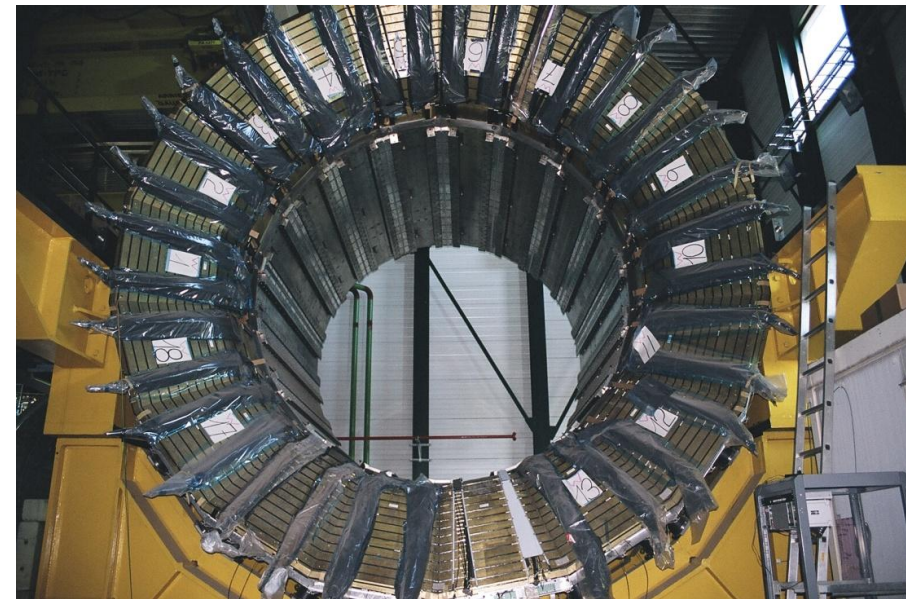
Better γ/π^0 separation ?
(electron pre-Id)

3D Granularity (6)

□ The HCAL consists of

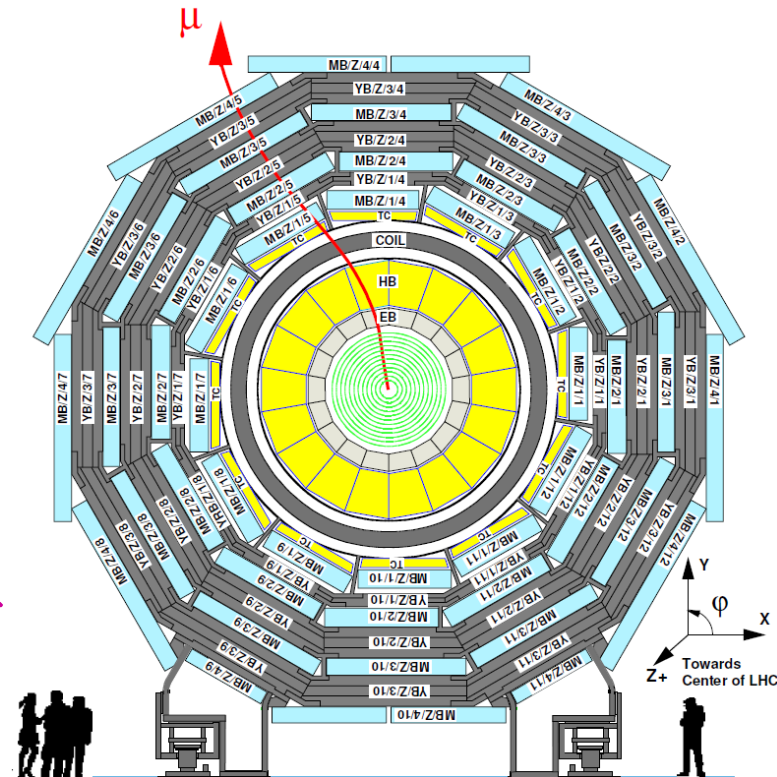
- ◆ A sandwich of 16 plates of brass and scintillator tiles, from 5.5 to 10 λ_0
 - Read out into towers of transverse size 10 x 10 cm
 - Similar transverse granularity to that of ALEPH
 - No longitudinal segmentation in the barrel (except HO), 2 depths in endcaps
 - Similar to ALEPH for hadrons, but no layer-by-layer readout for muons

Thin!



3D Granularity (1)

- ... but extensive 3D muon tracking behind the coil
 - ◆ For example, in the barrel ($|\eta| < 1.2$)
 - 4 muon stations, the first 3 with
 - 8 layers of drift tubes
 - 4 for the $r\phi$ coordinate
 - 4 for the z coordinate
 - And the fourth with only $r\phi$
 - For the momentum resolution
 - ◆ Similar design in end-caps ($|\eta| < 2.4$)
 - But cathode-strip chambers instead
 - with anode wire measurements for z
- Muons should not be a problem
 - ◆ But watch out HCAL thickness
 - Energetic pion punch-through likely



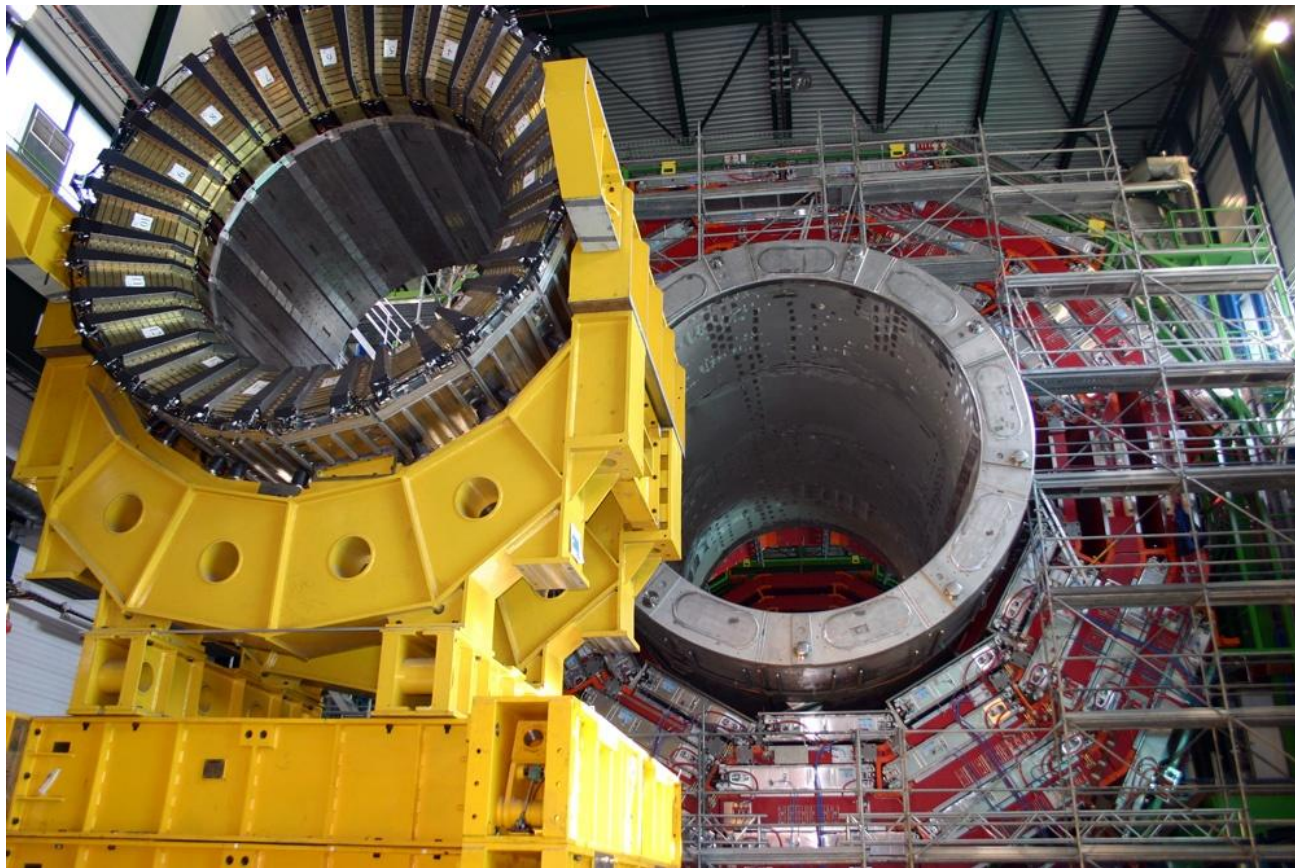
Redundancy : tough luck !

□ That one is easy

- ◆ There is essentially no redundancy for the energy measurements
 - Neither in the ECAL (only the crystal light is collected)
 - Nor in the HCAL (only the scintillator light is collected)
- ◆ Calorimeter cleaning might be challenging
 - And CMS calorimeters turned out to be more noisy than ALEPH's ...
 - Need to use isolation and timing of the signals to disentangle with noise
 - Careful with overcleaning !
- ◆ Particle identification might be complicated by the absence of redundancy
 - But let's see how it comes out.

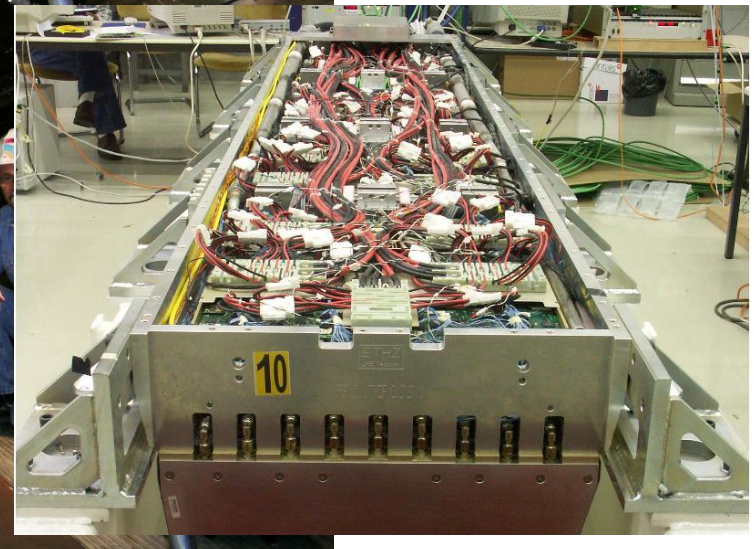
Material Budget (1)

- As we already mentioned : the coil is behind the HCAL
 - ◆ One λ_0 less in front of HCAL with respect to ALEPH
 - Less neutral hadrons lost, less charged hadrons showering



Material budget (2)

- But ECAL services ($0.5\lambda_0$) between ECAL and HCAL
 - ◆ Mostly back to ALEPH situation ... It was too good to be true !

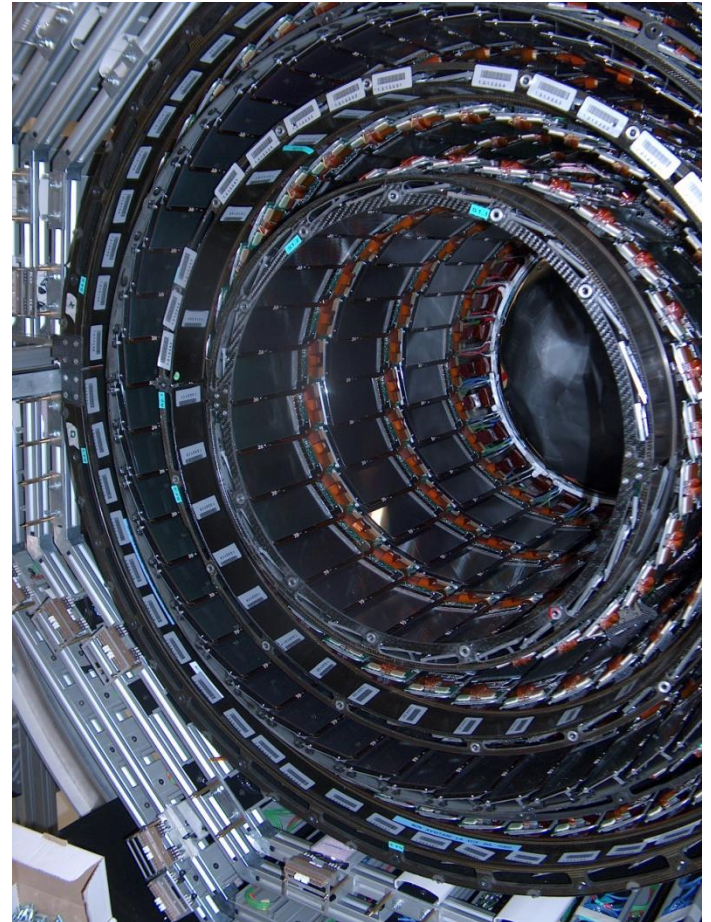


Material budget (3)

- The tracker is not a TPC (too slow for LHC), but a silicon tracker
 - ◆ Remember the TPC [mostly empty] ? And here we are now [mostly full] !



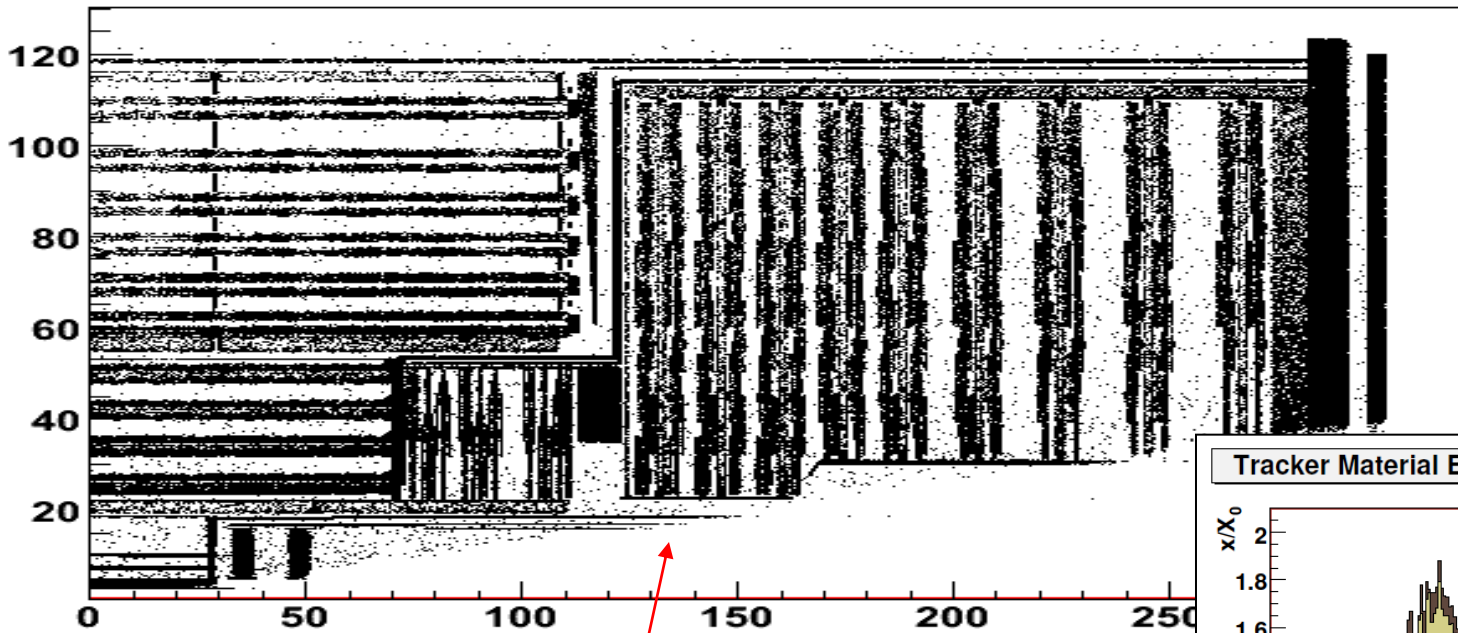
- The CMS tracker looks very thick



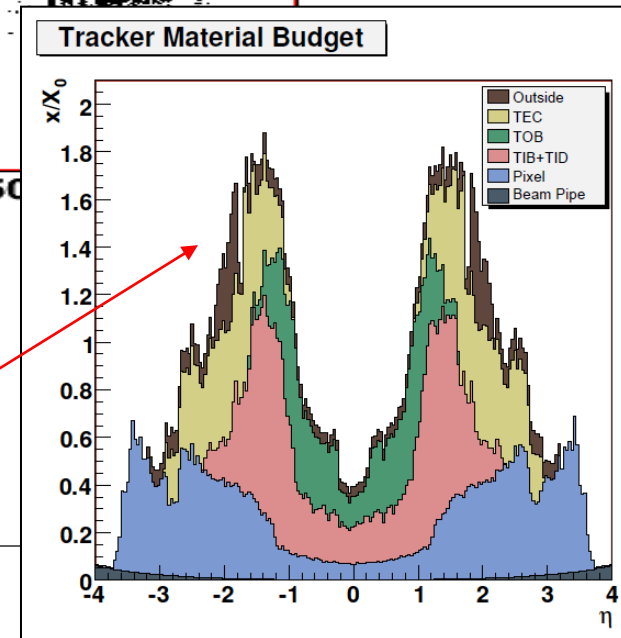
Material budget (4)

□ ... And it is thick indeed !

Full Tracker radiography (MC information)



- ◆ Vertices of converted single photons (100K photons)
 - Services account for 90% of the material
 - From $0.5X_0$ up to $2X_0$ in the tracker !
- (was predicted to be $0.8X_0$ in 2000)



CMS design and Particle Flow : Summary

- Situation so far, when compared to ALEPH
 - ◆ Simplicity : OK
 - (despite the presence of the pre-shower ?)
 - ◆ Hermeticity : OK
 - (did not talk of the forward HCAL, no tracker there)
 - ◆ Large magnetic field : OK
 - Actually much larger than ALEPH, may compensate granularity
 - ◆ 3D granularity : almost OK
 - Less tracker measurements, not all 3D; No ECAL longitudinal segmentation
 - Note: much better E resolution in ECAL; slightly worse in HCAL ($120\% \sqrt{E}$)
- So far so good - and even quite promising - but ...
 - ◆ No redundancy whatsoever
 - ◆ Tracker material might be a killer for particle-flow reconstruction
 - Note: what is the need of a pre-shower after $2X_0$ in the tracker ?

Overall CMS environment (1)

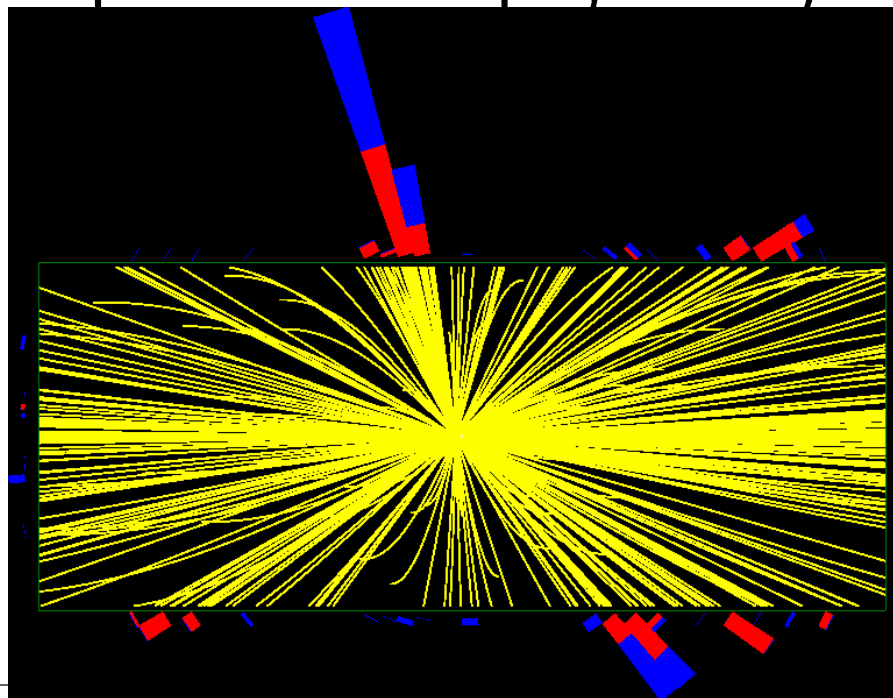
□ Two major differences with ALEPH

- ◆ The LHC is a pp collider (and from time to time a PbPb collider)

- Much larger particle multiplicity in the final state than at LEP

→ Not mentioning pile-up collisions, which enhance the problem

More confusion is possible, although granularity is mostly adequate to keep the detector occupancy low enough



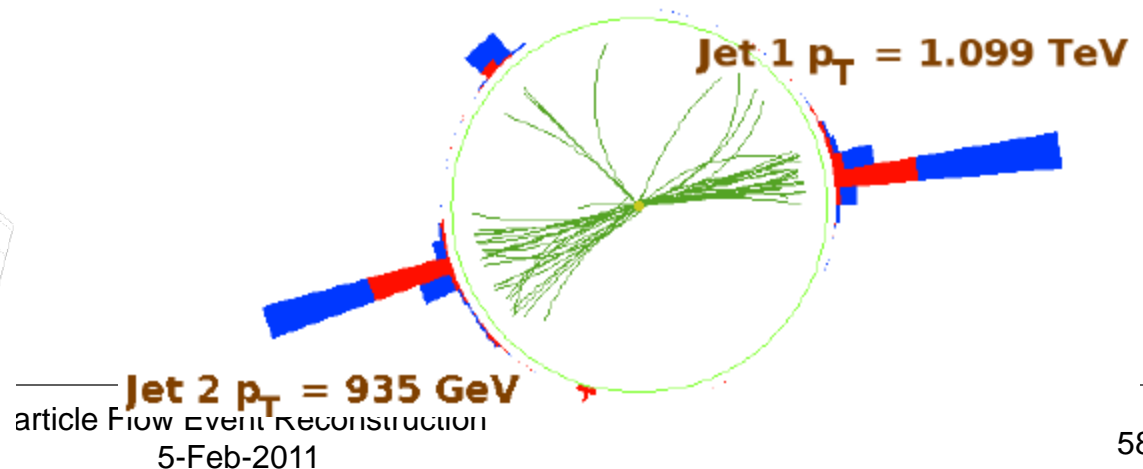
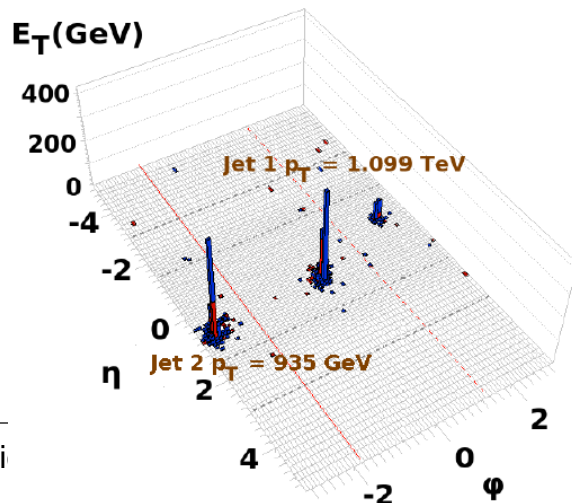
Overall CMS environment (2)

□ Two major differences with ALEPH (cont'd)

- ◆ Energy spectrum for reconstructed particles from 200 MeV to 5 TeV
 - Was 100 MeV to 100 GeV at LEP
- ◆ For jets with very large p_T , charged and neutrals are less separated
 - Separation becomes smaller than detector segmentation at one point;
 - Particle identification becomes less efficient; momenta less well measured;
 - The PF performance converges (for jets) to a pure calo determination ?

