

NANJING
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Craig Roberts ... <http://inp.nju.edu.cn/>

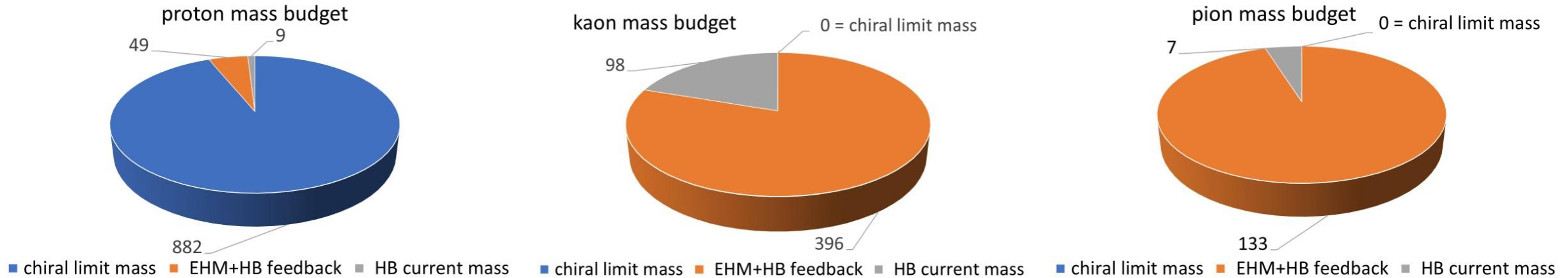
On Mass and Matter

Trace Anomaly

$$\Theta := T_{\mu\mu}^{\text{QCD}} = \frac{1}{4}\beta(\alpha(\zeta))G_{\mu\nu}^a G_{\mu\nu}^a + [1 + \gamma(\alpha(\zeta))] \sum_{i=u,d,\dots} m_i^\zeta \bar{q}_i q_i$$

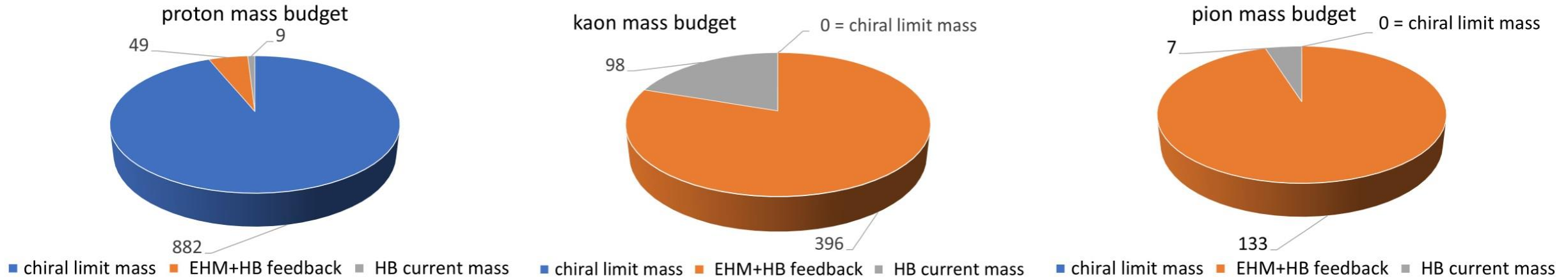
- It's not a miracle
- Straightforward consequence of quantum definition of interacting field theory
 - Electroweak theory has trace anomalies
 - Never talk about them because their consequences are expressed perturbatively
- The coupling runs if and only if there is a trace anomaly
- The masses and operators run if and only if the coupling runs
- Trace anomaly enables mass scales to emerge in the absence of a Higgs mechanism.
 - By itself, it doesn't guarantee such an outcome
- Magnitude of QCD trace anomaly in Nature is an empirical observable
 - Theory cannot predict it
 - But theory can elucidate its manifold connections, correlations, and manifestations

Mass Budgets



- Poincaré invariant and scale independent decompositions of masses of selected light-quark hadrons
- Proton
 - Higgs contribution
 - Renormalisation group invariant current-quark masses $\hat{m}_u \approx 4.5 \text{ MeV}$ & $\hat{m}_d \approx 2 \hat{m}_u$
 - Sum = 2% of proton mass
 - EHM+HB interference = constructive = 4% of proton mass
 - EHM alone = 94% of proton mass

