

HIGGS PHYSICS @ HIGH SCALES

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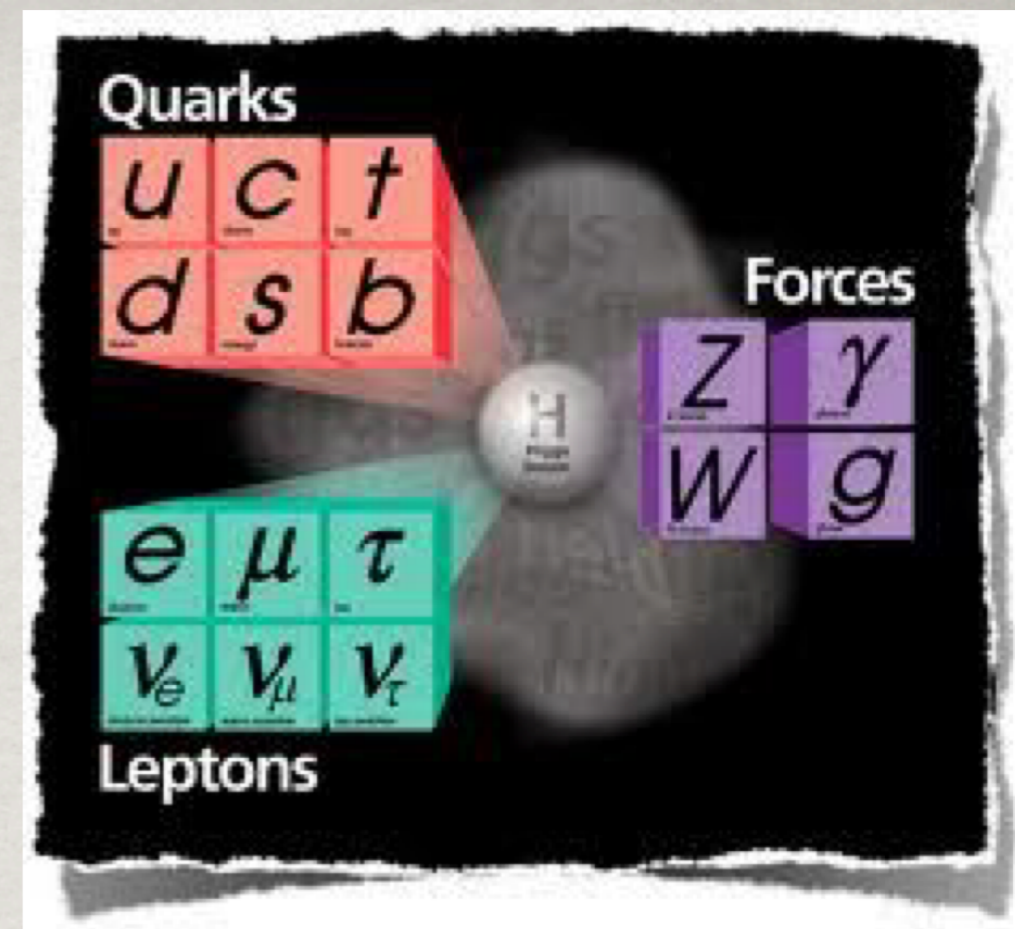
HPNP 2021, Osaka University

March 25, 2011

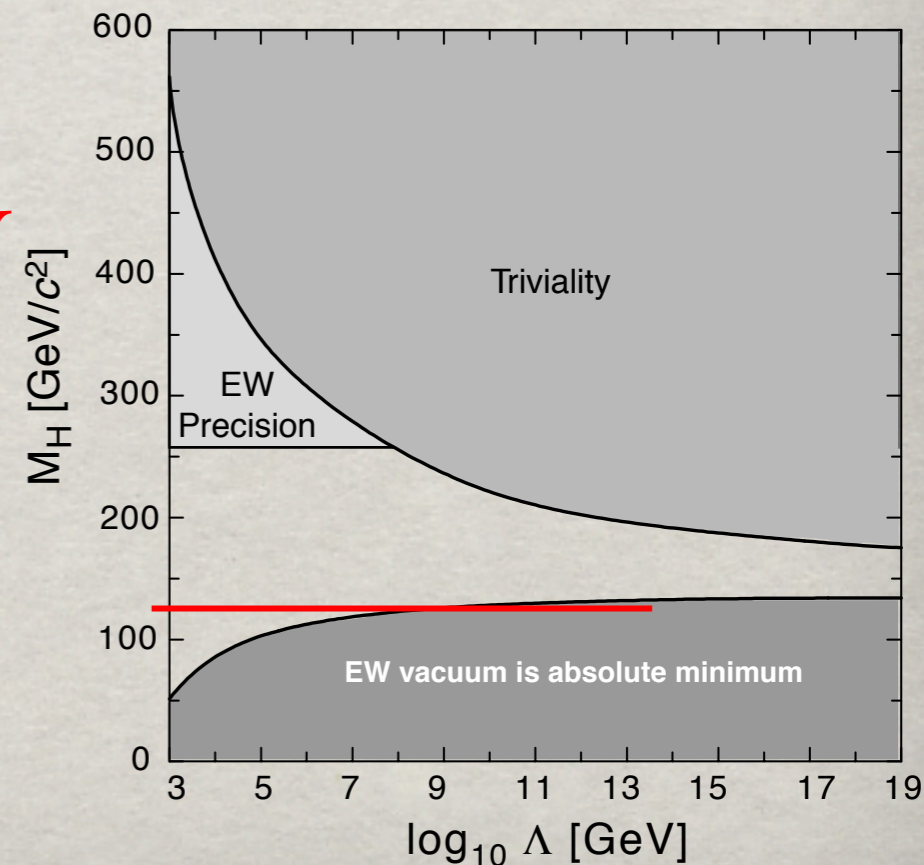


D. Goncalves, TH, S. Mukhopadhyay, arXiv:1710.02149 (PRL, 2017); arXiv:1803.09751; D. Goncalves, TH, S.C.I. Leung, H. Qin, arXiv:2012.05272; R. Abraham, D. Goncalves, TH, S.C.I. Leung, H. Qin, to appear.

The completion of the SM:
First time ever, we have a
 consistent relativistic/
 quantum mechanical theory:
 weakly coupled, unitary,
 renormalizeable, vacuum (quasi?) stable.



**Valid up to an exponentially
 high scale, perhaps to the
 Planck scale M_{Pl} !**



“... most of the grand underlying principles have been firmly established.

... the future truths of physical science are to be looked for in the sixth place of decimals. ”

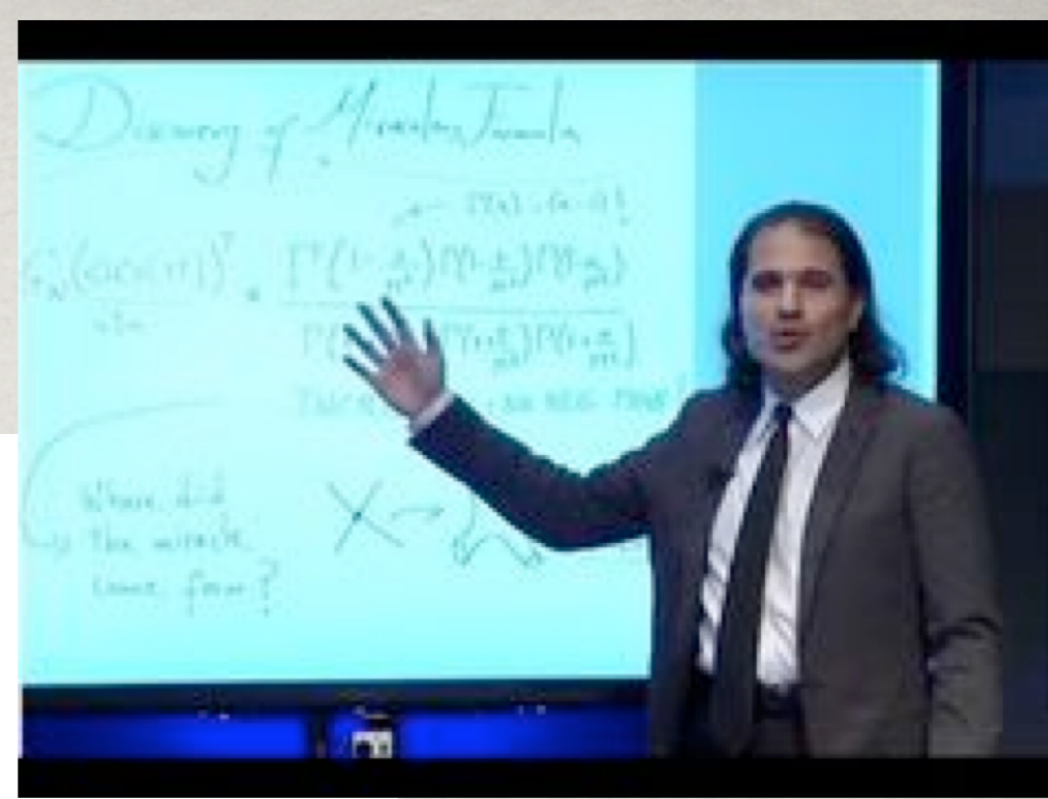
--- Albert Michelson (1894)

Michelson–Morley’s null experiment (1887):

“the moving-off point for the theoretical aspects of the second scientific revolution”

Will History repeat itself (soon)?

