

Charmonium production with Remler generalized coalescence model

Denys Yen Arrebato Villar

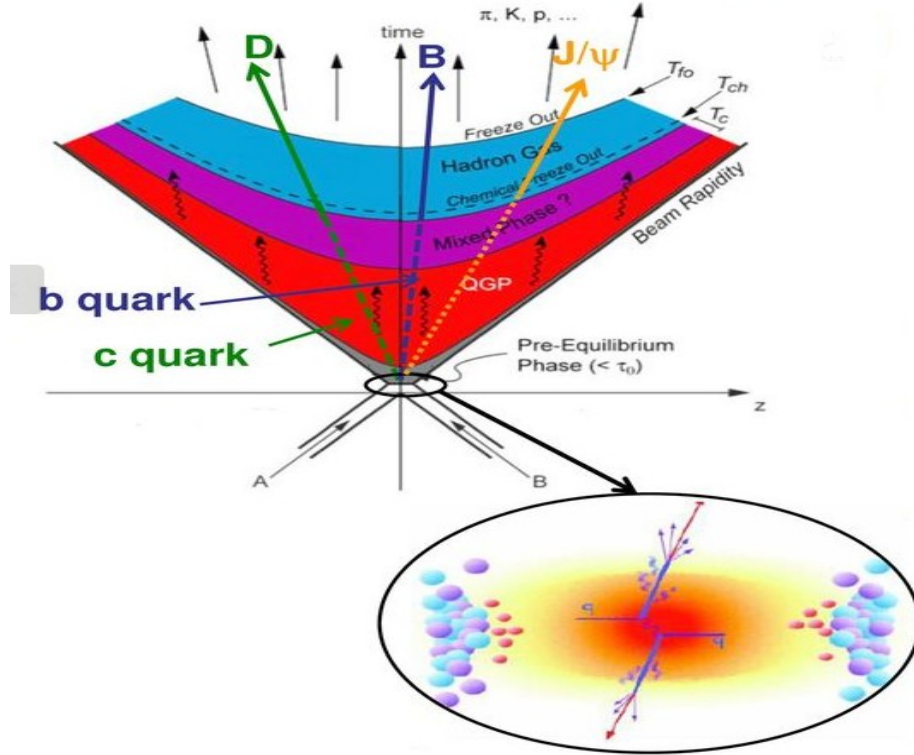
in collaboration with

Pol Bernard Gossiaux and Joerg

Aichelin



Background



For that the QGP is study by through different probes. Among the most promising probes are the heavy Quarkonium (bound state of a quark and its respective antimatter quark). Due to their heavy mass ($m_c = 1.5$ GeV and $m_b = 4.18$ GeV) they are only created at early stage of QGP in hard processes and travels through the medium interacting and leave without reaching thermal equilibrium.

Need for theoretical and phenomenological models to understand the experimental data

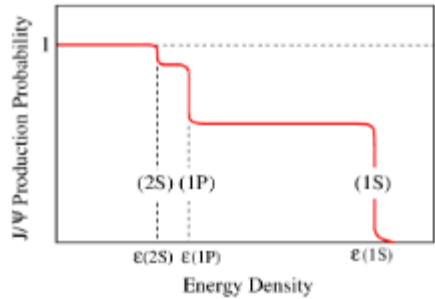
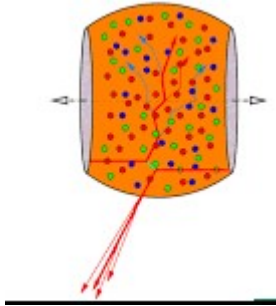
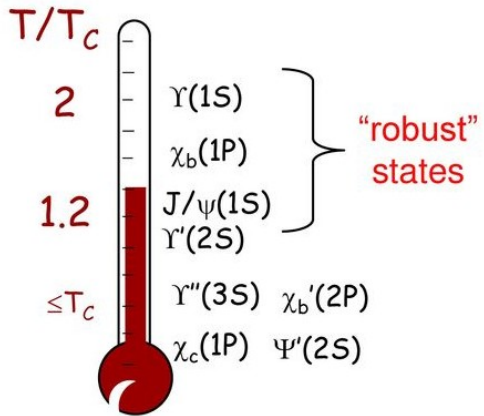
Background

