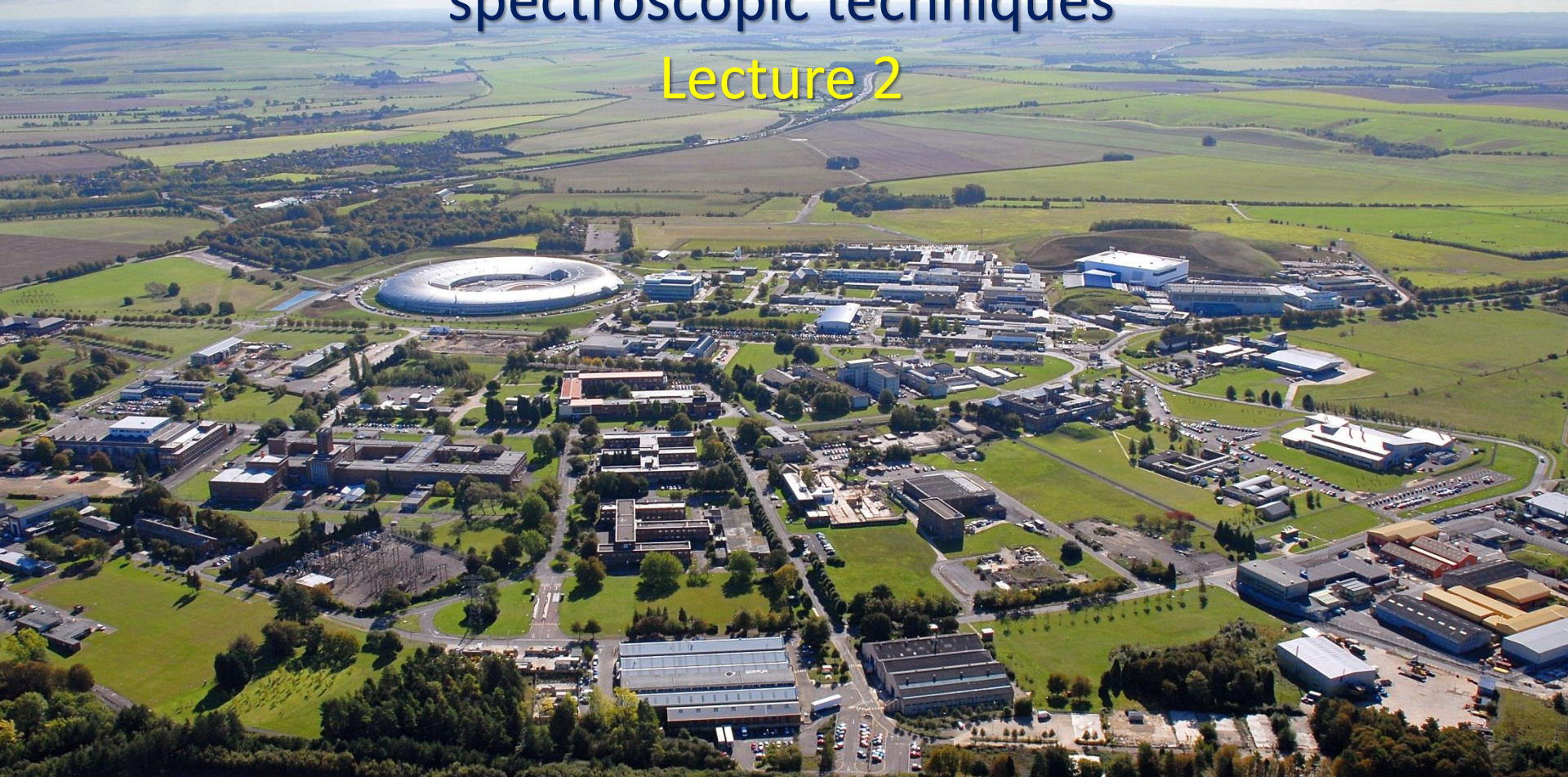


Synchrotron and neutron based diffraction and spectroscopic techniques

Lecture 2

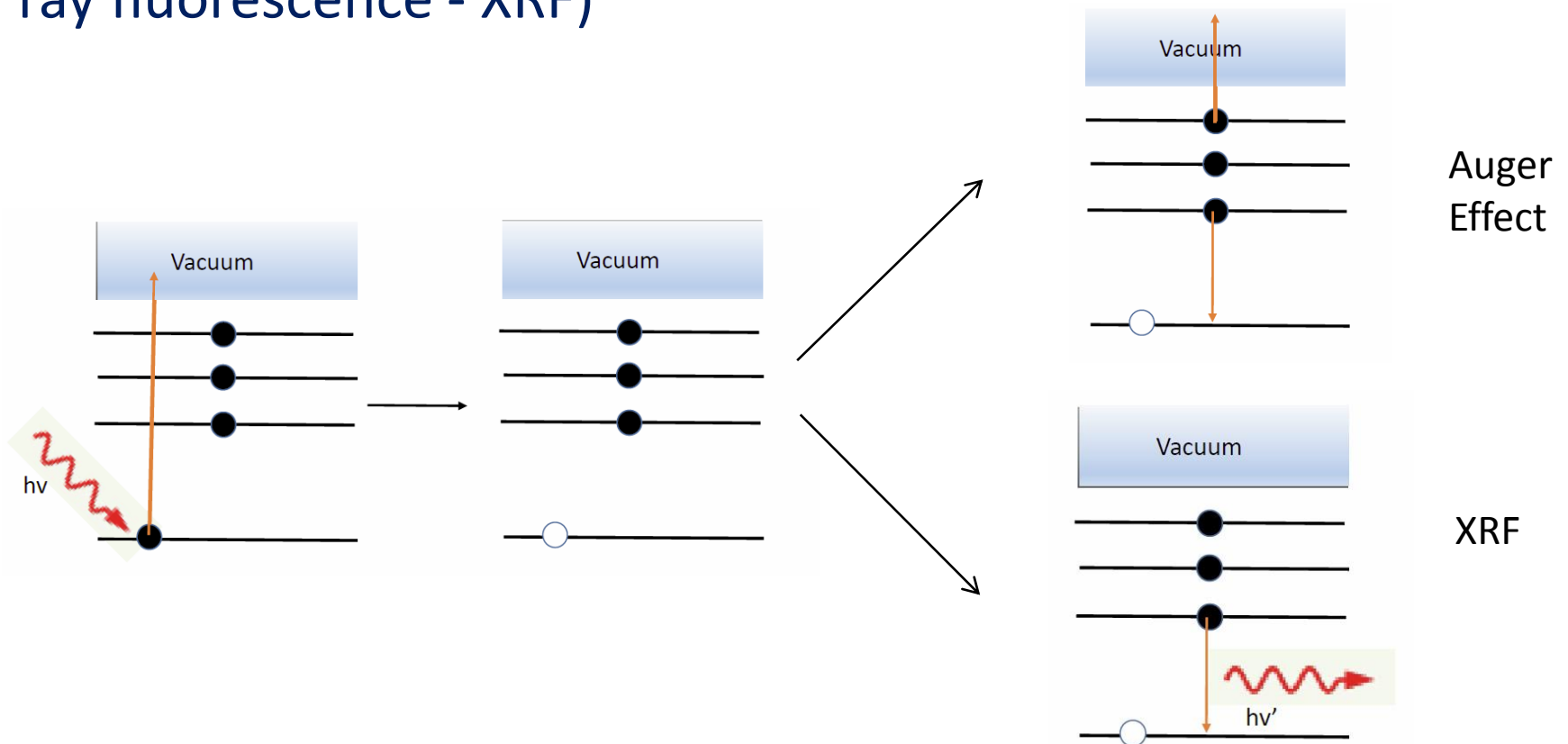


ASP Online Series 2020

Andrew Harrison, Diamond Light Source

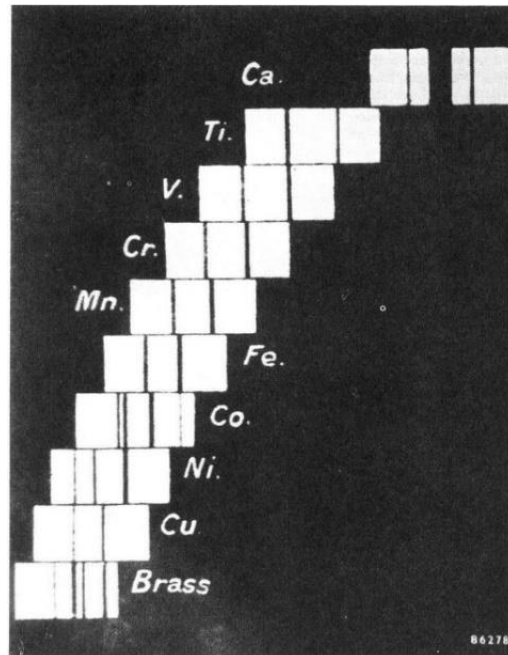
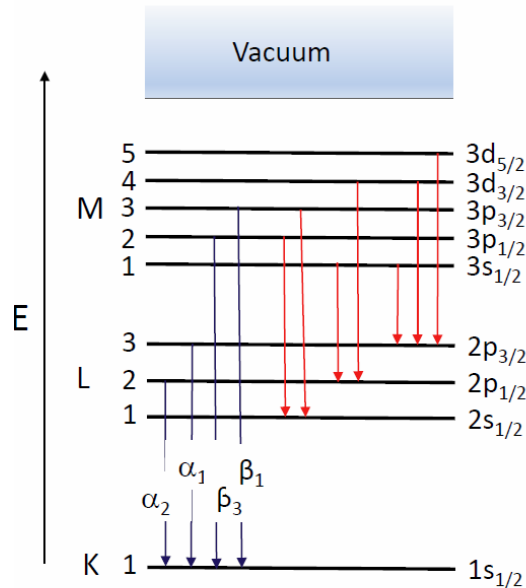
Relaxation after photoionisation ?

- Photoionisation leaves a hole behind which is electronically unstable and is rapidly filled by relaxation of higher-energy electrons and the additional energy is carried away either by a secondary photoionisation event (Auger effect) or photons (X-ray fluorescence - XRF)

