

Jet error parametrisations for kinematic fitting

CLICdp WG Analysis meeting

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December 14, 2020



CLUSTER OF EXCELLENCE
QUANTUM UNIVERSE



Outline

Introduction

- Higgs production and decays to heavy flavour jets
- Semi-leptonic decay of heavy flavour jets
- Kinematic fitting

Higgs mass reconstruction with semi-leptonic decays

- Simple neutrino correction
- Jet-specific energy resolution

Revisiting the $H \rightarrow b\bar{b}$ mass reconstruction

- Neutrino correction from first principles
- Improved jet error parametrisation
- Application to $ZH \rightarrow \mu\bar{\mu}b\bar{b}$

Summary

- Outlook
- Conclusions



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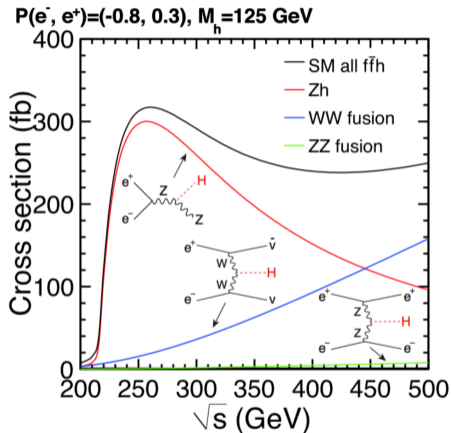
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Higgs production mechanisms at the ILC

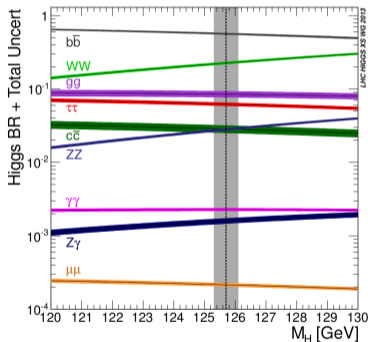


Higgs strahlung is dominant Higgs production mechanism at ILC 250 GeV

Higgs decays into $b\bar{b}$ and $c\bar{c}$

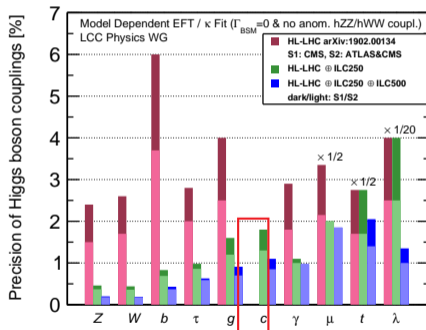
- ▶ Most frequent Higgs decay mode:

$$H \rightarrow b\bar{b}$$



- ▶ Extremely challenging in Hadron colliders:

$$H \rightarrow c\bar{c}$$



- ▶ Heavy flavour jets play key role in Higgs physics.

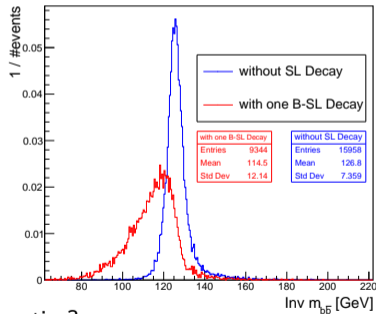
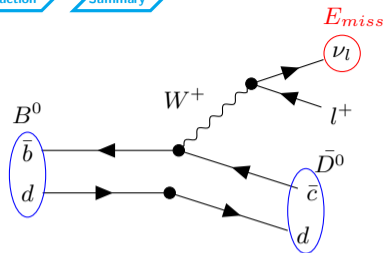


Semi-leptonic b / c decays

- Semi-leptonic decay (SLD) rate in $e^+e^- \rightarrow b\bar{b}$ samples

		nBSLD		
		0	1	2
nCSDL	0	34%	24%	4%
	1	18%	12%	2%
	2	3%	2%	0%

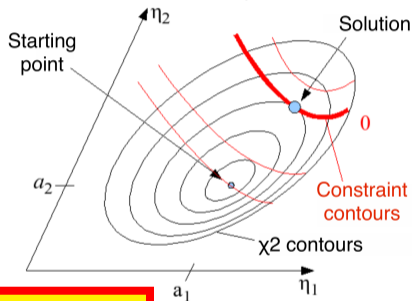
- nBSLD: number of B-hadron semi-leptonic decays
- nCSDL: number of C-hadron semi-leptonic decays
- Mis-reconstruction of $b\bar{b}$ invariant mass due to **missing neutrino energy** from semi-leptonic decays
- Can the **missing momentum** be retrieved from event and decay kinematics?



Kinematic fit

- ▶ Kinematic fit: adjustment of measured quantities under certain kinematic constraints:

- ▶ Energy and momentum conservation
- ▶ Invariant masses of particles



arXiv:0901.4656

Exploit well-known initial state in e^+e^- colliders

- ▶ Minimize χ^2 :

$$\chi^2(\mathbf{a}, \boldsymbol{\xi}, \mathbf{f}) = (\boldsymbol{\eta} - \mathbf{a})^T \mathbf{V}^{-1}(\boldsymbol{\eta} - \mathbf{a}) - 2\boldsymbol{\lambda}^T \mathbf{f}(\mathbf{a}, \boldsymbol{\xi})$$

$\boldsymbol{\eta}$: vector of measured kinematic variables (x)

\mathbf{a} : vector of fitted quantities

$\boldsymbol{\xi}$: vector of unmeasured kinematic variables

\mathbf{V} : covariance matrix

$\boldsymbol{\lambda}$: Lagrange multipliers

$\mathbf{f}(\mathbf{a}, \boldsymbol{\xi})$: vector of constraints

- ▶ Measures of performance:

$F(\chi^2; \text{ndf})$: cumulative χ^2 distribution for a certain ndf

$P(\chi^2)$: fit probability

$$\text{pull}(x) = \frac{x_{\text{fitted}} - x_{\text{measured}}}{\sqrt{\sigma_{\text{fitted}}^2 - \sigma_{\text{measured}}^2}}$$

$$P(\chi^2) = 1 - F(\chi^2; \text{ndf})$$

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Simple neutrino correction for Higgs mass reconstruction

- ▶ Neutrino energy correction:

Estimating neutrino energy as a fraction of corresponding lepton energy:

$$E_{jet}^{corr} = E_{jet} + E_{\nu} = E_{jet} + \left(\frac{1}{x} - 1\right)E_{lep}$$

