

# "What are the atoms of a BH?"

[Bekenstein-Hawking '70s]



Black Hole

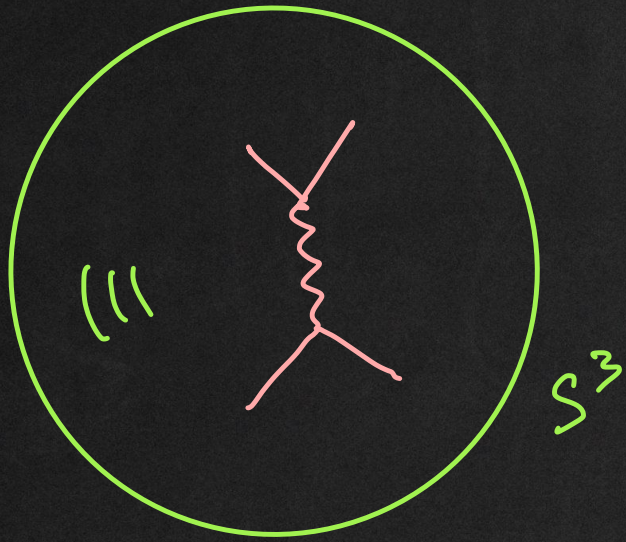
$$S_{BH} = \frac{1}{4} \frac{A_H}{l_p^2}$$

$$\left( l_p = \sqrt{\frac{\hbar G_N}{c^3}} \right)$$

$$S_{BH} \stackrel{?}{=} \log d_{micro}$$



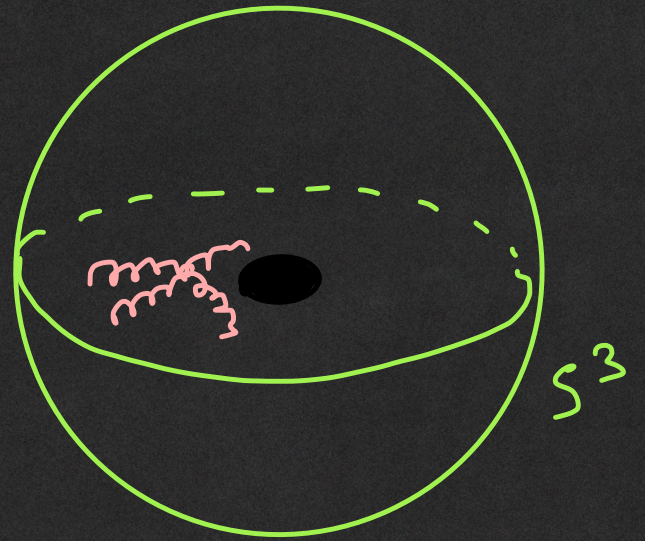
# AdS/CFT $\Rightarrow$ definition of quantum gravity



CFT<sub>4</sub> on  $S^3 \times \mathbb{R}_t$

(gauge theory)

=



AdS<sub>5</sub> - Quantum gravity

$\partial(\text{AdS}_5) = S^3 \times \mathbb{R}_t$

$\lceil S^3 \rightsquigarrow 3\text{-manifold} \rceil$

$\lceil \partial \equiv \text{conformal boundary} \rceil$



CFT

$N=4$  SYM  $U(N)$

$(N, \lambda \equiv g_{\text{YM}}^2 N)$

Gauge theory  $\xleftarrow{\lambda \ll 1}$

$\lambda = l^4$

$\xrightarrow{\lambda \gg 1}$

Gravity  
(semi classical)

AdS

$\text{IB} / \text{AdS}_5 \times S^5$

$(g_s, l = l_{\text{AdS}}/l_s)$

('t Hooft limit)

$\downarrow 1/N$

Non-planar

$1/N = g_s$   
 $\Downarrow$   
 $1/N^2 = G_N$

$\downarrow g_s$

Quantum gravity



# BH Thermodynamics in AdS/CFT



$$Z_{\text{CFT}}(\lambda, N; \beta) = Z_{\text{grav}}^{\text{AdS}}(l, g_S; \beta)$$

