

[arXiv:2012.14889]

# Radiative corrections to charged Higgs bosons decays in NMSSM

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# Introduction

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- There are several unsolved problems in the SM.

Unsolved phenomena:

- Tiny neutrino masses
- Dark matter
- Baryon asymmetry of the Universe
- etc.

Theoretical problems:

- Hierarchy problem
- Strong CP problem
- etc.

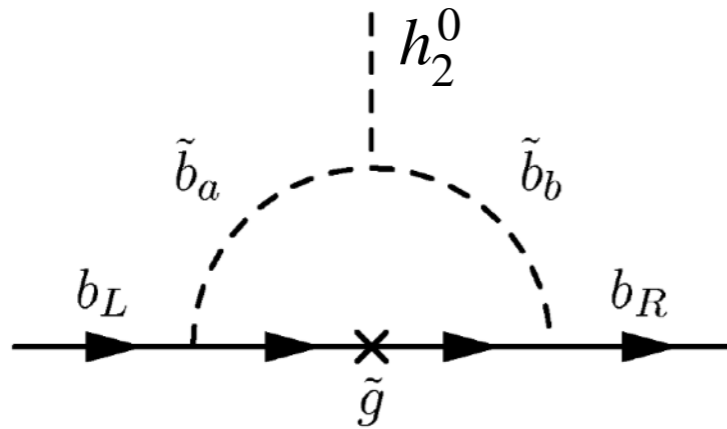
↔ This relates paradigm of BSM.

- Symmetry solve the Hierarchy problem? → Supersymmetry (SUSY)

# Indirect search of NP via precise calculations

- No direct evidence of new particles
  - In future collider experiments (HL-LHC, ILC, FCC-ee, the CEPC, etc.), the discovered Higgs boson will be precisely measured.
- It would be more important to focus on indirect searches of NP via precise measurement of Higgs bosons.
- Precise calculations for various new physics models are necessary in order to compare with precise measurement.

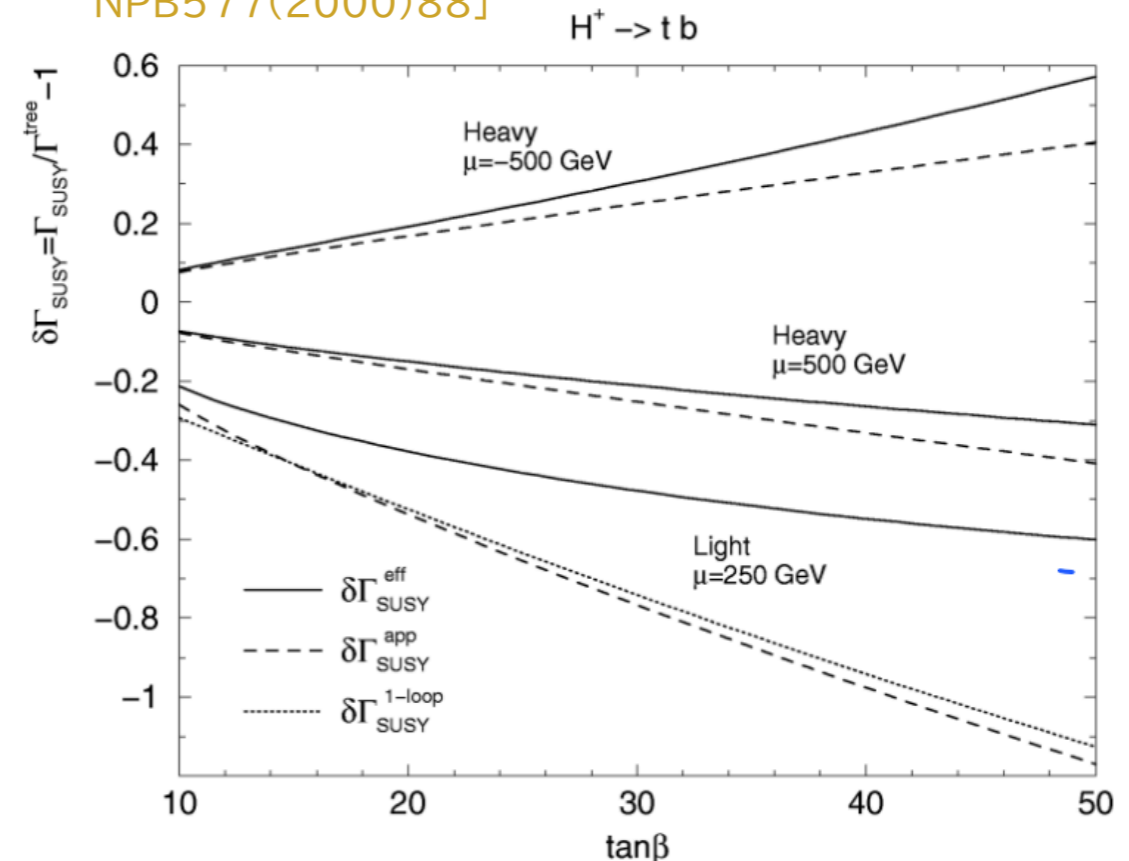
Ex.)  $\Delta_b$  corrections to  $H^+ \rightarrow tb$  in MSSM



