



DISCOVERING PARTONIC RESCATTERING IN LIGHT NUCLEUS COLLISIONS

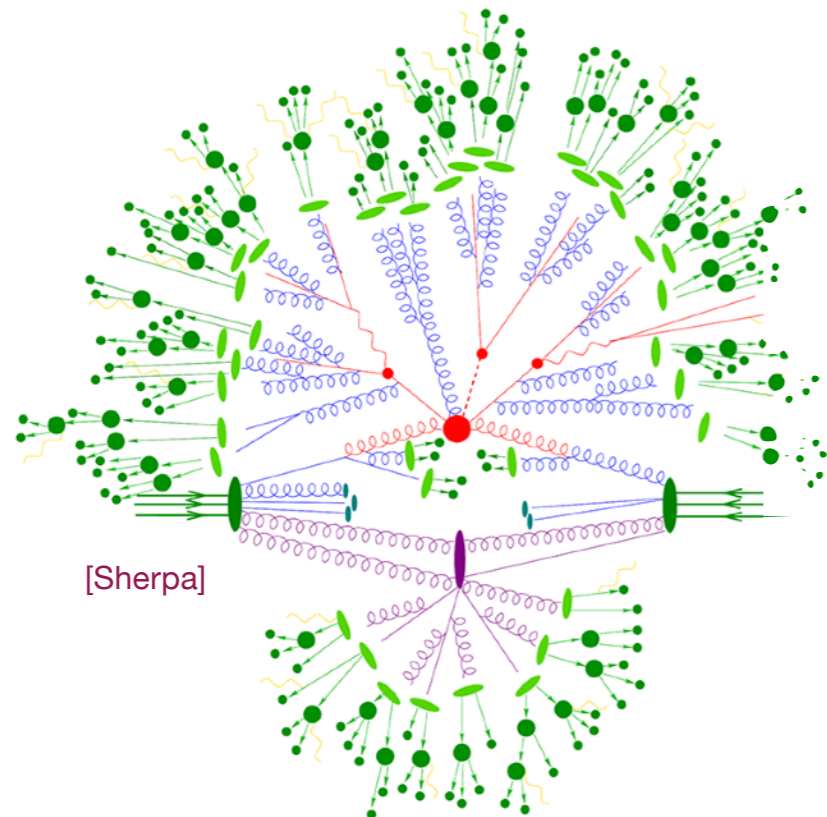
Alexander Huss

◉ based on: [\[2007.13754, 2007.13758\]](#)

AH, A. Kurkela, A. Mazeliauskas, R. Paatelainen, W. van der Schee, U. Wiedemann

DESCRIBING HADRON COLLISIONS — THE TWO PARADIGMS

HEP

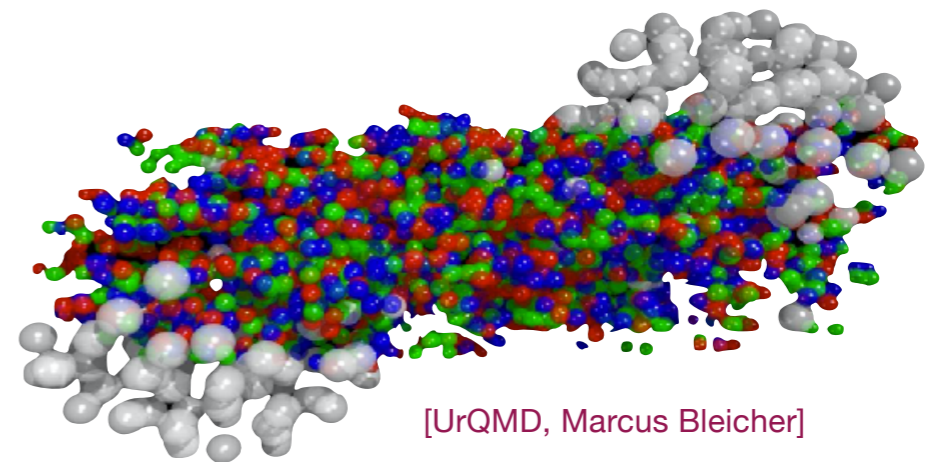


- proton-proton

↔ no re-scattering

“free streaming”

HIP



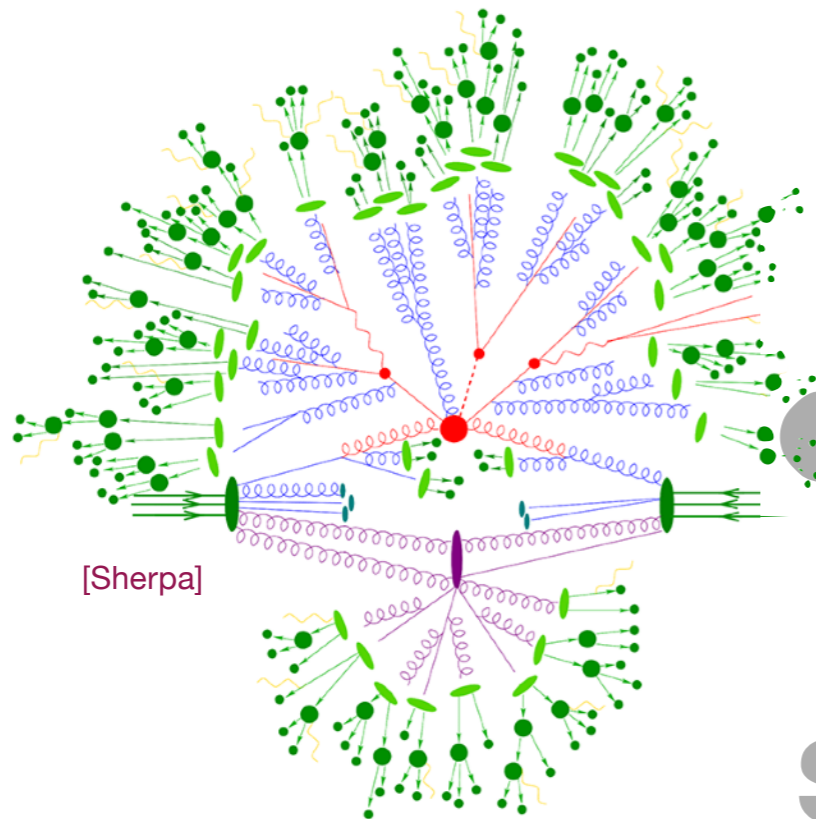
- nucleus-nucleus (AA)

↔ many re-scatterings

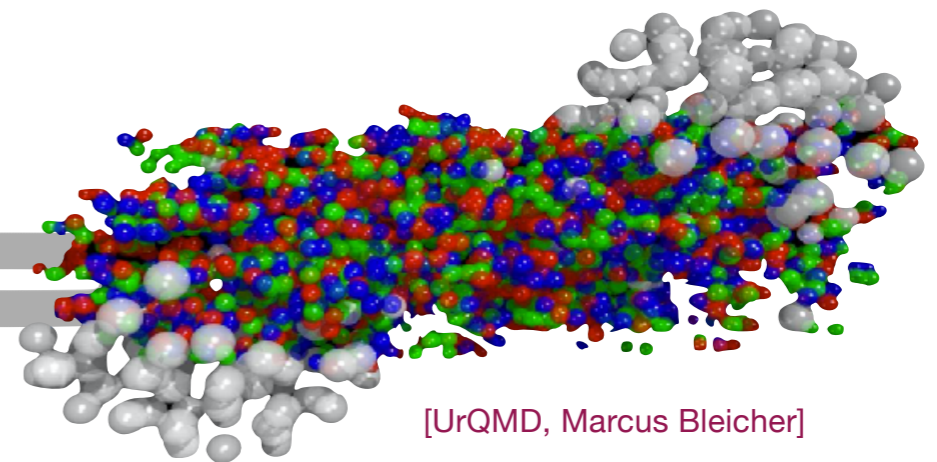
nearly perfect fluid

DESCRIBING HADRON COLLISIONS — THE TWO PARADIGMS

HEP



HIP



?

small systems

- proton-proton

↔ no re-scattering

“free streaming”

- nucleus-nucleus (AA)

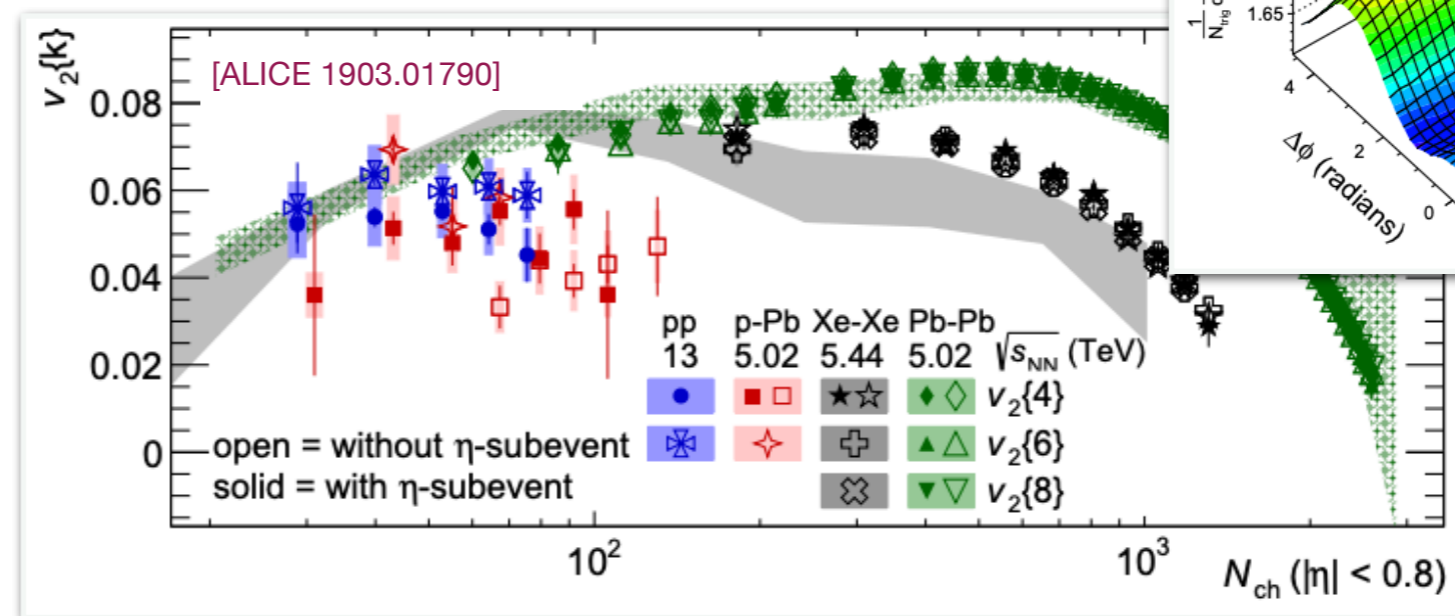
↔ many re-scatterings

nearly perfect fluid

SMALL SYSTEMS — CHALLENGING THE TWO PARADIGMS

(talk by Jacquelyn Noronha-Hostler)

- ✓ observation of collectivity in small systems:
(collective flow v_n 's)



✗ no sign of parton energy loss!

⚡ collectivity ↔ parton rescattering

- ⊙ some parton energy loss must also be present in small systems!

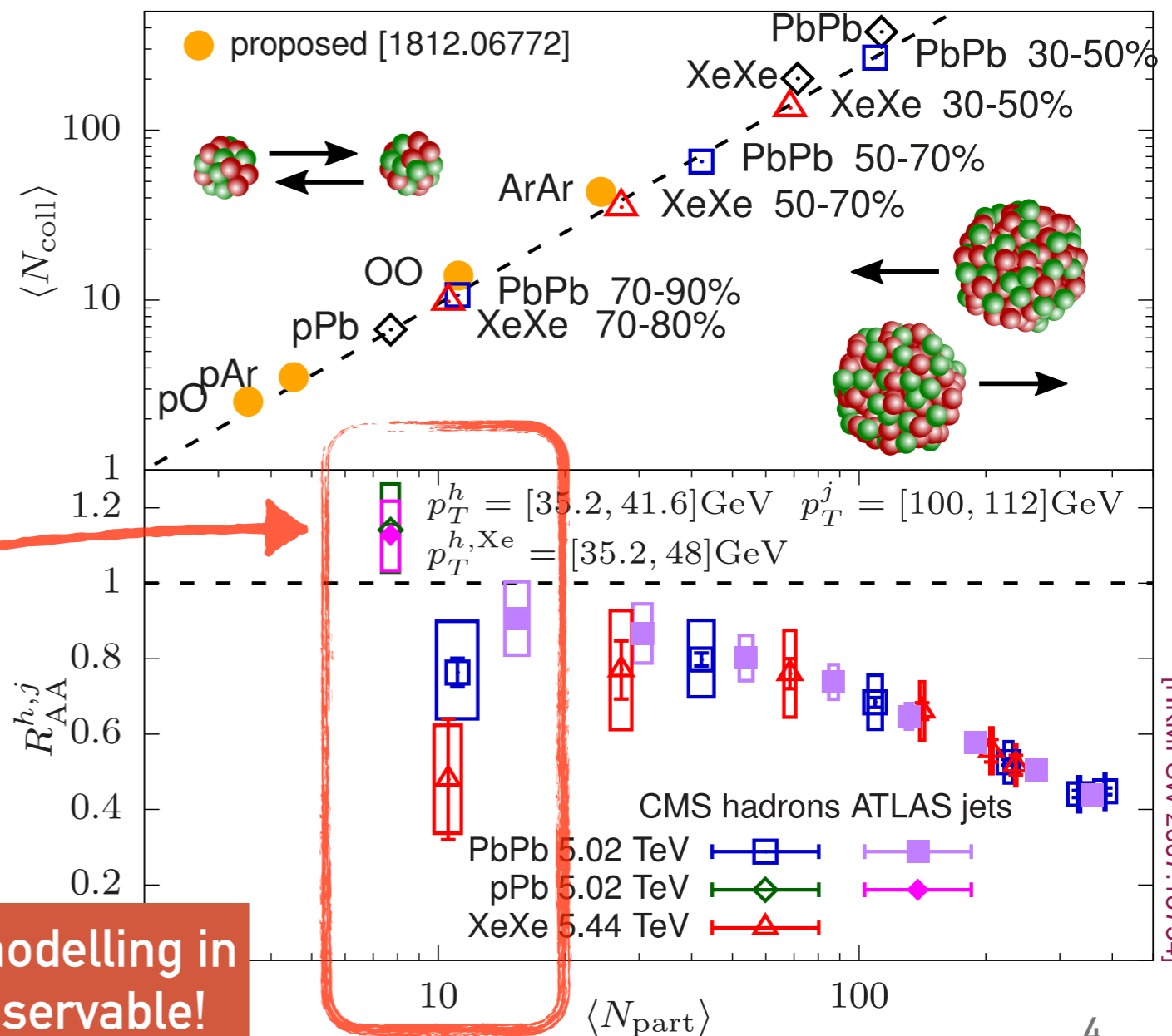
THE NUCLEAR MODIFICATION FACTOR $R_{AA}^{j,h}$

compare differential yield in AA collisions to an equivalent number of pp collisions

$$R_{AA}^{j,h}(p_T) = \frac{(1/N_{ev})}{\langle T_{AA} \rangle} \frac{dN_{AA}^{j,h}/dp_T}{d\sigma_{pp}^{j,h}/dp_T}$$