→ Given the B field, the Σp_T at which this happens is probably large

→ There are always low p_T particles in jets, for which PF will help anyway

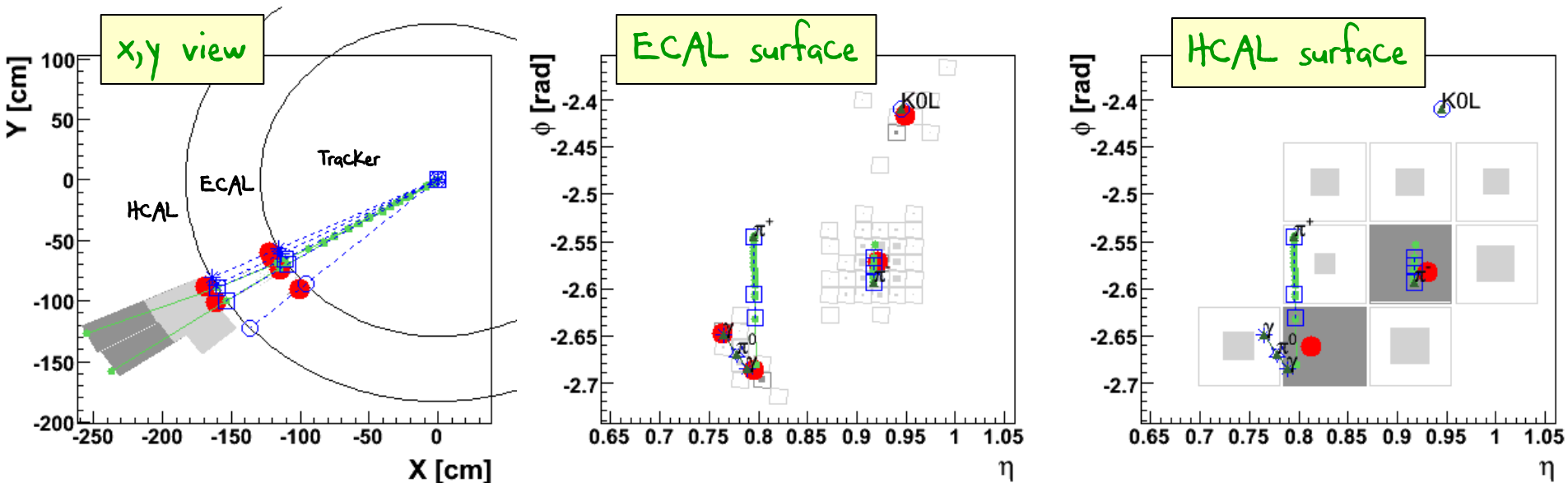


Back to 2005 : Checking CMS basic promises (1)

- Pre-requisite for particle-flow reconstruction
 - ◆ Reconstruct (some) charge particle tracks for charged hadrons
 - Seeded with at least two hits in the pixel detector
 - Originating from the beam axis within tight tolerances
 - Pattern recognition with a combinatorial Kalmann-Filter track finder
 - Request at least 8 hits and $p_T > 1 \text{ GeV}/c$
 - ◆ Reconstruct clusters in the ECAL and HCAL, for photons and hadrons
 - Seeded by cells above a given energy threshold
 - 230 MeV in ECAL, 1 GeV in HCAL
 - Surrounded by 4 direct neighbouring cells with smaller energies
 - Optimal use of the available granularity (better than ALEPH for HCAL)
 - There might be several seeds/sub-clusters in a given “topological” cluster
 - Share the cell energies among the sub-clusters (according to $d_{\text{seed},\text{cell}}$)
 - ◆ Ignore electrons and muons for the time being (for pedagogical purposes)

Back to 2005 : Checking CMS basic promises (2)

- First fully simulated jet, $p_T = 100$ GeV (Tracker thickness $< 1X_0$)
 - ◆ Four particles : π^+ , π^- , π^0 and K_L^0



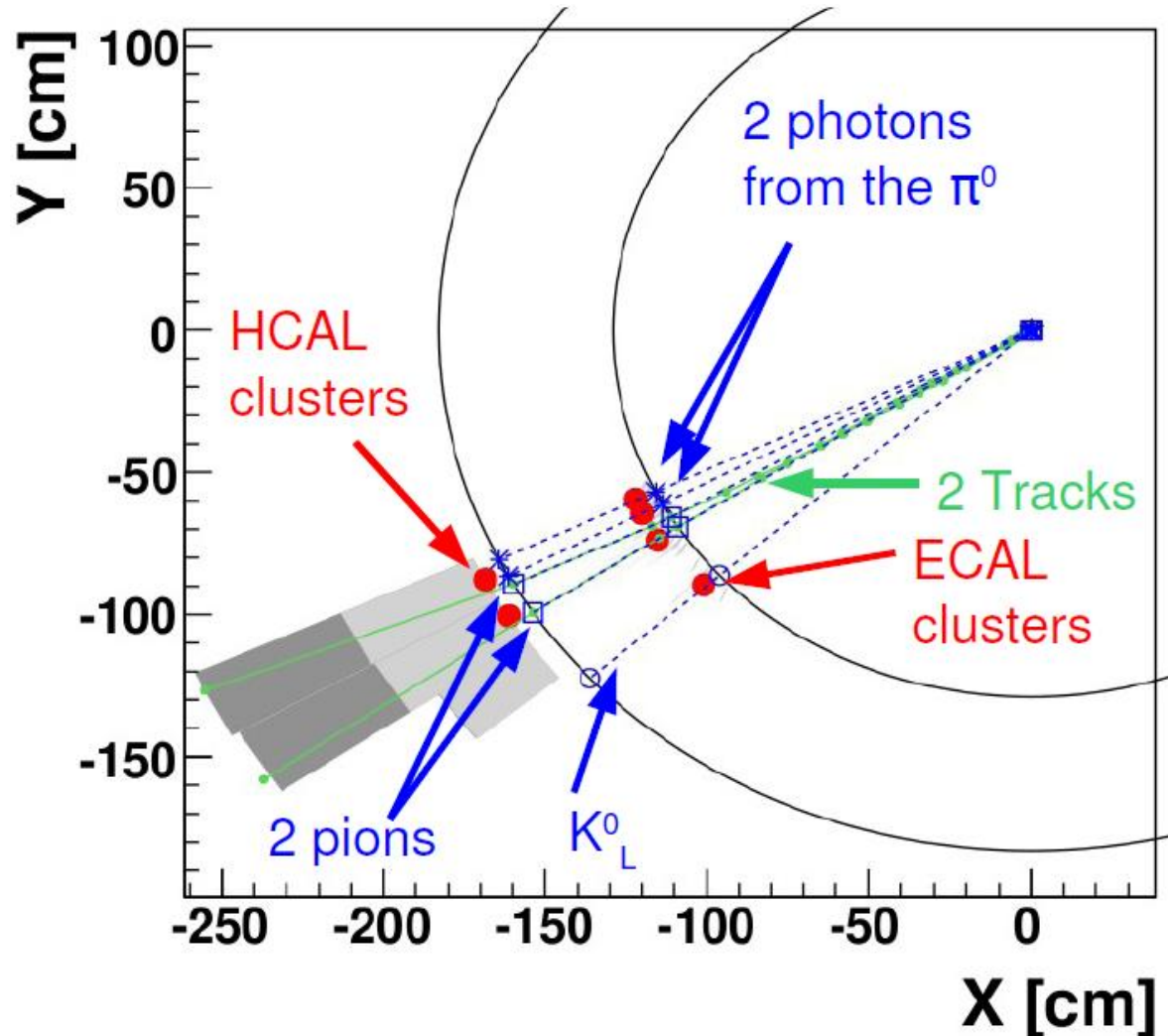
True particles in blue, Tracks and tracker hits in green

Cluster seeds in dark grey, cluster position in red, other cluster cells in light gray

- From this sole jet, PF feasibility with CMS was declared
 - ◆ Simplicity, magnetic field and granularity appeared to be adequate indeed.

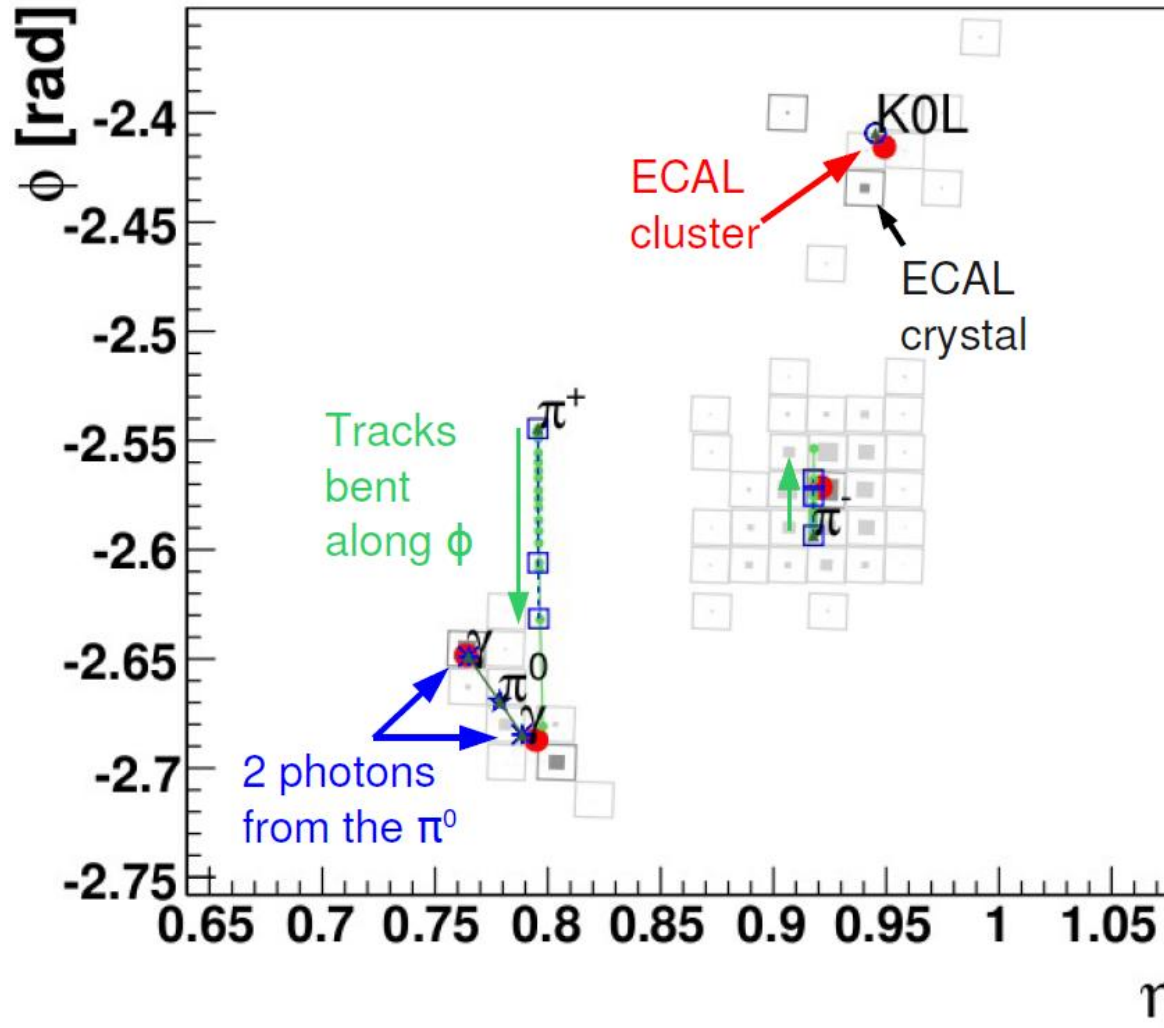
Back to 2005 : checking CMS basic promises (3)

- The x,y (or r,ϕ) view in more details



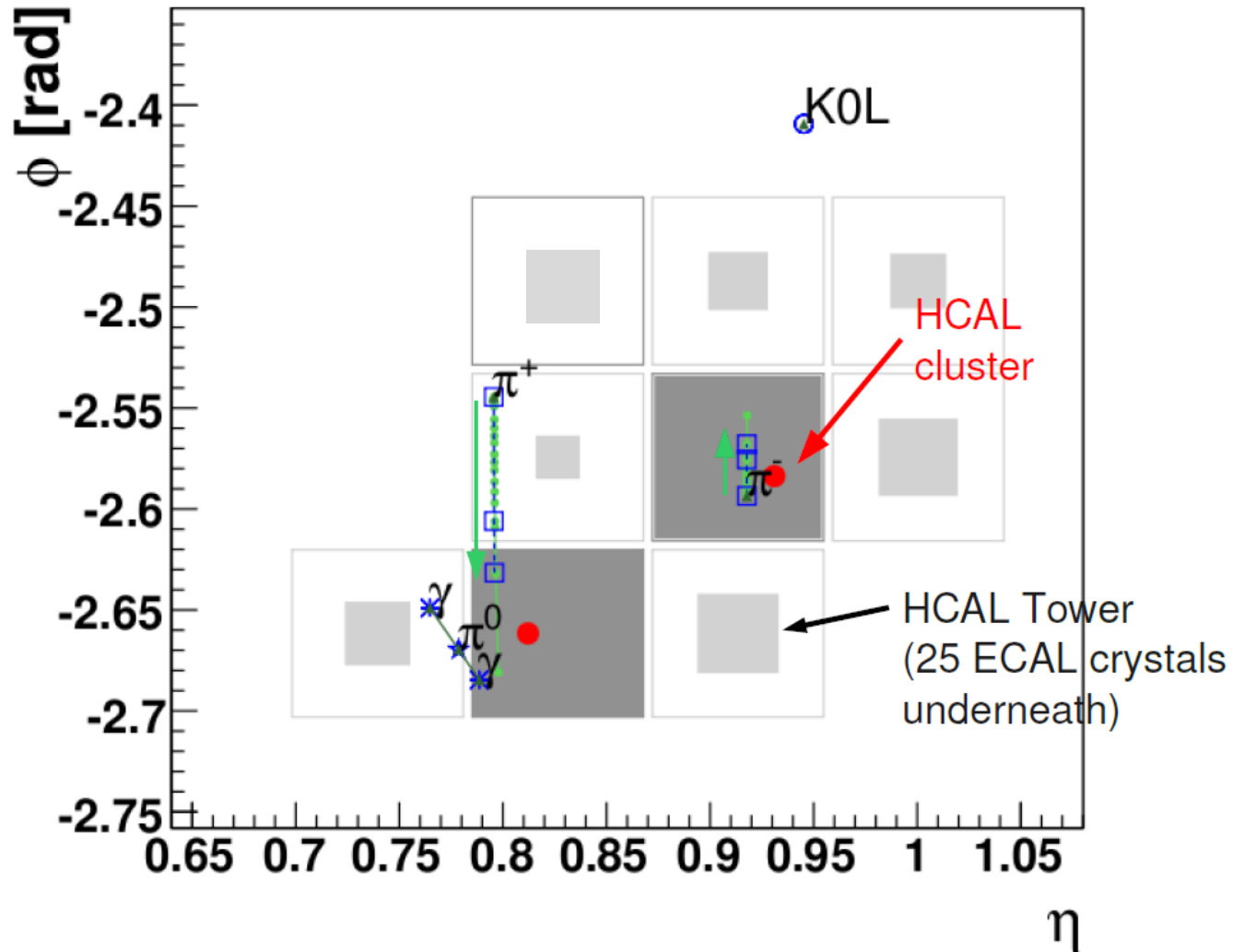
Back to 2005 : Checking CMS basic promises (4)

- The ECAL surface (η, ϕ) view in more details



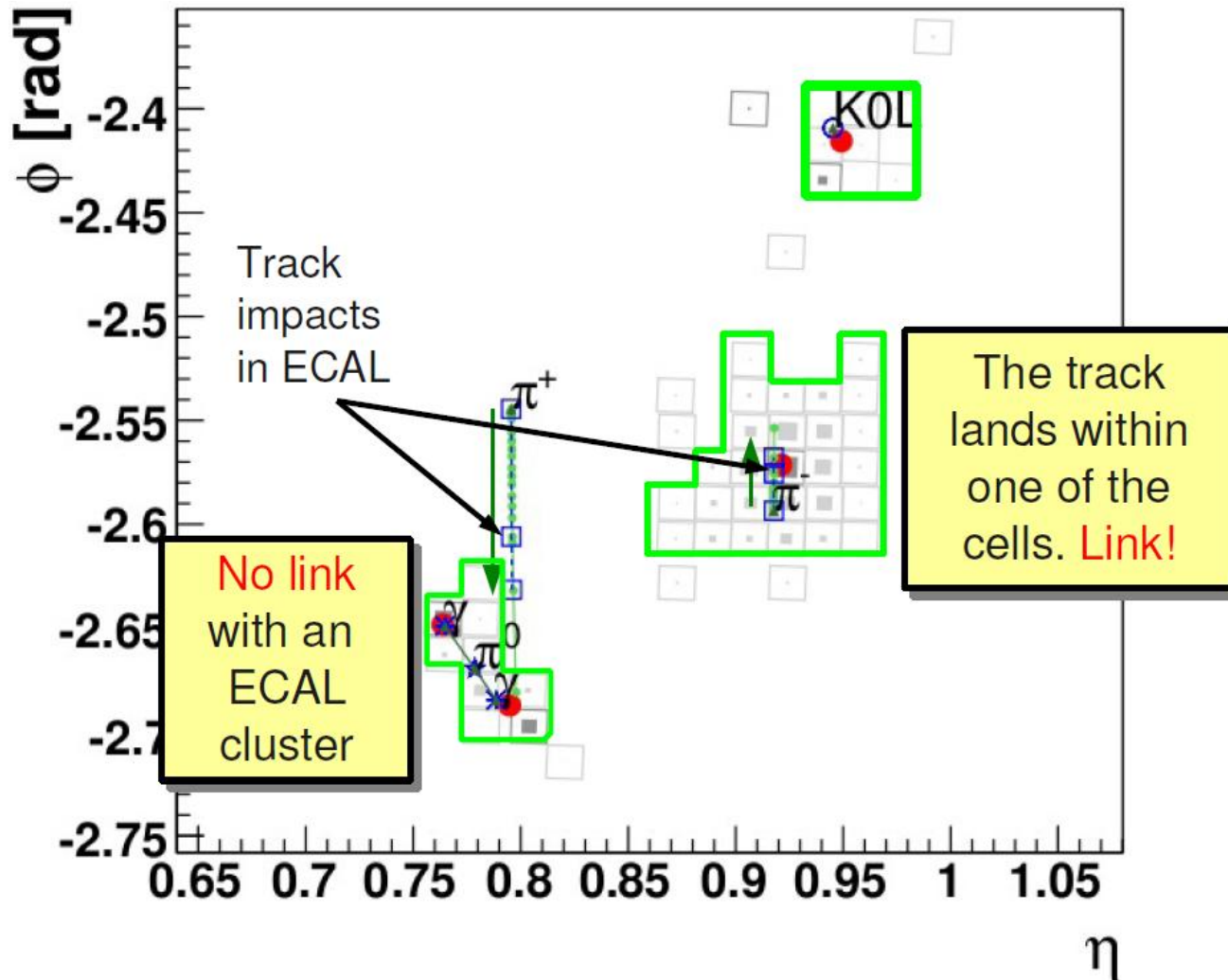
Back to 2005 : Checking CMS basic promises (5)

- The HCAL surface view (η, ϕ) in more details



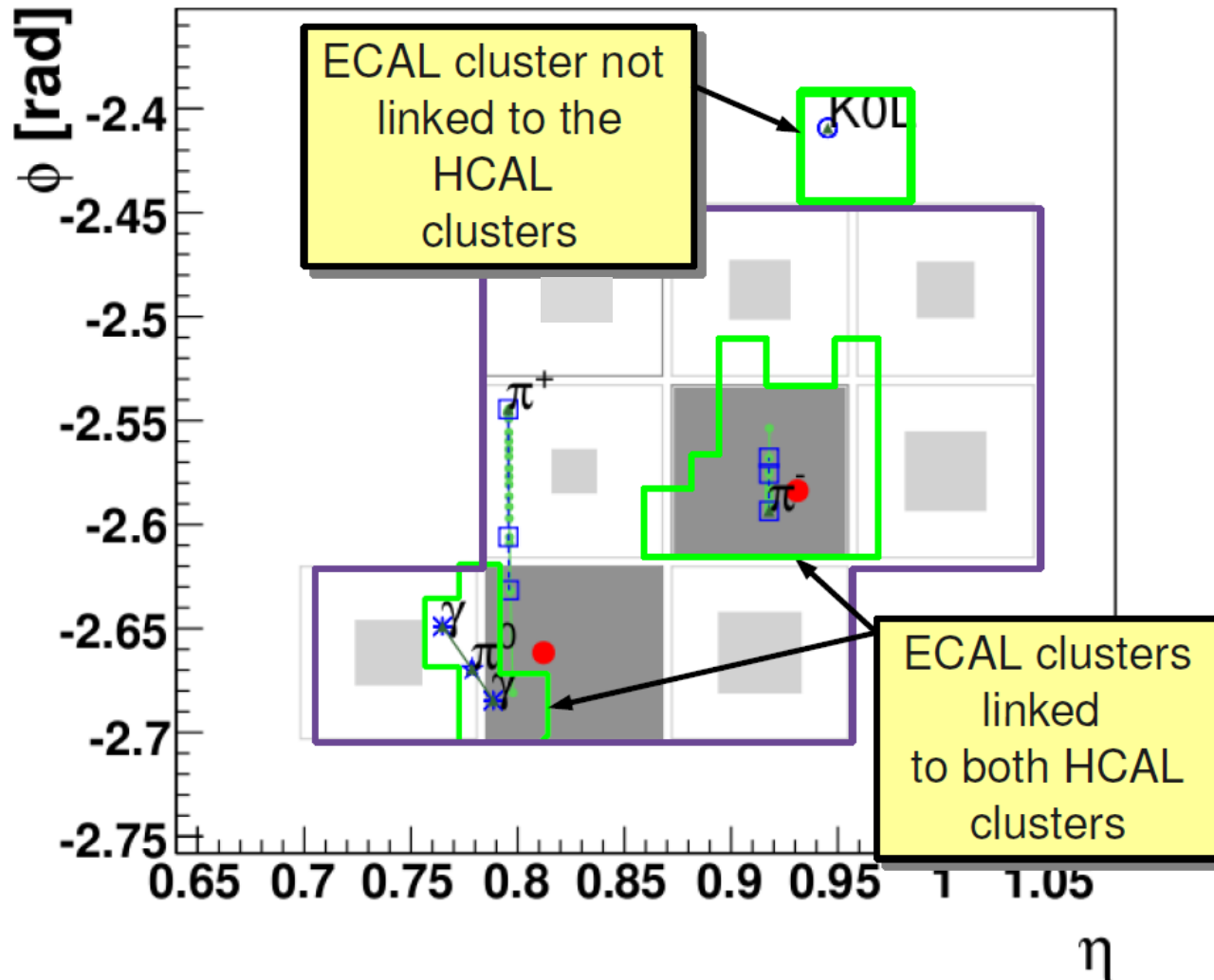
Back to 2005 : checking CMS basic promises (b)

□ Track - ECAL cluster link



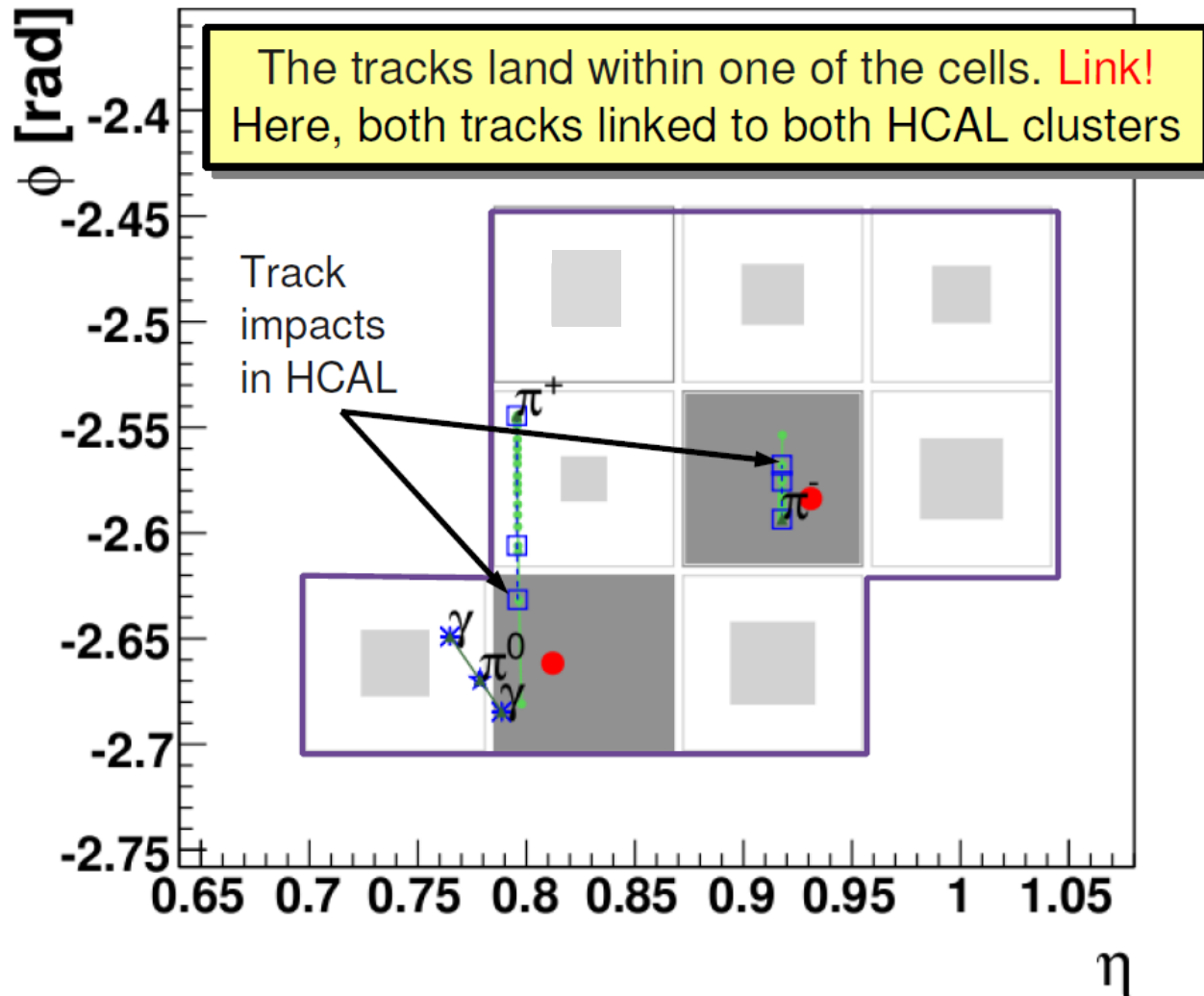
Back to 2005 : checking CMS basic promises (7)