Quarkonium as Hard Probe

Thermometer

Transport properties

Medium density



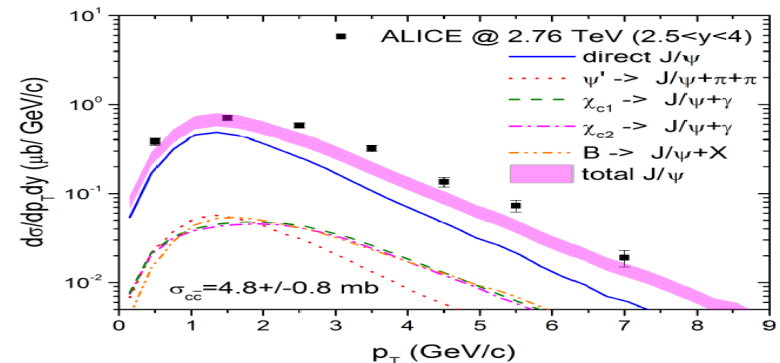
Model History

The idea of the formalism goes back to Remler's work in which a general formula connecting composite particle cross section with time-dependent density operators was presented. The formalism is able to deal with many particles (nucleons \rightarrow deuterium)

E.A. Remler, ANNALS OF PHYSICS 136, 293-316 (1981)

The model was also applied to Quarkonium production in pp and heavy ion collision(only primordial). And for the case of pp collision the model was able to reproduce the experimental data.

Taesoo .S, J.Aichelin and E.Bratkovskaya , Physical Review C 96. 014907 (2017)



Motivation

However in the same contribution was pointed out that for heavy ion collision a considerable enhancement of primordial (in the initial state) J/Ψ was found when QGP effects are ignored.

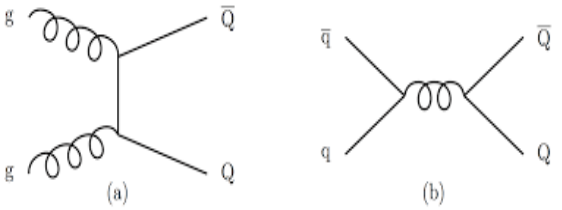


Apply the formalism to Quarkonium production in heavy ion collisions

- Interaction of heavy quarks with the bulk particles
- Expansion of the medium
- Off-diagonal contributions

How does it work

HQ production



The probability of a Quarkonium state formation in the medium is given by

$$P^\Psi(t) = \text{Tr}[\rho_{Q\bar{Q}}^\Psi \rho_N(t)]$$

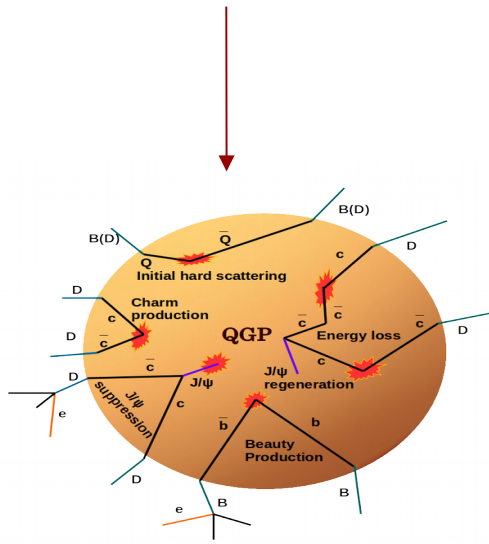
$$\rho_{Q\bar{Q}}^\Psi = \sum_i |\Psi_{Q\bar{Q}}^i\rangle \langle \Psi_{Q\bar{Q}}^i|$$

Two-body density matrix

$$\frac{\partial \rho_N(t)}{\partial t} = -i[H_N, \rho_N(t)]$$

N-body system density matrix

Quarkonium interaction with QGP



How does it work

The effective rate for Quarkonium state creation(dissociation) in the medium will be

$$\Gamma^\Psi(t) = \frac{\partial P^\Psi(t)}{\partial t} = \text{Tr}[\rho_{Q\bar{Q}}^\Psi \frac{\partial \rho_N(t)}{\partial t}]$$

