

$\Delta^{\pm\pm}$ (or $H^{\pm\pm}$) Production at Hadron Colliders

HPNP 2021 - Osaka University

Richard Ruiz¹

Institute of Nuclear Physics – Polish Academy of Science (IFJ PAN)

25 March 2021



¹with B. Fuks and M. Nemevšek [[1912.08975](#)] + review with Y. Cai, T. Han, and T. Li [[1711.02180](#)]

what is this talk about?

how to best make doubly charged Higgs bosons ($\Delta^{\pm\pm}$) at colliders

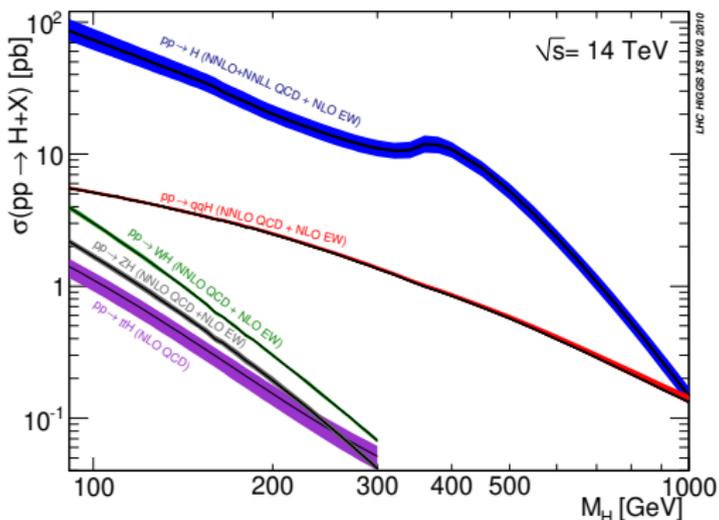
Cross section predictions not updated since publication of total NLO rates for the neutral current Drell-Yan process

Muhlleitner & Spira [[hep-ph/0305288](https://arxiv.org/abs/hep-ph/0305288)]

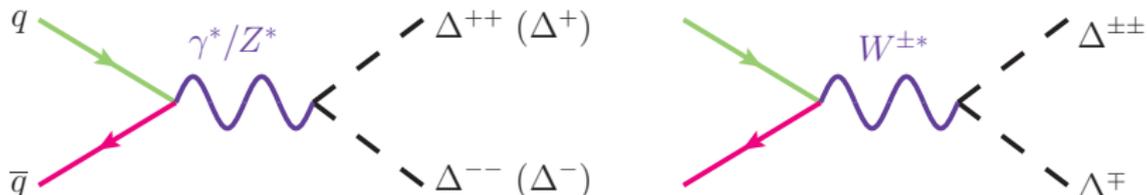
- K -factors for $m_{\Delta} < 1$ TeV at $\sqrt{s} = 14$ TeV with one PDF

Missing from the lit.: differential information at NLO, clarity (again) on size of $\gamma\gamma$ fusion, and systematic comparison of production channels

- Where is the $\Delta^{\pm\pm}$ equivalent of this????



why? because $\Delta^{\pm\pm}$, Δ^{\pm} are important²



²Konetschny and Kummer ('77); Schechter and Valle ('80); Cheng and Li ('80); Lazarides, et al ('81); Mohapatra and Senjanovic ('81)

Example: Type II Seesaw Mechanism for Neutrino Masses

Hypothesize a **scalar** $SU(2)_L$ triplet with **lepton number** $L = -2$

$$\hat{\Delta} = \frac{1}{\sqrt{2}} \begin{pmatrix} \Delta^+ & \sqrt{2}\Delta^{++} \\ \sqrt{2}\Delta^0 & -\Delta^+ \end{pmatrix}, \quad \text{with} \quad \mathcal{L}_{\Delta\Phi} \ni \mu_{h\Delta} \left(\Phi^\dagger \hat{\Delta} \cdot \Phi^\dagger + \text{H.c.} \right)$$

The mass scale $\mu_{h\Delta}$ **breaks lepton number**, and induces $\langle \hat{\Delta} \rangle \neq 0$:

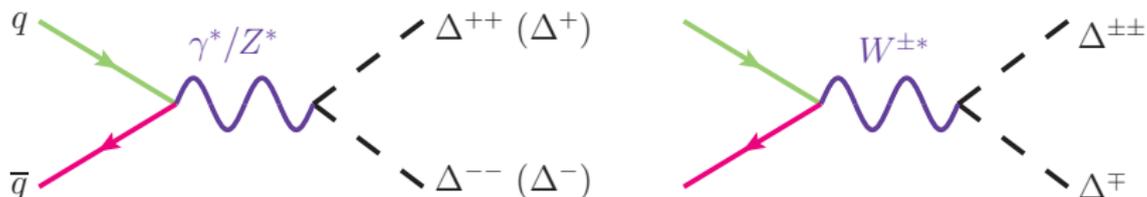
$$\sqrt{2}\langle \hat{\Delta} \rangle = v_\Delta \approx \frac{\mu_{h\Delta} v_{EW}^2}{\sqrt{2}m_\Delta^2}$$

