

Vector DM from hidden local SU(2)

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(arXiv:2012.11377)

1. Introduction

Existence and nature of DM is mystery in physics

DM as new particle is one of the attractive scenario

- ✧ **No electric charge, no color, Non-baryonic**
- ✧ **Weakly interacting**
- ✧ **Stable in cosmological scale**

A DM model with hidden gauge symmetry is interesting

- **DM stability from hidden gauge symmetry**
- **Dark gauge symmetry provide DM or Z' or dark photon**
- **Vector DM + Z' mediator from non-Abelian case**

Natural resonant annihilation is realized for vector DM

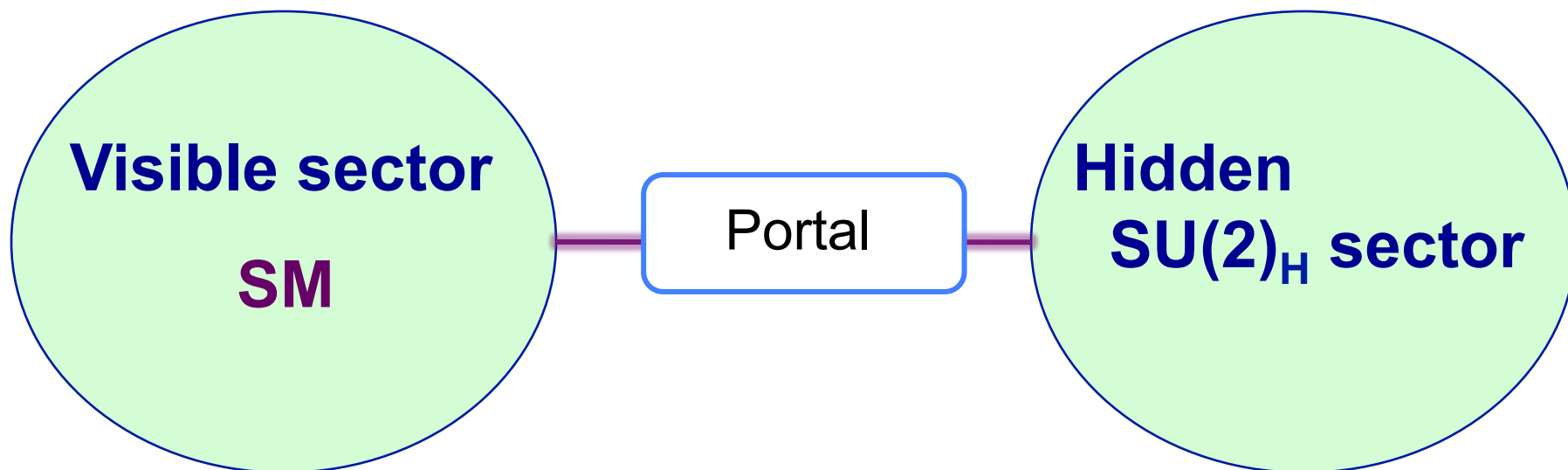
$2m_X = m_{Z'}$: when $SU(2)_H$ is broken by quintet scalar VEV

$SU(2)_H \times U(1)_{B-L}$ model C. W. Chiang, T. N, J. Tandean, (JHEP 01, 2014)

Structure of our scenario

❖ Singlet under $SU(2)_H$

❖ Singlet under G_{SM}



Two sector can interact via scalar mixing and/or gauge boson mixing

⇒ We consider kinetic mixing associated with $SU(2)_H$ and $U(1)_Y$

❖ In previous work we consider mass mixing in $U(1)_{B-L}$ and $SU(2)_H$

2. Model

A model

❖ $G_{SM} \times SU(2)_H$ gauge symmetry

❖ New field contents: $SU(2)_H$

*Scalar fields

Quintet: $\Phi = (\phi_2, \phi_1, \phi_0, \phi_{-1}, \phi_{-2})$ Triplet: $\varphi = \varphi^a \frac{\sigma^a}{2}$

*Gauge boson

$SU(2)_H : (X_\mu^1, X_\mu^2, X_\mu^3)$

❖ Scalar VEV and symmetry breaking

$$\langle \Phi \rangle = (v_\Phi / \sqrt{2}, 0, 0, 0, 0) \quad \langle \varphi \rangle = \frac{1}{2\sqrt{2}} \begin{pmatrix} v_\varphi & 0 \\ 0 & -v_\varphi \end{pmatrix}$$