$$\bar{t}P_R b H^+ : (y_b + \Delta y_b \tan \beta)$$

→ In case of  $\tan \beta \gg 1$ ,  $\Delta_b$  correction can be large.

[M. Carena, D. Garcia, U. Nierste, C. Wagner, NPB577(2000)88]



# Next-to Minimal supersymmetric model (NMSSM)

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- NMSSM is a simplest extension of the MSSM with a singlet field:

$$\mathcal{L}_{\text{NMSSM}} = \mathcal{L}_{\text{MSSM}} + \mathcal{S}$$

➔ Physical Higgs states:  $H_1, H_2, H_3, A_1, A_2, H^\pm$

- There exist 5 neutralinos due to the new degree of freedom, singlino.
- $\mu$  problem can be solved in this model:

$$W_{\text{MSSM}} \ni \mu H_u H_d$$

$$\text{NMSSM: } \mu_{\text{eff}} \equiv \lambda \langle S \rangle \leftarrow \lambda S H_u H_d$$

- Constraint of mass of the Higgs boson may be easily satisfied than MSSM:

$$\text{MSSM: } (m_{H_1}^{\text{tree}})^2 = M_Z^2 \cos 2\beta$$

$$\text{NMSSM: } (m_{H_1}^{\text{tree}})^2 = M_Z^2 \left( \cos 2\beta + \frac{\lambda^2}{g_1^2 + g_2^2} \sin 2\beta \right)$$

# Precise calculations of Higgs bosons

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Many studies of precise calculations of Higgs bosons have been performed in the NMSSM.