which leads to **left-handed Majorana masses** for neutrinos

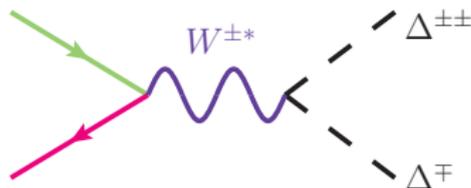
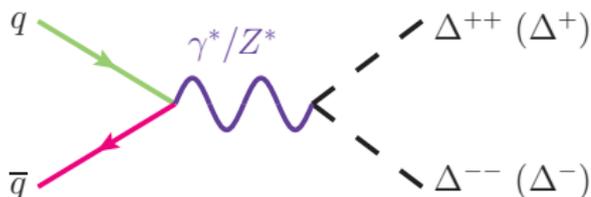
$$\begin{aligned} \Delta\mathcal{L} &= -\frac{y_\Delta^{ij}}{\sqrt{2}} \overline{L^c} \hat{\Delta} L = -\frac{y_\Delta^{ij}}{\sqrt{2}} \begin{pmatrix} \overline{\nu^{jc}} & \overline{\ell^{jc}} \end{pmatrix} \begin{pmatrix} 0 & 0 \\ v_\Delta & 0 \end{pmatrix} \begin{pmatrix} \nu^i \\ \ell^i \end{pmatrix} \\ &\ni -\frac{1}{2} \underbrace{\left(\sqrt{2} y_\Delta^{ij} v_\Delta \right)}_{=m_\nu^{ij}} \overline{\nu^{jc}} \nu^i \end{aligned}$$

Generates light ν_m masses via vev **WITHOUT** invoking sterile N_k !

$\Delta^{\pm\pm}$, Δ^{\pm} couple directly to W, Z, γ (unambiguous xsec prediction)



$\Delta^{\pm\pm}, \Delta^{\pm}$ couple directly to W, Z, γ (unambiguous xsec prediction)



“State of the Art” Event Generation

- Hard coded pair production of $\Delta^{\pm\pm}\Delta^{\mp}(\mp)$ in Pythia8 and CalcHEP

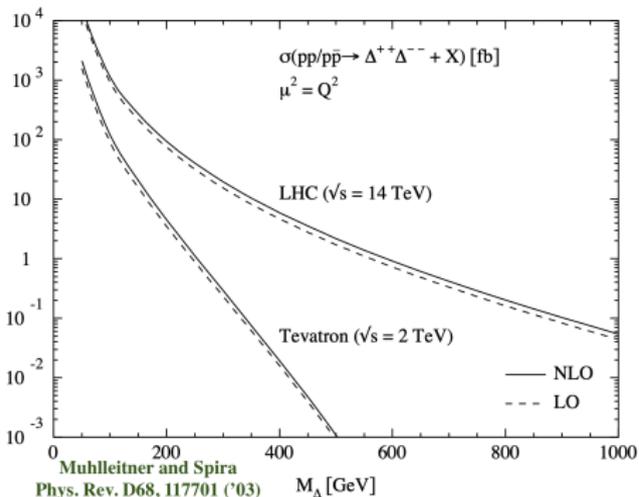
(no longer actively supported)

- Event generation at LO+PS

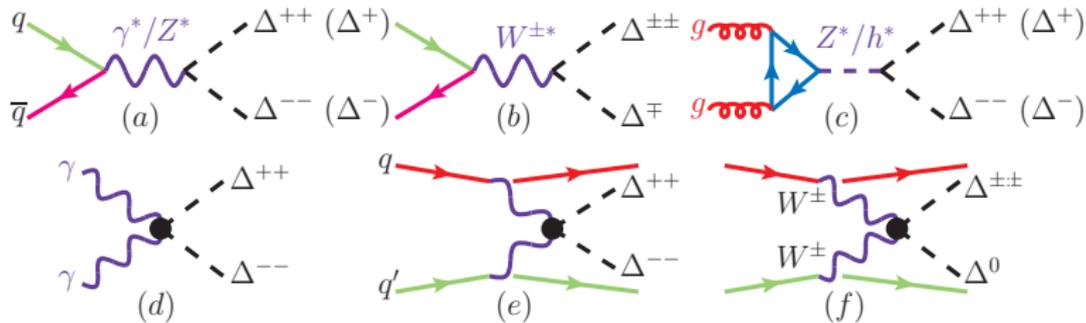
(no MLM matching, etc)

- Normalized to NLO with CTEQ6 \rightarrow

(sizable PDF uncertainties)



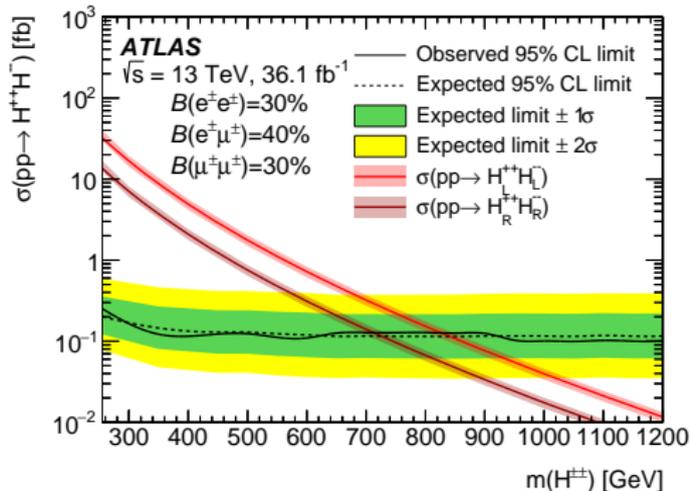
Limitations of “State of the Art”