- ECAL cluster - HCAL cluster link



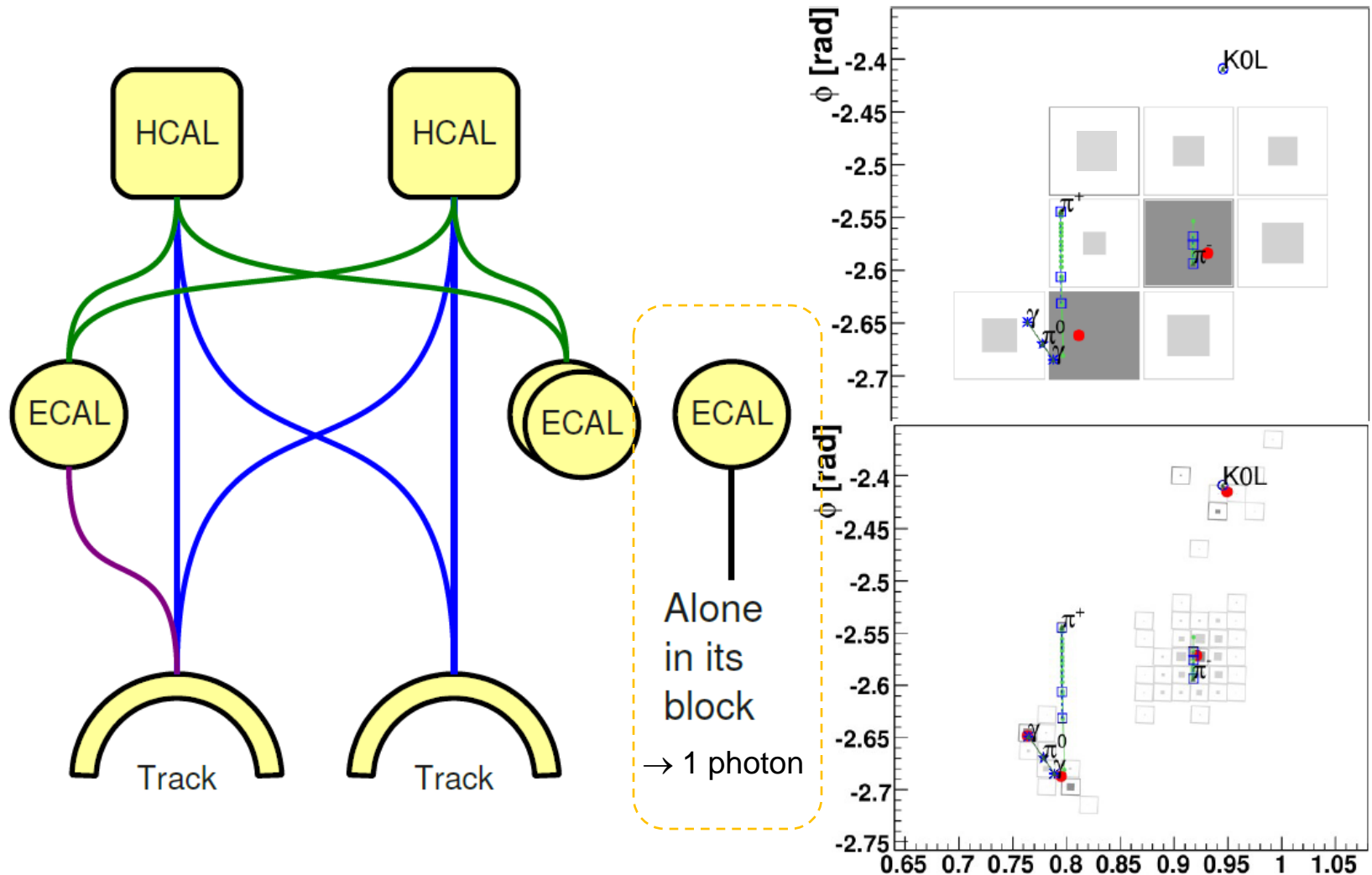
Back to 2005 : Checking CMS basic promises (8)

□ Track - HCAL cluster link



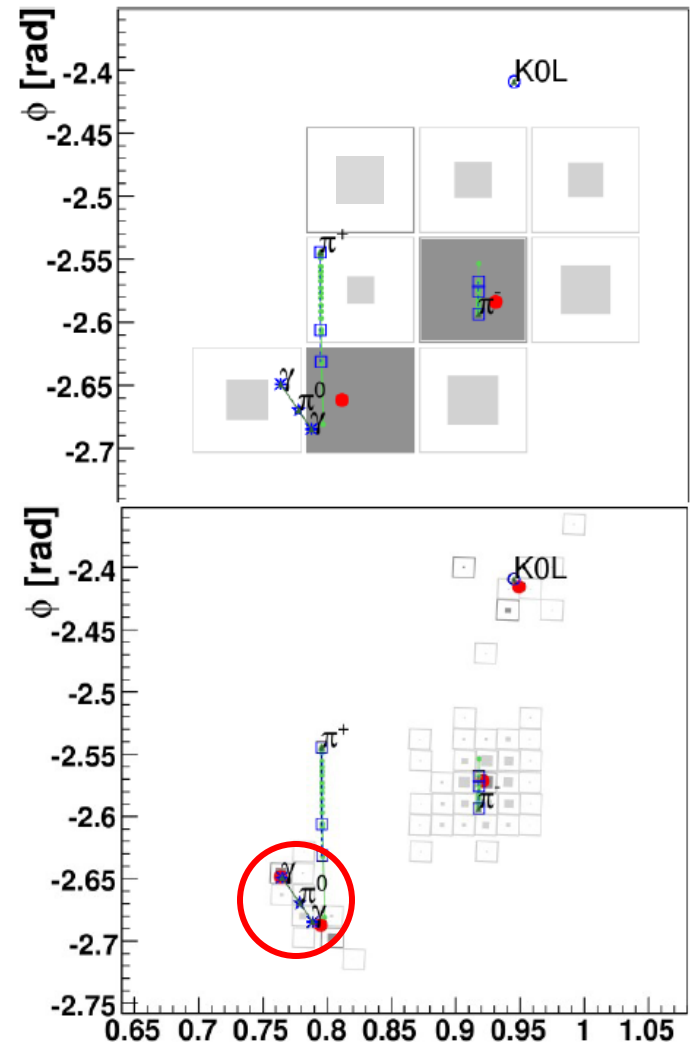
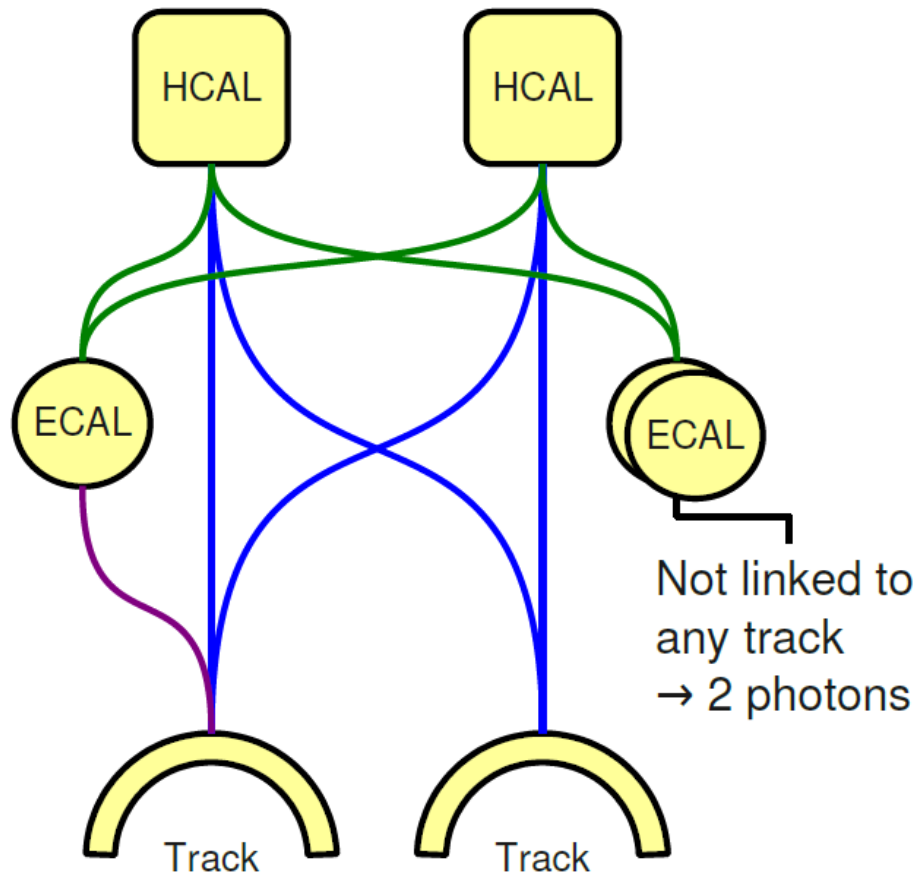
Back to 2005 : checking CMS basic promises (9)

- Build "blocks" of elements linked to each other



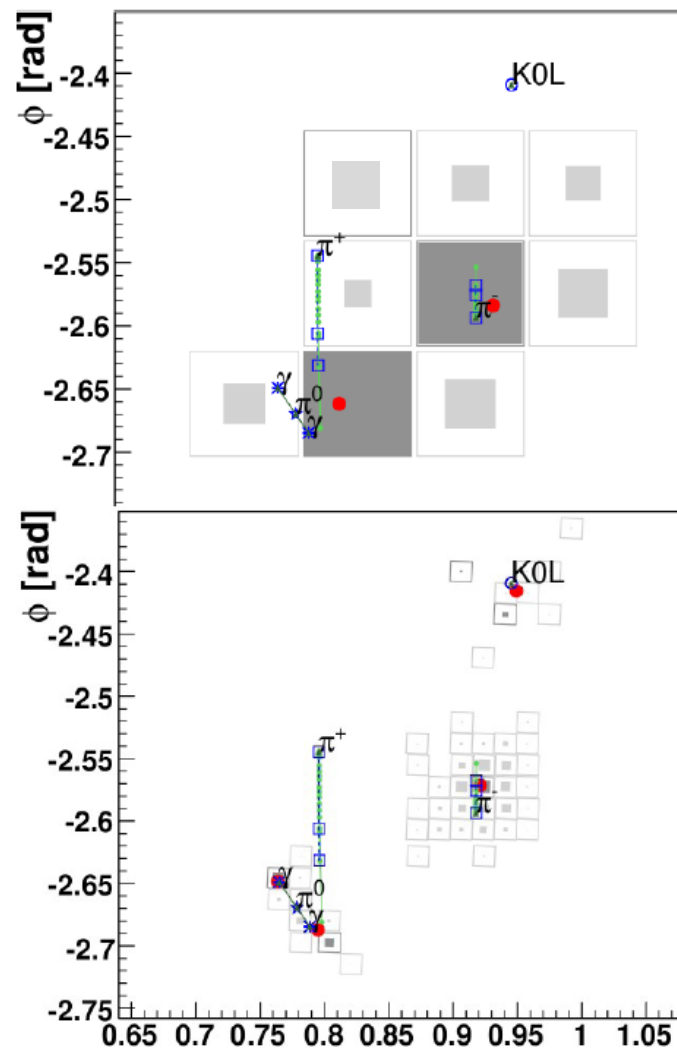
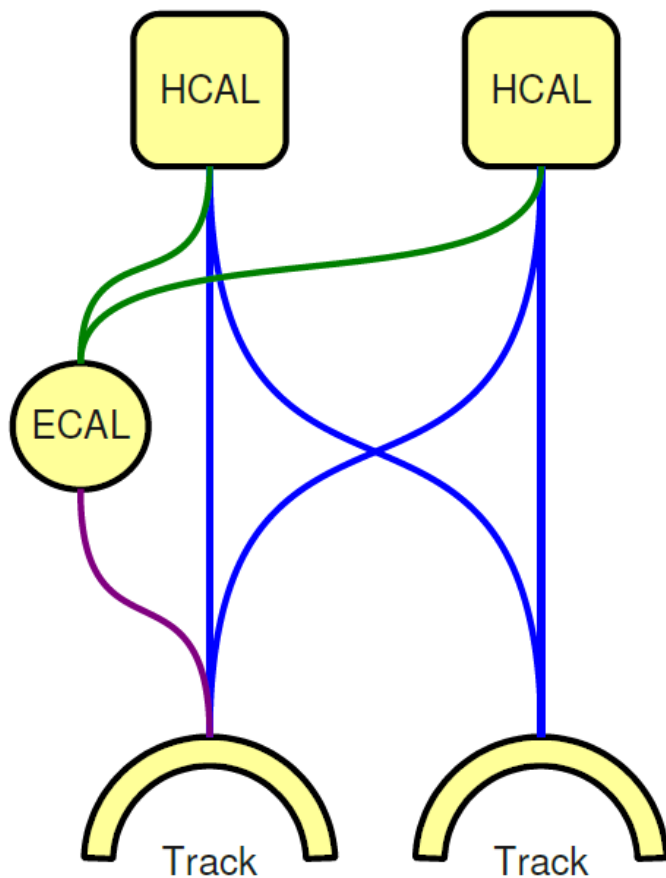
Back to 2005 : Checking CMS basic promises (10)

- Find isolated photons in the blocks



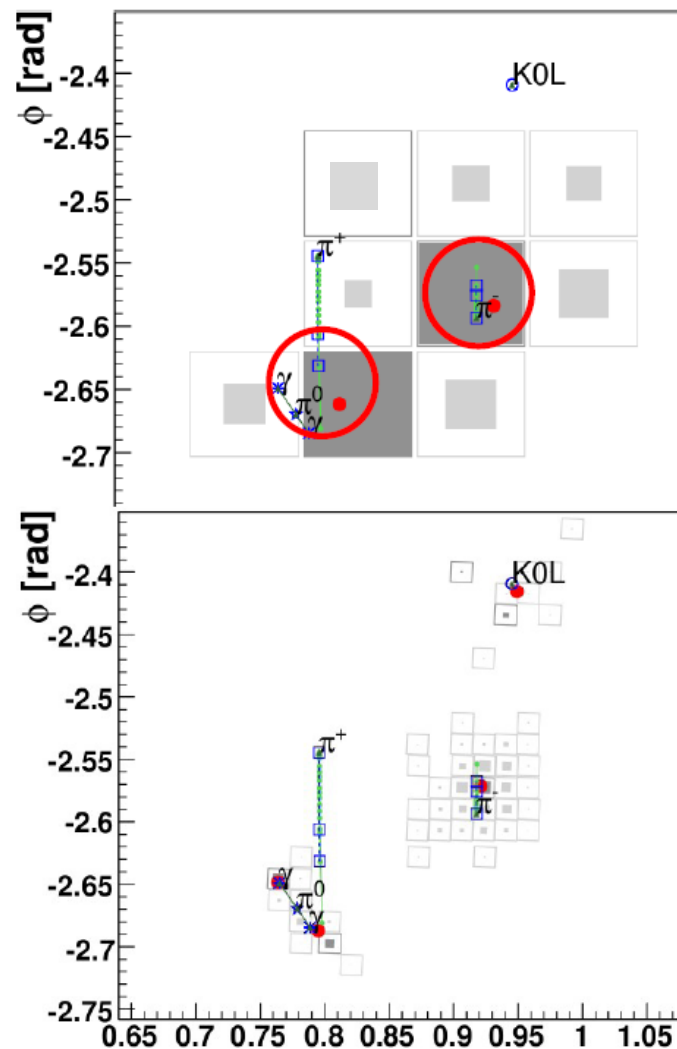
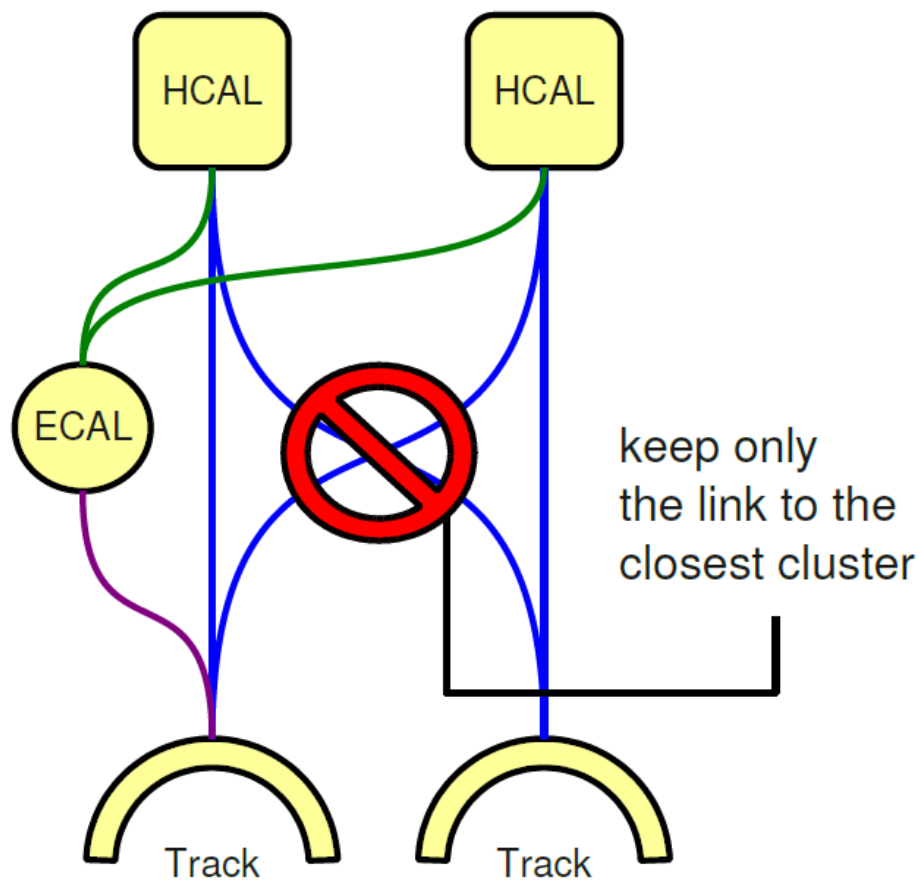
Back to 2005 : Checking CMS basic promises (II)

- Simplified block (1st step)



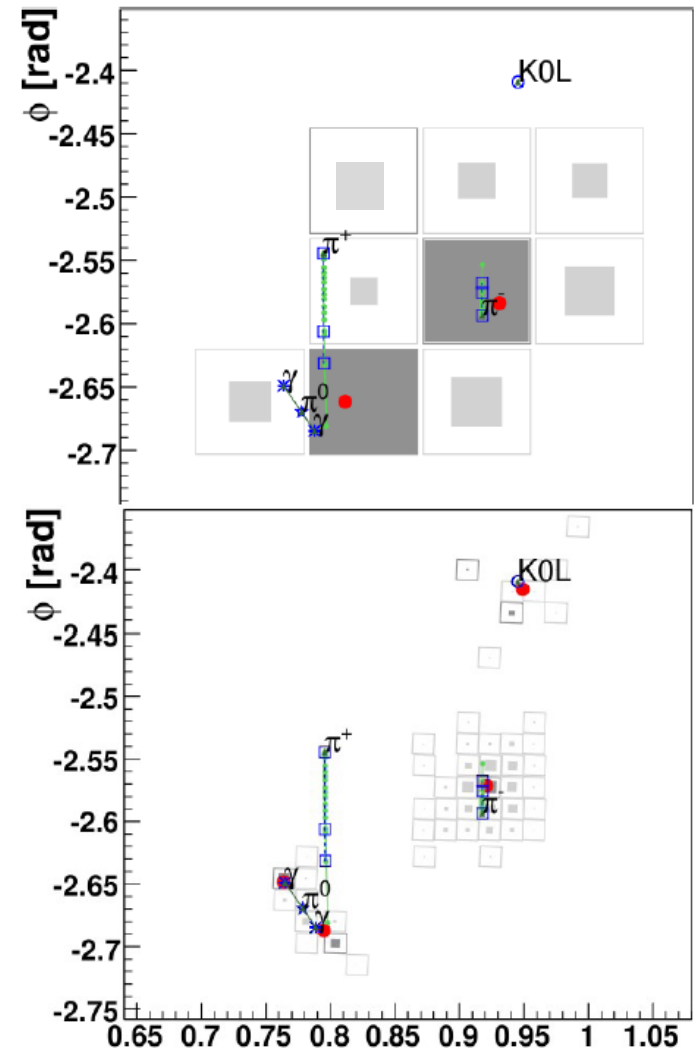
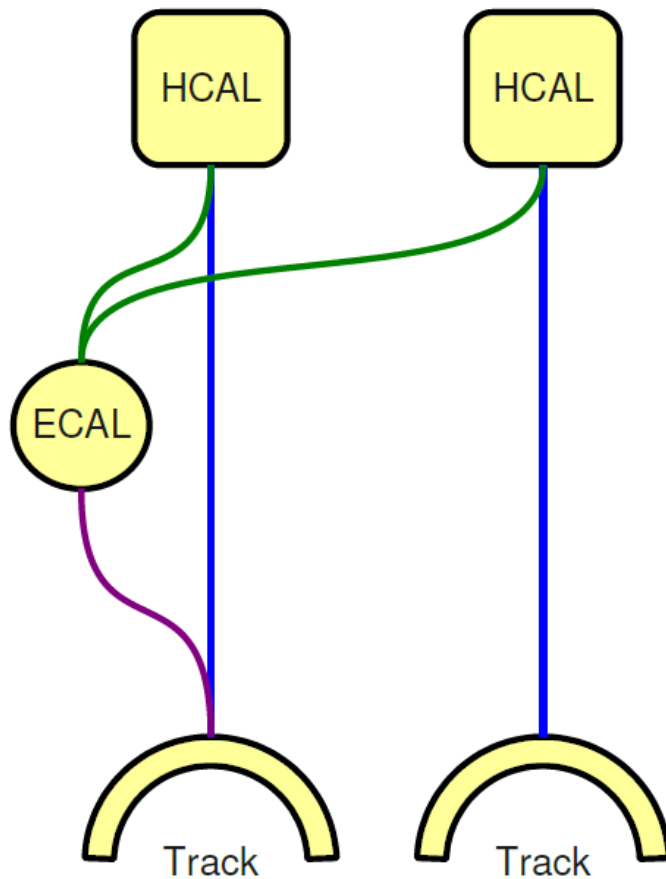
Back to 2005 : Checking CMS basic promises (12)

- Optimize the use of HCAL granularity



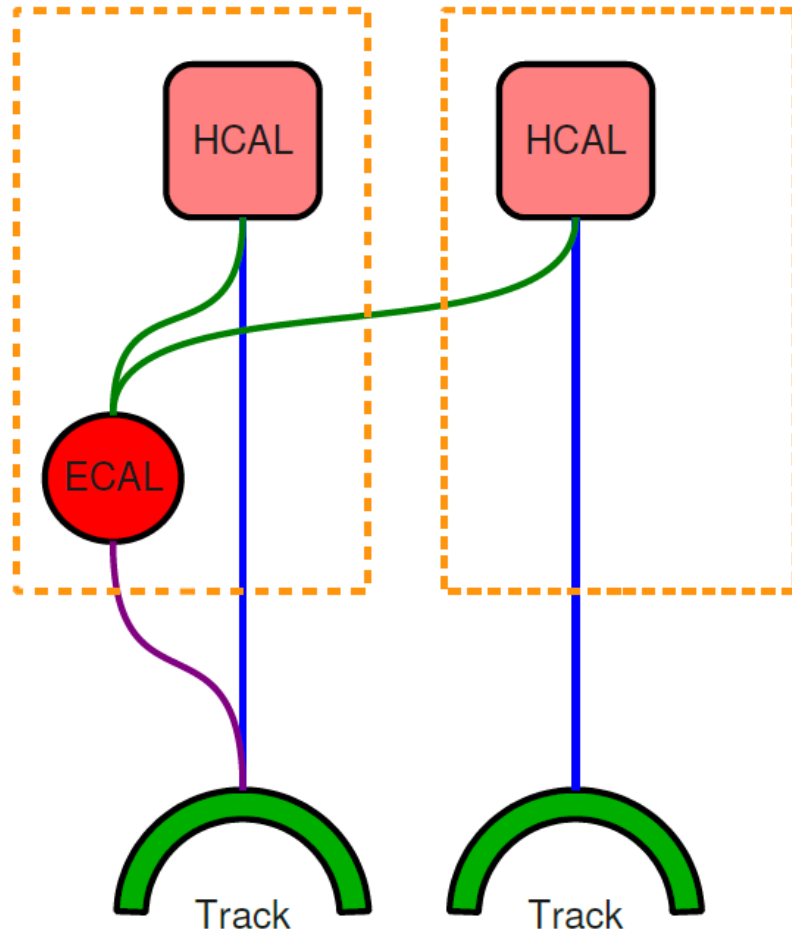
Back to 2005 : Checking CMS basic promises (13)

- Further simplified block (2nd step) : blocks are usually very small !