- Higgs boson masses ( Full 1-loop ) [G. Degrassi, P. Slavich, NPB 825 (2010)]  
(  $O(a_t a_s)$  2-loop ) [F. Staub, W. Porod, B. Herrmann, JHEP10(2010)], etc  
(  $O(a_t^2)$  2-loop ) [M. Mühlleitner, D. T. Nhung, H. Rzehak, K. Walz, JHEP05(2015)]  
[M. D. Goodsell, K. Nickel, F. Staub, PRD91(2015)], etc  
[T.N. Dao, R. Gröber, M. Krause, M. Mühlleitner, H. Rzehak, JHEP08(2019)]
- Higgs boson decays
  - $H_i, A_i \rightarrow ff, VV$  ( Full 1-loop ) [F. Domingo, S. Heinemeyer, S. Paßehr, G. Weiglein, Eur.Phys.J.C78(2018)]
  - $H_i \rightarrow H_j H_k$  ( Full 1-loop ) [D. T. Nhung, M. Mühlleitner, J. Streicher, K. Walz, JHEP11 (2013)]  
(  $O(a_t a_s)$  2-loop ) [G. Belanger, V. Bizouard, F. Boudjema, G. Chalons, PRD96 (2017)]  
[M. Mühlleitner, D. T. Nhung, H. Ziesche, JHEP 12 (2015)]
  - $A_i \rightarrow \tilde{t}\tilde{t}$  ( Full 1-loop ) [J. Baglio, C. Krauss, M. Mühlleitner, K. Walz, JHEP 10 (2015)]
  - All 2-body neutral Higgs decays, BRs [J. Baglio, T. N. Dao, M. Mühlleitner, 1907.12060]
  - $H_i^\pm \rightarrow WH_i$  ( Full 1-loop ) [T. N. Dao, L. Fritz, M. Krause, M. Mühlleitner, S. Patel, Eur.Phys.J. C80 (2020)]

→ Missing part for Higgs boson decays is charged Higgs boson decays.

# This talk

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- We calculated NLO corrections to various charged Higgs bosons in the NMSSM.
- We investigated the impact of the NLO corrections for each decay process.

Open questions:

- What is typical size of NLO corrections for each decay?
- How does CPV effects change the NLO corrections?

# Decay rates for charged Higgs (Leading order)

$H^+ \rightarrow tb$  :

$$\lambda(x, y) = (1 - x - y)^2 - 4xy$$

$$\Gamma(H^+ \rightarrow tb) = \frac{3M_{H^\pm}}{8\pi} |V_{tb}|^2 \lambda^{1/2} \left( \frac{M_t^2}{M_{H^\pm}^2}, \frac{M_b^2}{M_{H^\pm}^2} \right) \left[ \left( 1 - \frac{M_t^2}{M_{H^\pm}^2} - \frac{M_b^2}{M_{H^\pm}^2} \right) \left( \frac{M_t^2}{v^2} \frac{1}{\tan^2 \beta} + \frac{M_b^2}{v^2} \tan^2 \beta \right) - 4 \frac{M_t^2 M_b^2}{M_{H^\pm}^2 v^2} \right]$$

→ Depending on value of  $\tan\beta$ , top Yukawa or bottom Yukawa can dominate.

$H^+ \rightarrow \chi_i^0 \chi_j^+$  : (CP conserving case)

$U, V, N$  : mixing matrix for electroweakinos

$$\Gamma(H^+ \rightarrow \chi_i^0 \chi_j^+) = \frac{M_{H^\pm}}{16\pi} \lambda^{1/2} \left( \frac{M_{\chi_i^+}^2}{M_{H^\pm}^2}, \frac{M_{\chi_j^0}^2}{M_{H^\pm}^2} \right) \left[ \left( 1 - \frac{M_{\chi_i^+}^2}{M_{H^\pm}^2} - \frac{M_{\chi_j^0}^2}{M_{H^\pm}^2} \right) (g_L^2 + g_R^2) - 4 \frac{M_{\chi_i^+} M_{\chi_j^0}}{M_{H^\pm}^2} g_L g_R \right]$$

$$g_R = -c_\beta \left[ \frac{g_1}{\sqrt{2}} N_{i1} V_{j2} + \frac{g_2}{\sqrt{2}} (N_{i2} V_{j2} + \sqrt{2} N_{i4} V_{j1}) + \lambda N_{i5} V_{j2} \tan \beta \right]$$

$$g_L : g_R [U \rightarrow V, N_{i4} \rightarrow -N_{i3}, \tan \beta \rightarrow -\cot \beta, c_\beta \rightarrow -s_\beta]$$

→  $g_{L/R} \sim \lambda$  ( $\chi_i^0 \sim \tilde{S}$ ),  $g_{L/R} \sim g_1$  or  $g_2$  ( $\chi_i^0 \sim \tilde{B}, \tilde{W}^3, \tilde{H}_{u/d}$ )

