

KYUSHU SANGYO UNIVERSITY

Modification of critical current density properties in high-T_c superconductors by tuning columnar defect morphologies in different directions

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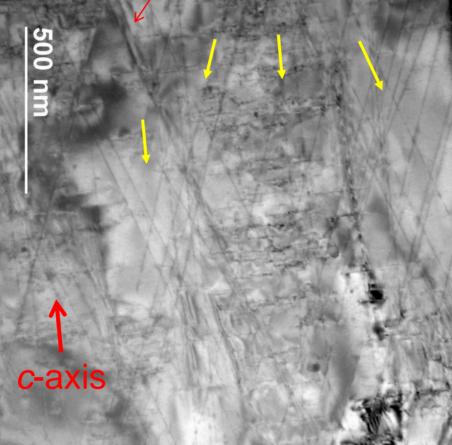
Supercurrent

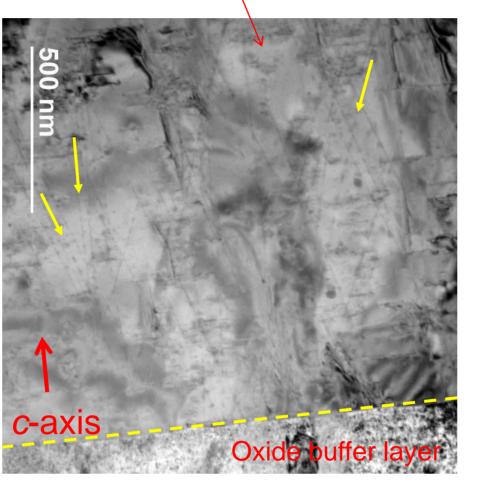
Mechanism of the formation of CDs by the thermal spike model wift heavy ions GdBCO 1 lon energy coated conductor **Electronic** excitation 2 The electronic 3 An amorphous region along the ion path is formed excitation energy Thermal energy CDs Discontinuous CDs of lattice

 $\theta_i = 0^\circ$ and $\pm 20^\circ$ relative to the *c*-axis

Near the surface

| 2.2 µm |





000 nm

2000 nm

14.9

Morphology of CDs at a depth from the surface

Electronic stopping power S_e [keV/nm] at a depth from the surface

at a depth nom the surface					
Depth from surface	$\theta_{\rm i} = 0^{\circ}$	$\theta_{\rm i} = \pm 20^{\circ}$	$\theta_{\rm i} = \pm 45^{\circ}$	$\theta_{\rm i} = \pm 80^{\circ}$	De s
500 nm	Discont.	Continuous	Continuous	Continuous	
1000 nm	Discont.	Discont.	Continuous	_	5
1500 nm	Discont.	Discont.	Discont.	_	1(
2000 nm	Discont.	Discont.	Discont.	_	1:

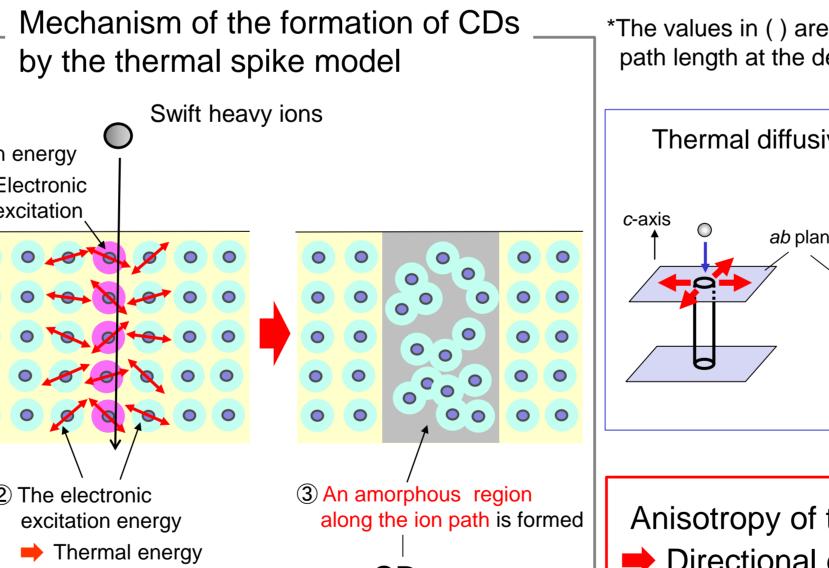


Table Samples in this study.