- system size
↔ centrality selection
- $\langle T_{AA} \rangle$: nuclear overlap function (from model calculations)
- $\langle N_{part} \rangle \sim 10$
systematics (boxes)
dominated by $\langle T_{AA} \rangle$

need to get rid of soft modelling in our nominally hard observable!

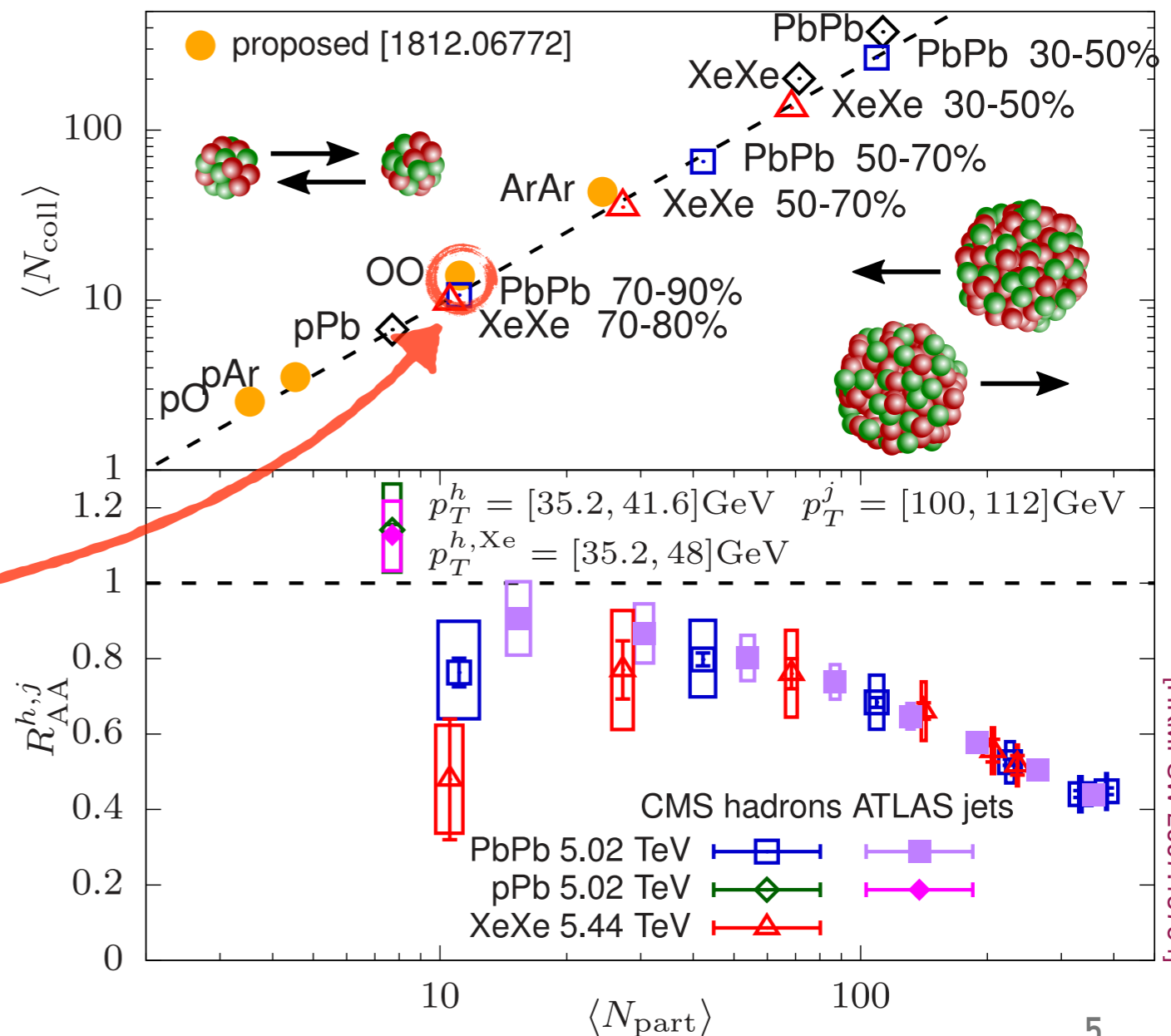


THE NUCLEAR MODIFICATION FACTOR $R_{AA}^{j,h}$

only in minimum-bias collisions

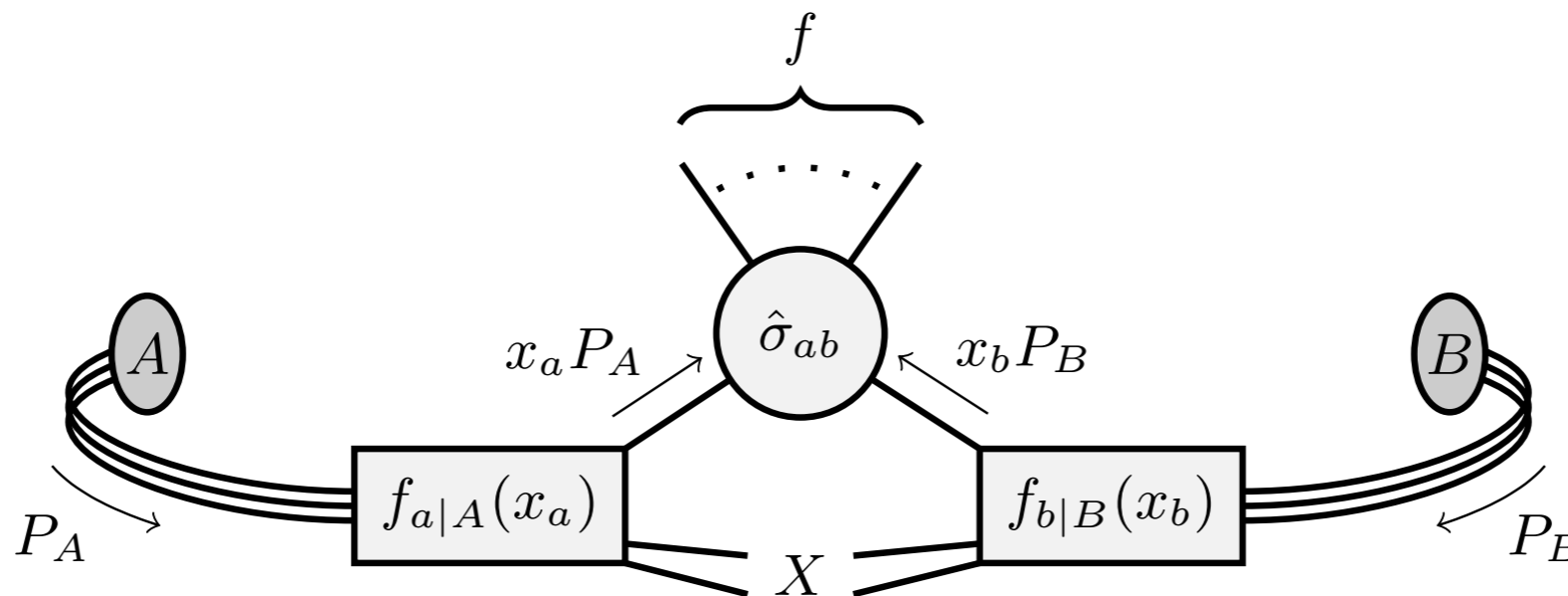
$$R_{AA, \text{min bias}}^{j,h}(p_T) = \frac{1}{A^2} \frac{d\sigma_{AA}^{j,h}/dp_T}{d\sigma_{pp}^{j,h}/dp_T}$$

- system size
↔ A : nucleon number
- $\langle T_{AA} \rangle$ replaced by:
beam luminosity
- $\langle N_{\text{part}} \rangle \sim 10$
collide light nuclei: OO



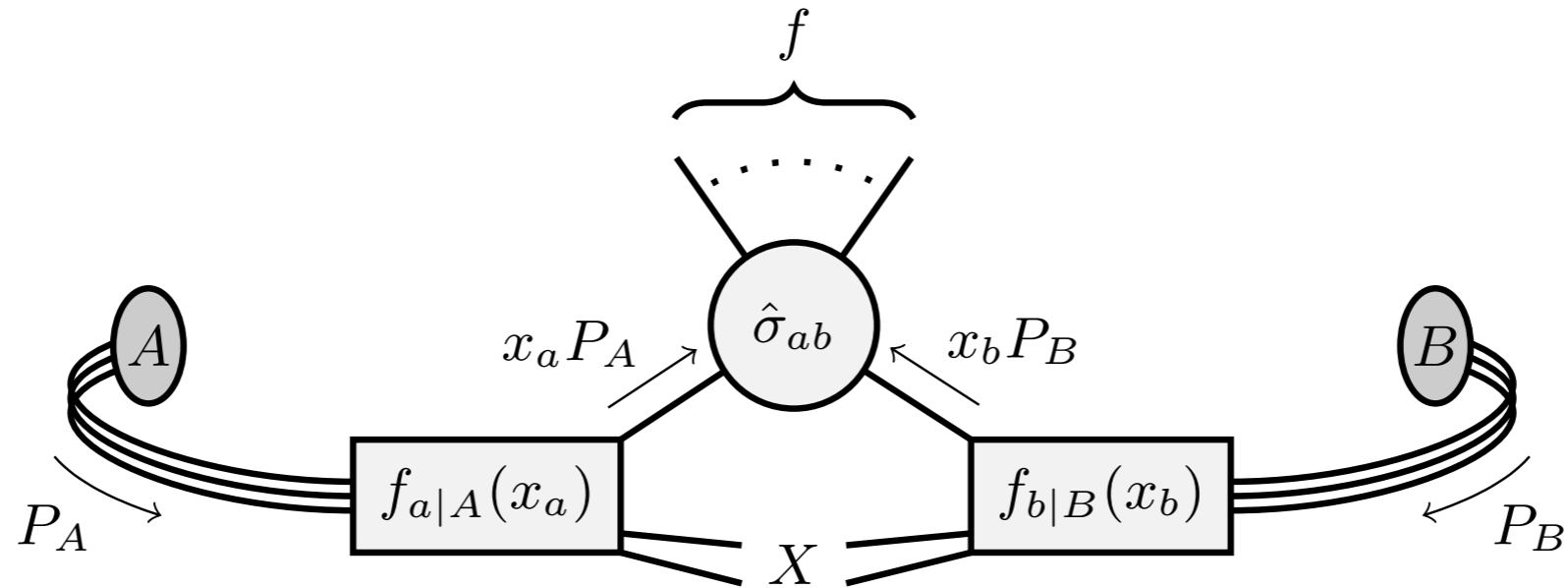
THE NULL HYPOTHESIS

- discovering a small energy loss signal requires a **precise theoretical baseline** (null = “no energy loss”)
- can compute $R_{AA,\text{min bias}}^{j,h}$ (null) using QCD factorization ($Q \gg \Lambda_{\text{QCD}}$):



$$\sigma_{AB} = \sum_{ab} \int_0^1 dx_a \int_0^1 dx_b f_{a|A}(x_a) f_{b|B}(x_b) \hat{\sigma}_{ab}(x_a, x_b) (1 + \mathcal{O}(\Lambda_{\text{QCD}}/Q))$$