Mass Budgets



➤ Proton

– Higgs contribution

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Assessment and estimate of size from low-energy ChPT analyses

– EHM+HB interference = constructive = 4% of proton mass

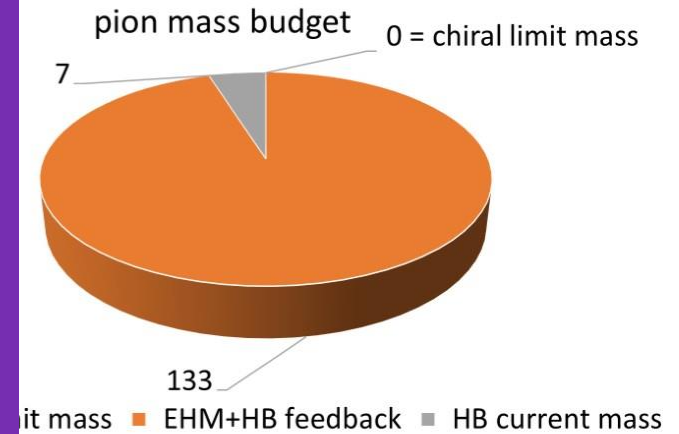
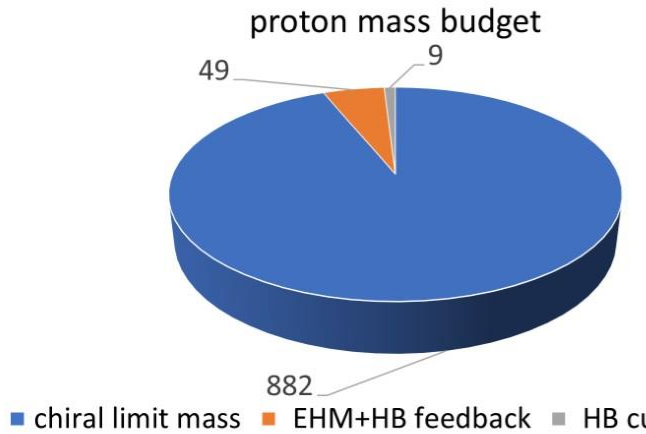
Assessment and estimate from analysis of low-energy π N scattering \sim nucleon σ -term

Mass Budgets

Assuming validity of all assumptions needed to justify formula for charmonium production, then analysis of GlueX data

$$\Rightarrow \zeta_h = 0.07(17) = \text{small}$$

Confirms that bulk of nucleon mass is emergent = EHM



➤ Proton

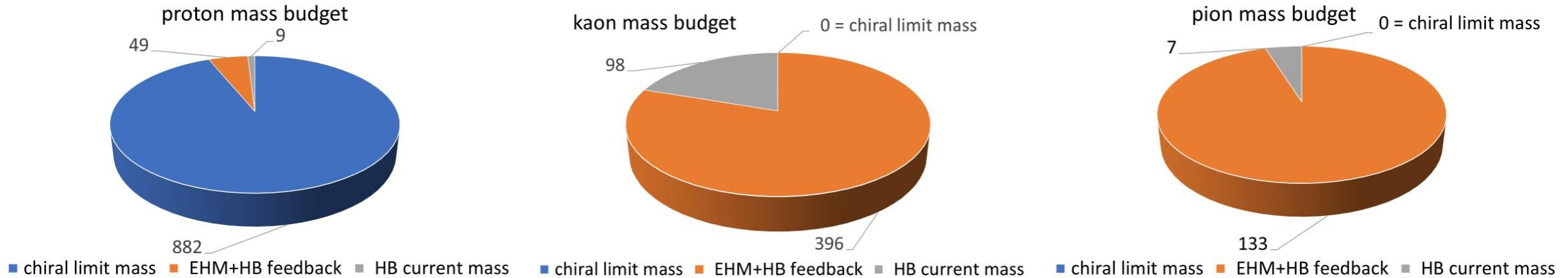
- EHM alone = 94% of proton mass
- Heavy-meson production at threshold – operator product expansion and low-energy theorems,

$$\mathcal{M}_{\Upsilon N} \simeq 2m_{\Upsilon} r_0^3 d_2 \frac{2\pi^2}{27} \left(2m_N^2 - \left\langle N \left| \sum_{i=u,d,s} m_i \bar{q}_i q_i \right| N \right\rangle \right) \simeq 2m_{\Upsilon} r_0^3 d_2 \frac{2\pi^2}{27} 2m_N^2 (1 - \zeta_h)$$

Such measurements do not provide news on the trace anomaly – instead, they may provide indirect information via the light-quark and s-quark in-nucleon σ -terms

Mass Budgets

How are these manifest differences expressed in parton distributions that characterise these hadrons?