Back to 2005 : Checking CMS basic promises (14)

- Find charged hadrons, and merged photons / neutral hadrons



For each HCAL cluster, compare :

The sum of the track momenta p
The sum of the cluster energies E

Linked to the tracks

In ECAL and in HCAL

“Hadron Calibrated” (see later)

If p and E are compatible

Charged hadrons only (1 per track)

If $E > p + 120\% \sqrt{p}$

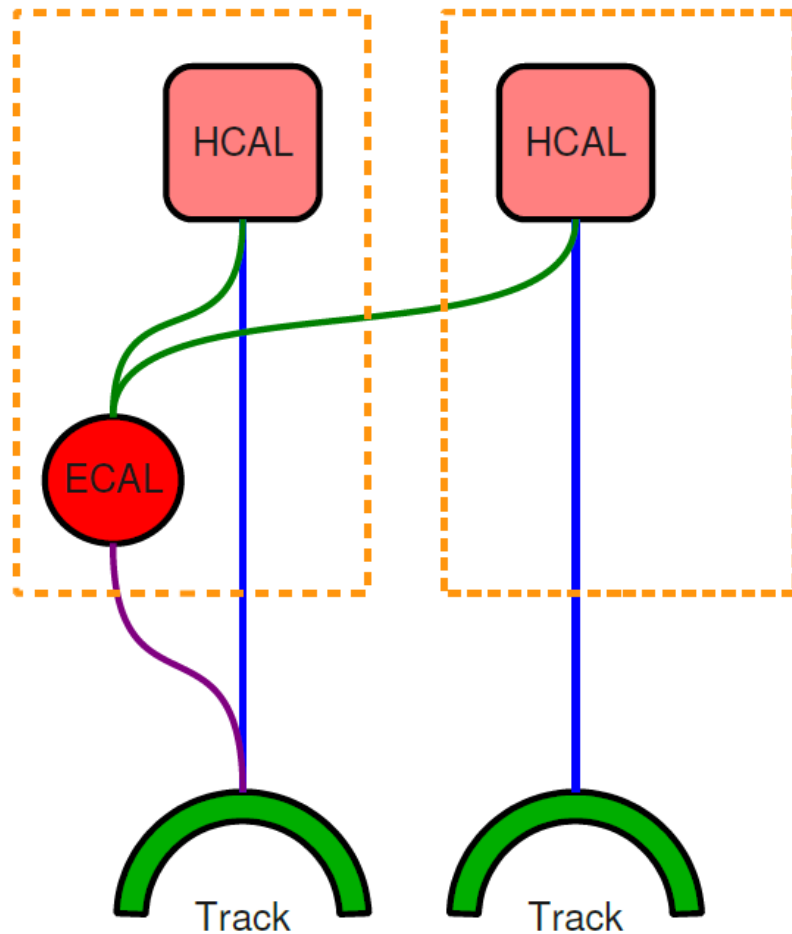
Charged hadrons + photon/neutral hadrons

If $E \ll p$

Something odd going on... Needs attention (doesn't happen often)

Back to 2005 : Checking CMS basic promises (15)

- ... and determine their energies



If p and E are compatible

Fit of p_i and E according to σ_{E,p_i}

Charged hadrons : p_i at small $p_{T,i}$

Converges to $p = E$ at high E

If E is in significant excess of p

Charged hadrons : p_i

If E is from HCAL or ECAL only

HCAL : Neutral hadron ($E-p$)

ECAL : Photon ($E_{ECAL}-p$)

If E is from ECAL and HCAL

If $(E-p) > E_{ECAL}$

Photon(E_{ECAL}) + neutral hadron

otherwise

Photon : $(E-p)/b$

(Always give precedence to photons in ECAL)

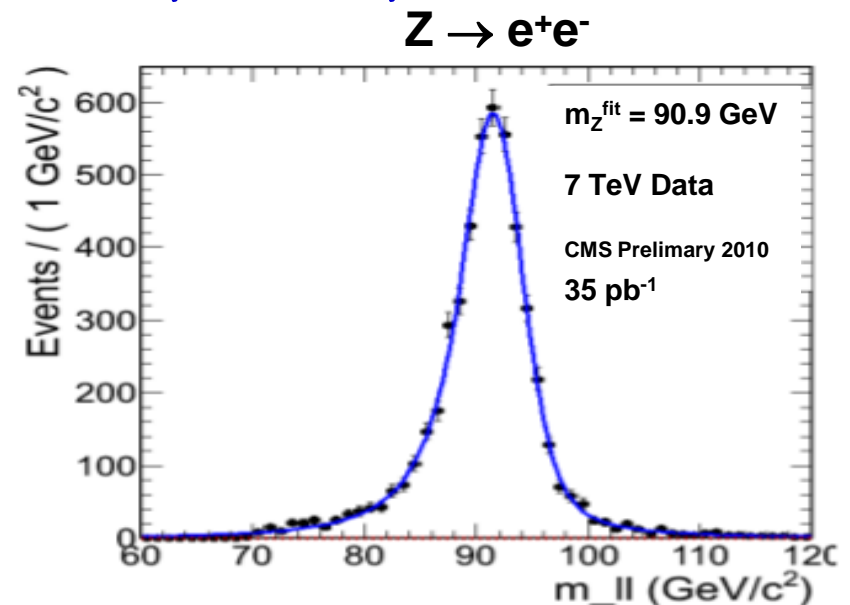
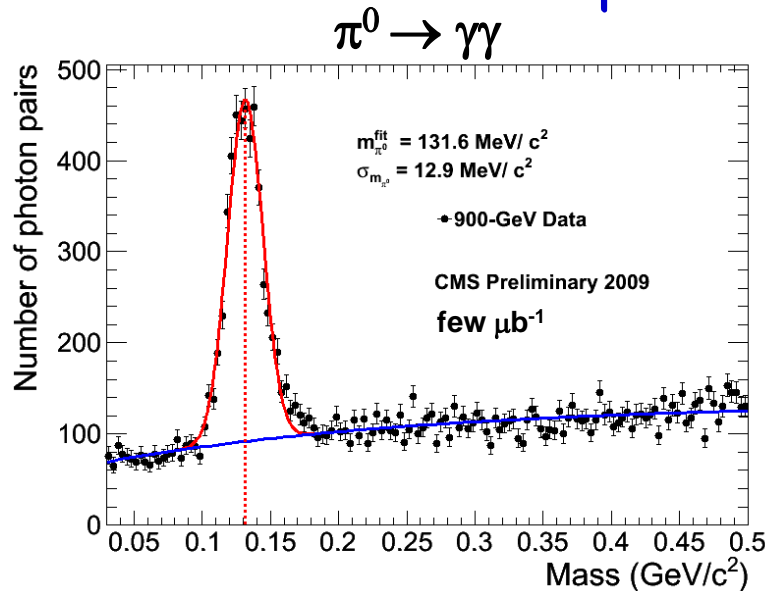
Back to 2005 : Checking CMS basic promises (16)

- In our famous first simulated jet ever :
 - ◆ Four particles generated : π^+ , π^- , π^0 and K^0_L
 - ◆ Five particles reconstructed :
 - Two oppositely-charged hadrons (π^+ and π^-)
 - Three photons
 - Two from the π^0 decay and one from the K^0_L energy deposit
 - No neutral hadron
 - Because for each of the two tracks, E was compatible with p
 - ◆ Note : the precedence given to photon identification in ECAL
 - May underestimate the ECAL energy deposits of neutral hadrons
 - But the neutral hadron energy deposited in ECAL corresponds to
10% neutral hadron x 30% ECAL fraction = 3% of event energy
May lose <0.5% of the event energy from this identification choice
(Was the same in ALEPH)

A few subtleties : the devil is in the details (I)

□ Calibration of ECAL and HCAL energies for charged hadrons

- ◆ ECAL is calibrated for photons (and electrons, see later), not for hadrons

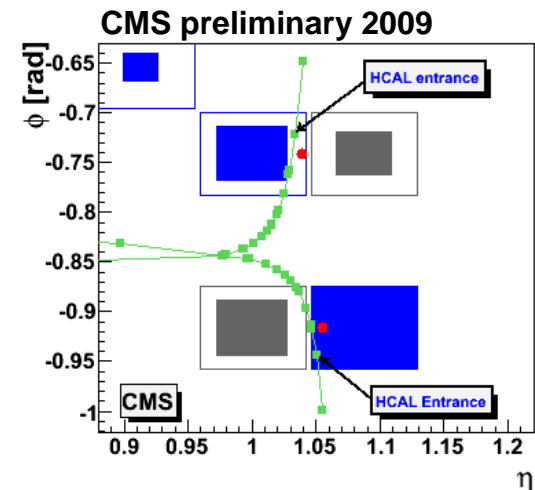
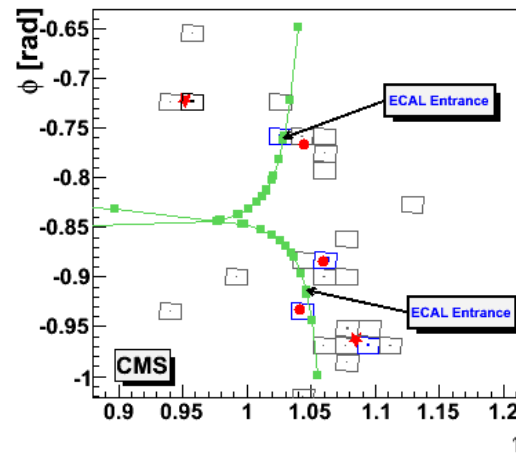


- ◆ HCAL is calibrated for 50 GeV charged pions at normal incidence
 - Test-beam calibration done without ECAL/Services in front of HCAL
- ◆ Hence, when a charged hadron (p) interacts with the calorimeters
 - $E_{\text{ECAL}} + E_{\text{HCAL}}$ does not equal p (in general significantly smaller)

A few subtleties : the devil is in the details (2)

- Calibration of ECAL and HCAL energies for charged hadrons (cont'd)
 - ◆ To optimize merged neutral hadron identification, need to calibrate $E_{\text{ECAL,HCAL}}$ as
 - $E = a + b(p, \eta) E_{\text{ECAL}} + c(p, \eta) E_{\text{HCAL}}$
 - Compensates for ECAL response and also for HCAL nonlinearities (!)
 - ◆ Charged hadrons and photons (90% event energy) insensitive to this calibration
 - Only the neutral hadron identification efficiency is
 - Calorimeter calibration for hadrons is a second order effect for PF
 - ◆ Use "isolated" tracks in minimum bias events for a , b and c determination

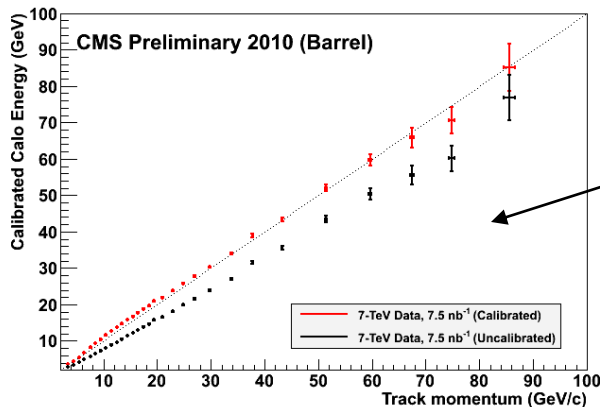
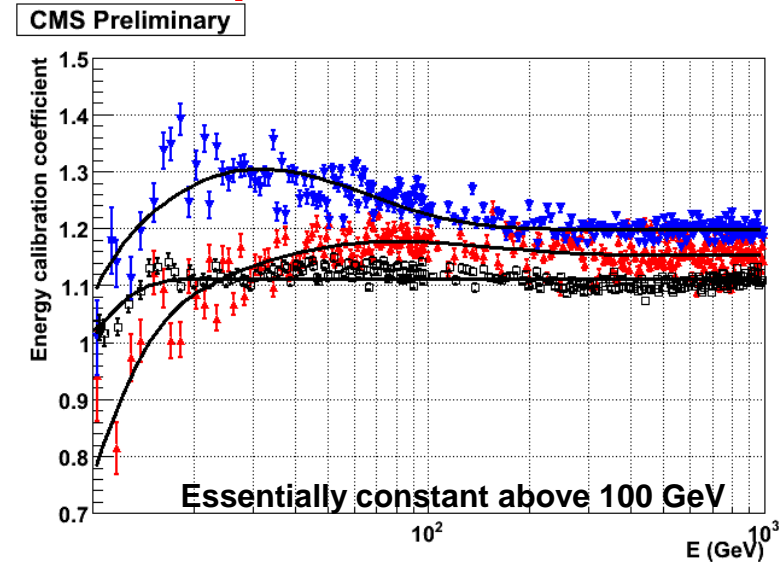
- Granularity Helps !
- One HCAL in the block
 - One Track in the block
 - High-quality fit
 - Fit a , b , c
- As function of p
- As function of η



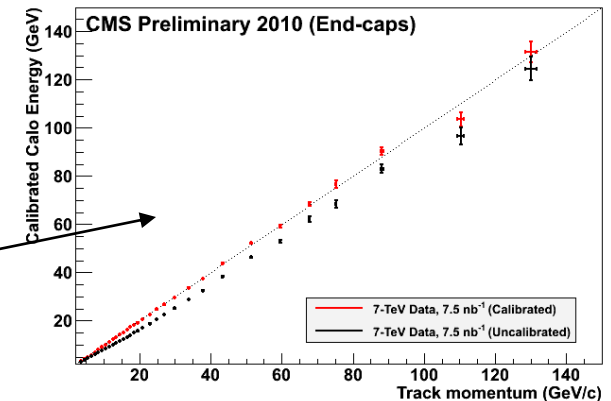
A few subtleties : the devil is in the details (3)

□ Calibration of ECAL and HCAL energies for charged hadrons (cont'd)

- ◆ Coefficients obtained from simulation
 - Blue : $b(E,0)$
 - Black : $c(E,0)$ / HCAL only
 - Red : $c(E,0)$ / ECAL + HCAL
 - a = threshold correction (couple GeV)
 - ◆ Can be applied directly on data
 - But can be /are obtained from data
- Seem fine all the way to 150 GeV

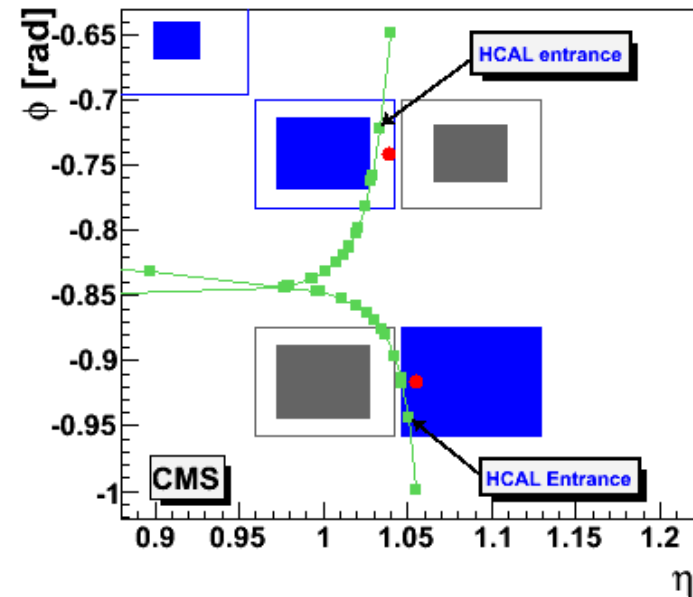
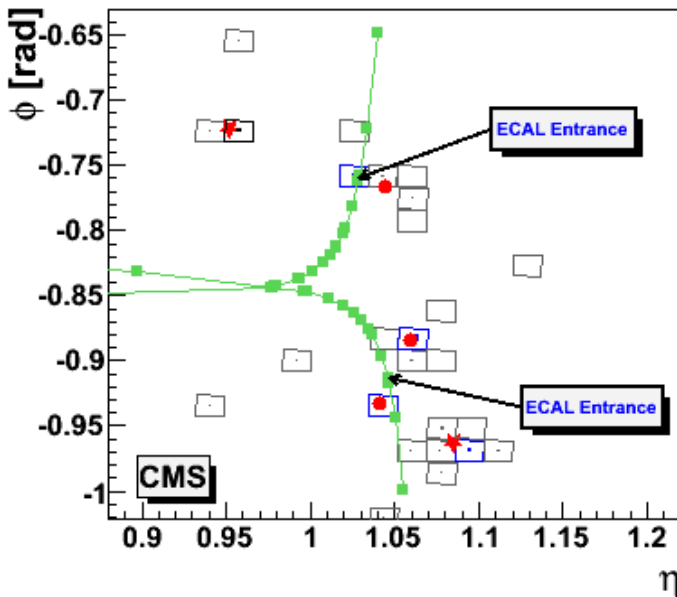


Only 7.5 nb⁻¹ data
 350K tracks in barrel
 1.7M tracks in end-caps



A few subtleties : the devil is in the details (4)

□ Calibration of ECAL and HCAL energies for charged hadrons (cont'd)



◆ Track pointing "downwards"

- $p = 14.64 \text{ GeV}/c$, $E_{\text{ECAL}} = 1.87 \text{ GeV}$, $E_{\text{HCAL}} = 7.35 \text{ GeV}$, $E_{\text{CALIB}} = 14.33 \text{ GeV}$

◆ Track pointing "upwards"

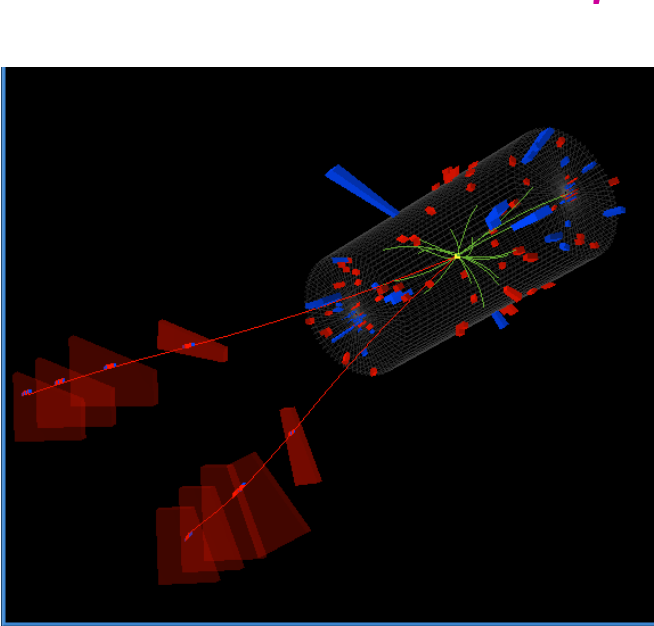
- $p = 10.94 \text{ GeV}/c$, $E_{\text{ECAL}} = 0.98 \text{ GeV}$, $E_{\text{HCAL}} = 6.77 \text{ GeV}$, $E_{\text{CALIB}} = 9.19 \text{ GeV}$

◆ Gives two charged hadrons of 14.64 GeV and 10.94 GeV in the particle list.

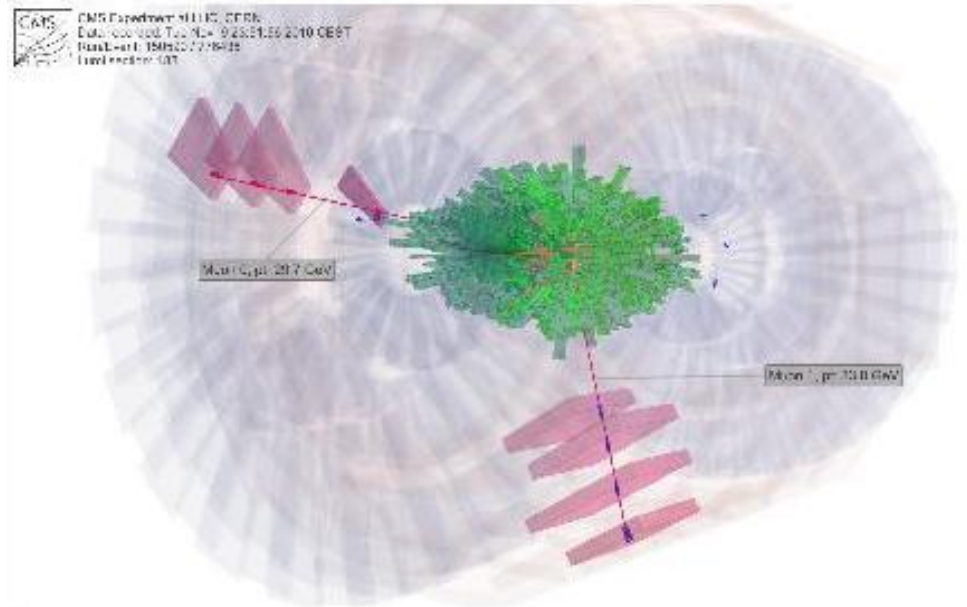
A few subtleties : the devil is in the details (5)

□ Cases $E \ll p$

- ◆ Arise from muons, in the majority of the cases
 - Muon reconstruction and identification is rather easy in CMS
 - Even in a very busy environment



$J/\psi \rightarrow \mu\mu$ (pp)



$Z \rightarrow \mu\mu$ (PbPb)

A few subtleties : the devil is in the details (6)

□ Cases with $E \ll p$: Muons (cont'd)

- ◆ Muon reconstruction/identification requires
 - A high-quality track in the silicon tracker
 - A high-quality track in the muon system
 - A global fit with a good χ^2
- ◆ Typically very efficient (95% in the muon system acceptance) and 99.5% pure
 - As in ALEPH, these muon tracks are removed from the block
 - Before the charged hadrons, photons, neutral hadron treatment
- ◆ The remaining 5% lead to $E \ll p$ cases
 - A much looser muon identification is used in this configuration
 - Isolated tracks (not likely to be a hadron)
 - Tracker track only (plus a few hits in the muon system)
 - Poor-quality tracker track (but good fit in the muon system)
 - Typical particle-flow attitude: use all detectors to improve particle ID/Reco

A few subtleties : the devil is in the details (7)

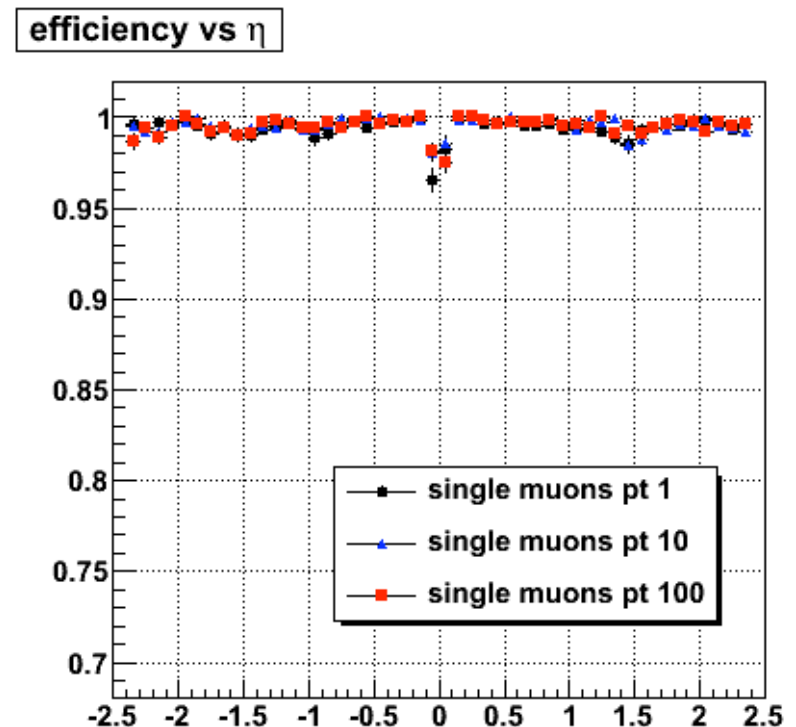
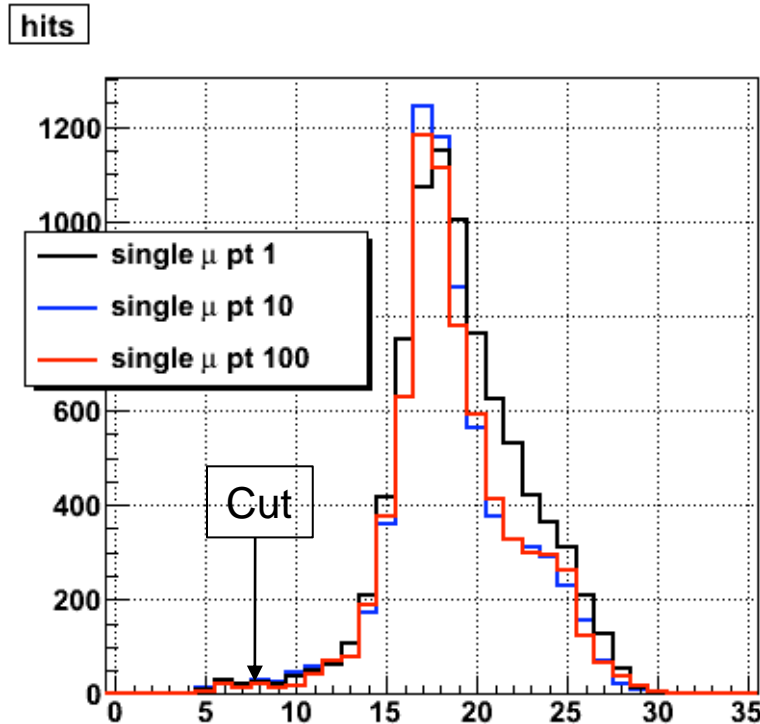
- Cases where $E \ll p$: Fake tracks, p mis-measurements
 - ◆ Despite the tight selection for charge particle tracks (reminder)
 - Seeded with at least two hits in the pixel detector
 - Originating from the beam axis within tight tolerances
 - Pattern recognition with a combinatorial Kalmann-Filter track finder
 - Request at least 8 hits and $p_T > 0.9 \text{ GeV}/c$
 - Leading to only 85% efficiency for primary charged pions (see just after)
 - ◆ The 3D-ish granularity of the tracker leads to ambiguities
 - Hits wrongly associated to tracks
 - Confusion in the pattern recognition
 - Leading to fake tracks with a 1-2% rate
 - Leading to incorrect momentum determination
 - Typically create high-momentum tracks
 - The problem increases with the charge multiplicity, i.e., with jet p_T

A few subtleties : the devil is in the details (8)

- Cases where $E \ll p$: Fake tracks, p mis-measurements (cont'd)
 - ◆ The problem of fake tracks was solved in the following way
 - Reject tracks obviously fake (bad quality fit, missing hits along the track)
 - For each block, rank the remaining tracks according to $\sigma_{\text{fit}}(p_T)$
 - In the very few blocks with more than one track (few %)
 - For the very few tracks with $\sigma_{\text{fit}}(p_T) > 1 \text{ GeV}$ (0.1%)
 - Remove the worse tracks until E compatible with p
 - The “last” track is simply rescaled to make $p = E$
 - 0.001 % of the tracks concerned by this procedure
 - Note : for these (small) blocks, the energy is given by the sole calorimeters
 - But photon identification (and energy determination) still holds
 - Individual charged particle multiplicity and momenta are still available
 - Momentum direction mostly from tracking and ECAL
 - Typical particle-flow attitude : use the redundancy of all measurements.

