



# High-temperature EW symmetry breaking: Reassessing the window for EW baryogenesis in Composite Higgs models

Oleksii Matsedonskyi

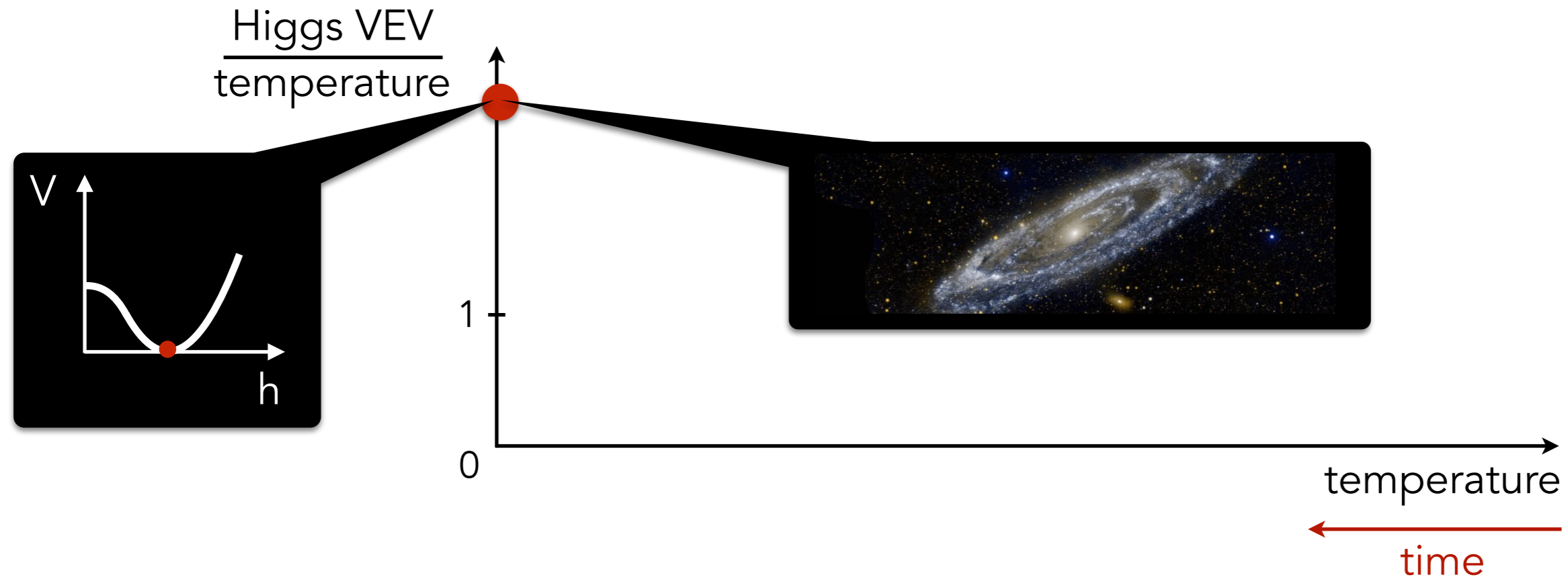
*HPNP 2021*

# Motivation

# Electroweak symmetry in the early Universe

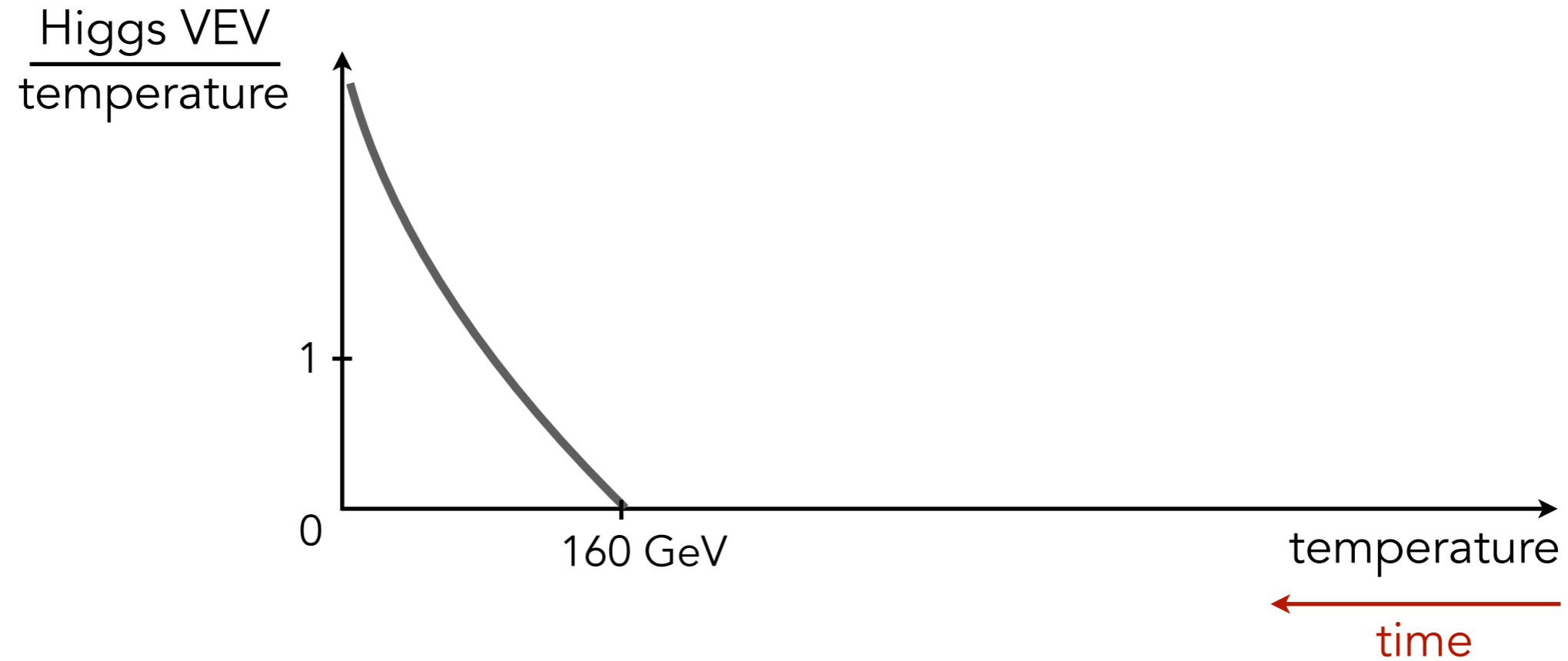


# Electroweak symmetry in the early Universe



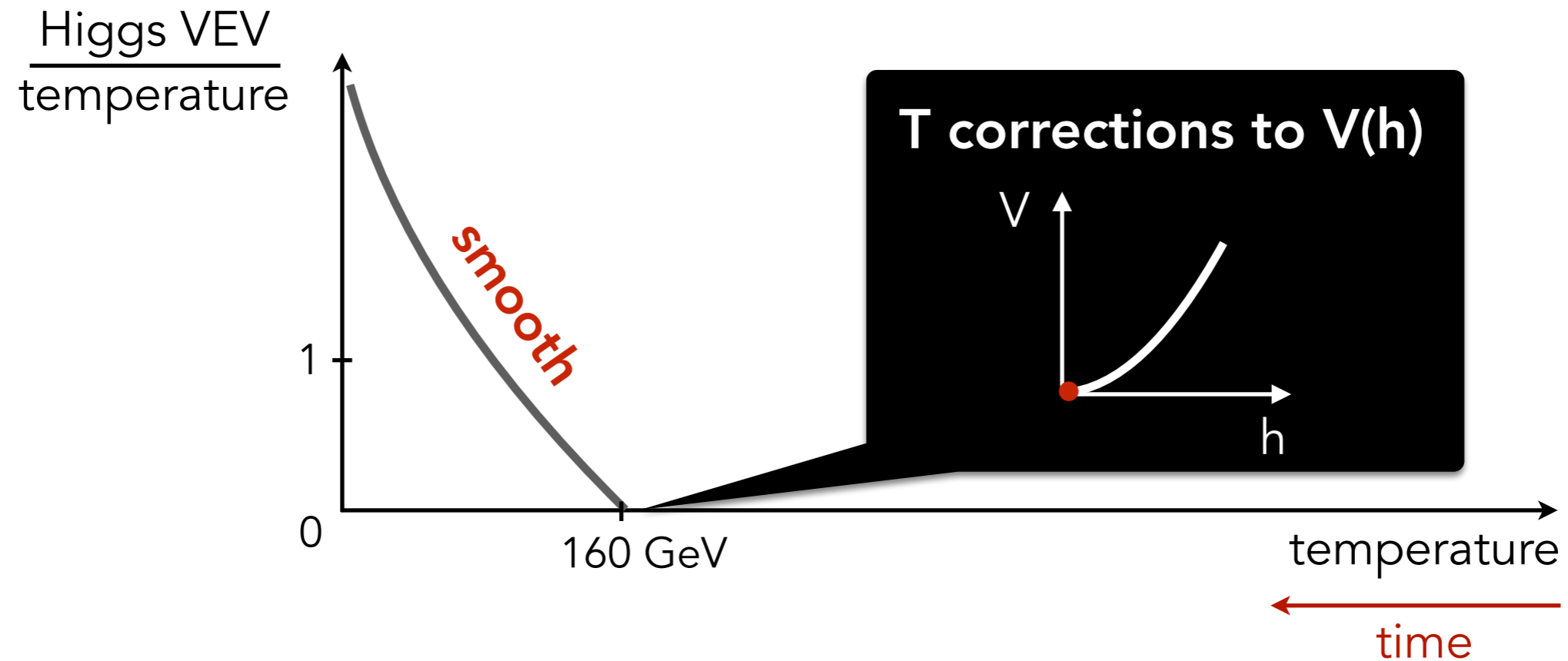
# Electroweak symmetry in the early Universe

in Standard Model: high-T symmetry restoration



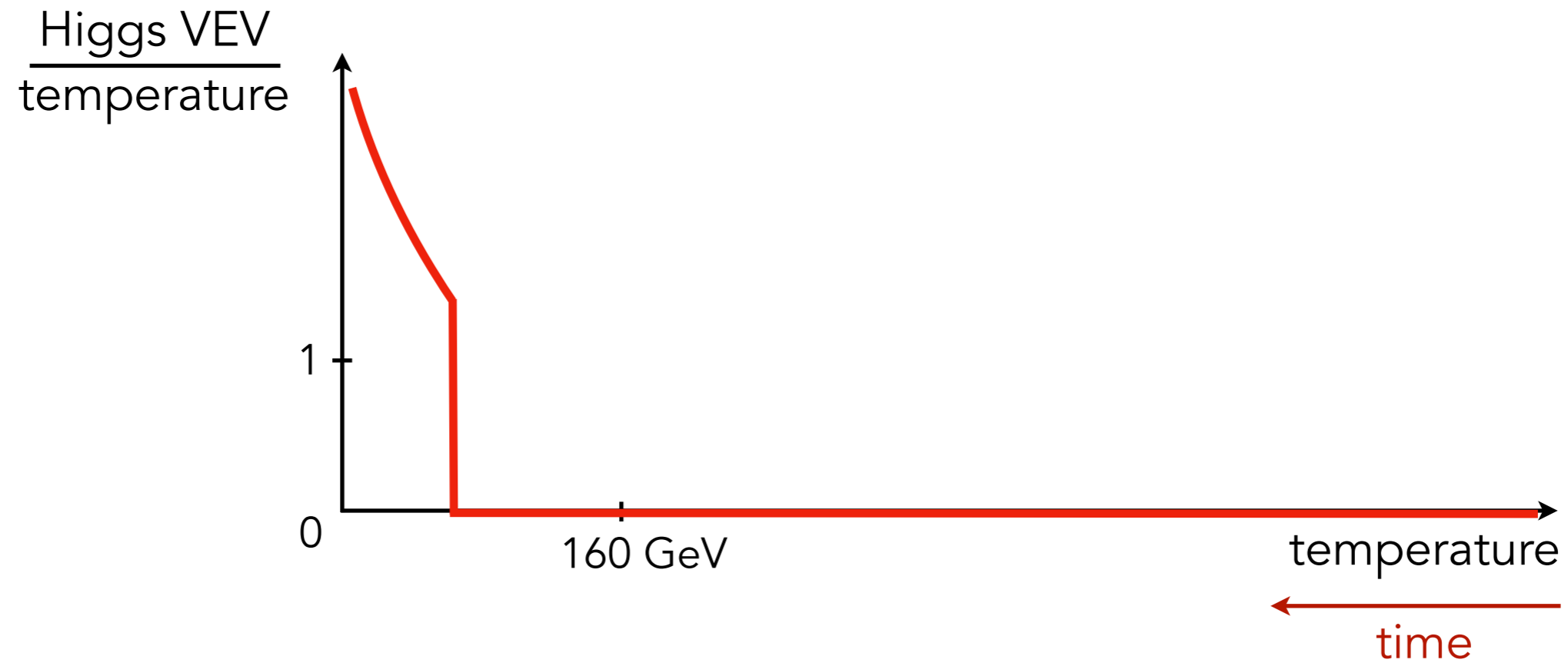
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in Standard Model: high-T symmetry restoration



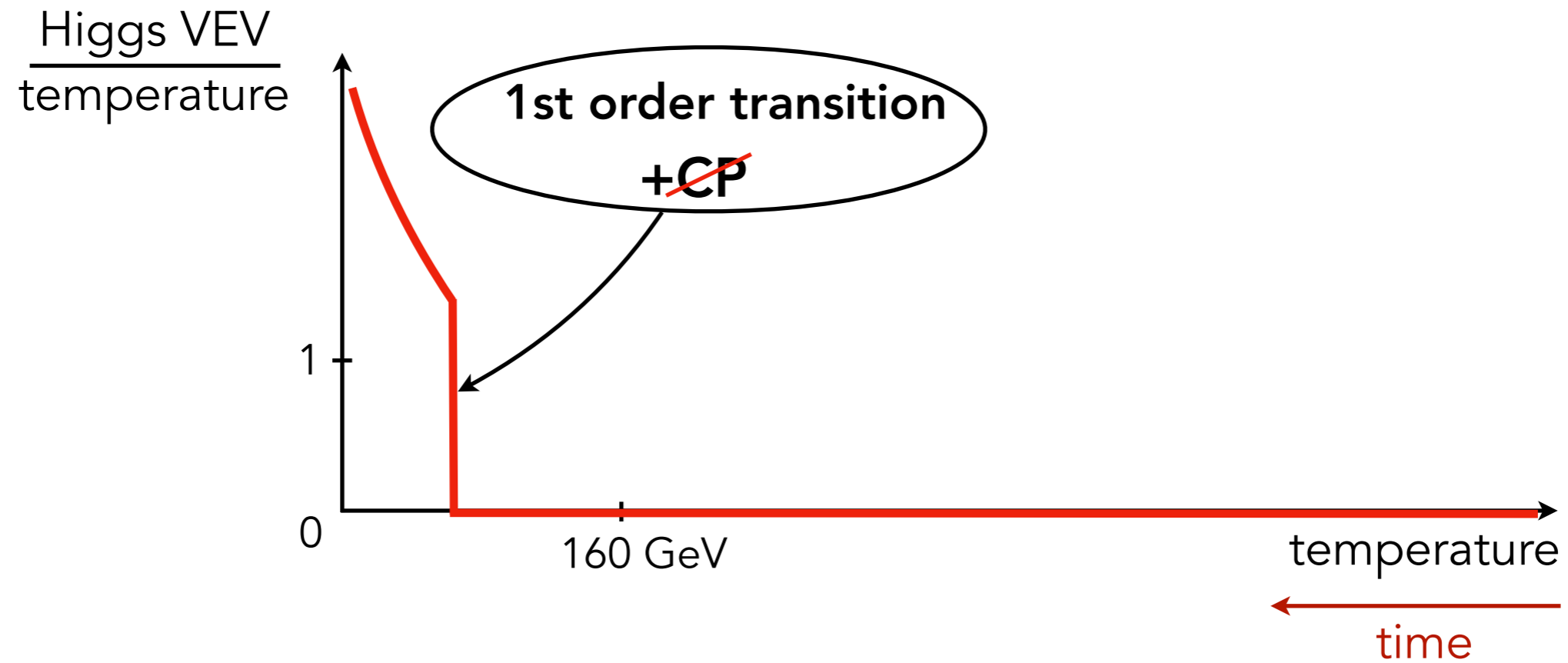
# Electroweak symmetry in the early Universe

in Electroweak Baryogenesis scenarios



# Electroweak symmetry in the early Universe

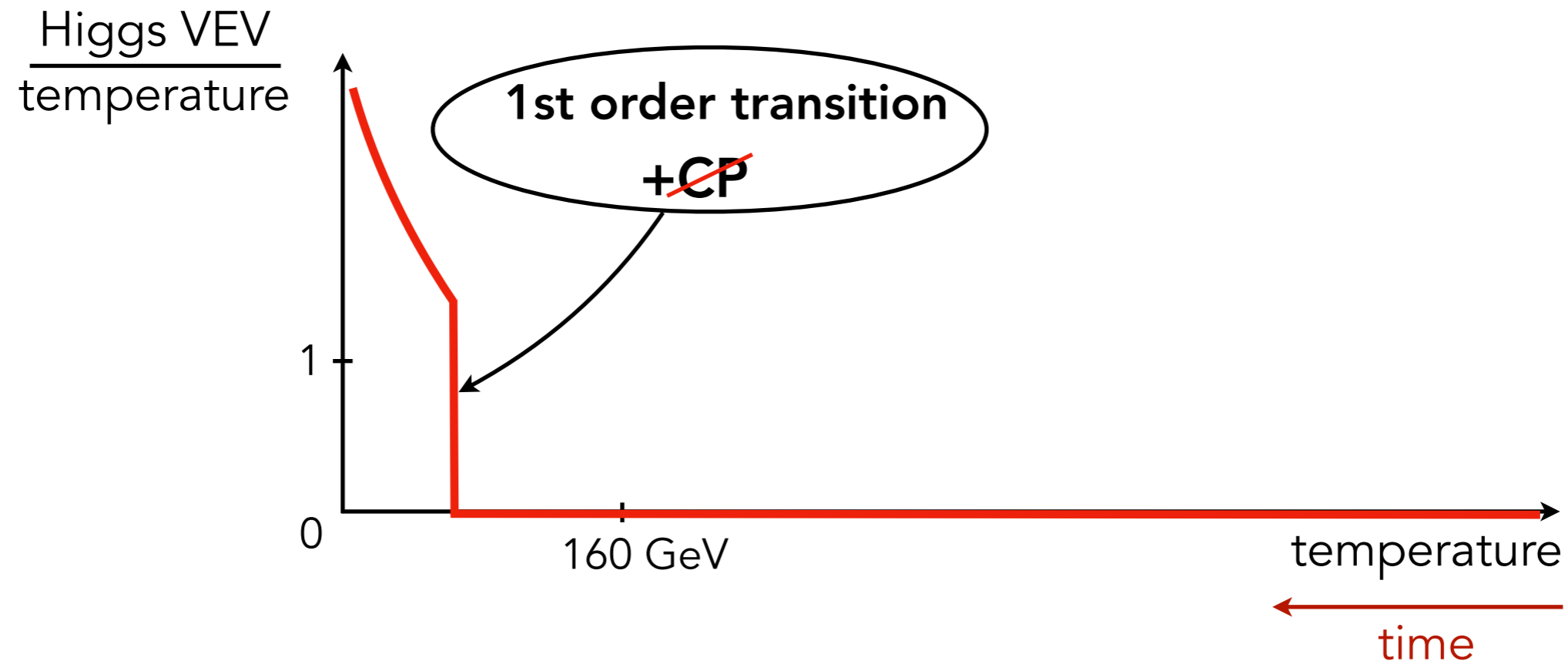
in Electroweak Baryogenesis scenarios





# Electroweak symmetry in the early Universe

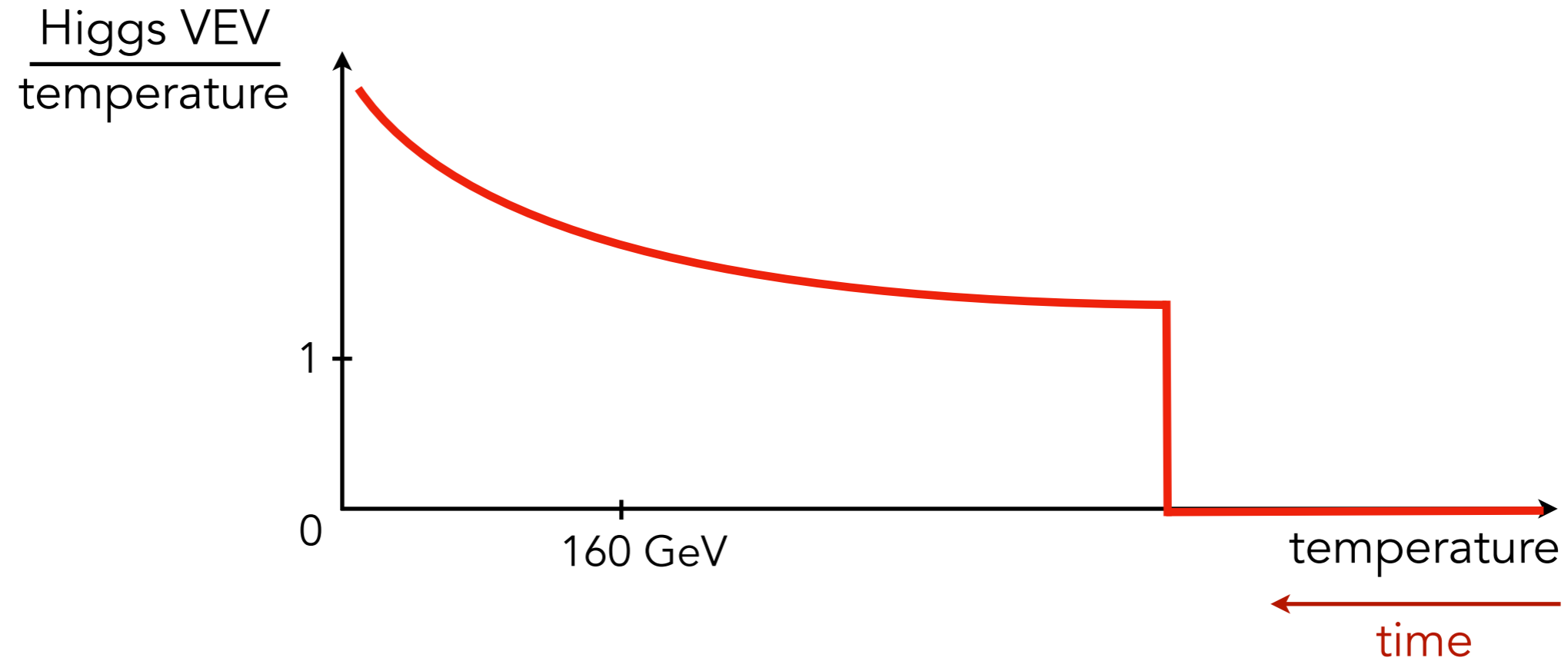
in Electroweak Baryogenesis scenarios



- new physics responsible for CP violation and first-order phase transition is at a few 100 GeV scale
- **good** because testable, **bad** (for some models) because overconstrained