➤ Kaon – highlights Higgs modulation

- Higgs contribution
 - Sum of current-quark mass = 20% of pion mass
- EHM+HB interference = constructive = 80% of kaon mass
- No Blue Zone because kaon, too, is NG mode \Rightarrow NO mass in chiral limit

➤ Pion

- Higgs contribution
 - Sum of current-quark mass = 5% of pion mass
- EHM+HB interference = constructive = 95% of pion mass
- No Blue Zone because pion = NG mode \Rightarrow there is NO mass in chiral limit

Trace Anomaly

$$\Theta := T_{\mu\mu}^{\text{QCD}} = \frac{1}{4}\beta(\alpha(\zeta))G_{\mu\nu}^a G_{\mu\nu}^a + [1 + \gamma(\alpha(\zeta))] \sum_{i=u,d,\dots} m_i^\zeta \bar{q}_i q_i$$

- Expressed in terms of the QCD Lagrangian degrees-of-freedom
- Poincaré invariant, but everything depends on the probe scale, ζ
- The trace anomaly is not homogeneous in the running coupling, $\alpha(\zeta)$; so, renormalisation-group-invariance does not imply form invariance of the right-hand-side.
- Material point because
 - Many discussions implicitly assume that all operators and associated identities are expressed with reference to a partonic basis, *i.e.* using elementary field operators that can be renormalised perturbatively
 - In this basis, the state-vector for any hadron is an extremely complicated wave function = intractable problem in any approach based on parton-basis Fock space expansion

Trace Anomaly

$$\Theta := T_{\mu\mu}^{\text{QCD}} = \frac{1}{4}\beta(\alpha(\zeta))G_{\mu\nu}^a G_{\mu\nu}^a + [1 + \gamma(\alpha(\zeta))] \sum_{i=u,d,\dots} m_i^\zeta \bar{q}_i q_i$$

- Different perspective is required at probe scales $\zeta < m_p$
- Metamorphosis from parton-basis to quasiparticle-basis occurs. Under such reductions in ζ ,
 - Light partons evolve into heavy dressed objects, corresponding to complex and highly nonlinear superpositions of partonic operators
 - Using the associated quasiparticle operators, the wave functions can be expressed in a relatively simple form.

Dressed-partons are quasiparticles ... they are real

- Lane:1974he & Politzer:1976tv
- QCD has a renormalization group invariant running quark mass

$$M(-Q^2) \underset{Q^2 \rightarrow \infty}{=} \frac{c}{Q^2} \left[\ln(Q^2 / \Lambda_{QCD}^2) \right]^{d_M - 1} + m \left[\frac{\ln(\mu^2 / \Lambda_{QCD}^2)}{\ln(Q^2 / \Lambda_{QCD}^2)} \right]^{d_M}$$

$$\frac{1}{Q^2} = \frac{1}{\Lambda_{QCD}^2}$$

$$e^{-\pi d_M / \alpha(Q^2)}$$

- ✓ Essentially nonperturbative
- ✓ Nonzero in chiral limit
- ✓ “c” is the pseudoscalar projection of the pion wave function onto the origin – direct measure of EHM = chiral condensate
 - ✓ Power-law × standard QFT scaling violation
 - ✓ Inaccessible in perturbation theory
- Essential singularity in the running coupling
- ✓ Now known to be origin of constituent-quark mass scale, introduced by Gell-Mann and Zweig

Higgs-induced
Runs as logarithm
Perturbative contribution
can be computed loop-
by-loop
BUT loop-by-loop always
gives contribution
proportional to “m”; so
vanishes in chiral limit