$SU(2)_H \rightarrow Z_2$

Eigen value of $\sigma^3/2$ is even(odd) \rightarrow Parity even(odd)

❖ X^1 and X^2 are Z_2 odd and our DM candidate

Hidden Gauge Sector with SM $U(1)_Y$

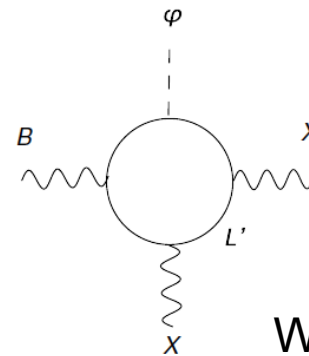
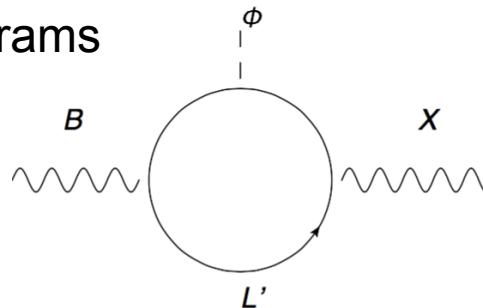
To connect $SU(2)_H$ and $U(1)_Y$ we introduce dim 5 operator

$$-\frac{C_\phi}{\Lambda} \varphi^a X^{a\mu\nu} B_{\mu\nu}$$

It can be generated introducing field charged under $SU(2)_H$ and $U(1)_Y$

e.g. $SU(2)_H$ doublet fermion L' with $U(1)_B$ charge -1

1-loop diagrams



Work in progress

We use the Lagrangian :

$$L = -\frac{1}{4} X^{a\mu\nu} X_{\mu\nu}^a - \frac{1}{4} B^{\mu\nu} B_{\mu\nu} - \frac{C_\phi}{\Lambda} \varphi^a X^{a\mu\nu} B_{\mu\nu}$$

Hidden Gauge Sector with SM $U(1)_Y$

$$L = -\frac{1}{4} X^{a\mu\nu} X_{\mu\nu}^a - \frac{1}{4} B^{\mu\nu} B_{\mu\nu} - \frac{C_\phi}{\Lambda} \varphi^a X^{a\mu\nu} B_{\mu\nu} \quad \left[\sin \chi \equiv \sqrt{2} C_\phi v_\phi / \Lambda \right]$$

$$+ D^\mu \Phi D_\mu \Phi + \text{Tr}[D^\mu \varphi D_\mu \varphi] \quad \Rightarrow \quad \mathcal{L}_{\text{KM}} = -\frac{1}{2} \sin \chi X_{\mu\nu}^3 \tilde{B}^{\mu\nu}$$

Diagonalizing kinetic terms

$$B_\mu \rightarrow B_\mu - \tan \chi X_\mu^3$$

$$X_\mu^3 \rightarrow \frac{1}{\cos \chi} X_\mu^3$$

For tiny χ

$$\Rightarrow \quad B_\mu \rightarrow B_\mu - \chi X_\mu^3$$

$$X_\mu^3 \rightarrow X_\mu^3$$

Gauge boson mass term after symmetry breaking $X_\mu^\pm = (X_\mu^1 \mp X_\mu^2) / \sqrt{2}$

$$L_M = \frac{1}{2} m_{Z_{SM}}^2 \tilde{Z}_\mu \tilde{Z}^\mu + m_{Z_{SM}}^2 \chi \sin \theta_W \tilde{Z}_\mu X^{3\mu} + \frac{1}{2} m_{X^3}^2 X_\mu^3 X^{3\mu} + m_{X^\pm}^2 X_\mu^+ X^{-\mu},$$

$$m_{Z_{SM}}^2 = \frac{v^2}{4} (g^2 + g_B^2), \quad m_{X^3}^2 = 4g_X^2 v_\Phi^2, \quad m_{X^\pm}^2 = g_X^2 v_\Phi^2 \left(1 + \frac{v_\phi^2}{v_\Phi^2} \right),$$