$$= \text{Tr}_{\mathcal{H}} e^{-\beta H}$$

$$= \sum_E d_{\text{micro}}(E) e^{-\beta E}$$

$$= \sum_{\alpha \in \text{Saddles}} \exp(-I^\alpha)$$

$$= e^{-I^{\text{AdS}}} + e^{-I^{\text{BH}}} + \dots$$



## E.g. Pure AdS<sub>5</sub> gravity

$$I[g] = -\frac{1}{16\pi G_N} \int d^5x \sqrt{g} \left( R + 6/l^2 \right)$$

Ansatz:  $(it = t_E \sim t_E + \beta)$

$$ds^2 = g(r) dt_E^2 + \frac{dr^2}{g(r)} + r^2 d\Omega_3^2$$



- $AdS_5 : g(r) = 1 + r^2/l^2$

- $(AdS_5) BH : g(r) = 1 + \frac{r^2}{l^2} - \frac{r_0^2}{r^2}$   
 (Schwarzschild)

$$(r_0^2 = \frac{8}{3\pi} G_N M)$$

Horizont :  $g(r_+) = 0$

$$T_{BH} = \frac{1}{\beta} = \frac{1}{4\pi} g'(r_+)$$



# Note: "Small" & "Large" BHs

$$T_{\text{BH}} = \frac{1}{2\pi l} \left( \frac{2r_+}{l} + \frac{l}{r_+} \right)$$

$$\Rightarrow T_{\text{min}} \text{ at } \frac{r_+}{l} = \frac{1}{\sqrt{2}}$$

Exercise  $\rightarrow$

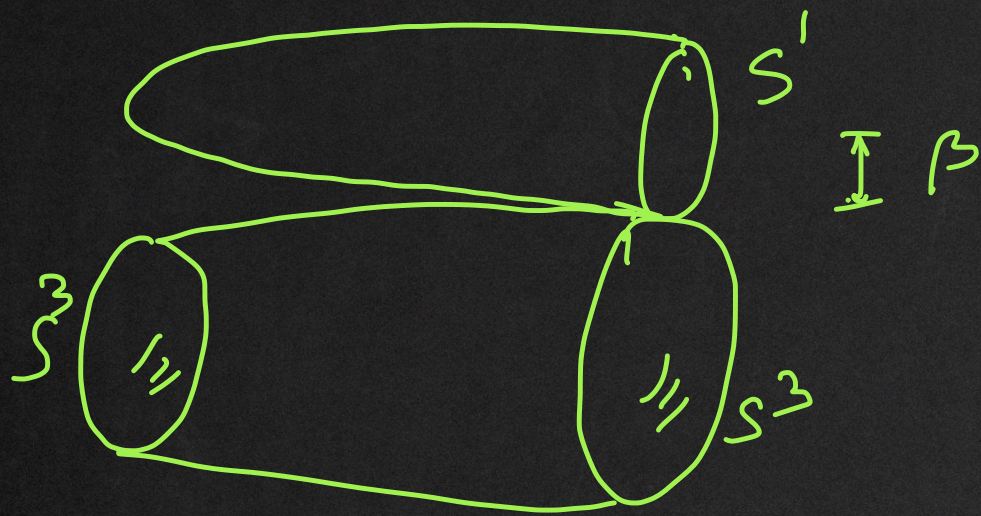
•  $\frac{r_+}{l} \ll 1 \Rightarrow T_{\text{BH}} \rightarrow \frac{1}{2\pi r_+}$  (specific heat  $\downarrow C_v < 0$ )  
(cf flat space)

•  $\frac{r_+}{l} \gg 1 \Rightarrow T_{\text{BH}} \rightarrow \frac{r_+}{\pi l^2}$   $C_v > 0$   
 $\Downarrow$   
stable

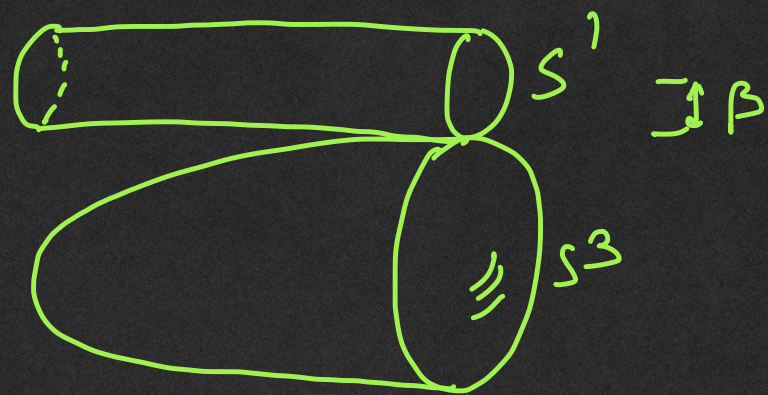


# What happens below $T_{min}$ ?

(E)BH



(E)AdS



$$I[g]_{\text{som}} = \frac{1}{2\pi G_N} \int d^5x \sqrt{g}$$

= "volume of spacetime"