Working in the phase space through Wigner function

Semi-classical approach

$$W^\Psi = \int d^3 y e^{ipy} \left\langle r - \frac{y}{2} \left| \Psi^i \right\rangle \left\langle \Psi^i \left| r + \frac{y}{2} \right. \right\rangle$$

Double Gaussian approximation
(Harmonic oscillator)

$$W_N = \prod_i \hbar^3 \delta(x_i - x_{i0}(t)) \delta(p_i - p_{i0}(t))$$

$$W_{Q\bar{Q}}^\Psi(r_{rel}, p_{rel}) = C e^{-r_{rel}^2/\sigma^2} e^{-\frac{p_{rel}^2}{\sigma^2}}$$

The Gaussian width σ

$$\frac{\int e^{-\frac{r^2}{\sigma^2}} r^2 d^3 r}{\int e^{-\frac{r^2}{\sigma^2}} d^3 r}$$



$$\langle r^2 \rangle$$



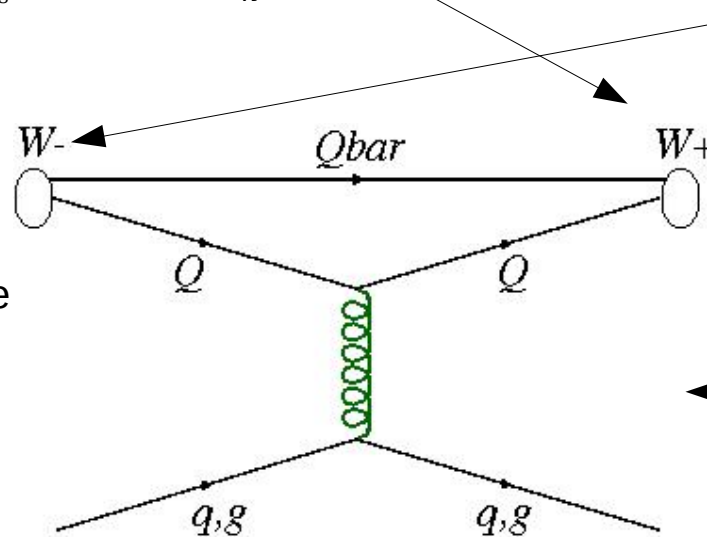
$$\left[\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right] \Psi_{Q\bar{Q}}(r) = E_{Q\bar{Q}} \Psi_{Q\bar{Q}}$$

How does it work

Substituting the expression for the Wigner functions associated to the density operator of Quarkonium state and the N-body system

$$\Gamma(t) = \sum_{i=1,2} \sum_{j \geq 3} \delta(t - t_{ij}(\epsilon)) \int \frac{d^3 p_i d^3 x_i}{h^3} [W_{Q\bar{Q}}^\Psi(p_1, x_1; p_2, x_2) W_N(t + \epsilon) - W_{Q\bar{Q}}^\Psi(p_1, x_1; p_2, x_2) W_N(t - \epsilon)]$$

- The Quarkonium production in our model is a three body process, the HQ(anti-quark) interact only by collision !!
- Very good results for D (bound state of c-quark or anti-quark + light quarks) and B (b-quark or anti-quark + light quarks) mesons production