A few subtleties : the devil is in the details (9)

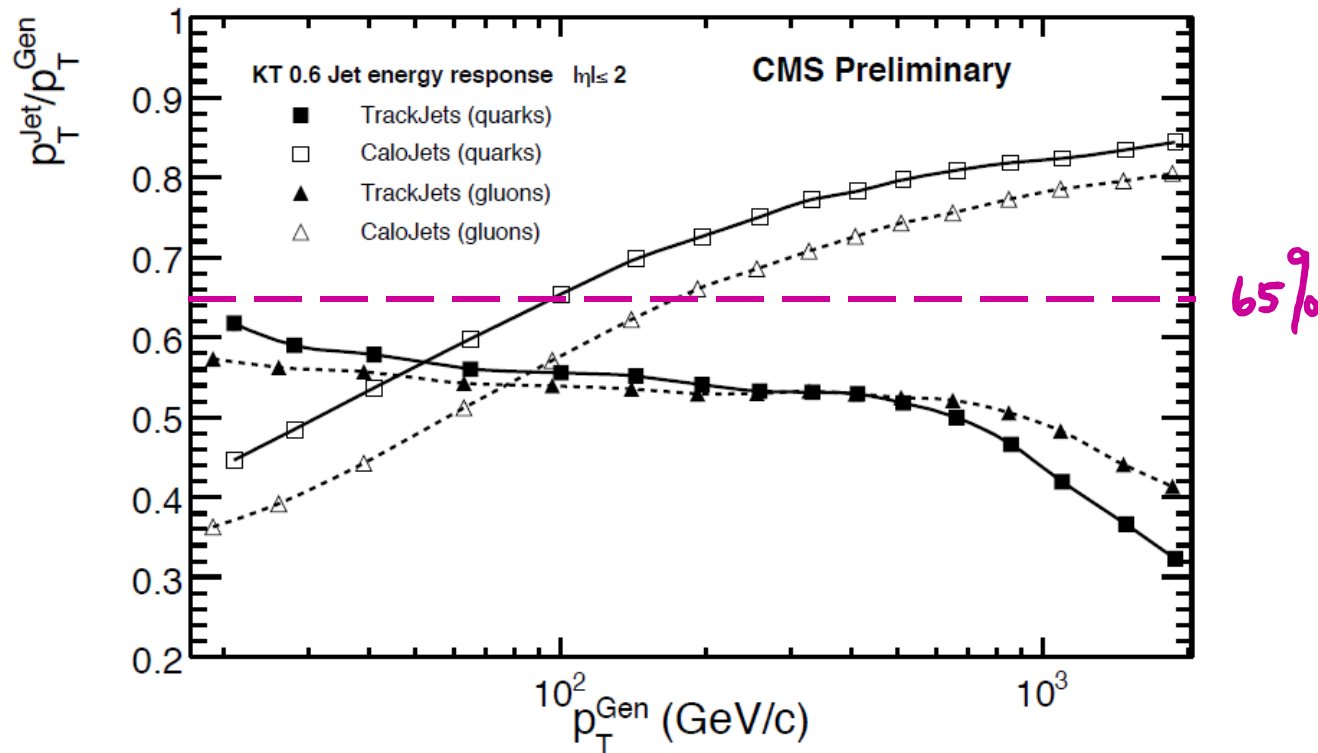
- A tracking efficiency of 85% : why ?
 - ◆ Silicon detectors are very efficient in finding hits (>99%)
 - As seen with simulated single muons (confirmed with data)



- Only slight drops of efficiency in the pixel detector cracks and overlaps

A few subtleties : the devil is in the details (10)

- A tracking efficiency of 85% : why ? (cont'd)
 - ◆ Question : Is it due to large occupancy in jets ?
 - Hence large overlap between tracks ?



- ◆ Answer : not for p_T up to 800 GeV/c, where this effect starts to show up

A few subtleties : the devil is in the details (II)

□ A tracking efficiency of 85% : why ? (cont'd)

◆ Remember : About $2X_0$, or $0.4 \lambda_0$, of material in the tracker

• About 20% of the hadrons interact in the tracker material

→ Sometimes in a spectacular manner

A single K_L^0 (270 GeV) interacting in the tracker after 15 cm:

Blue : true particles

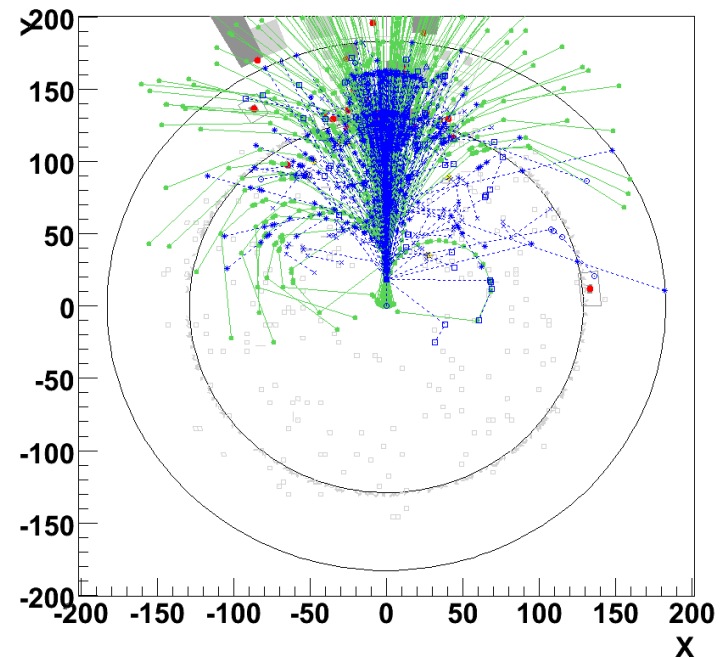
Green : reconstructed tracks

(well, not yet)

→ what happens to these tracks ?

Not enough hits on the primary

Bad origin for the secondaries

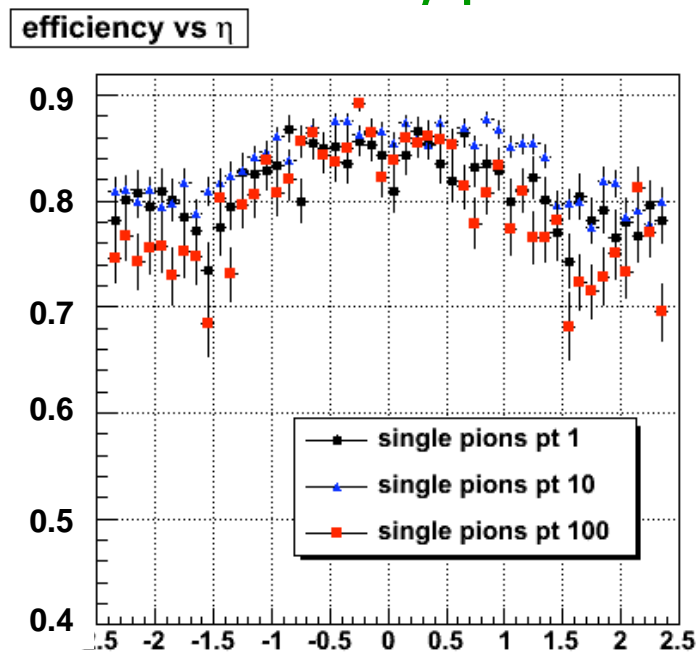
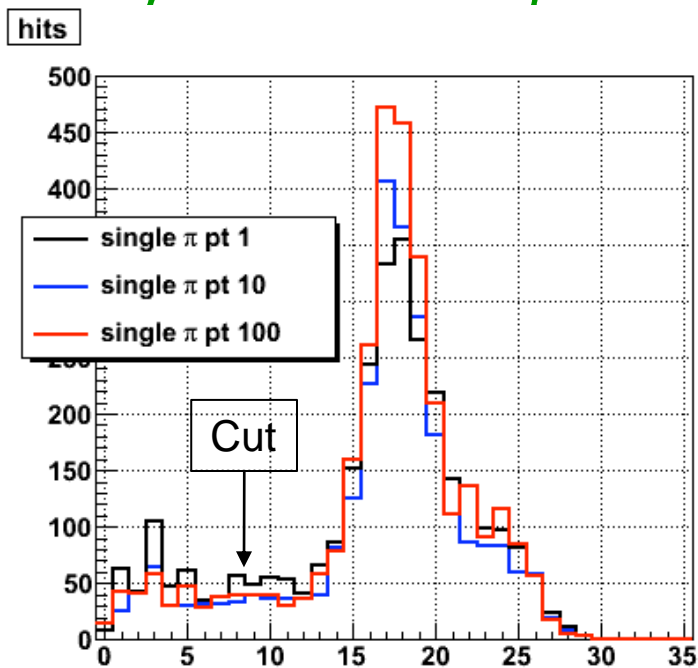


A few subtleties : the devil is in the details (12)

□ A tracking efficiency of 85% : why ? (cont'd)

◆ Indeed, for single pions :

- Large loss of efficiency due to short tracks from interacting pions



- With a $n_{\text{Hit}} > 5$ cut, the fake rate jumps to about 20%, not manageable.

◆ Can one increase the tracking efficiency without increasing the fake rate ?

A few subtleties : the devil is in the details (13)

□ Increasing the tracking efficiency : Iterative tracking

- ◆ Fake tracks come from wrong combinatorial hit association
 - Reduce the number of hits fed to the combinatorial track finder ?
- ◆ 1. Start from a very pure seeding
 - e.g., 3 pixel hits, very tight origin constraint, $p_T > 0.9 \text{ GeV}/c$
→ 75% efficiency, less than 1% fake rate
- ◆ 2. Reconstruct the corresponding tracks (≥ 3 hits) and “remove” the hits used
 - 40% of the hits in the tracker are removed in this first iteration
- ◆ 3. With the 60% remaining hits, try a looser seeding
 - e.g., 2 pixel hits, very tight origin constraint, $p_T > 0.9 \text{ GeV}/c$
→ Adds 15% efficiency, but still less than 1% fake rate
Because the combinatorial possibilities are much less
- ◆ 4. Reconstruct the corresponding tracks (≥ 3 hits) and “remove” the hits used
 - 10% of the hits are further removed in this second iteration

A few subtleties : the devil is in the details (14)

□ Increasing the tracking efficiency : Iterative tracking (cont'd)

◆ And so on with more iterations :

- Third iteration : 3 pixel hits, tight origin constraint, $p_T > 0.2 \text{ GeV}/c$
- Fourth iteration : 2 pixel hits, looser origin constraint, $p_T > 0.3 \text{ GeV}/c$

◆ And even try to catch secondary tracks (interactions, conversions, decays ...)

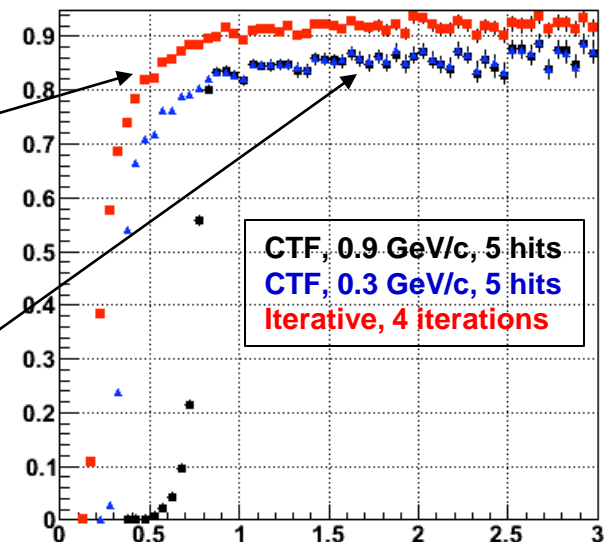
- Fifth : TIB/TID seeding, loose origin constraint, $p_T > 0.5 \text{ GeV}/c$
- Sixth : TOB/TEC seeding, very loose origin constraint, $p_T > 0.8 \text{ GeV}/c$

◆ After 4 iterations

- 93% efficiency, 1-2% fake rate
- Down to very low momentum

85% efficiency, 20% fake rate

efficiency vs p_T



A few subtleties : the devil is in the details (14)

Increasing the tracking efficiency : Iterative tracking (cont'd)

◆ And so on

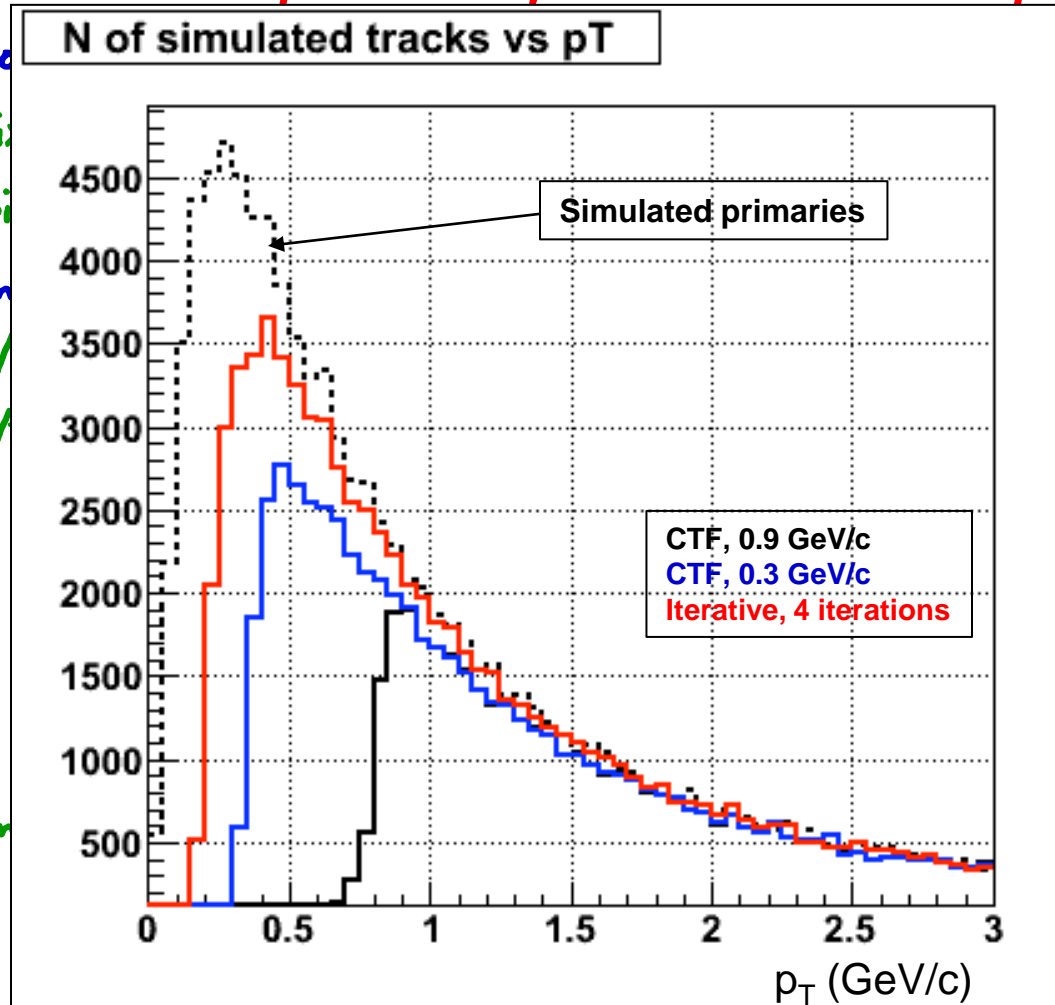
- 3 pi
- 2 pi

◆ And even

- TIB/
- TOB/

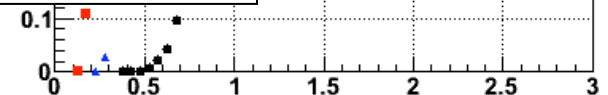
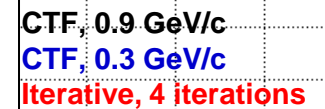
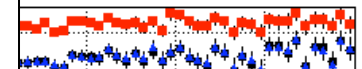
◆ After 4

- 93%
- Down



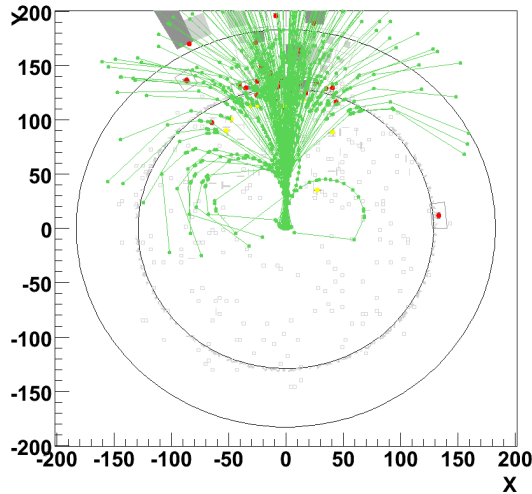
ons, decays ...)

GeV/c



A few subtleties : the devil is in the details (15)

- Increasing the tracking efficiency : Iterative tracking (cont'd)
 - ◆ Nuclear interactions, photon conversions, decays in flight :



Many tracks reconstructed
with TIB/TID and TOB/TEC seeding
Small fake rate dealt with by PF protections

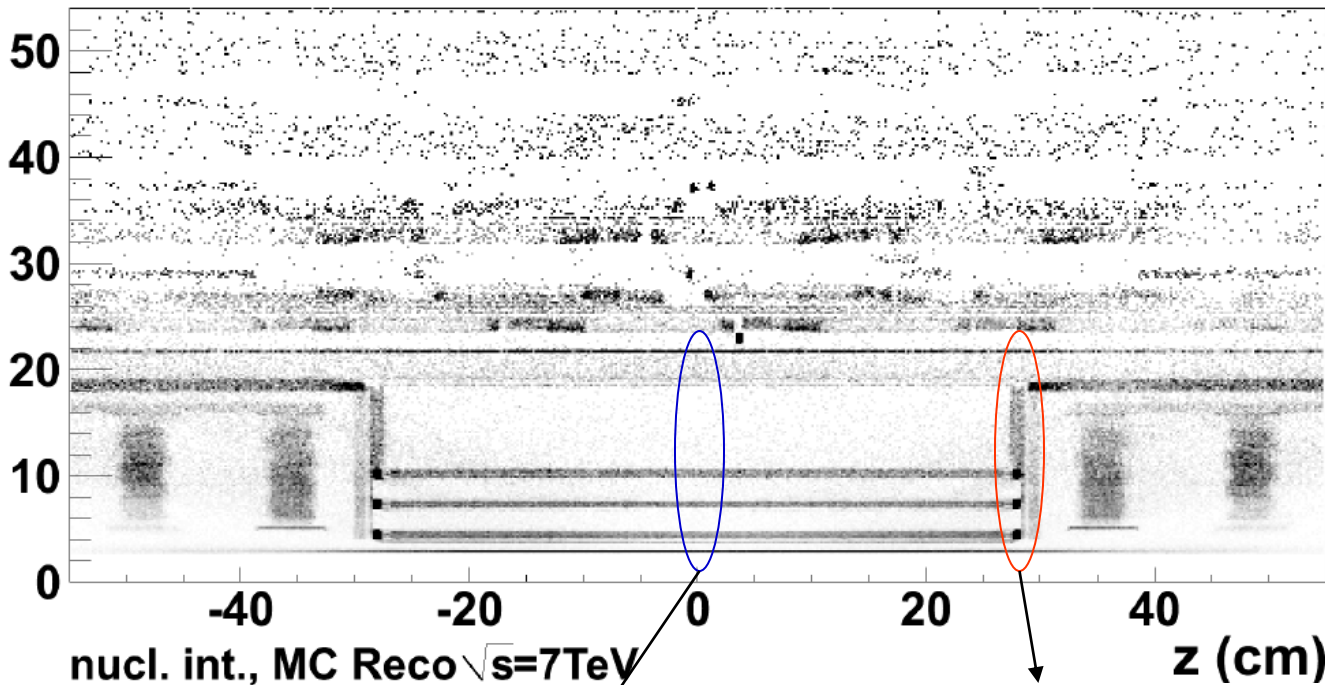
- ◆ Next problem : avoid double counting from primary vs secondaries
 - Create a “link by vertex” between primary and secondaries
 - And chose the best energy determination
 - Most likely the primary if more than 5-6 hits, secondaries otherwise
- ◆ Typical particle-flow attitude : reconstruct/identify as many particles as possible !

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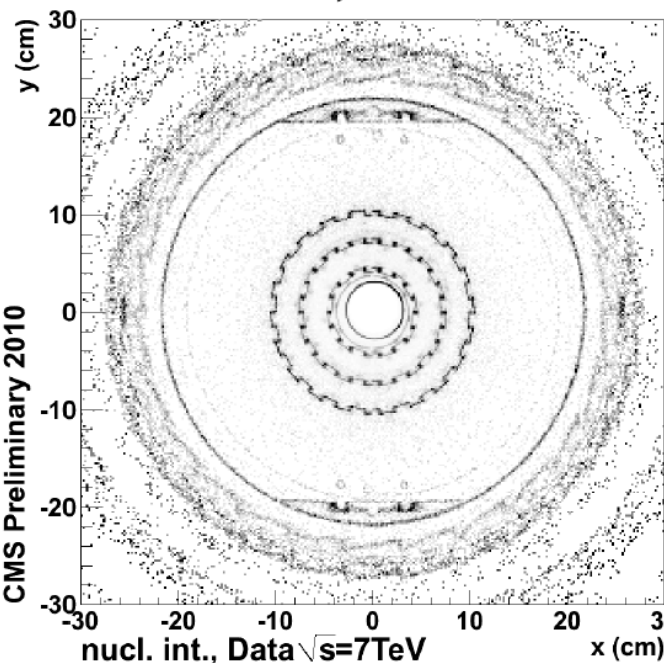
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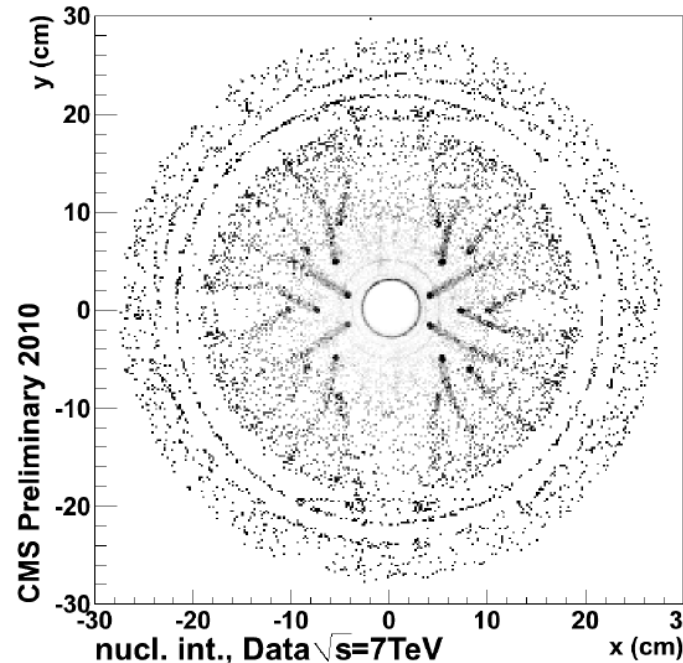
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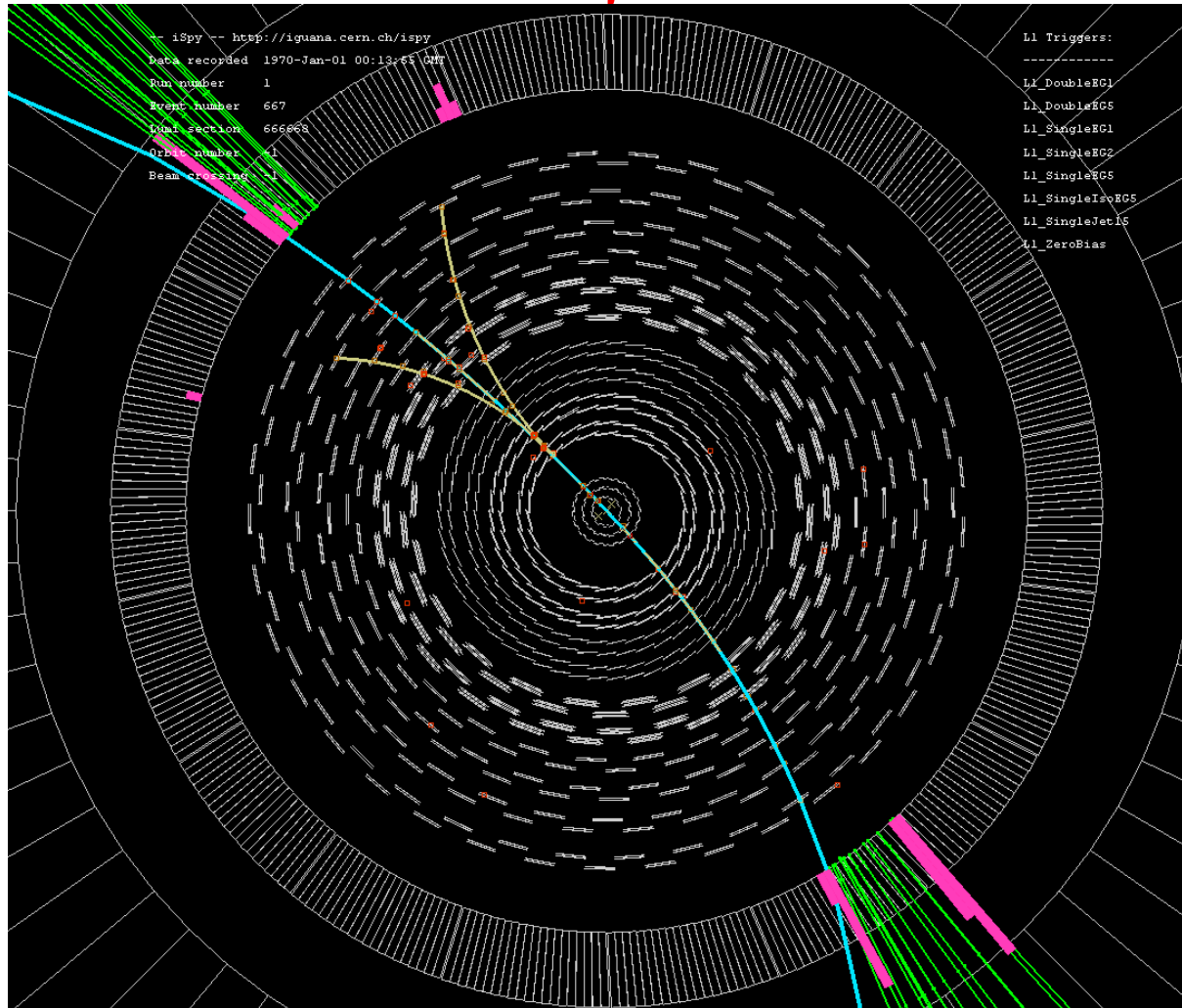
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A few subtleties : the devil is in the details (16)

