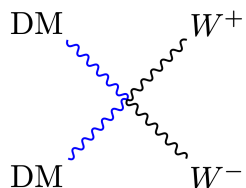


A model of electroweakly interacting non-abelian vector dark matter

Motivation & Summary

What is the nature of **electroweakly interacting spin-1 dark matter (DM)**?

- how to construct a spin-1 DM model w/ EW int.?
- What is a good probe for spin-1 DM?



Exchange symme. btw gauge groups \rightarrow **Z_2 -odd spin-1 DM w/ EW int.**
Next-generation direct detection can probe DM via **Higgs exchange**

Model

Exchange symmetry

field	spin	SU(3) _c	SU(2) ₀	SU(2) ₁	SU(2) ₂	U(1) _Y
q_L	$\frac{1}{2}$	3	1	2	1	$\frac{1}{6}$
u_R	$\frac{1}{2}$	3	1	1	1	$\frac{2}{3}$
d_R	$\frac{1}{2}$	3	1	1	1	$-\frac{1}{3}$
ℓ_L	$\frac{1}{2}$	1	1	2	1	$-\frac{1}{2}$
e_R	$\frac{1}{2}$	1	1	1	1	-1
Φ_1	0	1	2	2	1	0
Φ_2	0	1	1	2	2	0
H	0	1	1	2	1	$\frac{1}{2}$

$W_{0\mu}^a$ $W_{1\mu}^a$ $W_{2\mu}^a$

Symmetry trans. (after SSB)
 $\left[\sigma_1 \leftrightarrow \sigma_2, W_{0\mu}^a \leftrightarrow W_{2\mu}^a (a = 1, 2, 3) \right]$

Exchange symmetry
 \leftrightarrow **Z_2 -parity for physical states**

Z_2 -odd particles

$$V^0 = \frac{W_{0\mu}^3 - W_{2\mu}^3}{\sqrt{2}} \rightarrow \text{Spin-1 DM w/ EW int. } (m_V < m_{h_D})$$

$$V^\pm = \frac{W_{0\mu}^\pm - W_{2\mu}^\pm}{\sqrt{2}}$$

$$h_D = \frac{\sigma_1 - \sigma_2}{\sqrt{2}}$$

$$\left(\begin{array}{c} \sigma_1 + \sigma_2 \\ \sigma_3 \end{array} \right) \xrightarrow{\text{mixing } \phi_h} \begin{array}{c} h' \\ h \end{array}$$

(SM Higgs)

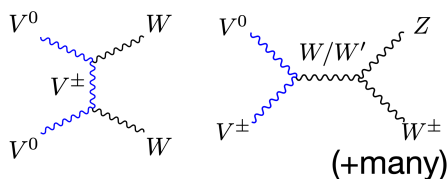
BSM spectrum

vector scalar Z_2 mass
 $Z' \ W'^\pm \ h'$ even } $\sim v_\Phi$
 $V^0 \ V^\pm \ h_D$ odd }

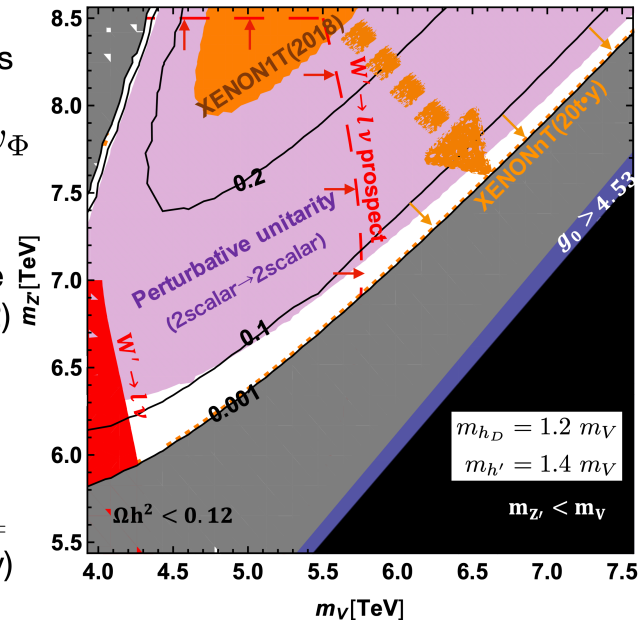
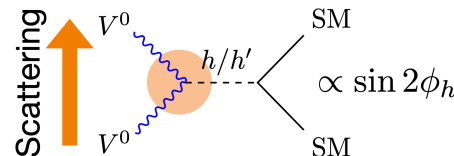
Results