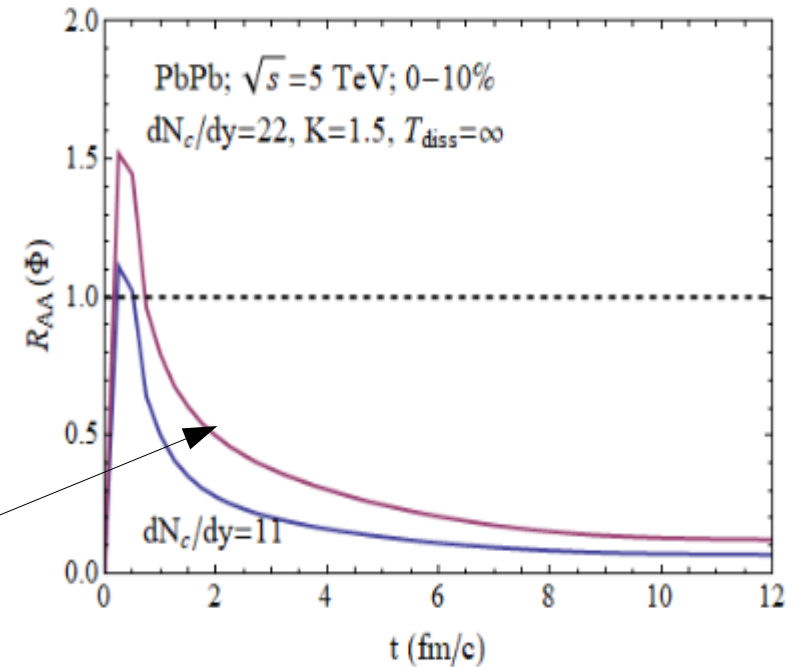
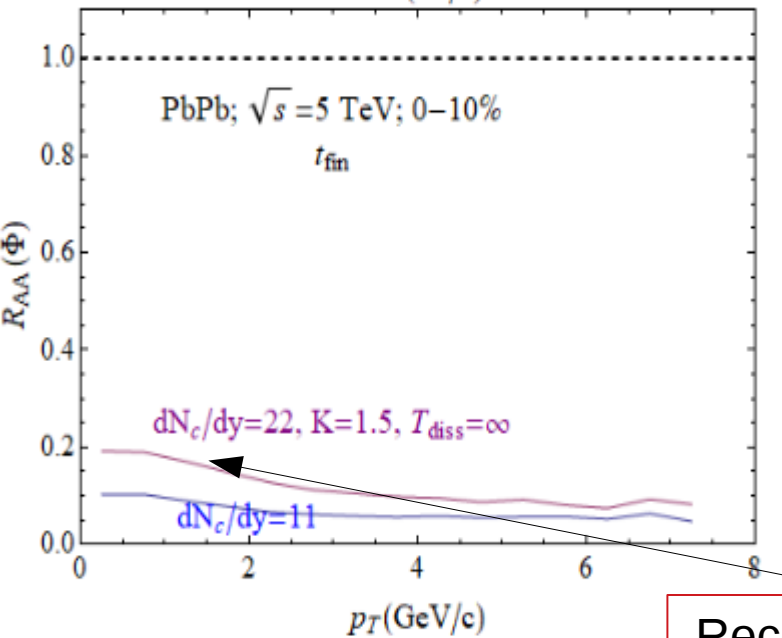
Interaction of HQ with the QGP were carried out by EPOSHQ

$$P^\Psi(t) = P^\Psi(t_{init}^\Psi) + \int_{t_{init}^\Psi}^t \Gamma(t) dt$$

Preliminary Results J/Psi

Effect of primordial production (uncorrelated $c\bar{c}$) $\frac{dN_c}{dy}$

$$R_{AA}(p_T, y) = \frac{\frac{dN^{AA}}{d^2 p_T dy}}{\langle N_{coll} \rangle \frac{dN^{pp}}{d^2 p_T dy}}$$