Nima Arkani-Hamed:



The central questions
today are not details —
but structural: origin of
spacetime, UV/IR connection,
standard model → real theory

SM AS AN EFFECTIVE FIELD THEORY

“The present educated view of the standard model, and of general relativity, is again that these are the leading terms in effective field theories.” S. Weinberg, hep-th/9702027

“We are all Wilsonians now.”

- J. Preskill, Quantum Frontier (2013)



In terms of a new physical scale Λ ,
below which the theory is valid:

$$\mathcal{L} = \sum c_i \Lambda^n \mathcal{O}_n = \underbrace{c_0 \Lambda^4 + c_2 \Lambda^2 \mathcal{O}_{\text{dim } 2} + c_3 \Lambda \mathcal{O}_{\text{dim } 3}}_{\text{(relevant operators)}} + \underbrace{c_4 \mathcal{O}_{\text{dim } 4}}_{\text{(marginal operators)}} + \underbrace{\frac{c_6}{\Lambda^2} \mathcal{O}_{\text{dim } 6} + \dots}_{\text{(irrelevant operators)}}$$

The 1st (most) “relevant operator”: $c_0\Lambda^4$

Known physics scales and the observation:

$$(M_{\text{PL}}/\Lambda_{\text{cosm}})^4 \sim 10^{120} ! \quad (\Lambda_{\text{QCD}}/\Lambda_{\text{cosm}})^4 \sim 10^{44} !$$

Wilsonian argument failed (badly)!

“... I do not understand (quantum) gravity” --- William Bardeen

The 2nd “relevant operator”: the Higgs boson mass

$$V = \underbrace{-\mu^2}_{\text{red circle}} |\phi|^2 + \lambda |\phi|^4$$

$$c_2\Lambda^2 \sim m_h^2 : \quad \lambda v^2 \sim \mu^2 \sim (100 \text{ GeV})^2 \sim (10^{-16} M_{\text{Planck}})^2$$

“... scalar particles are the only kind of free particles whose mass term does not break either an internal or a gauge symmetry.” Ken Wilson, 1970

The “Hierarchy problem” between m_h & M_{Planck}

→ Higgs as a Probe of New Physics !

Immediate target: Precision Higgs physics $v^2/\Lambda^2 \sim 2\%$

- Cross sections/branching fractions \rightarrow Higgs couplings
- Rare processes \rightarrow search new phenomena

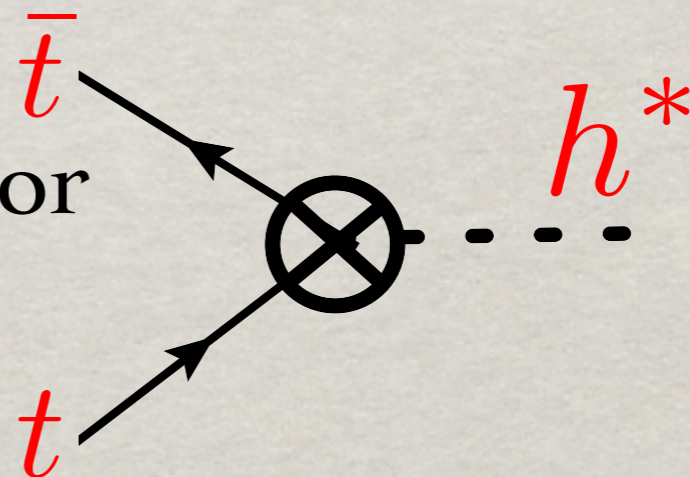
Much work/talks in this workshop ...

Complementary Approach: Higgs physics at high scales

- “Naturalness” is a UV problem
- Sensitive to new physics
- “Higgs portal” to a subtle sector.

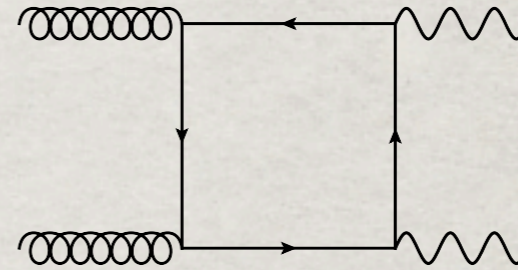
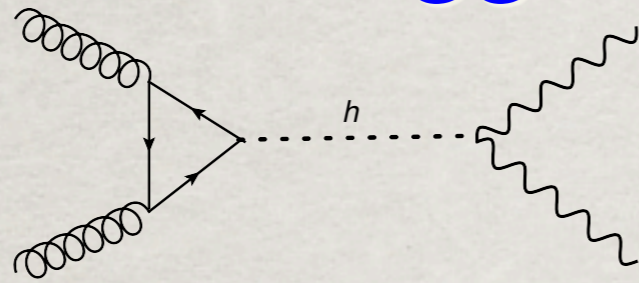
Focus on ttH coupling @ high scales:

- a. Yukawa $y_t(Q)$ RGE running
- b. EFT probe
- c. Composite form factor

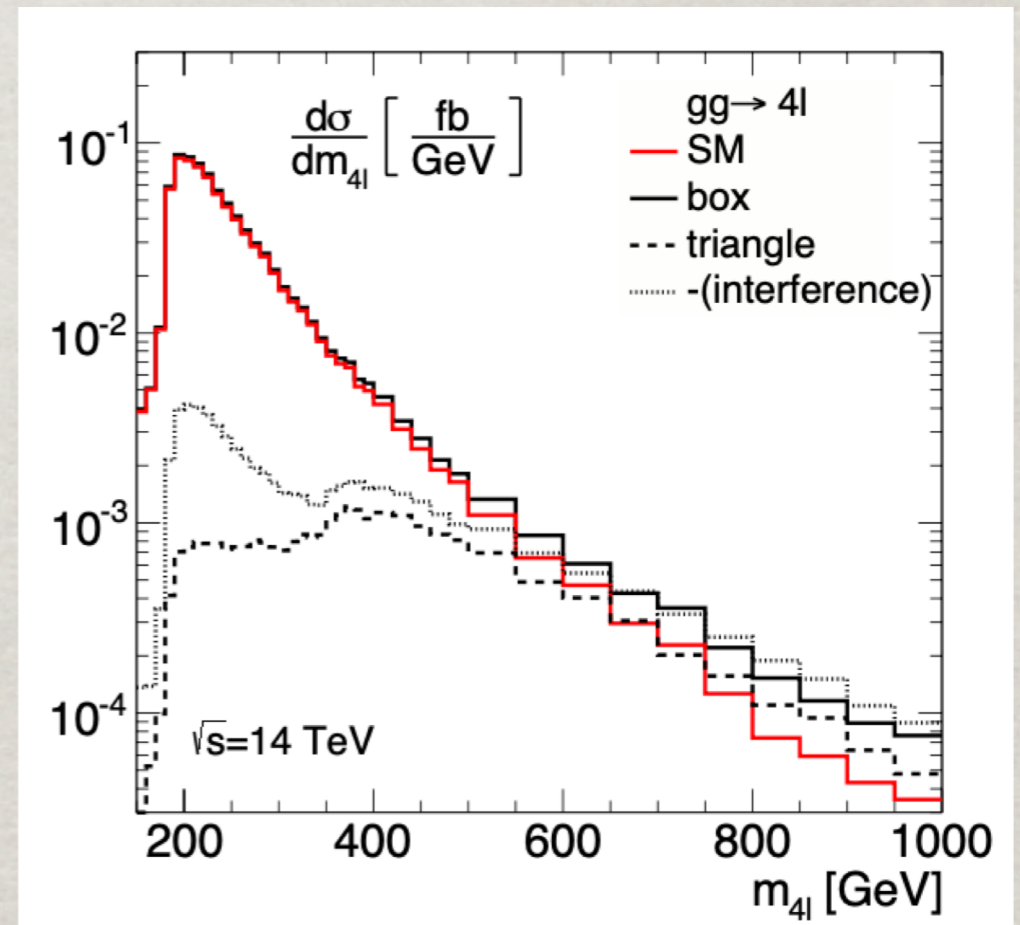
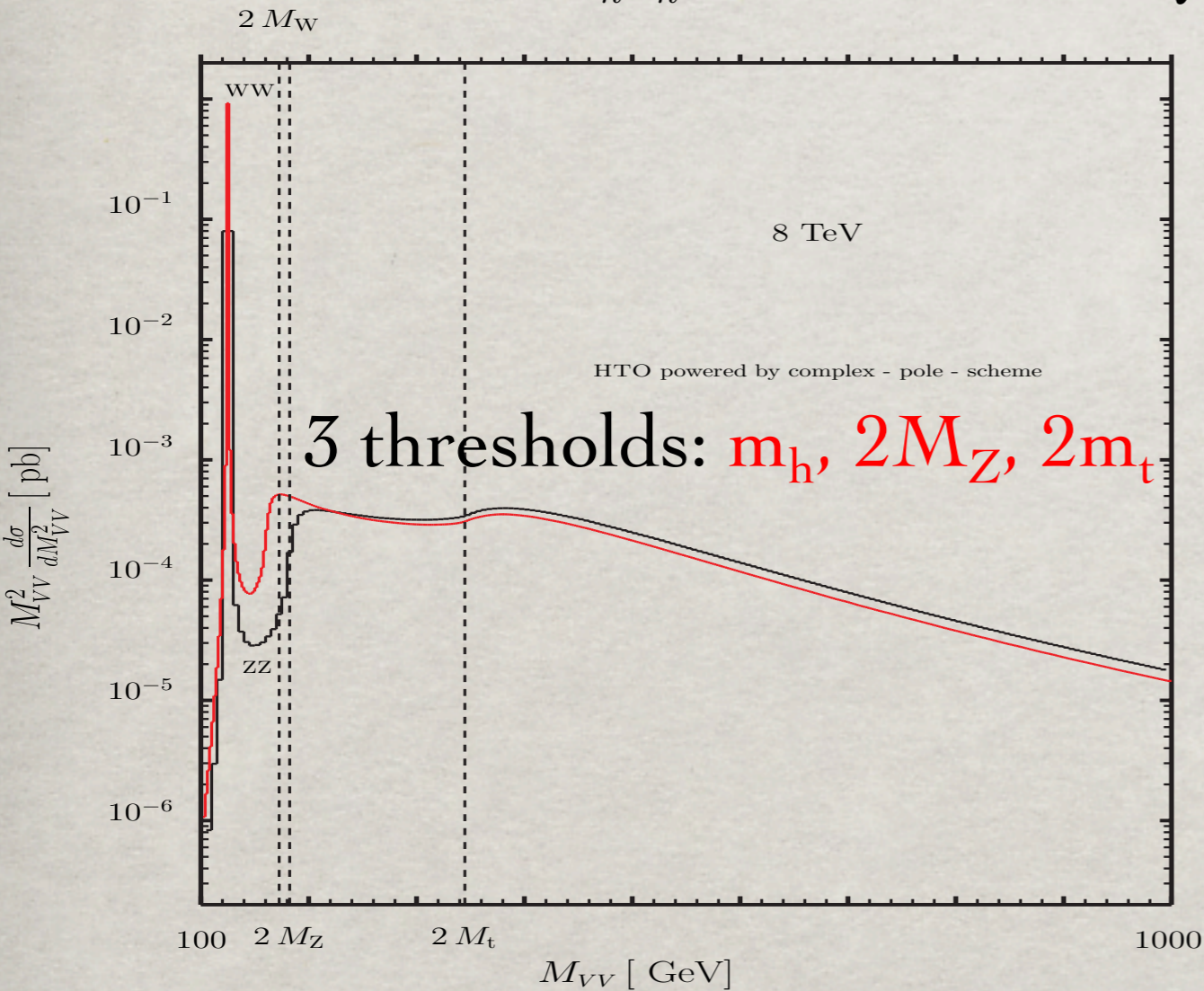