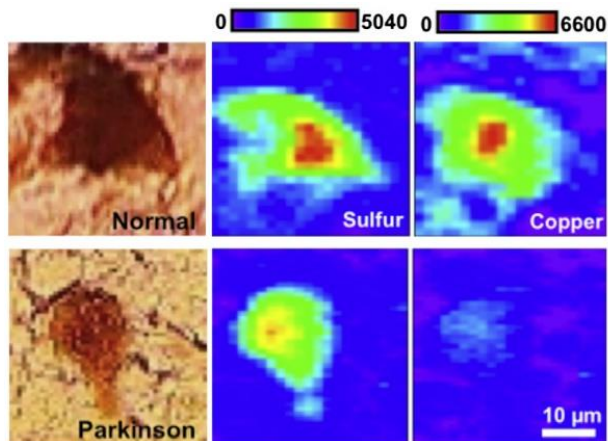
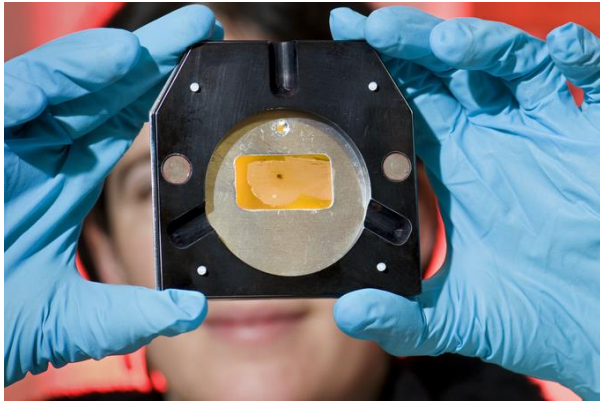


X-ray fluorescence (XRF)

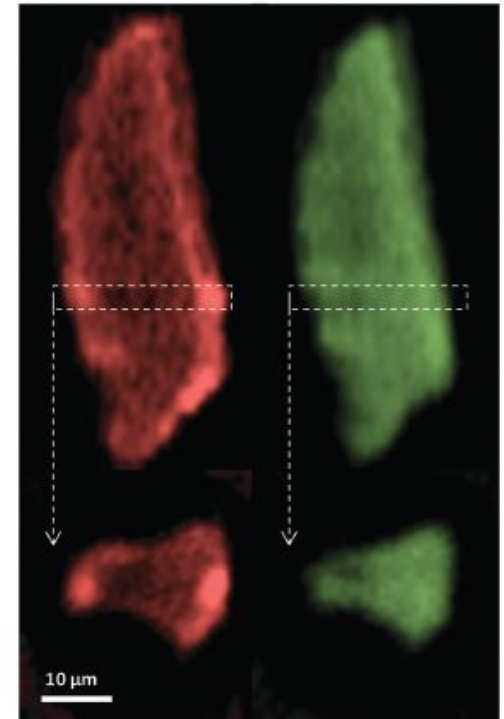
- XRF observed by Moseley who observed K transitions (relaxation into $n=1$ shell) and showed $\nu = c/\lambda = K(Z-1)^2$ where $K = 2.47 \times 10^{15}$ Hz
- Start and end states are quantised electronic states so (dipole) selection rules apply: $\Delta l = \pm 1$, $\Delta s = 0$, $\Delta j = 0, \pm 1$ but $j=0 \rightarrow j=0$ forbidden
- For K series, lowest energy transitions are α lines (split by spin orbit coupling to α_1, α_2) then β lines etc...



XRF – chemical mapping



Influence of elements on neurodegenerative disease



XRF tomography on supported catalyst particles
Mo (green) and Pt (red)

X-ray scattering

Scattering and diffraction

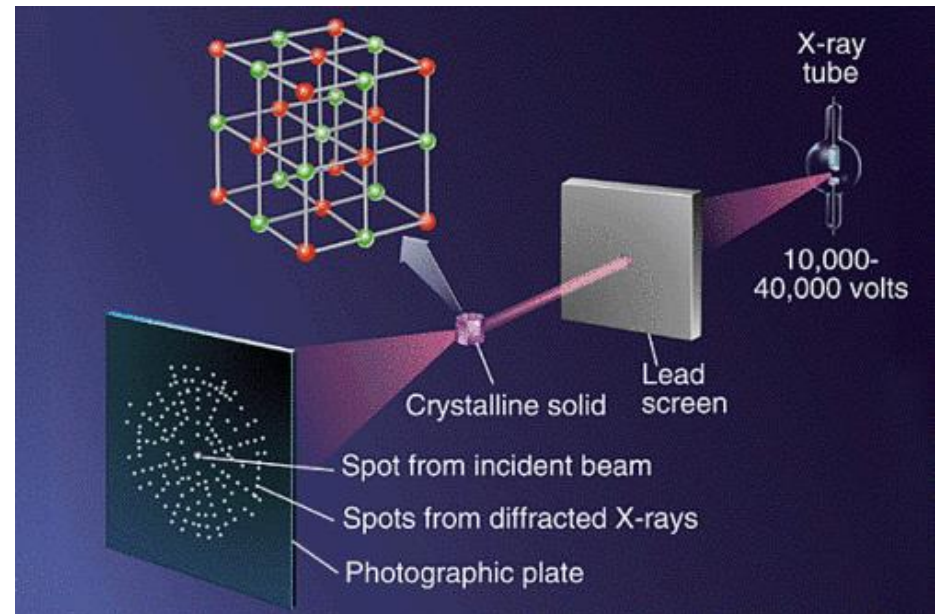
- X-ray diffraction is the most incisive technique to determine atomic structure of crystalline materials. Synchrotrons allow the technique to be exploited particularly effectively.
- Diffraction arises from the coherent interference between waves scattered by an array of objects whose spacing is comparable to wavelength (of the order of 1 \AA). In X-ray diffraction the scattering is by the electrons around the atoms in the structure.