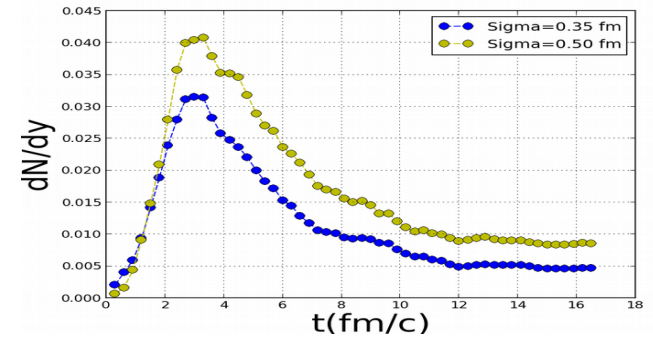
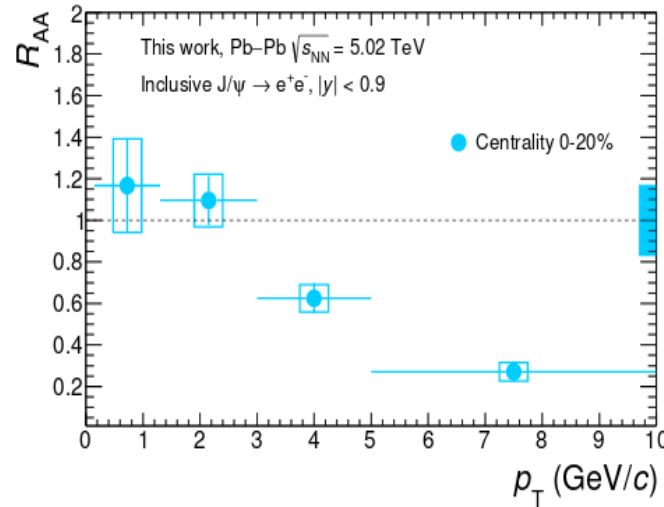
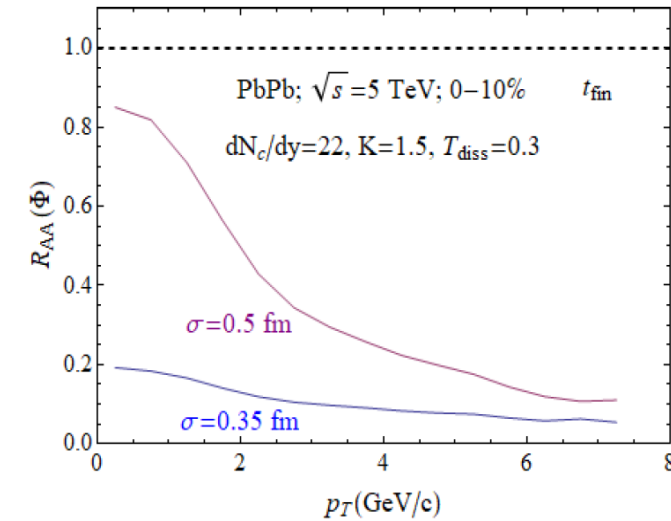
Recombination effects

Preliminary Results J/Psi

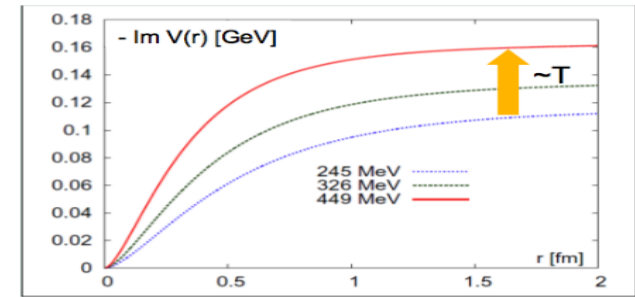
Temperature effects (uncorrelated $c\bar{c}$)