HIGH-SCALE HIGGS PHYSICS I:

far off-shell Higgs $gg \rightarrow h^* \rightarrow WW, ZZ$



$$\sigma_{\text{on}} \propto \frac{g_i^2(m_h^2)g_f^2(m_h^2)}{m_h\Gamma_h} \quad \text{and} \quad \sigma_{\text{off}} \propto \frac{g_i^2(Q^2)g_f^2(Q^2)}{Q^2}$$



Significant destructive interference between the box & triangle diagrams

N. Kauer, G. Passarino, arXiv:1206.4803
 F. Caola, K. Melnikov, arXiv:1307.4935
 Campbell, Ellis, Williams 1312.1628

a. Top quark Yukawa coupling $y_t(Q)$:

SM:
$$\frac{dy_t}{dt} = \beta_{y_t}^{\text{SM}} = \frac{y_t}{16\pi^2} \left(\frac{9}{2}y_t^2 - 8g_3^2 - \frac{9}{4}g_2^2 - \frac{17}{20}g_1^2 \right)$$

$$\beta_Q = \beta_Q^{\text{SM}} + \sum_{\text{s: massive new states}} \theta(\mu - M_s)(N_s \beta_{s,Q}^{\text{NP}})$$

MSSM:
$$\frac{dy_t}{dt} = \frac{y_t}{16\pi^2} \left(6y_t^2 - \frac{16}{3}g_3^2 - 3g_2^2 - \frac{13}{15}g_1^2 \right), \quad \text{MSSM}$$

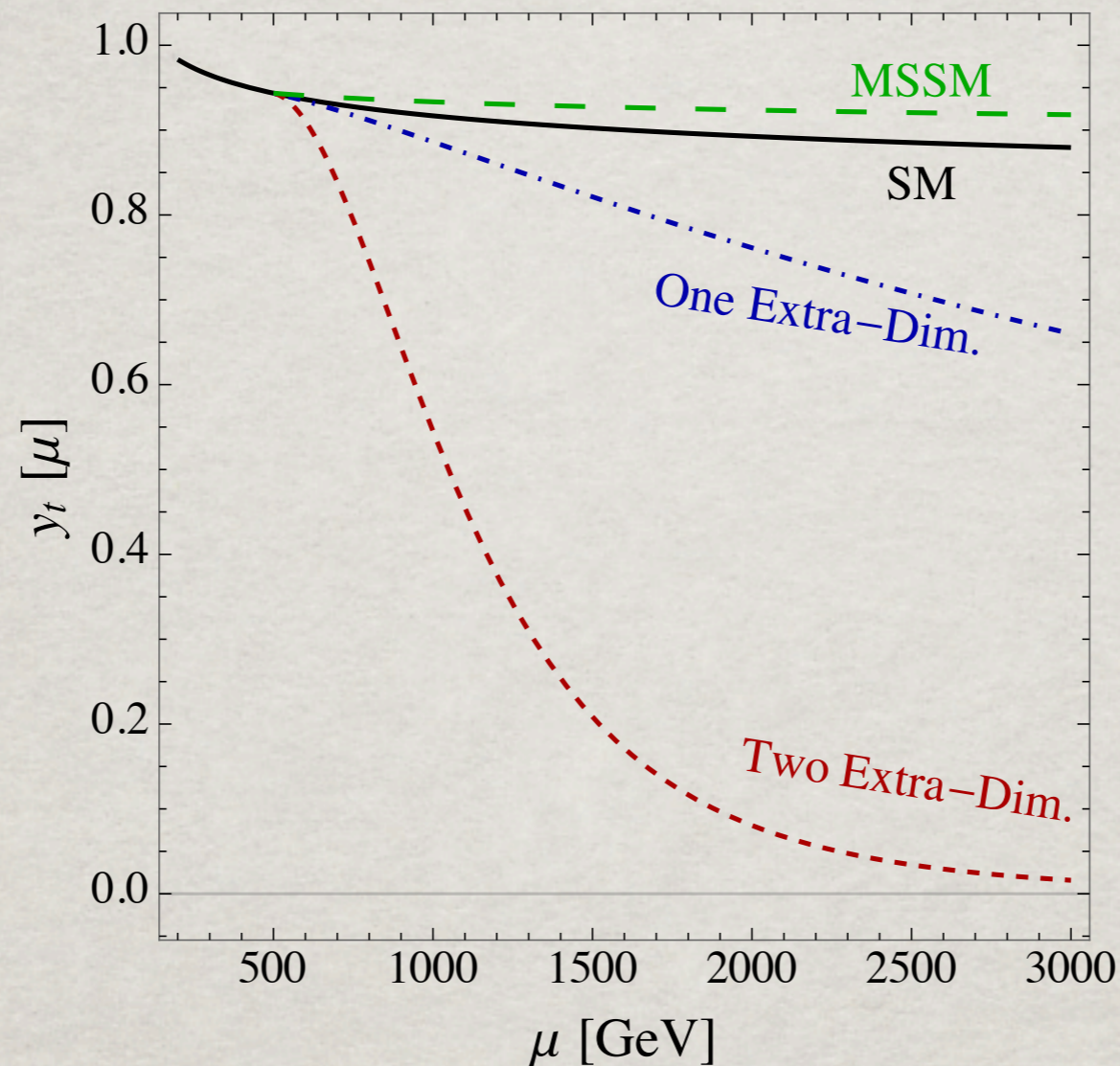
Gauge fields in extra-dimensions:

$$\frac{dy_t}{dt} = \beta_{y_t}^{\text{SM}} + \frac{y_t}{16\pi^2} 2(S(t) - 1) \left(\frac{3}{2}y_t^2 - 8g_3^2 - \frac{9}{4}g_2^2 - \frac{17}{20}g_1^2 \right), \quad \text{5D,}$$

$$\frac{dy_t}{dt} = \beta_{y_t}^{\text{SM}} + \frac{y_t}{16\pi^2} 4\pi(S(t)^2 - 1) \left(\frac{3}{2}y_t^2 - 8g_3^2 - \frac{9}{4}g_2^2 - \frac{17}{20}g_1^2 \right), \quad \text{6D.}$$

$S(t) = e^t R \sim \mu R$: counts the number of states.
 \rightarrow “volume”, power-law running!