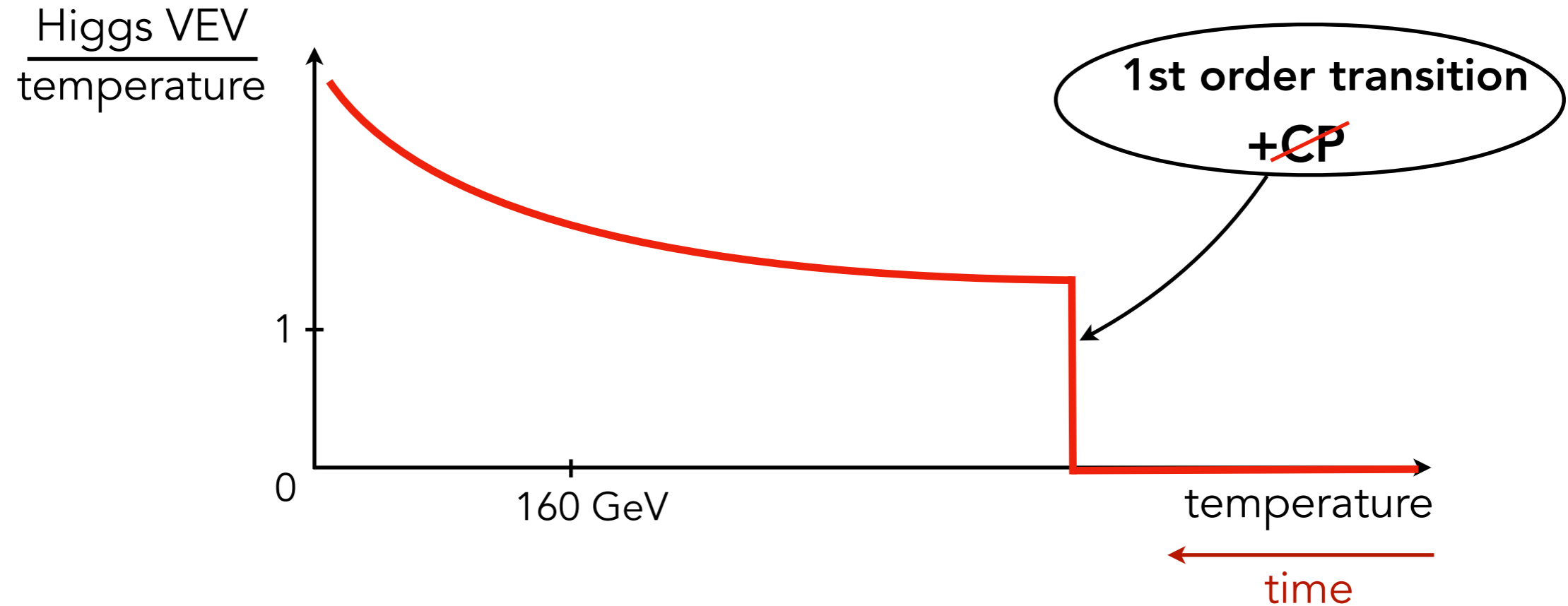
# Electroweak symmetry in the early Universe

What if?



# Electroweak symmetry in the early Universe

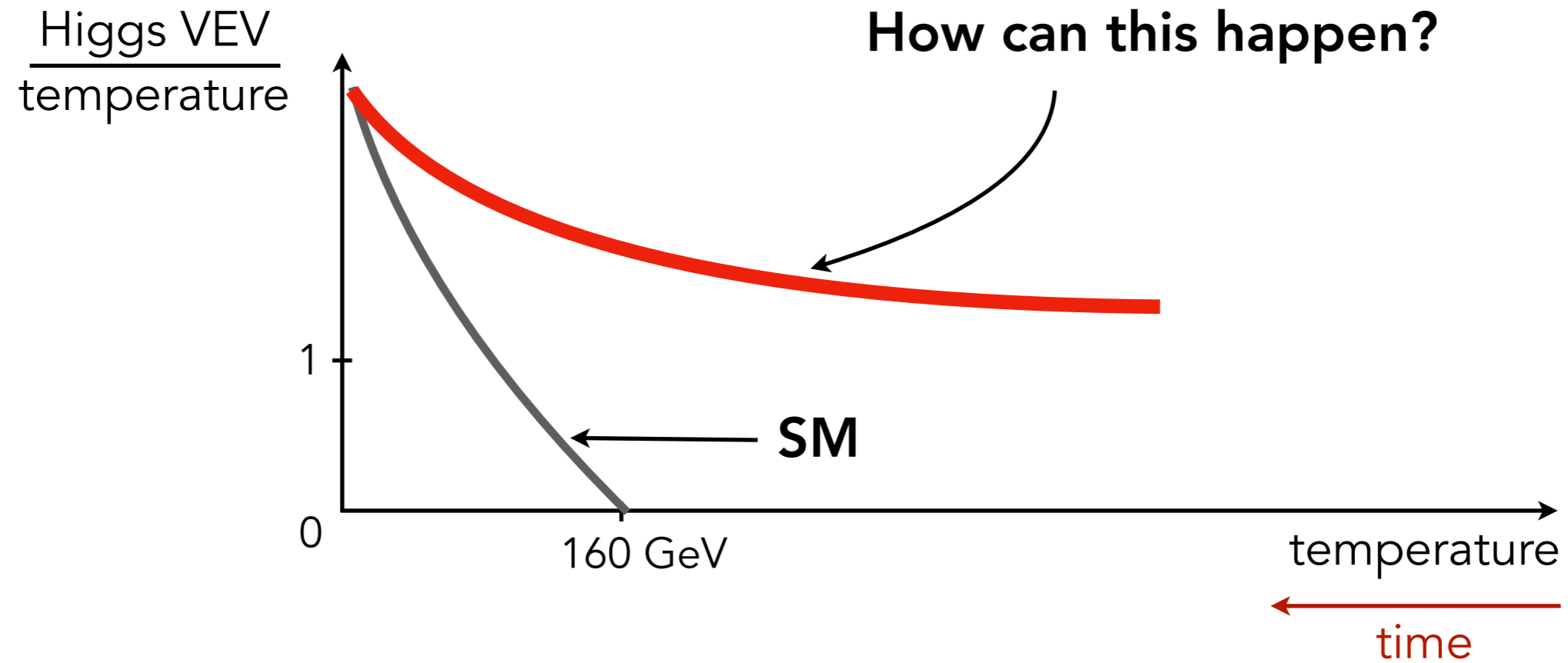
What if?



→ new physics responsible for CP violation and first-order phase transition is **far above** 100 GeV scale

→ **new phenomenology**

# Electroweak symmetry in the early Universe



EW Symmetry Breaking at High T

# Electroweak Symmetry Non-Restoration at High T

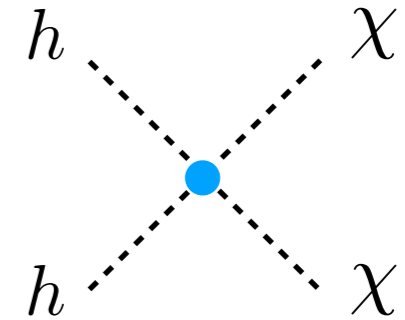
→ EW SNR was first proposed in:

Meade, Ramani, 1807.07578

Baldes, Servant, 1807.08770

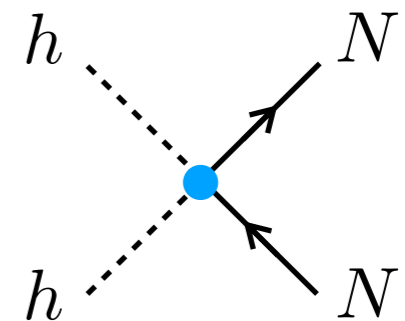
Glioti, Rattazzi, Vecchi, 1811.11740

with **new light scalars**



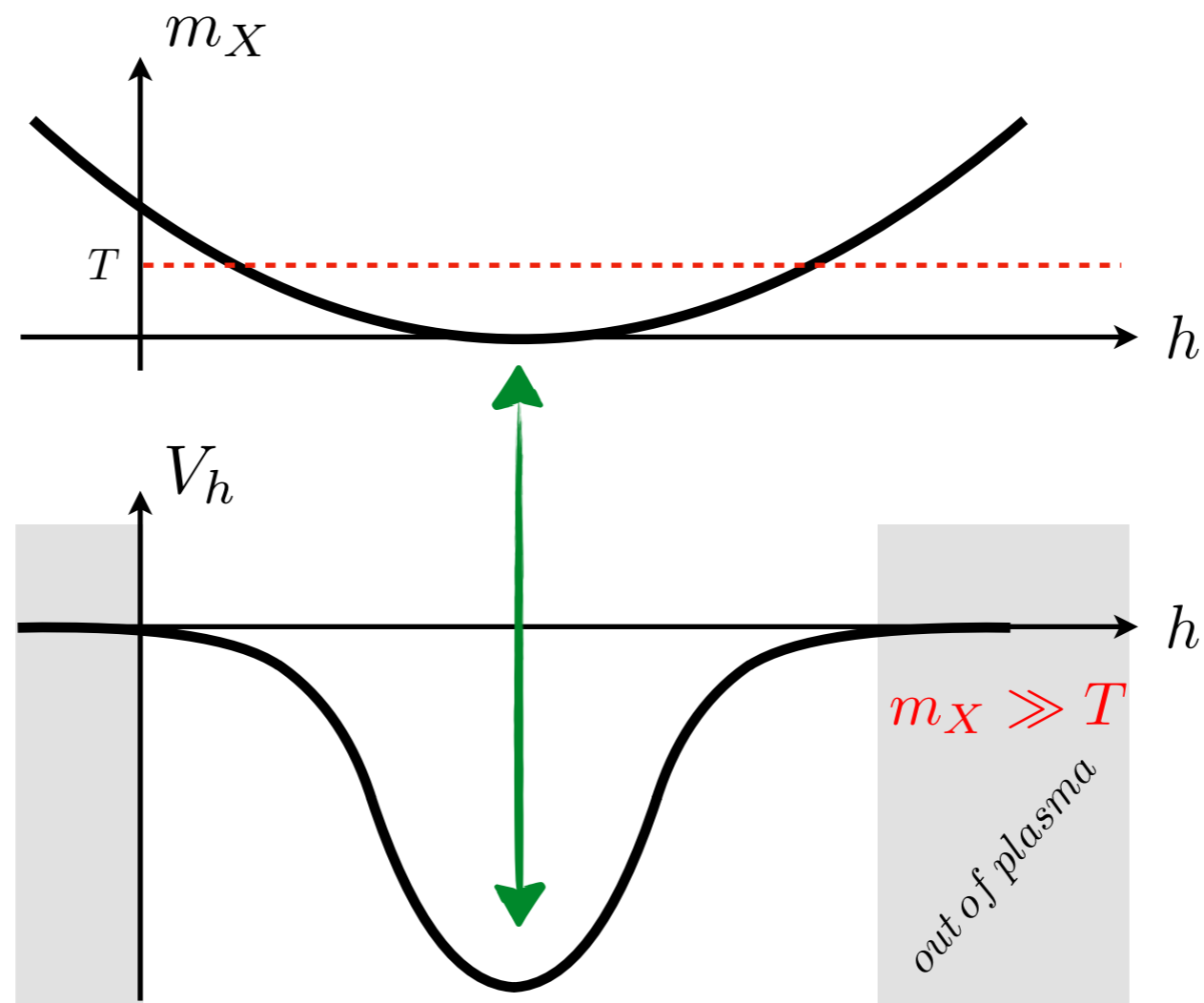
→ We showed how can be done with **new fermions**

OM, Servant, 2020.05174



# Electroweak Symmetry Non-Restoration at High T

→ If  $h$  potential is induced by plasma of particles with  $h$ -dependent mass:



# SNR with fermions

OM, Servant, 2020.05174

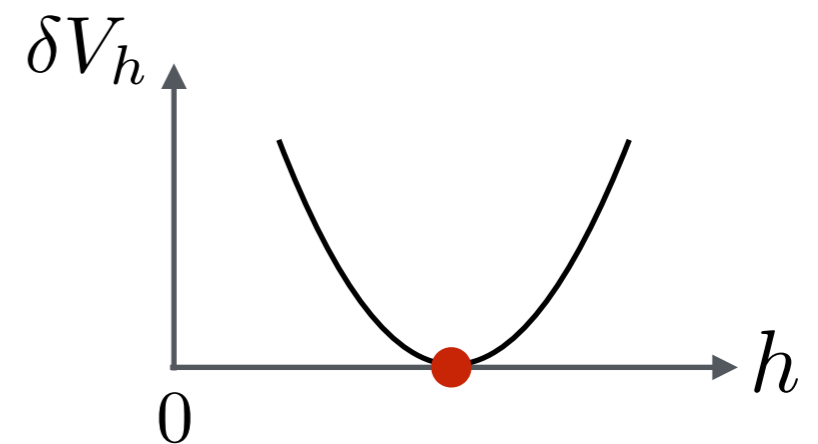
- Add  $n$  copies of new SM singlet Dirac fermion  $N$

$$\mathcal{L}_N = -m_N \bar{N}N + \lambda_N \bar{N}N h^2 / \Lambda$$

- $N$  mass is minimized at large  $h$

$$m_N(h) = m_N - \lambda_N h^2 / \Lambda = 0$$

$$\longrightarrow h^2 = m_N \Lambda / \lambda_N$$





# SNR with fermions

OM, Servant, 2020.05174

- Add  $n$  copies of new SM singlet Dirac fermion  $N$

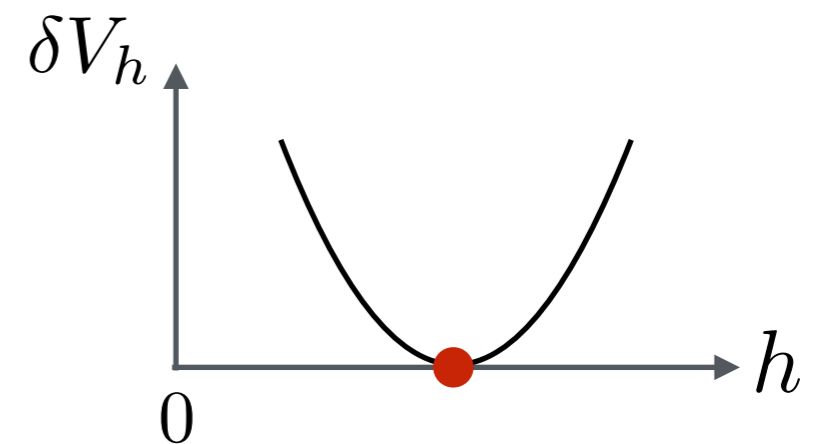
$$\mathcal{L}_N = -m_N \bar{N}N + \lambda_N \bar{N}N h^2 / \Lambda$$

↖ cutoff

- $N$  mass is minimized at large  $h$

$$m_N(h) = m_N - \lambda_N h^2 / \Lambda = 0$$

$$\longrightarrow h^2 = m_N \Lambda / \lambda_N$$



# SNR with fermions

OM, Servant, 2020.05174

why not 1?

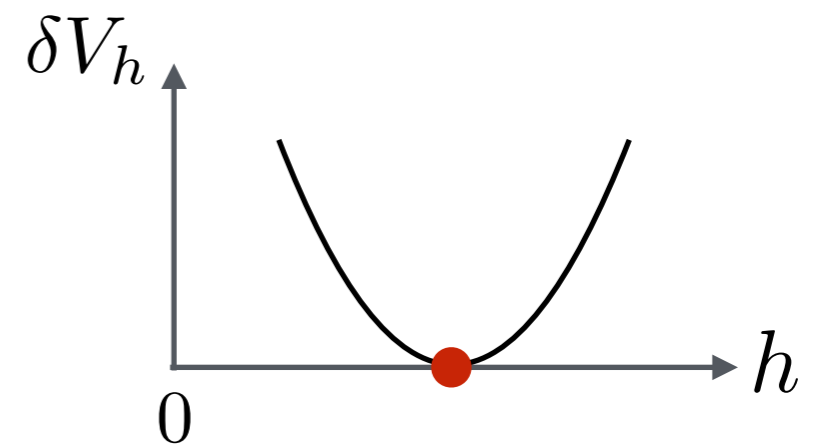
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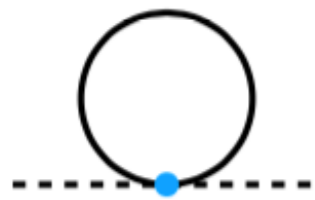


# SNR with fermions

OM, Servant, 2020.05174

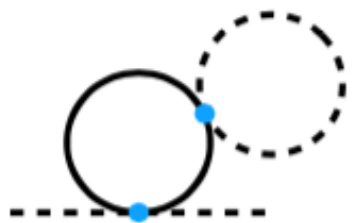
→ High-T perturbativity implies  $T_{\text{SNR}}^{\text{max}} \sim \sqrt{n} m_N$

calculable 1-loop (negative mass):