Hidden gauge boson masses and interactions

- Z,Z' masses and Z-Z' mixing:

$$m_{Z,Z'}^2 = \frac{1}{2}(m_{X^3}^2 + m_{Z_{SM}}^2) \mp \frac{1}{2}\sqrt{(m_{X^3}^2 - m_{Z_{SM}}^2)^2 + 4\chi^2 \sin^2 \theta_W m_{Z_{SM}}^4},$$

$$\tan 2\theta_{ZZ'} = \frac{2 \sin \theta_W \chi m_{Z_{SM}}^2}{m_{Z_{SM}}^2 - m_{X^3}^2},$$

- DM mass: $m_{X^\pm}^2 = g_X^2 v_\Phi^2 \left(1 + \frac{v_\varphi^2}{v_\Phi^2}\right) \Rightarrow m_{Z'} \simeq 2m_{X^\pm} (1 + R_M)^{-\frac{1}{2}}$
 $R_M \equiv v_\varphi^2/v_\Phi^2$

- Relevant DM interactions:

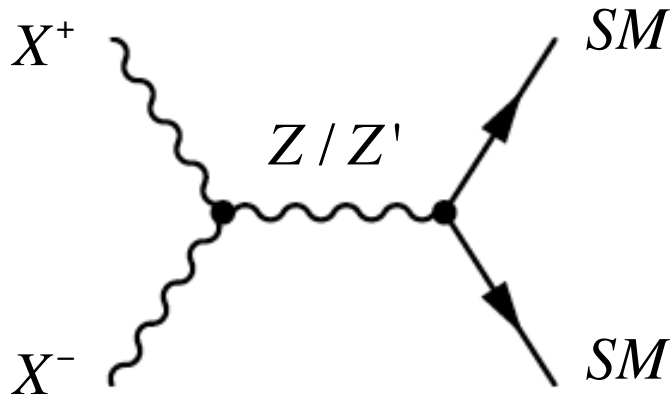
$$\mathcal{L} \supset ig_X C_{ZZ'} \left[(\partial_\mu X_\nu^+ - \partial_\nu X_\mu^+) X^{-\mu} Z'^{\nu} - (\partial_\mu X_\nu^- - \partial_\nu X_\mu^-) X^{+\mu} Z'^{\nu} + \frac{1}{2} (\partial_\mu Z'_\nu - \partial_\nu Z'_\mu) (X^{+\mu} X^{-\nu} - X^{-\mu} X^{+\nu}) \right],$$

[CZZ'(SZZ') = cosθ_{ZZ'}, (sinθ_{ZZ'})]

$$\mathcal{L}_{Z'ff} = \frac{g}{\cos \theta_W} Z'_\mu \bar{f} \gamma^\mu \left[-S_{ZZ'} (T_3 - Q \sin^2 \theta_W) + C_{ZZ'} \chi Y \sin \theta_W \right] f.$$