Top quark Yukawa coupling $y_t(Q)$: RGE



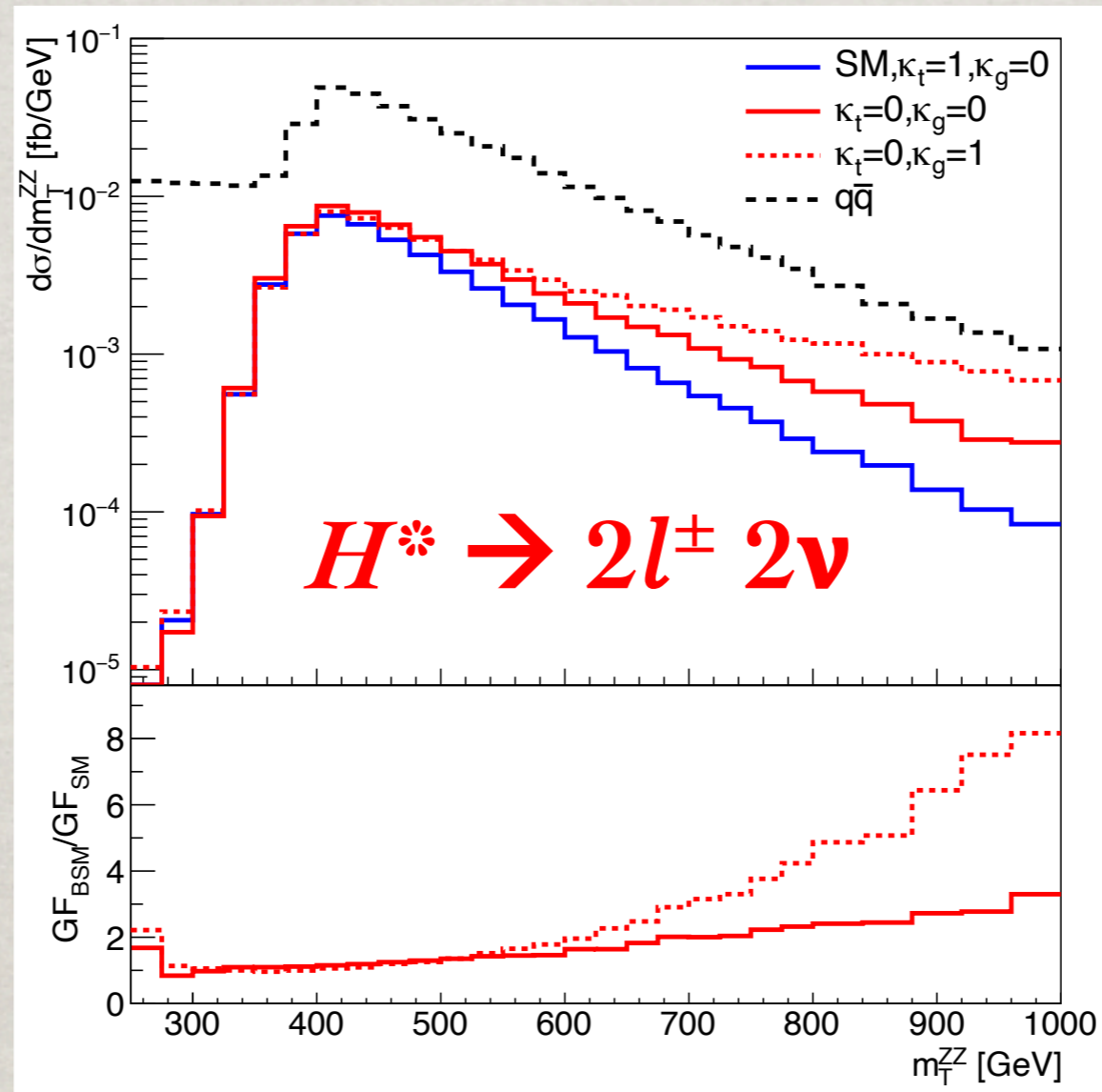
A.S. Cornell et al.,
arXiv:1209.6239.

- Suppressed Yukawa coupling \rightarrow larger ZZ signal
- HL-LHC: rather insensitive to SM/MSSM/5D
- HE-LHC/FCC-hh @ 6D \rightarrow Reach $1/R \sim 1$ TeV

D. Goncalves, TH, S. Mukhopadhyay, arXiv:1710.02149 (PRL, 2017); arXiv:1803.09751 (PRD, 2018).

b. Effective Field Theory description:

$$\mathcal{L} \supset c_g \frac{\alpha_s}{12\pi v^2} |\mathcal{H}|^2 G_{\mu\nu} G^{\mu\nu} + c_t \frac{y_t}{v^2} |\mathcal{H}|^2 \bar{Q}_L \tilde{\mathcal{H}} t_R + \text{h.c.} \quad M_{\text{BSM}} \sim M_{\text{SM}} (1 + m_{ZZ}^2 / \Lambda^2)$$



| | Γ_H / Γ_H^{SM} | Λ_{EFT} |
|---|----------------------------|-----------------|
| $H^* \rightarrow ZZ \rightarrow ll\nu\nu$ | 1.31 | 0.8 TeV |
| $H^* \rightarrow ZZ \rightarrow 4l$ | 1.3 (68% CL) [33] | 0.55 TeV [34] |

D. Goncalves, TH, S. Mukhopadhyay, arXiv:1710.02149 (PRL, 2017); arXiv:1803.09751 (PRD, 2018).
D. Goncalves, TH, S.C.I. Leung, H. Qin, arXiv:2012.05272.

c. The Momentum-dependent Form Factor:

Deviation from point-like interactions:

$$V_{ttH}(p^\mu, \bar{p}^\mu) = \frac{\sqrt{2}m_t}{v} \Gamma(p^2/\Lambda_c^2, \bar{p}^2/\Lambda_c^2, q^2/\Lambda_c^2)$$

Current 95% CL bound from the LHC Higgs signal:

$$|\Gamma(m_h^2/\Lambda^2)^2 - 1| < 0.1$$

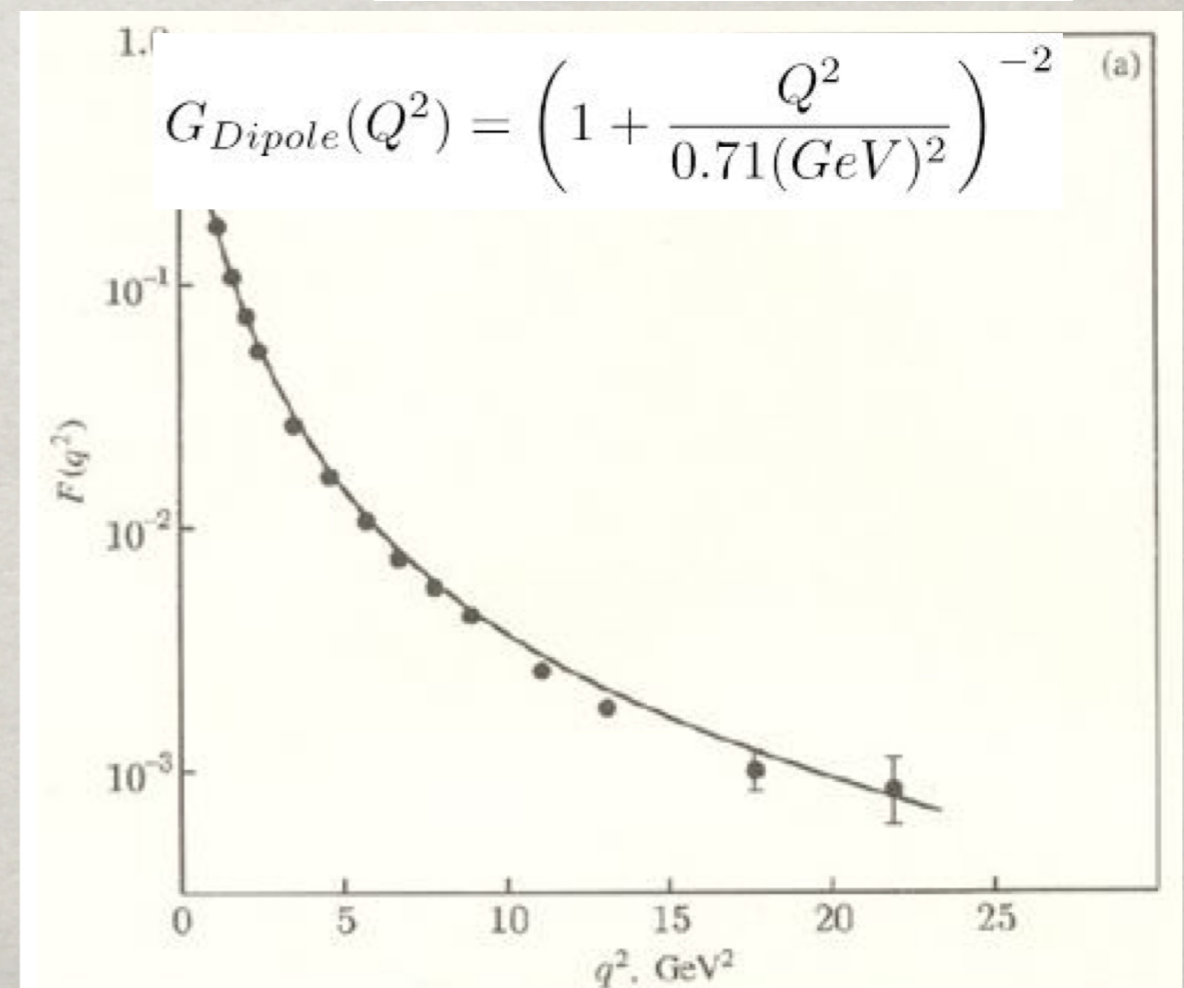
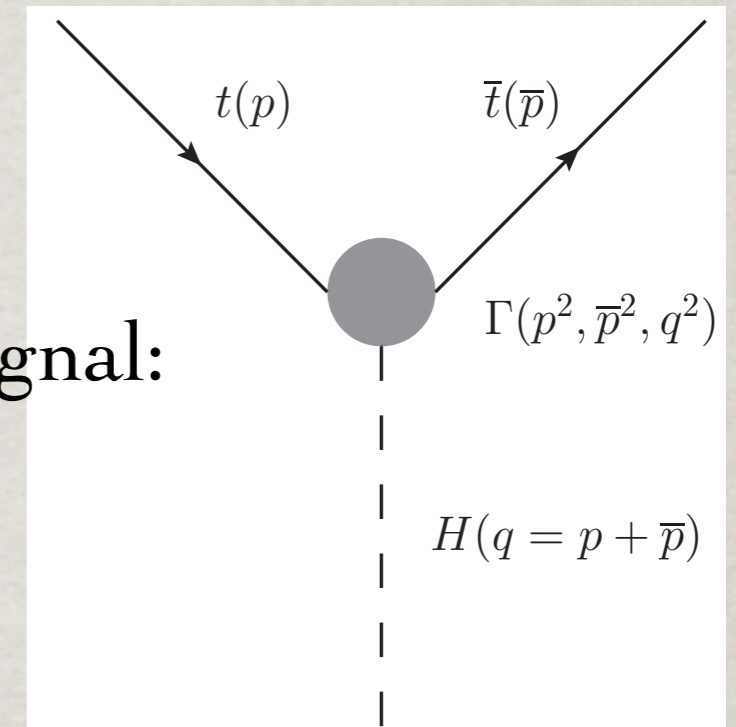
Analogue: Nucleon form factor