QCD FACTORIZATION

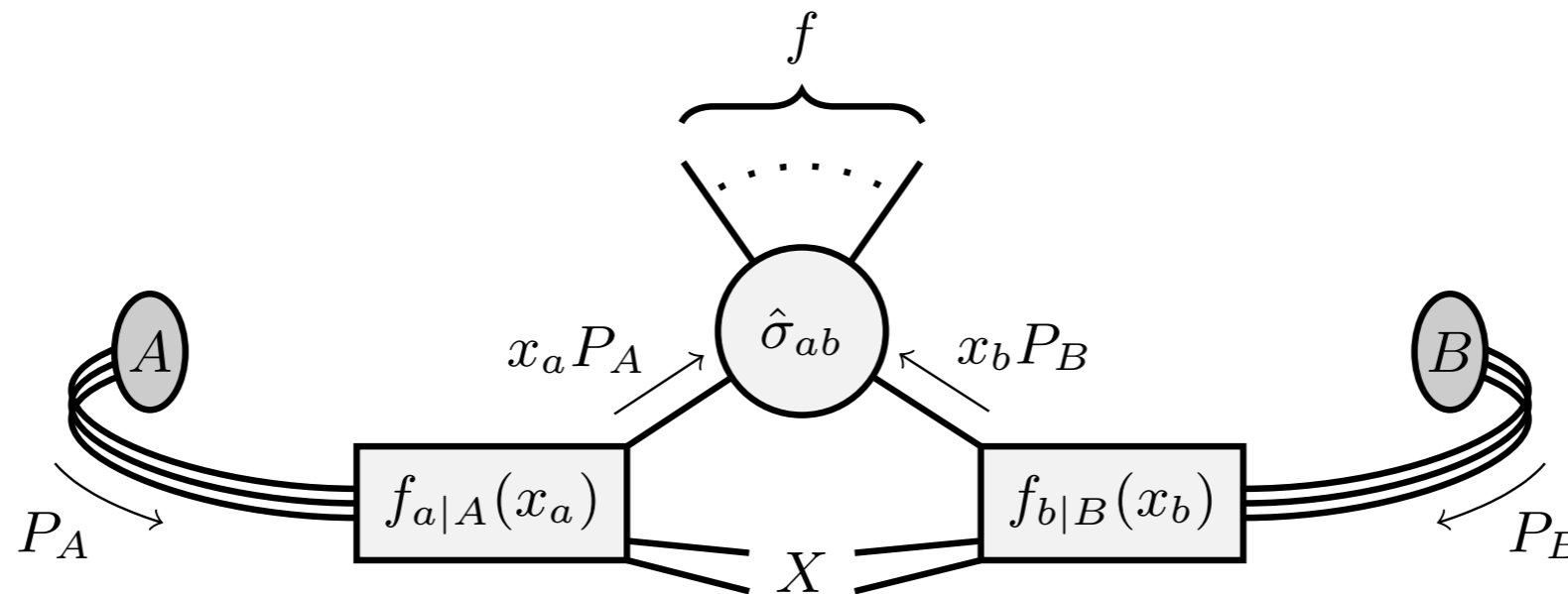


$$\sigma_{AB} = \sum_{ab} \int_0^1 dx_a \int_0^1 dx_b f_{a|A}(x_a) f_{b|B}(x_b) \hat{\sigma}_{ab}(x_a, x_b) (1 + \mathcal{O}(\Lambda_{\text{QCD}}/Q))$$

(nuclear) parton distribution functions
(non-perturbative, universal)

hard scattering
(perturbation theory)

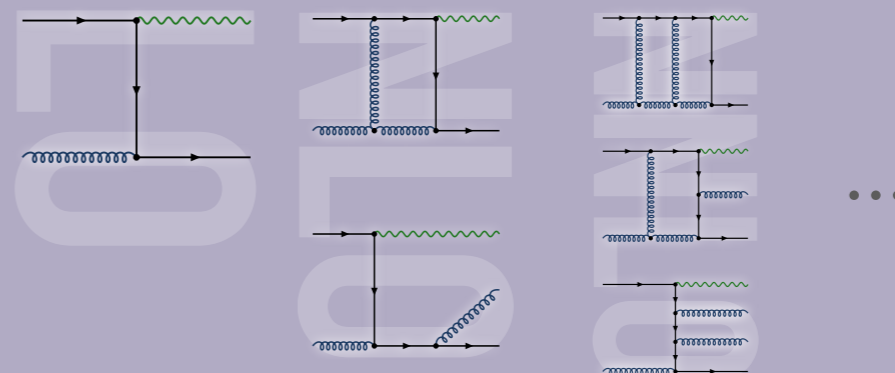
QCD FACTORISATION



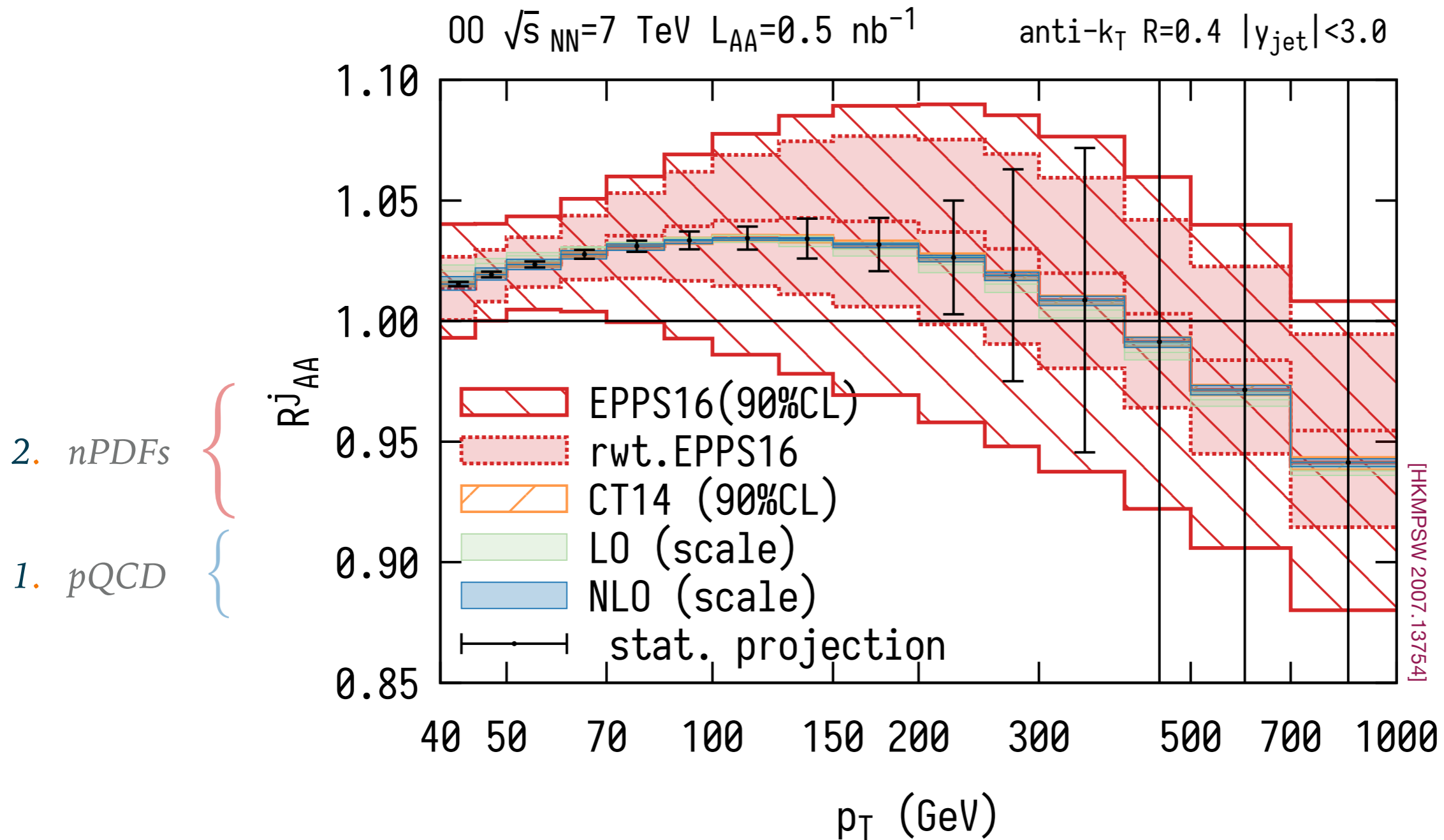
$$\sigma_{AB} = \sum_{ab} \int_0^1 dx_a \int_0^1 dx_b f_{a|A}(x_a) f_{b|B}(x_b) \hat{\sigma}_{ab}(x_a, x_b) (1 + \mathcal{O}(\Lambda_{\text{QCD}}/Q))$$

(n)PDFs
(in principle, improvable)
 (talk by Petja Paakkinen)

hard scattering
(systematically improvable)



MINIMUM-BIAS R_{AA}^j (null)

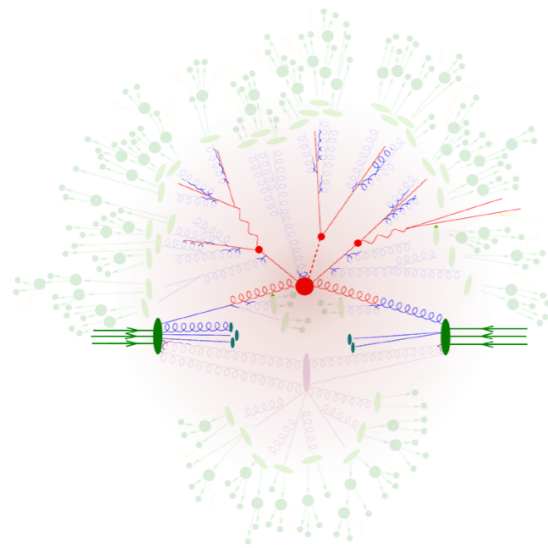


MINIMUM-BIAS R_{AA}^j (null) — SCALE UNCERTAINTIES

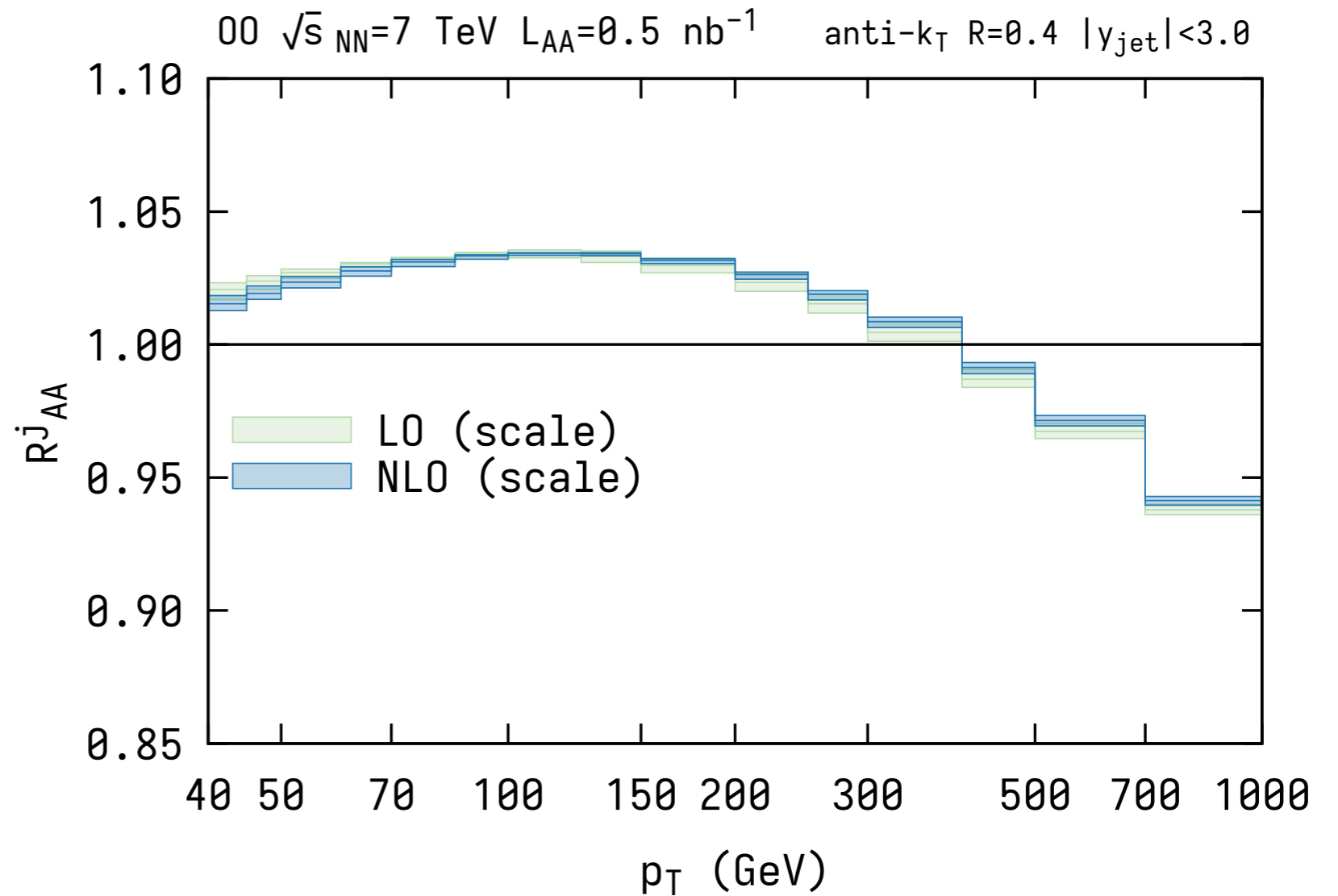
- truncation of series:
 - ↔ scale-dependence: μ_R & μ_F (unphysical) } $\sigma^{\text{N}^k\text{LO}}(\mu) - \sigma^{\text{N}^k\text{LO}}(\mu') = \mathcal{O}(\text{N}^{k+1}\text{LO})$
- estimate size of missing higher-order terms:
 - ↔ vary μ_R & μ_F by factors $(1/2, 2)$ ($1/2 \leq \mu_R/\mu_F \leq 2$)
- estimate for ratios:
 - ↔ same process \Rightarrow assume correlated μ_R & μ_F
 - ↔ test prescription by inspecting perturbative series

$$R^{A/B}(\mu_R, \mu_F) = \frac{\sigma^A(\mu_R, \mu_F)}{\sigma^B(\mu_R, \mu_F)} \quad \text{“7-point variation”}$$