3. DM phenomenology

DM annihilation and interaction with nucleon



Z' exchange is dominant for DM annihilation



Resonant enhancement

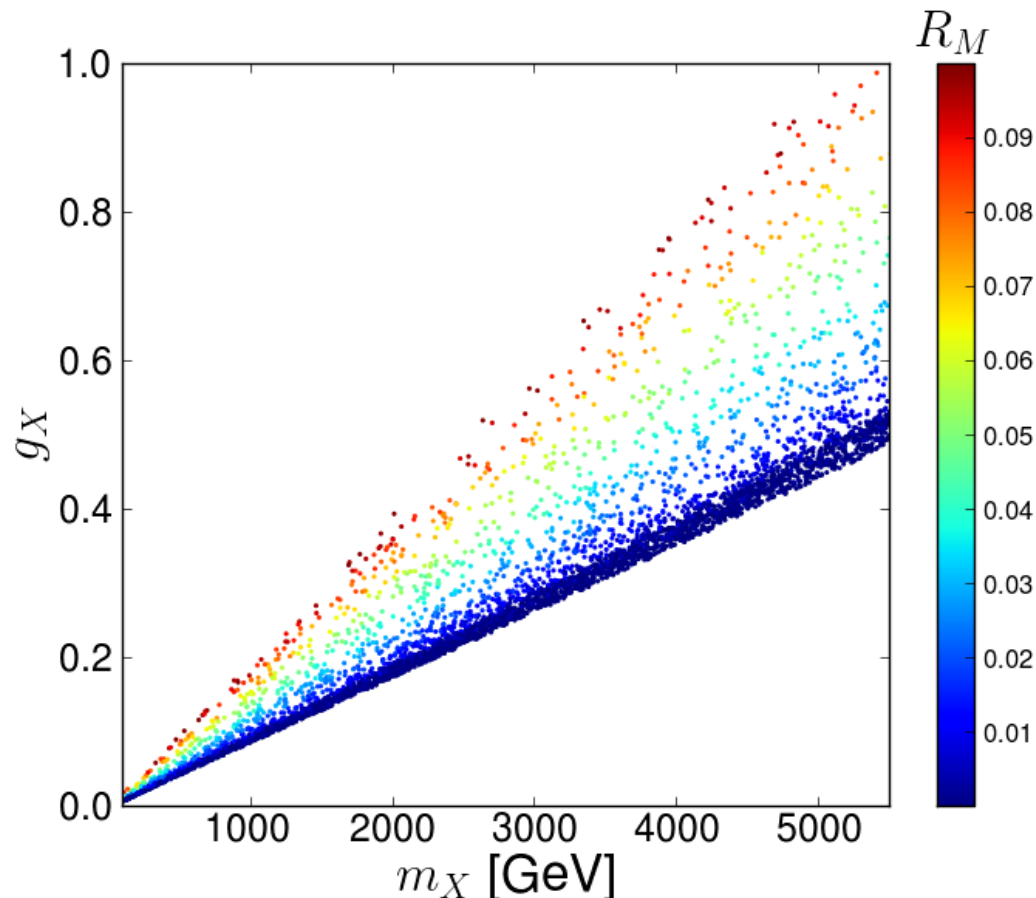
Relic density is estimated by micrOMEGAs

DM-Nucleon scattering cross section

$$\sigma_{NX} = \frac{2g_X^2 g^2}{\pi \cos^2 \theta_W} \left(\frac{m_{X^\pm} m_N}{m_X + m_N} \right)^2 \left(\frac{S_{ZZ'} C_{NNZ}^V}{m_Z^2} + \frac{C_{ZZ'} C_{NNZ'}^V}{m_{Z'}^2} \right)^2$$

$$\left(\begin{array}{l} C_{ppZ}^V = C_{ZZ'} \left(\frac{1}{4} - \sin^2 \theta_W \right), \quad C_{nnZ}^V = C_{ZZ'} \left(-\frac{1}{4} \right), \\ C_{ppZ'}^V = -S_{ZZ'} \left(\frac{1}{4} - \sin^2 \theta_W \right) + \frac{3}{2} C_{ZZ'} \sin \theta_W \delta, \quad C_{nnZ'}^V = \frac{1}{4} S_{ZZ'} + \frac{1}{2} C_{ZZ'} \sin \theta_W \delta. \end{array} \right)$$