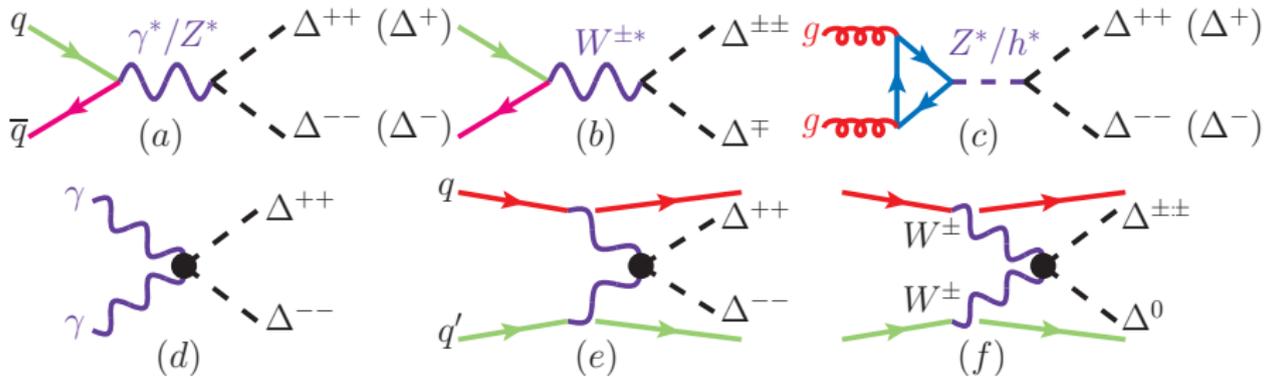
- Fixed production channels
do not reflect full parameter spaces

- Fixed decay modes do not utilize
potential of full LHC datasets

- does not facilitate newer analysis
techniques e.g., jet vetoes or fast
recasting tool chains



Type I Seesaw @ NLO³



³ with B. Fuks, O. Mattelaer, and M. Nemevšek [1912.08975], feynrules.irmp.ucl.ac.be/wiki/TypeISeesaw

The lack of Monte Carlo support for $\Delta^{\pm\pm}$, Δ^{\pm} was an oversight

- Universal FeynRules Object (UFO) libs exist for many other models

See FeynRules database feynrules.irmp.ucl.ac.be/wiki/ModelDatabaseMainPage

- UFOs encode Feynman rules for mainstream event generators, e.g. **MadGraph**, to simulate **BSM in HEP experiments** (not just colliders)

After a request by **hep-ex'ers**, we wrote **TypeIISeesaw@NLO** UFO libraries

available at feynrules.irmp.ucl.ac.be/wiki/TypeIISeesaw

```
> import model TypeII_NLO_UFO
> generate p p > d++ d-- [QCD]
> output TypeIInlo_DYX_DxxDxx_NLO
> launch
> order=NLO
> fixed_order=ON
> set mdpp scan1:range(100,2001,50)
> set dynamical_scale_choice -1
> set no_parton_cut
> set jetalgo -1
> set jetradius 0.4
```

TypeIISeesaw : Canonical type II Seesaw at NLO in QCD

Contact Information

Implementation author: Benjamin Fuks

- LPTHE / Sorbonne U.
- fuchs@...

In collaboration with Miha Nemešek and Richard Ruiz. See [arXiv:1911.NNNNN](https://arxiv.org/abs/1911.NNNNN) [hep-ph].

Model Description and FeynRules Implementation

We extend the Standard Model by adding a scalar field lying in the adjoint representation of the weak group Δ with an hypercharge equals to 1. After electroweak rise to one extra CP-even scalar Higgs Δ^0 , one charged scalar Δ^{\pm} and one doubly-charged scalar $\Delta^{\pm\pm}$ that are mostly of a triplet nature. The corresponding Lagrangian

$$\mathcal{L} = \mathcal{L}_{SM} + \text{Tr}[D_{\mu}\Delta^{\dagger}D^{\mu}\Delta] - V_{\Delta} + \mathcal{L}_{Y_{\Delta}} \quad (1)$$

where the extra terms respectively represents the gauge-invariant kinetic terms for the triplet, the extra pieces to the scalar potential and the Yukawa interaction. The above Lagrangian was implemented in the Feynman Gauge into FeynRules 2.3.35. QCD renormalisation and R_2 rational counterterms were determined using FeynRts 3.9. Feynman rules were collected into a single UFO, available below. In the normal hierarchy and inverted hierarchy UFO models, four massless quark diagonal CKM matrix entries. These additions permit tree-level calculations at LO and NLO in QCD and loop-induced calculations at LO in QCD using MadGraph_a