# Hawking - Page transition

[H-P '87]

$$I^{\text{BH}} - I^{\text{AdS}} = \frac{\pi^2}{4G_N} \frac{M_+^3 (l^2 - M_+^2)}{(2M_+^2 + l^2)}$$

→  
Exercise

⇒

High  $T$  : BH dominates

Low  $T$  : AdS "



Free energy = on-shell action

[Gibbons-Hawking '77]

$$S_{\text{BH}} = \frac{A_H}{4G_N} = \frac{\pi^2}{2} \frac{r_+^3}{G_N}$$

$$I := I_{\text{AdS BH}} - I_{\text{AdS}} \quad (= \text{reg. on-shell action of BH})$$

$$E = \partial I / \partial \beta, \quad S_{\text{BH}} = \beta E - \underbrace{I}_{\uparrow F}$$



True for more general ensembles

Rotating BH: Horizon generated by

$$V = \frac{\partial}{\partial t} + \Omega \frac{\partial}{\partial \phi}$$

angular  
velocity  
at  
horizon.

Charged BH:  $g_{\mu\nu}, A_\mu$

Electrostatic potential

$$\underline{\Phi} = V \cdot A \Big|_{r_+} - V \cdot A \Big|_{\infty}$$



# "Quantum Statistical Relation" [Gibbons-Hawking]

$$I = -S_{\text{BH}} + \beta E - \beta \Omega J - \beta \Phi Q$$

$$E = \frac{\partial I}{\partial \beta}, \quad J = -\frac{1}{\beta} \frac{\partial I}{\partial \Omega}, \quad Q = -\frac{1}{\beta} \frac{\partial I}{\partial \Phi}$$

AdS gravity  $\Rightarrow$  Thermodynamics  
Canonical ensemble



What happens in  $CFT_4 / S^3 \times S'_\beta$ ?

$CFT_4$  d.o.f

$$\Rightarrow E \sim \frac{1}{\beta^4} \quad \text{as } \underline{\beta \rightarrow 0}$$

$$\Rightarrow S \sim \frac{1}{\beta^3} \quad \text{" "}$$

$U(N)$  gauge theory  $\Rightarrow$

$$S_{CFT} \sim \frac{N^2}{\beta^3} \quad \begin{array}{l} \beta \rightarrow 0 \\ N \rightarrow \infty \end{array}$$
$$S_{BH} \sim \frac{1}{G_N} r_+^3 \quad \checkmark$$



# Basic picture of BH thermodynamics in AdS/CFT

	Gravity	Gauge theory	Entropy
Low temp $\beta \rightarrow \infty$	AdS (+ small fluc.)	<u>Confined</u> glueballs (singlets)	$1 + \dots$
High temp $\beta \rightarrow 0$	AdS BH (+ small fluc.)	<u>Deconfined</u> quark-gluon plasma	$\frac{N^2}{\beta^3} + \dots$



# Supersymmetric BHs in AdS space

Minimal sugra (gauged) 8 supercharges

$$\mathcal{L} = R + \frac{12}{l^2} - \frac{2}{3} l^2 F^2 - \frac{1}{e} \frac{8}{27} A \wedge F \wedge F$$

$$\text{AdS}_5 \supset \underbrace{\text{su}(2)}_J \times \underbrace{\widetilde{\text{su}(2)}}_{\tilde{J}} \quad \text{U(1) gauge } R$$