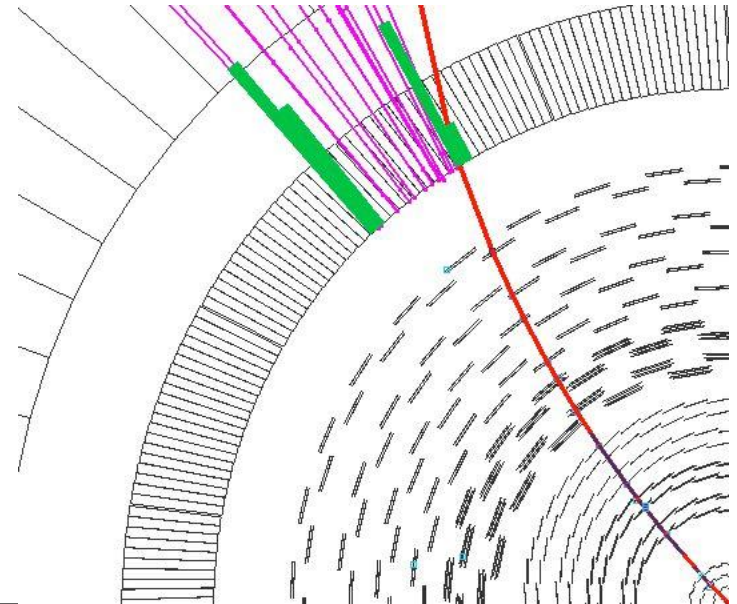
- And what about electrons ? They radiate, and the brem γ 's convert !



A few subtleties : the devil is in the details (π)

□ Electrons : Tracking

- ◆ Because they radiate, many electrons would have failed the $N_{\text{hits}} > 8$ cut
 - The Kalman filter pattern recognition quickly gives up
- ◆ The iterative tracking was initially meant at solving this issue for PF
 - Tracks with at least 3 hits are used as seed
 - Use a Gaussian-Sum filter to follow the electron track all the way to ECAL
- ◆ Issue : GSF tracking is slow
 - Use it only for pre-identified tracks
 - Small number of hits
 - Or : Poor quality fit
 - Or : p/E_{ECAL} not far from unity
 - Linked to pre-shower hits
 - Concerns only 5% of the tracks
 - with 95% efficiency on electrons



A few subtleties : the devil is in the details (18)

- **Electrons : Recovering the Bremsstrahlung photons**
 - ◆ If nothing is done, radiated photon energy is counted twice
 - Once from the electron initial momentum, p_{IN}
 - Twice from the energy from corresponding ECAL cluster(s), E_{BREM}
 - ◆ Create a new type of track - ECAL cluster link
 - The link "by tangent"

→ Purple lines are tangents to the GSF track

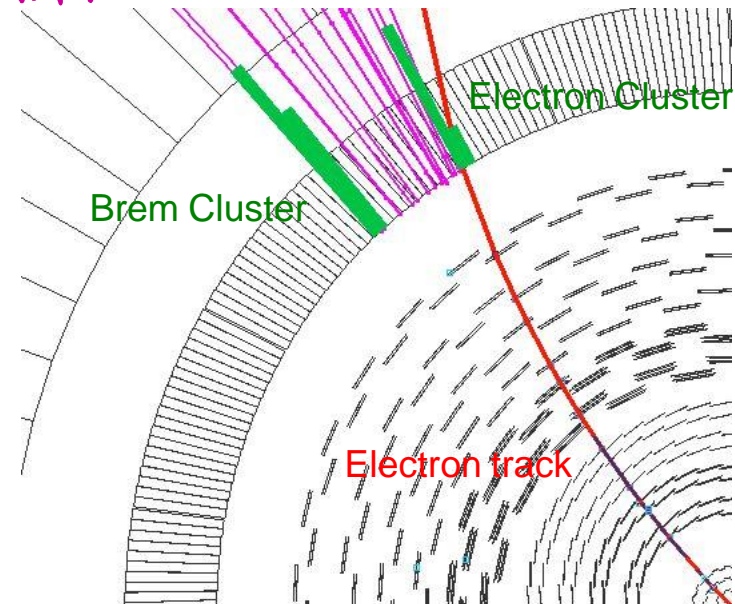
Starting from each tracker layer

→ If the tangent points to a cluster

Link the cluster to the track

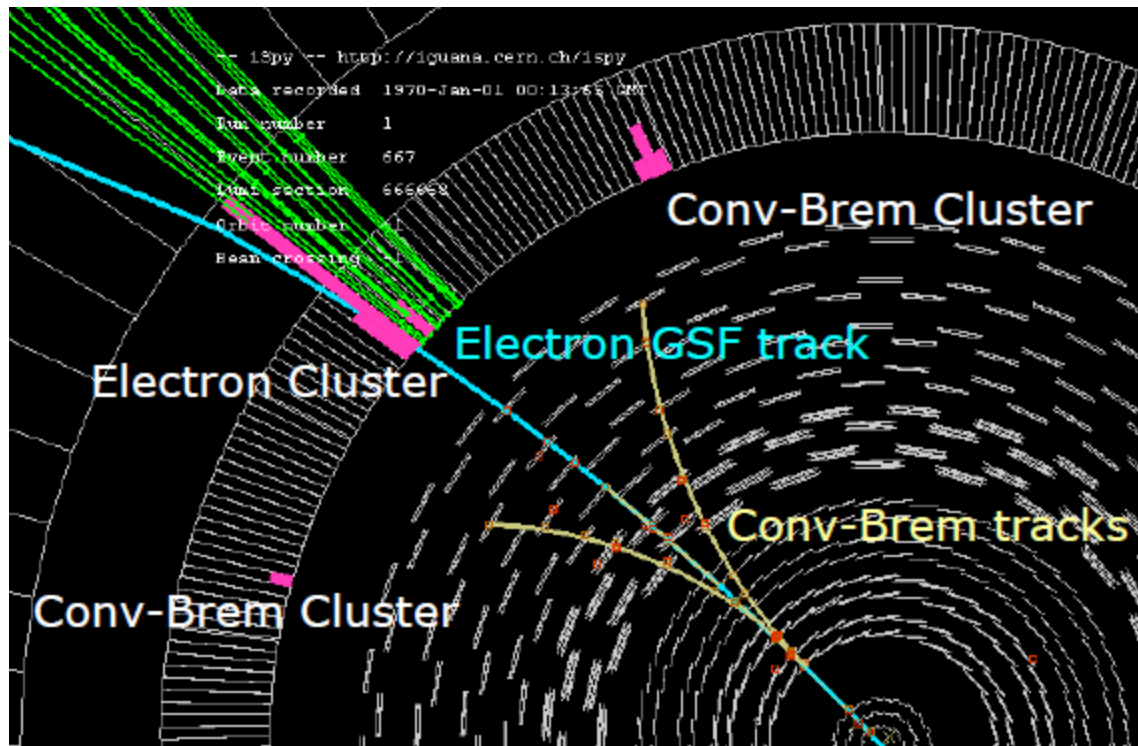
→ Another handle (not used) :

Compatibility between Ecluster
and ΔP along the GSF track



A few subtleties : the devil is in the details (19)

- **Electrons : Recovering the converted Bremsstrahlung photons**
 - ◆ Tracks recovered by the 5th and 6th step of the iterative tracking
 - Linked by vertex to the original electron tracks
 - Linked to ECAL clusters in a classical way
- And added to the block ...



A few subtleties : the devil is in the details (20)

□ Electrons : Identification

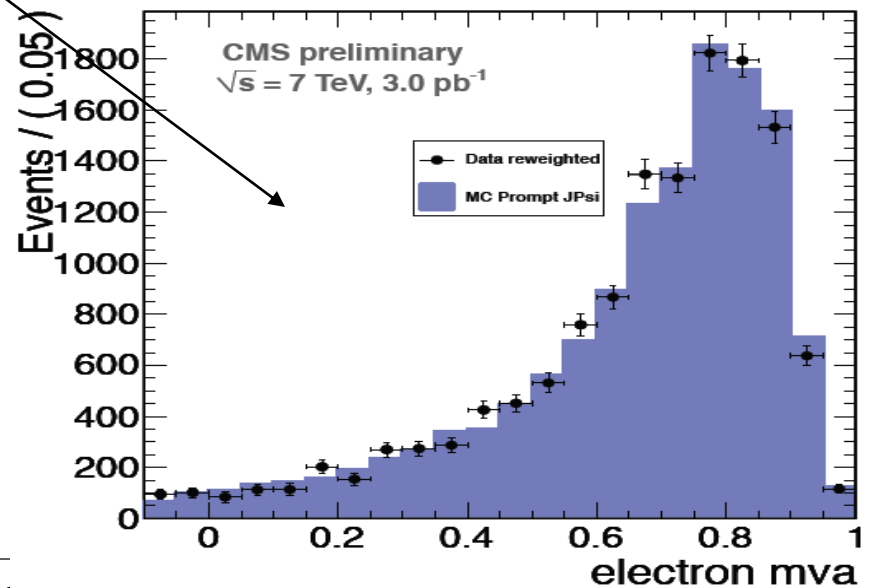
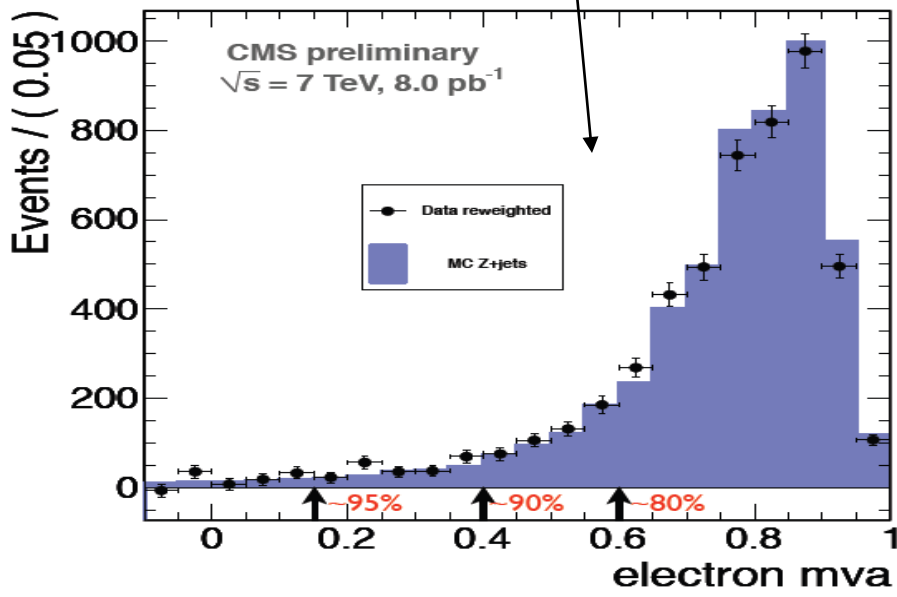
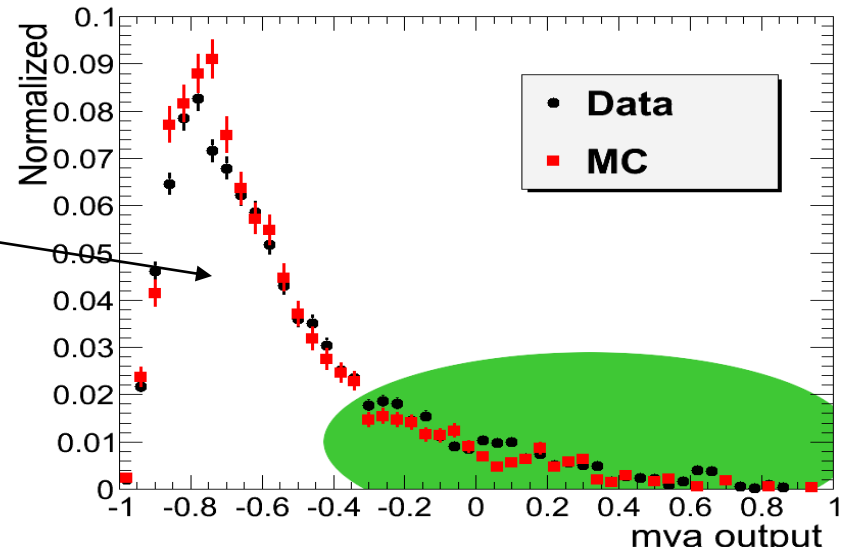
- ◆ Use the tracker as a preshower !
 - Number of hits of the KF tracks
 - Energy loss along the GSF track : $P_{IN} - P_{OUT}$
 - Number of Bremsstrahlung photons associated to the track
 - Comparison of E_{BREM} and $P_{IN} - P_{OUT}$
 - Comparison of $E_{ELECTRON} + E_{BREM}$ and P_{IN}
 - ...
- ◆ Plus some calorimeter-only quantities
 - Shower width along η
 - Linked HCAL energy
 - ...
- ◆ And combine in a boosted decision tree (could be any MVA tool)
 - 70-80% efficiency in jets, 95% for isolated electrons

A few subtleties : the devil is in the details (21)

Electrons : Identification, cont'd

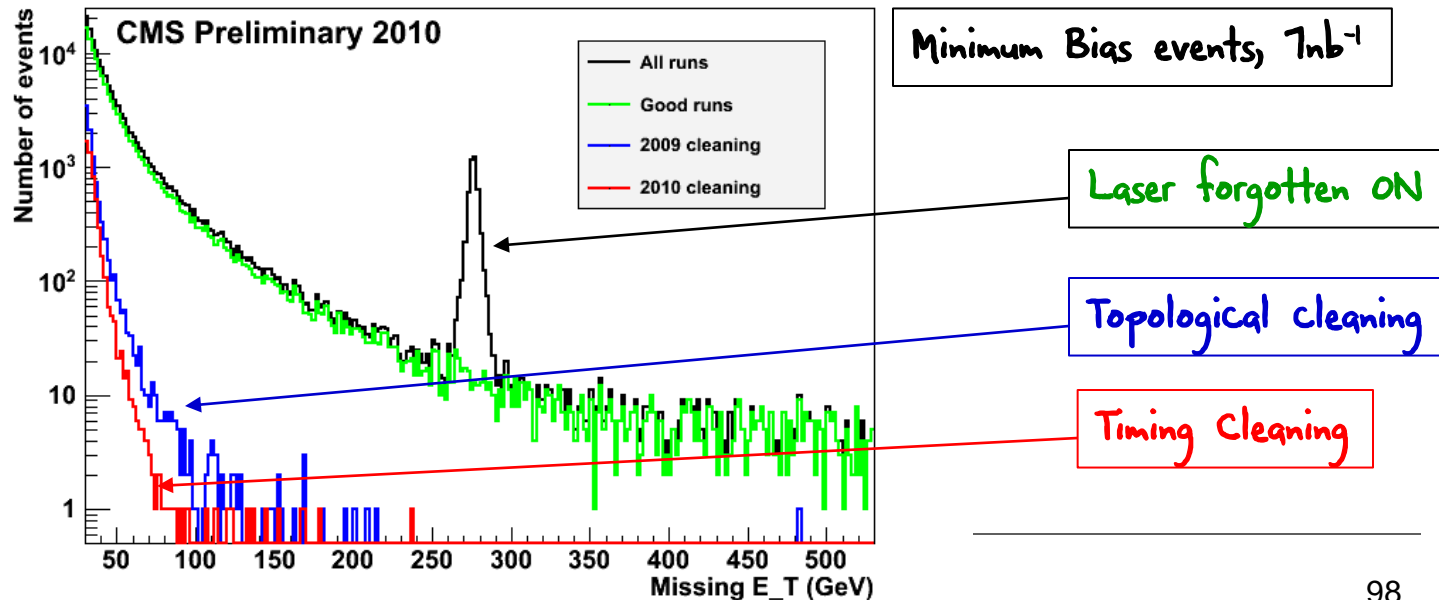
It works with data as in simulation !

- $K_S^0 \rightarrow \pi^+\pi^-$
- $J/\psi \rightarrow e^+e^-$
- $Z \rightarrow e^+e^-$



A few subtleties : the devil is in the details (22)

- The tracker material was the cause of most devilish details
 - ◆ The lack of redundancy was also a high price to pay for calorimeter cleaning
 - Needed to use information on timing, pulse shape, isolation, ... to get rid of
 - ECAL spikes (due to slow neutrons hitting APD's)
 - HF spikes (due to Cerenkov light from hadron-shower muons in PMT's)
 - HCAL spikes (due to ion feedback in HPD's)
 - ◆ But efficient cleaning could be achieved (so far) - without over-cleaning.



Particle-Flow performance in CMS (I)

- We now have a list of particles to work with
 - ◆ Charged hadrons, photons, neutral hadrons, (electrons, muons) :

- An event at $\sqrt{s}=2.36$ TeV

Full lines : charged hadrons

Dashed lines : photons

Dotted lines : neutral hadrons

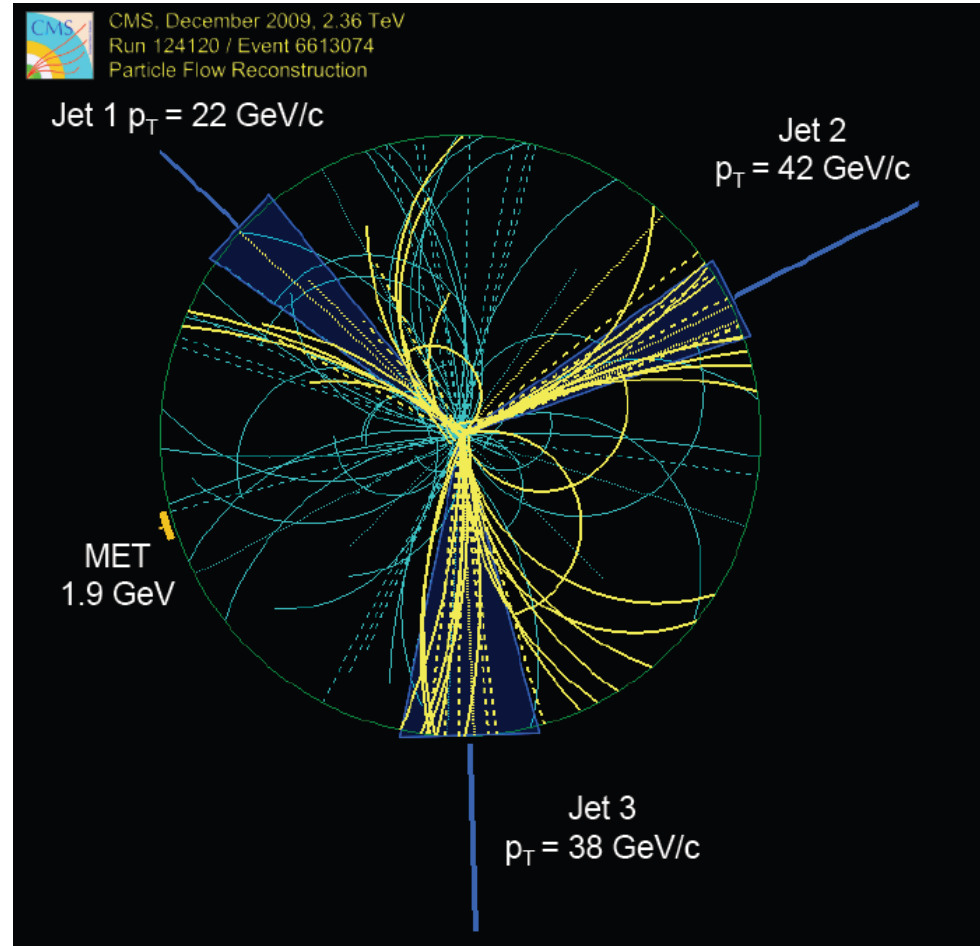
- Jets with $p_T > 20$ GeV/c

Yellow lines : Jet constituents

Blue areas : Jet "Cone"