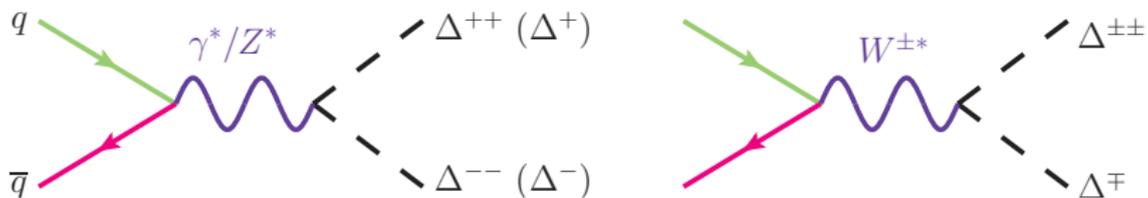
In the case of a normal neutrino mass hierarchy ($M_{\nu 1} < M_{\nu 2} < M_{\nu 3}$), the model contains 16 free parameters (on top of the Standard Model ones):

- the mass of the first neutrino $M_{\nu 1}$ (the other neutrino masses are internal parameters);
- the squared neutrino mass differences dmsq21 ($\Delta m_{21}^2 > 0$) and dmsq31 ($\Delta m_{31}^2 > 0$) collected in the LH block MNLU;
- the oscillation parameters θ_{12} (θ_{12}), θ_{13} (θ_{13}), θ_{23} (θ_{23}), δCP (φ_{CP}), phIM1 (φ_{12}), phIM2 (φ_{13}) collected in the LH block PMNS;
- the masses of the Higgs (PDG 125), Δ^0 (PDG 44), Δ^{\pm} (PDG 38) and $\Delta^{\pm\pm}$ (PDG 44) fields (in the LH block MASS);
- the triplet vev v_{Δ} (block VEVDELTA)
- two quartic couplings collected in the block QUARTICS

some newish results⁴

⁴ with B. Fuks and M. Nemevšek [[1912.08975](#)]

Neutral Current and Charged Current Drell-Yan



Charged Current and Neutral Current Drell-Yan

Drell-Yan production at NLO(+PS) in QCD is sanity check

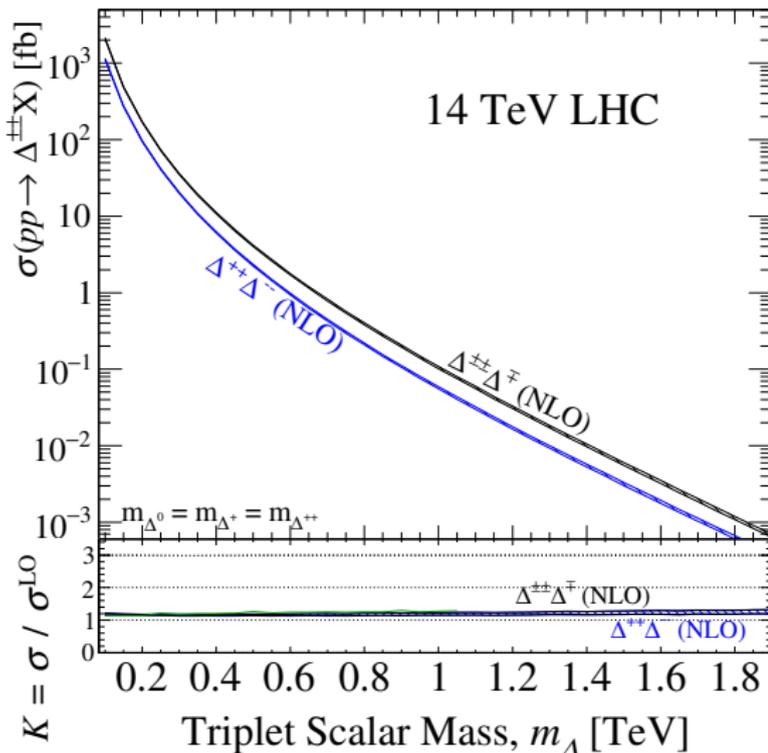
only other group to compute this is Muhlleitner and Spira ('03)

- With one event generator, all the following can be computed:

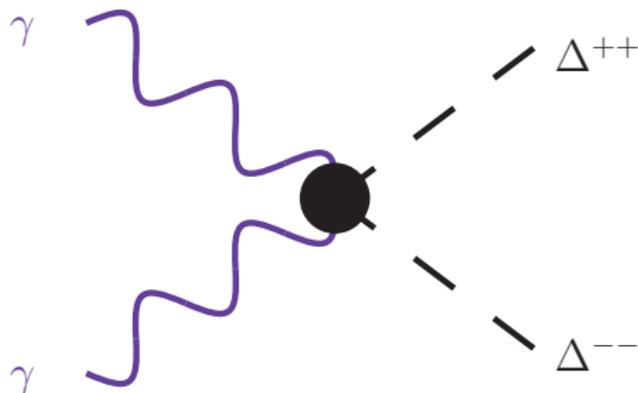
$$pp \rightarrow \Delta^{++}\Delta^{--}, \\ \Delta^{\pm\pm}\Delta^{\mp}, \Delta^+\Delta^-, \Delta^{\pm}\Delta^0, \\ \Delta^{\pm}\chi^0, \Delta^0\chi^0 \quad (\leftarrow \text{unexplored})$$

- These other channels needed to check **quantum numbers** and discriminate against **Georgi-Machacek, etc.**

- At NLO, jet vetoes can be implemented (**yay!**) (see backup)