• BH charges:  $(R, J_1, J_2)$

$$\begin{aligned} J_1 &= J + \tilde{J} \\ J_2 &= J - \tilde{J} \end{aligned}$$

•  $\{Q, Q^\dagger\} = E - J_1 - J_2 - \frac{3}{2} R$



$$(J_1 = J_2 = J)$$

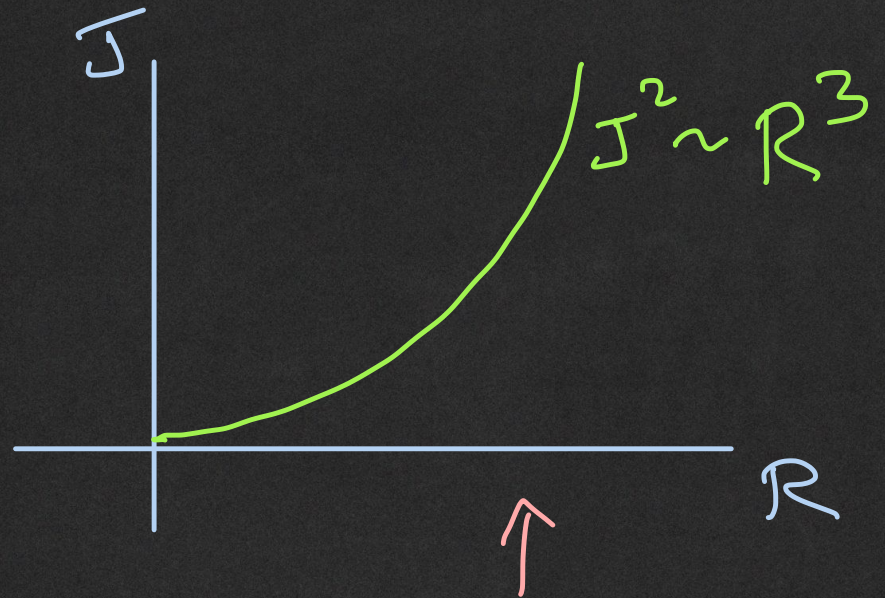
$$J_* = \frac{\pi a^2 (a+3)}{2(1-a)^3}$$

Susy

$$R_* = \frac{\pi a}{(1-a)^2}$$

$$(0 < a < 1)$$

$$l=1$$



Exercise

Note:  $J=0 \Rightarrow a=0 \Rightarrow$  SUSY BH  
 $\Rightarrow R=0, S_{BH}=0 \Rightarrow$  must rotate!



# Thermodynamics?

$$\beta \rightarrow \infty, \quad \Omega_1^* = \Omega_2^* = 1, \quad \Phi^* = 3/2$$

• Frozen! Nothing to vary!

cf  $\Omega(J, Q), \Phi(J, Q)$

• Infinite throat in the interior

$$\Rightarrow I^{BH} \rightarrow \infty \quad (e^{-I^{BH}} = 0)$$



Extremal BH = lim (non-extremal BH)

4-parameter family of BHs

[Chong-Greig-Lu-Pope '05]

$$(J_1, J_2, E, R) \iff (a, b, m, q)$$

$$\text{Horizon } g(r) = 0$$

$$\implies (r^2 - r_+^2)(r^2 - r_0^2)(r^2 - r_-^2) = 0$$

$$(r_+ \geq r_0 \geq r_- : \text{fns of } (a, b, m, q))$$



# Thermodynamic relations (QSR)

hold as expected

$$I = \beta E - S_{\text{BH}} - \beta \Omega_i J_i - \beta \Phi R$$

$$E = \frac{\partial I}{\partial \beta}, \quad J_i = -\frac{1}{\beta} \frac{\partial I}{\partial \Omega_i},$$

