

Scalar and fermion on-shell amplitudes in generalized Higgs effective field theory

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Based on **1904.07618** (PRD)
and **2102.08519**

in collaboration with

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Big questions

- SM cannot explain
 - Why EWSB scale $\sim O(100 \text{ GeV})$
 - What dark matter is
 - \vdots
- Physics beyond SM is needed.

Searches for BSM signals

- Direct searches
- Indirect searches

How to organize BSM signals ?

Can we be systematic ?

Higgs EFT

- BSM effects in the SM processes can be described by

$$\begin{aligned}\mathcal{L}_{\text{HEFT}} = & \frac{1}{2}(\partial_\mu h)(\partial^\mu h) + i\psi^\dagger \bar{\sigma}^\mu \partial_\mu \psi \\ & + \frac{v^2}{4} F(h) \text{Tr}[\partial_\mu U^\dagger \partial^\mu U] - V(h) \\ & + M(h) \psi_L^\dagger U \psi_R + h.c.\end{aligned}$$

Feruglio (1993), ++

- F , V , M parameterize BSM contributions.

Beyond HEFT ?

- We **cannot** compute production/decay rate of BSM particles in HEFT framework.
- HEFT cannot be applied to direct searches.

How can we extend HEFT to include BSM particles?

Talk plan

- Adding spin-0
1904.07618 (PRD)
- Adding spin-1/2
2102.08519

Scalar sector of Higgs EFT

- Scalar sector of HEFT

$$\mathcal{L}_{\text{HEFT}} = \frac{1}{2}(\partial_\mu h)(\partial^\mu h) + \frac{v^2}{4}F(h) \text{Tr}[\partial_\mu U^\dagger \partial^\mu U] - V(h)$$

\uparrow 125GeV higgs \uparrow NGBs (π^1, π^2, π^3)

Scalar sector of Higgs EFT

- Scalar sector of HEFT

$$\mathcal{L}_{\text{HEFT}} = \frac{1}{2}(\partial_\mu h)(\partial^\mu h) + \frac{v^2}{4}F(h) \text{Tr}[\partial_\mu U^\dagger \partial^\mu U] - V(h)$$

- Geometrical form $\phi = (h, \pi^1, \pi^2, \pi^3)$

$$\mathcal{L}_{\text{HEFT}} = \frac{1}{2}g_{ij}(\phi) \partial_\mu \phi^i \partial^\mu \phi^j - V(\phi)$$

Alonso, Jenkins, Manohar
1511.00724, 1605.03602

Adding spin-0

- Generalization with spin-0 particles:

$$\mathcal{L}_{\text{HEFT}} = \frac{1}{2} g_{ij}(\phi) \partial_{\mu} \phi^i \partial^{\mu} \phi^j - V(\phi)$$

$$\phi = (h, \pi^1, \pi^2, \pi^3, S, H^+, H^-, \dots)$$

- The “charged” scalar manifold can be constructed by Callen-Coleman-Wess-Zumino (CCWZ) method.

Nagai, Tanabashi, Tsumura, Uchida,
1904.07618 (PRD)

Talk plan

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Fermionic extension

$$\mathcal{L}_{\text{HEFT}} = \frac{1}{2} g_{ij}(\phi) \partial_\mu \phi^i \partial^\mu \phi^j - V(\phi)$$

- SUSY !

$$\phi^i \quad \longrightarrow \quad (\phi^i, \phi^{\dagger i*}, \psi^i, \psi^{\dagger i*})$$

(Weyl) fermion
↙ ↘

$$g_{ij}(\phi) \quad \longrightarrow \quad g_{ij^*}(\phi, \phi^\dagger) \quad \text{Kähler metric}$$

$$V(\phi) \quad \longrightarrow \quad P(\phi) \quad P^\dagger(\phi^\dagger) \quad \text{Superpotential}$$

SUSY HEFT

$$\mathcal{L}_{\text{scalar}} = g_{ij^*} (\partial_\mu \phi^i) (\partial^\mu \phi^{\dagger j^*}) - g^{ij^*} P_{,i} P_{,j^*}^\dagger$$

$$\begin{aligned} \mathcal{L}_{\text{fermion}} &= \frac{i}{2} g_{ij^*} \psi^{\dagger j^*} \bar{\sigma}^\mu \overset{\leftrightarrow}{\partial}_\mu \psi^i \\ &+ \frac{i}{2} (g_{ij^*,k} \partial_\mu \phi^k - g_{ij^*,k^*} \partial_\mu \phi^{\dagger k^*}) (\psi^{\dagger j^*} \bar{\sigma}^\mu \psi^i) \\ &- P_{;ij} (\psi^i \psi^j) + h.c. \\ &- \frac{1}{8} R_{ik^*jl^*} (\psi^i \psi^j) (\psi^{\dagger i^*} \psi^{\dagger j^*}) \end{aligned}$$