- ▶ Uncertainty on jet energy parametrised as:

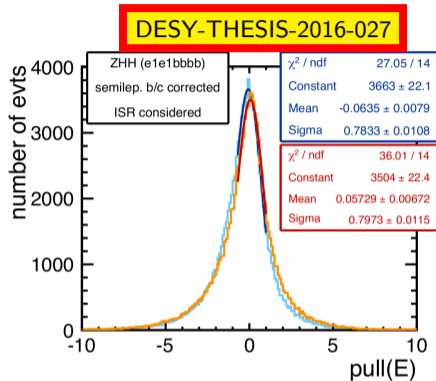
$$\sigma_{E_{jet}}^{corr} = \frac{100\%}{\sqrt{E_{jet}}} \oplus \sigma_{\nu}$$

$$\sigma_{\nu}^2 = \left(\frac{\sigma_{\langle x \rangle}}{\langle x \rangle^2}\right)^2 E_{lep}^2 + \left(\frac{1}{x} - 1\right)\Delta E_{lep}^2$$

- ▶ Fixed uncertainties on angles:

$$\Delta\theta_{jet} = \Delta\phi_{jet} = 100 \text{ mrad}$$

Simple correction to jet energy improves jet energy pull distribution as a measure of fit performance.



Blue: before neutrino energy correction

Orange: After neutrino energy correction

Jet specific energy resolution

Parametrize sources of uncertainties (assumed uncorrelated) in jet energy measurements (`ErrorFlow`):

$$\sigma_{E_{jet}} = \sigma_{Had} \oplus \sigma_{Det} \oplus \sigma_{Conf} \oplus \sigma_{\nu} \oplus \sigma_{Clus}$$

DESY-THESIS-2017-045

- ▶ QCD effects in parton shower and hadronization, σ_{Had}
- ▶ Detector resolution using track and cluster parameters, σ_{Det}
MarlinReco: `AddClusterProperties` and `FourMomentumCovMat`
- ▶ Particle confusion in PFA, σ_{Conf}
Estimated based on jet energy and neutral hadron / photon energy fractions, based on
- ▶ Semi-leptonic decays: now treated in analysis level, ν -correction
in previous: $\sigma_{\nu}^2 = \left(\frac{\sigma_{\langle x \rangle}}{\langle x \rangle^2}\right)^2 E_{lep}^2 + \left(\frac{1}{x} - 1\right) \Delta E_{lep}^2$
- ▶ Misassignment of particles in the jet clustering, σ_{Clus}

arXiv:0907.3577



Error flow and application in kinematic fit

DESY-THESIS-2017-045

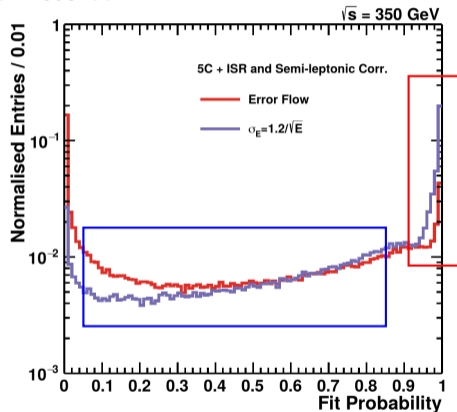
Jet specific energy resolution for $e^+e^- \rightarrow ZH \rightarrow q\bar{q}b\bar{b}$ process at $\sqrt{s} = 350$ GeV

- ▶ Full 4×4 CovMatrix on 4-momentum of jets $\sigma(\vec{p}, E)$:
 - ▶ σ_{Det} : computed using subdetector momentum/energy resolution
 - ▶ σ_{Conf} : computed using jet energy and particle content (charged, neutral and photon)
 - ▶ $\sigma_\nu = 0.73 \cdot E_l$
 - ▶ $\sigma_{Had}, \sigma_{Clus}$ are not accounted for error flow procedure yet.
- ▶ **Fixed (and wide) angular resolution:** $\sigma_\theta = \sigma_\phi = 100$ mrad

Kinematic fit: vary jet quantities (E, θ, ϕ) within uncertainties $(\sigma_E, \sigma_\theta, \sigma_\phi)$

Improved fit probability by applying Error Flow on jet energy

⇒ Further improvements by error parametrization and handling sl-decays



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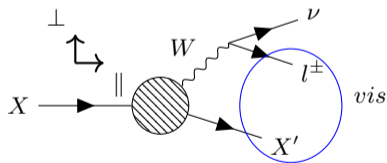
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Concept of correction in a semi-leptonic decay

- ▶ Find heavy-quark jets: Identify b or c jet \rightarrow flavour tag
- ▶ Find semi-leptonic decay(s): Identify lepton in jet if present \rightarrow possible using detector's high granularity
- ▶ Estimate neutrino energy from decay kinematics:
 - ▶ Assign B^0 or D^0 meson mass to mother hadron.
 - ▶ Reconstruct flight direction of mother hadron from position of primary and secondary vertex.
 - ▶ Calculate neutrino momentum: up to 3-fold ambiguity.
- ▶ As proof-of-principle: CHEAT from MC truth
- ▶ Lepton ID
- ▶ Flavour tag
- ▶ Mother hadron mass
- ▶ **Associate of reconstructed particles to secondary vertex**
- ▶ Momenta of visible decay products



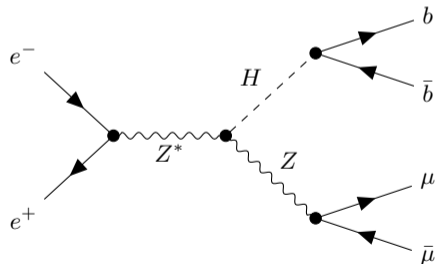
The neutrino momentum can be determined up to a two-fold ambiguity

Can we use overall event kinematics to decide between solutions? \Rightarrow kinematic fit!

