

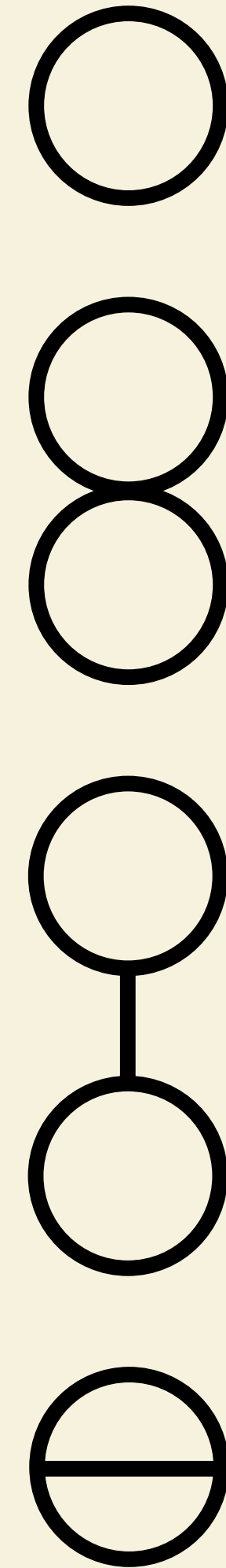


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# A Critical Look at the Electroweak Phase Transition