Double Slit Diffraction



Scattering and diffraction

- X-ray diffraction is the most incisive technique to determine atomic structure of crystalline materials. Synchrotrons allow the technique to be exploited particularly effectively.
- Diffraction arises from the coherent interference between waves scattered by an array of objects whose spacing is comparable to wavelength (of the order of 1 \AA). In X-ray diffraction the scattering is by the electrons around the atoms in the structure.



Diffraction – building blocks

Recall crystallography courses you may have taken

- Scattering of X-rays from a crystal produces a set of Bragg spots (reflections) whose **position** is determined by the size and structure of the unit cell and whose **intensity** is determined by the (Fourier transform of) the distribution of electron density in the unit cell, the minimum repeat unit
- There are 7 crystal systems which may be modified further by centring to produce the 14 distinct Bravais lattices.
- Decorate the cell with atoms i whose position is defined by fractional coordinates (x_i, y_i, z_i)

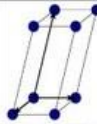
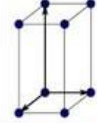
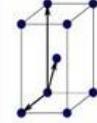
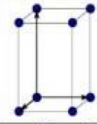
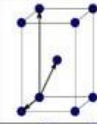
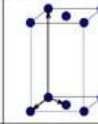

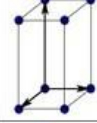

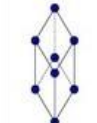
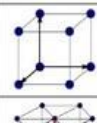
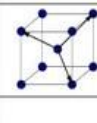
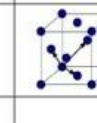
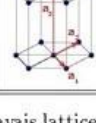
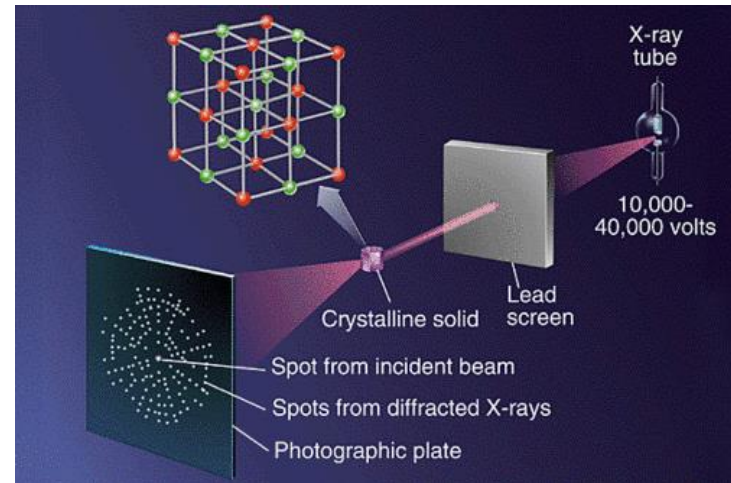
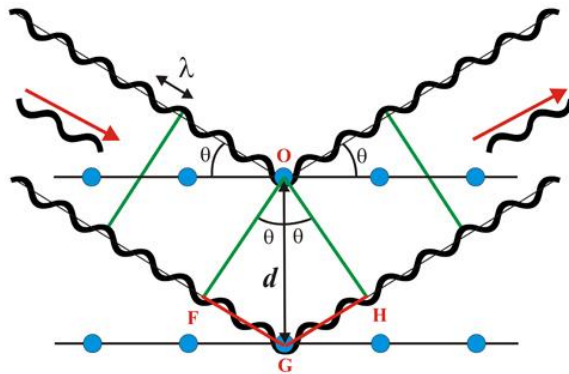
Bravais lattice	Parameters	Simple (P)	Volume centered (I)	Base centered (C)	Face centered (F)
Triclinic	$a_1 \neq a_2 \neq a_3$ $\alpha_{12} \neq \alpha_{23} \neq \alpha_{31}$				
Monoclinic	$a_1 \neq a_2 \neq a_3$ $\alpha_{23} = \alpha_{31} = 90^\circ$ $\alpha_{12} \neq 90^\circ$				
Orthorhombic	$a_1 \neq a_2 \neq a_3$ $\alpha_{12} = \alpha_{23} = \alpha_{31} = 90^\circ$				
Tetragonal	$a_1 = a_2 \neq a_3$ $\alpha_{12} = \alpha_{23} = \alpha_{31} = 90^\circ$				
Trigonal	$a_1 = a_2 = a_3$ $\alpha_{12} = \alpha_{23} = \alpha_{31} < 120^\circ$				
Cubic	$a_1 = a_2 = a_3$ $\alpha_{12} = \alpha_{23} = \alpha_{31} = 90^\circ$				
Hexagonal	$a_1 = a_2 \neq a_3$ $\alpha_{12} = 120^\circ$ $\alpha_{23} = \alpha_{31} = 90^\circ$				

Table 1.1: Bravais lattices in three-dimensions.