$ZH \rightarrow \mu\bar{\mu}b\bar{b}$ event preparation

Overview

- ▶ Test sample for upcoming 250 GeV MC production: $e^+e^- \rightarrow ZH$ samples with $Z \rightarrow \mu\bar{\mu}$ and $H \rightarrow b\bar{b}$ (57016 events, no $\gamma\gamma \rightarrow hadrons$ overlay):
 - ▶ Event generator: WHIZARD-2
 - ▶ Software version: iLCSoft v02-01
 - ▶ Detector model: ILD_I5_v02
- ▶ Event selection:
 - ▶ $Z \rightarrow \mu\bar{\mu}$: IsolatedLeptonTagging
 - ▶ $H \rightarrow b\bar{b}$: FastJetProcessor: Durham jet clustering
- ▶ CovMatrix of jets: ErrorFlow
- ▶ Neutrino correction from semi-leptonic decays ((E, \vec{p})-based approach, 3-fold ambiguity: $E_\nu^+, E_\nu^-, 0$)
- ▶ Kinematic fit (2-jets + 2-leptons)
 - ▶ Add constraints to fit
 - ▶ Choose best ν -correction (highest fit probability)



Error parametrisations for jets and leptons

ZH11qq5CFit

- ▶ Jets parametrized by $(E_{jet}, \theta_{jet}, \phi_{jet}, \sigma_{E_{jet}}, \sigma_{\theta_{jet}}, \sigma_{\phi_{jet}}, m_{jet})$:

$$(p_x, p_y, p_z, E, \text{CovMat}(\vec{p}, E)) \rightarrow (E, \theta, \phi, \sigma_E, \sigma_\theta, \sigma_\phi, m)$$

$$\tan \theta = \frac{\sqrt{p_x^2 + p_y^2}}{p_z}, \quad \tan \phi = \frac{p_y}{p_x}, \quad m = \sqrt{E^2 - (p_x^2 + p_y^2 + p_z^2)}$$

σ_θ^2 : Error propagation of $(\sigma_{p_x}^2, \sigma_{p_z}^2, \sigma_{p_z}^2, \dots)$, σ_ϕ^2 : Error propagation of $(\sigma_{p_x}^2, \sigma_{p_z}^2, \dots)$

σ_θ and σ_ϕ are not fixed \Rightarrow Angular resolutions are calculated for each individual jet

Planned TO DO: input full CovMatrix to fit objects

- ▶ Leptons are parametrized by $(\frac{1}{p_T}, \theta_{lep}, \phi_{lep}, \sigma_{\frac{1}{p_T}}, \sigma_{\theta_{lep}}, \sigma_{\phi_{lep}}, m_{lep})$:

$$(\Omega, \tan \lambda, \phi, \text{CovMatrix}(\text{track}), m) \rightarrow (\frac{1}{p_T}, \theta, \phi, \sigma_{\frac{1}{p_T}}, \sigma_\theta, \sigma_\phi, m)$$

$$\frac{1}{p_T} = \frac{\Omega}{eB}, \quad \theta = \frac{\pi}{2} - \lambda \Rightarrow \sigma_{\frac{1}{p_T}} = \frac{1}{eB} \sigma_\Omega, \quad \sigma_\theta = \frac{\sigma_{\tan \lambda}}{1 + \tan^2 \lambda}$$

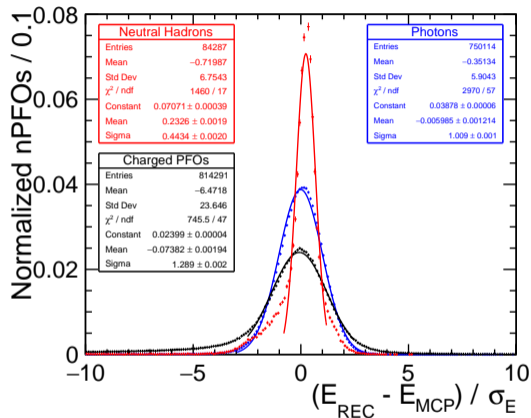
Parameters are obtained from tracks \Rightarrow precisely measured parameters



Uncertainty in PFO level (Energy)

Situation of error estimation in PFO level

- ▶ Photons: energy error is perfectly modeled.
- ▶ Charged PFOs: possible improvement by tracks refitted with true mass hypothesis (need to be checked!).
Charged PFOs are investigated through track parameters
agenda.linearcollider.org/event/8341
- ▶ Neutral Hadrons: energy and energy error are overestimated.



The energy error of the neutral PFOs is directly obtained from cluster energy error. \Rightarrow Not easy to maneuver!

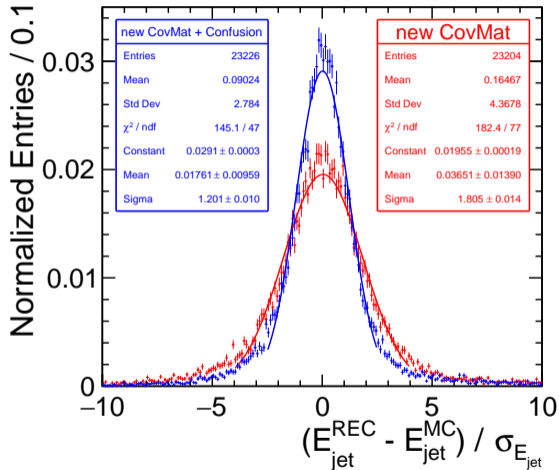
The PFOs energy error quadratically summed up \Rightarrow jet energy error: $\sigma_{E_{jet}}^2 = \sum \sigma_{E_{PFO}}^2 + [\sigma_{conf}^2]$

Uncertainties in jet-level: Energy

Jet-level truth:

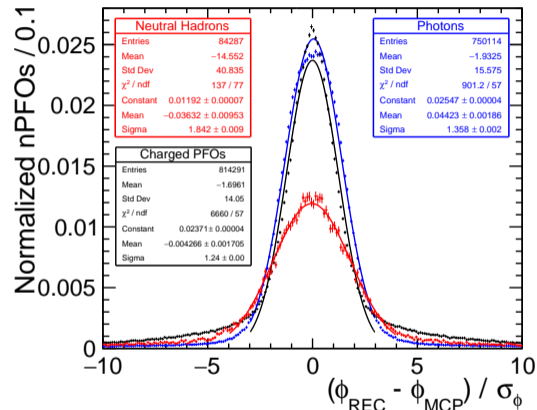
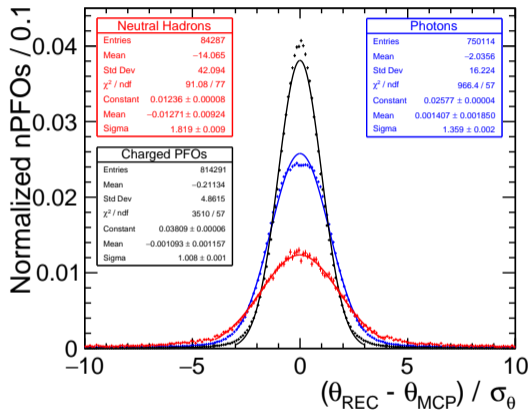
- ▶ All stable particles in MCParticleSkimmed collection (generator status = 1)
- ▶ Remove ISR photons
- ▶ Find and remove isolated muons
look for stable $\mu/\bar{\mu}$ in daughter chain of initial muon pair
- ▶ include/exclude ν_s from semileptonic decay (optional)
- ▶ Run jet clustering algorithm

Confusion term improves estimating uncertainty on the jet energy



Uncertainty in PFO level (Angles)

Obtained directly from track parameters / cluster position errors



The angular uncertainties of the neutral PFOs are obtained from cluster position error.