Based on work with Andreas Ekstedt  
arXiv: 2006.12614

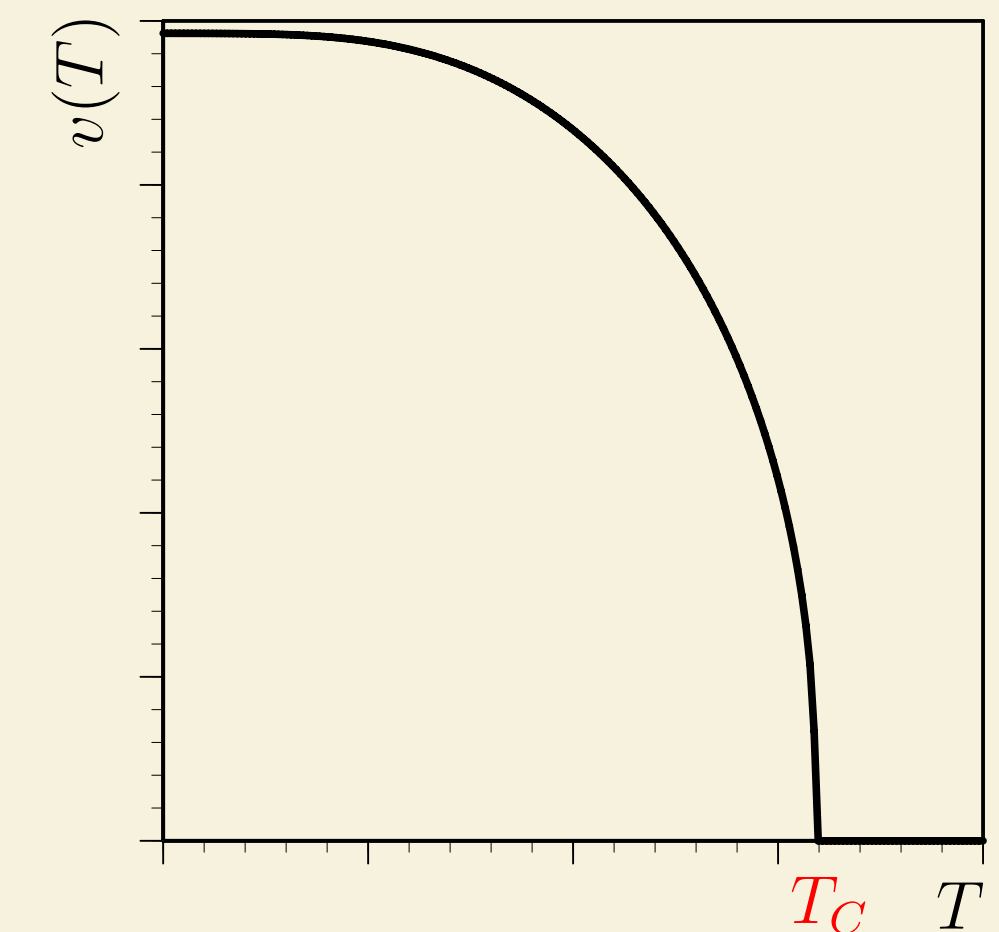