measurement	Sample	$ heta_{i}$	Irradiated Ions, [MeV] / <i>B</i> _{\u036} [T]	Measurement(s)
ai, Japan	GdBCO	0°, ±20°	Xe, 80 / 0.5 × 3	TEM
urement)	GdBCO	0°, ±45°	Xe, 80 / 0.5 × 3	TEM
-AXIS Irradiated	GdBCO	0°, ±80°	Xe, 80 / 0.5 × 3	TEM
θ_i	YBCO	0 °	Xe, 200 / 1.5	SQUID- $J_{\rm c}$
	YBCO	±45°	Xe, 200 / 0.75 \times 2	SQUID- $J_{\rm c}$
	YBCO	0 °	Kr, 50 / 1.5	TEM / SQUID- $J_{ m c}$
	YBCO	$\pm 45^{\circ}$	Kr, 50 / 0.75 $ imes$ 2	TEM / SQUID- $J_{\rm c}$
<i>B ab</i> plane =				

3. Results

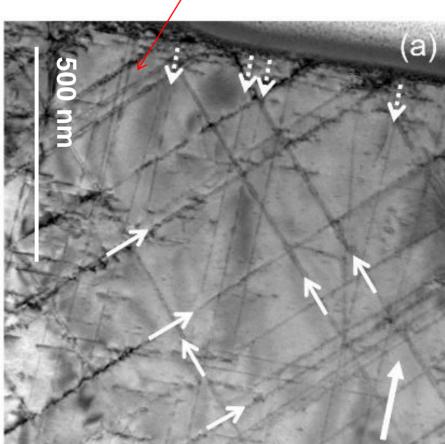
Microstructures of GdBCO coated conductors irradiated with 80MeV Xe ions at different angles

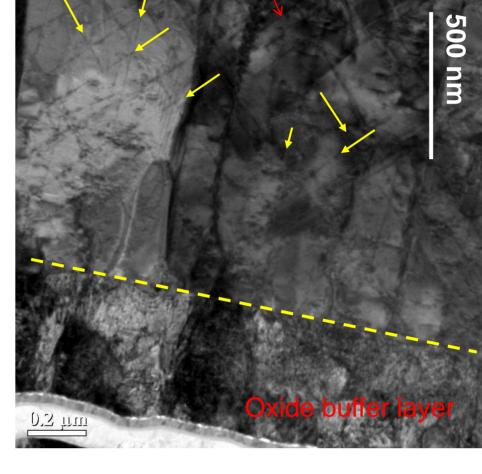
$\theta_i = 0^\circ$ and $\pm 45^\circ$ relative to the *c*-axis

