## Affine Gravitational Scenario for Dark Matter Decays

Hemza Azrí Department of Physics, UAEU

5th International Workshop on "Higgs as a Probe of New Physics" 25-27 March 2021, Osaka University, Japan

A talk based on: H. Azrí, A. Jueíd, C. Karahan, and S. Nasrí, Phys. Rev. D102. 8, 084036 (2020)

# Gravity: Metrical or affine?

- of distances and angles...etc !
- Gravity is a "manifestation" of spacetime curvature, hence, what is needed is the concept of connection.
- fields are described by connections (arising through gauge invariance...)
- simple and possibly viable description for gravity !

• GR is based on the metric tensor which is only of secondary importance: measurements

property of parallel displacements of tensors. This is achieved by introducing the

• The SM interactions are written in the framework of gauge field theory where the gauge

• Affine spacetime, with its geodesics and no prior concept of metric, provides a very

## Attine space structure

• Affine space is endowed with an <u>affine</u> connection from which we construct covariant derivatives:

$$\nabla_{\mu}A^{\alpha} = \partial_{\mu}A^{\alpha} + \Gamma^{\alpha}_{\mu\beta}A^{\beta}$$

It vanishes locally,

$$\Gamma^{\alpha}_{\mu\beta}(P) = 0$$

• It transforms as  

$$\hat{\Gamma}^{\alpha}_{\lambda\mu} = \frac{\partial \hat{x}^{\alpha}}{\partial x^{\beta}} \frac{\partial x^{\sigma}}{\partial \hat{x}^{\lambda}} \frac{\partial x^{\tau}}{\partial \hat{x}^{\mu}} \Gamma^{\beta}_{\sigma\tau} + \frac{\partial^{2}}{\partial \hat{x}^{\lambda}} \frac{\partial x^{\sigma}}{\partial \hat{x}^{\mu}} \Gamma^{\beta}_{\sigma\tau} + \frac{\partial^{2}}{\partial \hat{x}^{\mu}} + \frac{\partial^{2}}{\partial \hat{x}^{\mu}} + \frac{\partial^{2}}{\partial \hat{x}^{\mu}} + \frac{\partial^{2}}{\partial \hat{x}^{\mu}} + \frac{\partial^{2}}{\partial \hat{x}^{\mu}$$

• <u>Curvature</u> can be defined from the "failure" of the commutativity:  $[\nabla_{\mu}, \nabla_{\nu}]A^{\alpha} = R^{\alpha}_{\beta\mu\nu}(\Gamma)A^{\beta}$ 



 $\partial_{\mu}\partial_{\nu}A^{\alpha} - \partial_{\nu}\partial_{\mu}A^{\alpha} = 0$ Flat background

 $\nabla_{\mu}\nabla_{\nu}A^{\alpha} - \nabla_{\nu}\nabla_{\mu}A^{\alpha} \neq 0$ 

Curved background

When  $\Gamma$  is symmetric

### • Gravitational action with the SM fields:

- This involves Einstein-Hilbert action and the SM Lagrangian
- The SM Lagrangian is written by simply replacing the flat metric by the general curved metric

$$L[g, SM] = -g^{\mu\nu}(D_{\mu})$$

• How about dark matter?

Its existence is revealed only through its gravitational influence!

 $S = \left[ d^4 x \sqrt{|g|} \left\{ \frac{M_{Pl}^2}{2} g^{\mu\nu} R_{\mu\nu}(g) + L[g, SM] \right\} \right]$ 

 ${}_{\mu}H)^{\dagger}(D_{\nu}H) - V(H) - i\bar{\psi}g^{\mu\nu}\gamma_{\mu}\nabla_{\nu}\psi - \frac{1}{4}g^{\mu\nu}g^{\alpha\beta}F^{a}_{\mu\alpha}F^{a}_{\nu\beta}$ Higgs Fermions Gauge bosons



• Decoupling dark matter from the metric:

Unlike the SM particle fields, DM may couple only to the spacetime connection!

- From Palatíní to purely affine action:
  - Take the following metric-affine model in which metric and connection are independent fields!

$$\mathcal{L}[\mathcal{I}] = \sqrt{|g|} \left\{ \frac{M^2}{2} g^{\mu\nu} R_{\mu\nu}(\Gamma) - \frac{1}{2} g^{\mu\nu} X_{\mu\nu} - U(\mathcal{I}) \right\}$$

Kinetic terms! 
$$X_{\mu\nu}(\mathcal{I}) = \begin{cases} 
abla_{\mu}\phi 
abla_{\nu}\phi \\ 
\bar{\chi}\gamma_{\mu} \nabla_{\nu}\chi \\ 
\frac{1}{2}g^{\alpha\beta} \mathbb{F}_{\alpha\mu} \mathbb{F}_{\beta\nu} \end{cases}$$



In this case, how to decouple it from the metríc?

for scalar DM for fermionic DM for vector DM

Fermions and vector fields involve metric in their kinetic terms. Scalars do not!