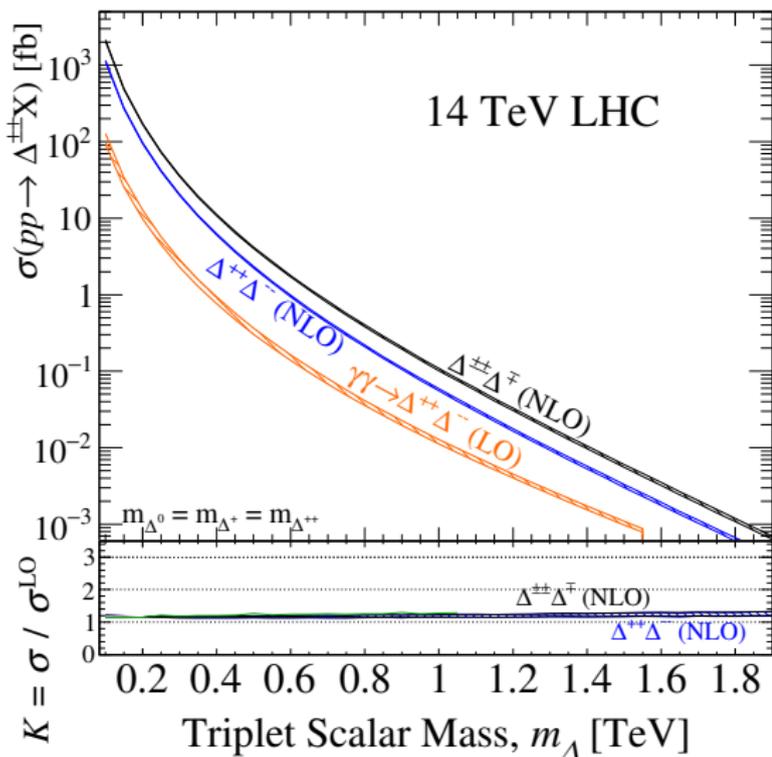


photon fusion



$\gamma\gamma \rightarrow \Delta^{++}\Delta^{--}$ is wickedly cool!

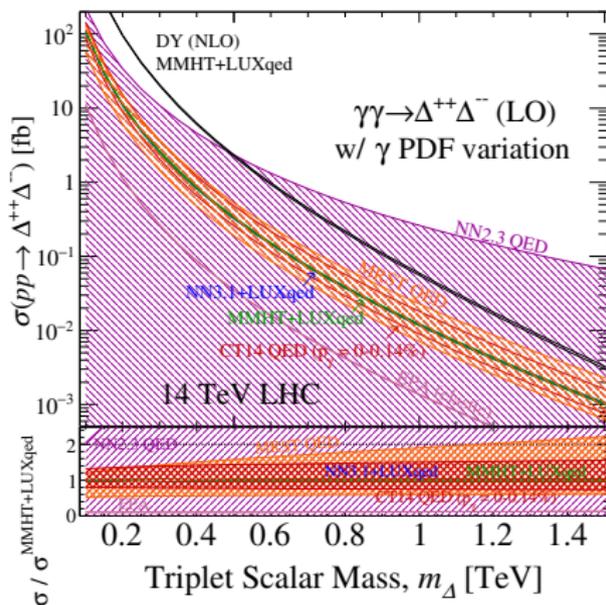
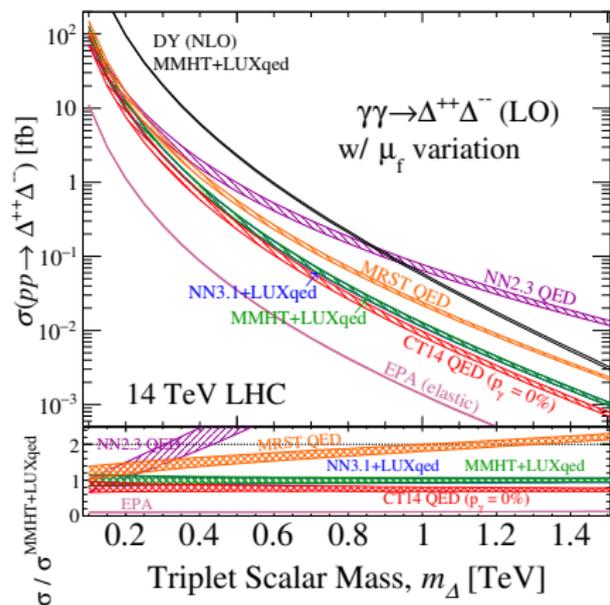
- **Subleading** but ultra peripheral collisions have *tiny* backgrounds
- **LO+Pythia8*** can match γ to $q \rightarrow \gamma q$ splitting, i.e., match to forward jets
- Many modern γ PDFs on the market these days. How do different γ PDFs compare?



*For details, see home.thep.lu.se/~torbjorn/pythia81html/SpacelikeShowers.html

Historically, back-and-forth about importance of $\gamma\gamma \rightarrow \Delta^{++}\Delta^{--}$

- How do different γ PDFs compare?