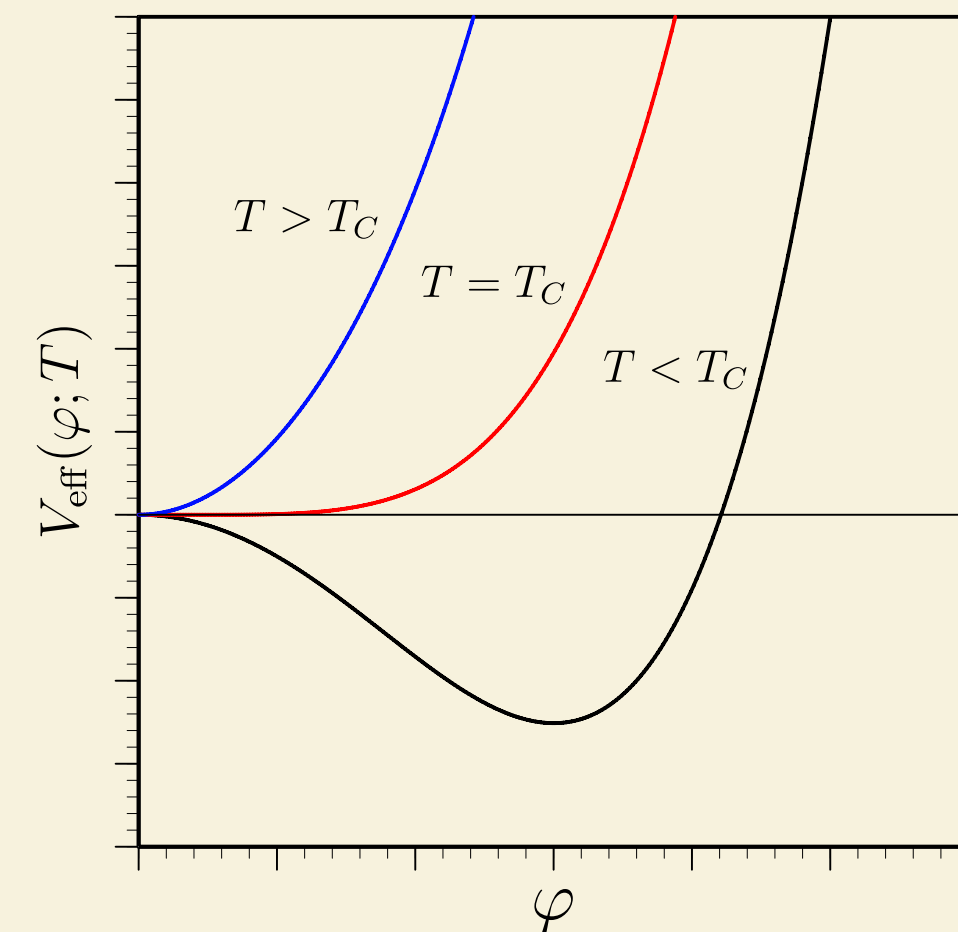
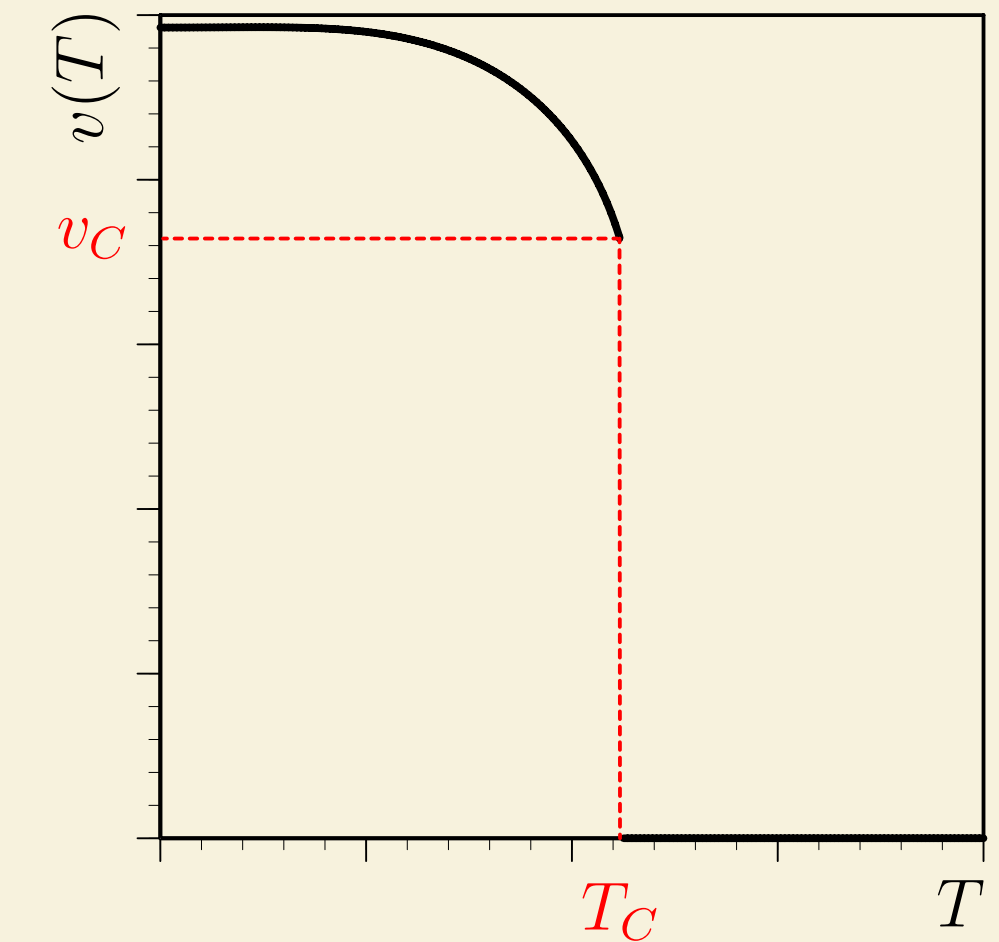
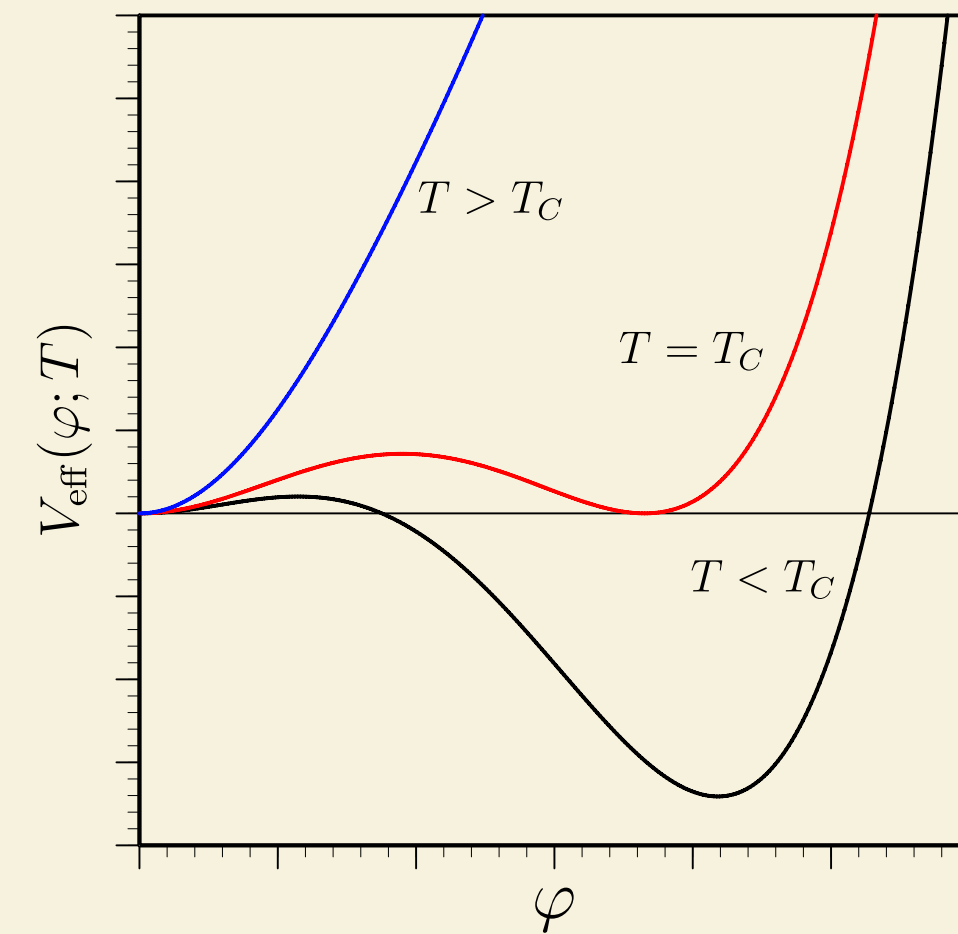
Johan Löfgren  
HPNP 2021  
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# Phase transitions