Underestimated σ_θ and σ_ϕ for photons (factor ~ 1.3) and neutral hadrons (factor ~ 1.8)

CovMat of Jets

- ▶ AddClusterProperties/FourMomentumCovMat: CovMat(cluster/track) \rightarrow CovMat(\vec{p} , E)

- ▶ Current CovMat calculation (inconsistent 4-momentum of neutral hadrons):

$$E_{PFO} = |\vec{p}_{PFO}| = E_{clu}, \quad p_x = E_{clu} \frac{x}{r}, \quad p_y = E_{clu} \frac{y}{r}, \quad p_z = E_{clu} \frac{z}{r}, \quad m_{PFO} = m_n$$

- ▶ Alternative CovMat calculation (taking consistent 4-momentum of neutral hadrons)

$$E_{PFO} = \sqrt{|\vec{p}_{PFO}|^2 + m_{PFO}^2} = \sqrt{E_{clu}^2 + m_n^2}$$

$J(\text{wrong}) \rightarrow J(\text{right})$

$$\begin{pmatrix} E_{clu} \frac{r^2 - x^2}{r^3} & -E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xz}{r^3} & 0 \\ -E_{clu} \frac{xy}{r^3} & E_{clu} \frac{r^2 - y^2}{r^3} & -E_{clu} \frac{yz}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ \frac{x}{r} & \frac{y}{r} & \frac{z}{r} & 1 \end{pmatrix}_{\text{wrong}} \rightarrow \begin{pmatrix} E_{clu} \frac{r^2 - x^2}{r^3} & -E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xz}{r^3} & 0 \\ -E_{clu} \frac{xy}{r^3} & E_{clu} \frac{r^2 - y^2}{r^3} & -E_{clu} \frac{yz}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ \frac{E}{E_{clu}} \cdot \frac{x}{r} & \frac{E}{E_{clu}} \cdot \frac{y}{r} & \frac{E}{E_{clu}} \cdot \frac{z}{r} & 1 \end{pmatrix}_{\text{right}}$$

- ▶ ErrorFlow:

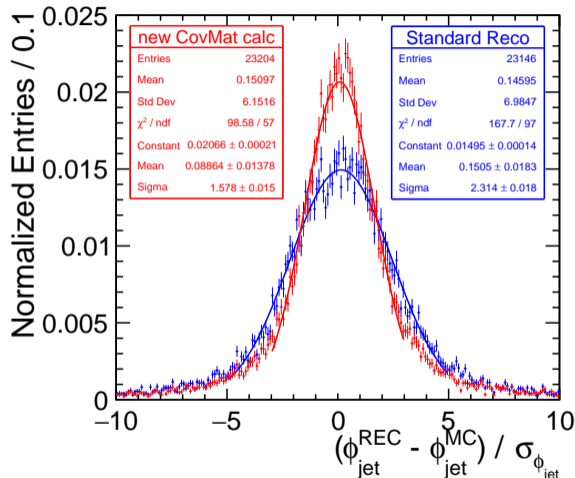
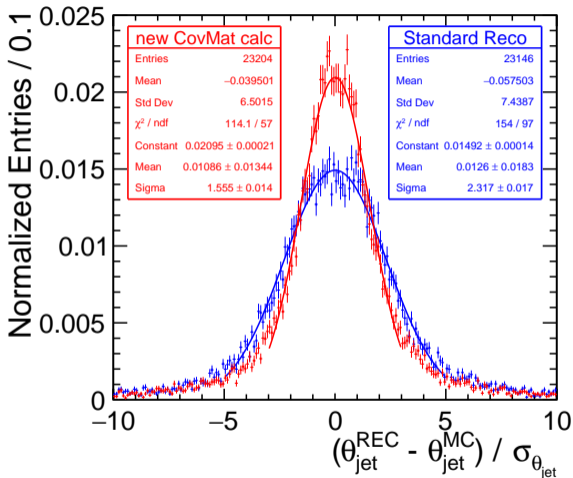
$$\text{CovMat}(\vec{p}_{jet}, E_{jet}) = \sum_{PFO} \text{CovMat}(\vec{p}, E) \quad : \quad \sigma_{E_{jet}}^2 = \sigma_{conf}^2 + \sum_{PFO} \sigma_{E_{PFO}}^2$$

- ▶ MarlinKinfitProcessors:

$$\text{CovMat}(\vec{p}_{jet}, E_{jet}) \rightarrow (\sigma_{\theta_{jet}}, \sigma_{\phi_{jet}}, \sigma_{E_{jet}})$$



Uncertainties in jet-level: θ & ϕ



By new CovMat, normalized residuals of jet angles improved by 50%! \Rightarrow improvedCovMat

5C Kinematic fit

in 2-leptons + 2-jets events

Parameters of jets and leptons are varied within their uncertainties to satisfy 5 constraints:
Conservation of momentum (hard constraints):

- ▶ p_x : e^+e^- crossing angle: 14 mrad
 $\Sigma p_x = \sqrt{s} \times \sin 0.007 \approx 1.75 \text{ GeV}$
- ▶ p_y : $\Sigma p_y = 0$
- ▶ p_z : $\Sigma p_z = 0$

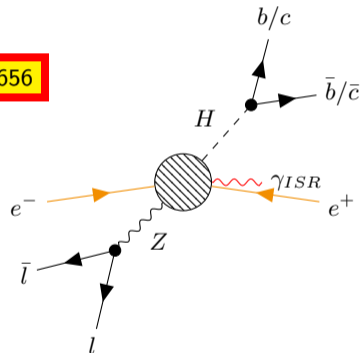
arXiv:0901.4656

Conservation of total energy (hard constraint):