PhD Thesis : Raúl Tonaituh Jiménez
ALICE Collaboration

Mid-rapidity and $\sqrt{s}=2.76\text{ TeV}$



PhD Thesis : Nirupam Dutta

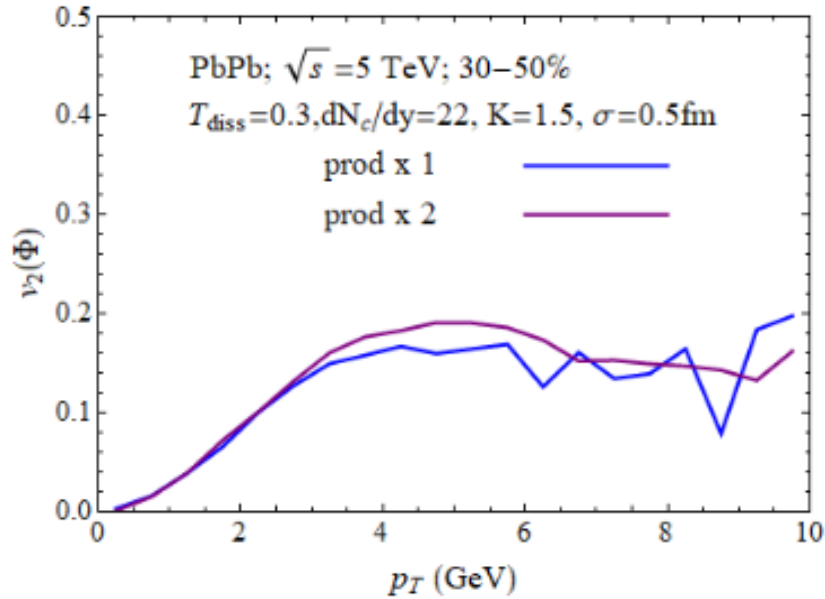


Need for temperature dependent Wigner function

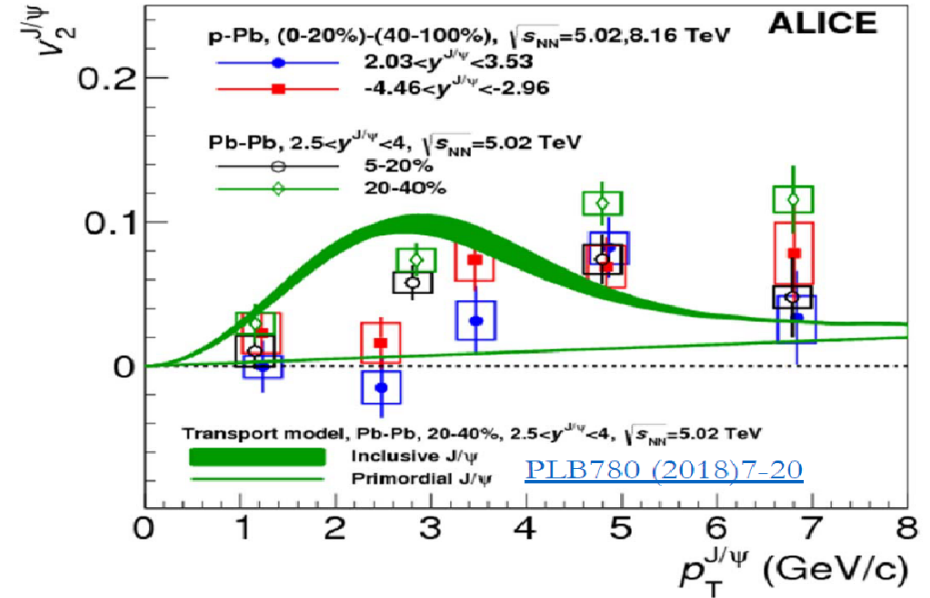


Preliminary Results J/Psi

Elliptic flow (uncorrelated c cbar)

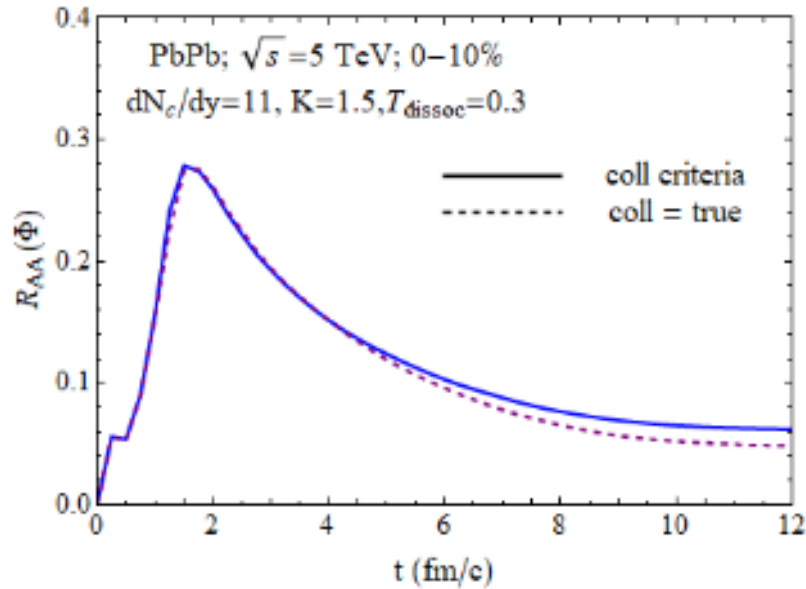


v_2 that extend at rather large p_T

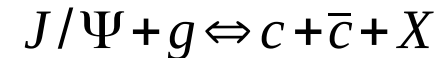


Preliminary Results J/Psi

Effects of collision criteria (uncorrelated $c\bar{c}$)



- The effect of collision increase with time together with the thermalization degree
- As the fireball expand in time most of the collision leads to a dissociation process



Conclusions

We have presented a model based on the probability density operator that allows us to obtain the time evolution of the formation probability of Quarkonium through an effective rate in QGP.