Senaha (DOI: 10.3390/sym12050733)

- Two different classes
- Electroweak phase transition might have detectable remnants
- **Two opposing problems:** gauge dependence and breakdown of the loop expansion

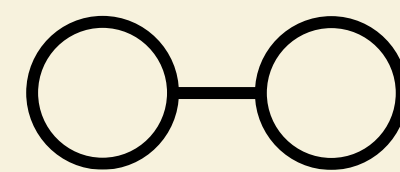
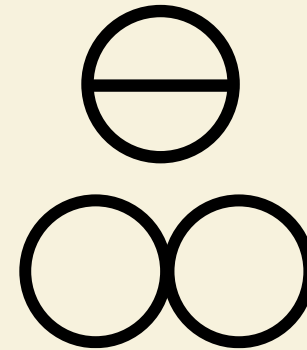
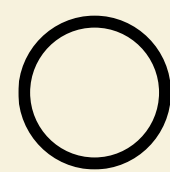


# Gauge dependence

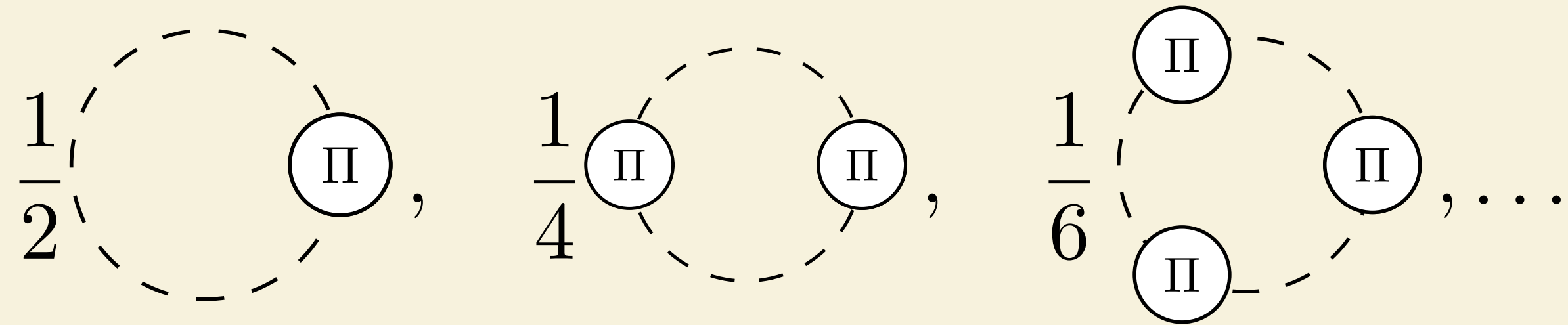
- The effective potential is gauge dependent
- Resolution: use the  $\hbar$ -expansion
- Be pedantic about loop counting!

S. Weinberg (DOI: [10.1103/PhysRevD.7.2887](https://doi.org/10.1103/PhysRevD.7.2887))  
 Fukuda and Kugo, (DOI: [10.1103/PhysRevD.13.3469](https://doi.org/10.1103/PhysRevD.13.3469))  
 Laine (arXiv: 9411252)  
 Patel and Ramsey-Musolf (arXiv: 1101.4665)

$$V(\phi^M) = V_0|_{\phi_0^M} + \hbar V_1|_{\phi_0^M} + \hbar^2 \left( V_2 - \frac{1}{2} (\phi_1^M)^2 \partial^2 V_0 \right) \Big|_{\phi_0^M} + \dots$$



# Resummation



$$\frac{d_n}{d_{n-1}} \sim e^{\frac{2T^2}{m^2}}$$

- The perturbative expansion breaks down
- Cannot order by loops anymore: a resummation must be made
- Then what about gauge invariance?

# Loop-induced symmetry breaking

$T = 0$  : *Coleman-Weinberg model*

- Regular loop counting:  $\lambda \sim e^2$

- Radiative symmetry breaking:  $\lambda \sim e^4$  ←

⋮ S. Coleman and E. Weinberg (DOI:10.1103/PhysRevD.7.1888) ⋮  
⋮ Andreassen, Frost, and Schwartz (1408.0287) ⋮

$$V^{\text{LO}} = \frac{\lambda}{24} \phi^4 + \frac{e^4}{16\pi^2} \phi^4 \left( -\frac{5}{8} + \frac{3}{2} \ln \frac{e\phi}{\mu} \right)$$

- The new **power counting** “forces” a resummation.

e.g. the Higgs mass must be resummed

- Perturbative expansions ordered according to a consistent counting are gauge invariant.

# Loop-induced barriers

$T > 0$  : *Abelian-Higgs*

- For phase transitions, think hard about the **power counting**
- $\lambda \sim e^3$  gives a barrier

Arnold and Espinosa (9212235)