$$\Gamma(q^2/\Lambda_c^2) = \frac{1}{(1 + q^2/\Lambda_c^2)^n}$$

$n=2 \rightarrow$ “Dipole FF”

1-particle exchange potential;
a spatial exponential distribution.

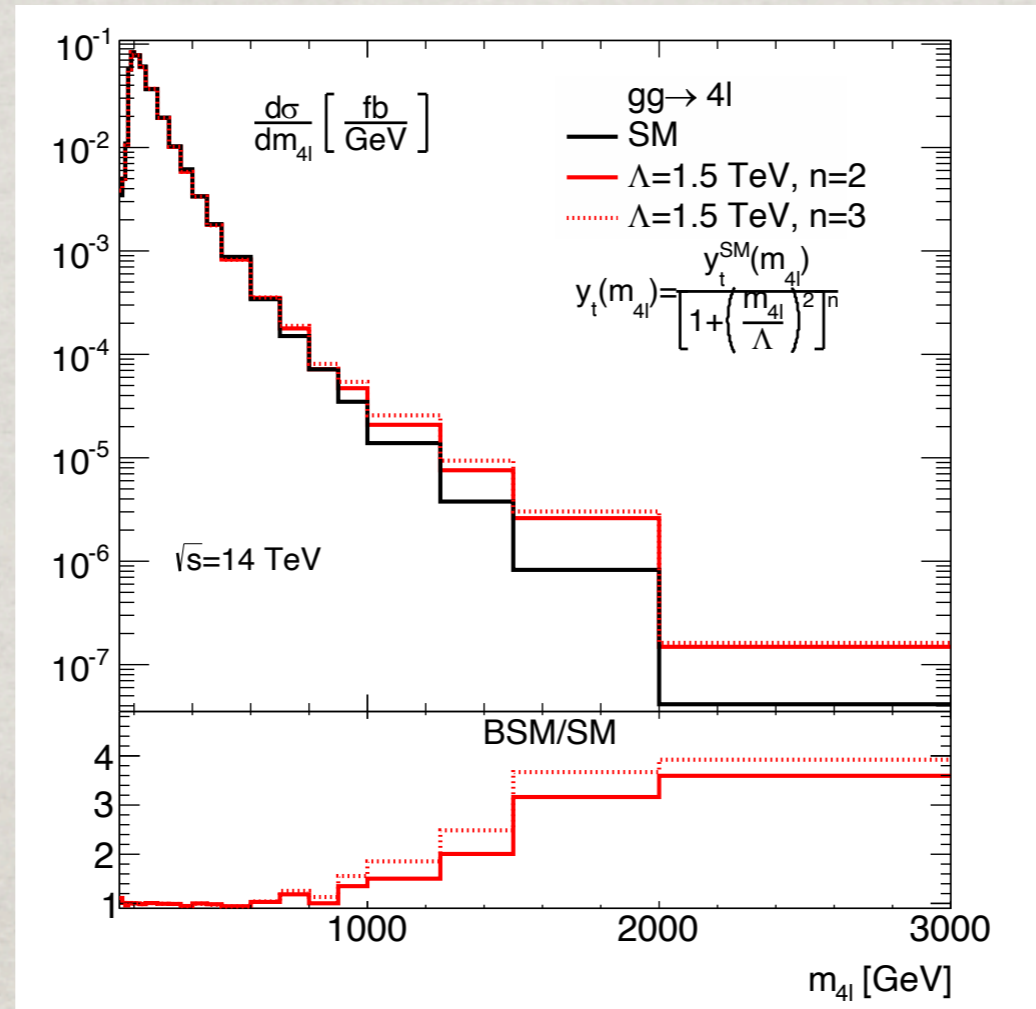
Leading to a suppressed ttH
 \rightarrow Enhanced ZZ signal!