Dressed-partons are quasiparticles ... they are real

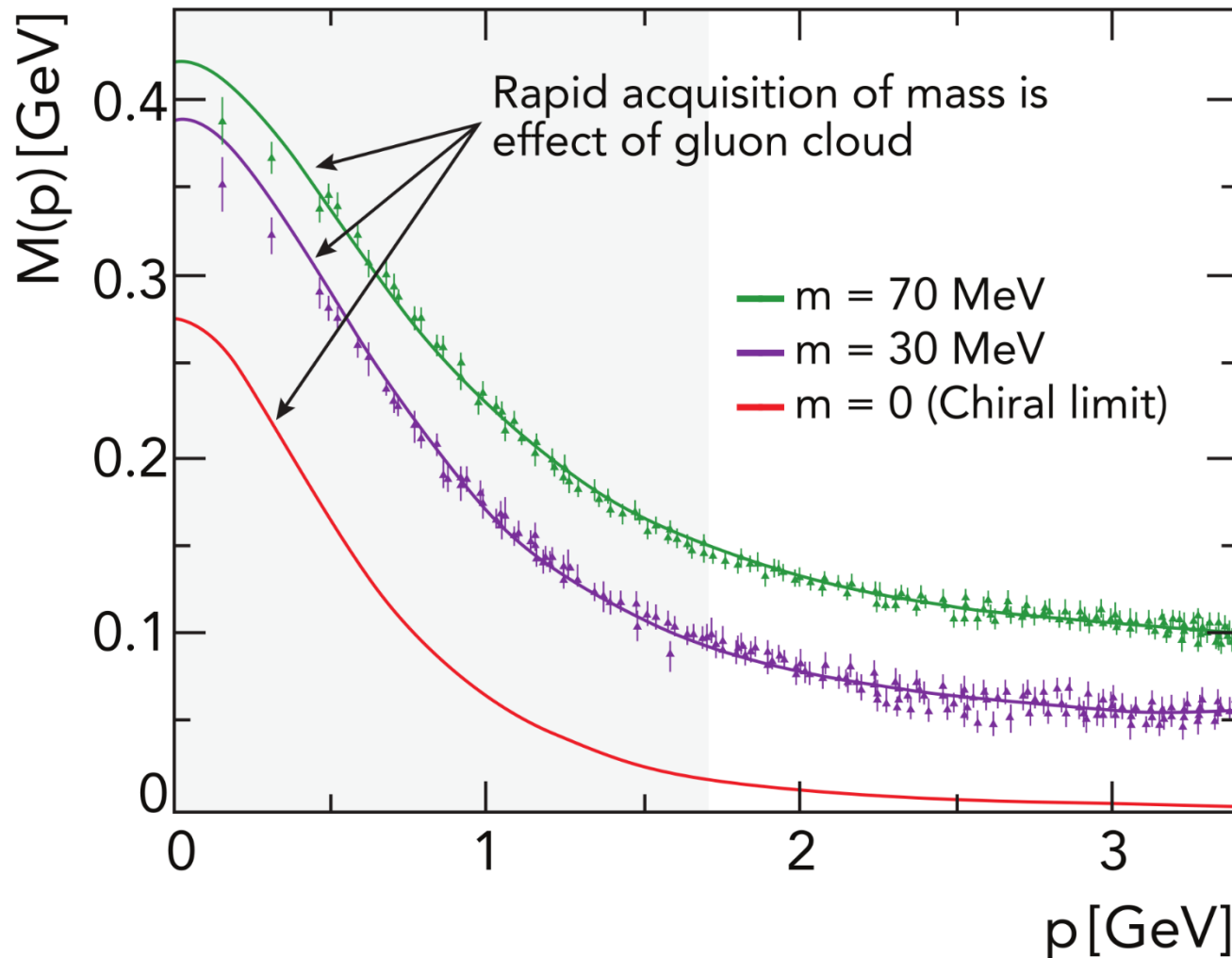
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$$M(-Q^2) \underset{Q^2 \rightarrow \infty}{=} \frac{c}{Q^2} \left[\ln \right]$$

$$\frac{1}{Q^2} = \frac{1}{\Lambda_{QCD}^2} e^{-\pi d_M / \alpha(Q^2)} \quad \checkmark \text{ Ess}$$