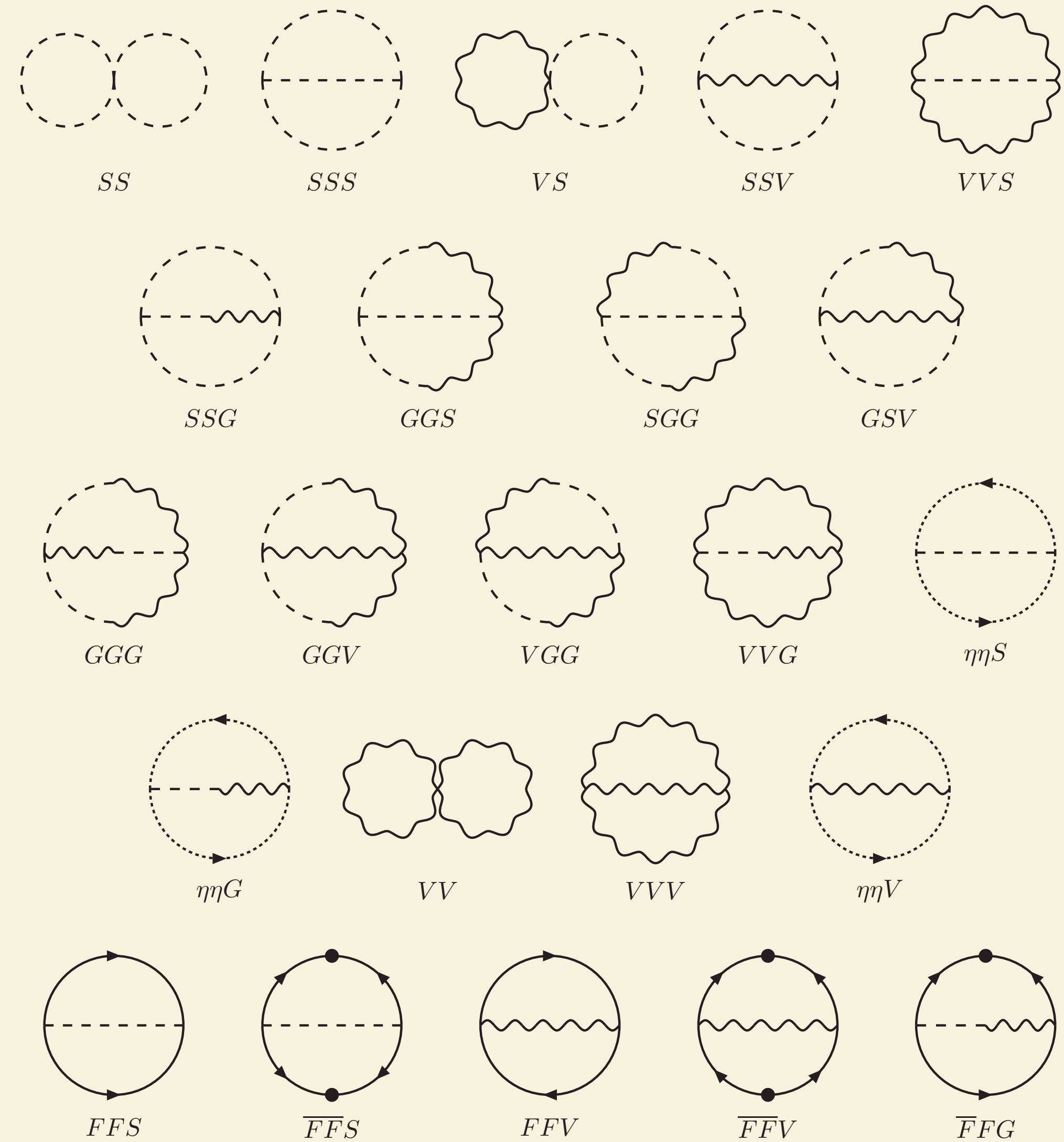
Konstandin and Garry (1205.3392)

$$V_{\text{LO}}(\phi, T) = -\frac{1}{2}m_{\text{eff}}^2(T)\phi^2 - e^3 \frac{T}{12\pi}\phi^3 + \frac{\lambda}{8}\phi^4.$$

- Resummation and gauge invariance at the same time

Ekstedt, JL: 2006.12614

Martin and Patel (1808.07615)



# Power-counting

Ekstedt, JL: 2006.12614

$$V_{\min} = e^{-1} V_{\text{LO}}|_{\phi_{\text{LO}}} + V_{\text{NLO}}|_{\phi_{\text{LO}}} + e^{1/2} V_{\text{NNLO}}|_{\phi_{\text{LO}}} \\ + e \left( V_{\text{N}^3\text{LO}} - \frac{1}{2} \phi_{\text{NLO}}^2 \partial^2 V_{\text{LO}} \right) |_{\phi_{\text{LO}}} + \dots$$

$$V_{\text{LO}}(\phi) \sim V_0(\phi) + \kappa T^2 \phi^2 (e^2 + \lambda^2) + \kappa T (2Z^{3/2} + Z_L^{3/2}),$$