LHC distribution: top-Higgs Form Factor

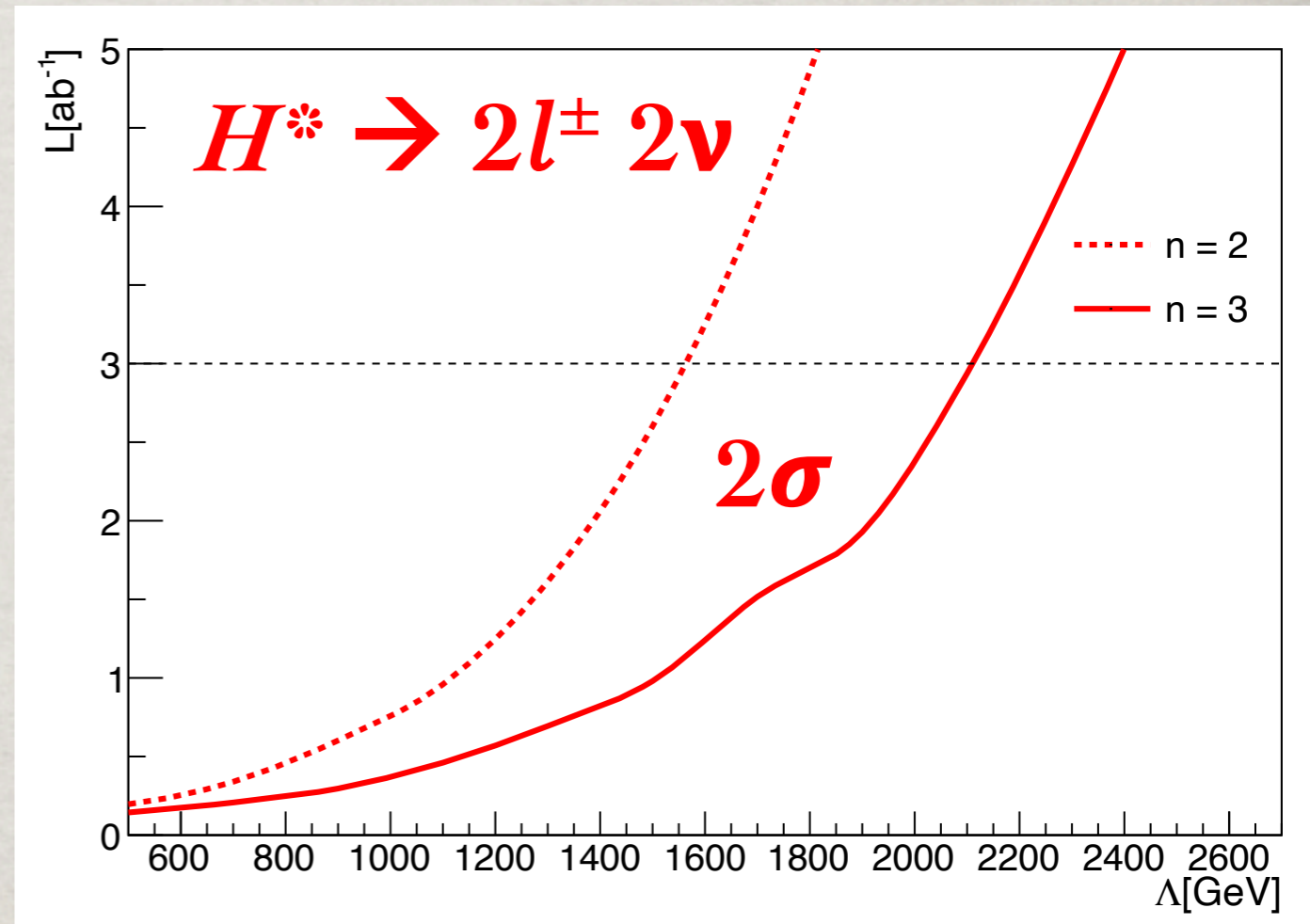
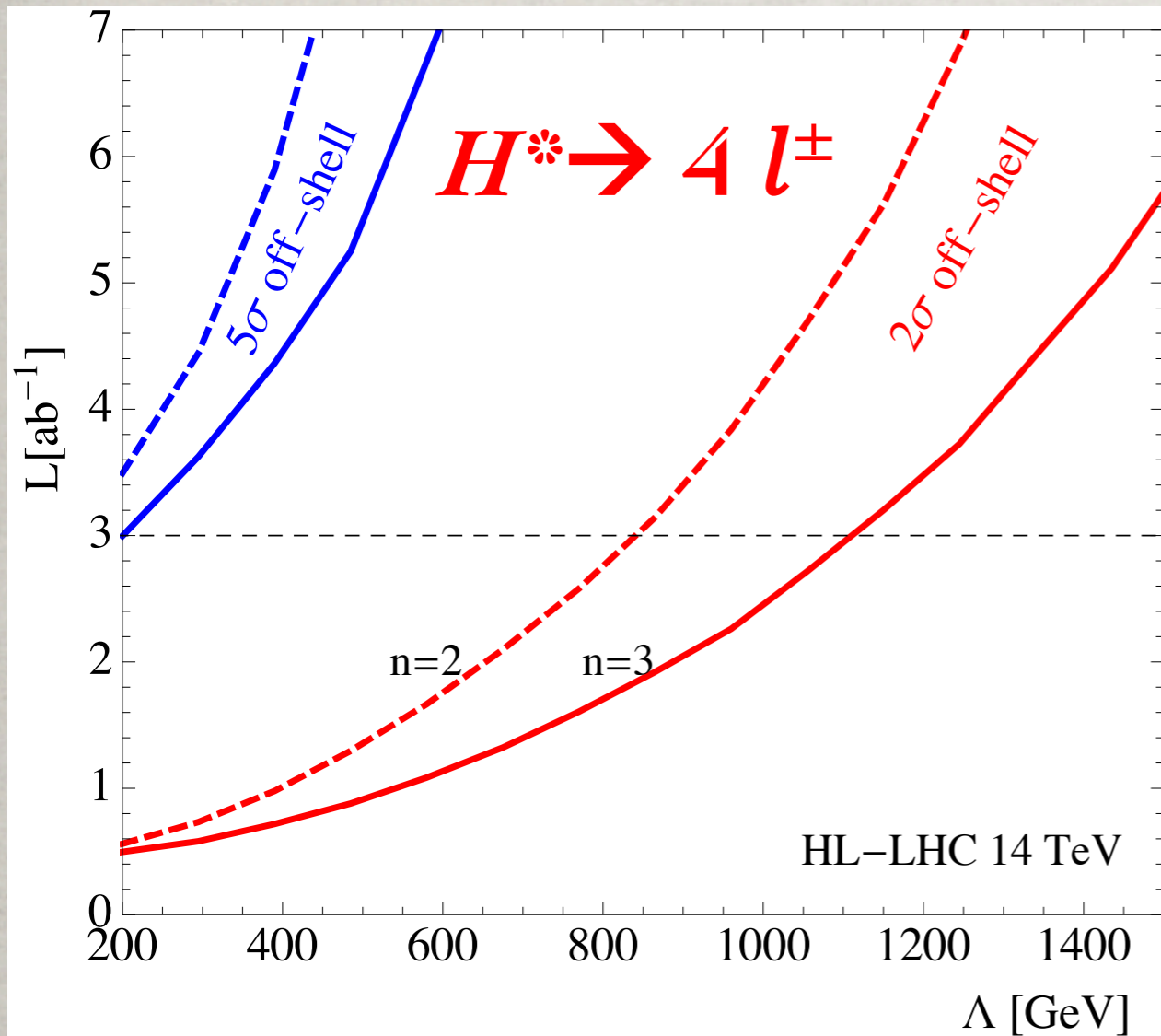
(similar effects for $n=2,3$)

$$H^* \rightarrow 4 l^\pm$$



- Form Factor suppression
- smaller off-shell Higgs signal
 - weaker interference
 - larger ZZ signal: a factor of 3-4!

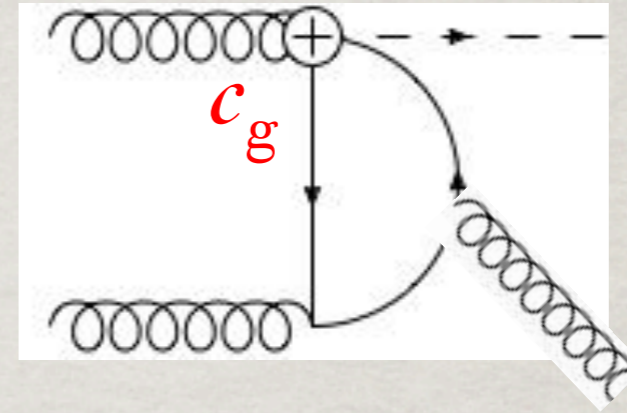
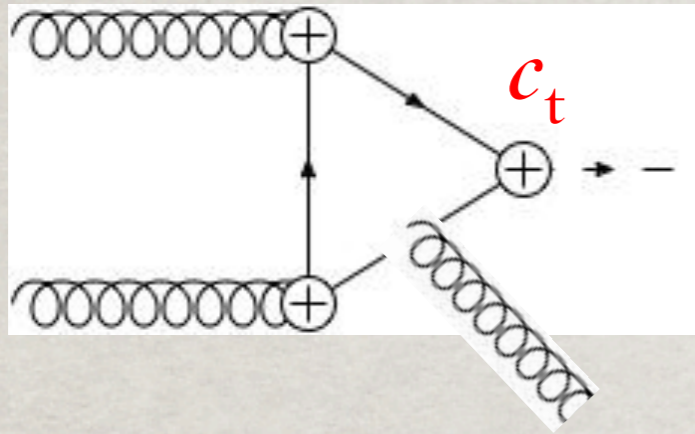
LHC Sensitivity: top-Higgs Form Factor



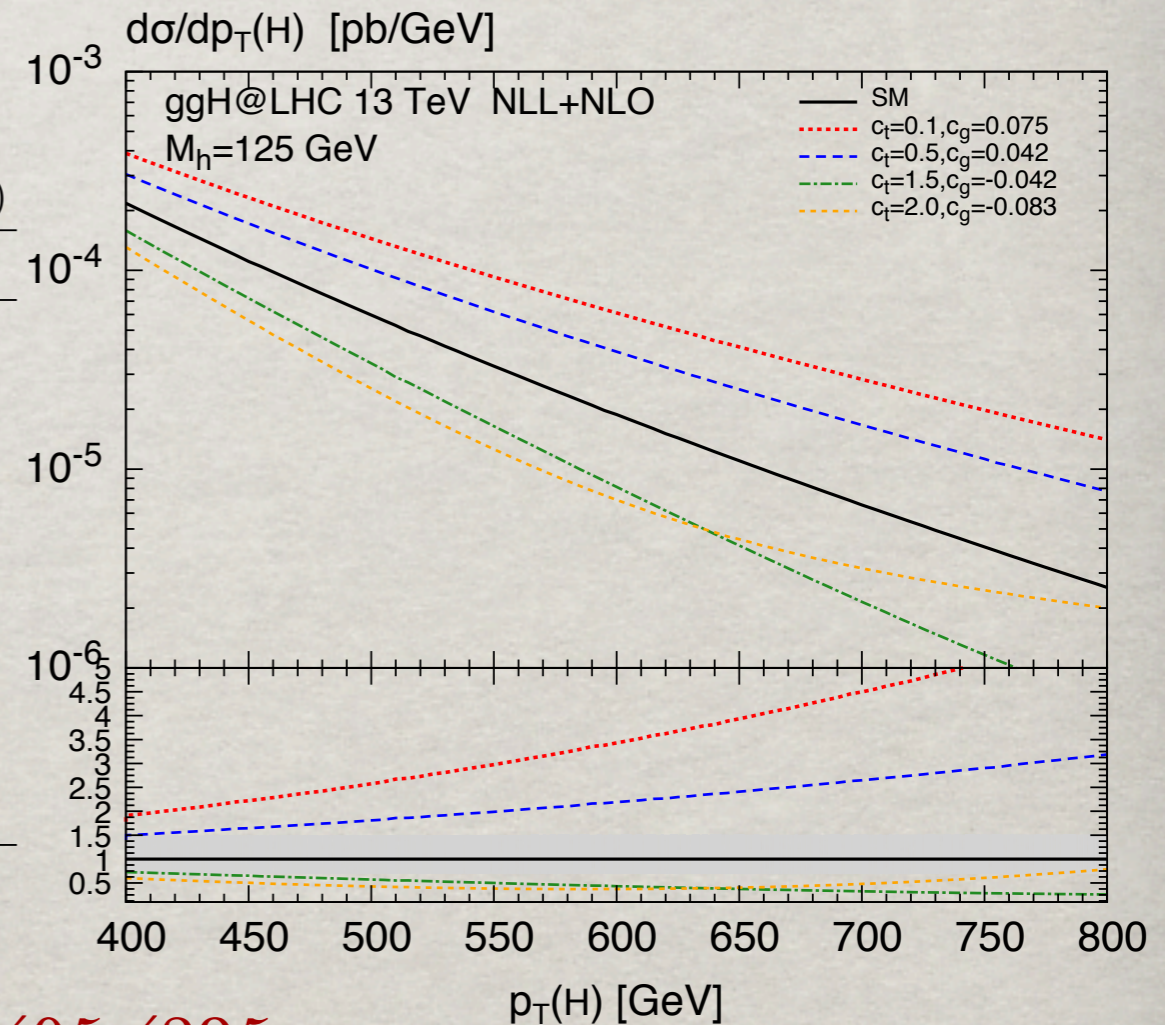
| | Γ_H / Γ_H^{SM} | Λ_{EFT} | $\Lambda_{Composite}^{n=2}$ |
|---|----------------------------|-----------------|-----------------------------|
| $H^* \rightarrow ZZ \rightarrow ll\nu\nu$ | 1.31 | 0.8 TeV | 1.5 TeV |
| $H^* \rightarrow ZZ \rightarrow 4l$ | 1.3 (68% CL) [33] | 0.55 TeV [34] | 0.8 TeV [18] |