Diffraction - measurements

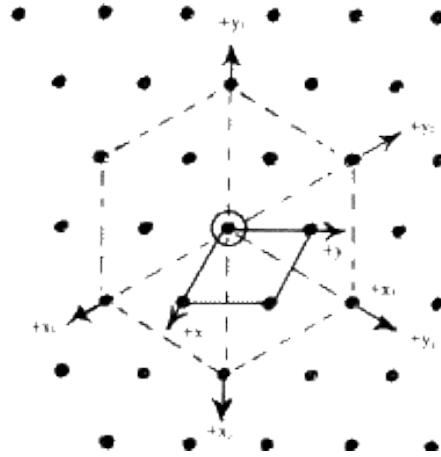
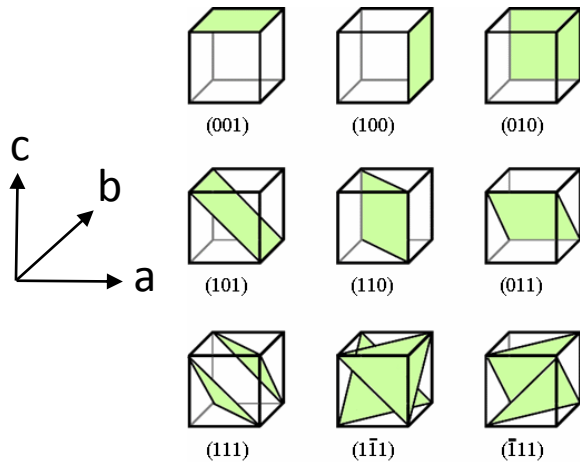
- The **position** of a Bragg spot can be determined simply by considering diffraction to be the consequence of coherent scattering of X-rays of wavelength λ from identical planes of spacing d at angle θ using the Bragg equation: $n \lambda = 2d \sin \theta$



- If the X-rays are **monochromatic**, the crystal has to be rotated so that the angle between the planes and the incident X-rays is θ .
- If the beam is **polychromatic**, containing a band of wavelengths, Bragg condition may be met for an arbitrary crystal orientation for different values of λ for different planes at different values of θ - Laue diffraction – more complex but more information and can be unravelled with detectors that discriminate between energy and hence λ

Defining and labelling planes

- The planes are defined and labelled using the so-called Miller indices



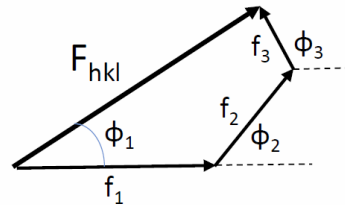
- A plane with Miller indices h , k and l makes intercepts a/h , b/k , c/l with the unit cell axes a , b and c . This is called the (hkl) plane. For crystals where a , b and c are perpendicular, $1/d^2 = h^2/a^2 + k^2/b^2 + l^2/c^2$
- An alternative representation that enables reflections to be located rapidly is to define a lattice in reciprocal (Fourier) space where the coordinate system is such that \mathbf{a}^* is \perp to both \mathbf{b} and \mathbf{c} , of length $2\pi/a$, etc. The line joining the origin to the point (hkl) in reciprocal space is perpendicular to the (hkl) plane and bisects the scattering angle 2θ

Scattering intensity

- The **amplitude** of the scattering from the (hkl) plane is then proportional to the sum of the scattering f_j from every atom j in the unit cell, called the structure factor F_{hkl} :

$$F_{hkl} = \sum_j f_j \exp[2\pi i(hx_j + ky_j + lz_j)] = \sum_j f_j \exp(i\phi_j)$$

Where $\exp(i\phi_j)$ is the phase of the scattering from atom j .



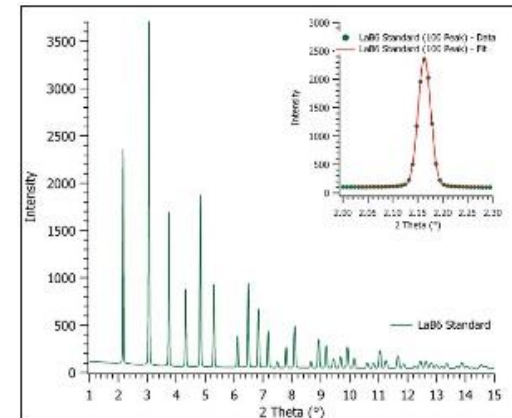
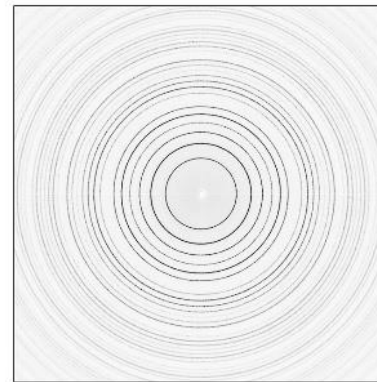
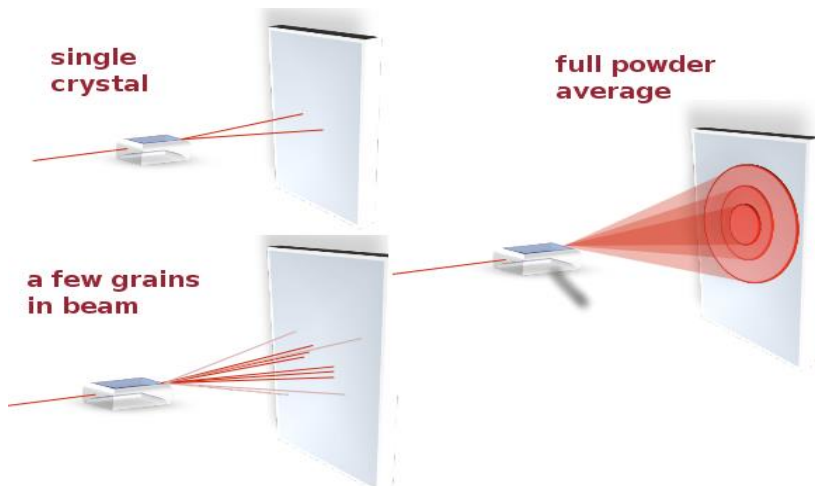
We measure the intensity in diffraction experiment, This is proportional to the square of the amplitude so we lose the information about the phase and cannot solve the structure directly – the **phase problem**