Parameter region realizing observed relic density



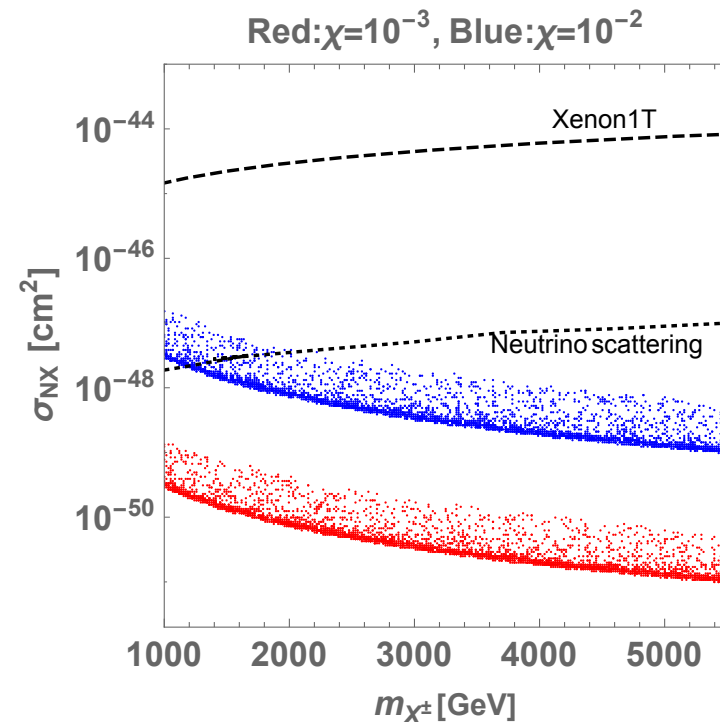
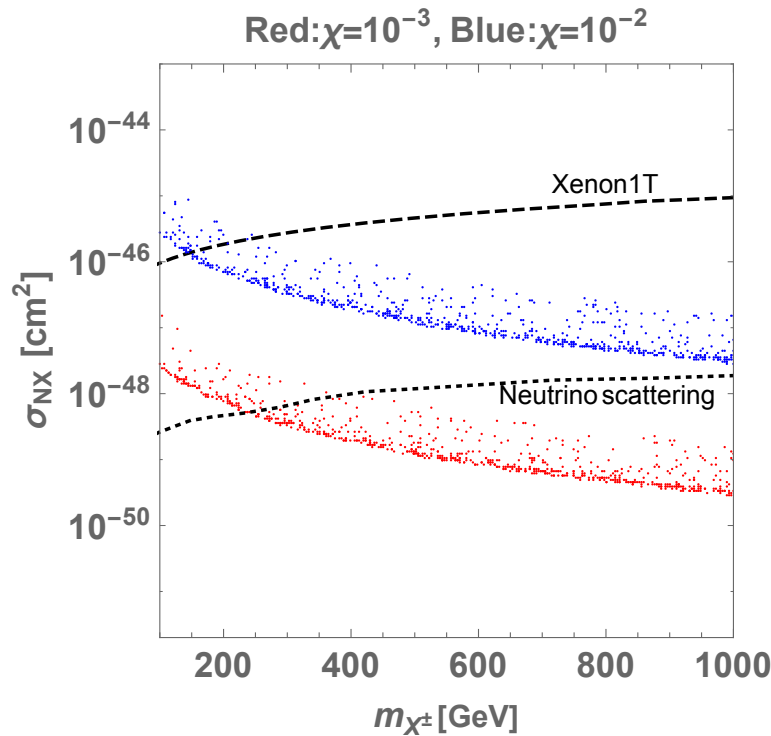
Parameters are scanned
fixing $\chi = 10^{-3}$

Approximated range
 $0.11 < \Omega h^2 < 0.13$

- O(0.01)-O(1) gauge coupling for 100 GeV – 5500 GeV DM mass
- Larger $R_M \rightarrow$ larger gauge coupling

3. DM phenomenology

DM-Nucleon scattering cross section



Compared with Xenon1T constraint

XENON (2018)

- Most of parameter region is allowed
- It is difficult to detect DM when kinetic mixing parameter is small

Summary and discussion

Construction of model with dark SU(2) model

- ✧ Remnant Z_2 symmetry as a subgroup of SU(2)
- ✧ DM candidate is hidden SU(2) gauge bosons
- ✧ DM interaction with SM through kinetic mixing

Application to DM phenomenology

- ✧ Relic density of DM
- ✧ Constraints from direct detection

Appendix

Z_2 symmetry as a remnant of $SU(2)_H$

Suppose $SU(2)$ quintet and triplet scalars Φ, Φ' get VEVs as

$\langle \phi_0 \rangle \neq 0, \langle \phi'_1 \rangle \neq 0$ For components with eigenvalue of T_3 is 0 and 2

$$\exp[iT_3\pi] |Vacuum\rangle = |Vacuum\rangle \quad \text{Remaining discrete symmetry}$$