HIGH-SCALE HIGGS PHYSICS II

Higgs @ high- P_T in $gg \rightarrow H+j$



| | (0.5,0.5) | (0.7,0.3) | (0.9,0.1) | SM | (1.1,-0.1) | (1.3,-0.3) | (1.5,-0.5) |
|---------------------|-----------|-----------|-----------|--------|------------|------------|------------|
| $p_{T,H} > 10$ GeV | 0.9733 | 0.9838 | 0.9945 | 1.0000 | 1.0055 | 1.0167 | 1.0281 |
| $p_{T,H} > 100$ GeV | 1.0044 | 1.0012 | 0.9999 | 1.0000 | 1.0006 | 1.0031 | 1.0076 |
| $p_{T,H} > 200$ GeV | 1.1166 | 1.0646 | 1.0198 | 1.0000 | 0.9820 | 0.9513 | 0.9277 |
| $p_{T,H} > 300$ GeV | 1.3450 | 1.1921 | 1.0591 | 1.0000 | 0.9459 | 0.8526 | 0.7791 |
| $p_{T,H} > 400$ GeV | 1.6531 | 1.3590 | 1.1087 | 1.0000 | 0.9023 | 0.7397 | 0.6210 |
| $p_{T,H} > 500$ GeV | 2.0233 | 1.5520 | 1.1633 | 1.0000 | 0.8573 | 0.6340 | 0.4932 |
| $p_{T,H} > 600$ GeV | 2.4869 | 1.7871 | 1.2274 | 1.0000 | 0.8076 | 0.5279 | 0.3882 |
| $p_{T,H} > 700$ GeV | 3.1213 | 2.1003 | 1.3093 | 1.0000 | 0.7482 | 0.4172 | 0.3161 |
| $p_{T,H} > 800$ GeV | 3.7427 | 2.3989 | 1.3841 | 1.0000 | 0.6981 | 0.3411 | 0.3129 |

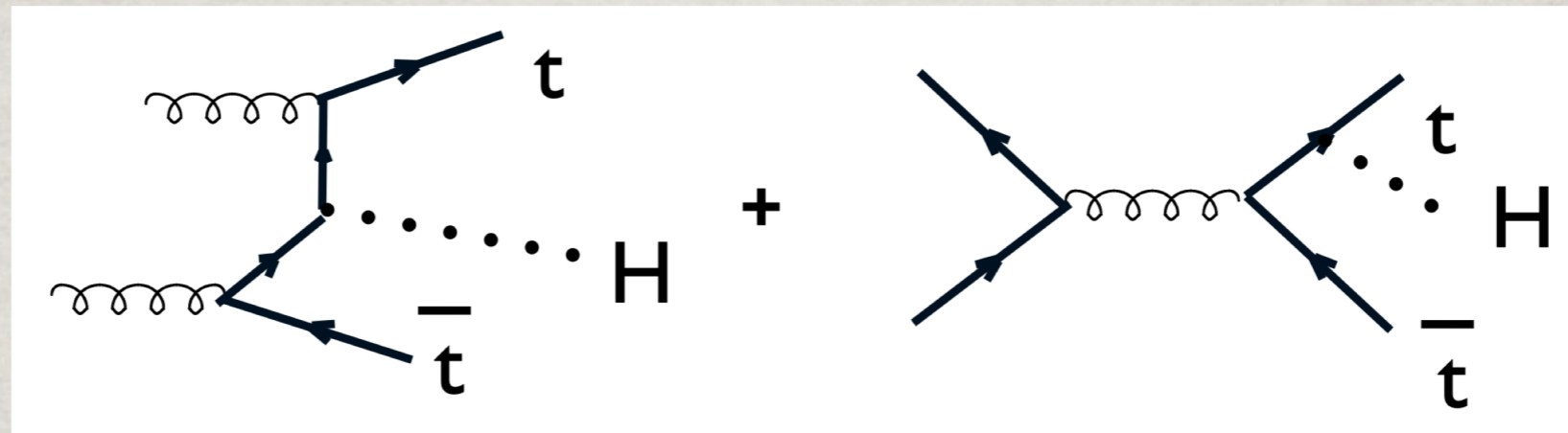


M. Schlaffer, M. Spannowsky et al., arXiv:1405.4295;

M. Grazzini, A. Ilnicka, M. Spira, M. Wiesemann, arXiv:1612.00283 ...

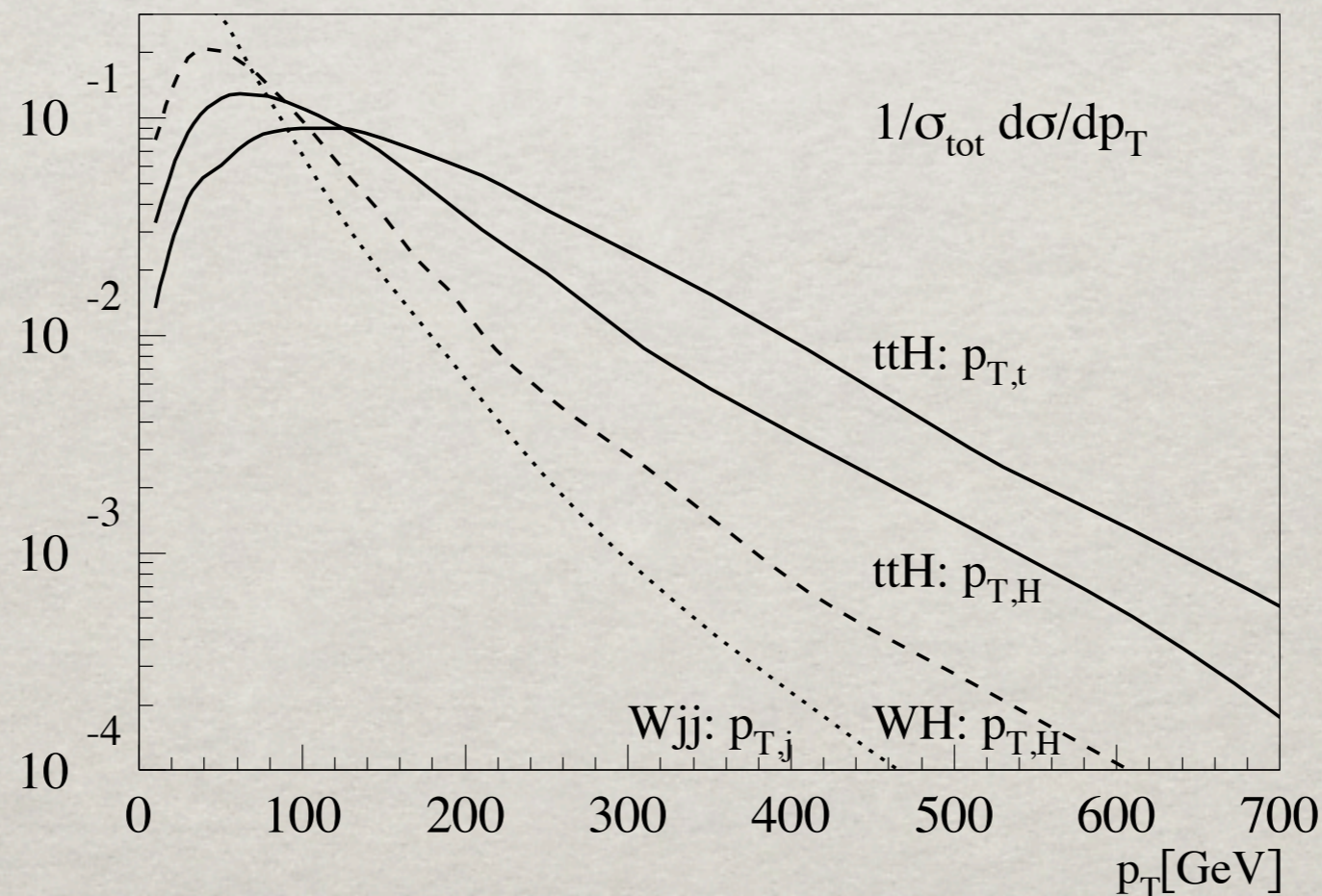
HIGH-SCALE HIGGS PHYSICS III

Higgs @ high- P_T in $gg \rightarrow ttH$



Event rate for $gg \rightarrow ttH$ comparable to $gg \rightarrow b$ at high p_T
at HL-LHC / FCC-hh