X-ray diffraction at synchrotrons

- The same principles determining position and intensity of scattering still apply to synchrotron X-rays, but they bring advantages
 - Exceptionally bright – focus down to very small beams ($< 1 \mu\text{m}$) and look at very small samples or regions of samples – though care with radiation damage (biological samples fried in seconds)
 - Follow structural features in seconds, and in some cases fractions of a second – look at dynamics
 - More energetic synchrotron X-rays (100 KeV and more) can penetrate materials much more deeply than the lower energy radiation more typical of a laboratory-based instrument (typically $\sim 8 \text{ KeV}$).
 - Energy of synchrotron X-rays may be tuned through the absorption edge of an element in the sample, altering the nature of the scattering from such elements and enable the phase problem to be solved.
 - Coherence reveals small structural deformation in nm – 10's nm regime, complementing EM techniques where samples are too thick for electron microscopy

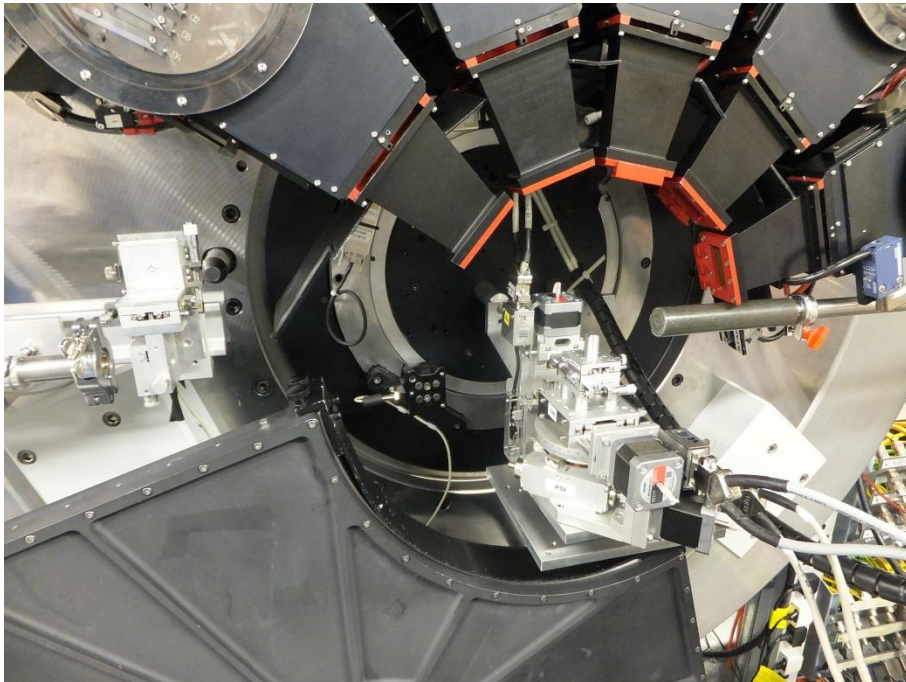
Powder diffraction

- Scattering from randomly oriented crystallites leads to an average over all possible orientations and reduces sets of discrete spots to rings characterised by scattering angle θ_{hkl} (where the indices hkl have to be determined too).
- Average over rings to produce powder diffraction pattern: I vs $2\theta_{hkl}$
 $n\lambda = 2d \sin\theta_{hkl}$ $1/d^2 = h^2/a^2 + k^2/b^2 + l^2/c^2$ for orthorhombic and intensity $\propto F_{hkl}^2$



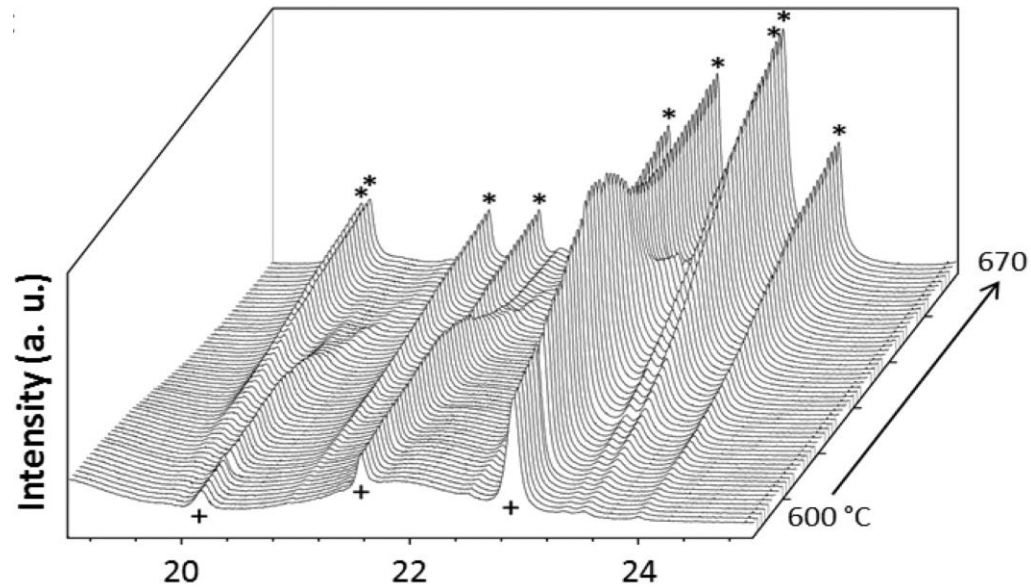
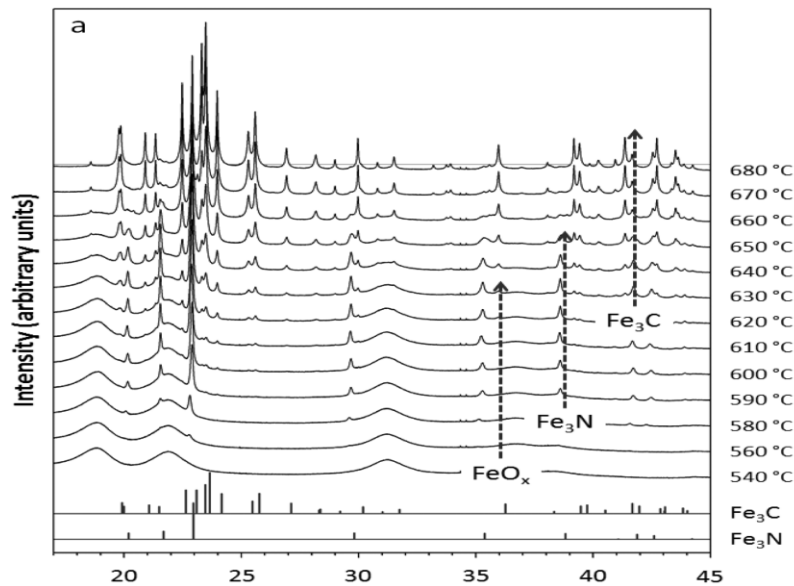
Powder diffraction