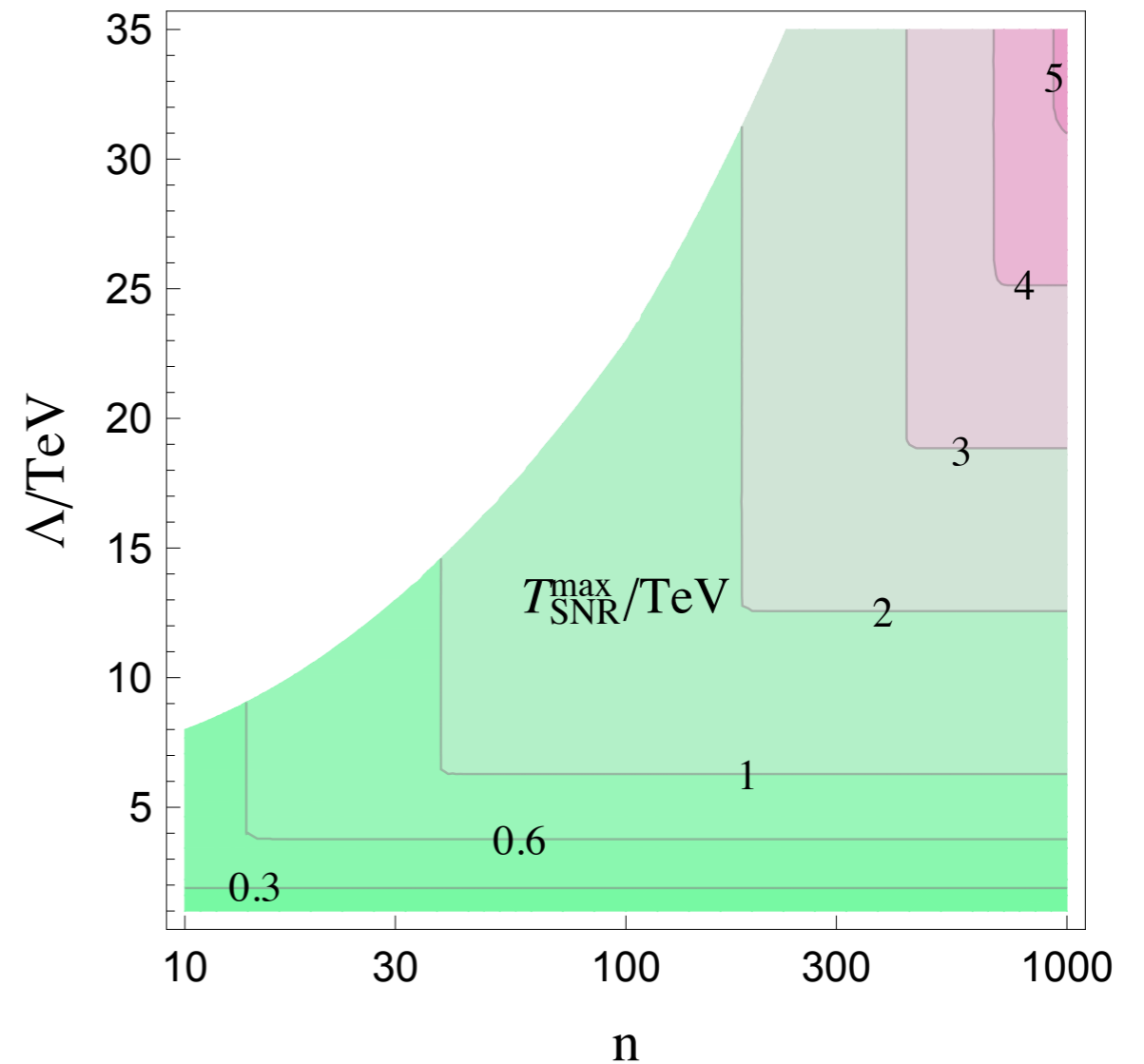


$$\propto n \lambda_N \frac{m_N}{\Lambda} \simeq 1$$

higher loops (out of control):



$$\beta = n \lambda_N^2 \frac{T^2}{\Lambda^2} \ll 1$$



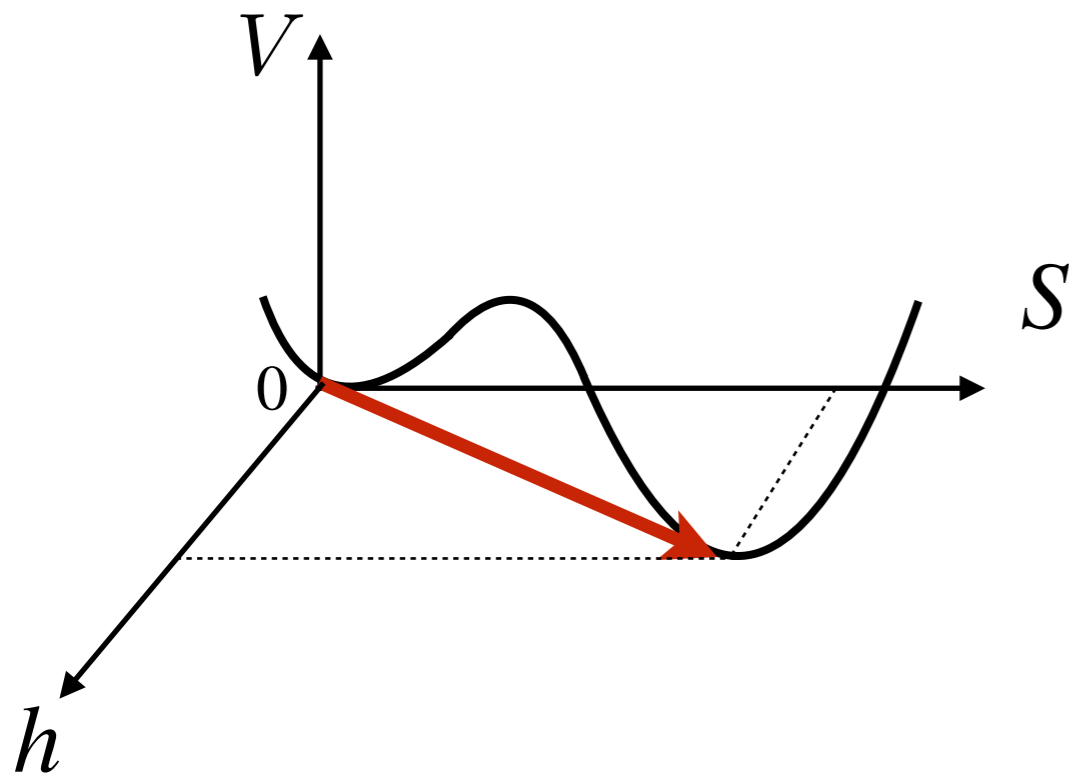
# Phenomenological Impact

# Phase Transition Temperature

work in progress

Bruggisser, VonHarling, OM, Servant

→ Usual way to get 1st order EW phase transition: add a new scalar  $S$



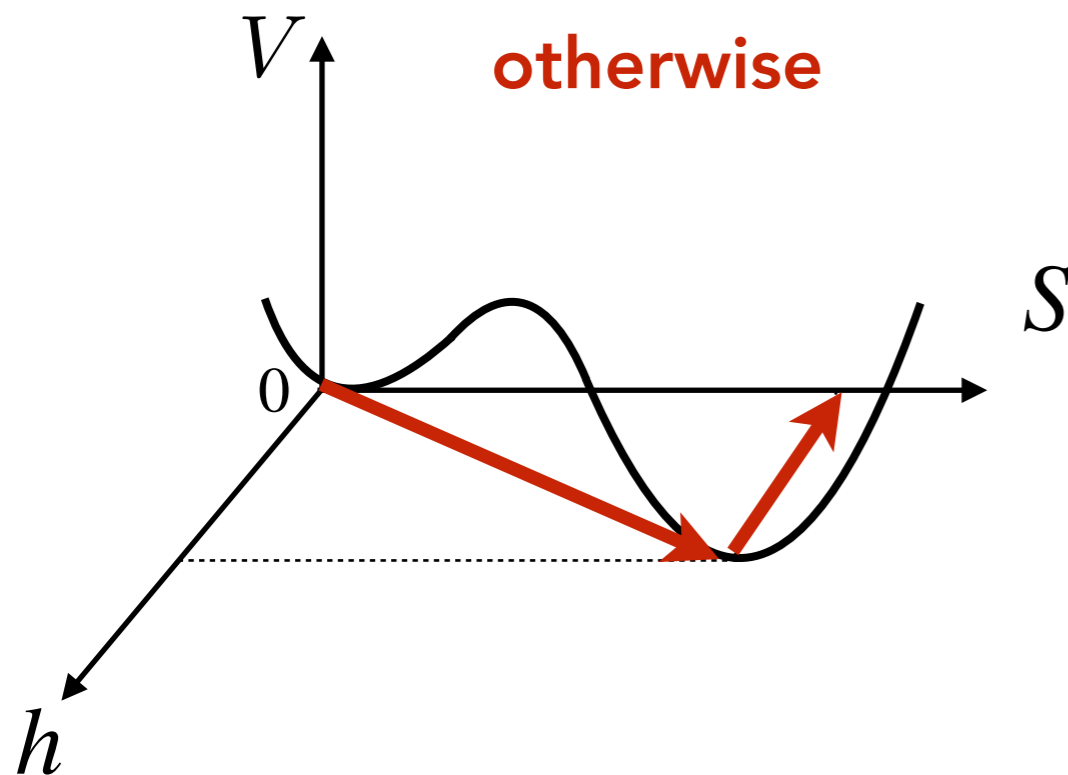
→ EWBG needs  $T < T$  of EW restoration

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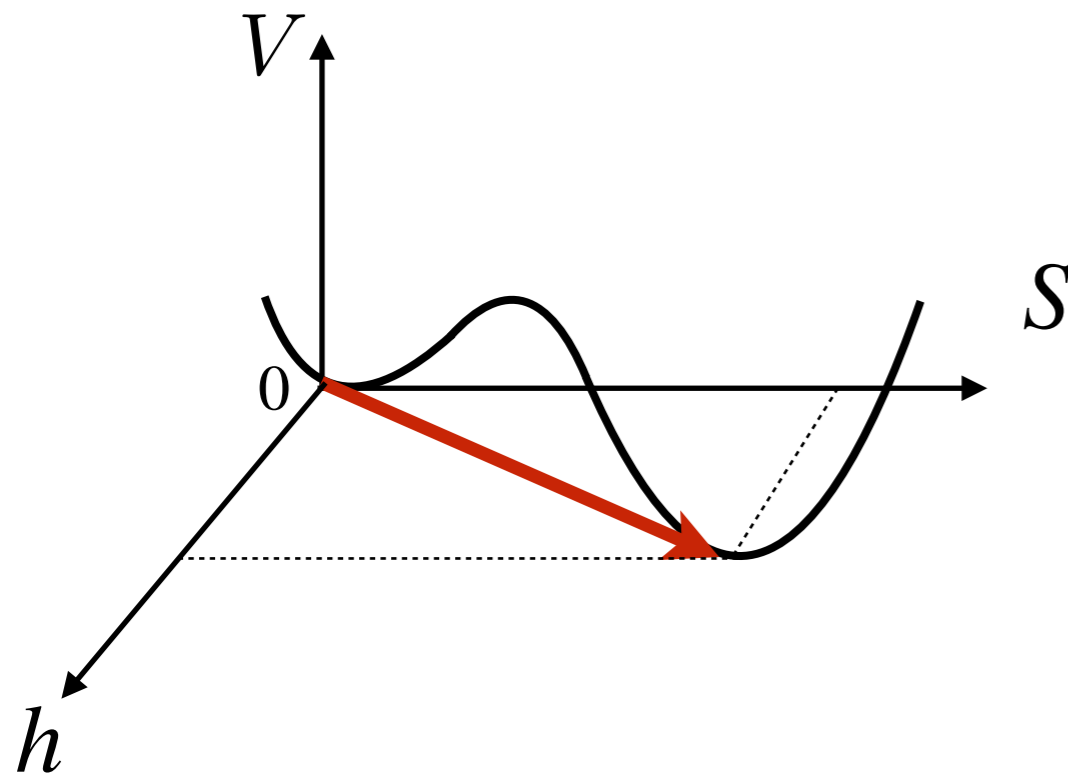
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→ EWBG needs  $T < T$  of EW restoration

→  $S$  phase transition releases latent heat

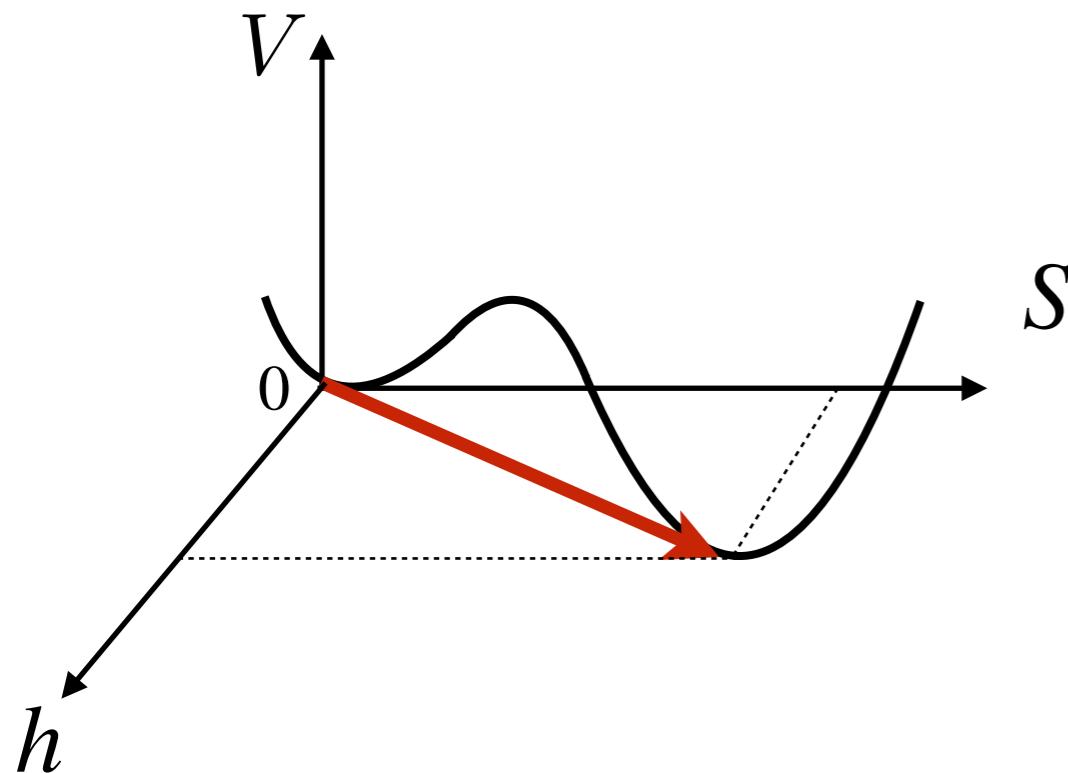
$$T^4 \propto m_S^2$$

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$$T^4 \propto m_S^2$$

⇒ for  $T$  restoration  $\sim 130$  GeV

$$m_S \lesssim \mathcal{O}(100 \text{ GeV})$$

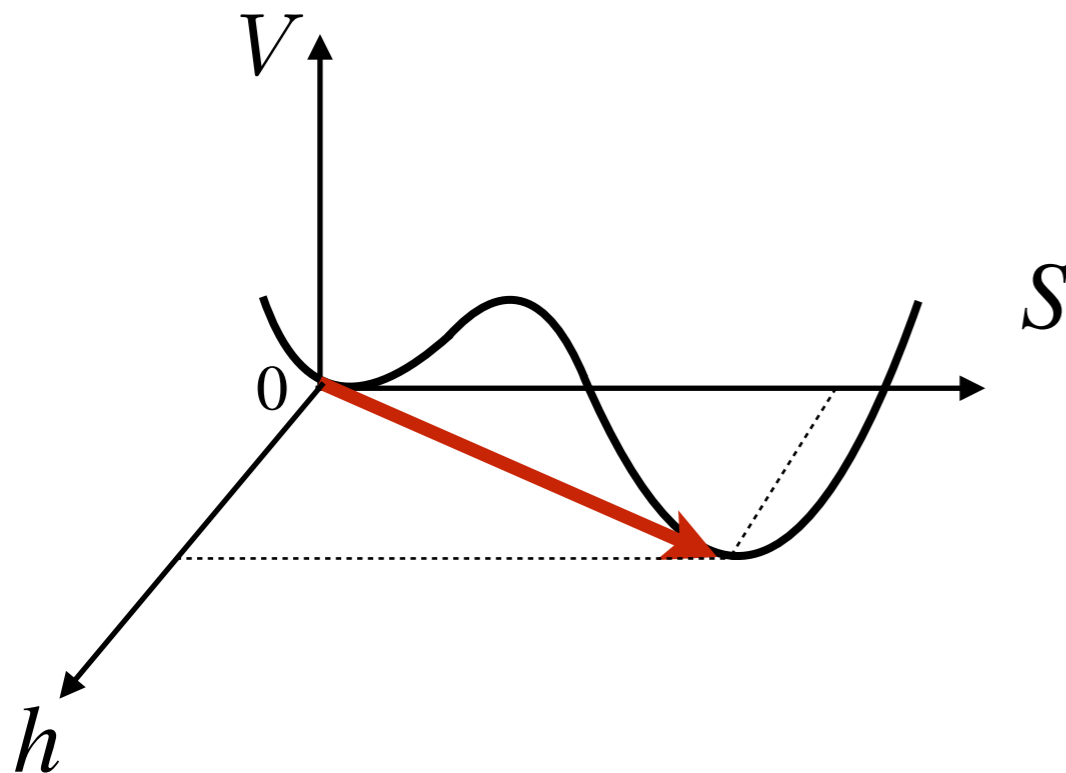


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Bruggisser, VonHarling, OM, Servant

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→ EWBG needs  $T < T$  of EW restoration

→  $S$  phase transition releases latent heat

$$T^4 \propto m_S^2$$

⇒ for  $T$  restoration  $\sim 1$  TeV

$$m_S \lesssim \text{O(few TeV)}$$

# Concrete Example: EWBG in Composite Higgs

work in progress

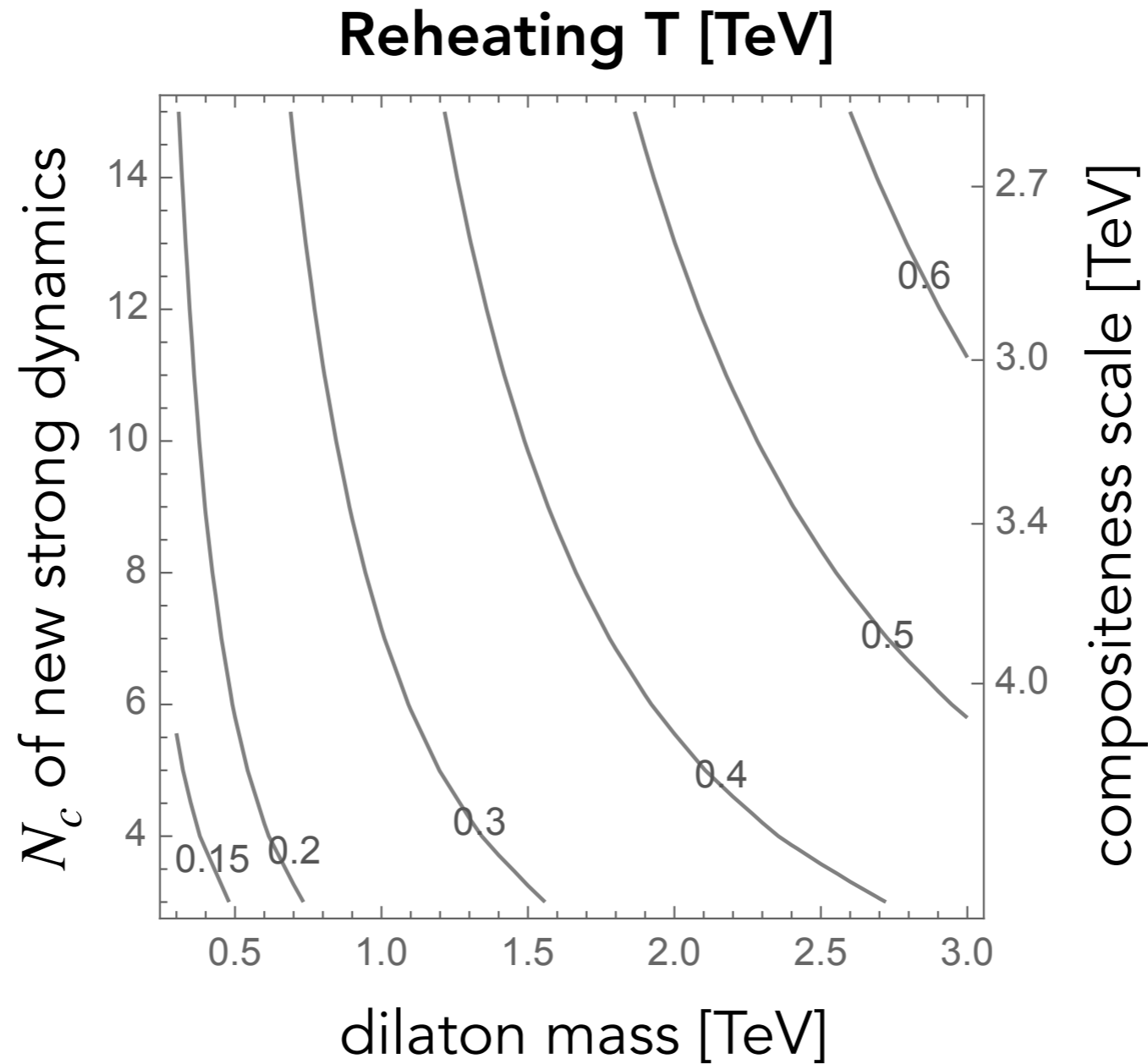
Bruggisser, VonHarling, OM, Servant

- Higgs is a bound state of new strong interactions
- New scalar triggering the first order phase transition
  - dilaton (PNGB of approximate conformal invariance)