- ▶ $E_{lab} = 2\sqrt{\left(\frac{\sqrt{s}}{2}\right)^2 + (\Sigma p_x)^2}$

Constrain di-muon mass to agree with m_Z within its natural width (soft constraint):

- ▶ $m_Z = 91.2 \text{ GeV}$, $\sigma_{m_Z} = \frac{2.4952}{2}$

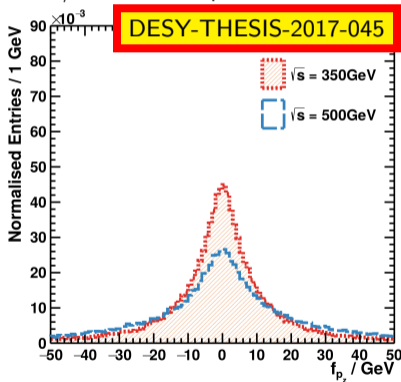


Fit constraints

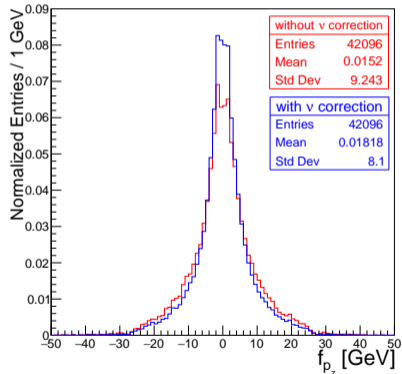
Momentum conservation: p_z

Adding 4-momentum of neutrino improves jet fit object initialization

► DBD 350/500 GeV samples



► MC-2020 250 GeV prod. samples



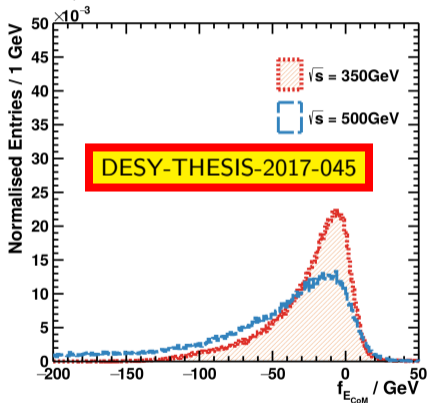
Proper neutrino correction for jets: improved constraint on momentum

fit constraints

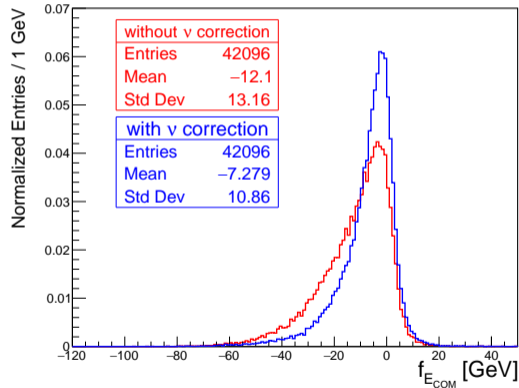
energy conservation: E

Neutrino correction (best pre-fit \vec{p}_ν for succesful fits) improves start values \Rightarrow better fit object initialization

- ▶ DBD 350/500 GeV samples



- ▶ MC-2020 250 GeV prod. samples



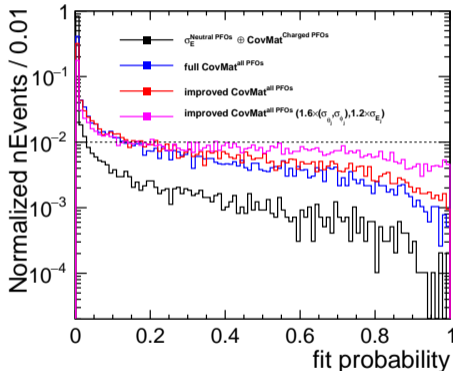
By neutrino correction, initial value of constraint function closer to target \Rightarrow fit should work better!

Kinematic fit performance

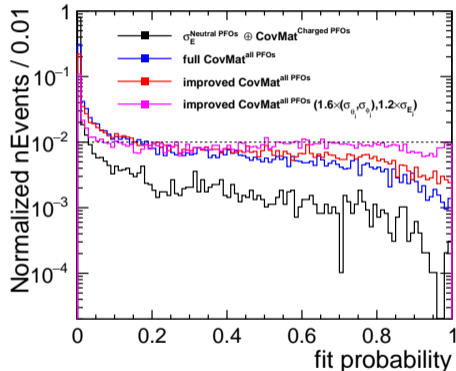
fit probability

$e^+e^- \rightarrow ZH \rightarrow \mu\bar{\mu}b\bar{b}$ at $\sqrt{s} = 250$ GeV (without semi-leptonic decays)

▶ without confusion term in ErrorFlow



▶ with confusion term in ErrorFlow



Improved fit probability with: right Jacobian (improved CovMat of Neutral PFOs) + including confusion term
 flat fit probability with scaling jet uncertainties \Rightarrow further improvement with tracks refitted with true mass!

Higgs mass with ν -correction

Reconstructed m_H :

Higgs mass by standard reconstruction: $m_H = \text{Inv. } M_{b\bar{b}}$

ν correction:

Higgs mass by standard reconstruction with

$$\vec{p}_{jet}^{corr} = \vec{p}_{jet} + \vec{p}_\nu$$

Run fit for all \vec{p}_ν solutions and use best pre-fit \vec{p}_ν

Kinfit (with improved CovMat + confusion):

Kinematic fit without ν correction

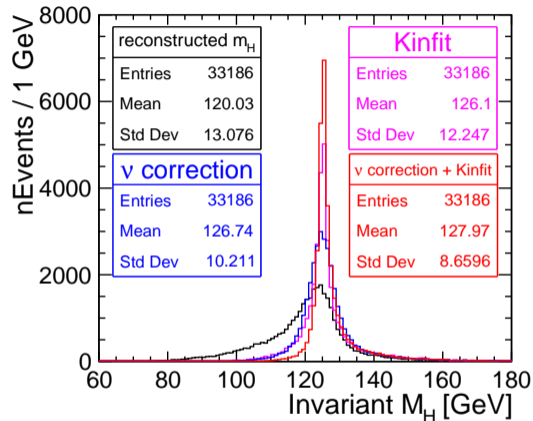
ν correction + Kinfit:

Kinematic fit with corrected jet momenta

$$\vec{p}_{jet}^{corr} = \vec{p}_{jet} + \vec{p}_\nu$$

- ▶ Improvement by individual ν -correction or kinematic fit is good but not enough
- ▶ Applying kinematic fit and ν -correction gives huge improvement on Higgs mass reconstruction
- ▶ ν correction still partially cheated

all $H \rightarrow b\bar{b}$'s



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Outlook

open issues:

- ▶ Further improve error estimation:
 - ▶ Full (E, \vec{p}) covariance matrix for JetFitObject (include $\sigma_{p_x E}, \sigma_{p_y E}, \sigma_{p_z E}$)
 - ▶ Estimate parton shower & hadronisation effects
 - ▶ Use proper masses, momenta and CovMatrices of PFOs from tracks refitted with correct mass hypothesis
 - ▶ Move neutrino correction from cheated to fully-reconstructed

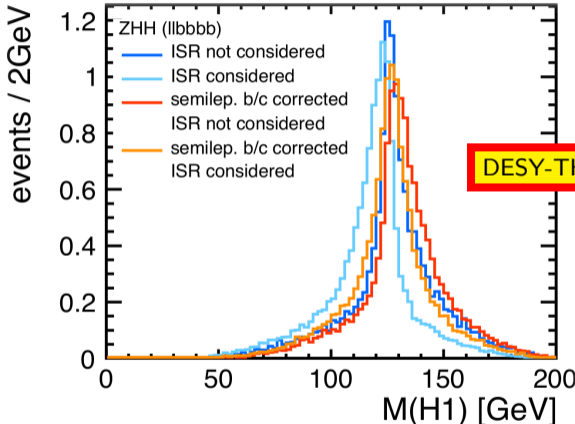
Conclusions

- ▶ Heavy flavour jets are essential for Higgs physics
- ▶ Correction of semi-leptonic decays of heavy flavour jets is important for Higgs mass reconstruction
- ▶ Kinematic fit provides a useful framework for:
 - ▶ Improving Higgs mass reconstruction / resolution
 - ▶ Better understanding on source of uncertainties / errors
- ▶ New in MarlinKinFit: uncertainties on jet angles are parametrized for individual jets
- ▶ Cheated ν correction shows very interesting potential \Rightarrow try on reco level!

BACKUP



Simple neutrino correction for Higgs mass reconstruction (Cntd.)



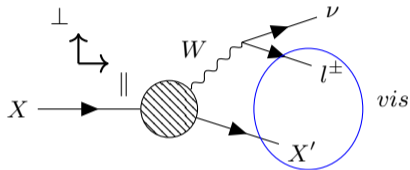
- Bias and assymetry in m_H is removed by correcting jet energy and adding ISR



correcting neutrino energy

4-vector based approach

- ▶ (E, \vec{p}) -based approach



$$\vec{p}_{\nu, \perp} = -\vec{p}_{vis, \perp}$$

$$\vec{p}_{\nu, \parallel} = \frac{1}{2D} (-A \pm \sqrt{A^2 - BD}) \hat{n}$$

$$A = p_{vis, \parallel} (2p_{vis, \perp}^2 + m_{vis}^2 - m_X^2)$$

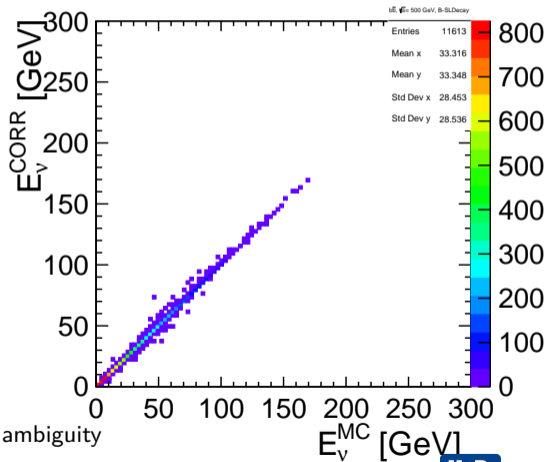
$$B = 4p_{vis, \perp}^2 E_{vis}^2 - (2p_{vis', \perp}^2 + m_{vis}^2 - m_X^2)^2$$

$$D = E_{vis}^2 - p_{vis, \parallel}^2$$

$$\hat{n} = \frac{p_{vis, \parallel}}{|p_{vis, \parallel}|}$$

The neutrino momentum can be determined up to a two-fold ambiguity

- ▶ closure test: apply correction with fully cheated information and compare with true neutrino energy

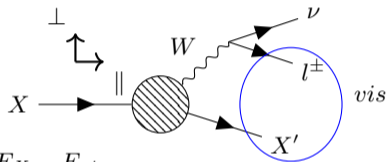


Correcting neutrino energy

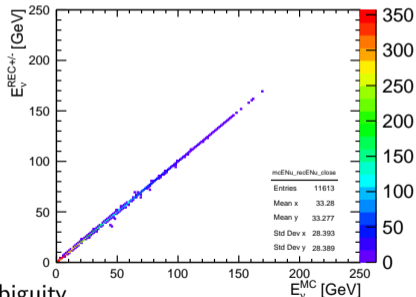
Rapidity based approach

Rapidity under Lorentz-transformations \sim velocity under Galileo-transformations: $\omega = \omega_X + \omega'$; $\omega = \frac{1}{2} \ln \frac{E+p'_{\parallel}}{E-p'_{\parallel}}$

ω : rapdity in lab frame , ω' : rapdity in rest frame of X , ω_X : rapdity of X in lab frame



- Closure test: fully cheated information ($e^+e^- \rightarrow b\bar{b}$ at $\sqrt{s} = 500$ GeV)



$$E_\nu = E_X - E_{vis}$$

$$E_X = \frac{E_{vis}E'_{vis} - p_{vis\parallel}p'_{vis\parallel}}{m_{vis}^2 + p_{vis\perp}^2} m_X$$

$$E'_{vis} = \frac{m_X^2 + m_{vis}^2}{2m_X}$$

$$p'_{vis\parallel} = \pm \sqrt{\left(\frac{m_X^2 - m_{vis}^2}{2m_X}\right)^2 - p_{vis\perp}^2}$$

The neutrino momentum can be determined up to a two-fold ambiguity

Can we use overall event kinematics to decide between solutions? \Rightarrow kinematic fit!