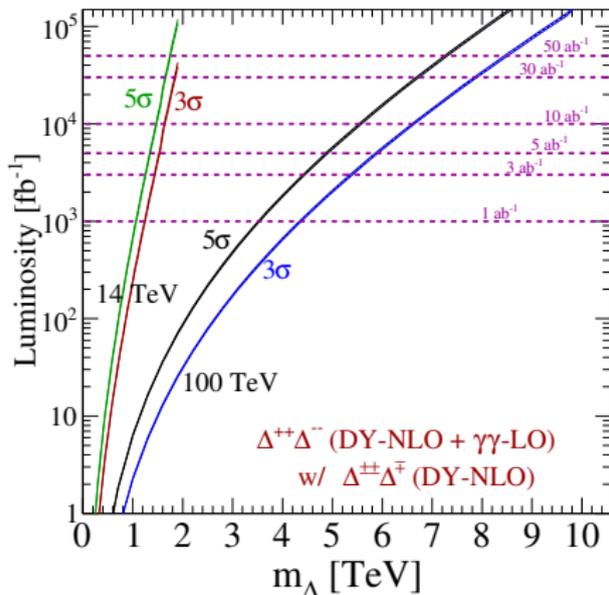
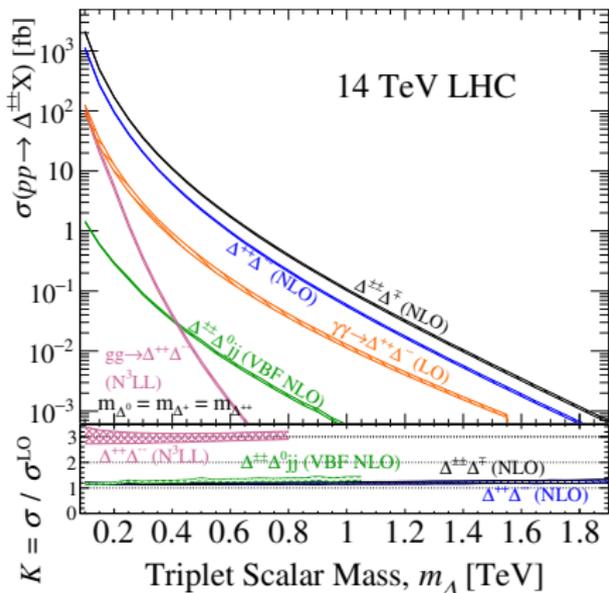


Most PDFs agree and discrepancies have disappeared (again)

- Warning:** NNPDF 2.3 QED (default in MG5aMC) has poor description of γ PDF; well-known and due to limited statistics (not method)

outlook (see backup for more channels!)

NEW: a revised outlook for both $\sqrt{s} = 14$ TeV and 100 TeV!



- At LHC with $\mathcal{L} = 5 \text{ ab}^{-1}$, 3 σ sensitivity up to $m_{\Delta} \sim 1.5 \text{ TeV}$
- At $\sqrt{s} = 100 \text{ TeV}$ with $\mathcal{L} = 30 - 50 \text{ ab}^{-1} \implies m_{\Delta} \approx 8 - 9 \text{ TeV}$
- **Warning:** can be improve for specialized final state / parameter space

Lots not covered!

Theory uncertainties (see the paper! [1912.08975])

- Uncertainties at LO, NLO, NLO+NNLL(veto), NNLL(threshold)
- PDF uncertainties (statistical PDF and modeling)
- 4FS vs 5FS scheme, with selection cuts

Monte Carlo recommendations (see the paper! [1912.08975])

- choices of scales, generator cuts, PDFs, etc. (care always needed)
- explicit MG5aMC@NLO instructions and lots tables for checks

Sensitivity beyond pp colliders (reader exercise)

- TypeIISeesaw UFO + MG5aMC@NLO enables beam dump, DIS/ion, electron/muon colliders, astroparticle studies
- Useful since Δ^0 (CP even) and ξ^0 (CP odd) are under-studied, long-lived particles and stable, dark matter candidates

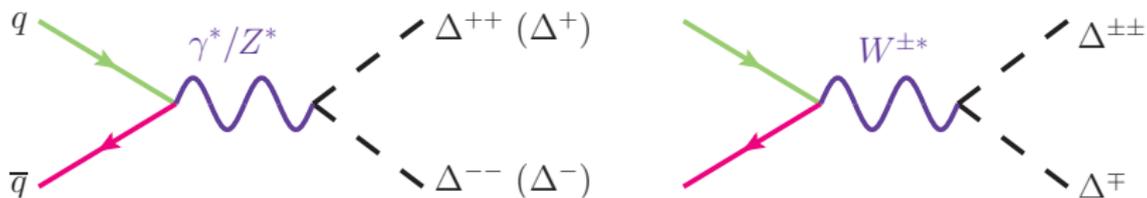
Designed to be user friendly but bother us with Q's or collaboration ideas!



Thank you.