$$R = -\frac{1}{\beta} \frac{\partial I}{\partial \Phi}$$

∴

✓



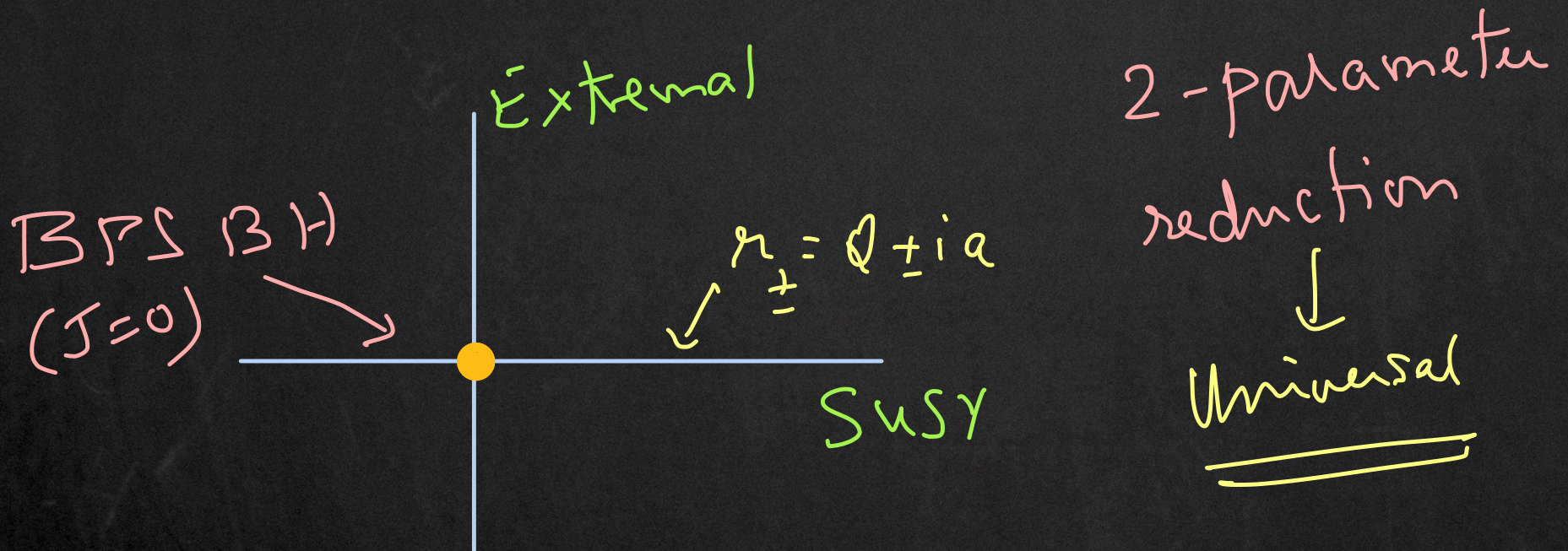
# SUSY vs Extremality (Reminder) [h/t J. Lucietti]

4d Kerr-Newman solution  $(M, Q, J)$

Regularity  $\Rightarrow f(M, Q, J) = M^2 - Q^2 - J^2/M^2 \geq 0$

Extremal:  $f(M, Q, J) = 0$  ( $r_+ \rightarrow r_-$ )

SUSY:  $M = Q$  + regular  $\Rightarrow J = 0$ .