MINIMUM-BIAS R_{AA}^j (null) — SCALE UNCERTAINTIES

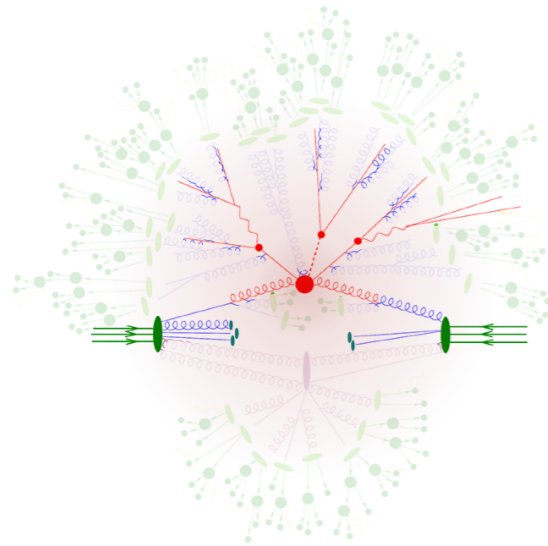


*fixed-order
hard scattering*

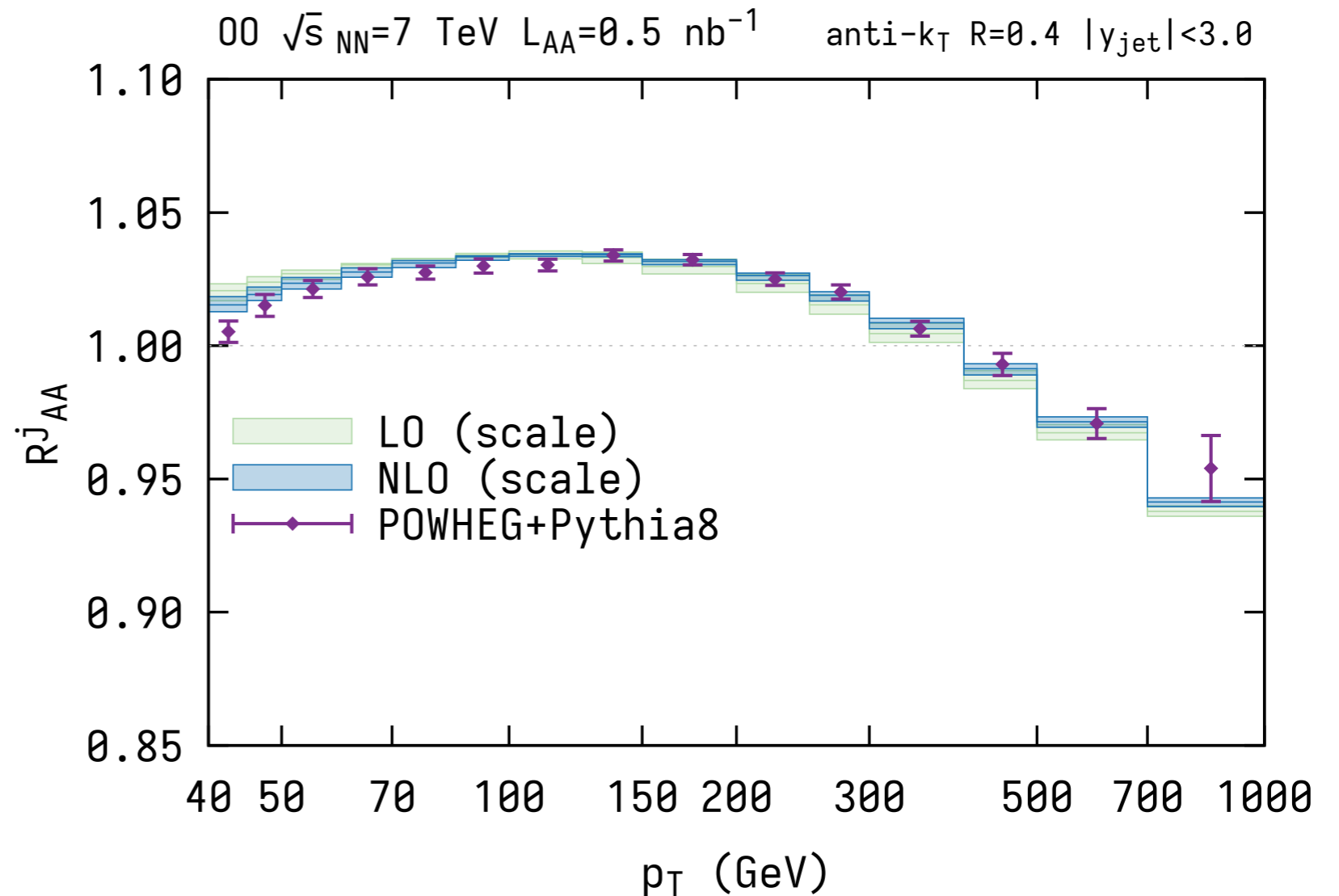
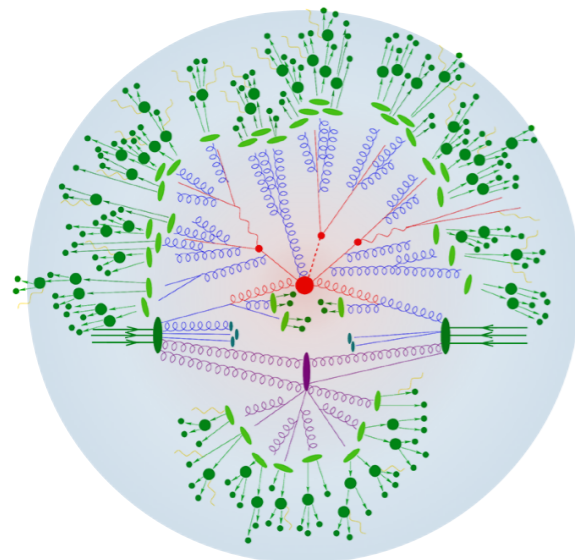


- observable under **good perturbative control**.

MINIMUM-BIAS R_{AA}^j (null) — PARTON SHOWER & HADRONIZATION



parton shower
 \oplus
 hadronization



- observable under **good perturbative control**.
- **parton shower** and **hadronization effects** largely cancel in the ratio

MINIMUM-BIAS R_{AA}^j (null) — (N)PDF UNCERTAINTIES

● EPPS16 for ^{16}O

→ proton baseline: CT14

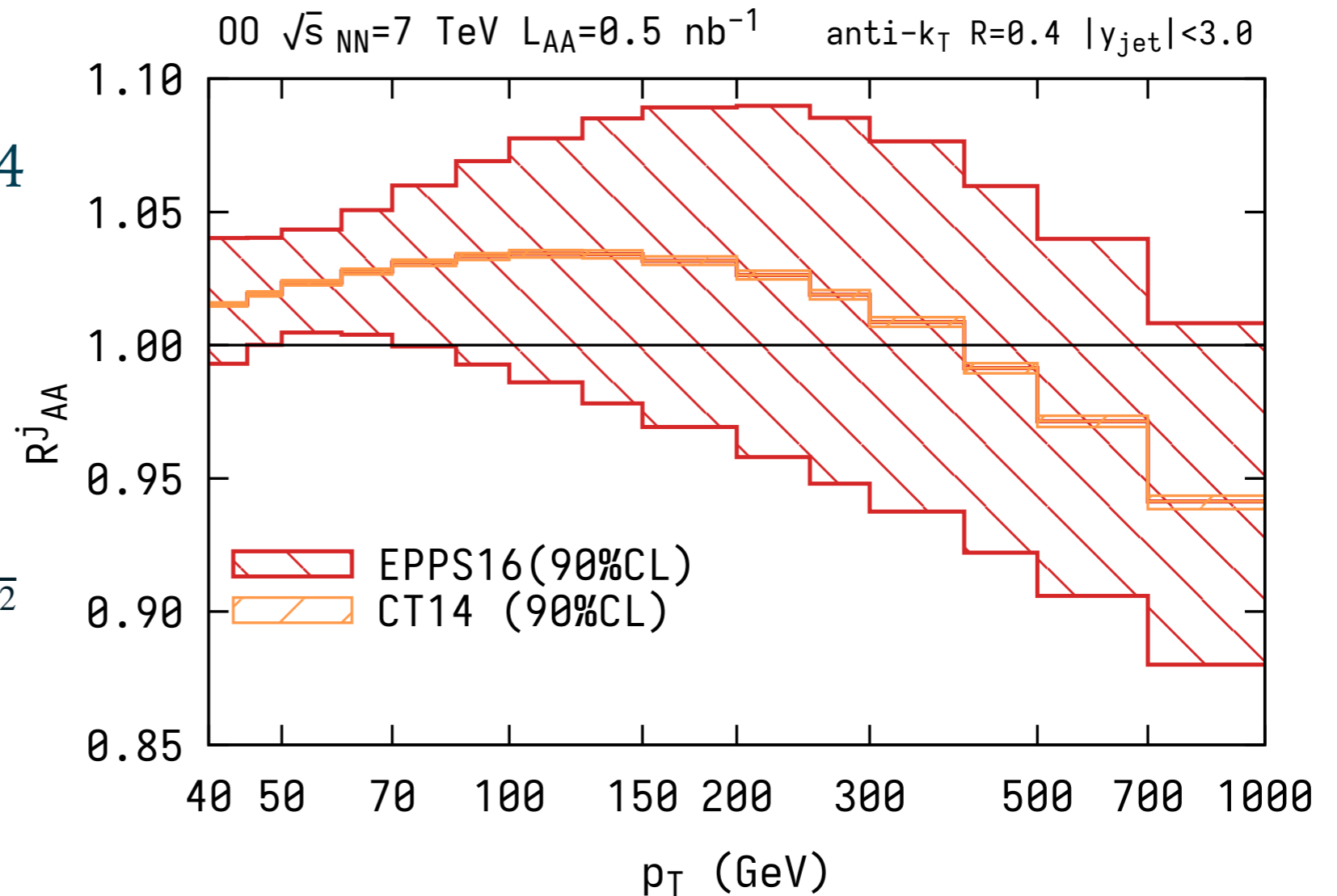
(56 members: $S_{i=21,\dots,48}^{\pm}$)

→ nuclear modification

(40 members: $S_{i=1,\dots,20}^{\pm}$)

→ errors: (90% CL)

$$\Delta\mathcal{O} = \frac{1}{2}\sqrt{\sum_i [\mathcal{O}(S_i^+) - \mathcal{O}(S_i^-)]^2}$$



● proton variation cancel

● nPDF uncertainties remain

can be improved by including more data!

$$R_{AA} = \frac{d\sigma_{AA}}{d\sigma_{pp}} \leftarrow \begin{array}{l} \text{EPPS16 (+CT14)} \\ \text{CT14} \end{array}$$

MINIMUM-BIAS R_{AA}^j (null) — (N)PDF UNCERTAINTIES

● EPPS16 for ^{16}O

→ proton baseline: CT14

(56 members: $S_{i=21,\dots,48}^{\pm}$)