# Branching ratios for charged Higgs (Leading order)

Charged Higgs boson decays with the state of the art QCD corrections by using NMSSMCALCEW (w/o EW corrections)

- Definition of BRs

$$\text{BR}(H^+ \rightarrow WH) = \sum_{i=1,\dots,3} \text{BR}(H^+ \rightarrow WH_i) + \sum_{i=1,2} \text{BR}(H^+ \rightarrow WA_i)$$

$$\text{BR}(H^+ \rightarrow \chi_{1/2}^+ \chi^0) = \sum_{i=1,\dots,5} \text{BR}(H^+ \rightarrow \chi_{1/2}^+ \chi^0)$$

- Inputs

$$\tan \beta = 10, \lambda = 0.09$$

- Mass spectrum

$$m_{\chi_1^0} = 748 \text{ GeV}, m_{\chi_1^+} = 819 \text{ GeV}$$

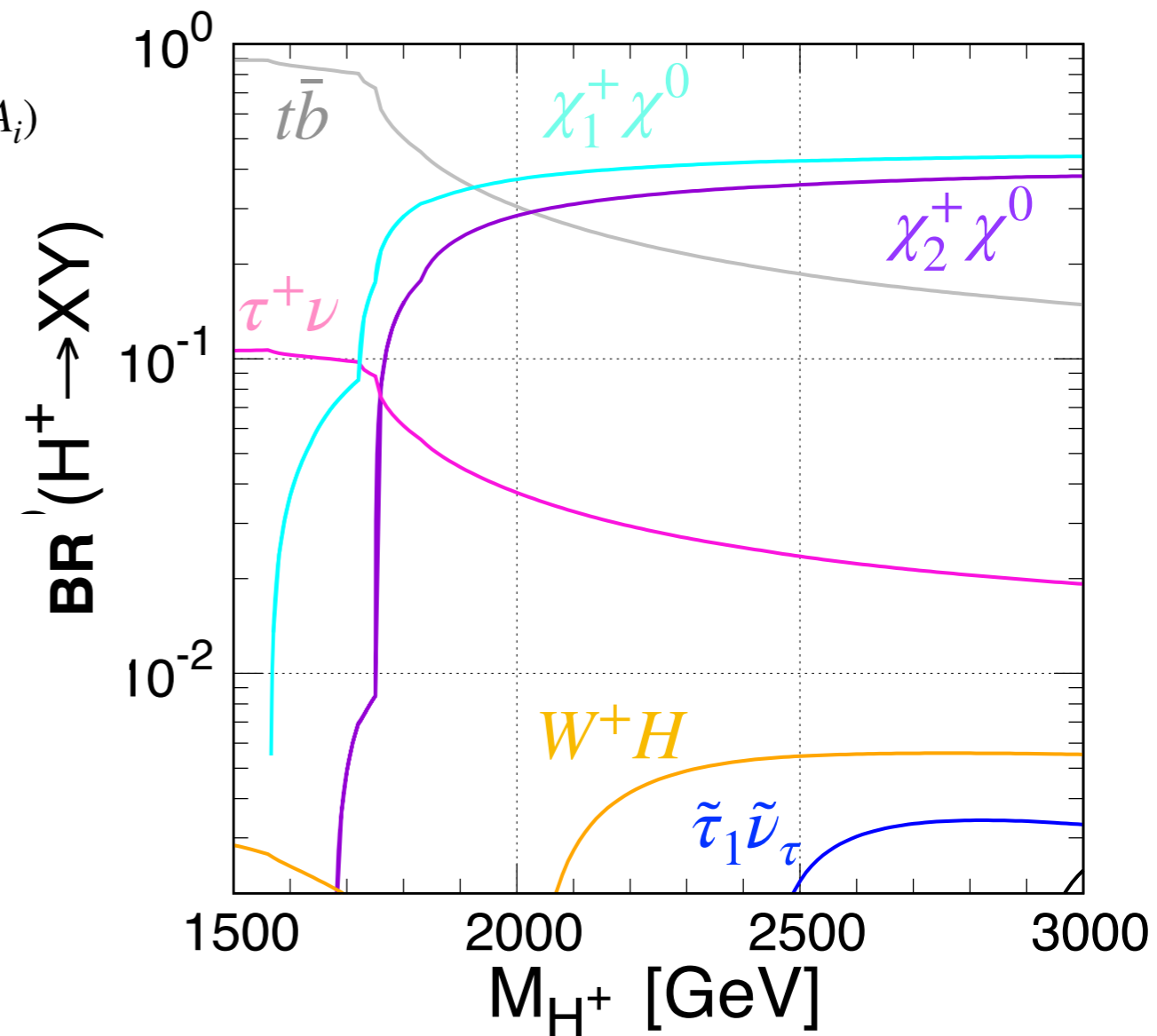
$$m_{\tilde{t}_1} = m_{\tilde{b}_1} = 1.4 \text{ TeV}$$