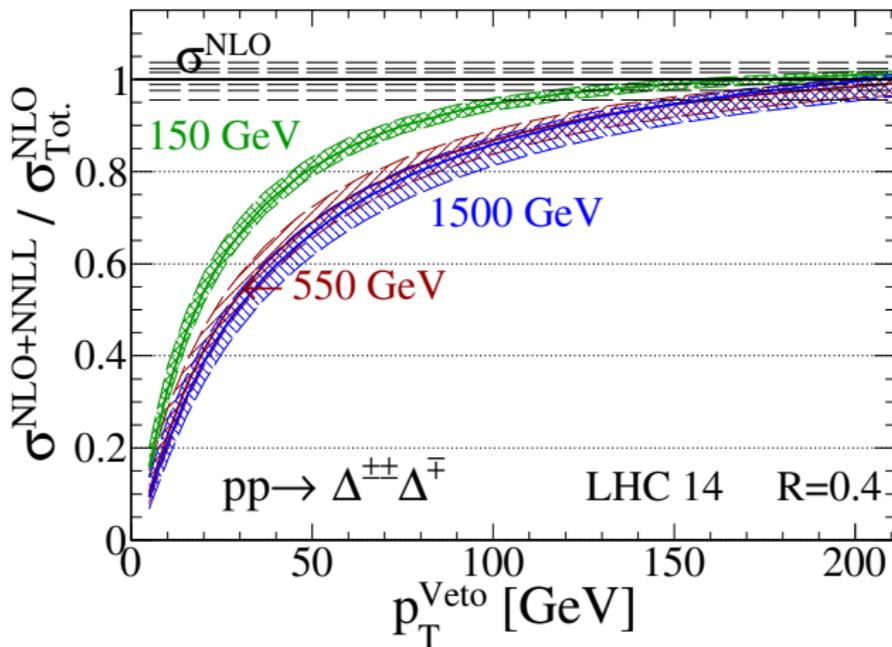
Backup

Jet Vetoes in Neutral Current and Charged Current Drell-Yan



Plotted: jet veto efficiency $\varepsilon = \sigma^{\text{NLO+NNLL}}(p_T^j < p_T^{\text{Veto}}) / \sigma_{\text{Tot.}}^{\text{NLO}}$

using MG5aMC@NLO+SCET. See Becher, et al [1412.8408] for details!

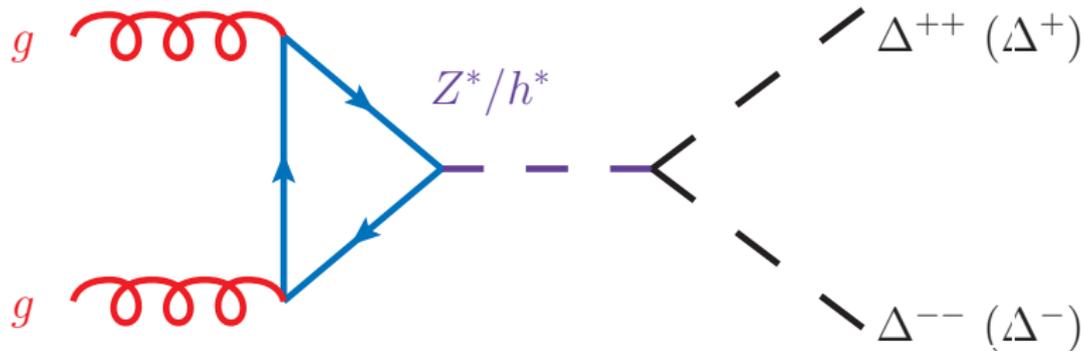


Signal veto efficiency is pretty terrible for $p_T^{\text{Veto}} = 20 - 50 \text{ GeV}$

- Clear that dynamic / safe jet veto schemes are meritted

see, e.g., Pascoli, RR, et al [1805.09335, 1812.08750], as well as Fuks, Nordstrom, RR, and Williamson [1901.09937]

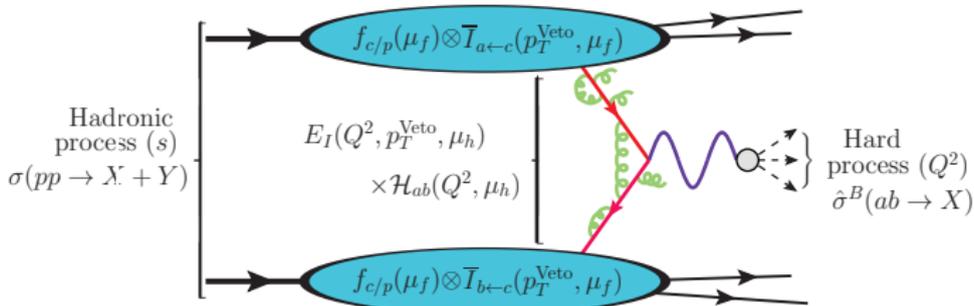
Gluon Fusion



Threshold Resummation in SCET

Each component of the Factorization Theorem obeys $\mathcal{D}X = \Gamma X$
 Collecting like-terms shows scattering rates have the form⁵:

$$d\sigma(pp \rightarrow \mathcal{B} + X) = f(\mu_f) \times f(\mu_f) \times \Delta(\mu_s) \times \underbrace{U(\mu_f, \mu_s, \mu_h)}_{\text{Messy, but universal}} \times d\hat{\sigma}(\mu_h)$$



This expression is known as “RG-improved” factorization theorem (SCET)

- Sudakov-like factors act as “dressing” functions, in QFT sense

⁵Contopanagos, et al; Becher, Neubrt etc; Bauer, Stewart, Tackmann etc

Threshold resummation at (approximately) N^3LL captures leading normalization at N^2LO in QCD