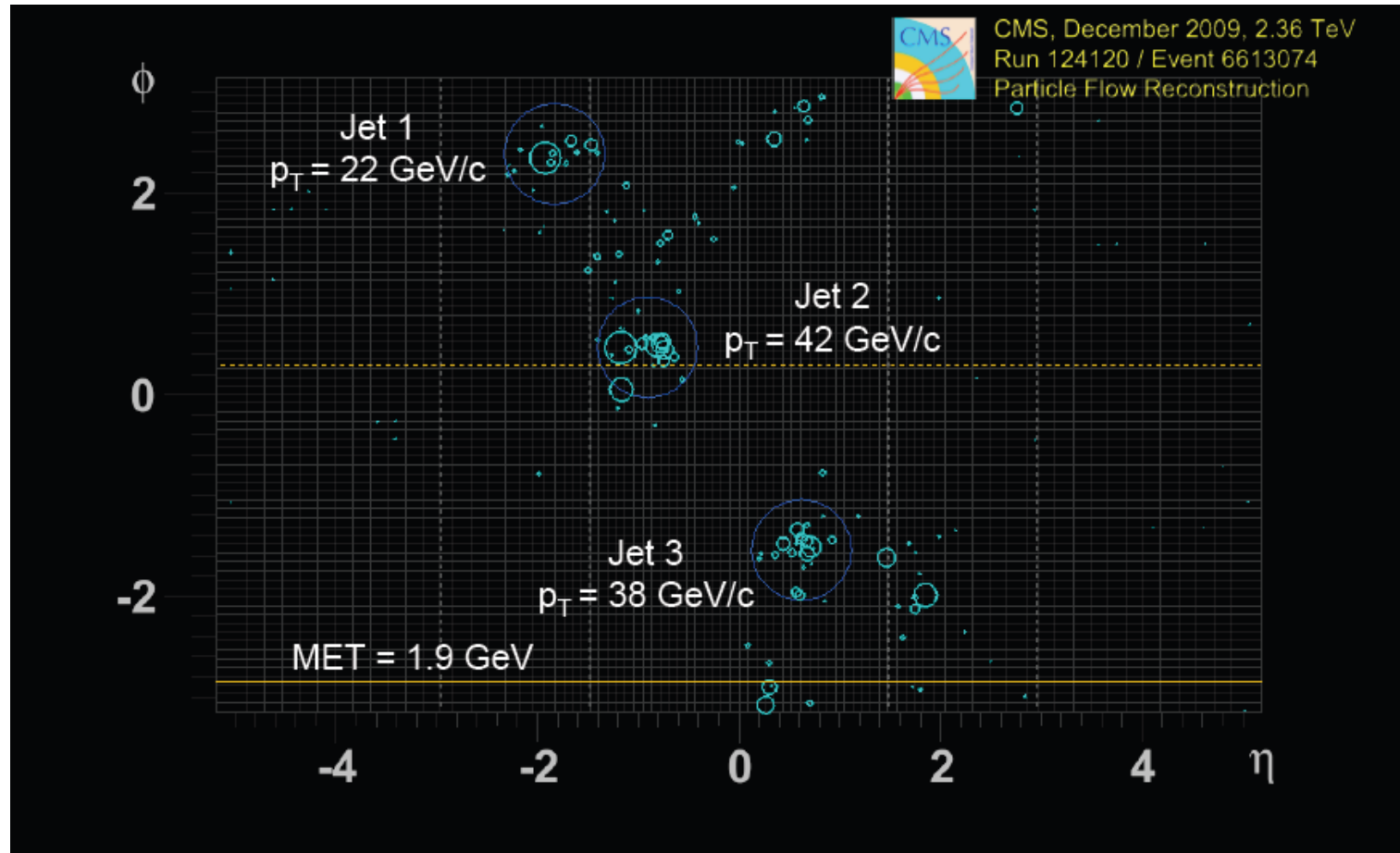
- Tastes like generated particles

But reconstructed particles, really



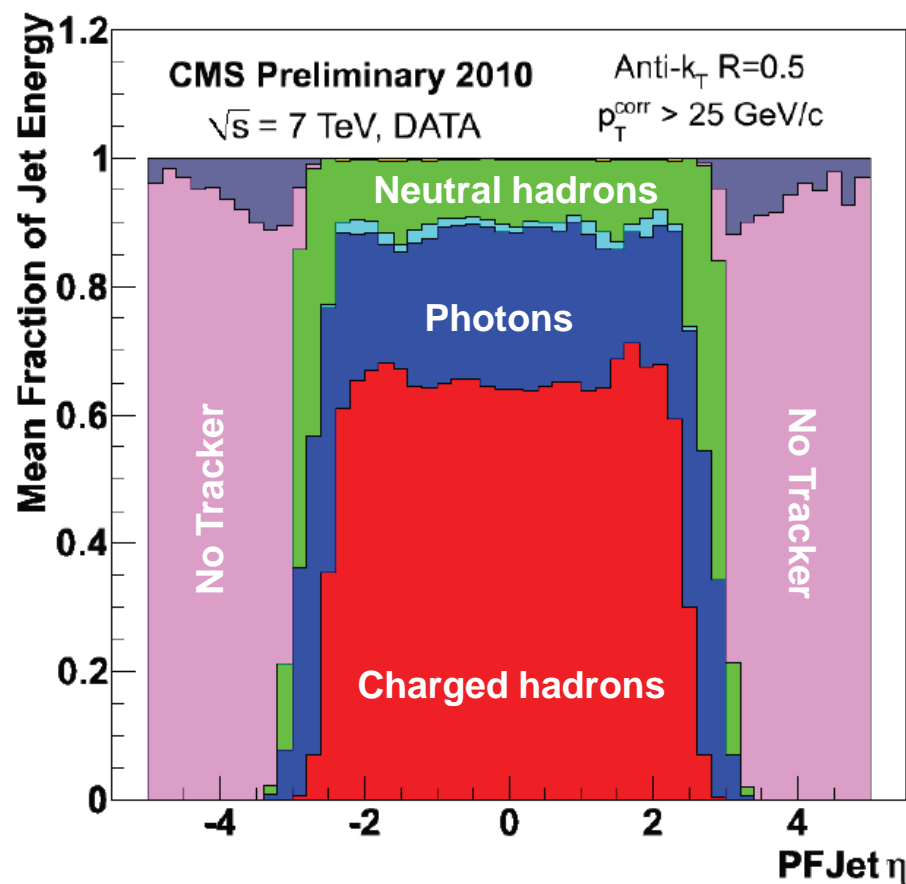
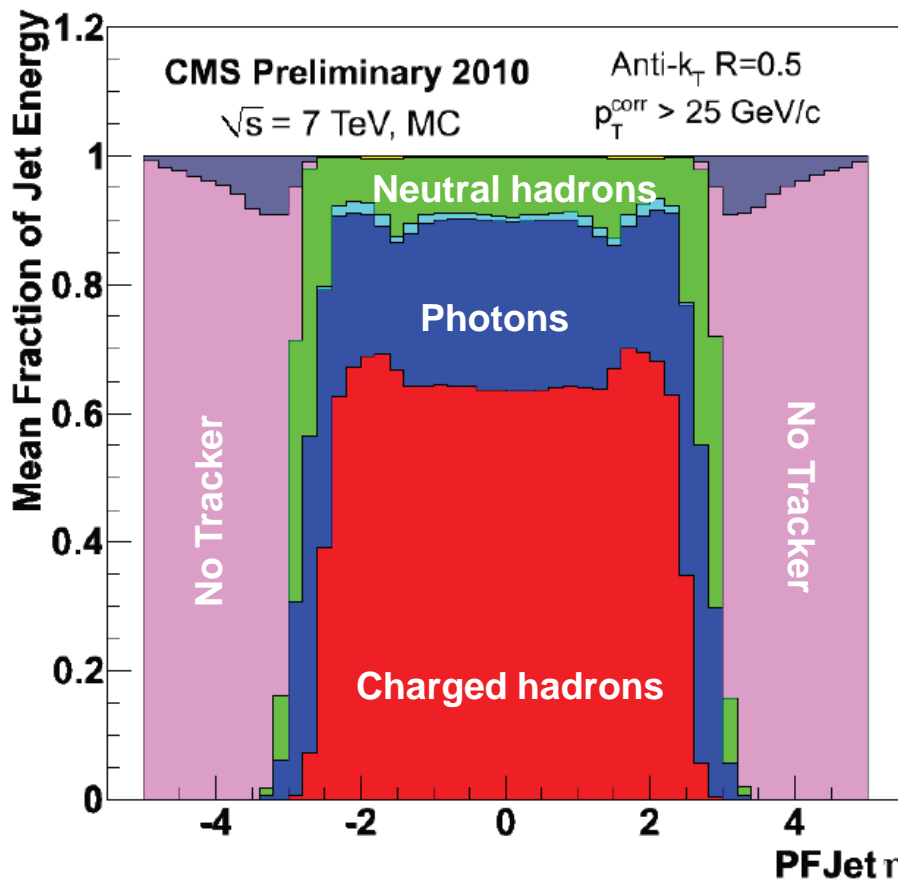
Particle-Flow performance in CMS (2)

- We now have a list of particles to work with (cont'd)
 - ◆ with optimal particle-to-particle granularity + collimated jets at the lowest p_T 's



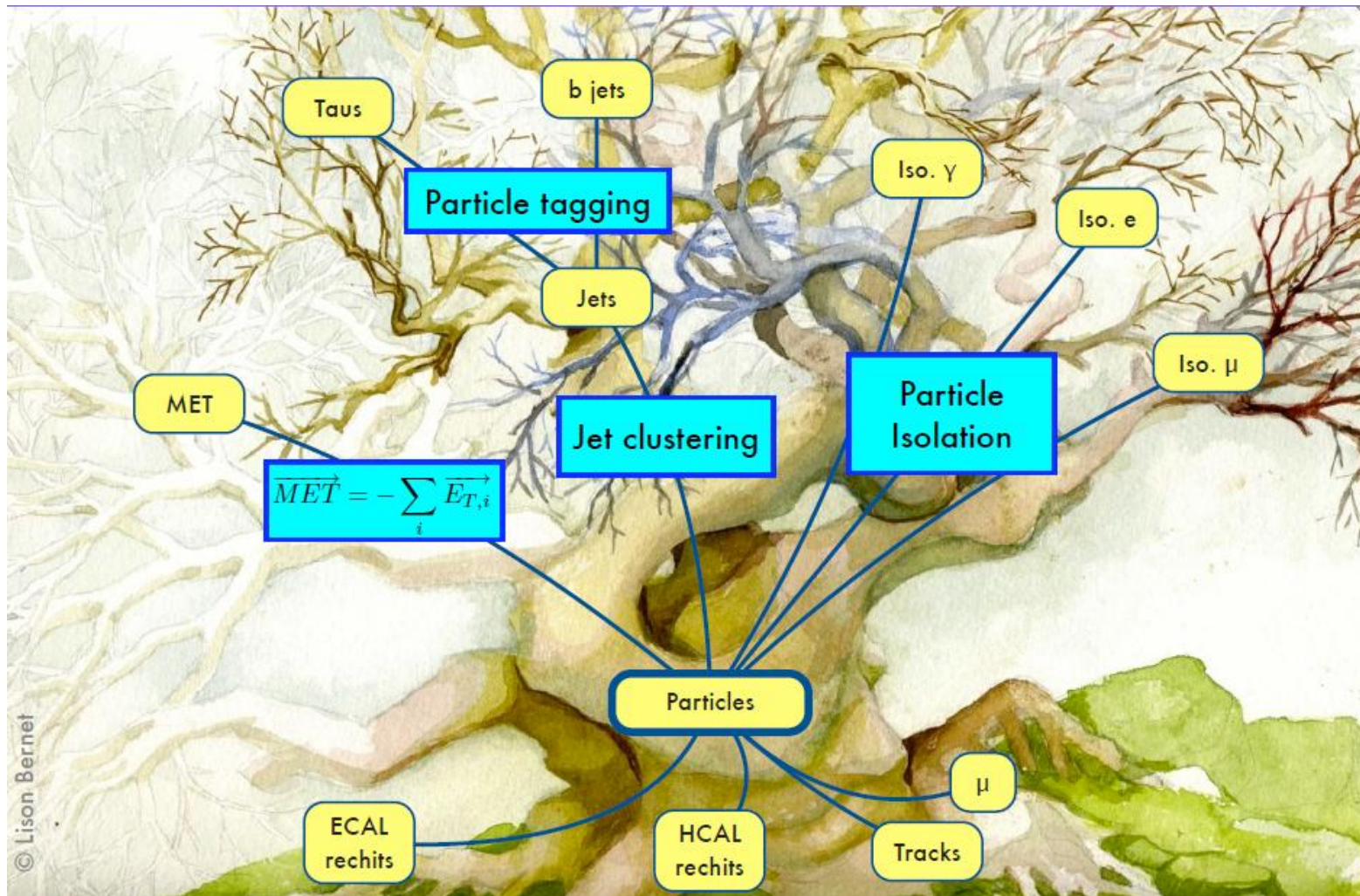
Particle-Flow performance in CMS (3)

- We now have a list of particles to work with (cont'd)
 - ◆ with the expected fractions in jets (65% h^+ , 25% g and 10% h^0)
 - what is data, what is simulation ?



Particle-Flow performance in CMS (4)

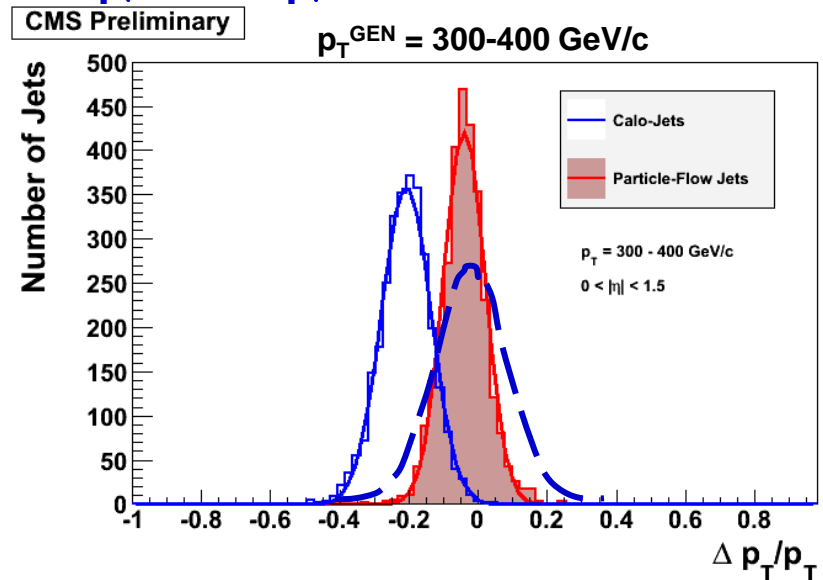
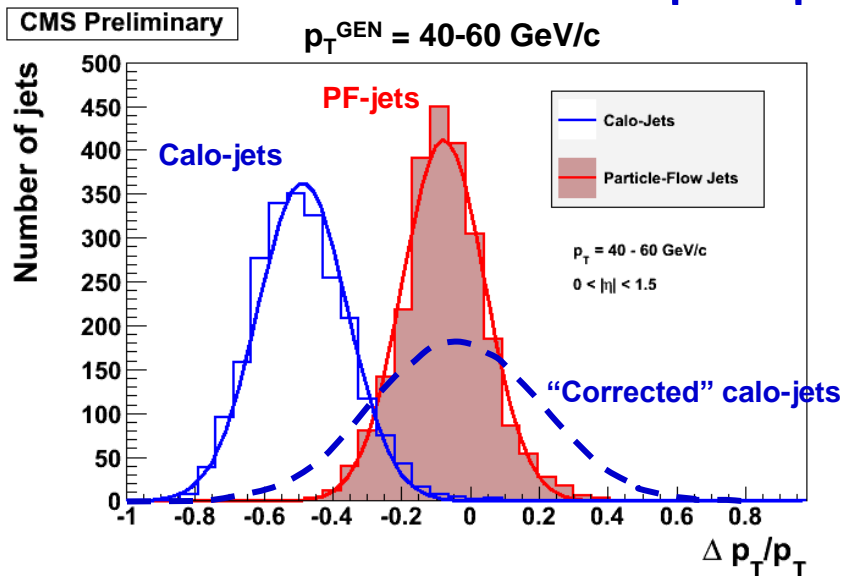
- Physics objects from the global event description with particles



Particle-Flow performance in CMS (5)

Expected performance for jets

Reconstruction of the jet p_T : $(p_T^{\text{RECO}} - p_T^{\text{GEN}}) / p_T^{\text{GEN}}$



Comparison with calorimetric jets

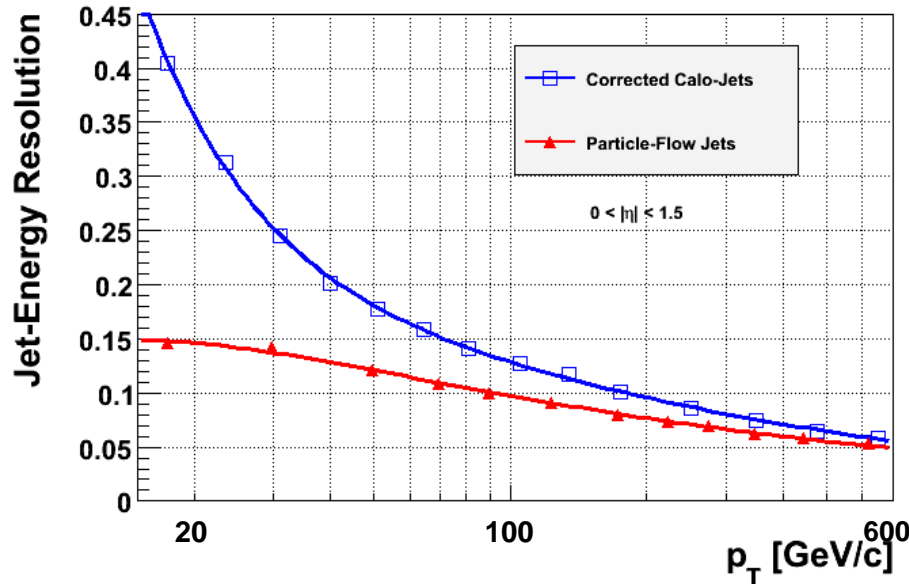
- Response larger than 95% of the original jet energy
- Almost no need for a posteriori corrections (systematic uncertainties!)
- Much better energy resolution
- Similar Gaussian behaviour despite the large number of sub-detectors used

Particle-Flow performance in CMS (6)

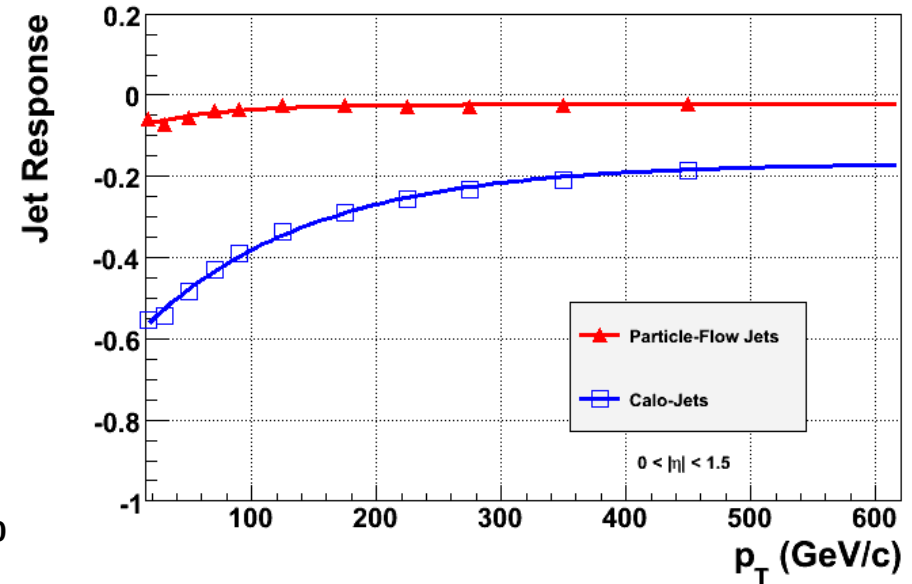
Expected performance for jets (cont'd)

Resolution and response as a function of the jet p_T :

CMS Preliminary



CMS Preliminary



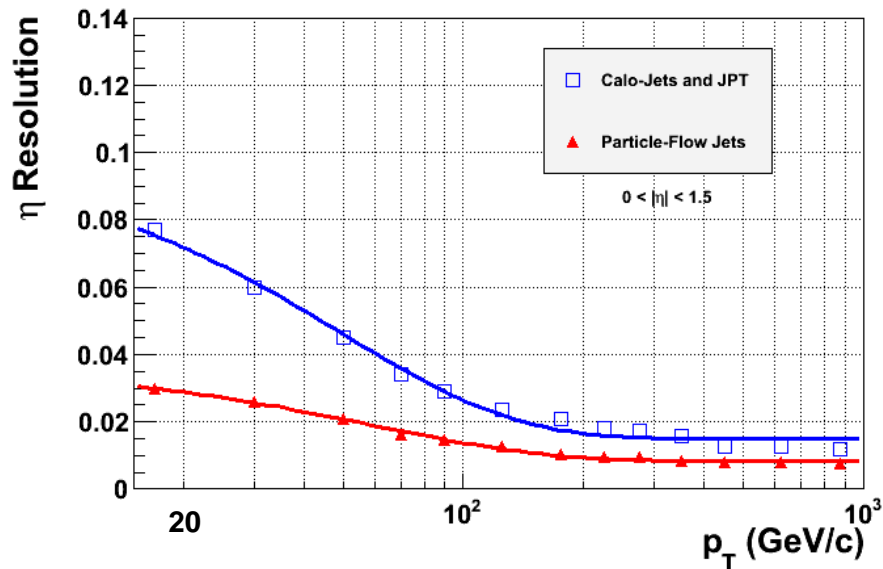
- Better resolution even at large p_T 's, where calorimeters are well behaved
 - Linear response, between 95% and 97% of the true jet energy
- No need for large, non-linear, corrections, down to $p_T = 10$ GeV/c
(Calo nonlinearities : Magnetic field, thresholds, photons, muons, ...)

Particle-Flow performance in CMS (7)

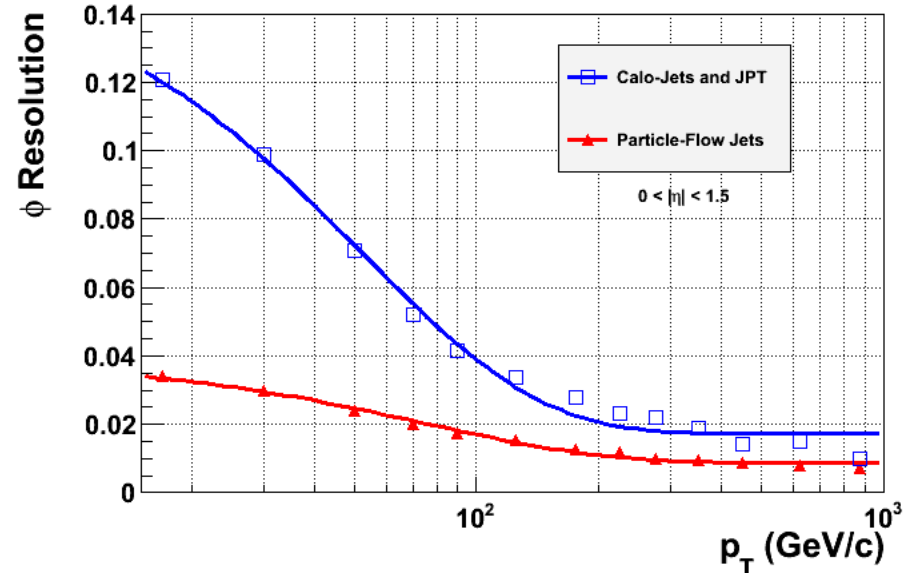
Expected performance for jets (cont'd)

Angular resolution

CMS Preliminary



CMS Preliminary

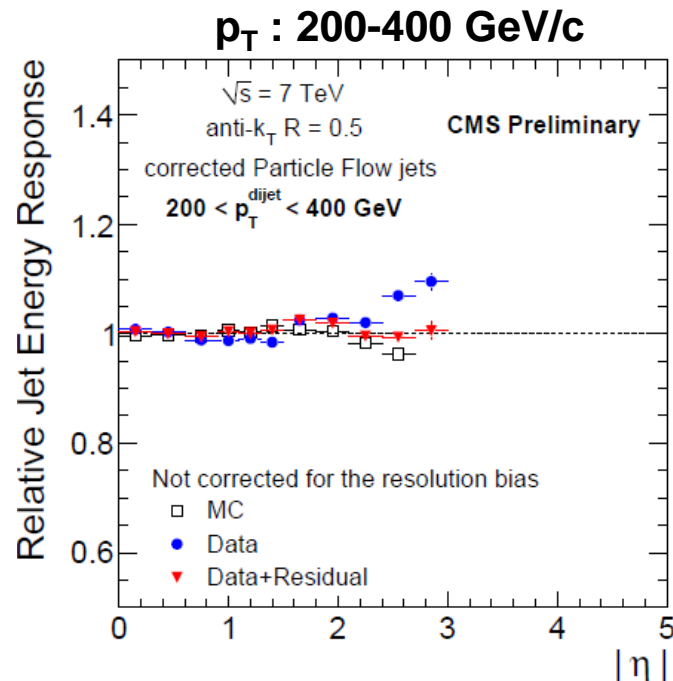
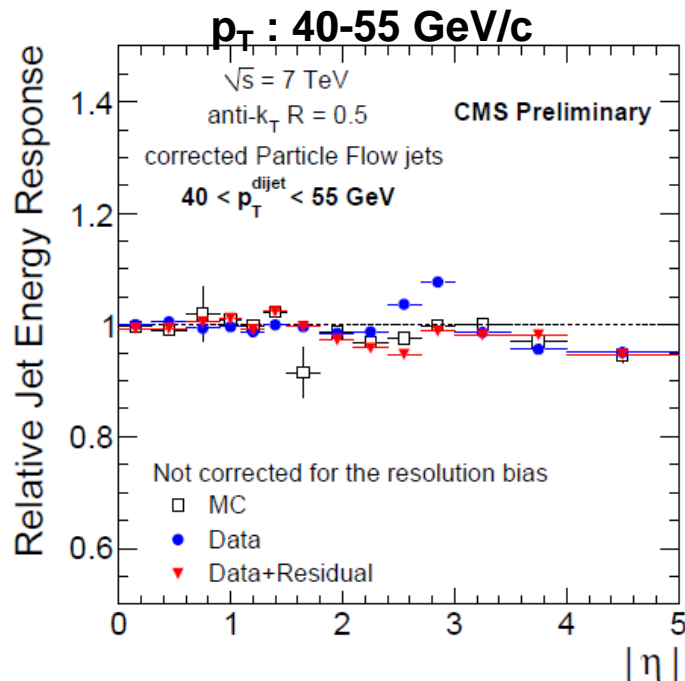


- A factor of 2.0 to 4.0 improvement in ϕ (B Field degrades calo jets)
 - A factor of 1.5 to 3.0 improvement in η (tracker and ECAL granularity)
- Large improvements expected in jet-jet mass resolution
with both improvements in jet energies and jet directions

Particle-Flow performance in CMS (8)

Measured jet response in data

- From γ jet events (as in ALEPH) : use the γ p_T to predict the jet p_T



- Impressive agreement between data and simulation
- Slight HCAL end-cap over-calibration, visible only beyond tracker acceptance
→ Being calibrated away for 2011 data taking

Particle-Flow performance in CMS (9)

□ Systematic uncertainty on the jet energy scale (JES)