→ nuclear modification

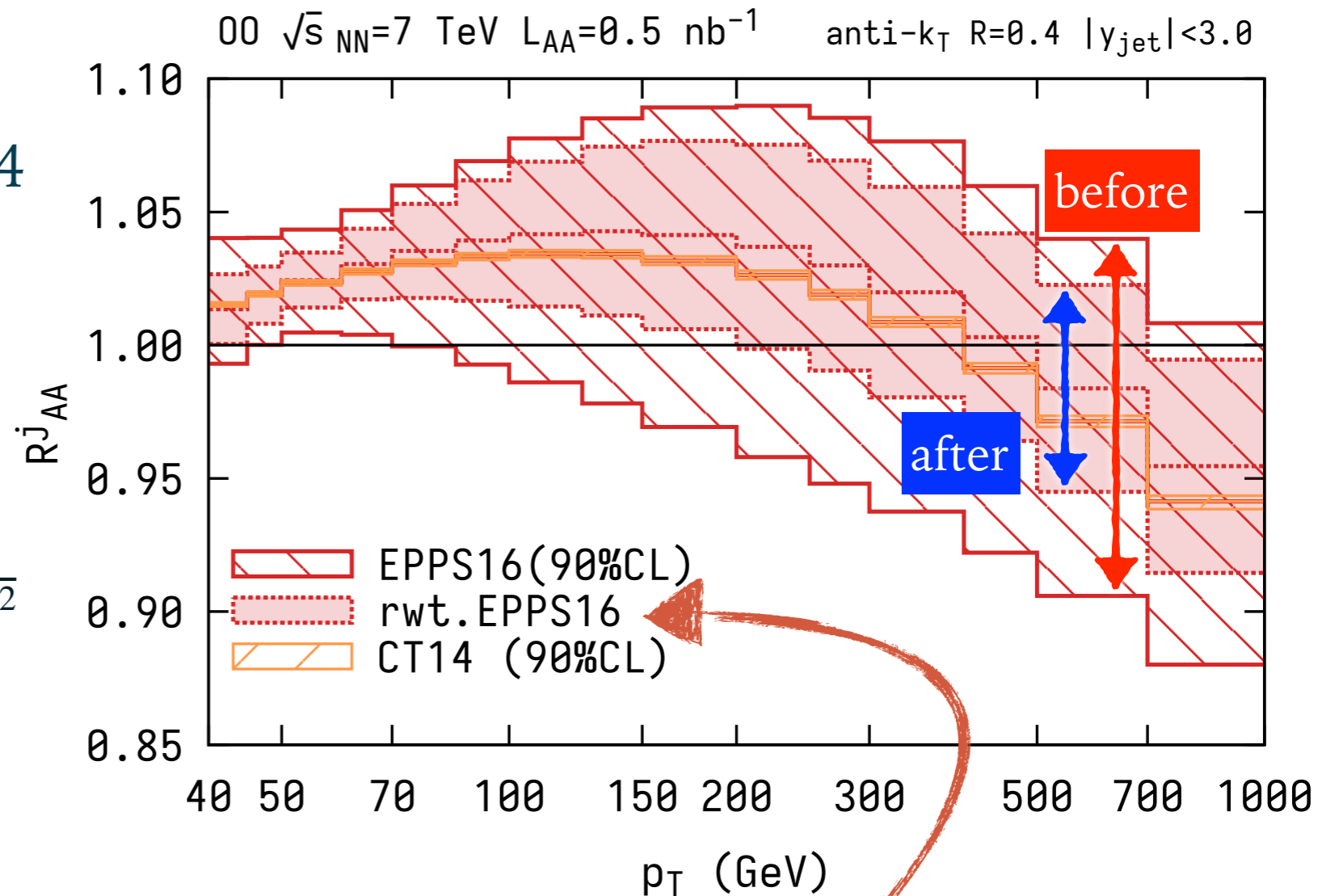
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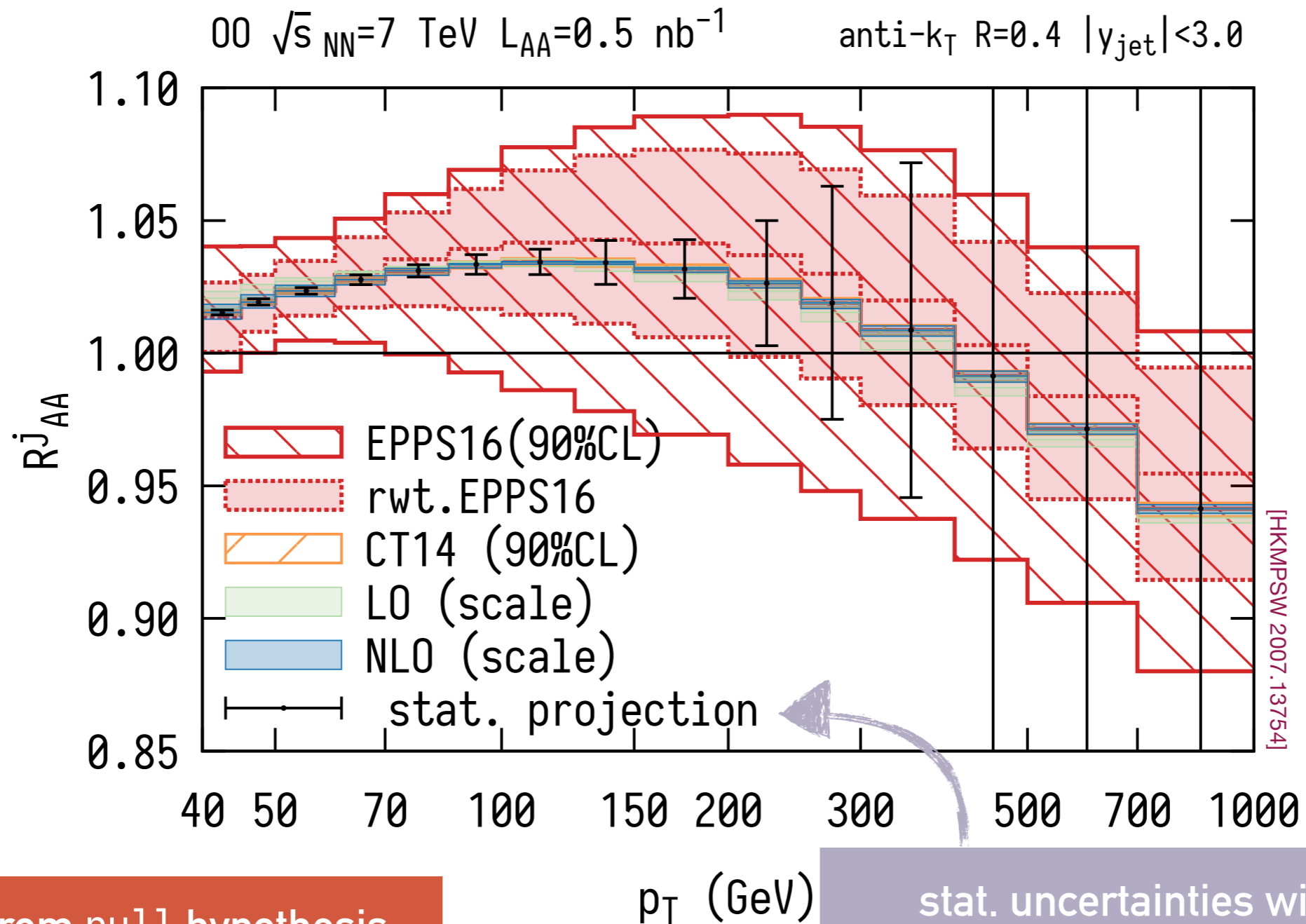
● proton variation cancel

● nPDF uncertainties remain



reweight nPDFs using
CMS pPb dijet data

MINIMUM-BIAS R_{AA}^j (null)



deviation from null hypothesis
 \Rightarrow discovery of energy loss!

stat. uncertainties with 0.5 nb $^{-1}$
 few hours “moderately optimistic” scenario

MINIMUM-BIAS R_{AA}^h (null) v.s. ENERGY-LOSS MODELS

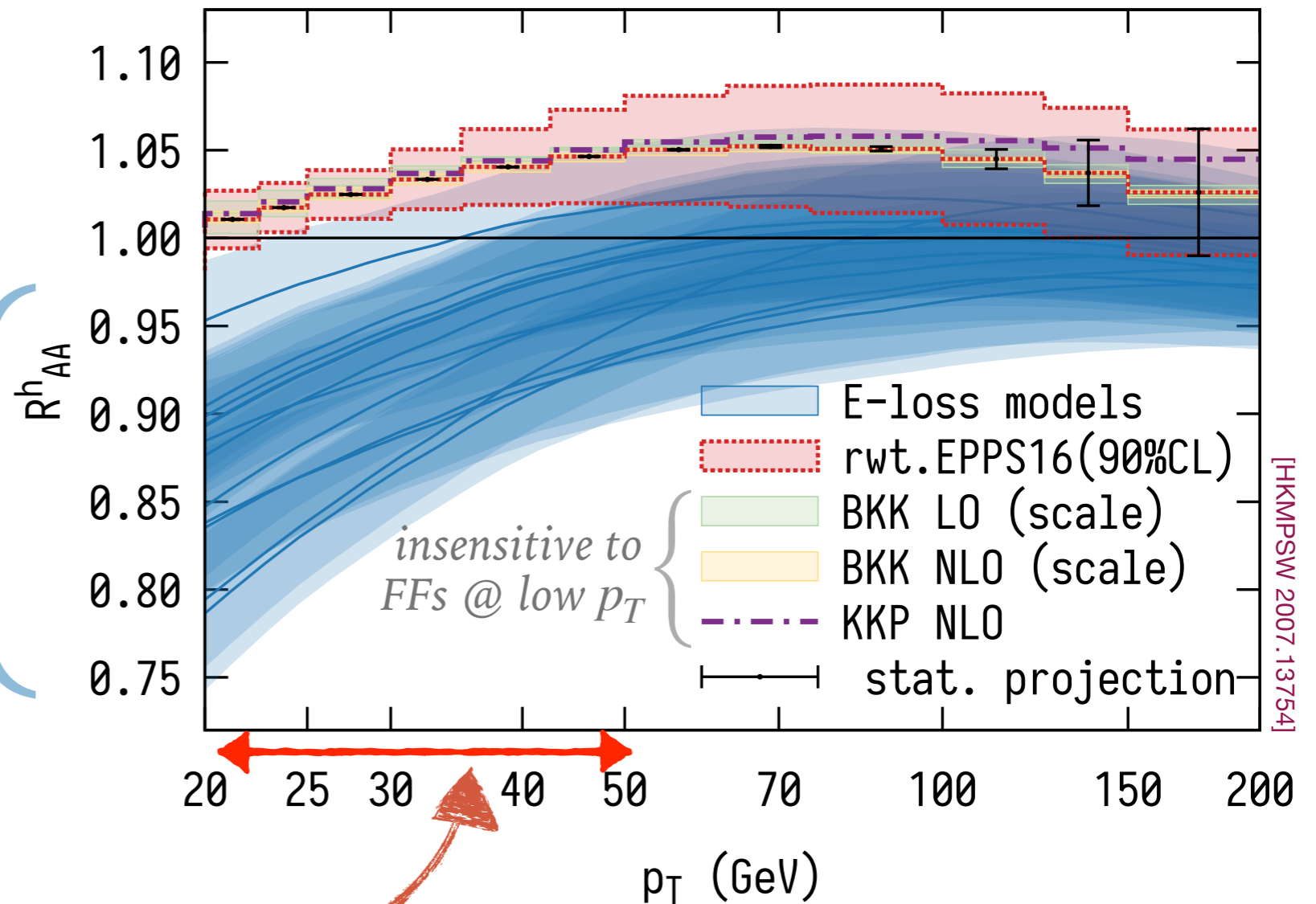
$$\sigma^h \sim (n)\text{PDFs} \otimes \hat{\sigma}_{q,g} \otimes \text{FFs}$$

fragmentation functions

- energy-loss models
- extrapolate to ∞
- estimate for **signal**

[HKMPSW 2007.13758]

$\infty \sqrt{s}_{NN}=7 \text{ TeV } L_{AA}=0.5 \text{ nb}^{-1} \quad |y_h| < 1.0$



potentially measurable signal in
 $20 \text{ GeV} \lesssim p_T \lesssim 50 \text{ GeV}$

* no signal \Rightarrow rule out models / scenarios

HADRON ENERGY-LOSS IN LIGHT-ION COLLISIONS

[HKMPSW 2007.13758]

- simple & modular framework to explore different model setups

↪ background $T(\tau, r)$ profile:

from hydro-like to free streaming

conformal EoS v.s. lattice EoS

dynamical geometry v.s. Bjorken $\tau^{-1/3}$

isotropic geometry v.s. azimuthal anisotropy ϵ_2

↪ energy-loss models:

BDMPS-Z [Arnold 0808.2767]

v.s. AdS/CFT inspired

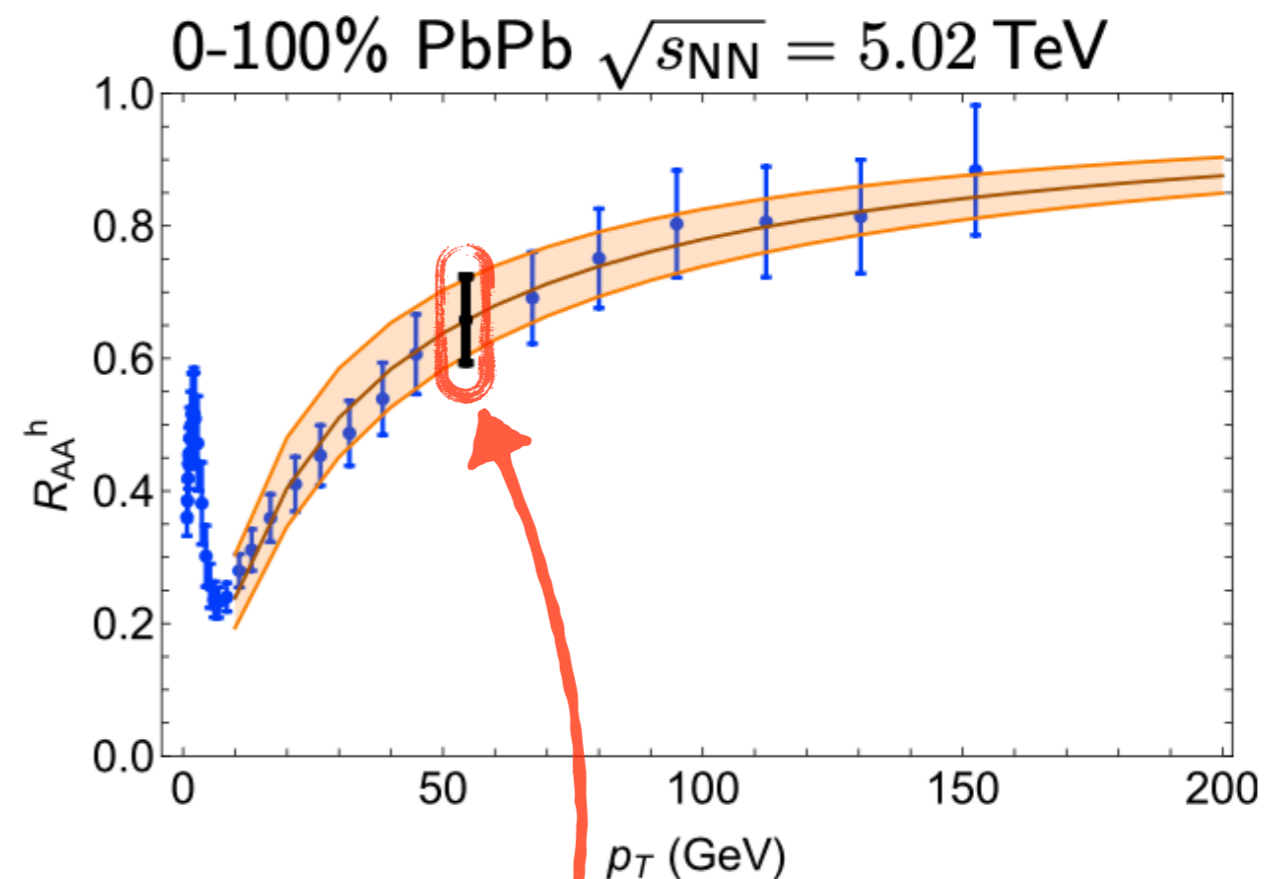
v.s. $dE/dL \sim L^{0.4}T^{1.2}$

v.s. $dE/dL \sim LT^3$

varying starting time of energy loss $\tau_0 = 0.05 - 0.5$ fm/c

with or without nPDFs and FFs

↪ model parameter (\hat{q}) fitted to single data point @ $R_{\text{PbPb}}^h(p_T = 54.4 \text{ GeV})$