Contours of Higgs mixing angle (determined to reach $\Omega h^2 = 0.12$)

•EW channel

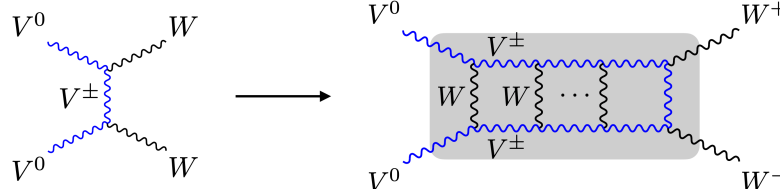


•Higgs channel



Next-generation direct detection can probe Higgs contributions in Ωh^2 determination

Future Work



A chance to discriminate DM spins in indirect detection!
 (Sommerfeld effect is relevant in annihilation processes)

Scalar sector

$$\Phi_j = \begin{pmatrix} \frac{v_\Phi + \sigma_j + i\pi_j^0}{\sqrt{2}} & i\pi_j^+ \\ i\pi_j^- & \frac{v_\Phi + \sigma_j - i\pi_j^0}{\sqrt{2}} \end{pmatrix}, H = \begin{pmatrix} i\pi_3^+ \\ \frac{v + \sigma_3 - i\pi_3^0}{\sqrt{2}} \end{pmatrix} \quad (j = 1, 2)$$

3 neutral scalars are physical after SSB

Model & Symmetry breaking structure

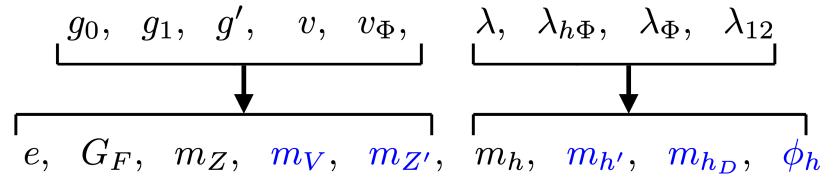
BSM Lagrangian

$$\mathcal{L} \supset -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \sum_{j=0}^2 \sum_{a=1}^3 \frac{1}{4} W_{j\mu\nu}^a W_j^{a\mu\nu} + D_\mu H^\dagger D^\mu H + \frac{1}{2} \text{tr} D_\mu \Phi_1^\dagger D^\mu \Phi_1 + \frac{1}{2} \text{tr} D_\mu \Phi_2^\dagger D^\mu \Phi_2 - V_{\text{scalar}},$$

$$V_{\text{scalar}} = m^2 H^\dagger H + m_\Phi^2 \text{tr} (\Phi_1^\dagger \Phi_1) + m_\Phi^2 \text{tr} (\Phi_2^\dagger \Phi_2) + \lambda (H^\dagger H)^2 + \lambda_\Phi (\text{tr} (\Phi_1^\dagger \Phi_1))^2 + \lambda_\Phi (\text{tr} (\Phi_2^\dagger \Phi_2))^2 + \lambda_{h\Phi} H^\dagger H \text{tr} (\Phi_1^\dagger \Phi_1) + \lambda_{h\Phi} H^\dagger H \text{tr} (\Phi_2^\dagger \Phi_2) + \lambda_{12} \text{tr} (\Phi_1^\dagger \Phi_1) \text{tr} (\Phi_2^\dagger \Phi_2).$$

g_0 : gauge coupling for $SU(2)_0$ & $SU(2)_2$
 g_1 : gauge coupling for $SU(2)_1$