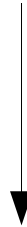
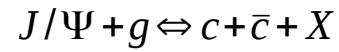
According with the results, the temperature dependence of the Quarkonium radius needs to be taken into account via optical potential with a temperature dependence (currently on going)

While trying to go for higher p_T (higher energy Quarkonium formation) relativistic effects such retarded potential should also be implemented

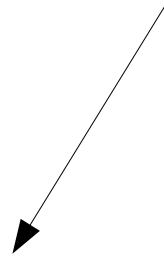
Thanks!!!

Back up

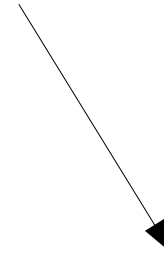
The detail balance law



Rates Equation

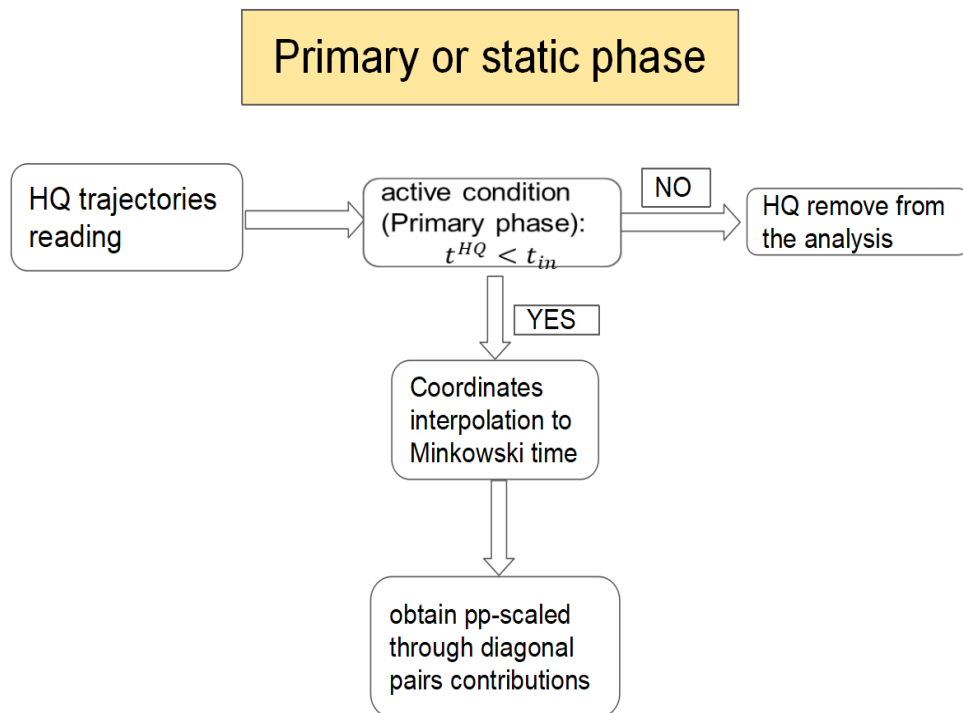


$$\frac{dN_{\Psi}}{d\tau} = \Gamma_{recomb} N_c N_{\bar{c}} [V_{FB}(\tau)]^{-1} - \Gamma_{diss} N_{\Psi}$$



$$\frac{dN_{\Psi}}{d\tau} = -N_{\Psi} L\tau + G(\Psi) = \frac{-1}{\tau_{\Psi}} [N_{\Psi}(\tau) - N_{\Psi}^{eq}(\tau)]$$

Analysis implementation



Analysis implementation