$$m_{\tilde{\nu}_\tau} = 1.2 \text{ TeV}, m_{\tilde{\tau}_1} = 1.1 \text{ TeV}$$

$$m_{H_2} = 1.1 \text{ TeV}, m_{A_1} = 1.9 \text{ TeV}$$

→  $H^+ \rightarrow t\bar{b}$  is the main decay mode in  $m_{H^+} \lesssim 2 \text{ TeV}$ .

→  $H^+ \rightarrow \chi_1^+ \chi^0$  dominates in  $2 \text{ TeV} \lesssim m_{H^+}$ .



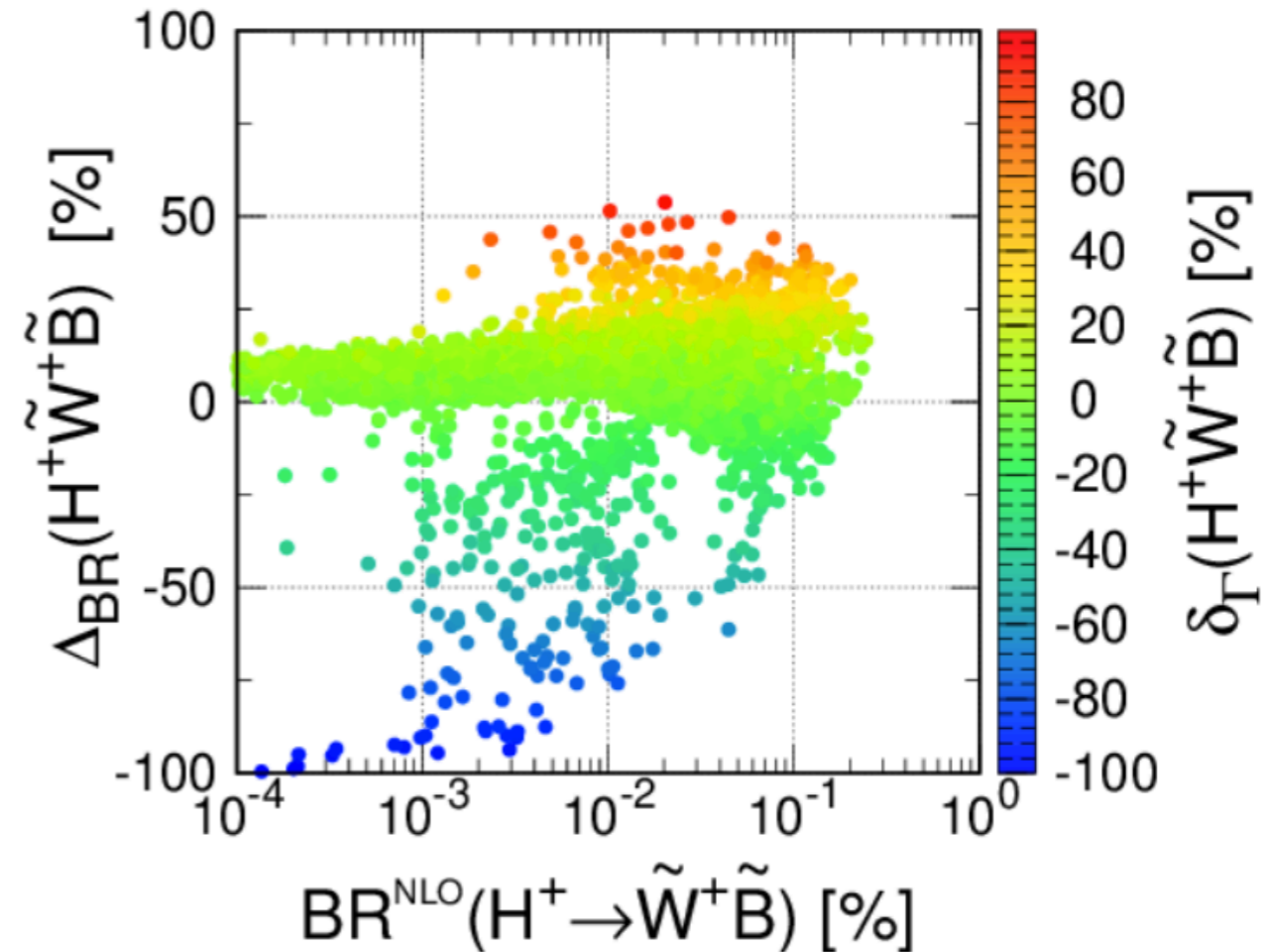
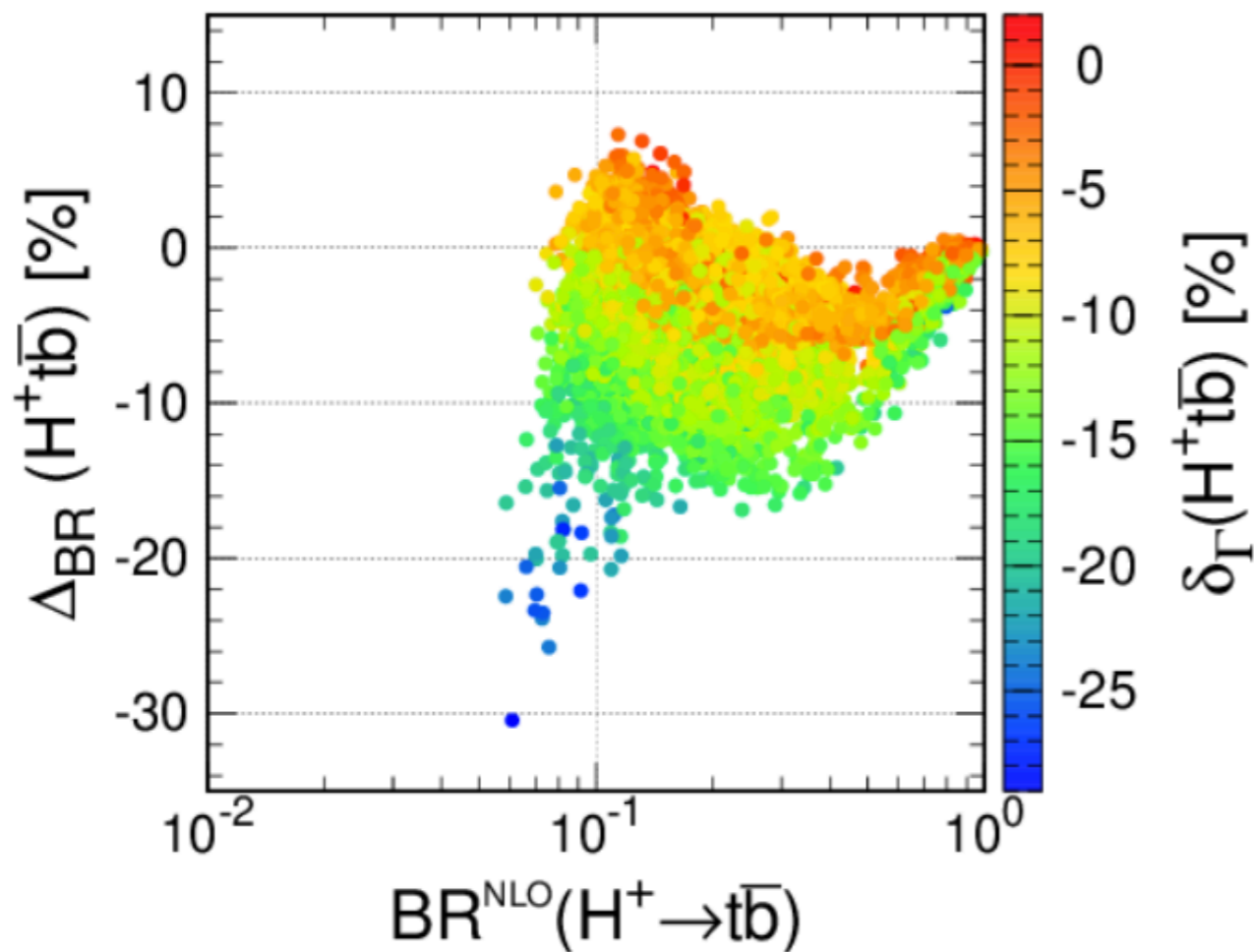