### From Palatíní to purely affine actíon:

Field equations arise from variation with respect to the (non-dynamical) metric

$$\mathcal{L}[\mathcal{I}] = \sqrt{|g|} \left\{ \frac{M^2}{2} g^{\mu\nu} R_{\mu\nu}(\Gamma) - \frac{1}{2} g^{\mu\nu} X_{\mu\nu} - U(\mathcal{I}) \right\}$$

For scalar DM

Hence, for nonzero potential, the metric can be easily integrated out!

$$g_{\mu\nu} = \frac{M^2 R_{\mu\nu}(\Gamma) - X_{\mu\nu}(\phi)}{U(\phi)}$$

$$M^{2}R_{\mu\nu}(\Gamma) = \nabla_{\mu}\phi \nabla_{\nu}\phi + U(\mathcal{F})g_{\mu\nu}$$
  
Metric arises  
only through the  
potential term!



The Scalar DM Lagrangian can be written in terms of the connection and its curvature only!



• If dark matter is to decouple totally from the metric, it can simply be a scalar

$$\mathscr{L}[\Gamma, \mathsf{DM}] = \frac{\sqrt{\left|M^2 R_{\mu\nu}(\Gamma) - \nabla_{\mu}\phi \nabla_{\nu}\phi\right|}}{U(\phi)}$$

- 1. The model describes DM in curved spacetime interacting only with the connection (notice the absence of metric!)
- 2. Gravitational field equation are obtained by variation w.r.t the connection

$$\nabla_{\alpha} \begin{cases} \frac{\sqrt{|M^2 R_{\mu\nu} (\Gamma) - \nabla_{\mu} \phi \nabla_{\nu} \phi|}}{U(\phi)} \left( [M^2 R (\Gamma) - \nabla \phi \nabla \phi] \right) \end{cases}$$

$$\sqrt{|g|} g^{\mu\nu}$$

This action can also be constructed based only on general covariance consideration!



 $M^2 R_{\mu\nu}(\Gamma) = \nabla_{\mu} \phi \nabla_{\nu} \phi + U(\mathcal{I})g_{\mu\nu}$ Einstein's field equations!





Nonmínímal couplings
 (DM\_Curvature interaction)

- 1. The  $Z_2$  symmetry enjoyed by the previous action is broken by an interaction between DM and the affine curvature like  $\xi \phi R_{\mu\nu}(\Gamma)$
- 2. Fields that enjoy a complete structure in purely affine spacetime are scalars, hence, the model can be extended by including the SM-Higgs potential

$$\mathscr{L}[\Gamma, \mathsf{DM}] = \frac{\sqrt{\left|M^2 R_{\mu\nu}(\Gamma) - \nabla_{\mu}\phi \nabla_{\nu}\phi + \xi M\phi R_{\mu\nu}(\Gamma)\right|}}{W(I, II)}$$

Where  $V(\phi, H) \supset U(\phi) + m_H^2 H^{\dagger} H + \lambda_H (H^{\dagger} H)^2$ .

 $V(\phi, H)$ 

Notice that there is No DM-SM coupling!

• Dark matter decays through the Higgs portal

> • The affine gravitational sector is recast to its canonical form as

$$\mathcal{L}[\Gamma, \mathsf{DM}] = \frac{\sqrt{\left|M^2 R_{\mu\nu}(\Gamma) - \alpha^{-1}(\phi) \nabla_{\mu} \phi \nabla_{\nu} \phi\right|}}{\alpha^{-2}(\phi) V(\phi, H)}$$

• As a result, the first order  $Z_2$ -symmetry breaking term implies DM decays to SM-Higgs through the induced interactions  $\alpha^{-2}(\phi)V(\phi,H)$ 



The interactions are a direct effects of the nonminimal coupling to affine gravity even in the absence of any explicit coupling of DM with SM particles within the potential !

With the dimensionless parameter  $\alpha(\phi) = 1 + \xi \phi/M$ 



• The resulting Decay modes

$$\phi \rightarrow hh$$
  
 $\phi \rightarrow hh^* \rightarrow hf\bar{f}, fVV, hhh$  Where  $V = W^{\pm}/Z$   
 $\phi \rightarrow h^*h^* \rightarrow VVVV, VVf\bar{f}, VVhh, f\bar{f}f'\bar{f}', f\bar{f}hh, hhhh$   
The decay rates are suppressed by  
the factor  $\xi^2/M^2$ , leading to a large  
enough life time when  $M = M_{Pl}$ !

• Relevant DM Mass

ranges:

- $m_{\phi} \sim 10^{-3} 125 \ GeV$
- $m_{\phi} \sim 125 250 \ GeV$
- $m_{\phi} \sim 250 10^6 \; GeV$











### Conclusion

- We have investigated the possibility that, unlike ordinary matter, DM interacts to gravity through only the affine connection of spacetime!
- We saw that DM decays, automatically, into the SM Higgs via a nonminimal coupling to gravity that breaks the Z<sub>2</sub>-symmetry! This decay occurs without imposing any prior coupling between the two!
- The present scenario for DM decay is fully gravitational and it stands on a formulation of gravity, different from GR. It is expected to reveal more interesting features of the physics of DM!