# Concrete Example: EWBG in Composite Higgs

work in progress

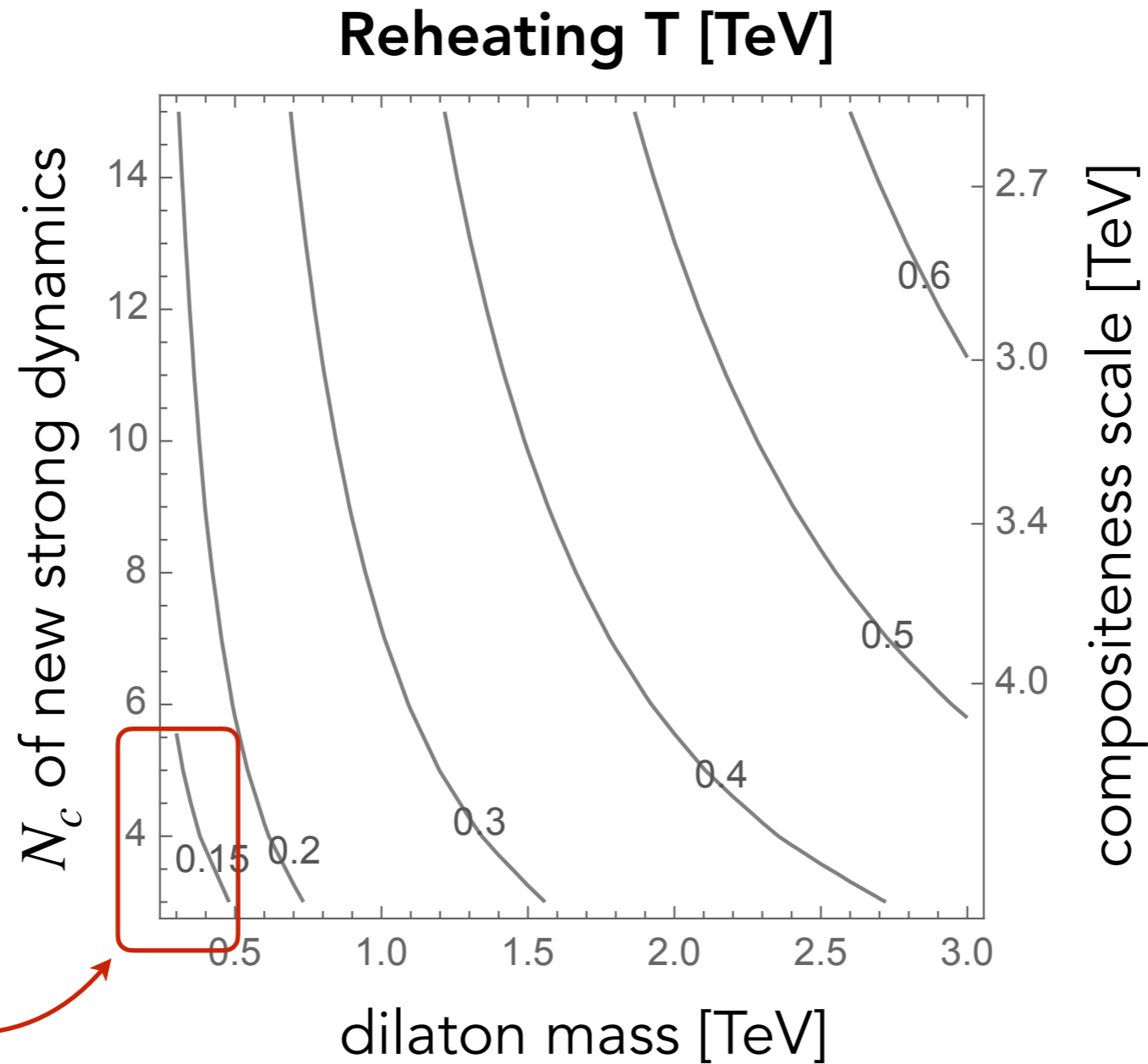
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# Concrete Example: EWBG in Composite Higgs

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Bruggisser, VonHarling, OM, Servant



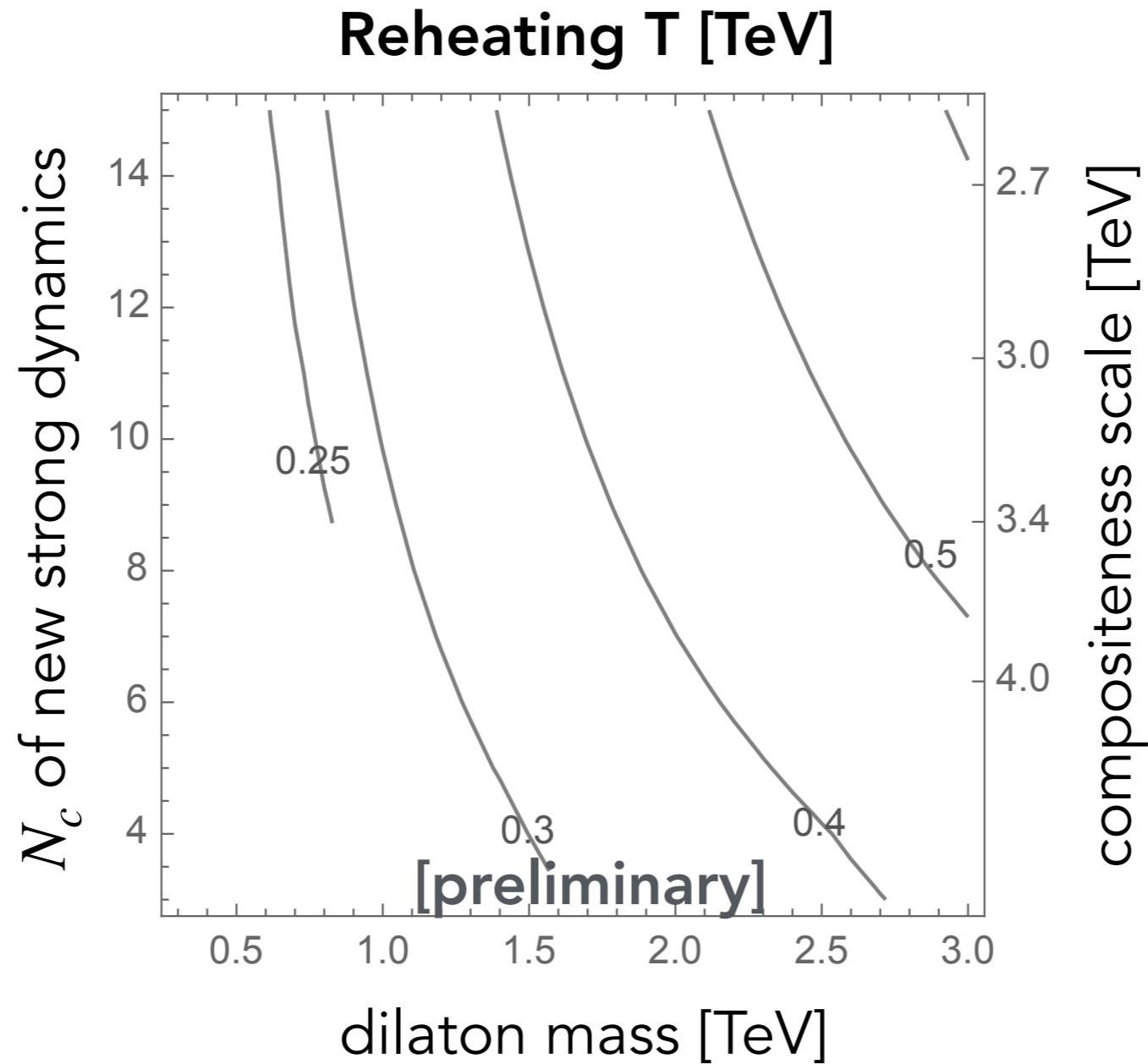
**tension with  
collider searches**

# Concrete Example: EWBG in Composite Higgs

work in progress

Bruggisser, VonHarling, OM, Servant

+ 10 SNR  
fermions

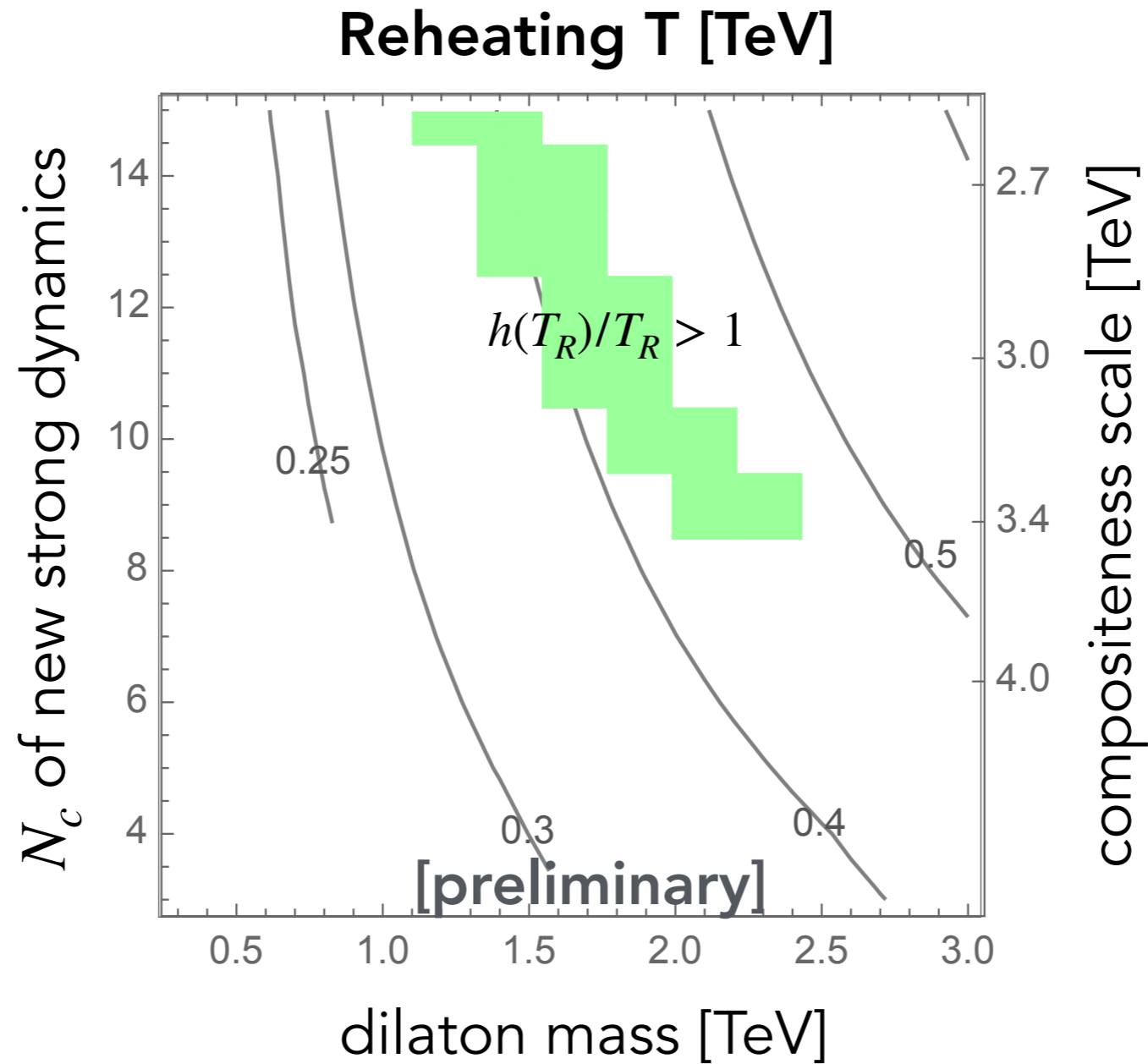


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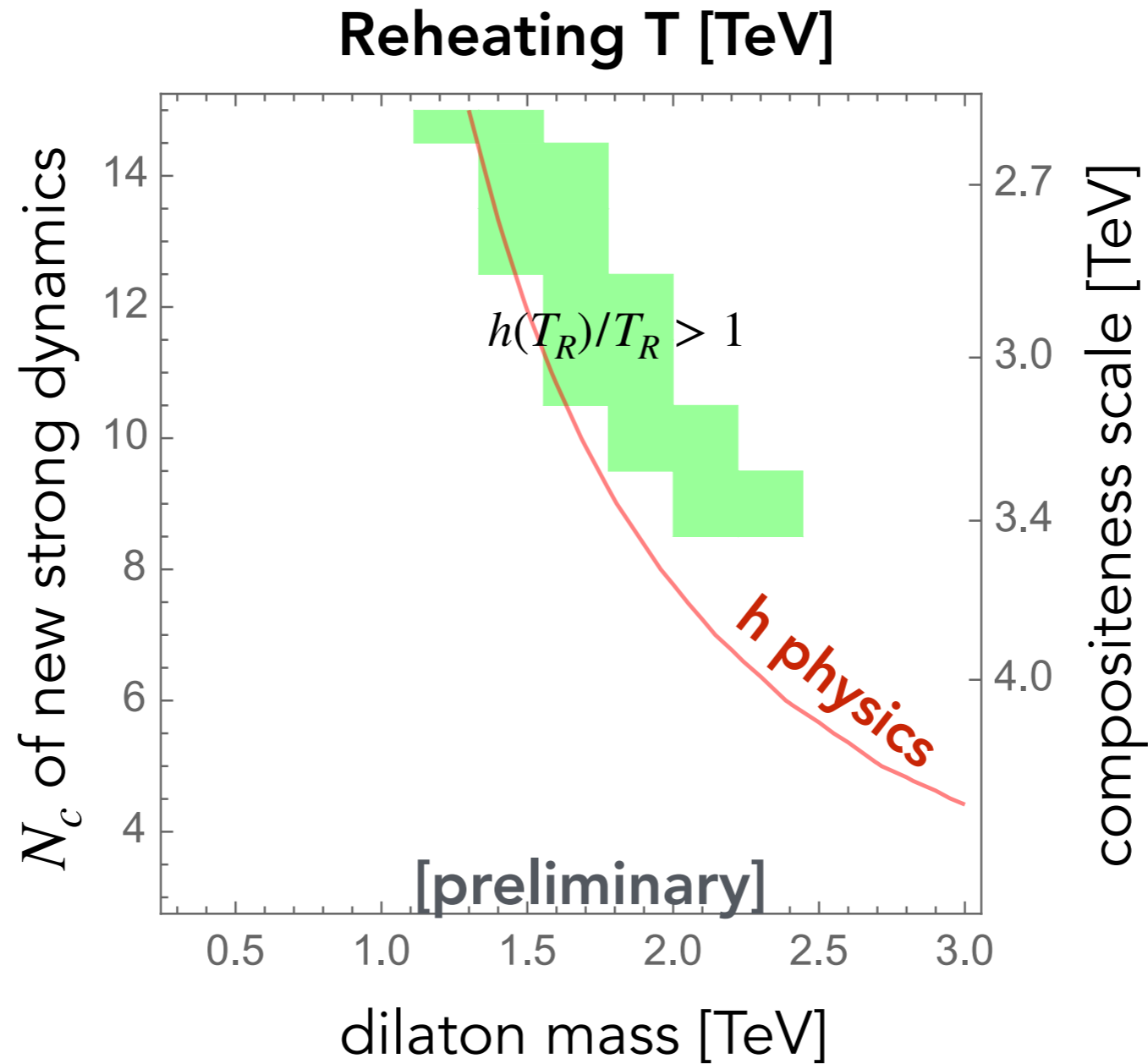


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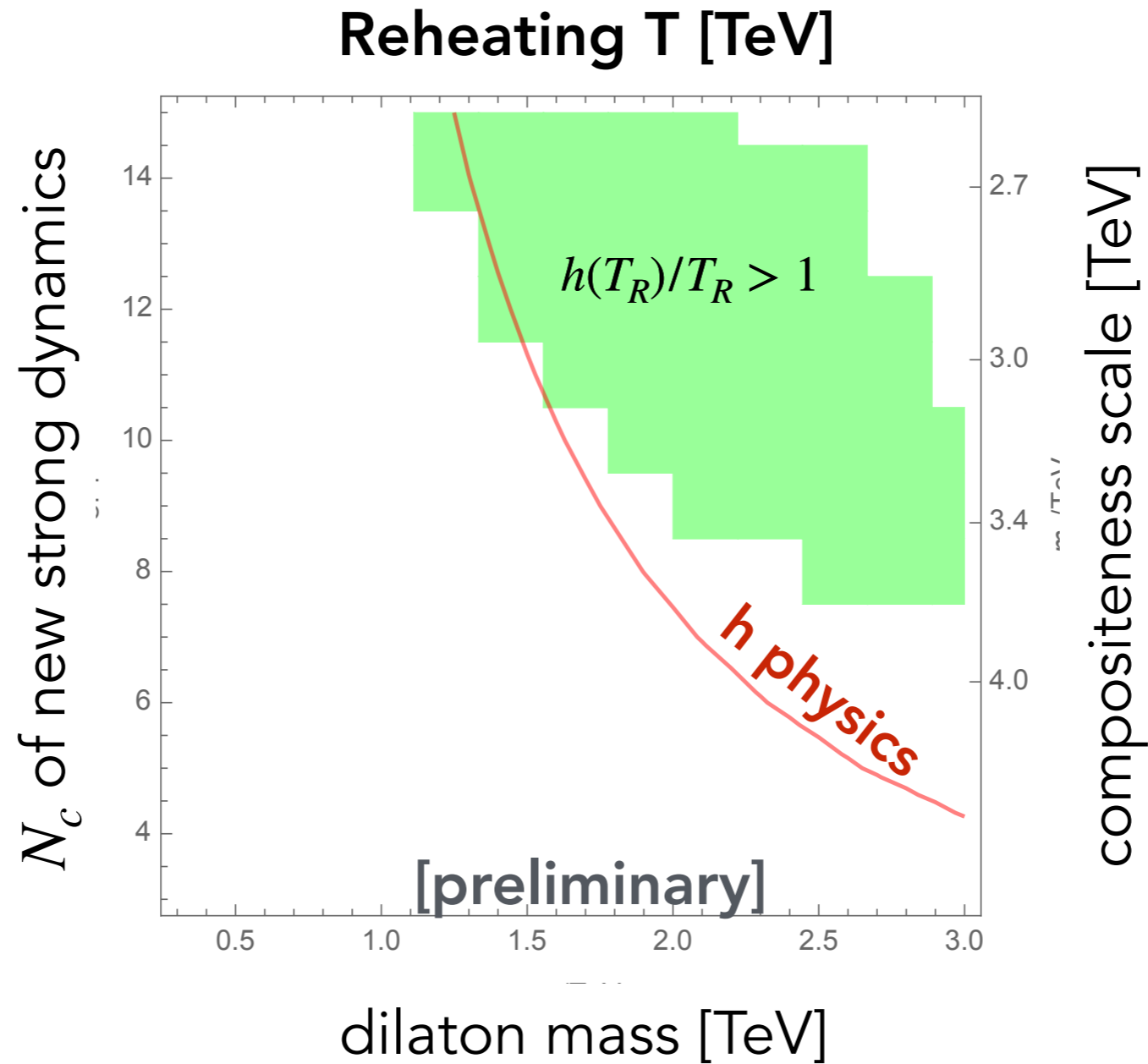


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work in progress

Bruggisser, VonHarling, OM, Servant

+ 12 SNR  
fermions





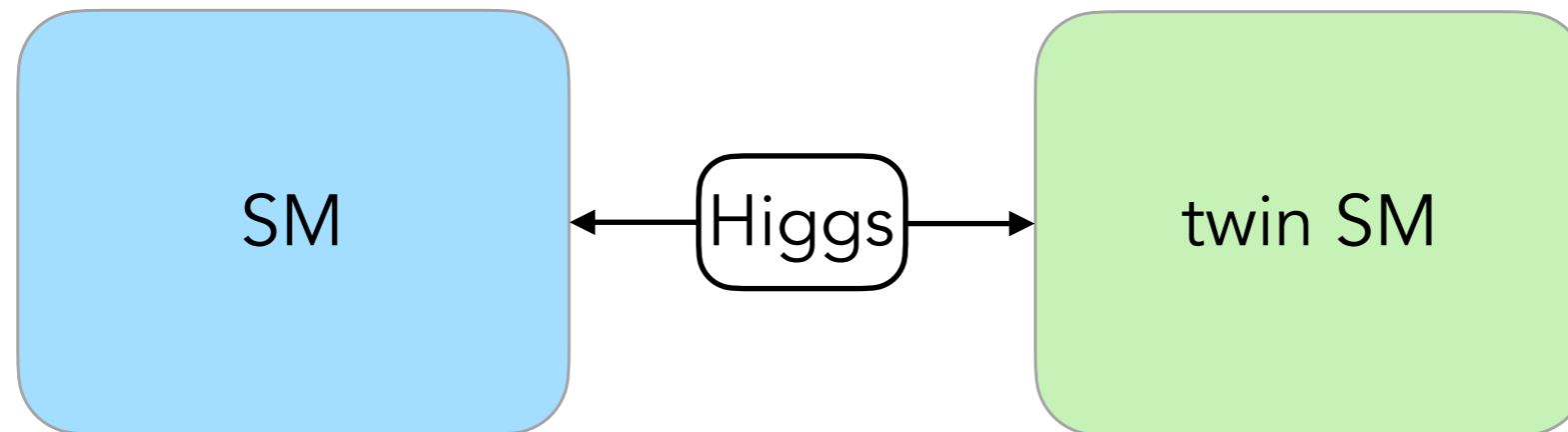
# "Automatic" Realizations

# Automatic SNR

→ Are there scenarios where the new SNR states appear automatically?