- dependence on p_T & system size are predictions of the model

HADRON ENERGY-LOSS IN LIGHT-ION COLLISIONS

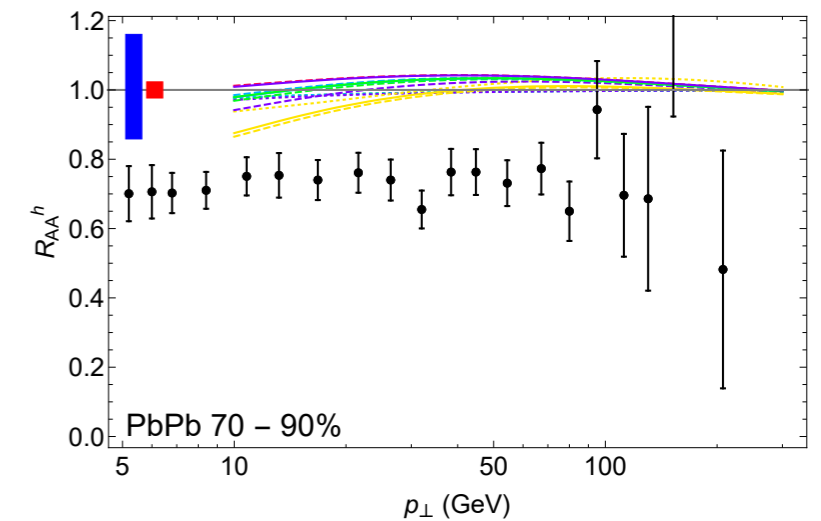
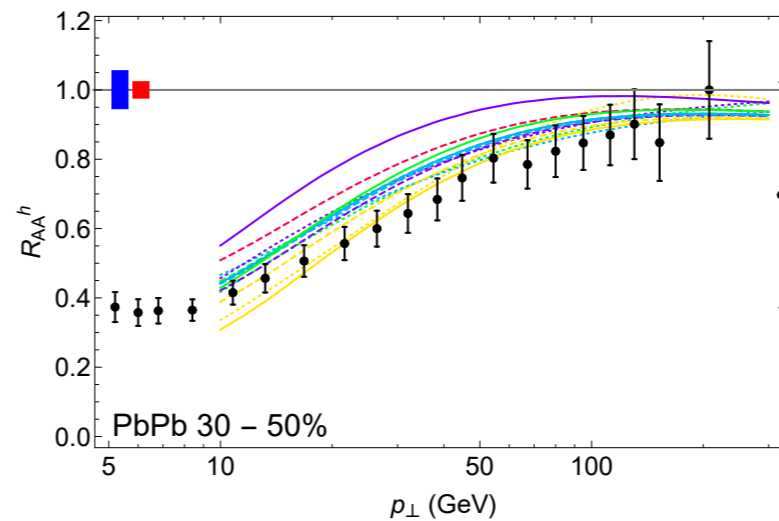
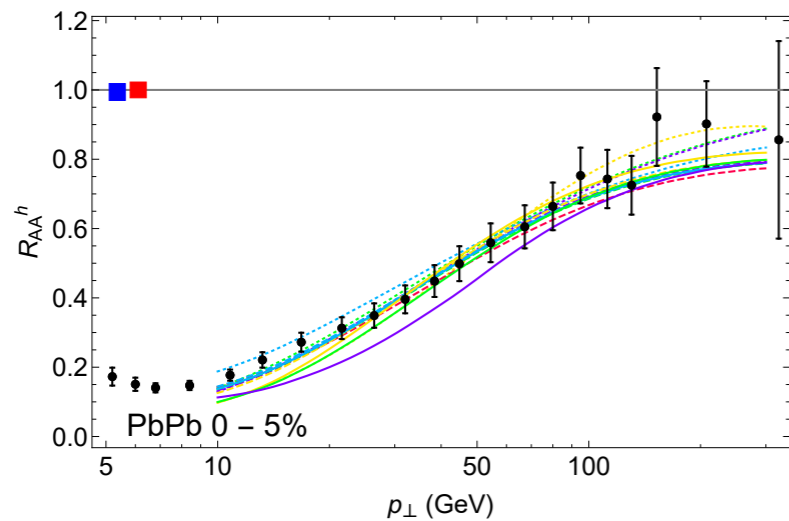
[HKMPSW 2007.13758]

0-5%

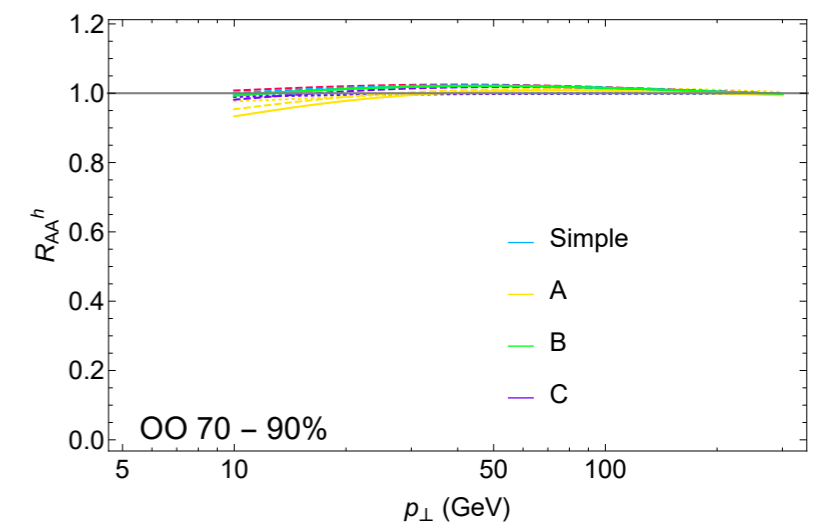
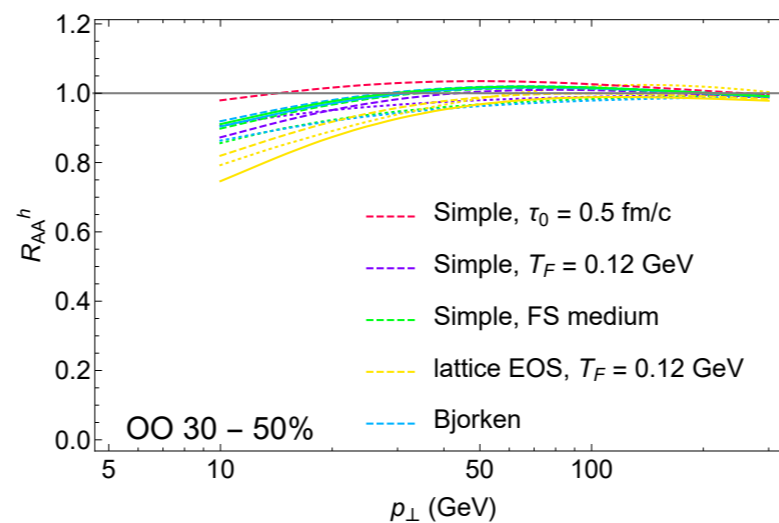
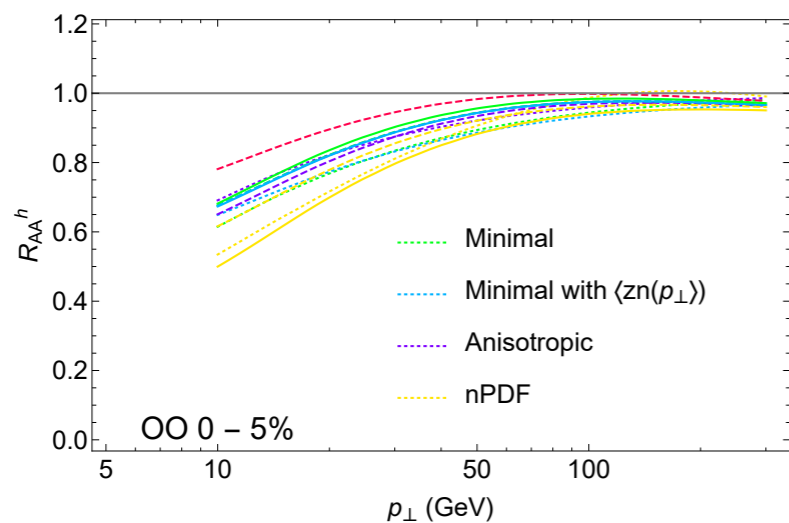
30-50%

70-90%

PbPb



OO



- good mid-central description
- theory uncertainty (conservative) — spread of different scenarios

Z - BOSONS AS A STANDARD CANDLE

- ◉ $\langle T_{AA} \rangle$ replaced by **beam luminosity** in $R_{AA, \text{min bias}}^{j,h}$
 - ↪ requires van der Meer scan (**session on Thursday**)
- ◉ use Drell-Yan as luminosity meter
 - ↪ cancel luminosity uncertainties in **cross-ratio**:

$$R_{AA,Z}^{j,h}(p_T) = \frac{\sigma_{pp}^Z \frac{d\sigma_{AA}^{j,h}}{dp_T}}{\sigma_{AA}^Z \frac{d\sigma_{pp}^{j,h}}{dp_T}}$$

? maybe even get some cancellation of nPDF uncertainties...

Z - BOSON NORMALISED $R_{AA,Z}^j$ (null)