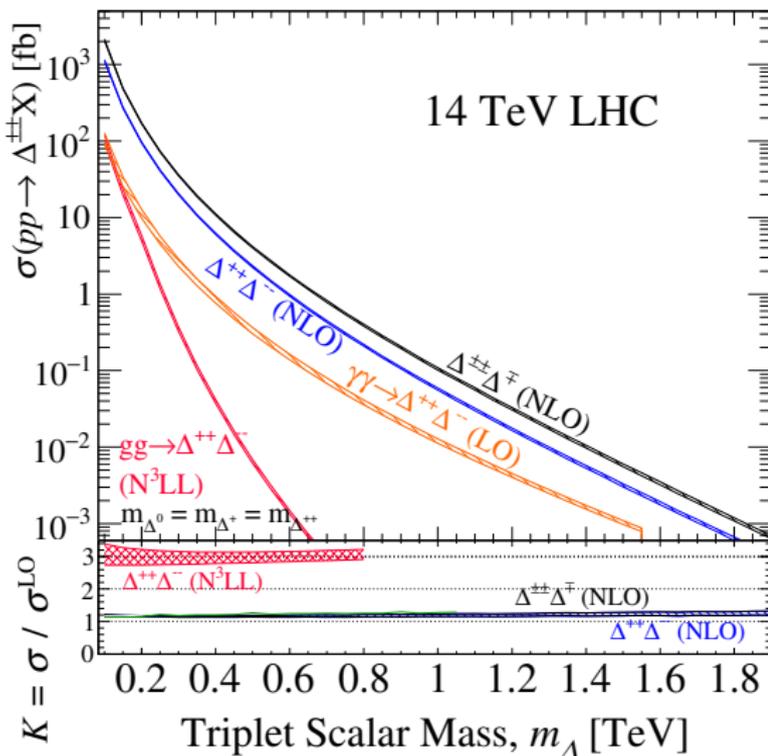
for N^3LL details, see RR, Spannowsky, Waite [1706.02298]

Loop-induced $gg \rightarrow \Delta\Delta$ is possible at LO with new UFO
 - N^3LL a bit more work

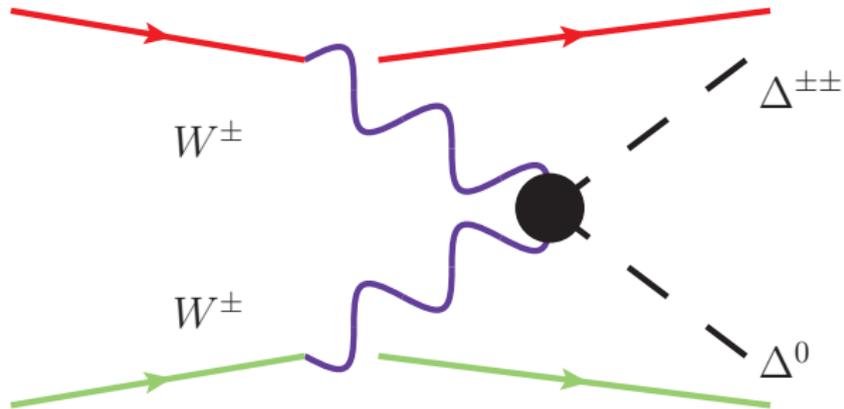
For lower masses, gg and $\gamma\gamma$ are competitive, but $\gamma\gamma$ becomes more relevant

(again, cool)

Caveat: we assume all Δ masses degenerate
 \implies scalar-scalar couplings are small (probably not justified)



Weak Boson Fusion⁶

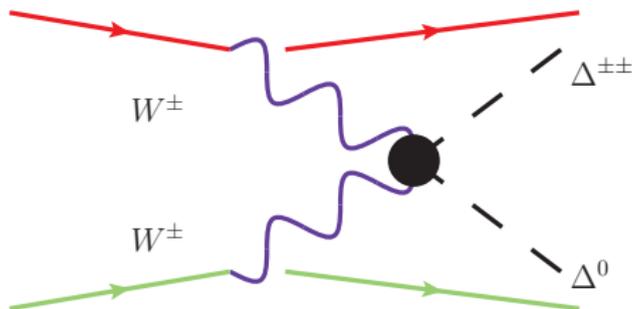


⁶many details being omitted!

$\Delta^{\pm\pm}$ from Vector Boson Scattering

Weak boson fusion is a useful production mechanism of $\Delta^{\pm\pm}$

- Except $W^\pm W^\pm \rightarrow \Delta^{\pm\pm}$ is vev/coupling-suppressed
- Except $VV \rightarrow \Delta^{++}\Delta^{--}$ is mass-suppressed



$W^\pm W^\pm \rightarrow \Delta^{\pm\pm}\Delta^0$ for $m_{\Delta^0} \ll m_{\Delta^{++}}$ keeps important features

- **LNV/LFV** dilepton signature with resonant invariant mass
- **Difficulty** is modeling beyond NLO in QCD (needed for jet veto modeling)

$pp \rightarrow \Delta^{\pm\pm} \Delta^0 jj + \text{VBS Cuts at NLO is under control!}$

$\sigma_{\text{DY}}/\sigma_{\text{VBF}} \sim 10^3$ difference
hides three things:

- VBF cuts **already applied at generator-level**

- Scalar-scalar couplings are **small since mass degenerate**

- Larger mass splittings
 \Rightarrow **higher mass reach**