- Applications (photos show powder diffractometer I11 @Diamond:
 - Fingerprint or match with known material with very high spatial (d) resolution
 - Refine structure based on known structure (Rietveld process) until model predicts the observed powder pattern
 - Solve structure *ab initio* (rare)



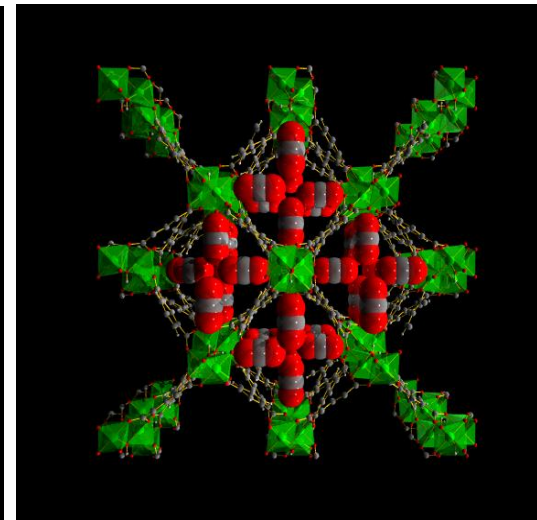
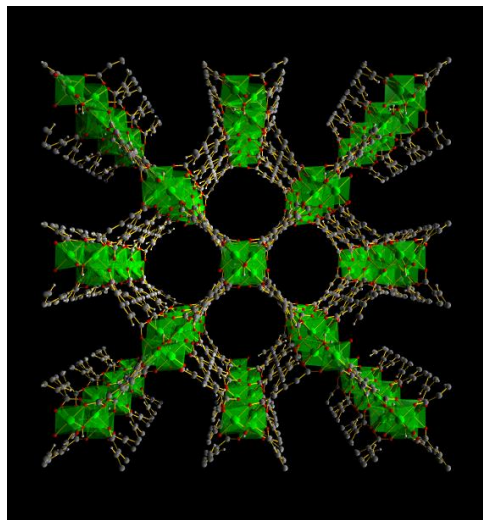
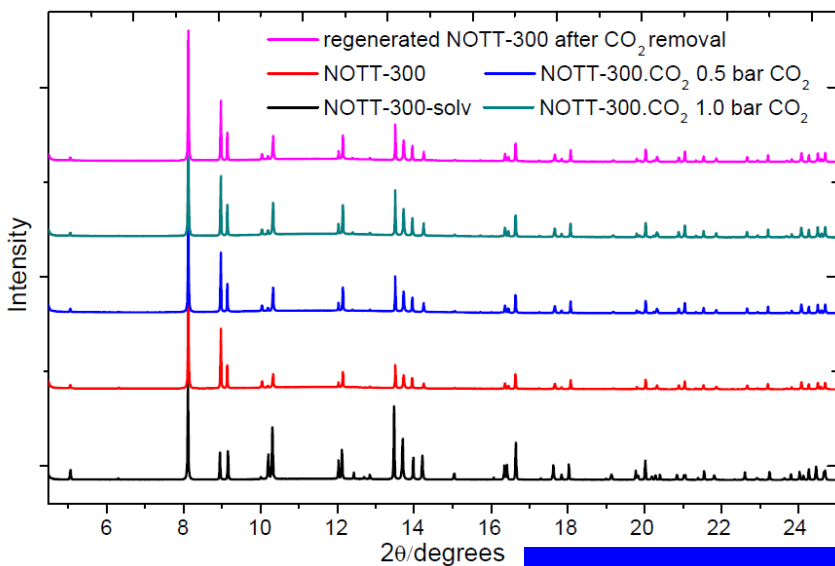
Follow processes *in operando*