- ✓ "c" is the pseudoscalar projection onto the origin – direct measure of I
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✓ Now known to be origin of constituent-quark mass scale, introduced by Gell-Mann and Zweig



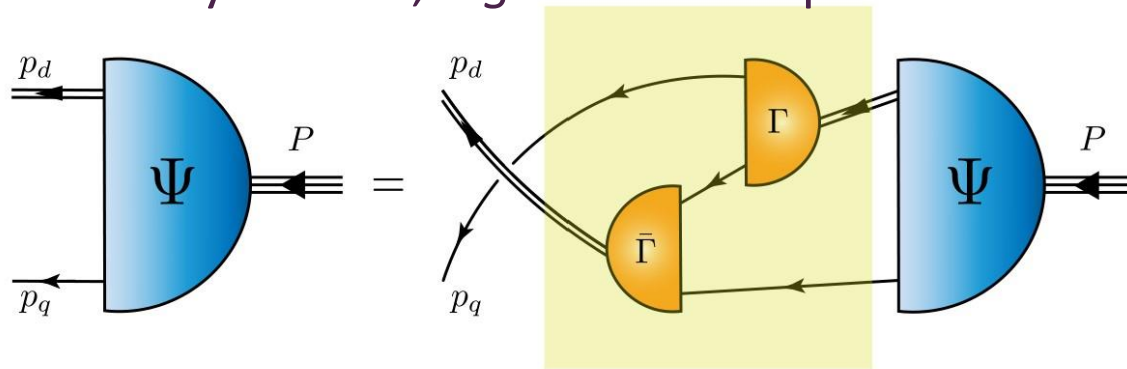
Metamorphosis expressed in Schwinger function field equations

Parton basis $\langle p(P) | \Theta | p(P) \rangle \stackrel{\zeta \gg m_p}{=} \langle p(P) | \left[\frac{1}{4} \beta(\alpha(\zeta)) G_{\mu\nu}^a G_{\mu\nu}^a + [1 + \gamma(\alpha(\zeta))] \sum_{i=u,u,d} m_i^\zeta \bar{q}_i q_i \right] | p(P) \rangle$

Quasiparticle basis $\stackrel{\zeta \lesssim m_p}{=} \langle \tilde{p}(P) | [\mathcal{D}_3 + I_3] | \tilde{p}(P) \rangle$

$$\mathcal{D}_3 = \sum_{i=u,u,d} M_f(\zeta) \bar{Q}_f(\zeta) Q_f(\zeta), \quad I_3 = \frac{1}{4} [\beta(\alpha(\zeta)) \mathcal{G}_{\mu\nu}^a(\zeta) \mathcal{G}_{\mu\nu}^a(\zeta)]_{2PI}$$

- Achieved mathematically by resumming infinite towers of diagrams in one-body, two-body, three-body sectors, e.g. Faddeev equation for the proton mass and wave function



Merit of EHM resummed basis = the wave function can be calculated

Enables verifiable predictions to be made and connections to be drawn with other applicable computational schemes

- **Parton basis** \Rightarrow almost all proton mass carried (somehow) by partonic/free-field gluons
- **EHM-resummed basis** \Rightarrow bulk of proton mass stored in dressed-quark mass-scale with residual inter-particle interactions providing the refinements

Avoid a Mass Crisis

- Steer away from decompositions whose interpretations are observer (frame) specific and/or whose interpretations change “weirdly” with scale
 - *e.g.*, exists at least one decomposition and scheme in which gluons contribute more than 100% of the proton’s mass and quarks a negative amount
- Focus on the Poincaré invariant form factors of the stress-energy tensor in hadron states
 - Beware ... like spin and orbital angular momentum, any attempt to decompose the form factors into contributions from different parton species is frame and scale dependent
 - OK, if one is using a mathematical tool that can draw reliable connections between the frames and scales

- Think before you leap

$$\langle \pi(q) | \Theta_0 | \pi(q) \rangle = 0$$

can only mean that pion is empty of glue in a theory that doesn’t produce a pion

- All bound states possess an interaction energy
- Rigorous mathematical statements in QCD reveal how Dyson and Bethe-Salpeter equations conspire to make a massless bound state from emergent massive constituent quasiparticles, themselves possessing a traceable, nonperturbative connection to partons in QCD’s Lagrangian