Parameters



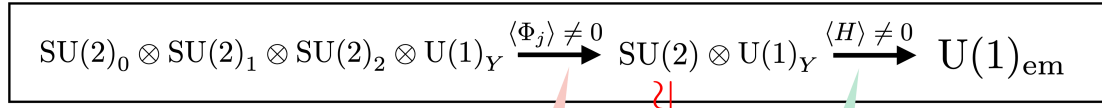
Free parameters (TeV scale)
 $\{ m_V, m_{Z'}, m_{h'}, m_{h_D}, \phi_h \}$

Gauge transformation

$$U_n = \exp[i\theta_n(x)] \quad (n = 0, 1, 2)$$

$$\Phi_1 \mapsto U_0 \Phi_1 U_1^\dagger \quad \Phi_2 \mapsto U_2 \Phi_2 U_1^\dagger \quad H \mapsto U_1 H$$

Symmetry Breaking



$$\langle \Phi_1 \rangle = \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} v_\Phi & 0 \\ 0 & v_\Phi \end{pmatrix}$$

$$\langle H \rangle = \begin{pmatrix} 0 \\ \frac{v}{\sqrt{2}} \end{pmatrix}$$

$(v_\Phi \gg v)$
 \uparrow \uparrow
 $\mathcal{O}(1) \text{ TeV}$ $\mathcal{O}(100) \text{ GeV}$

$\langle \Phi_1 \rangle, \langle \Phi_2 \rangle$ are invariant under the following transformations

- (1) Gauge trans. w/ $U_0 = U_1 = U_2$
 $U_0 \langle \Phi_1 \rangle U_1^\dagger = \langle \Phi_1 \rangle$
 $U_2 \langle \Phi_2 \rangle U_1^\dagger = \langle \Phi_2 \rangle$
- (2) Exchange trans.
 $\langle \Phi_1 \rangle \leftrightarrow \langle \Phi_2 \rangle$

Generators of $SU(2)_{0,1,2}$ are identified
 \rightsquigarrow **$SU(2)_L$ gauge symmetry**

Exchange symmetry still alive
 \rightarrow **Z_2 parity structure**

Features of EW spin-1 DM (Compared w/ Wino DM)

Z₂-odd vectors

$$V^0 = \frac{W_{0\mu}^3 - W_{2\mu}^3}{\sqrt{2}} \text{ (neutral)}$$

$$V^\pm = \frac{W_{0\mu}^\pm - W_{2\mu}^\pm}{\sqrt{2}} \text{ (charged)}$$

DM

“V-particles”
SU(2)_L triplet-like features

Mass relations

Mass degeneracy btw V-particles is broken by EW radiative corrections

$$\begin{cases} \text{tree: } m_{V^0}^2 = m_{V^\pm}^2 = \frac{g_0^2 v_\Phi^2}{4} \quad (\equiv m_V^2) \\ \text{1-loop: } \delta_{m_V} \equiv m_{V^\pm} - m_{V^0} \simeq 168 \text{ MeV} \end{cases}$$

Almost same as Wino!

→ Coannihilation is relevant (efficient annihilation through EW int.)

Wino(Spin-1/2) vs V-particles(Spin-1)

		vs	
Spin	1/2 (Majorana fermion) [SU(2) _L triplet, Y=0]		1 (vector)
Mass difference	~ 166 MeV		~ 168 MeV
Annihilation	EW		EW + Higgs exchange
Scattering	tree-level: None loop-level: EW		tree-level: Higgs exchange loop-level: EW
Z ₂ -even vectors	—		Z', W'

Features of electroweakly interacting DM?
(spin-1/2 vs spin-1)

Common properties
Coannihilation is efficient
→ **Thermal relic mass**
≈ O(1) TeV

Different properties
good probes for spin-1 DM
→ **•Direct detection**
•W' search @LHC