...nope ;(

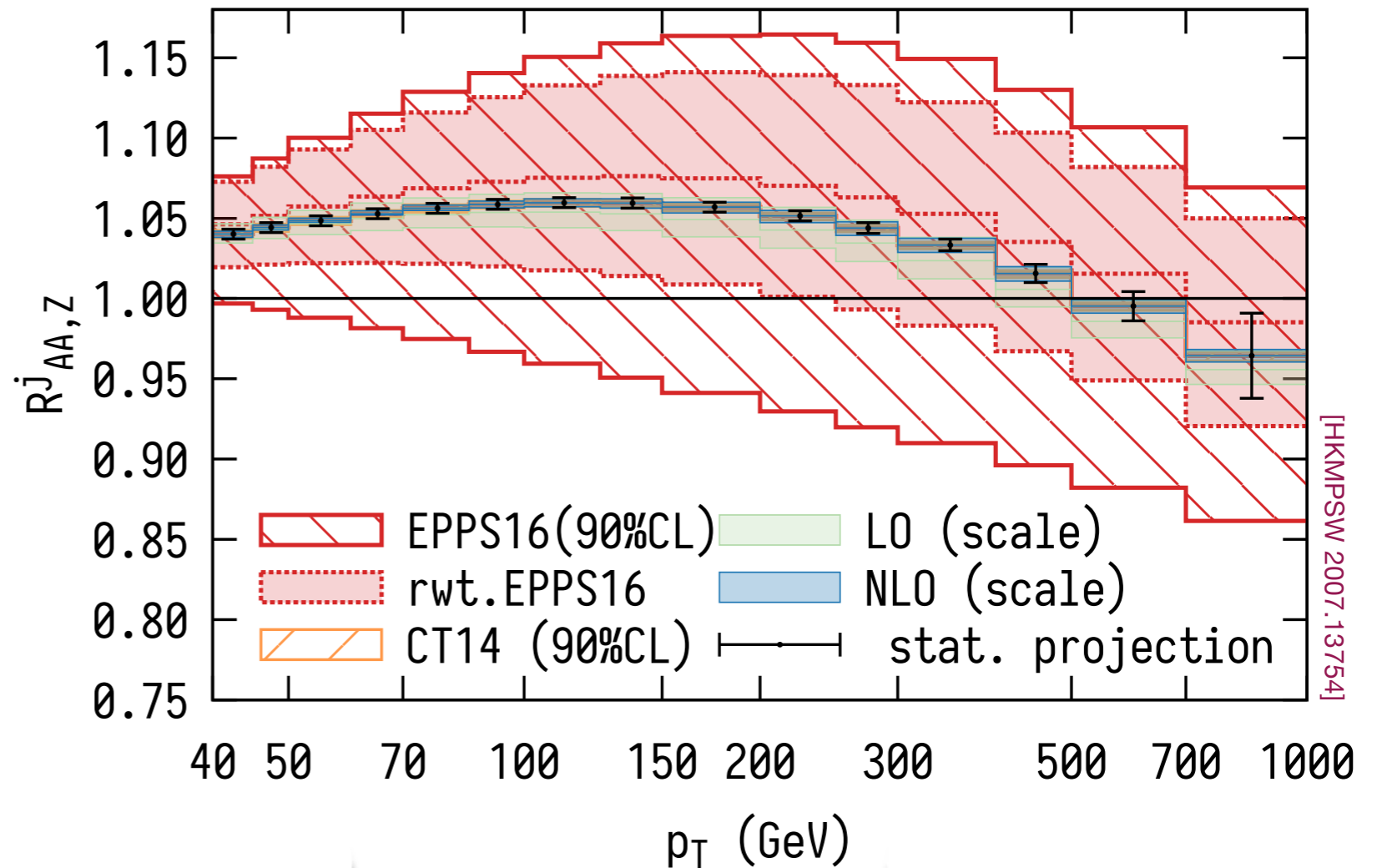
+ luminosity uncertainties cancel

- requires larger statistics $\mathcal{O}(1 \text{ pb}^{-1})$

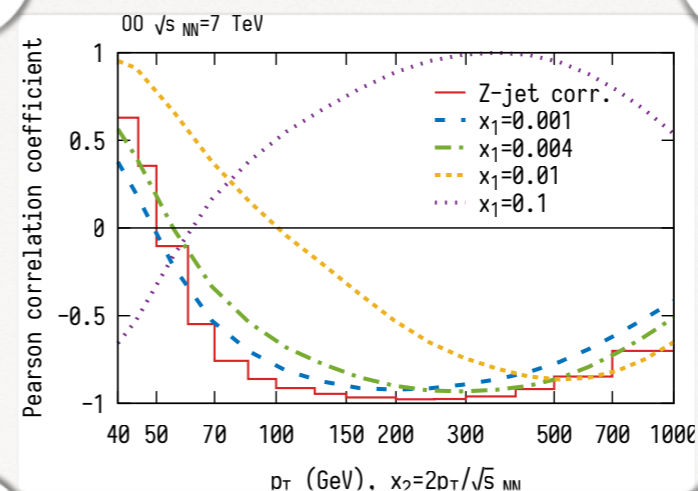
- nPDF uncertainties larger than for R_{AA}^j

↔ anti-correlation in Bjorken- x between Z - & jet-production:

$\sqrt{s}_{NN}=7 \text{ TeV}$ $L_{AA}=0.5 \text{ pb}^{-1}$ anti- k_T $R=0.4$ $|y_{jet}| < 3.0$



[HKMPSW 2007.13754]



TRIGGERED JET SPECTRUM

- trigger on Z-boson:
 - $\rightsquigarrow p_T^Z > 30 \text{ GeV}$ (\Rightarrow also inclusive case, effectively “Z+jets”)
 - $\rightsquigarrow m_{\ell\ell} \in [76, 106] \text{ GeV}, p_T^{\ell^\pm} > 20 \text{ GeV}, |y^{\ell^\pm}| < 2.5$
- measure inclusive jet p_T^j spectrum in triggered events
& normalize w.r.t. $N_Z \sim \sigma^Z(p_T^Z > 30 \text{ GeV})$

$$I_{AA}^{Z+\text{jets}}(p_T^j) \equiv \left. \frac{R_{AA}^{Z+\text{jets}}(p_T^j)}{R_{AA}^Z} \right|_{Z\text{-trig.}} = \frac{\sigma_{pp}^Z}{\sigma_{AA}^Z} \frac{d\sigma_{AA}^{Z+\text{jets}}/dp_T^j}{d\sigma_{pp}^{Z+\text{jets}}/dp_T^j} \Bigg|_{Z\text{-trig.}}$$

! surely, nPDF uncertainties cancel in this case...

Z - BOSON TRIGGERED R_{AA}^{Z+jets}

...still no



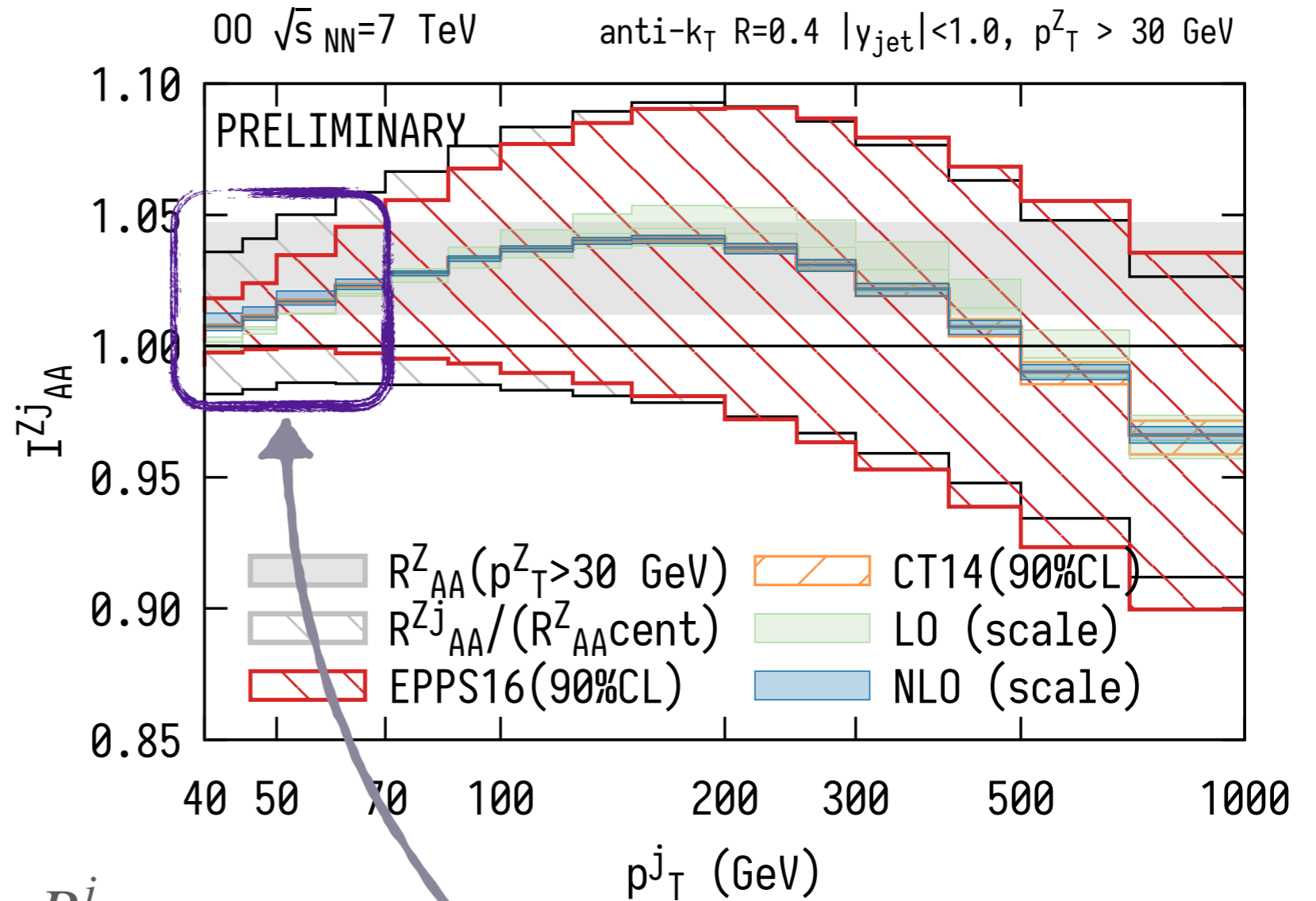
+ luminosity uncertainties cancel

- higher luminosity required

o nPDF uncertainties

$p_T > 100$ GeV: similar to R_{AA}^j

↔ denominator inclusive in jet(s)
smears out x -dependence



some cancellation $p_T^j \simeq p_T^Z$
maybe better observables?!

* hadron+jets, work in progress: smearing due to FFs

CONCLUSIONS & OUTLOOK

- ○○ — explore the intersection between HEP & HIP
 - inclusive measurement (minimum bias) for R_{AA}
 - ↔ eliminates soft-physics dependence
(source of large uncertainties)
 - ↔ accurate theoretical baseline possible
 - ↔ main residual uncertainties from nPDFs (improvable)
 - opportunity to establish:
collectivity ↔ parton rescattering ↔ energy loss
- possibility to eliminate/reduce luminosity/nPDF uncertainties
 - ❖ Drell-Yan as luminosity meter: nPDF anti-correlations
 - ❖ triggered observables: still sizeable residual nPDF uncertainties

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THANK YOU!

BACKUP.

BACKUP.

SCALE DEPENDENCE

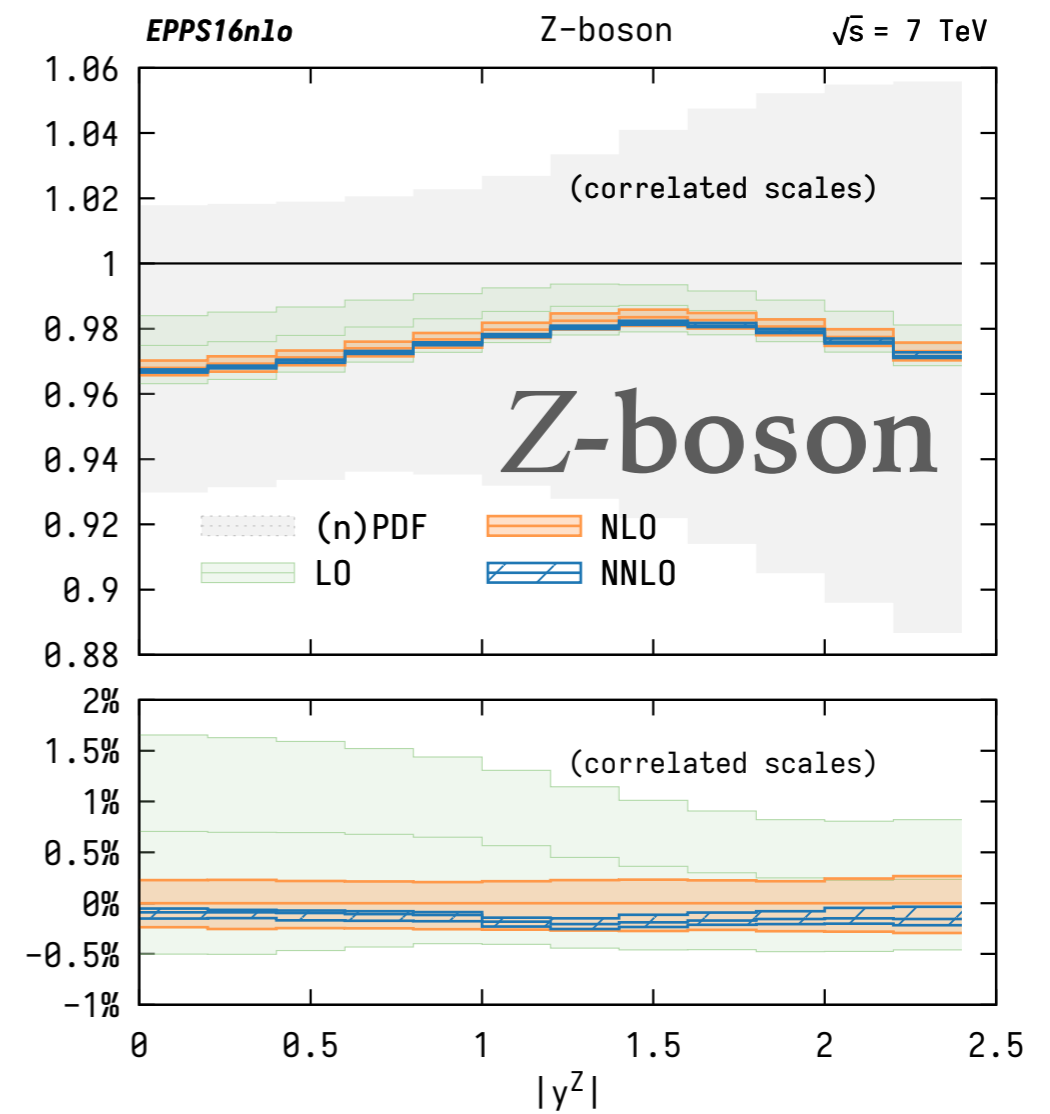
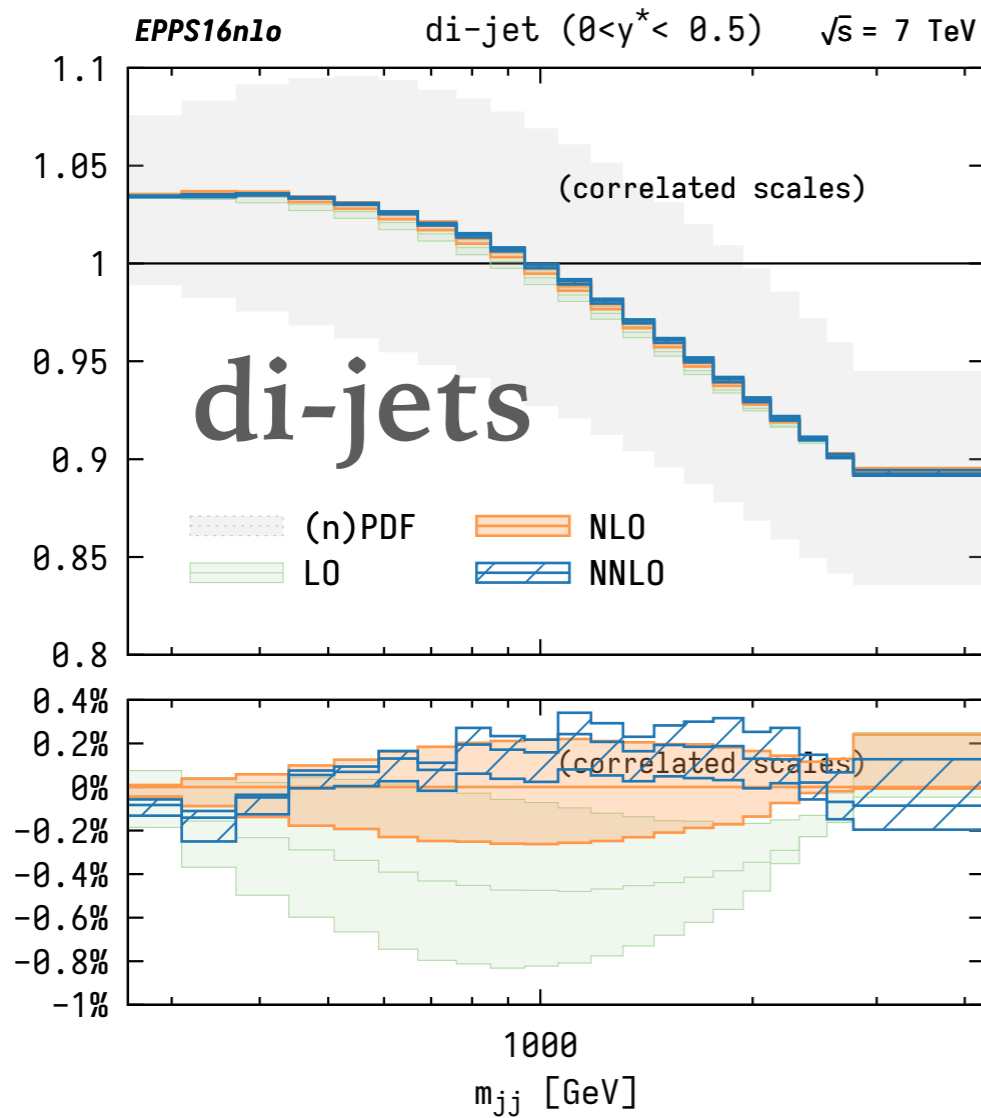
$$\sigma(\mu_0, \alpha_s(\mu_0)) = \left(\frac{\alpha_s(\mu_0)}{2\pi}\right)^n \sigma^{(0)} + \left(\frac{\alpha_s(\mu_0)}{2\pi}\right)^{n+1} \sigma^{(1)} + \left(\frac{\alpha_s(\mu_0)}{2\pi}\right)^{n+2} \sigma^{(2)} + \left(\frac{\alpha_s(\mu_0)}{2\pi}\right)^{n+3} \sigma^{(3)} + \mathcal{O}(\alpha_s^{n+4}).$$