Event selection

Select $e^+e^- \rightarrow ZH \rightarrow \mu\bar{\mu}b\bar{b}$ events at $\sqrt{s} = 250$ GeV with (exactly) 2-leptons + 2-jets final state:

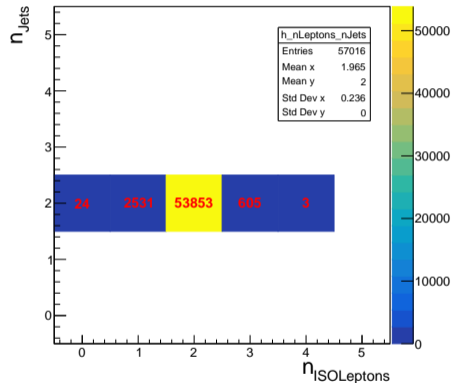
► IsolatedLeptonTagging

Training for the IDR 500 GeV samples is used,

1. Lepton ID: μ^\pm
Deposited energy in subdetectors
2. Vertex: primary or secondary
Significance of impact parameters (d_0 , z_0)
3. Isolated: not belong to jets

► FastJetProcessor

- Exclusive k_t (Durham) algorithm (no overlay)
- Find smallest of (d_{ij}, d_{iB})
$$d_{ij} = 2 \min(E_i^2, E_j^2)(1 - \cos \theta_{ij})$$
 i, j : particles, B : Beam
- $d_{ij} < d_{iB}$: combine i & j as pseudojet(p): $p_i + p_j$
- $d_{iB} < d_{ij}$: remove particle i from list
- Repeat iteration until d_{ij} or $d_{iB} > d_{cut}$ (threshold)

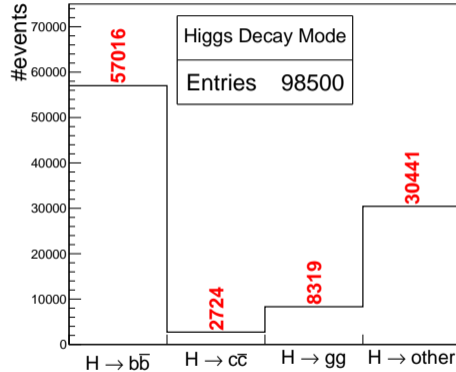


IsolatedLeptonTagging has not been trained for new software at 250 GeV yet!



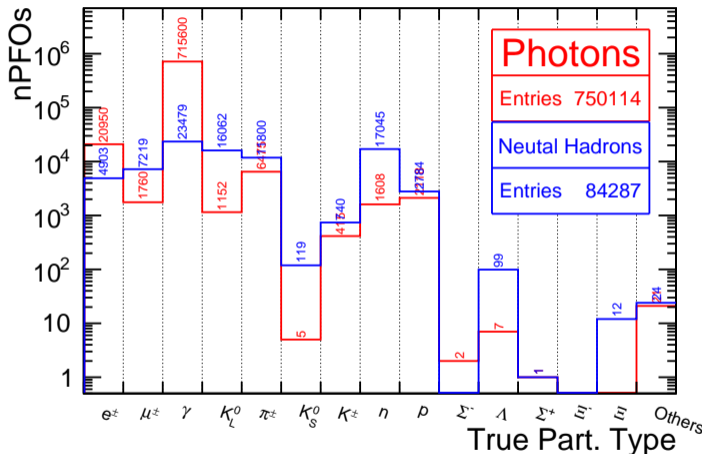
event selection

separate Higgs decay modes: $H \rightarrow b\bar{b}$, cheat from MCTruth



$\frac{2}{3}$ of $b\bar{b}$ jets contain at-least one semi-leptonic decay \Rightarrow Frequent $H \rightarrow b\bar{b}$ needs neutrino correction.

Neutral PFO identification by Pandora



Majority of identified photons are true photons.

No explicit decision for mass of identified neutral hadrons due to their multiplicity.

Pandora treatment with Neutral Hadrons

What Pandora does:

- ▶ Cluster energy is assigned to PFO(massless) energy
 $E_{PFO} = |\vec{p}_{PFO}| = E_{cluster}$
- ▶ Neutral Hadrons are identified as neutron
- ▶ neutron mass is set for PFO \Rightarrow **inconsistent 4-momentum!**
- ▶ CovMat of Neutral PFO is calculated (using inconsistent 4-momentum):
 $\text{CovMat}(\vec{p}, E) = J^T \text{CovMat}(\vec{x}_{clu}, E_{clu}) J$

$$J = \begin{pmatrix} \frac{\partial p_x}{\partial x_c} & \frac{\partial p_y}{\partial x_c} & \frac{\partial p_z}{\partial x_c} & \frac{\partial E}{\partial x_c} \\ \frac{\partial p_x}{\partial y_c} & \frac{\partial p_y}{\partial y_c} & \frac{\partial p_z}{\partial y_c} & \frac{\partial E}{\partial y_c} \\ \frac{\partial p_x}{\partial z_c} & \frac{\partial p_y}{\partial z_c} & \frac{\partial p_z}{\partial z_c} & \frac{\partial E}{\partial z_c} \\ \frac{\partial p_x}{\partial E_c} & \frac{\partial p_y}{\partial E_c} & \frac{\partial p_z}{\partial E_c} & \frac{\partial E}{\partial E_c} \end{pmatrix}$$

CovMat(\vec{p}, E) of Neutral PFOs depend on the mass assumption.

Suggestion: Take consistent 4-momentum of massive neutral hadrons for CovMat calculations.