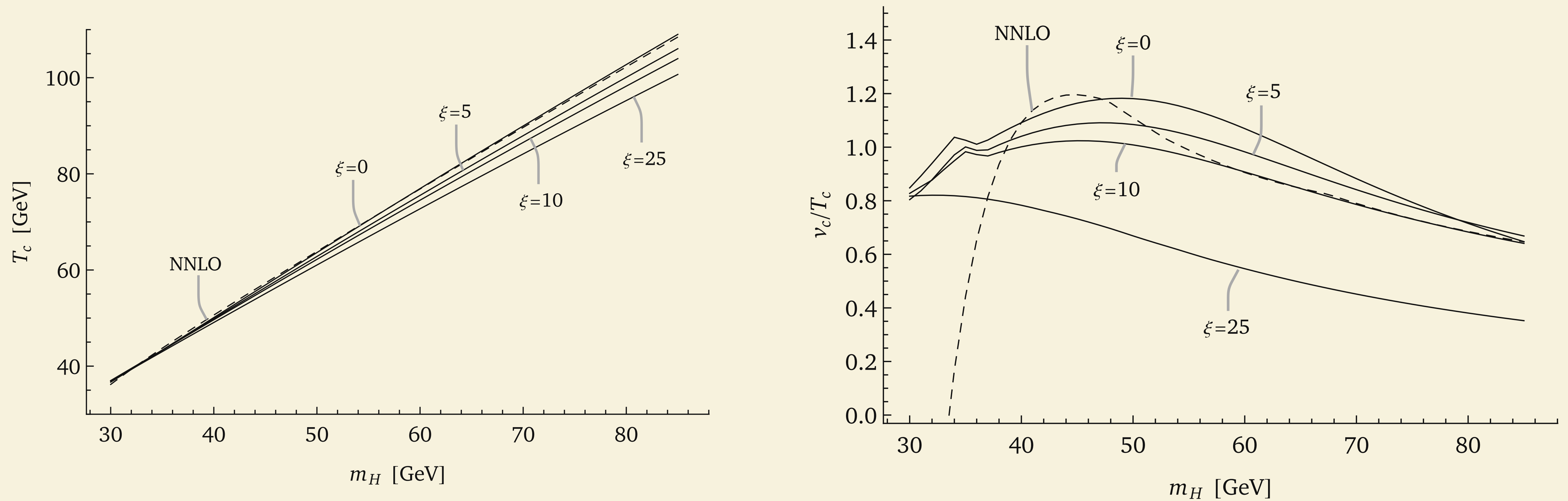
$$V_{\text{NLO}}(\phi) \sim \kappa Z^2 + \kappa^2 e^2 T^2 Z,$$

$$V_{\text{NNLO}}(\phi) \sim \kappa T (\bar{G}^{3/2} + \bar{H}^{3/2}),$$

$$V_{\text{N}^3\text{LO}}(\phi) \sim \kappa^2 e^2 T (Z^{3/2}) + \kappa^3 e^4 T^3 (Z^{1/2}) + \dots,$$

# Results

Ekstedt, JL: 2006.12614



Comparisons of **power counting** method and standard approach (gauge dependent numerical minimization), for some various values of the gauge fixing parameter, in the SM. (JL, Ekstedt: 2006.12614)



# Conclusions

- Gauge independence and resummation simultaneously with “ $\hbar$ -expansions” that employ proper **power counting**

# Future Work

- Further theoretical uncertainties  
Croon, Gould, Schicho, Tenkanen, White (2009.10080)
- Other observables related to tunneling
- Other models  
Camargo-Molina, Enberg, JL: 2103.14022