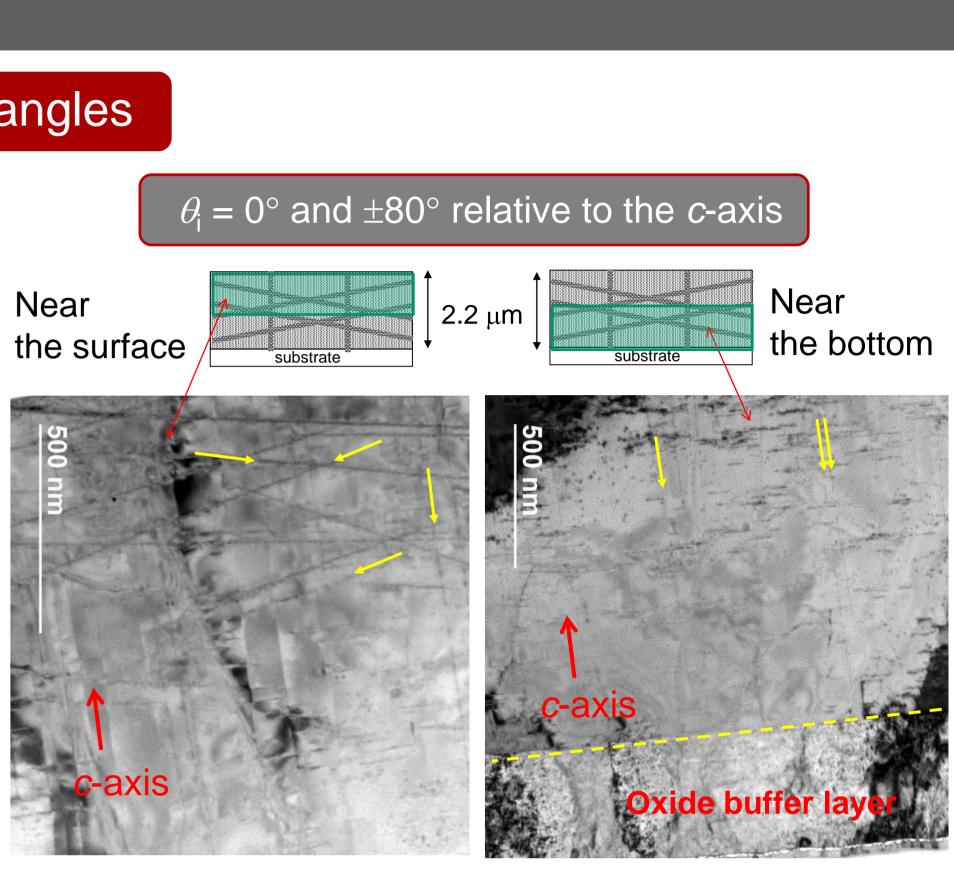
Near the bottom the surface

2.2 μm

the bottom







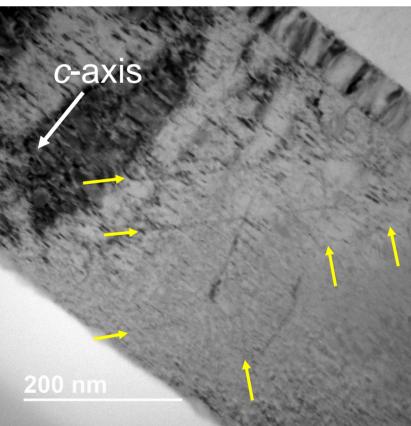
	$\theta_{\rm i}$ = ±20°	$\theta_{\rm i} = \pm 45^{\circ}$	$\theta_{\rm i}$ = ±80°						
	20.3	20.3	20.3						
	19.0 (532)	18.5 (707)	11.8 (2879)						
	17.5 (1064)	16.5 (1414)	— (5759)						
	<mark>16.0</mark> (1596)	<mark>14.3</mark> (2121)	— (8638)						
	<mark>14.3</mark> (2128)	<mark>12.0</mark> (2828)	—(11518)						
e the ion lepth [nm].									
ivity D in high- T_c superconductors $D_c < D_{ab}$ ne c-axis A radial temperature spread from the ion path is more suppressed in the plane containing the <i>c</i> -axis									
More severe irradiation damage									