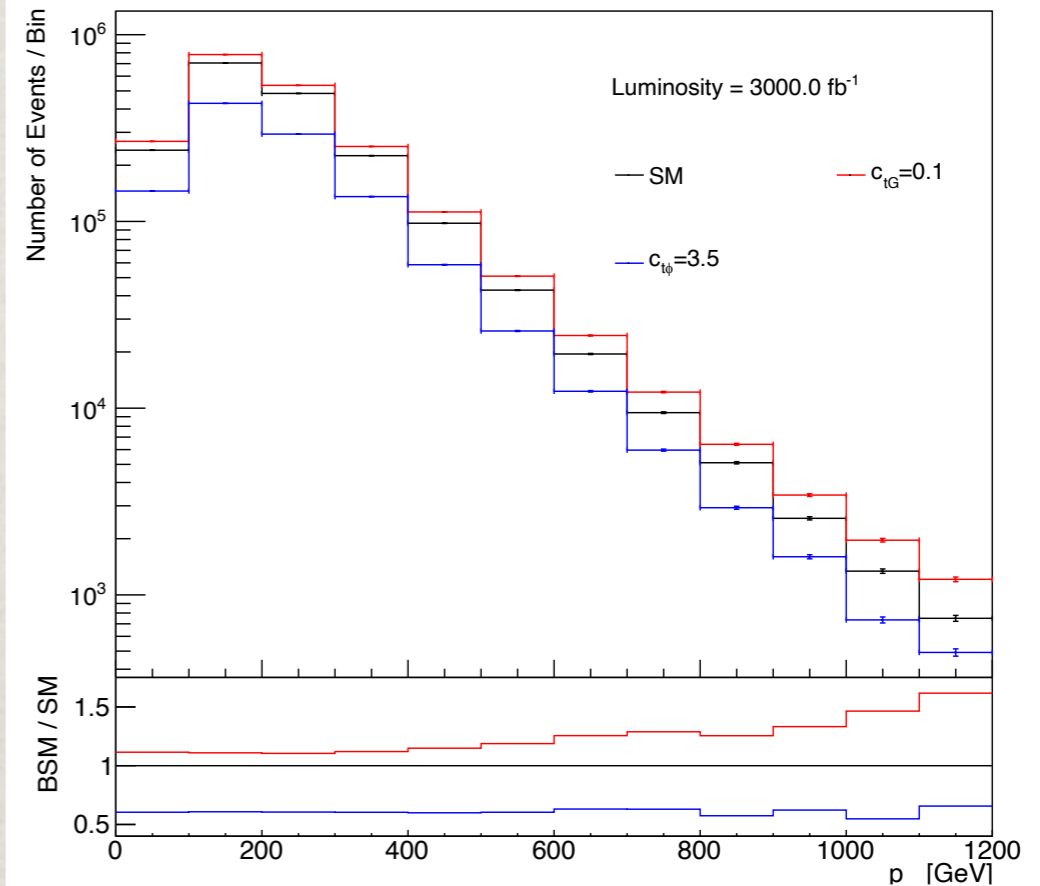
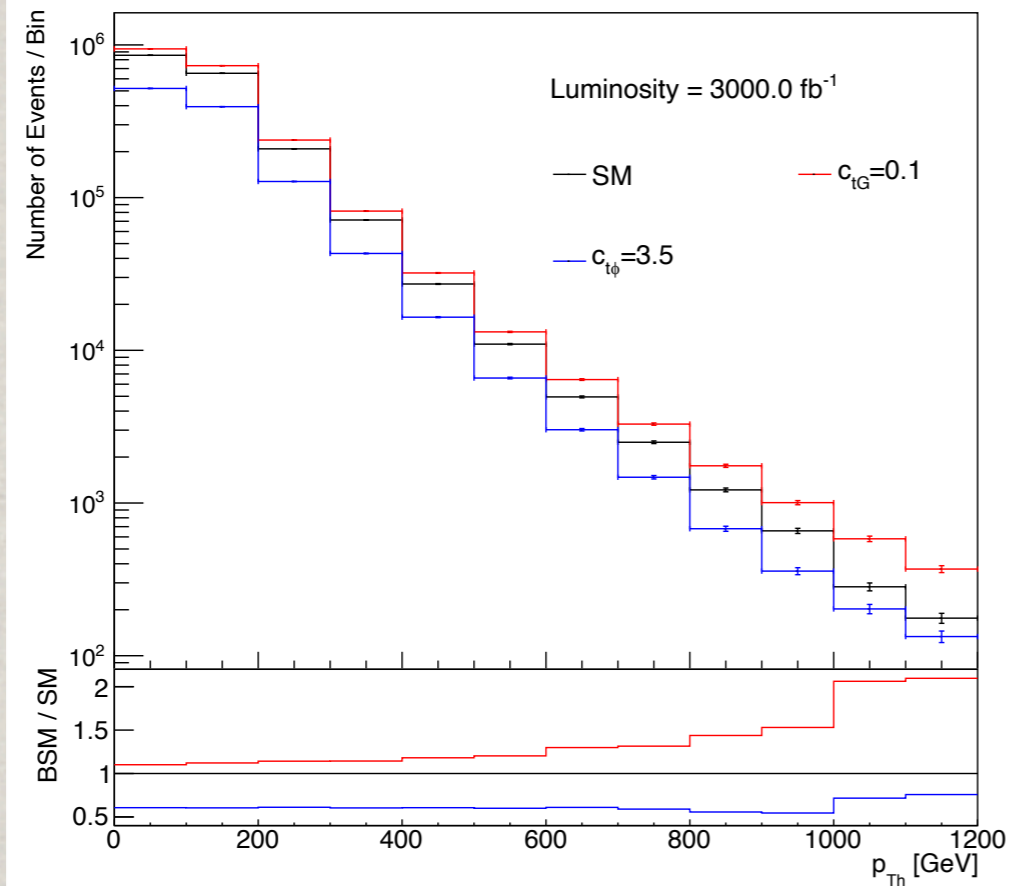
M. Mangano, 2018
FCC Report



T. Plehn, G. Salam, M₁₆Spannowsky, arXiv:0910.5472

$$\mathcal{O}_{t\phi} = (H^\dagger H)(\bar{Q}t)\tilde{H}$$

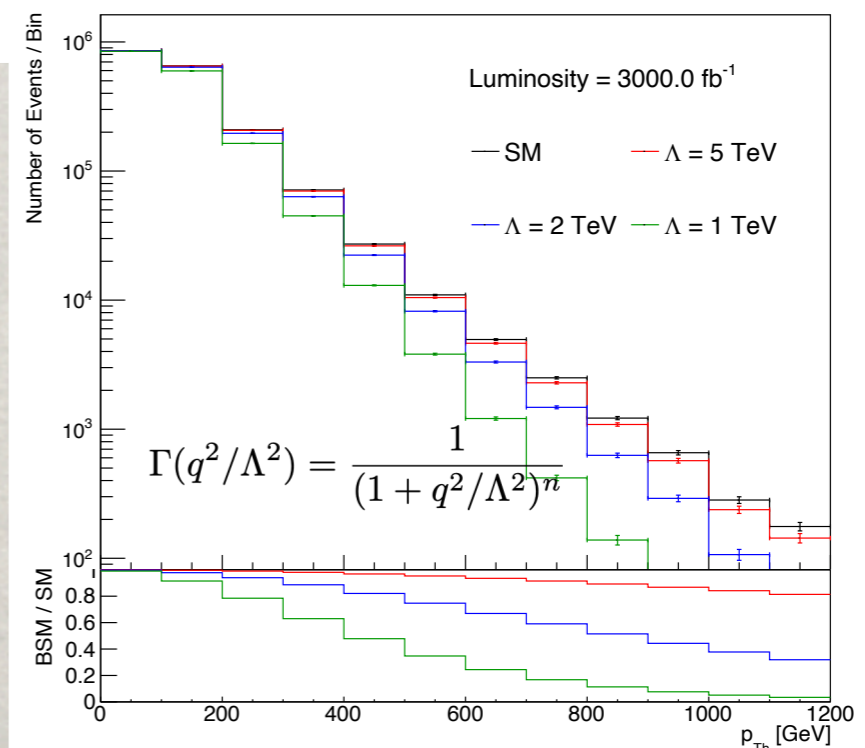
$$\mathcal{O}_{tG} = (\bar{Q}\sigma^{\mu\nu}T^A t)\tilde{H}G_{\mu\nu}^A$$



gg → *ttH*

Direct probe of space-like regime
(no BR's & cuts)

R. Abraham, D. Goncalves, TH,
S.C.I. Leung, H. Qin, to appear.



Summary:

Under the Higgs lamppost
→ HPNP



Prevision Higgs measurements → 1% for $\Lambda \sim 1 \text{ TeV} - 2 \text{ TeV}$

Strong motivation to target on ttH coupling @ High scales

- Weakly coupled: RGE for SM; MSSM: weak effects.
- Extra dimensions → Power-law running!
- Strong dynamics / composite Higgs / EFT / Form Factor
- New states coming into Higgs propagation:
scalar singlets; a continuum spectrum ...

Modifications may be observable at the HL-LHC/FCC
 $2\sigma - 5\sigma$ level sensitivity.

Exciting journey on HPNP !