Then for any particle with odd T_3 value

$$\exp[iT_3\pi] X = \exp[i(2n+1)\pi] X = -X$$

 Z_2 symmetry: odd/even for field with odd/even T_3 value

The lightest Z_2 odd particle in dark sector is DM

Scalar VEV and hidden gauge boson masses

When a component of scalar $SU(2)_H$ multiplet (characterized by l, m) gets VEV

$$\begin{aligned} & \frac{g_X^2}{4} [(C_{\ell, m}^+)^2 + (C_{\ell, m}^-)^2] v_\Phi^2 X_\mu^+ X^{-\mu} + \frac{g_X^2}{2} m^2 v_\Phi^2 X_\mu^3 X^{3\mu} \\ & \equiv m_{X^\pm}^2 X_\mu^+ X^{-\mu} + \frac{1}{2} m_{X^3}^2 X_\mu^3 X^{3\mu}. \end{aligned} \quad X_\mu^\pm \equiv (X_\mu^1 \mp iX_\mu^2)/\sqrt{2}$$

$$C_{\ell, m}^+ \equiv \sqrt{(\ell - m)(\ell + m + 1)} \text{ and } C_{\ell, m}^- \equiv \sqrt{(\ell + m)(\ell - m + 1)}.$$

Mass relation

$$\frac{m_{X^3}^2}{m_{X^\pm}^2} = \frac{4m^2}{(C_{\ell, m}^+)^2 + (C_{\ell, m}^-)^2} = \frac{4m^2}{(\ell - m)(\ell + m + 1) + (\ell + m)(\ell - m + 1)}.$$

Z-Z' mixing

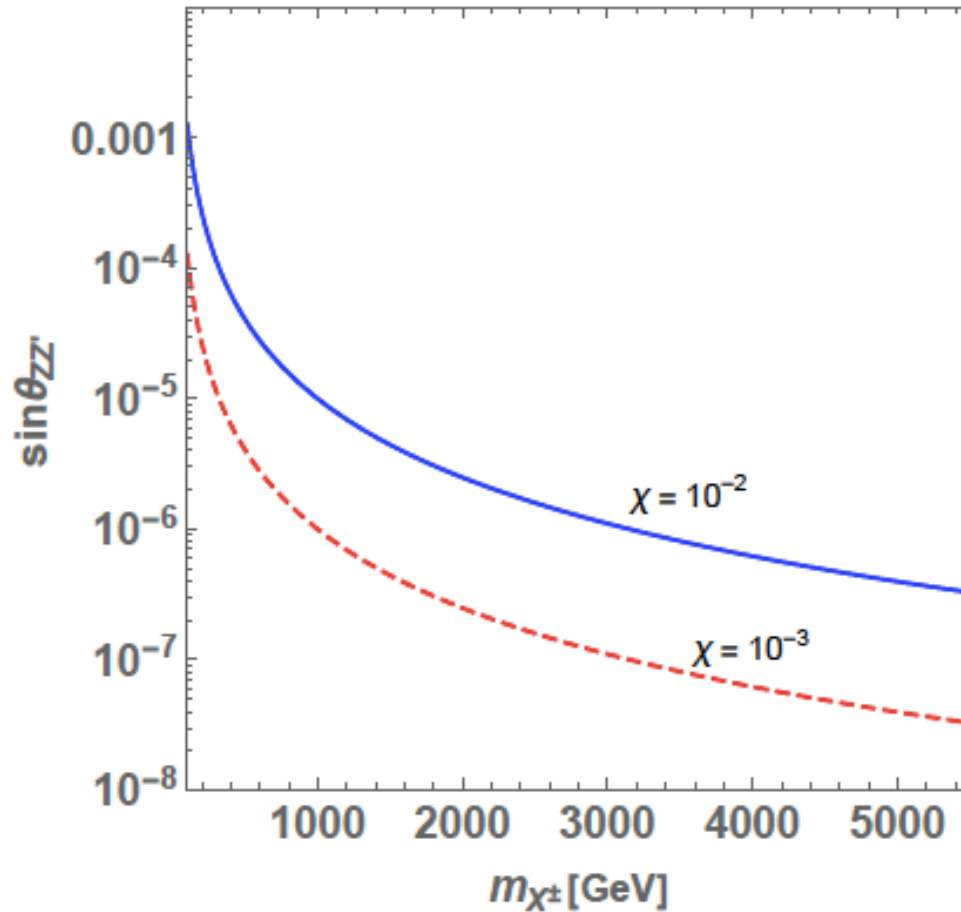


FIG. 1: $\sin\theta_{ZZ'}$ as a function of m_{χ^\pm} for $\chi = 10^{-2}$ and 10^{-3} where we fix $2v_\varphi^2/v_\Phi^2 = 0.01$.