First suspect - Twin Higgs models ([Chacko et al, hep-ph/0506256](#))

Constructed to **solve the Higgs mass naturalness** problem.



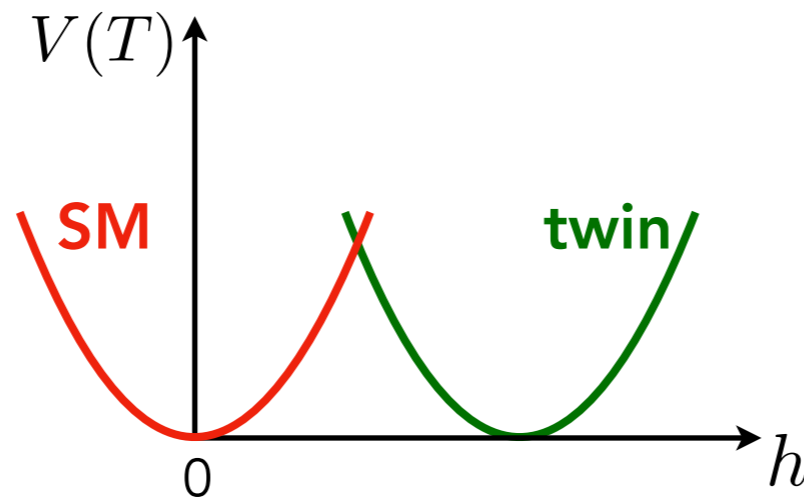
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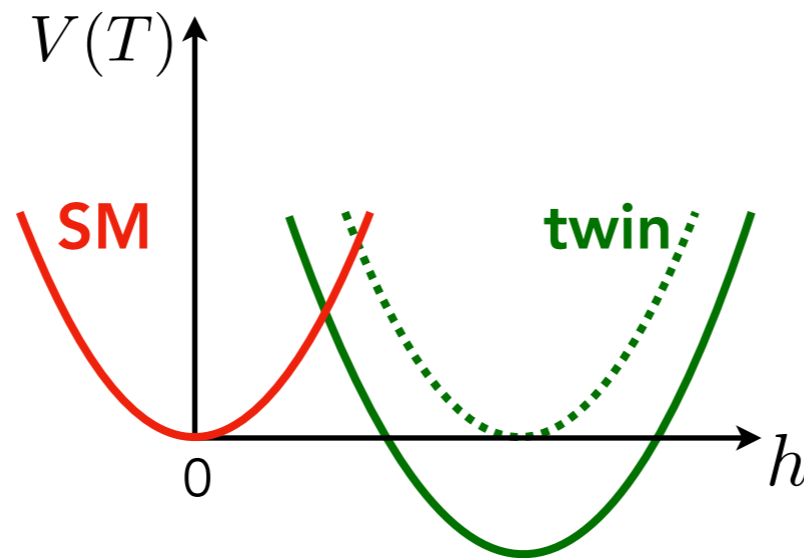
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The (anyway necessary) Twin symmetry breaking by Yukawas of light quarks and/or leptons can tip the balance to the non-restoration.

(OM, 2008.13725)

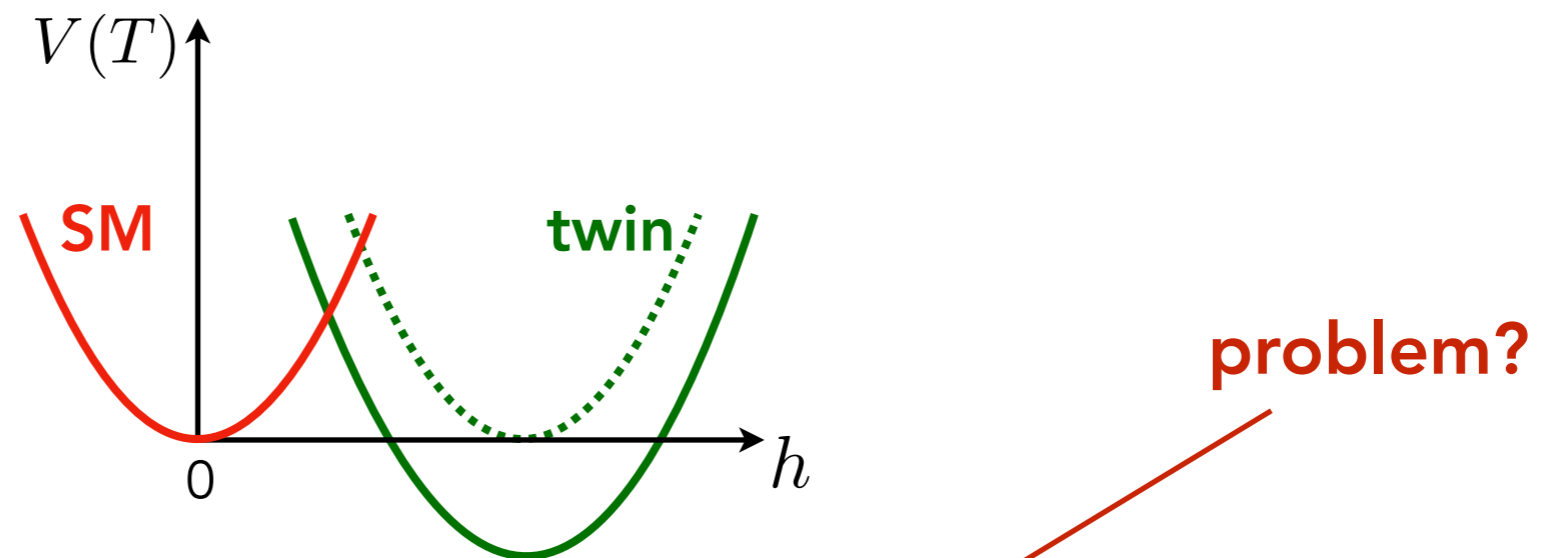
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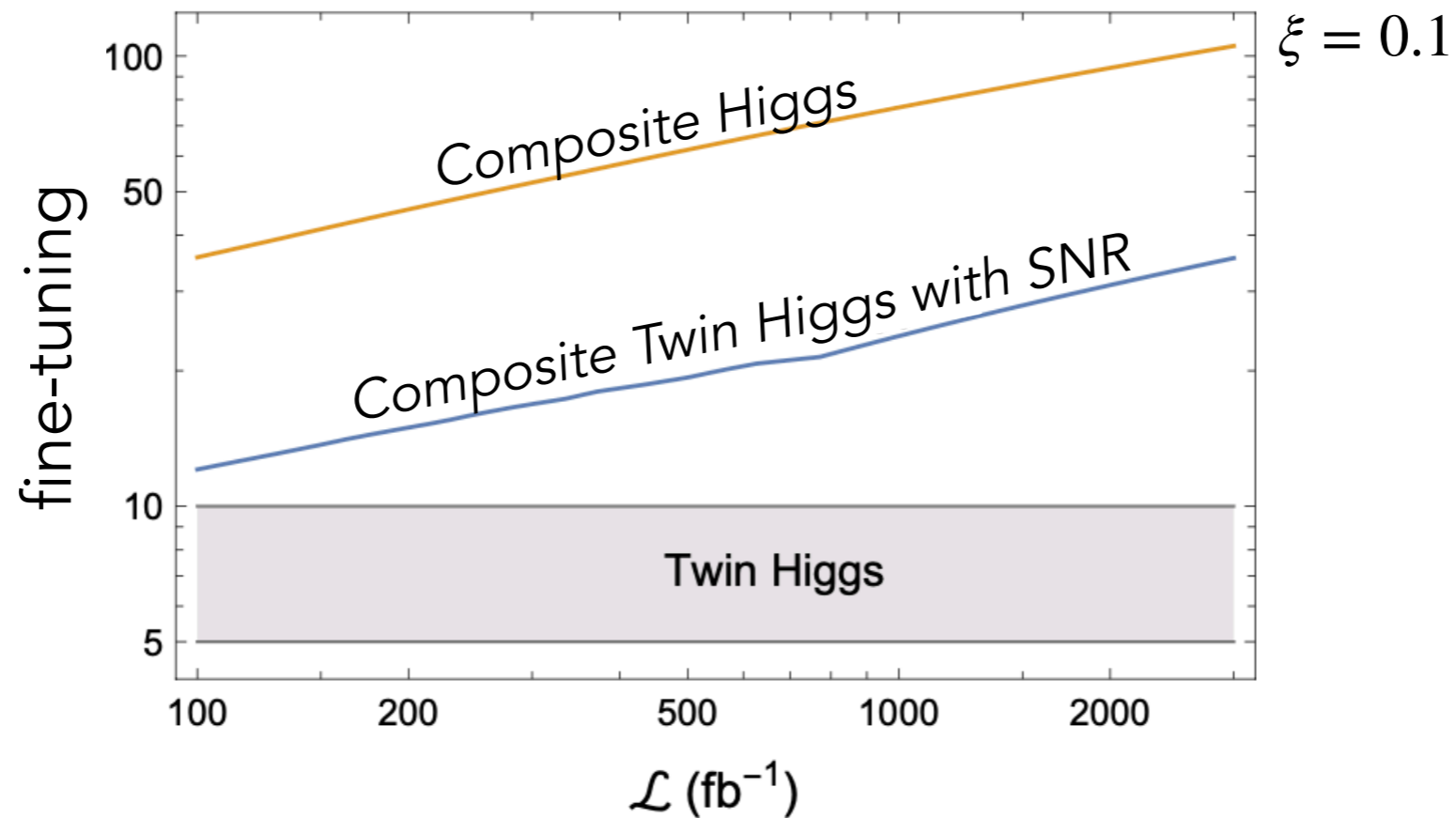


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# Automatic SNR

→ Higgs potential **fine-tuning** in Composite Twin Higgs scenario with SNR



LHC luminosity projection

# Summary

- The typical energy scales associated with EW baryogenesis can change substantially in the presence of new physics triggering high-T EW symmetry non-restoration.
- From the collider physics point of view, this opens up a broad parameter space to search for EWBG traces.
- Perturbative models with EW SNR require a sizeable number of new states. One example of new physics scenarios which provides them “naturally” is the Twin Higgs models.

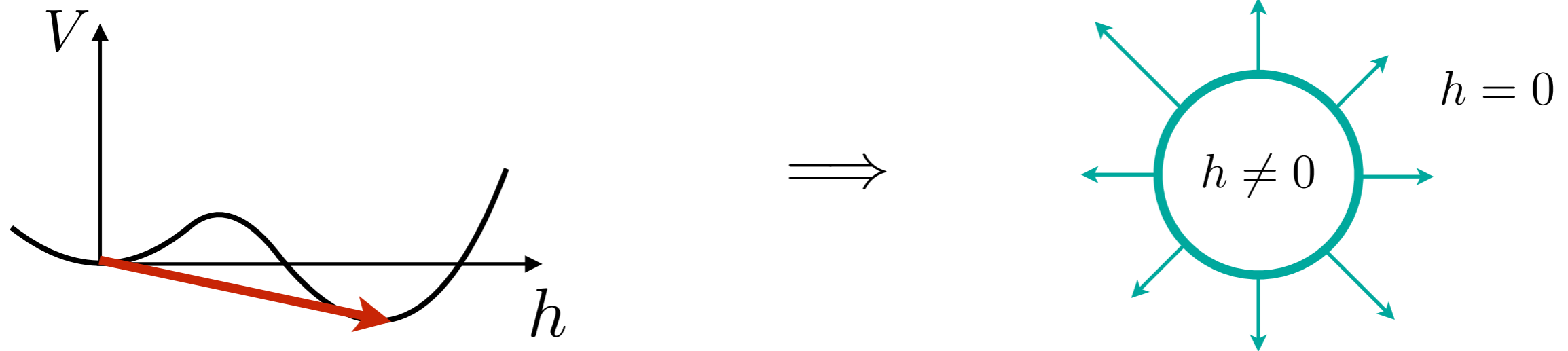
Thank you!



# Backup slides

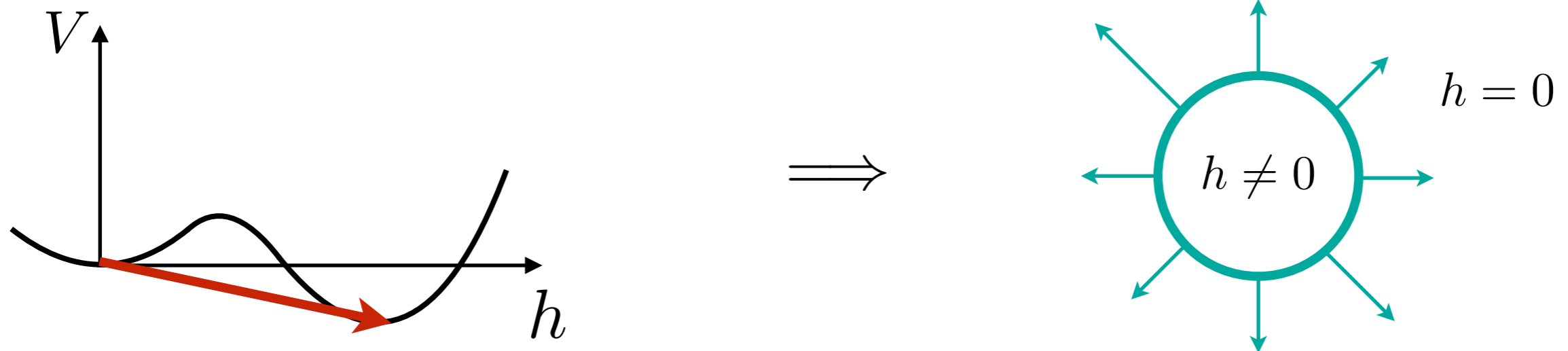
# Electroweak Baryogenesis

First order EW phase transition proceeds through bubble nucleation:

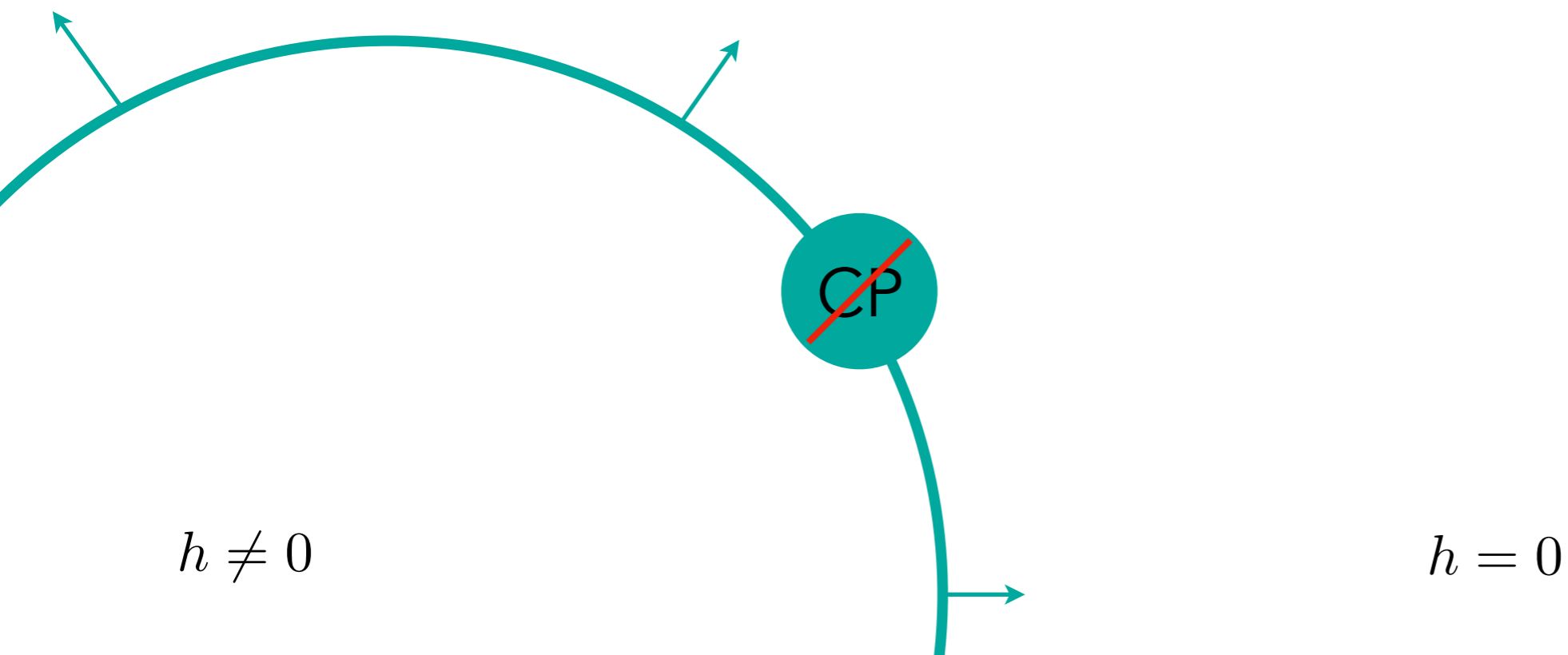


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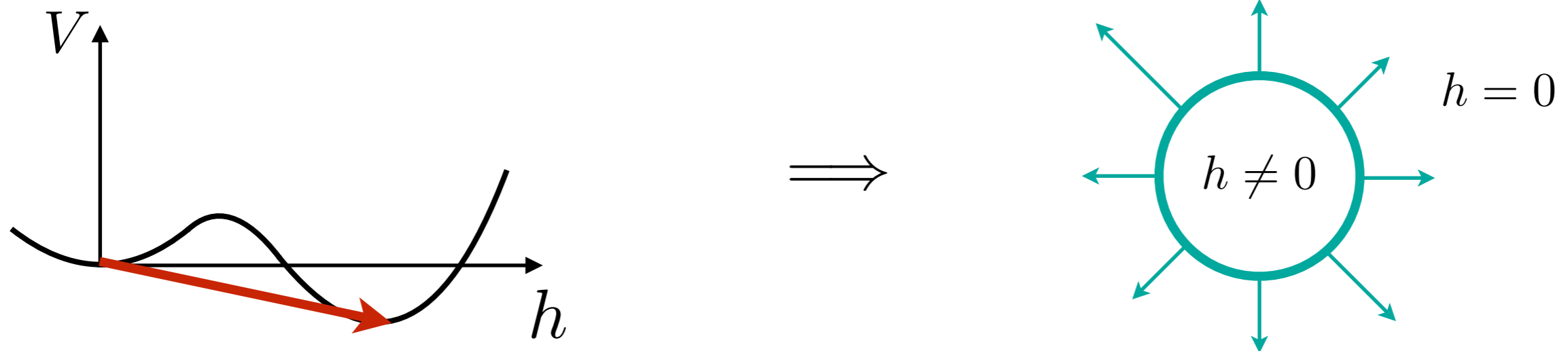


Baryon asymmetry is created close to bubble walls:

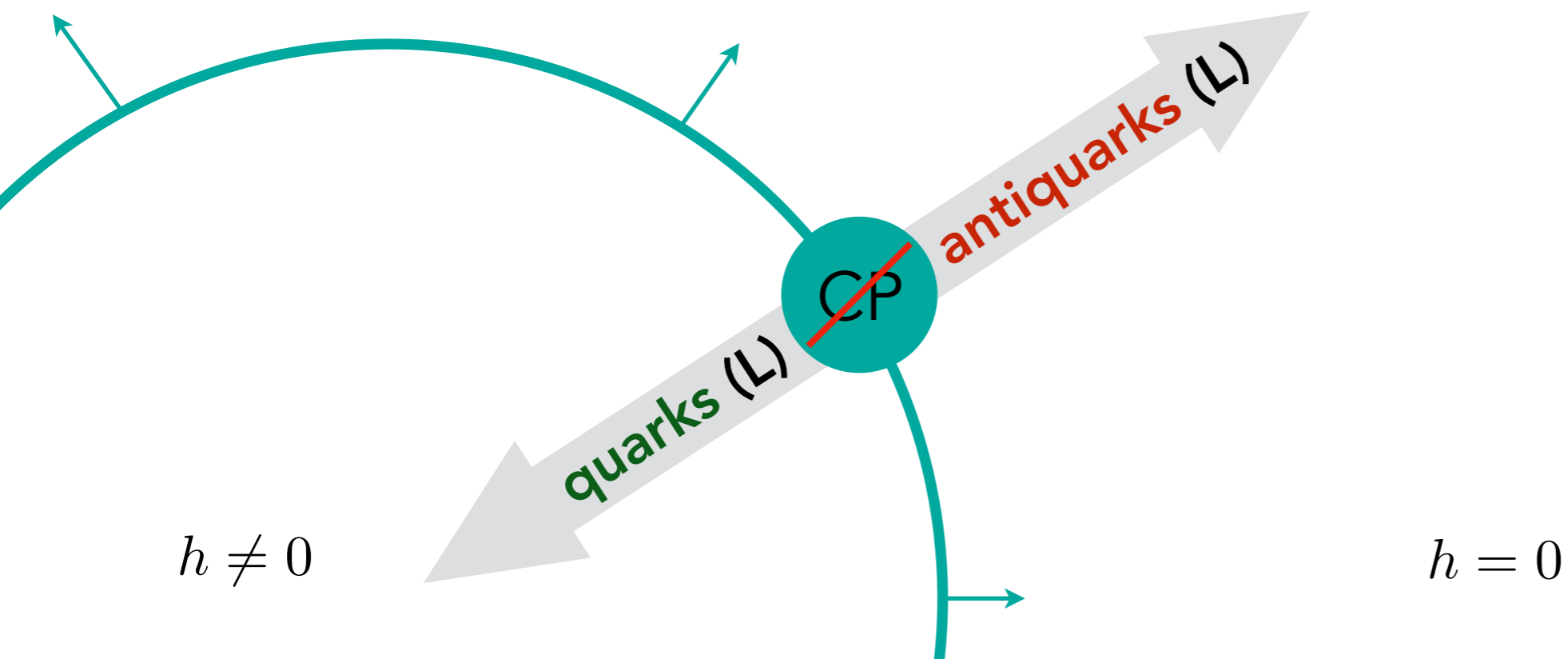


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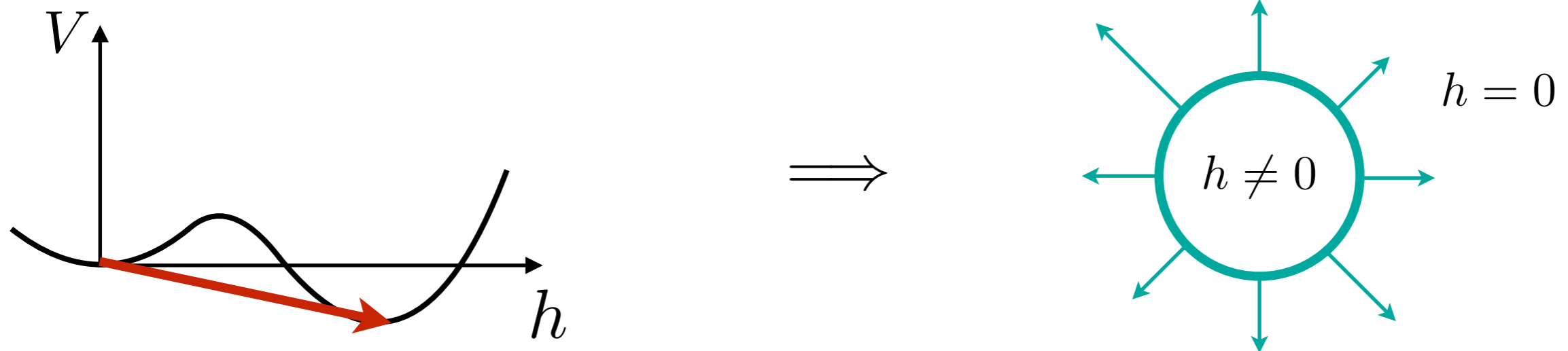


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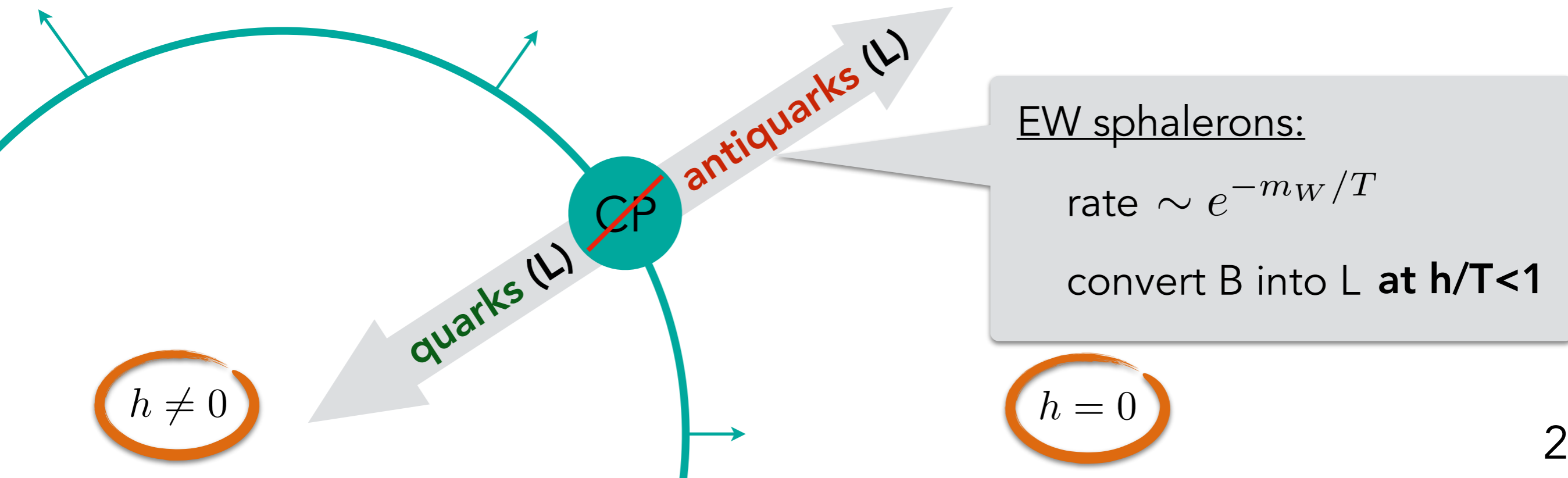


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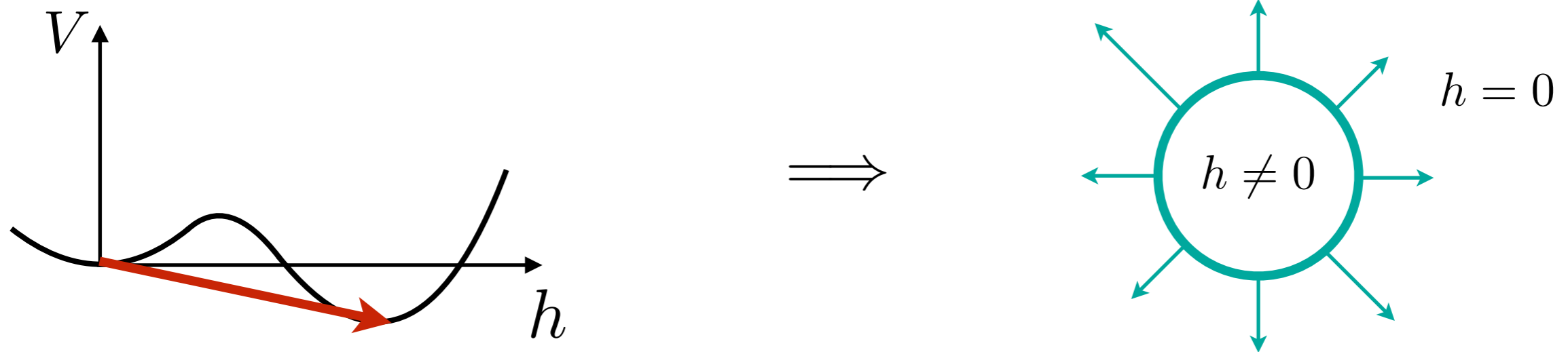


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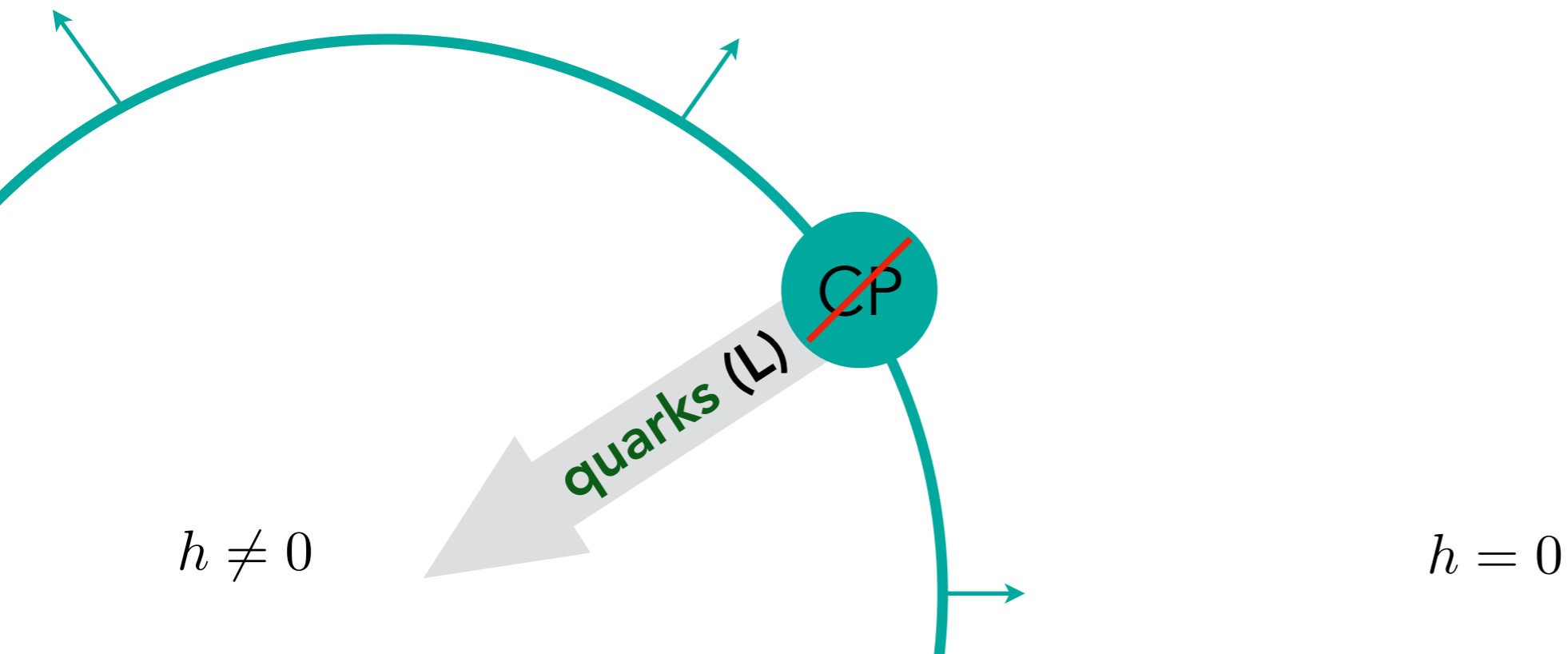


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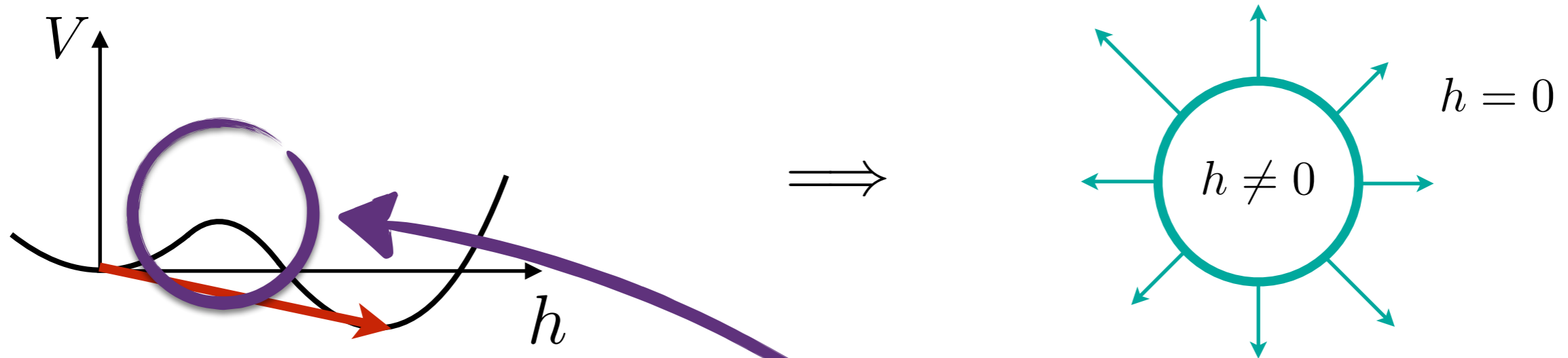


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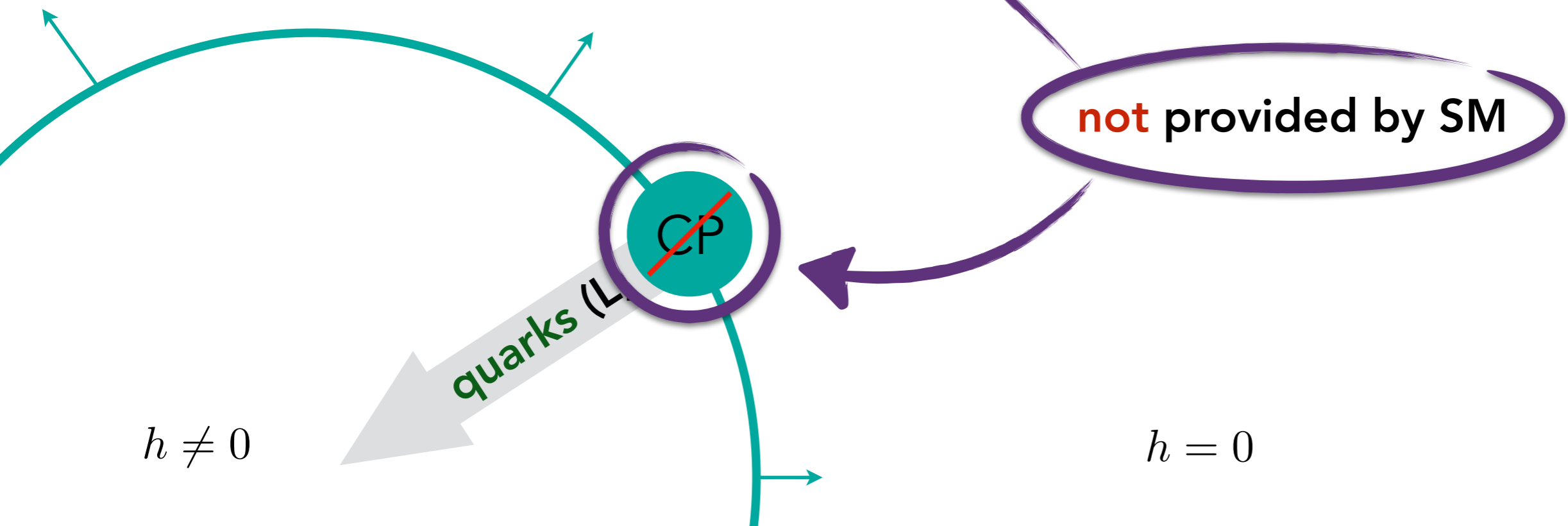


# Electroweak Baryogenesis

First order EW phase transition proceeds through bubble nucleation:



Baryon asymmetry is created close to bubble walls:



# SNR with fermions

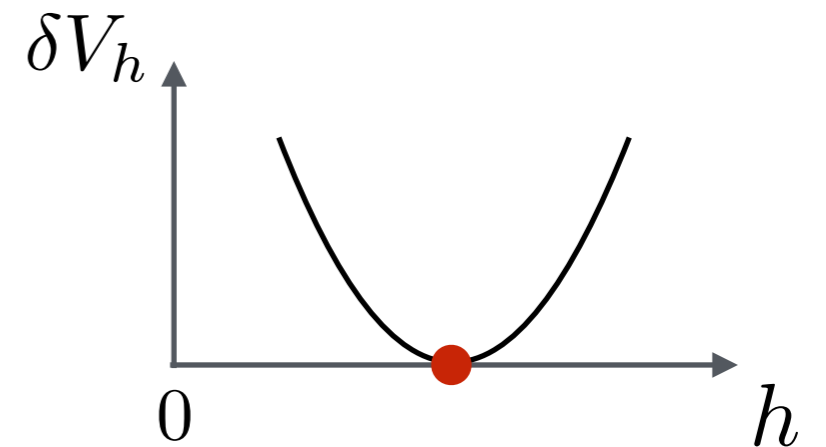
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$$m_N(h) = m_N^{(0)} - \lambda_N h^2 / \Lambda = 0$$

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- can't be done with renormalizable interactions:

$$\delta V_h \propto \sum_i m_i^2 = \text{Tr}[\mathcal{M}^\dagger \mathcal{M}] = \sum_{i,j} |\mathcal{M}_{ij}|^2$$

- in **renorm.** theories  $|\mathcal{M}_{ij}|$  are  $h$ -independent or linear in  $h$

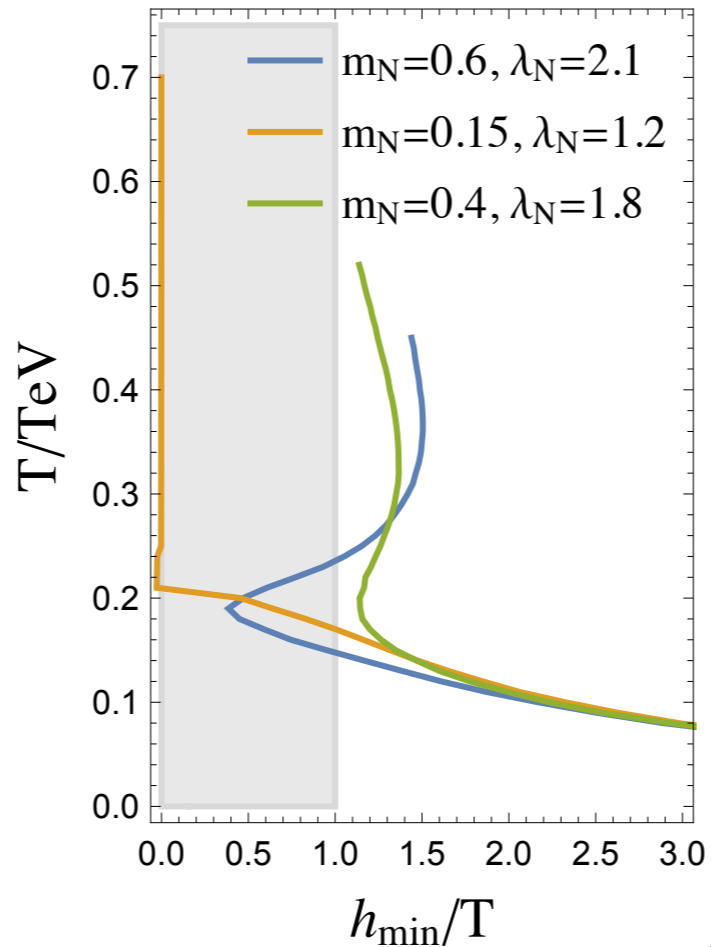
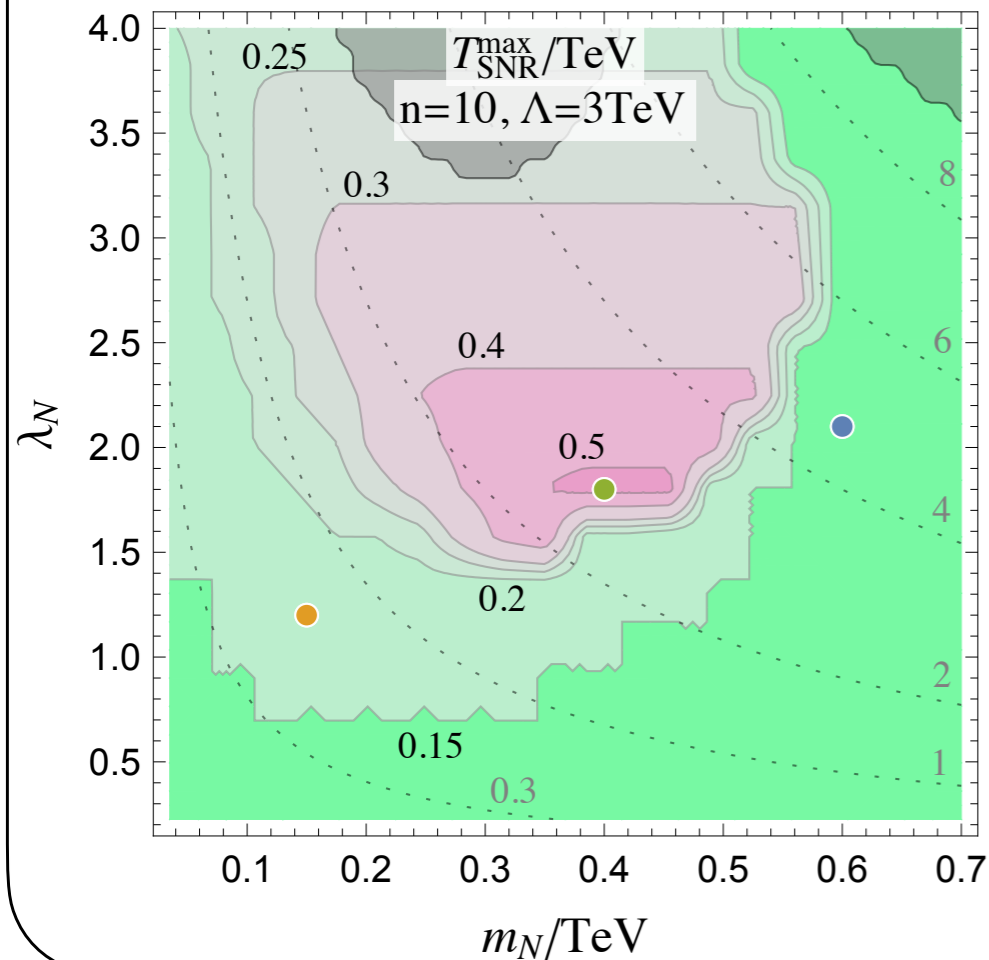
$$\sum_{i,j} |\mathcal{M}_{ij}|^2 = \text{const}_1^2 + \text{const}_2^2 h^2 \quad \Longrightarrow \quad \text{min at } h=0$$