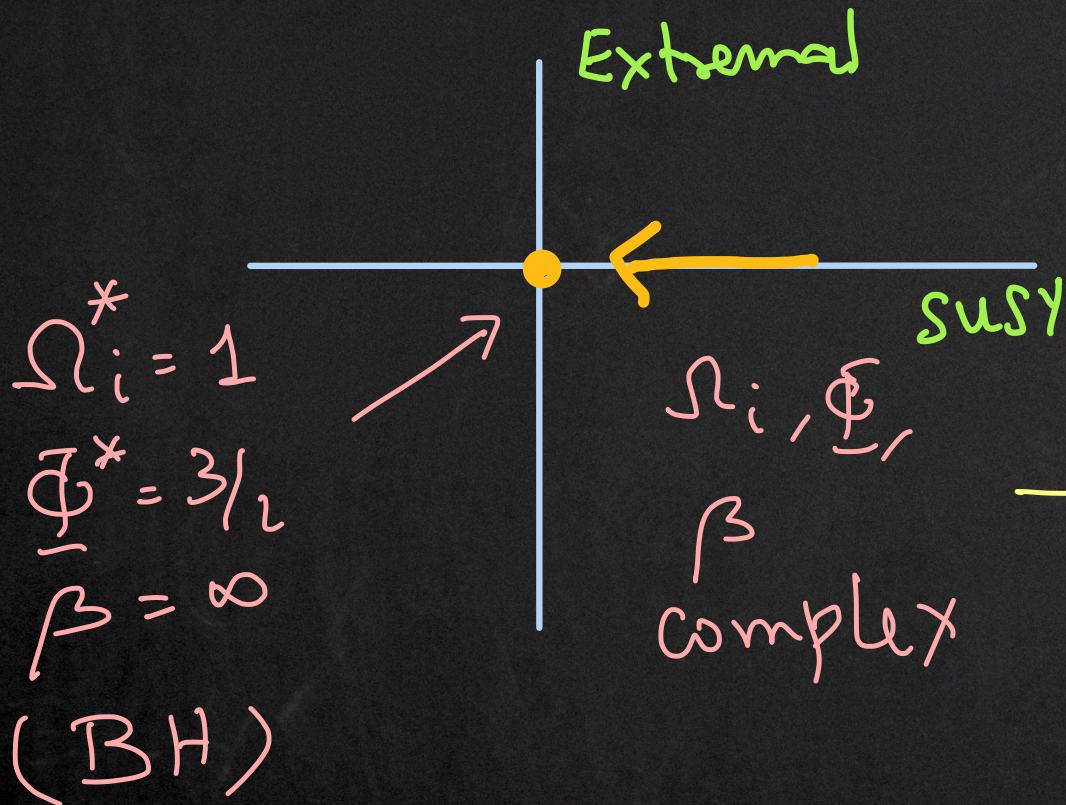


# AdS<sub>5</sub> SUSY BH as a limit

• Extremal  $\Rightarrow r_+ = r_0$

[CCMM '18]  
(cf Silva '06)

• SUSY  $\Rightarrow q = \frac{m}{1+a+b}$



$$\beta (1 + \Omega_1 + \Omega_2 - 2\bar{\Phi}) = 2\pi i$$



Meaning of constraint : global regularity

$$\underline{\text{KSE}}: \left( \nabla_{\mu} - i (\Gamma \cdot F)_{\mu} - \frac{1}{l} \Gamma_{\mu} - i A_{\mu}^R \right) \xi = 0$$

$$\begin{aligned} \text{V} \cdot \text{KSE} &\Rightarrow \frac{1}{2} (1 + \Omega_1 + \Omega_2 - 2 \underline{\Phi}) = \frac{\pi i}{\beta} \\ \parallel \\ &\partial_t + \Omega_1 \partial_{\phi} + \Omega_2 \partial_{\psi} \end{aligned}$$

Euclidean  
BH



Smoothness

$\Rightarrow \xi$  anti-periodic

$\Rightarrow$  Need  $A_{\mu}^R \neq 0$  at  $\infty$ !



# Supersymmetric "Thermodynamics"

$$\omega_i := \beta (\Omega_i - \Omega_i^*) \quad \beta \rightarrow \infty, \quad \Omega_i \rightarrow \Omega_i^* \quad \Phi \rightarrow \Phi^*$$

$$\varphi := \beta (\Phi - \Phi^*) \quad \underline{\omega_i, \varphi \text{ finite.}}$$

$$\Rightarrow \boxed{\omega_1 + \omega_2 - 2\varphi = 2\pi i}$$

• Action: 
$$\boxed{I = \frac{2\pi}{27} \frac{\varphi^3}{\omega_1 \omega_2}}$$
 (Note:  $\frac{\partial I}{\partial \beta} = 0!$ )

• QSR:

$$\begin{aligned} I &= \beta E - S_{\text{BH}} - \sum_i \beta \Omega_i J_i - \beta \Phi R \\ &= \beta \left( \underbrace{E - \Omega_i^* J_i - \Phi^* R}_{\{Q, Q^+\} = 0!} \right) - S_{\text{BH}} - \omega_i J_i - \varphi R \end{aligned}$$



# SUSY BH Thermodynamics

[CCMM '18]

BH entropy follows from on-shell action  $I$   
as constrained Legendre transform

$$I(\omega_i, \varphi) = -S_{\text{BH}} - \omega_i J_i - \varphi R - \Lambda(\omega_1 + \omega_2 - 2\varphi - 2\pi i) \quad (\otimes)$$

- $S_{\text{BH}} = A/4 = \pi \sqrt{3R^2 - N^2 J}$  (Exercise)

(\*) Inspiring observation of  
Hosseini-Hristov-Zaffaroni 2017



# What have we achieved?

- Non-linear constraint  $(J_i, R)$

$\rightsquigarrow$  linear constraint  $(w_i, \varphi)$

- Entropy  $\rightsquigarrow$  Canonical ensemble

$$S_{BH} = \text{extr} \left( -I(\varphi, w_i) - w_i J_i - \varphi R \right.$$

$$\left. - \Lambda (w_1 + w_2 - 2\varphi - 2\pi i) \right)$$

(reg) on-shell action of  $AdS_5$  BH



$$Z_{\text{Ads}}(\omega_i, \varphi) = e^{-\underline{I_{\text{gran}}(\omega_i, \varphi)}}$$

$$\begin{array}{c} \parallel \quad \downarrow \\ \boxed{Z_{\text{SCFT}}(\omega_i, \varphi)} \end{array}$$