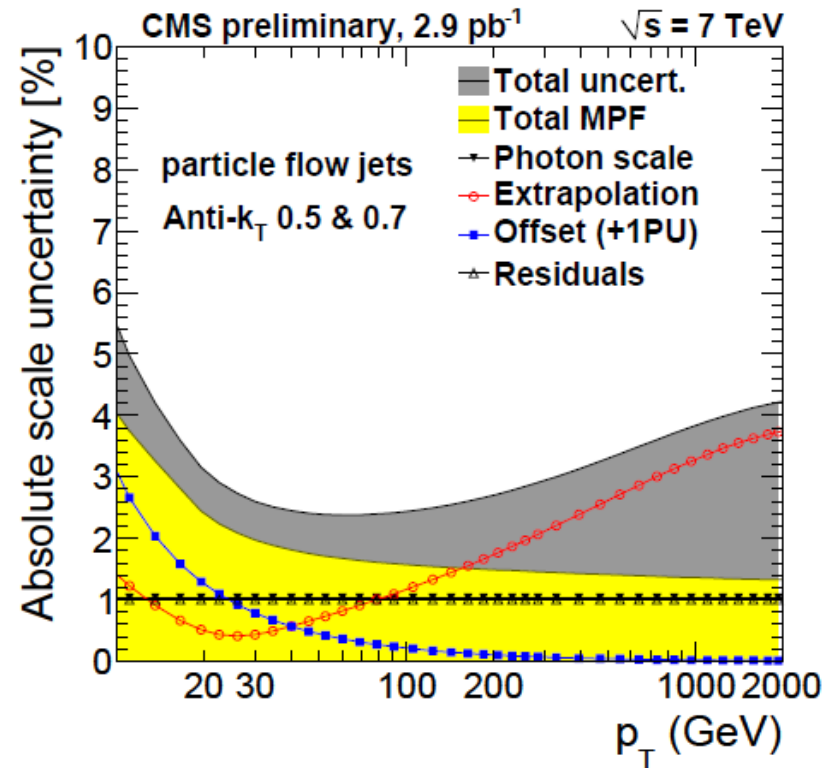
- ◆ From data : ECAL scale Known to 1% (π^0 , Z). HCAL scale Known to 5%
 - with 25% γ and 10% h^0 , the ultimate JES uncertainty is about 1%

◆ with γ jet events, including

- Flavour uncertainty
- Photon scale
- Method uncertainty
- Statistical uncertainty
- Extrapolation uncertainty

◆ About 3% JES uncertainty for PF jets

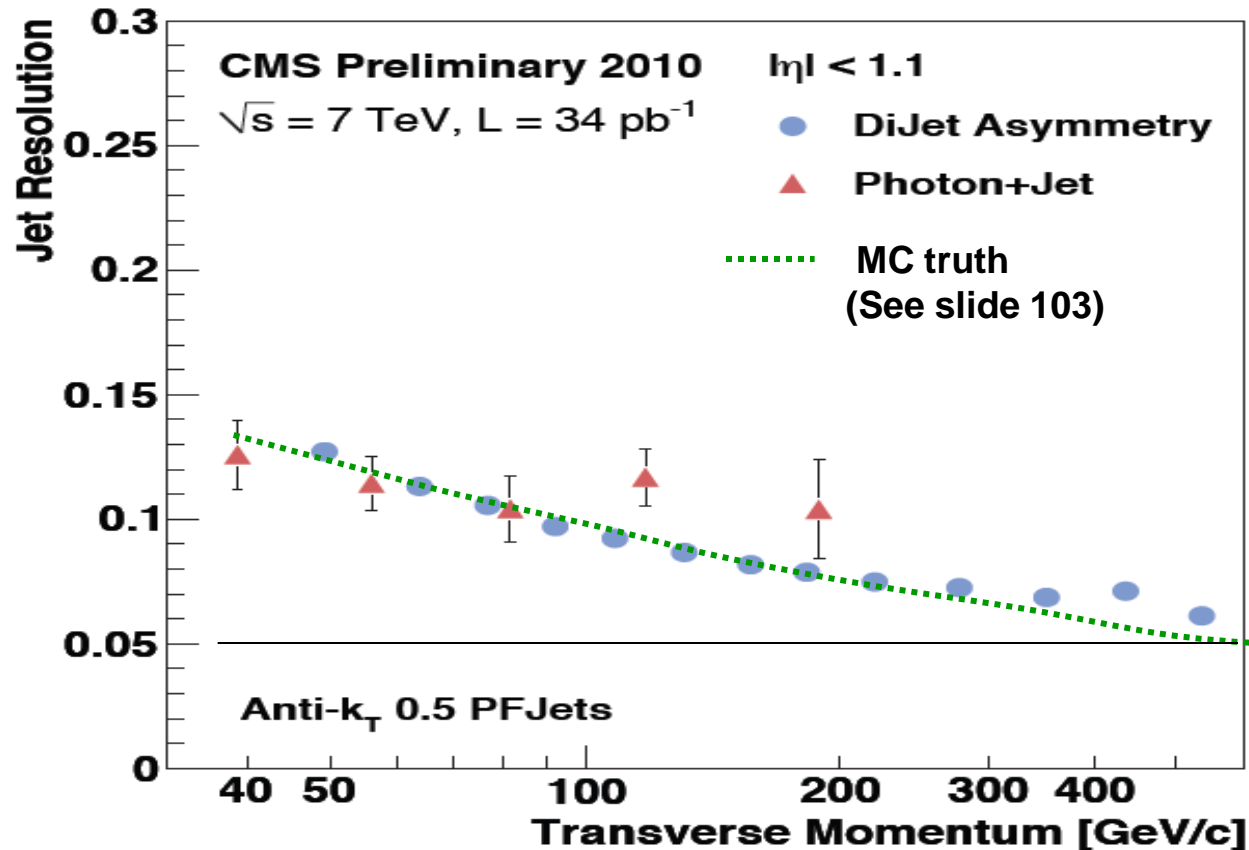
- with only 3 pb^{-1} of data
- would be 10% for calorimetric jets



Particle-Flow performance in CMS (10)


Measured jet resolution in data

- Still with γ jet events, but also with di-jet events (p_T imbalance)



Particle-Flow performance in CMS (II)

Expected performance for Missing Transverse Energy, MET

- ◆ Aimed at measuring the p_T of neutrinos or 
- MET from particle flow is very intuitive (from p_T conservation):
 - Sum over all particles from particle flow reconstruction

$$MET = - \sum_{i=1}^{N_{\text{Particles}}} \vec{E}_T^i$$

→ The importance of having all particles reconstructed appears clearly

Down to the smallest possible p_T
with the best possible energy determination

- Another important variable, SET, measures the overall activity in the event

$$SET = \sum_{i=1}^{N_{\text{Particles}}} E_T^i$$

Particle-Flow performance in CMS (12)

Expected performance for Missing Transverse Energy, MET (cont'd)

◆ Aimed at measuring the p_T of neutrinos or 

• MET from calorimetry-alone is, instead, a complicated beast :

→ Raw MET :

Sum over all calo towers

$$MET = - \sum_{i=1}^{N_{\text{Towers}}} \vec{E}_T^i$$

→ Corrected for muons :

Muon-corrected MET

$$- \sum_{i=1}^{N_{\text{muons}}} \vec{E}_T^i$$

→ Corrected for jet response :

Type-I corrected MET

$$- \sum_{i=1}^{N_{\text{Jets}}} \left(\vec{E}_{T,\text{corr}}^i - \vec{E}_{T,\text{raw}}^i \right)$$

→ Corrected for un-clustered energy response

Type-II corrected MET

$$- \alpha \sum_{i=1}^{N_{\text{Unclustered Towers}}} \vec{E}_T^i$$

Particle-Flow performance in CMS (13)

Expected performance for Missing Transverse Energy, MET (cont'd)

◆ Aimed at measuring the p_T of neutrinos or 

• MET from calorimetry-alone can also be corrected for tracks:

→ Raw MET:

Sum over all calo towers

$$MET = - \sum_{i=1}^{N_{\text{Towers}}} \vec{E}_T^i$$

→ Corrected for muons:

Muon-corrected MET

$$- \sum_{i=1}^{N_{\text{muons}}} \vec{E}_T^i$$

→ Corrected for track expected response in a calorimetric road:

Track-corrected MET

$$- \sum_{i=1}^{N_{\text{Tracks}}} \left(\vec{p}_T^i - \vec{E}_{T,\text{expected}}^i \right)$$

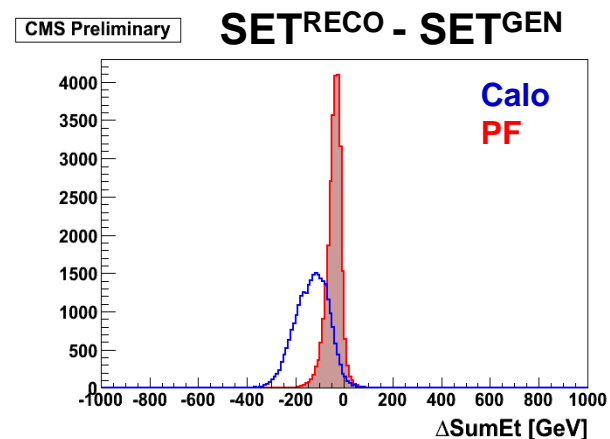
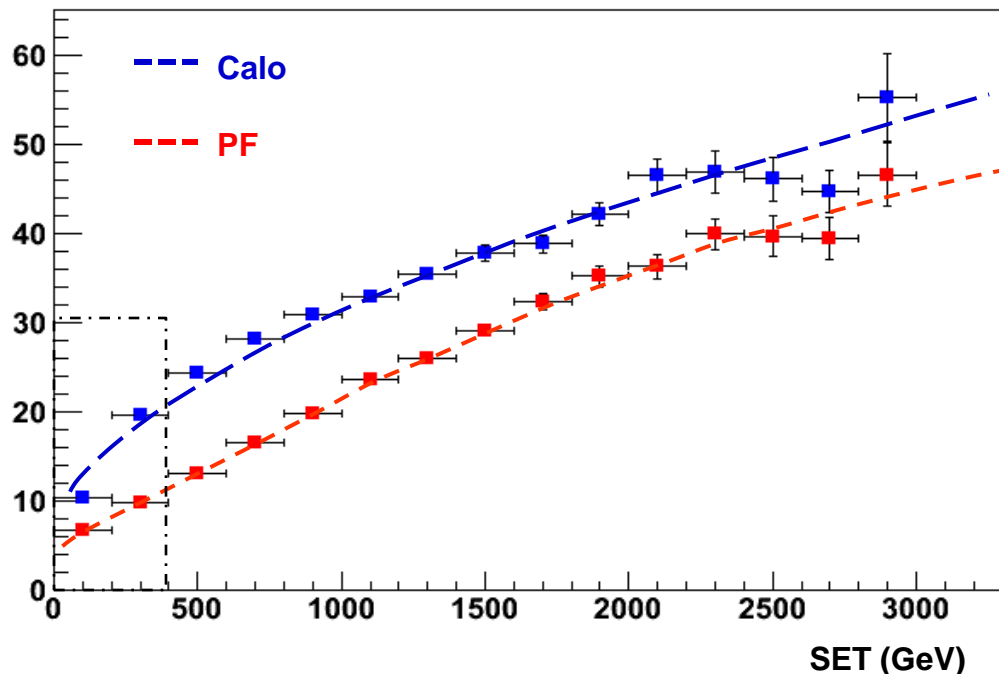
• ... which gives an "Energy-Flow" type of MET

→ An intermediate step, if one is too shy to go to full Particle Flow

Particle-Flow performance in CMS (14)

- Expected MET resolution for multi-jet events
 - No MET expected in these events (no neutrinos):
 - $\sigma(\text{MET}_{x,y})$ is the resolution to measure 0

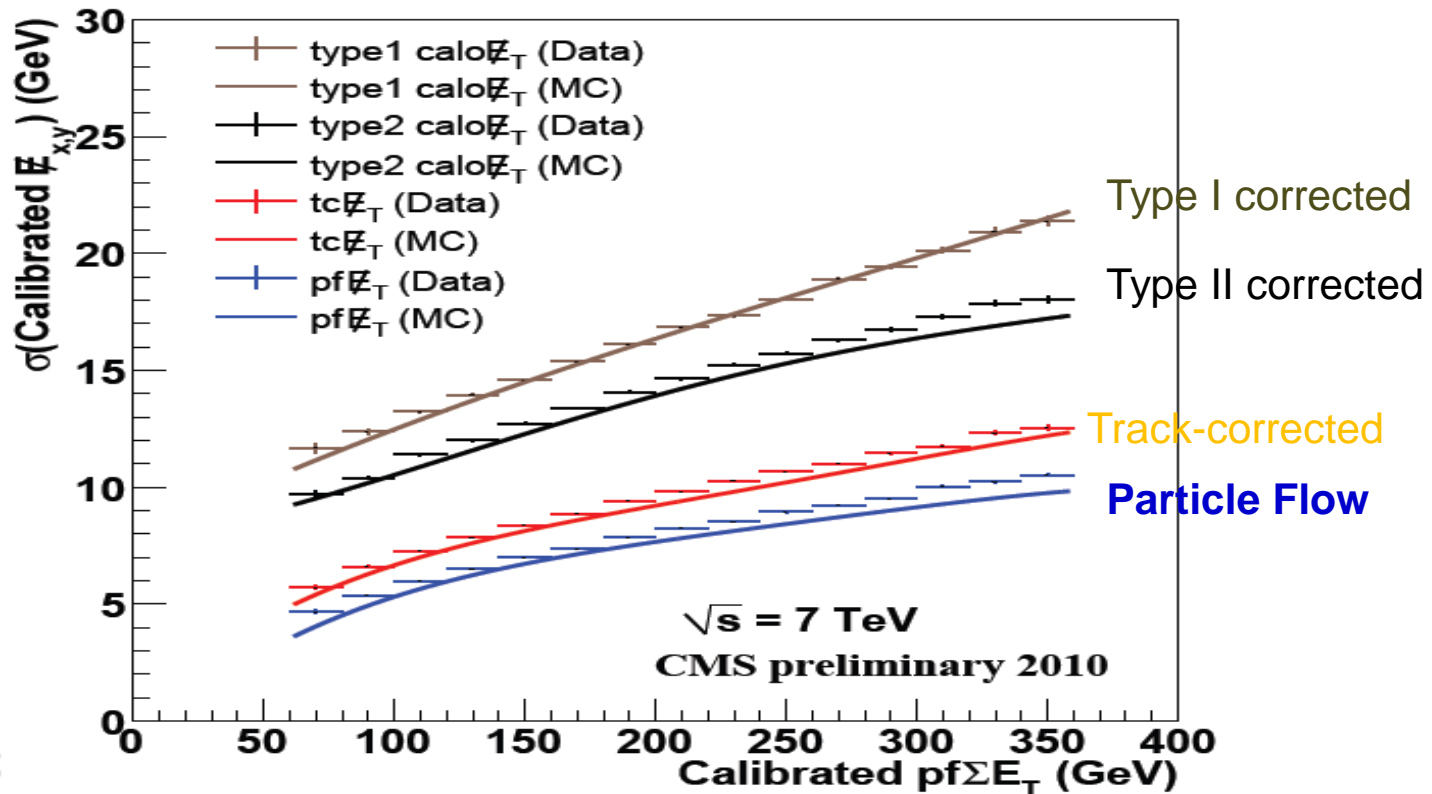
$\sigma(\text{MET}_{x,y})$ (GeV)



- Improved resolution by a factor > 2 all the way to SET = 1 TeV
 - Significant expected improvement up to the largest SET values

Particle-Flow performance in CMS (15)

- Measured MET resolution for multi-jet events
 - Factor of 2 improvement confirmed with data, on both x,y components

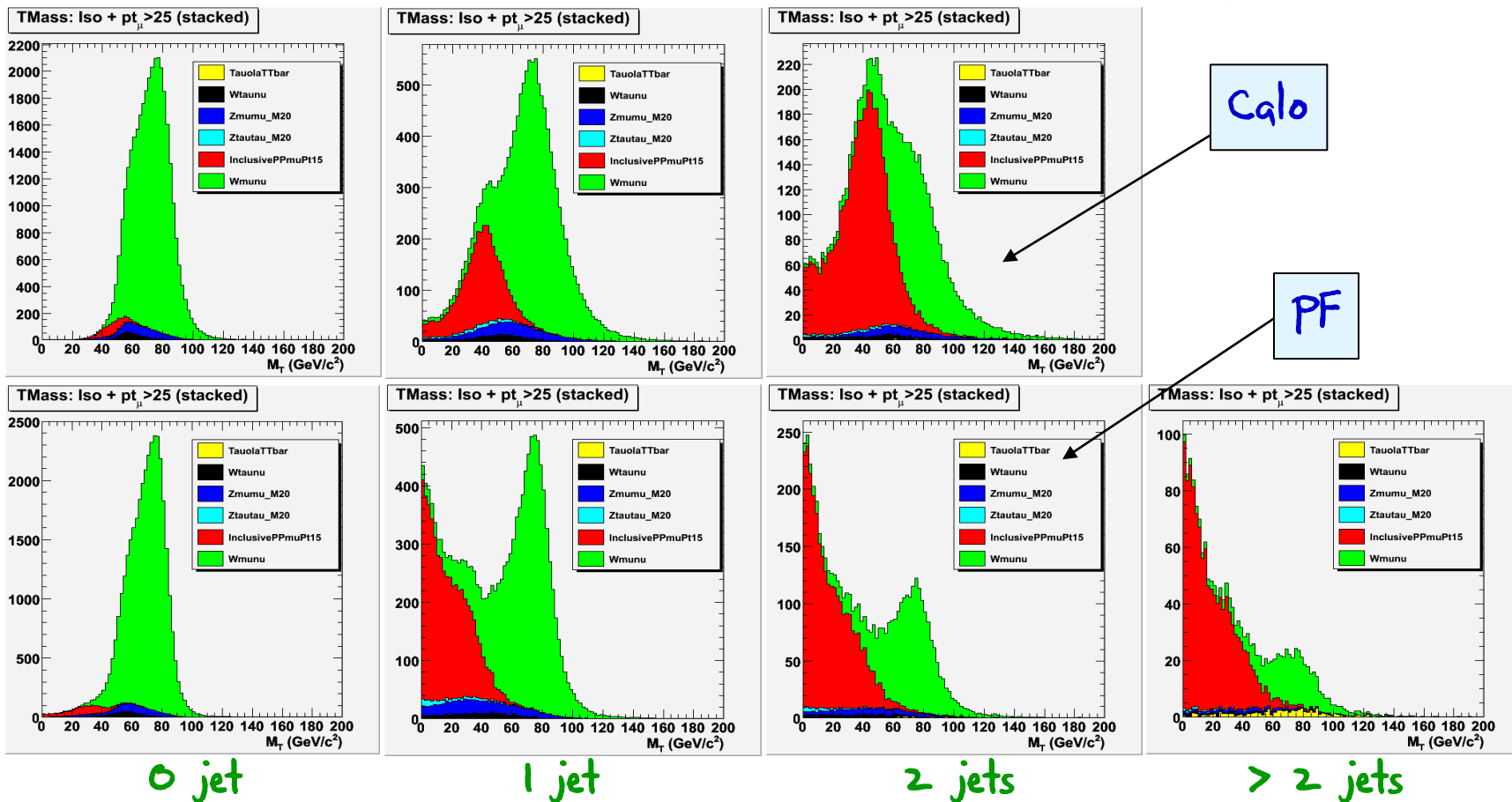


- Residual disagreement with simulation will disappear with HCAL calibration

Particle-Flow performance in CMS (16)

Expected MET resolution for events with missing energy

♦ $w + \text{jets}$ events, with $w \rightarrow \mu\nu$: Transverse mass made with $p_T(\mu)$ and MET



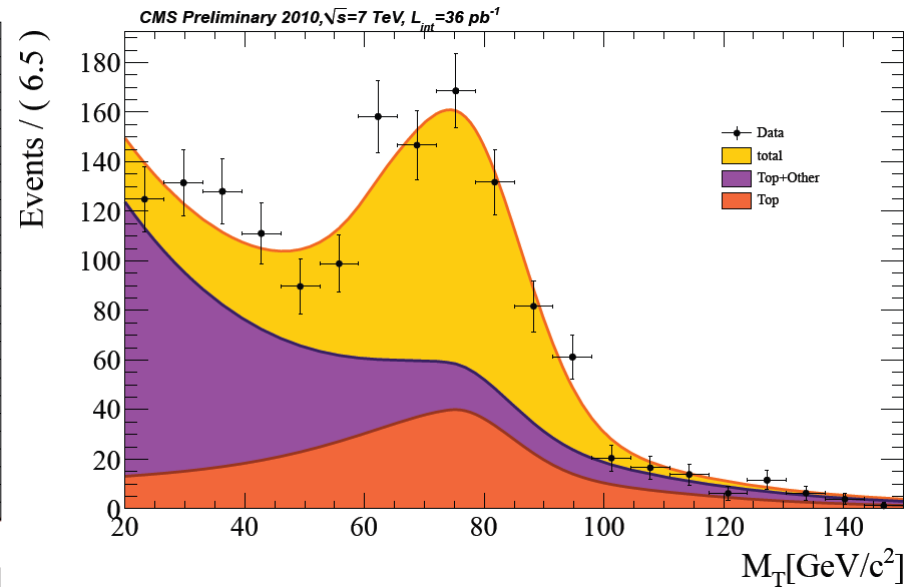
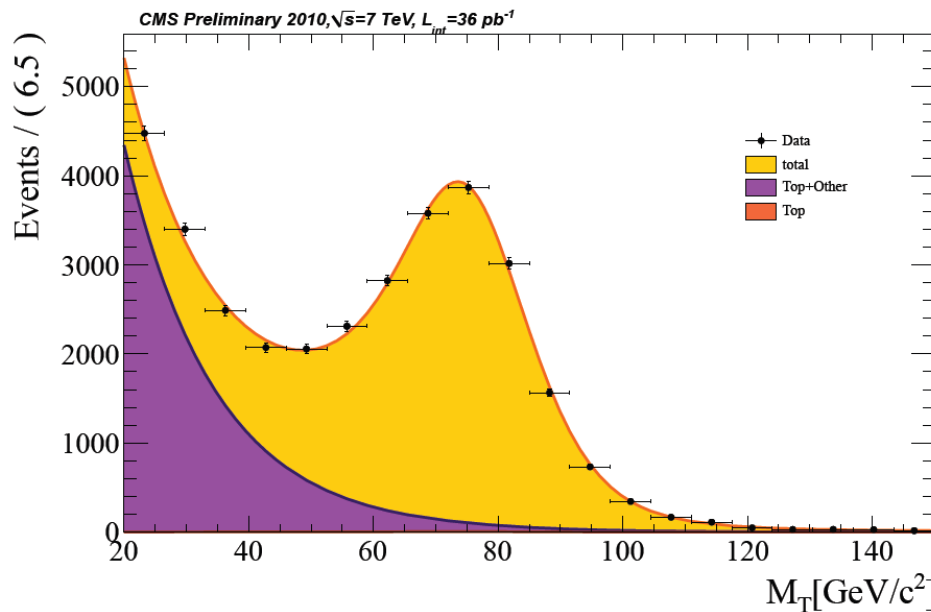
Particle-Flow performance in CMS (π)

Measured MET resolution for events with missing energy

◆ $w + \text{jets}$ events, with $w \rightarrow \mu\nu$: Transverse mass made with $p_T(\mu)$ and MET

$w + \text{one jet}$

$w + \text{three jets}$



- Missing energy resolution and response for neutrino confirmed
→ Both for events without and events with missing energy

Particle-Flow performance in CMS (18)

- CMS was (by chance) almost designed for particle-flow reconstruction
 - ◆ Hermeticity, Simplicity, Granularity, Magnetic Field [Slide 35]
 - Found to be just adequate
 - Even if one could have done better (tracker, HCAL)
 - ◆ Material : External constraints led to a (very) thick tracker [Slide 35]
 - Caused a lot of work for years to deal with it
 - Photon conversions, Electron Bremsstrahlung, Hadron nuclear interactions
 - ◆ Redundancy : No energy measurement redundancy in the calorimeters [Slide 35]
 - Source of continuous nightmares for calorimeter cleaning
 - ◆ But there was a handful of enthusiastic and knowledgeable people [Slide 36]
 - who believed in the success of the project from first principles
 - And who completed it in about four years of ceaseless work
 - (Remember : four months in ALEPH, because of the thoughtful design)
 - ◆ Particle Flow is now used in most CMS physics analyses

Bad/suboptimal design is not the worse enemy !

□ People are worse ...

- ◆ Religious war from calorimeter experts / users (50% of the CMS collaboration)
 - "Particle Flow never worked at a hadron collider" 2007
 - "Particle Flow is too complicated - calorimeters are simple" 2008
 - "It may work on simulation, but it won't work on data" 2009
 - "It may work on data, but it does not bring much ..." 2010
- ◆ Religious war from tracking experts / users (50% of the CMS collaboration)
 - "Iterative tracking is hopeless in such an environment" 2006
 - "It may work for electrons, but what else do we need it for?" 2007
 - "It works, but it does not bring that much" 2008
- ◆ Iterative tracking has become the official CMS tracking 2009
- ◆ Particle Flow is now used in most analyses in CMS 2011
 - But still people think it's magic inside

Concluding remark

- Extracts from Richard Wigmans' lecture on Monday, Jan. 31st
 - ◆ A number of religious statements, typical of a calorimeter expert
 - "Calorimeter granularity brings only confusion"
 - "The fact that 65% of the energy is perfectly measured in tracker is irrelevant"
 - "GEANT4 has never predicted anything correctly concerning hadron calorimetry"
 - "Advocates of particle-flow reconstruction use phony statistics"
 - "Because they are not happy of the results"

- Conclusion
 - ◆ once you have convinced yourself that you have a good idea
 - Based on scientific observations (first principles, feasibility study, ...)
 - Don't let experts (of something else) Kill your enthusiasm !

Just think, move forward, and FOLLOW YOUR DREAMS