# SNR with fermions

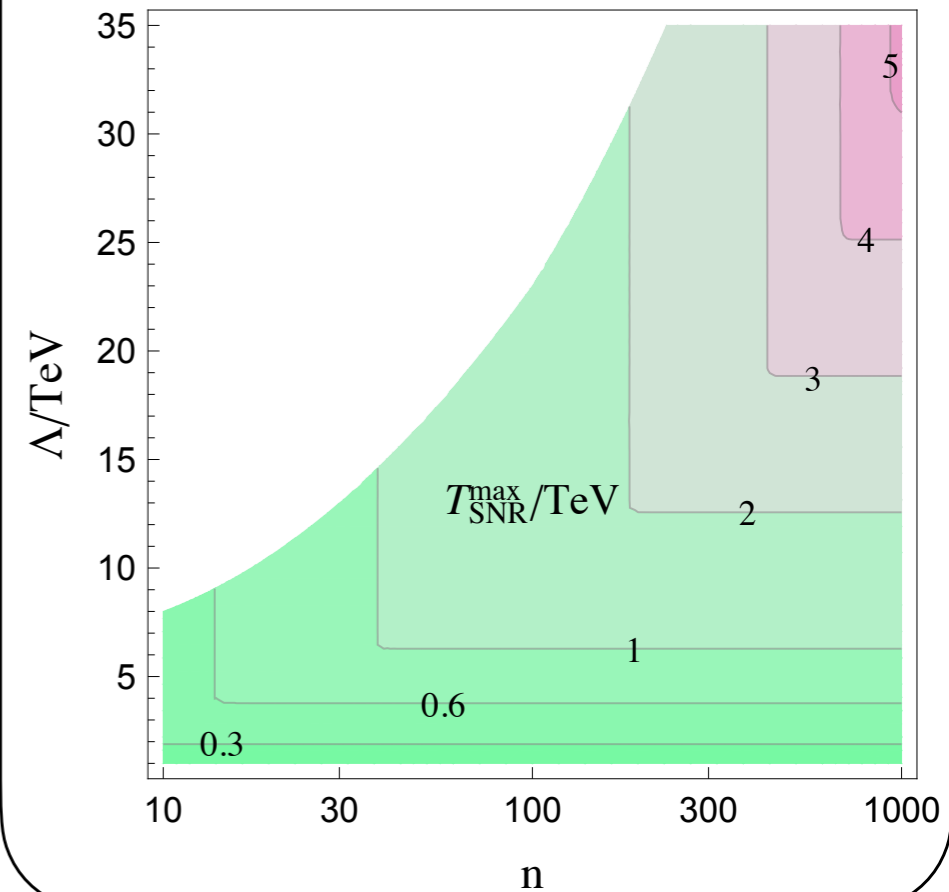
■ minimal and "natural" SM deformation:

$$n = 10, \Lambda = 3\text{TeV}$$



■ going to higher temperatures

$$T_{\text{SNR}}^{\text{max}} \sim \sqrt{n} m_N$$

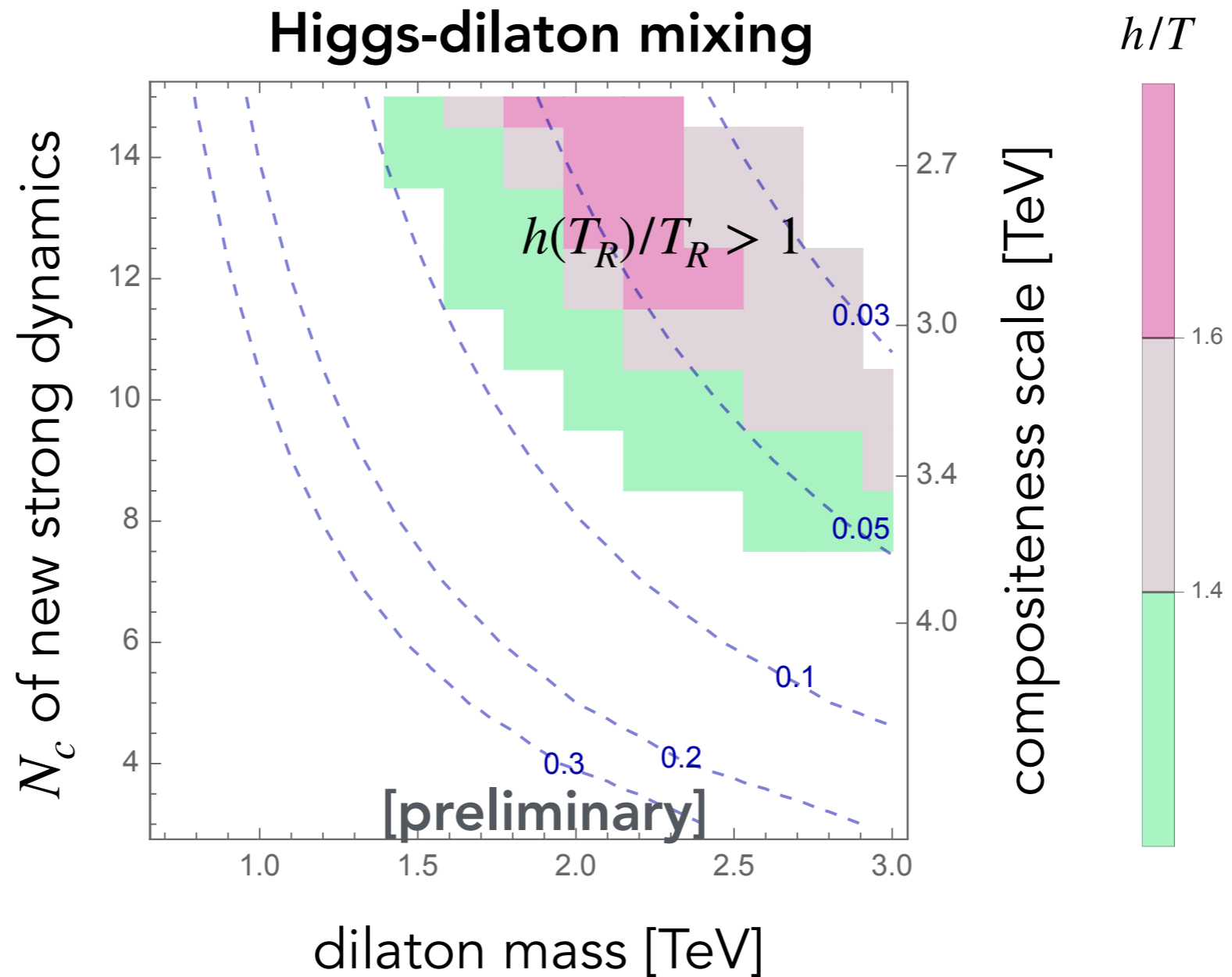


# Concrete Example: EWBG in Composite Higgs

work in progress

Bruggisser, VonHarling, OM, Servant

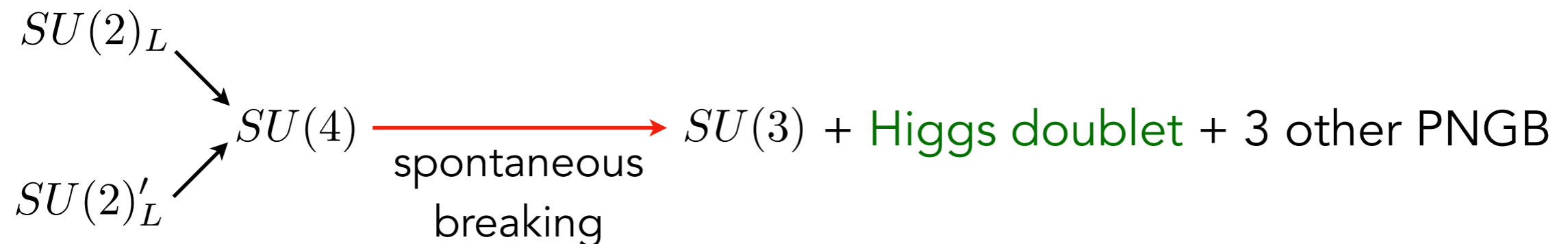
+ 12 SNR  
fermions



# SNR: Twin Higgs

- proposed in [Chacko,Goh,Harnic, hep-ph/0506256](#)  
provides a **solution to the Higgs mass naturalness problem** with **no QCD-charged new physics** up to  $\sim 10$  TeV
- Assumes a twin sector - an approximate  $Z_2$  copy of SM
- Higgs is a pseudo-Nambu-Goldstone boson of a large symmetry, which accommodates both SM and Twin  $SU(2)_L$

*example:*



# SNR: Twin Higgs

→ SM states couplings to the Higgs  $\propto \sin h/f$

Twin states couplings to the Higgs  $\propto \cos h/f$

→ Quadratic divergences  $\propto \Lambda^2$  in the Higgs mass cancel

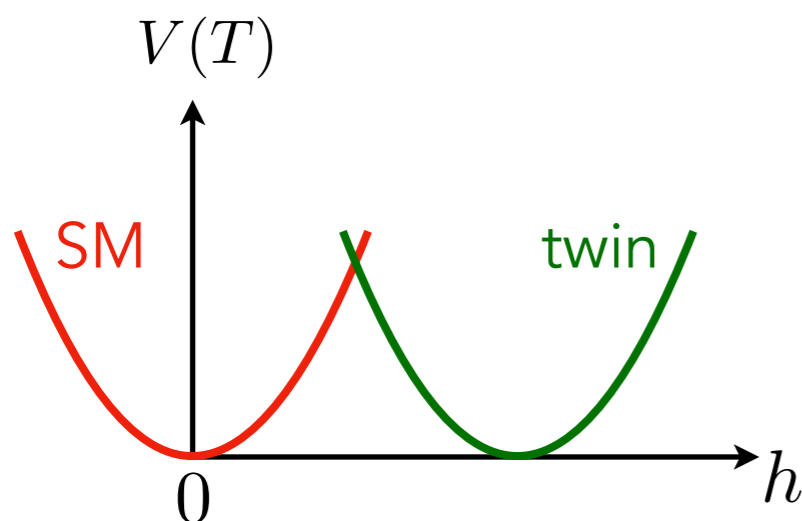
$$V \sim f^2 \Lambda^2 (\sin^2 h/f + \cos^2 h/f) = f^2 \Lambda^2$$

no Higgs dependence  
no contribution to the Higgs mass  $\propto \Lambda^2$

SM contribution

Twin contribution

→ The EW symmetry-restoring high-T SM correction is also canceled (balanced)



The resulting position of the minimum depends on weak  $Z_2$  breaking

→ analyze which type of  $Z_2$  can tilt high-T potential to SNR minimum

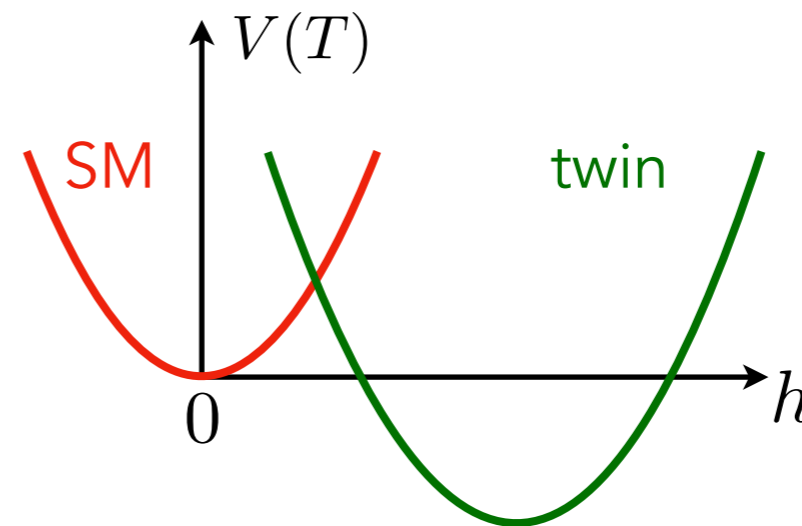
# SNR: Twin Higgs

→ Sources of  ~~$Z_2$~~

necessarily broken in the light fermion sector: eg twin neutrinos cannot be light.

simplest realization: larger Yukawas  $\tilde{\lambda}_q$  for light twin fermions

$$\tilde{\lambda}_q f \bar{q} q \cos h/f$$

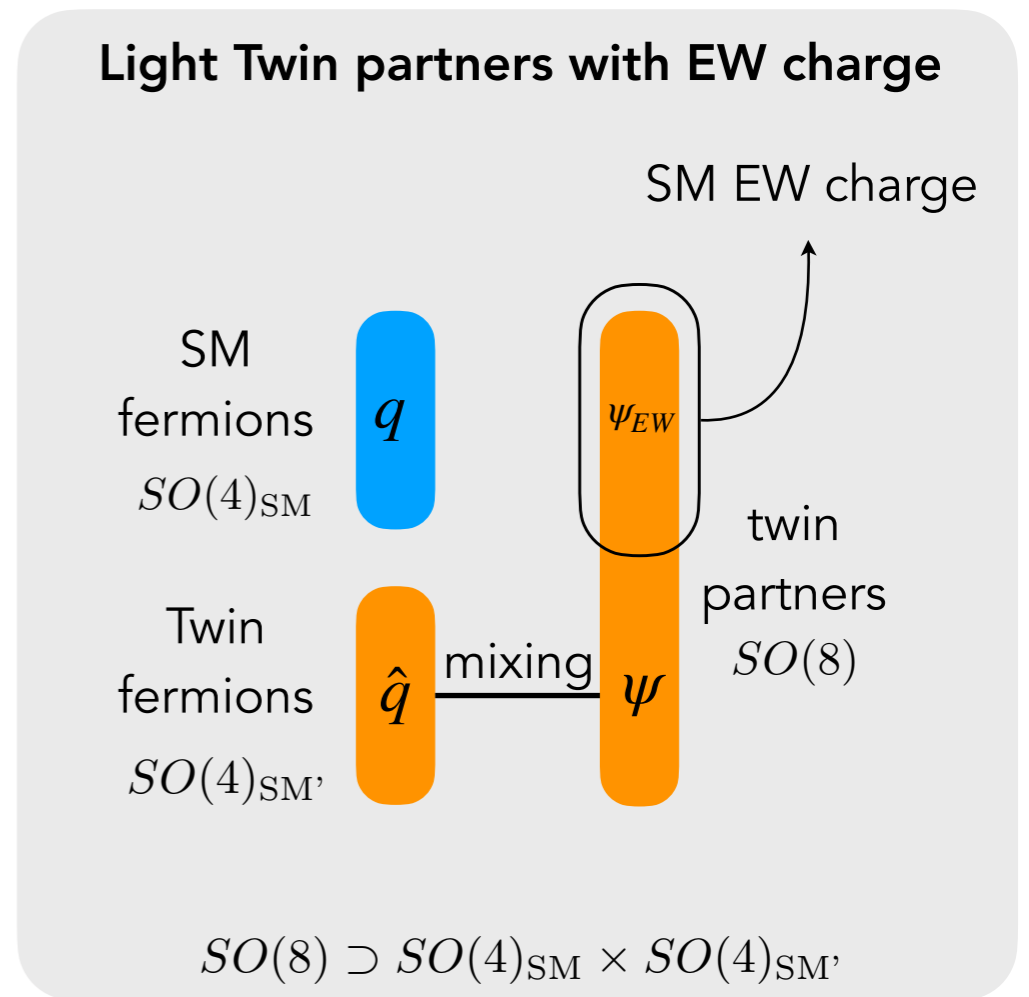
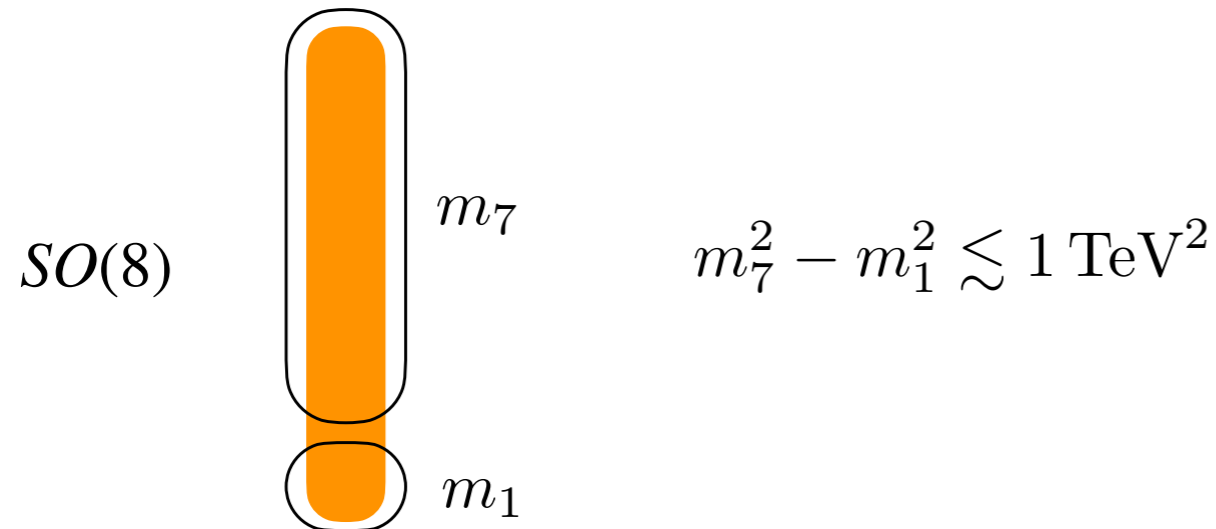


→ This also spoils the cancellation of  $T=0$  quadratic divergences to the Higgs mass:  $\delta m_h^2 \sim (\tilde{\lambda}_q^2/16\pi^2)\Lambda^2$

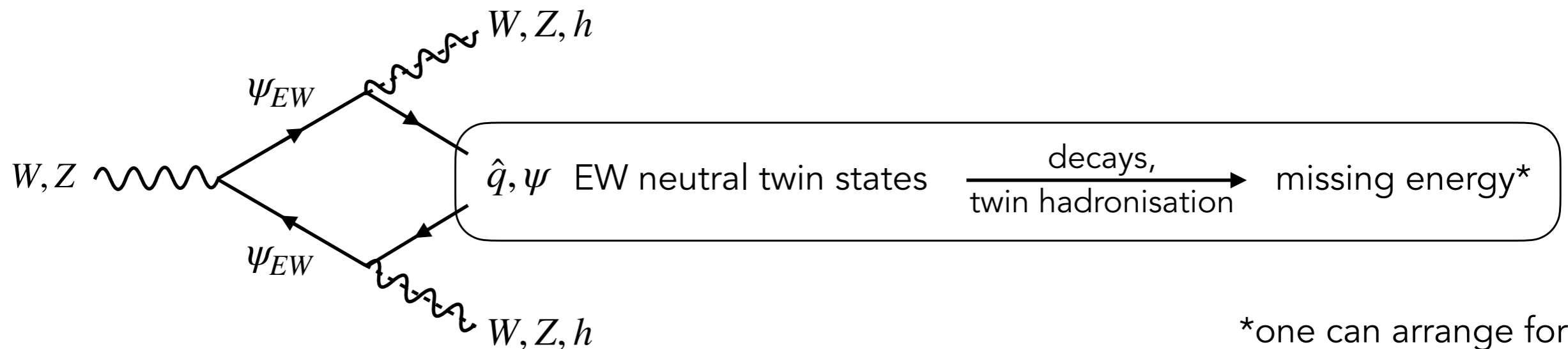
take a lower cutoff  $\Lambda$  in  
 $\Rightarrow$  the twin light quark sector  $\Rightarrow$  light **twin partners**  
(no SM QCD charge)

# SNR: Twin Higgs

→ To not spoil the Higgs mass, we need:



→ Collider signal:



\*one can arrange for visible signatures too

# SNR: Twin Higgs

→ numerical results

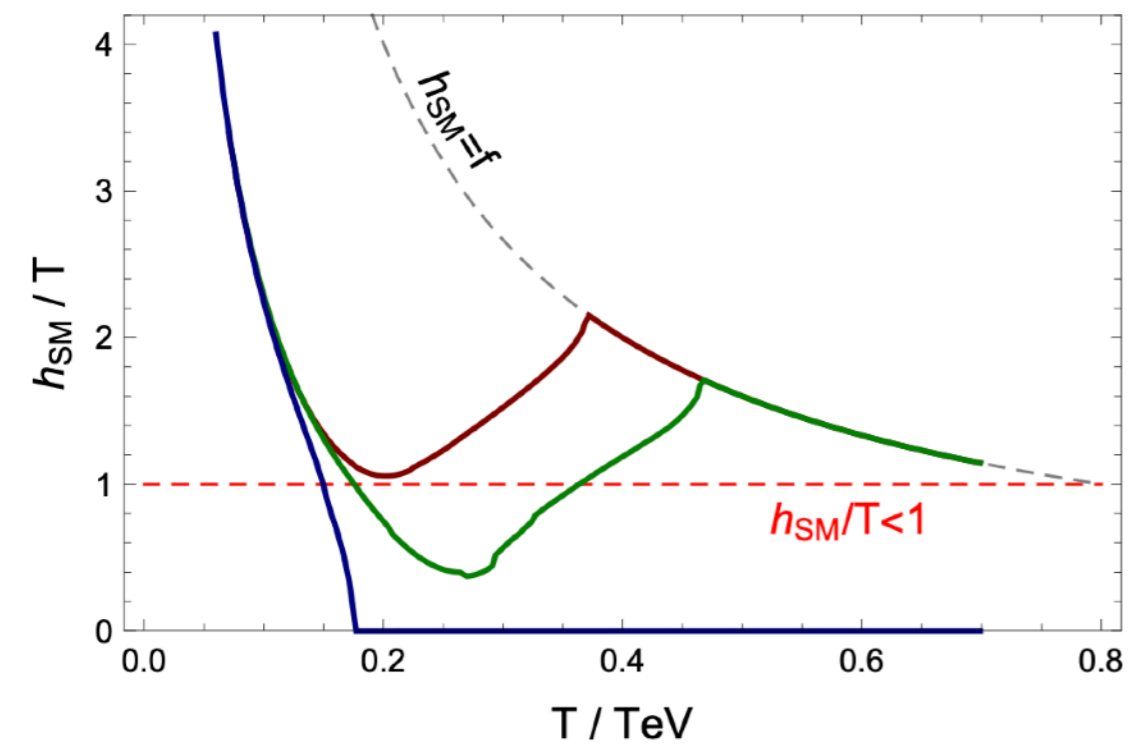
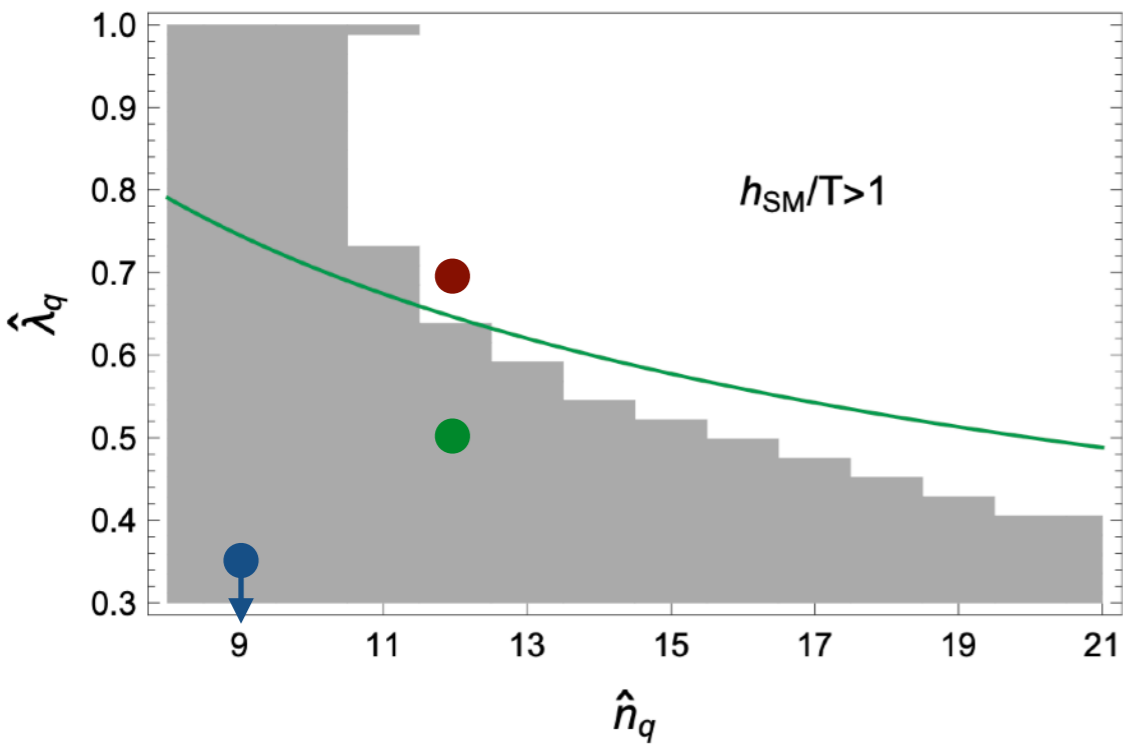


Figure 2: Solid lines show the evolution of  $h_{SM}/T$  in the minimum of the Higgs potential depending on the temperature, for three choices of parameters:  $\hat{\lambda}_q = 0.2, \hat{\nu}_q = 9$  (blue),  $\hat{\lambda}_q = 0.5, \hat{\nu}_q = 12$  (green), and  $\hat{\lambda}_q = 0.7, \hat{\nu}_q = 12$  (red).

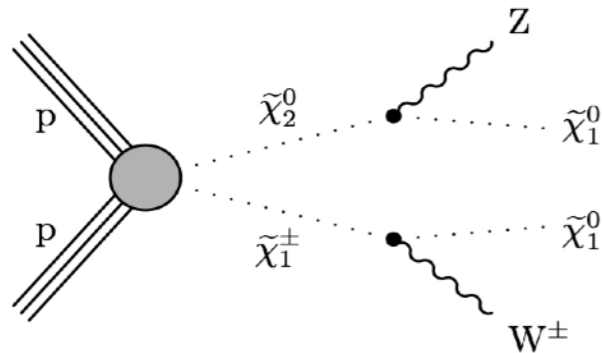
# Automatic SNR

→  $Z_2$  breaking leads to additional contributions to the Higgs mass

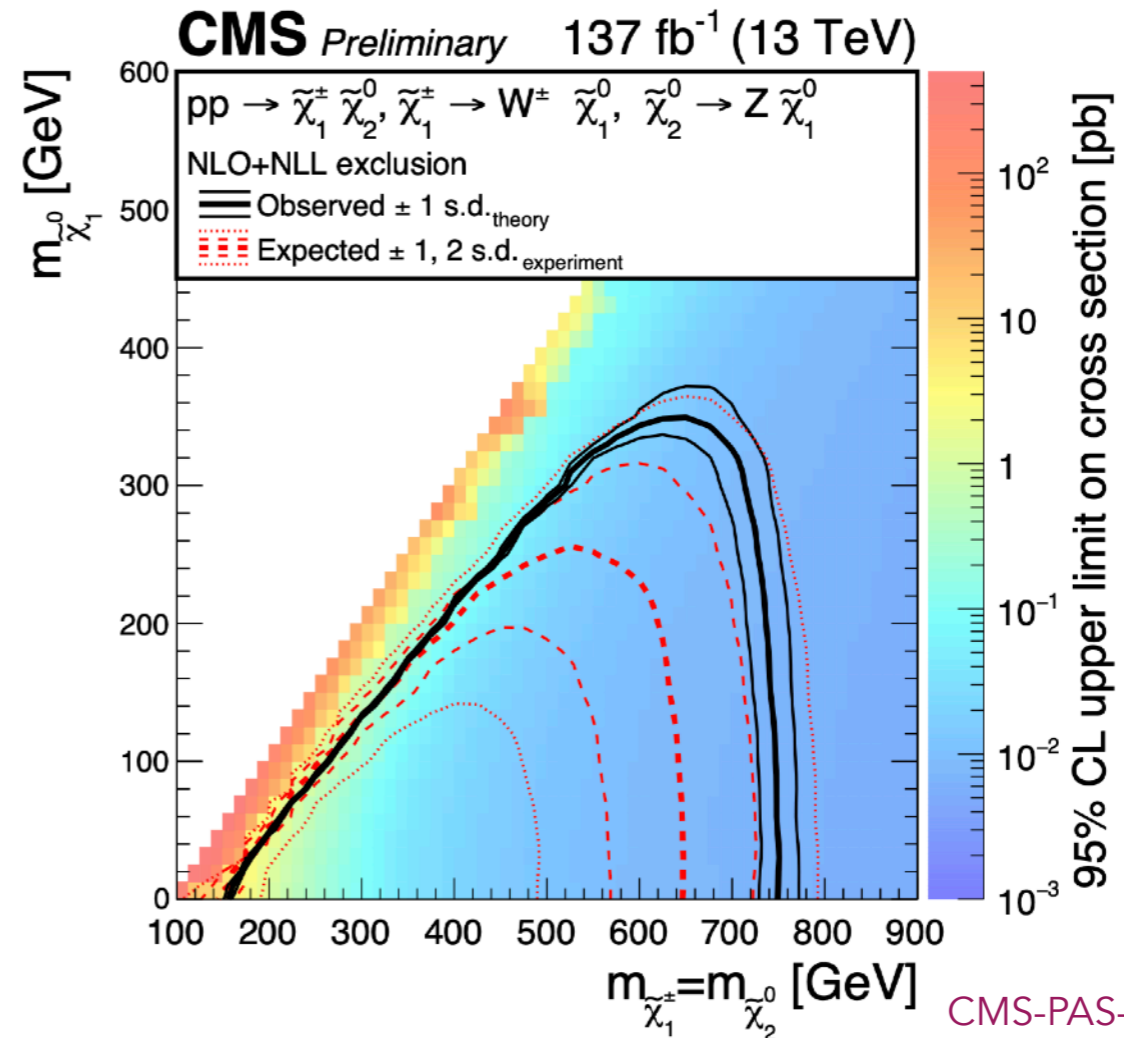
$$\delta m_h^2 \propto n_q \lambda_q^2 \Lambda^2 \left[ \Lambda^2 \right]$$

← mass of EW-charged twin partners

→ chargino/neutralino-like signal



for twin hadron decays see: Craig, Katz, Strassler, Sundrum  
1501.05310



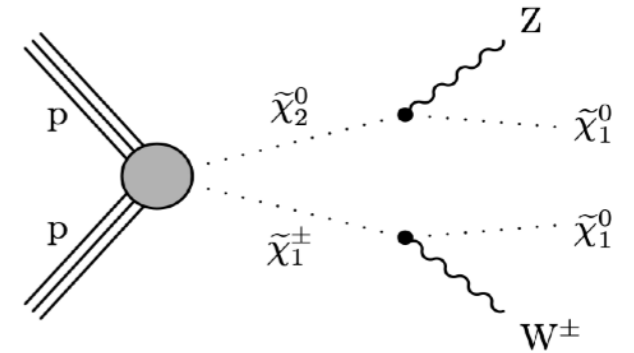


# SNR: Twin Higgs

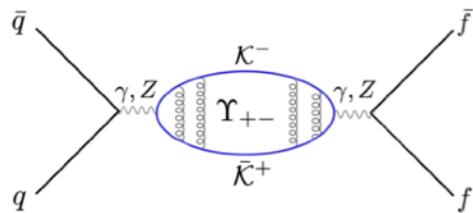
→ experimental bounds on EW-charged twin partners

chargino/neutralino pair production with a decay to a W,Z,h and neutralino (missing Et)

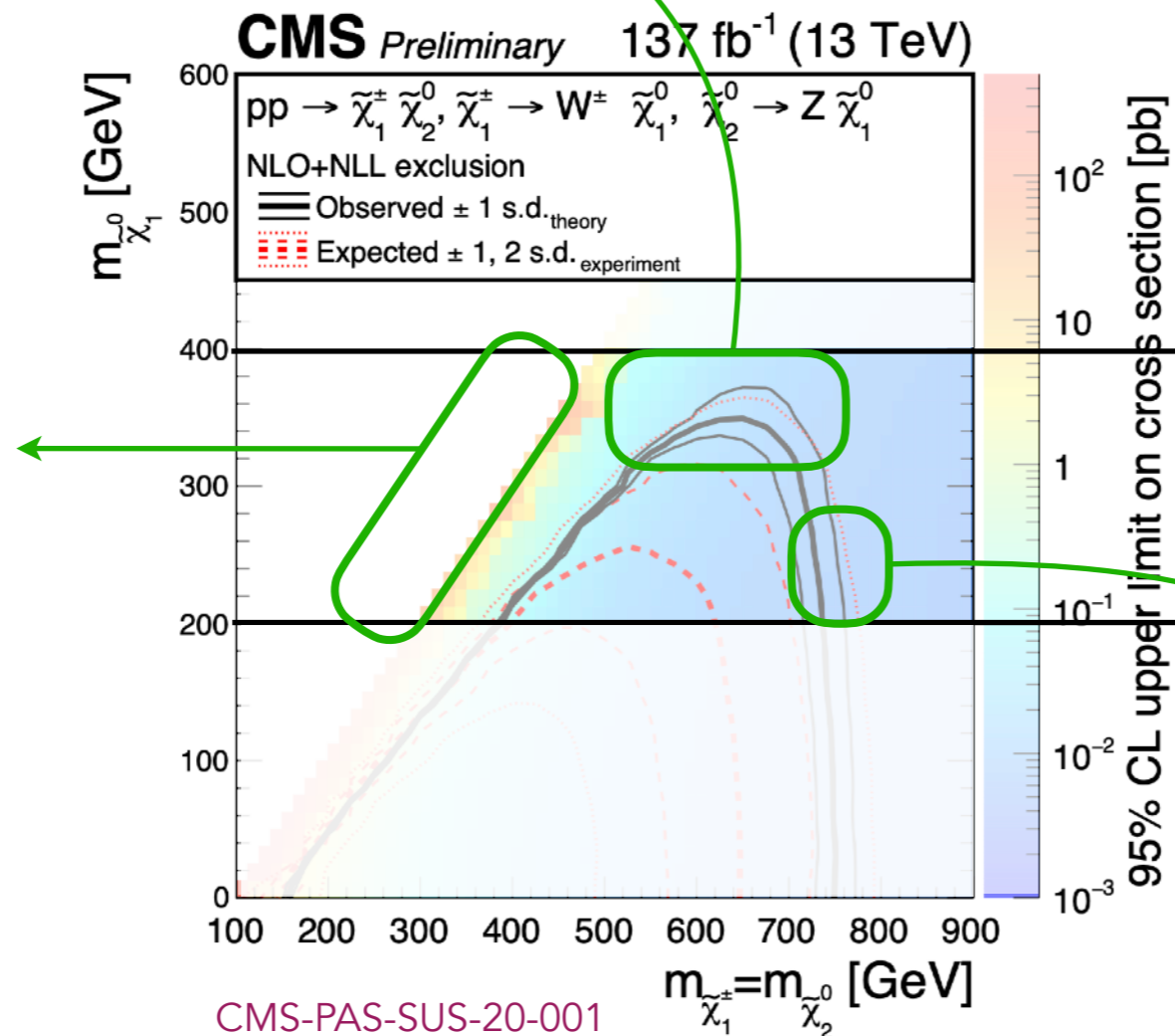
bounds significantly depend on kinematical distributions



if charged under twin QCD, partners hadronise, and decay back into SM



Cheng, Salvioni, Tsai  
1612.03176



CMS-PAS-SUS-20-001

optimal mass of twin fermions  $\hat{q}$

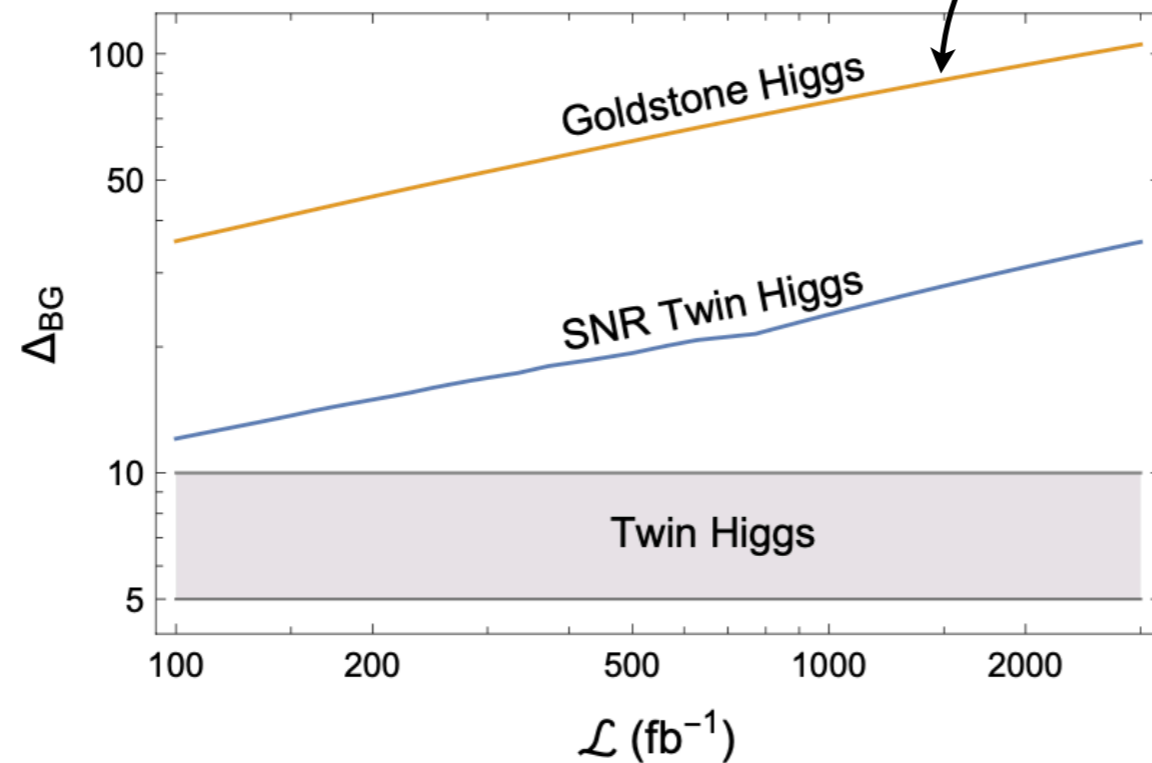
eventually used for the derivation of bounds

# SNR: Twin Higgs

→ fine-tuning

$$\Delta_{\text{BG}} = \frac{\hat{n}_q y_L^2 m_7^2 \cos^2 v / f}{2\pi^2 m_h^2}$$

lower bound from a pair production of one top partner, does not include single production or pile-up from several partners



# SNR: thermal corrections

- In thermal bath of particles  $X$ , what is their effect on the Higgs potential?

bosons:

$$\Delta V_b^T = \frac{T^4}{2\pi^2} J_b[m^2/T^2],$$

$$J_b[x] = \int_0^\infty dk k^2 \log \left[ 1 - e^{-\sqrt{k^2+x}} \right],$$

fermions:

$$\Delta V_f^T = -\frac{2T^4}{\pi^2} J_f[m^2/T^2]$$

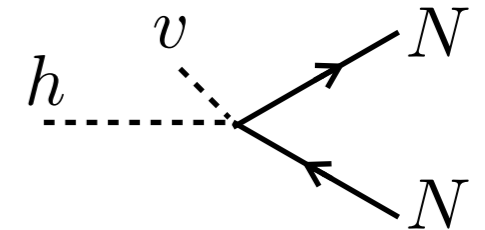
$$J_f[x] = \int_0^\infty dk k^2 \log \left[ 1 + e^{-\sqrt{k^2+x}} \right]$$

# SNR: collider pheno

→ invisible Higgs decays

$$\text{BR}_{h \rightarrow NN} \sim \frac{1}{n} \left( n \lambda_N \frac{m_N}{\Lambda} \right)^2 \frac{v_{\text{SM}}^2}{m_h \Gamma_h} \quad \text{where} \quad \Gamma_h \simeq 6 \text{ MeV}$$

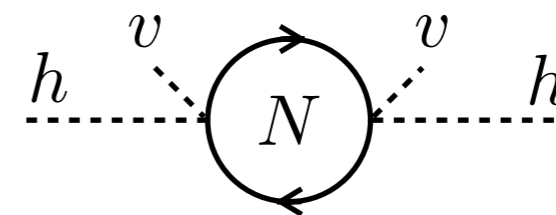
requires either  $n \gtrsim 10^6$  or  $m_N > m_h/2$



→ Universal Higgs couplings modification

from Higgs wave function renormalization

$$\delta Z_h \sim \frac{1}{(4\pi)^2} \frac{1}{n} \left( \frac{n \lambda_N m_N}{\Lambda} \right)^2 \frac{v^2}{m_N^2} \quad \longrightarrow \quad \text{future lepton colliders?}$$



→ Production via off-shell Higgs

D.Curtin et al 1409.0005

if no decay within detector, may be testable at 100TeV collider  
for low  $m_N$  and  $n$