- E.g. gel containing iron nitrate and gelatin heated under N_2 in a glass capillary in the X-ray beam. PSRD shows evolution of phases, identified against known samples – oxide then nitride then carbide (C from gelatine) forms



Follow processes *in operando*

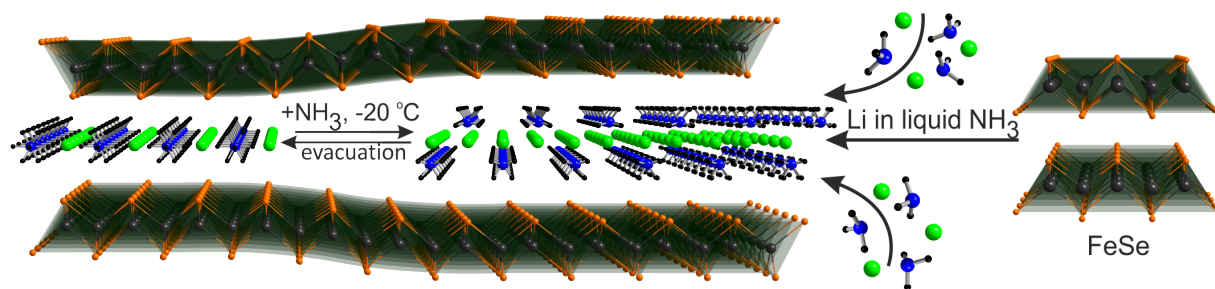
- e.g. follow insertion of gases such as CO₂ in porous materials such as MOFs *in situ*, refining structure to locate CO₂ and observe change in host material



In situ Powder Diffraction (CO₂)

Watching chemical reactions

e.g. study the chemical reaction between a host, layered chalcogenide FeSe and Li in ammonia to produce new layered superconductor

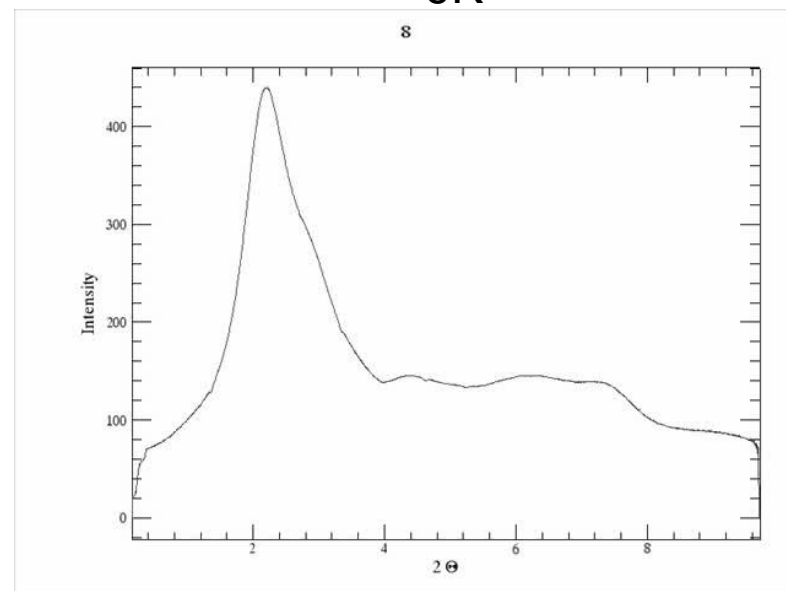


Tc

44 K

39 K

9K



Single crystal diffraction

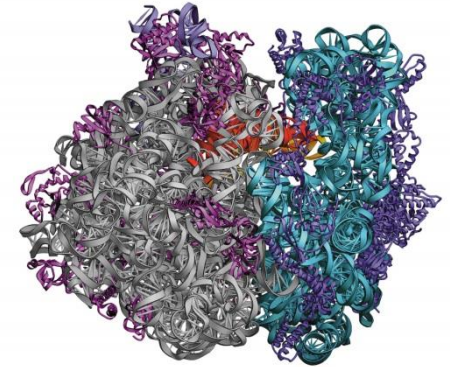
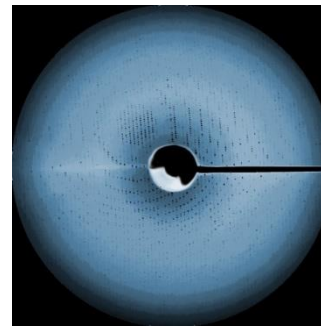
- Return to the phase problem.
- Wish to determine the electron density distribution in the unit cell:

$$\rho(xyz) = \sum_{hkl} F_{hkl} \exp[-2\pi i(hx + ky + lz)]$$

- However, we measure the intensity, which is proportional to the square of the structure factor F_{hkl} so we lose the phase information (+ F_{hkl} or - F_{hkl} ?)
- For relatively small unit cells with relatively few atoms there are fewer variables and we can sometimes crack the structure using brute force computational methods and other constraints such as common bond lengths. These are called '**direct methods**'

Failure of direct methods

- Direct methods do not work for the large macromolecules that govern life and are implicated in disease and drug discovery and can contain 100's of thousands of atoms – e.g. the work on the structure of the ribosome which was recognised by the award of the 2009 Nobel Prize in Chemistry to Ada Yonath, Venki Ramakrishnan and Thomas Steitz



- Note that scattering from an individual atom $\propto Z^2$ so heavy atoms can dominate scattering. If a natural atom such as S is substituted in the structure by Se this can help us determine the phase.