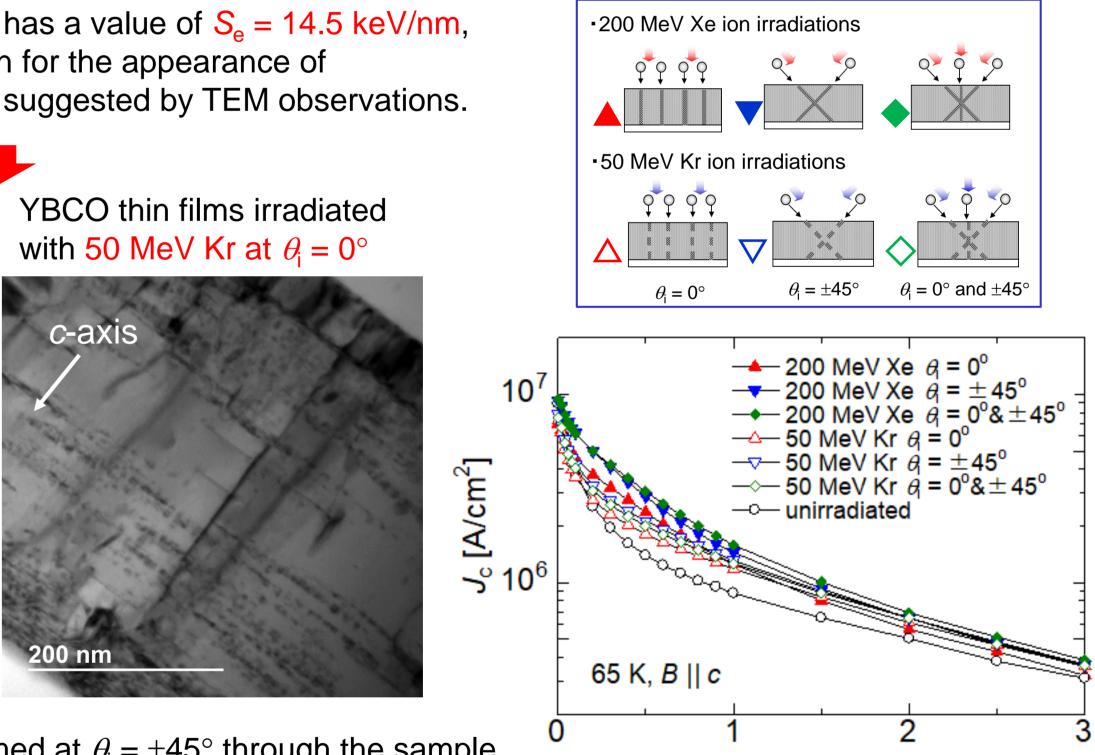
Anisotropy of thermal diffusivity Directional dependence of CD morphology

Introduction of discontinuous CDs at $\theta i = \pm 45^{\circ}$ by using 50 MeV Kr ions

50 MeV Kr irradiation to YBCO has a value of $S_e = 14.5$ keV/nm, which is the irradiation condition for the appearance of discontinuous CDs at $\theta_i = \pm 45^\circ$ suggested by TEM observations.

YBCO thin films irradiated with 50 MeV Kr at $\theta_i = \pm 45^\circ$





 \checkmark Discontinuous CDs are formed at $\theta_i = \pm 45^\circ$ through the sample. \checkmark No visible irradiation defects exists for the irradiation at $\theta_i = 0^\circ$.

4. Conclusions

Microstructures of GdBCO irradiated by 80 MeV Xe ions at various angles θ_i relative to the c-axis

• Not Discontinuous CDs but continuous ones are easily formed for the irradiation more tilted from the *c*-axis. • CDs are turned from continuous shape to discontinuous one at a deeper depth for the irradiation more tilted from the *c*-axis. • The TEM observations enabled us to identify the irradiation condition for the formation of discontinuous CDs at $\theta_i = 45^\circ$.

Effect of direction-dispersion of discontinuous CDs on J_c properties in YBCO thin films

• The flux pinning by discontinuous CDs is improved by the direction-dispersion of the ion tracks. • The irradiation defects by 50 MeV Kr were too small to provide superior J_c properties to continuous CDs by 200 MeV Xe.

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- \checkmark Direction-dispersion of the ion tracks is effective even for the flux pinning by discontinuous CDs.

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