Generalized HEFT

$$\mathcal{L}_{\text{scalar}} = \frac{1}{2} g_{ij} (\partial_\mu \phi^i) (\partial^\mu \phi^j) - V(\phi)$$

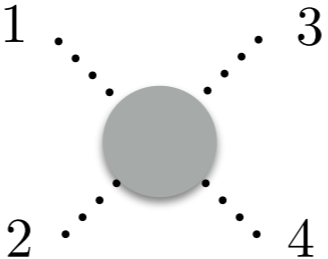
$$\begin{aligned} \mathcal{L}_{\text{fermion}} = & \frac{i}{2} g_{\hat{i}\hat{j}^*} \psi^{\dagger\hat{j}^*} \bar{\sigma}^\mu \overleftrightarrow{\partial}_\mu \psi^{\hat{i}} \\ & + \frac{i}{2} (g_{k\hat{j}^*,\hat{i}} - g_{k\hat{i},\hat{j}^*}) (\psi^{\dagger\hat{j}^*} \bar{\sigma}^\mu \psi^{\hat{i}}) (\partial_\mu \phi^k) \\ & - \frac{1}{2} M_{\hat{i}\hat{j}} (\psi^{\hat{i}} \psi^{\hat{j}}) + h.c. \\ & - \frac{1}{4} R_{\hat{i}\hat{k}^*\hat{j}\hat{l}^*} (\psi^{\hat{i}} \psi^{\hat{j}}) (\psi^{\dagger\hat{k}^*} \psi^{\dagger\hat{l}^*}) \\ & - \frac{1}{12} R_{\hat{i}\hat{l}\hat{j}\hat{k}} (\psi^{\hat{i}} \psi^{\hat{j}}) (\psi^{\hat{k}} \psi^{\hat{l}}) + h.c. \end{aligned}$$

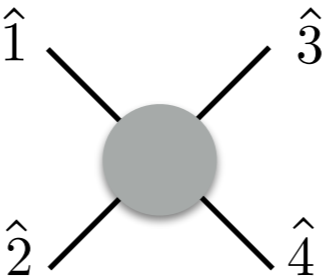
NEW!

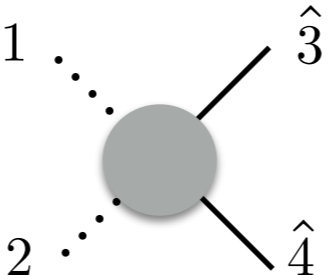
- can be constructed by CCWZ method. **Nagai**, Tanabashi, Tsumura, Uchida, 2102.08519

On-shell amplitudes

- High-energy on-shell amplitudes probe the “curvature”.

SSSS  $\sim R_{1234} s$

FFFF  $\sim R_{\hat{1}\hat{2}\hat{3}\hat{4}} s$

SSFF  $\sim R_{12\hat{3}\hat{4}} s$

* See [2102.08519](#) for complete expressions.

Summary

- We have generalized HEFT by extending the field space.

$$\begin{aligned}\mathcal{L}_{\text{scalar}} &= \frac{1}{2}g_{ij}(\partial_\mu\phi^i)(\partial^\mu\phi^j) - V(\phi) \\ \mathcal{L}_{\text{fermion}} &= \frac{i}{2}g_{\hat{i}\hat{j}^*}\psi^{\dagger\hat{j}^*}\bar{\sigma}^\mu\overset{\leftrightarrow}{\partial}_\mu\psi^{\hat{i}} \\ &\quad + \frac{i}{2}(g_{k\hat{j}^*,\hat{i}} - g_{k\hat{i},\hat{j}^*})(\psi^{\dagger\hat{j}^*}\bar{\sigma}^\mu\psi^{\hat{i}})(\partial_\mu\phi^k) \\ &\quad - \frac{1}{2}M_{\hat{i}\hat{j}}(\psi^{\hat{i}}\psi^{\hat{j}}) + h.c. \\ &\quad - \frac{1}{4}R_{\hat{i}\hat{k}^*\hat{j}\hat{l}^*}(\psi^{\hat{i}}\psi^{\hat{j}})(\psi^{\dagger\hat{k}^*}\psi^{\dagger\hat{l}^*}) \\ &\quad - \frac{1}{12}R_{\hat{i}\hat{l}\hat{j}\hat{k}}(\psi^{\hat{i}}\psi^{\hat{j}})(\psi^{\hat{k}}\psi^{\hat{l}}) + h.c.\end{aligned}$$

- On-shell amplitudes probes the “geometry”.