CovMat of Neutral PFOs

- ▶ Current CovMat calculation (MarlinReco/Analysis/AddClusterProperties)

$$E_{PFO} = |\vec{p}_{PFO}| = E_{clu}, \quad p_x = E_{clu} \frac{x}{r}, \quad p_y = E_{clu} \frac{y}{r}, \quad p_z = E_{clu} \frac{z}{r}$$

- ▶ Alternative CovMat calculation (taking consistent 4-momentum of neutral hadrons)

$$E_{PFO} = \sqrt{|\vec{p}_{PFO}|^2 + m_{PFO}^2} = \sqrt{E_{clu}^2 + m_n^2}$$

$$J = \begin{pmatrix} E_{clu} \frac{r^2 - x^2}{r^3} & -E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xz}{r^3} & 0 \\ -E_{clu} \frac{xy}{r^3} & E_{clu} \frac{r^2 - y^2}{r^3} & -E_{clu} \frac{yz}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ \frac{x}{r} & \frac{y}{r} & \frac{z}{r} & 1 \end{pmatrix} \rightarrow J = \begin{pmatrix} E_{clu} \frac{r^2 - x^2}{r^3} & -E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xz}{r^3} & 0 \\ -E_{clu} \frac{xy}{r^3} & E_{clu} \frac{r^2 - y^2}{r^3} & -E_{clu} \frac{yz}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ \frac{E}{E_{clu}} \cdot \frac{x}{r} & \frac{E}{E_{clu}} \cdot \frac{y}{r} & \frac{E}{E_{clu}} \cdot \frac{z}{r} & 1 \end{pmatrix}$$

using error propagation, PFO angular uncertainties are calculated directly from cluster position error:

$$\sigma_\theta^2 = \left(\frac{\partial\theta}{\partial x}\right)^2 \sigma_x^2 + \left(\frac{\partial\theta}{\partial y}\right)^2 \sigma_y^2 + \left(\frac{\partial\theta}{\partial z}\right)^2 \sigma_z^2 + \frac{\partial\theta}{\partial x} \frac{\partial\theta}{\partial y} \sigma_{xy} + \frac{\partial\theta}{\partial x} \frac{\partial\theta}{\partial z} \sigma_{xz} + \frac{\partial\theta}{\partial y} \frac{\partial\theta}{\partial z} \sigma_{yz}$$

$$\sigma_\phi^2 = \left(\frac{\partial\phi}{\partial x}\right)^2 \sigma_x^2 + \left(\frac{\partial\phi}{\partial y}\right)^2 \sigma_y^2 + \frac{\partial\phi}{\partial x} \frac{\partial\phi}{\partial y} \sigma_{xy}$$

MUST: angular and energy uncertainties remain unchanged!

Neutrino correction hypothesis

- ▶ Assign semi-leptonic decays to jets

- ▶ Add neutrino momentum to 4-momentum of assigned jet:

Test three hypothesis for neutrino energy in each semi-leptonic decay: E_ν^+ , E_ν^- , 0
 3^{nSLD} combination of E_ν 's for adding to jet 4-momentum:

Number of semileptonic decays in a jet: $nSLD = nSLDB + nSLDC$

Example:

If an event contains two jets: jet-1 contains 2 semi-leptonic decays and jet-2 contains 1 semi-leptonic decay, **27**(= $3^2 \times 3^1$) combinations of E_ν 's are available for neutrino correction in the event:

- ▶ jet-1:

comb.	1	2	3	4	5	6	7	8	9
$\vec{p}_{\nu,1}$	-	+	0	-	+	0	-	+	0
$\vec{p}_{\nu,2}$	-	-	-	+	+	+	0	0	0

- ▶ jet-2:

comb.	1	2	3
$\vec{p}_{\nu,3}$	-	+	0

$\vec{p}_{\nu,1} + \vec{p}_{\nu,2}$ is added to 4-momentum of jet-1 and $\vec{p}_{\nu,3}$ is added to 4-momentum of jet-2.

$\vec{p}_{\nu,1} + \vec{p}_{\nu,2} + \vec{p}_{\nu,3} = 0$ allows fitter to neglect neutrino correction

Combination with highest fit probability is chosen as best neutrino correction.

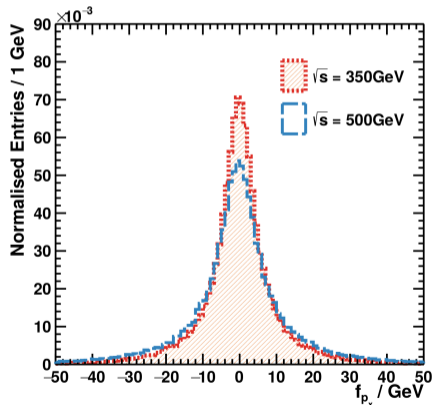


fit constraints

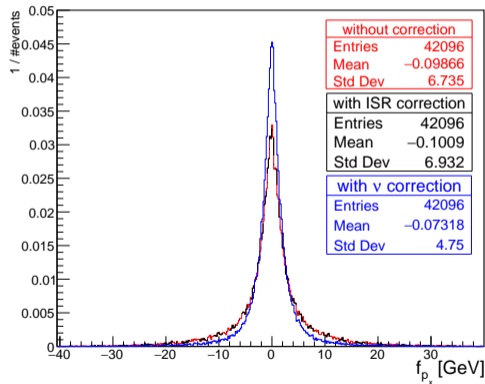
momentum conservation: p_x

ISR is initialized to satisfy momentum conservation on x direction

► by error flow on jet energy (Ali)



► by error flow on CovMatrix (new)



angular resolution for individual jets: improved constraint on momentum conservation

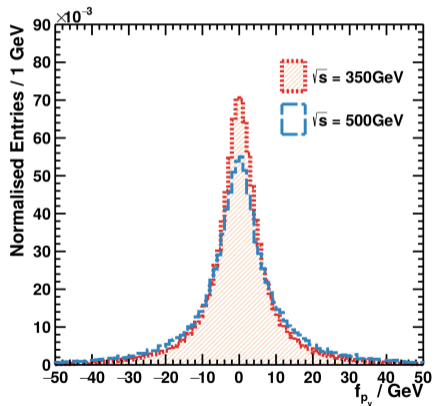


fit constraints

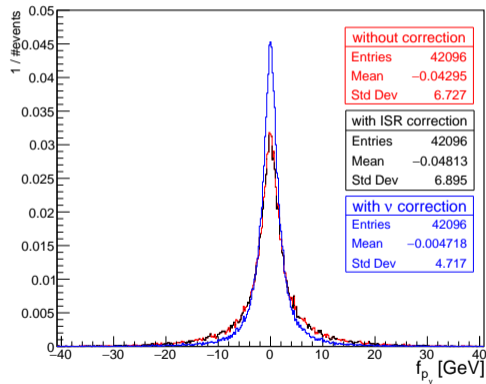
momentum conservation: p_y

ISR is initialized to satisfy momentum conservation on z direction

► by error flow on jet energy (Ali)



► by error flow on CovMatrix (new)



angular resolution for individual jets: improved constraint on momentum conservation



fit probability