# NLO corrections: $BR(H^+ \rightarrow tb)$ , $BR(H^+ \rightarrow \tilde{W}^+ \tilde{B})$

$$\Delta_{\text{BR}} = \frac{\text{BR}^{\text{NLO}} - \text{BR}^{\text{LO}}}{\max(\text{BR}^{\text{NLO}}, \text{BR}^{\text{LO}})}, \quad \delta_{\Gamma} = \frac{\Gamma^{\text{NLO}}}{\Gamma^{\text{LO}}} - 1$$

[T. N. Dao, M. Mühlleitner, S. Patel, KS]

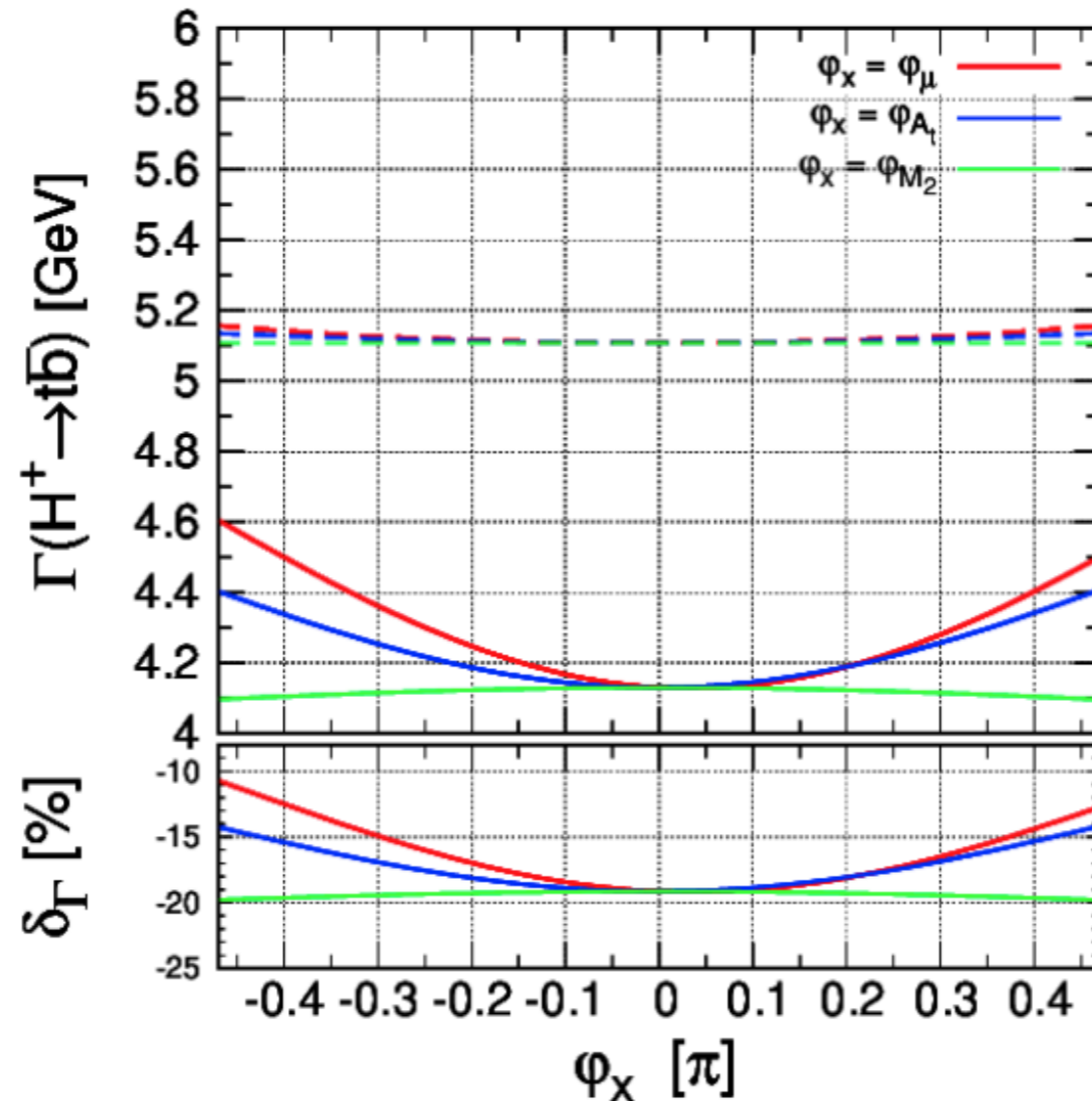


- For  $H^+ \rightarrow tb$ , maximum size of NLO corrections is  $\sim -30\%$ .
- For  $H^+ \rightarrow \tilde{W}^+ \tilde{B}$ , large corrections,  $|\Delta_{\text{BR}}| \sim 100\%$ , can appear.  
 → The main contributions is wave function renormalizations for electroweakinos.

# Effect of CP violation

Ex.)  $H^+ \rightarrow tb$

[T. N. Dao, M. Mühlleitner, S. Patel, KS]



- Phase of  $\mu_{\text{eff}}, A_t, M_2$  is varied.
- At LO, slight phase dependence appear for  $\mu_{\text{eff}}, A_t$ .
- At NLO, decay rates charges match compared with those of LO.

# Summary

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- The Higgs boson will be precisely measured at the future collider experiments, such as the HL-LHC, the ILC, FCC-ee and the CEPC.
  - This means that the theoretical predictions should be accurately evaluated.
- We study NLO (SUSY +)EW and SUSY QCD corrections for various charged Higgs bosons decays in the complex NMSSM.
  - $H^+ \rightarrow tb$  : NLO corrections with  $\sim 30\%$  can be obtained.
  - $H^+ \rightarrow \tilde{W}^+ \tilde{B}$  : large corrections due to mixing of electroweakinos can be found in  $BR < 1\%$ .
  - CPV effect : NLO corrections can deviate from those of CP conserving case within  $\sim 10\%$ .