$$\begin{aligned} \sigma(\mu_R, \mu_F, \alpha_s(\mu_R), L_R, L_F) = & \left(\frac{\alpha_s(\mu_R)}{2\pi}\right)^n \hat{\sigma}_{ij}^{(0)} \otimes f_i(\mu_F) \otimes f_j(\mu_F) \\ & + \left(\frac{\alpha_s(\mu_R)}{2\pi}\right)^{n+1} \hat{\sigma}_{ij}^{(1)} \otimes f_i(\mu_F) \otimes f_j(\mu_F) \\ & + L_R \left(\frac{\alpha_s(\mu_R)}{2\pi}\right)^{n+1} n \beta_0 \hat{\sigma}_{ij}^{(0)} \otimes f_i(\mu_F) \otimes f_j(\mu_F) \quad \leftarrow \text{Renormalization group equation} \\ & + L_F \left(\frac{\alpha_s(\mu_R)}{2\pi}\right)^{n+1} \left[-\hat{\sigma}_{ij}^{(0)} \otimes f_i(\mu_F) \otimes \left(P_{jk}^{(0)} \otimes f_k(\mu_F)\right) \right. \\ & \quad \left. - \hat{\sigma}_{ij}^{(0)} \otimes \left(P_{ik}^{(0)} \otimes f_k(\mu_F)\right) \otimes f_j(\mu_F) \right] \quad \leftarrow \text{DLAGP evolution} \end{aligned}$$

SCALE UNCERTAINTIES — CORRELATED

$$R^{A/B}(\mu_R, \mu_F) = \frac{\sigma^A(\mu_R, \mu_F)}{\sigma^B(\mu_R, \mu_F)}$$

7-point

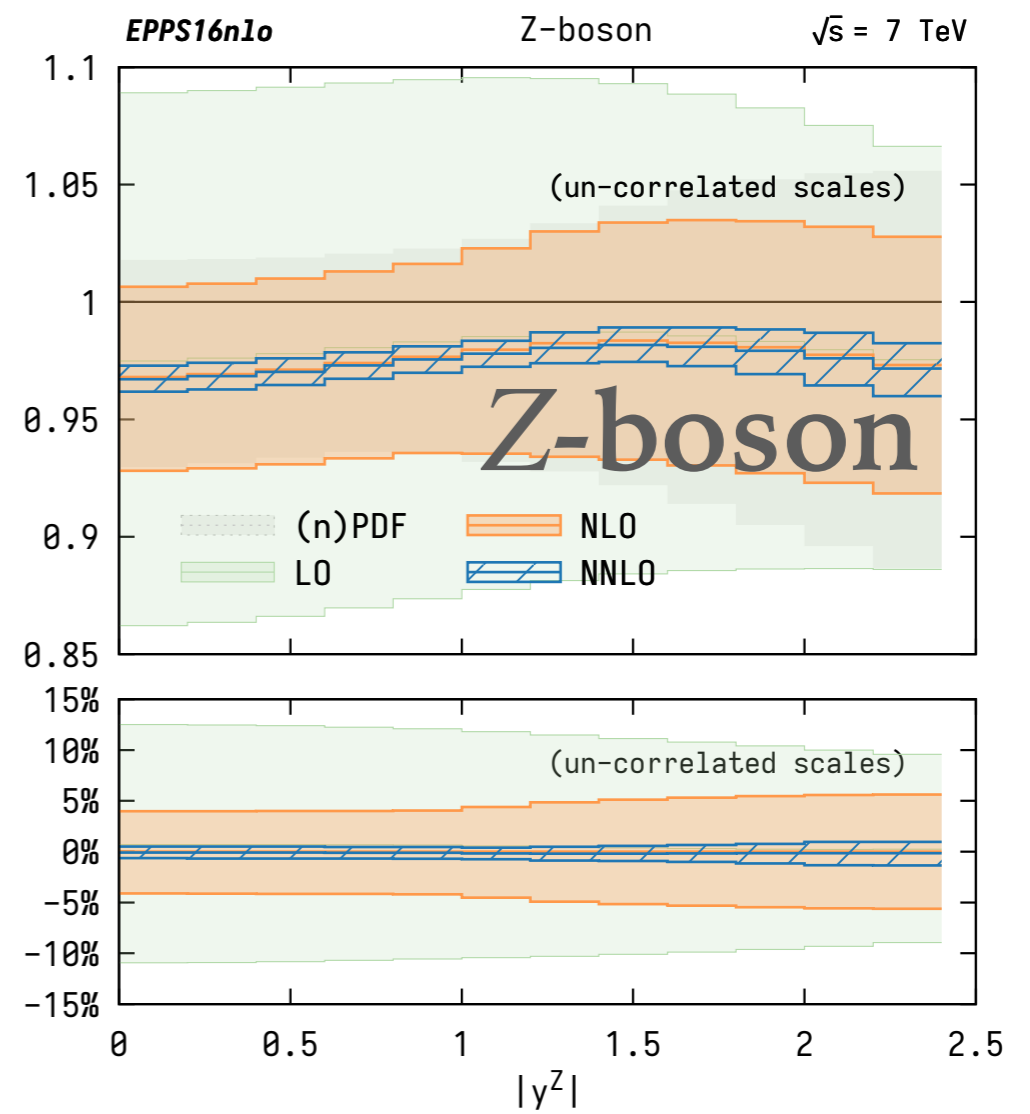
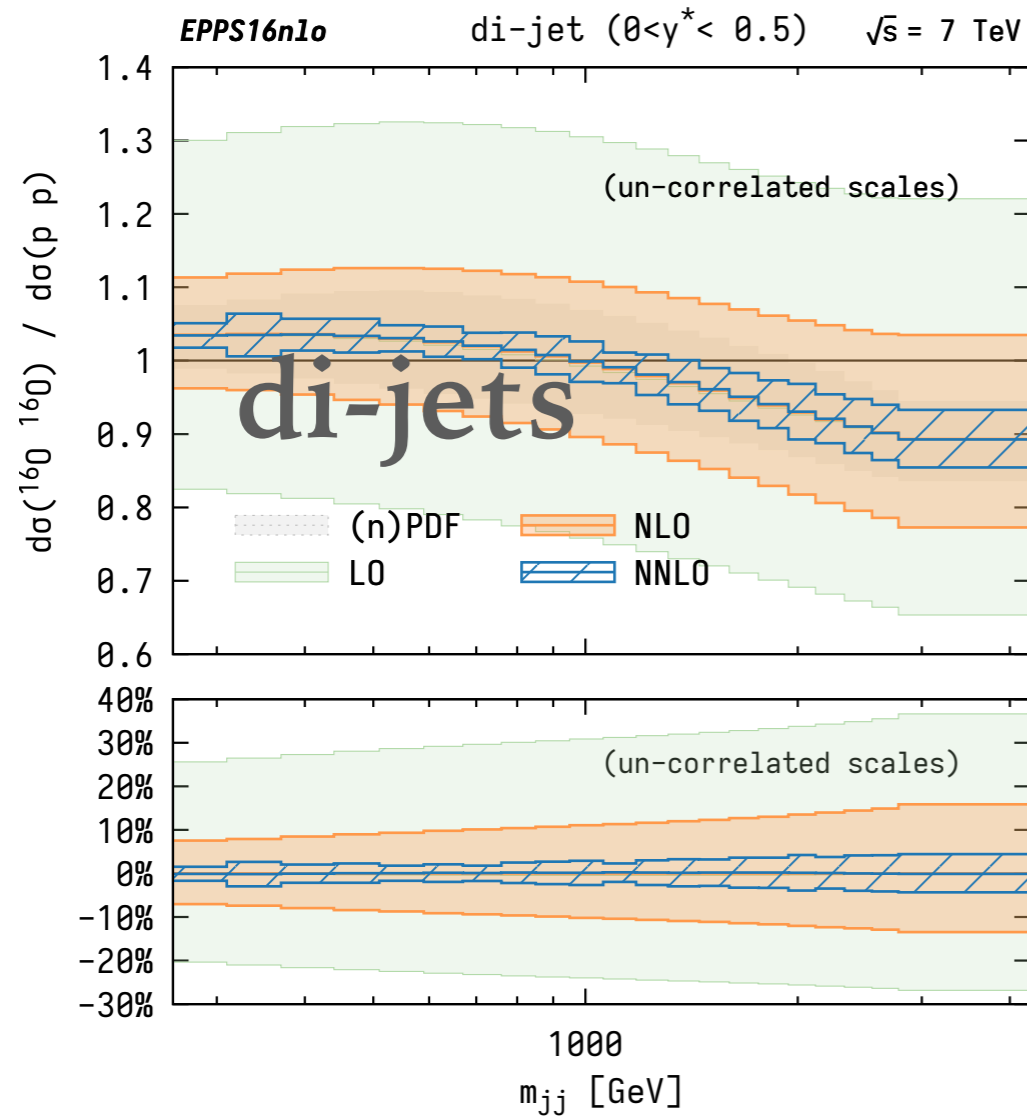


- higher-order corrections $< 1\%$
- **realistically** captured by uncertainty estimates

SCALE UNCERTAINTIES — UN-CORRELATED

$$R^{A/B}(\mu_R^A, \mu_B^B, \mu_F^A, \mu_F^B) = \frac{\sigma^A(\mu_R^A, \mu_F^A)}{\sigma^B(\mu_R^B, \mu_F^B)}$$

31-point



- higher-order corrections $< 1\%$
- **over-estimated** uncertainties: $\pm \mathcal{O}(10\%)$

NPDFs — IMPACT OF NEW DATA

following [Eskola, Paakkinen, Paukkunen '19]

→ reweight nPDFs using

CMS pPb dijet data →

[CMS 1805.04736]

→ **substantial reduction in uncertainties!**

danger of “fitting away” energy loss effects?!

→ observable insensitive:

$$R_{pPb}^{\text{norm}} = \left(\frac{1}{d\sigma_{pPb}/dp_T} \frac{d\sigma_{pPb}}{dp_T d\eta} \right) / \left(\frac{1}{d\sigma_{pp}/dp_T} \frac{d\sigma_{pp}}{dp_T d\eta} \right)$$

