



Novel probes of Higgs couplings to light SM quarks

Jose Miguel No IFT-UAM/CSIC, Madrid

Based on 2008.12538 (w. Aguilar-Saavedra & Cano) 2011.09551 (w. Falkowski, Ganguly, Gras, Tobioka, Vignaroli, You)



HPNP 2021





Higgs Yukawa couplings to light SM fermions

• Establish role of Higgs in mass generation of 1st & 2nd generation SM fermions



Higgs Yukawa couplings to light SM fermions

• Establish role of Higgs in mass generation of 1st & 2nd generation SM fermions

"Higgs flavour"



[source: nobelprize.org]

Leptons:

• Recent ATLAS & CMS evidence of Higgs coupling to muons



CERN-EP-2020-164 2020/09/10

EXPERIMENT Phys. Lett. B 812 (2021) 135980 DOI: 10.1016/j.physletb.2020.135980



A search for the dimuon decay of the Standard Model Higgs boson with the ATLAS detector

EUROPEAN ORGANISATION FOR NUCLEAR RESEARCH (CERN)

The ATLAS Collaboration

Evidence for Higgs boson decay to a pair of muons

The CMS Collaboration*

2009.04363

2007.07830

Quarks: (more complicated...)

O Strategies (non-exhaustive!) to probe light quark Yukawas @ LHC



- Charge asymmetry in W[±]h production Yu. JHEP 02 (2017) 083 (1609.06592)
- Double Higgs production

Alasfar, Corral Lopez, Grober. JHEP 11 (2019) 088 (1909.05279) Egana-Ugrinovic, Hollimer, Meade. 2101.04119

Quarks: (more complicated...)

O Strategies (non-exhaustive!) to probe light quark Yukawas @ LHC

Higgs + charm production Brivio, Isidori, Goertz. PRL 115, 211801 (1507.02916)

Higgs + jet production Bishara, Haisch, Monni, Re. PRL 118, 121801 (1606.09253) All these needed? Why look for new strategies?

Exotic Higgs decays

Bodwin, Petriello, Stoynev, Velasco. PRD 88, 053003 (1306.5770) Kagan, Perez, Petriello, Soreq, Stoynev, Zupan. PRL 114, 101802 (1406.1722) Konig, Neubert. JHEP 08 (2015) 012 (1505.03870) ...

• Higgs $\eta \in p_T$ distributions

Soreq, Zhu, Zupan. JHEP 12 (2016) 045 (1606.09621)

Charge asymmetry in W[±]h production Yu. JHEP 02 (2017) 083 (1609.06592)

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"Different probes sensitive to different sets of couplings/EFT operators"

New strategies to constrain Higgs Yukawas @LHC

• Higgs + photon production

Aguilar-Saavedra, Cano, No. 2008.12538

 $pp \to h \gamma$

2 Triple gauge boson production

Falkowski, Ganguly, Gras, No, Tobioka, Vignaroli, You. 2011.09551

$$pp \to V V V$$

Higgs + photon

Would-be leading contribution vanishes (Furry Theorem)





- $\Rightarrow (Q_u/Q_d)^2 = 4 \rightarrow$ more sensitive to up-type quark Yukawa deviations
- $\Rightarrow \text{ SM cross section small, enhanced for BSM Higgs Yukawas: } \kappa_q > 1$ (e.g. $\sigma_{u\bar{u}} = 1.3 \text{ fb for } y_u(m_h) \sim y_c^{\text{SM}}(m_h)$) $\kappa_q = y_q(m_h)/y_q^{\text{SM}}(m_h)$
- $\Rightarrow \text{ Most promising final state: } h \rightarrow WW^* \rightarrow \ell \nu \ell \nu$ (Clean + sufficient XS)

Higgs + photon $h \to WW^* \to \ell \nu \ell \nu$

⇒ Dominant SM backgrounds:

$$pp \rightarrow \ell^+ \nu \ell^- \bar{\nu} \gamma$$

$$pp \rightarrow Z\gamma, \ Z \rightarrow \tau^+ \tau^-$$

$$pp \rightarrow t\bar{t}\gamma \qquad t \rightarrow b\ell^+ \nu$$

$$\bar{t} \rightarrow \bar{b}\ell^- \bar{\nu}$$

$h \to WW^* \to \ell \nu \ell \nu$

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 $\begin{array}{ll} p_T^{\gamma} > 25 \ {\rm GeV} \\ p_T^{\ell_1} > 18 \ {\rm GeV}, \, p_T^{\ell_2} > 15 \ \ {\rm GeV} \ \ \left| \begin{array}{c} \right| & p_T^{\ell_1} > 23 \ \ {\rm GeV}, \, p_T^{\ell_2} > 9 \ \ {\rm GeV} \\ \\ {E_T} > 35 \ \ {\rm GeV} \end{array}$

Initial selection



$$h \to WW^* \to \ell \nu \ell \nu$$

⇒ Dominant SM backgrounds:

Higgs + photon

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Kinematically rich final state

 $M_T, M_{\ell\ell}, M_{\ell\ell\gamma}, \vec{p}_T^{\ell_1}, \vec{p}_T^{\ell_2}, \vec{p}_T^{\gamma}, \not\!\!\!E_T, \\ \Delta \phi^{\ell\ell}, \Delta \phi^{\ell_1\gamma}, \Delta \phi^{\ell_2\gamma}, \Delta \eta^{\ell\ell}, \eta^{\ell_1}, \eta^{\ell_2}, \eta^{\gamma}$

Strong correlations among variables

Multivariate analysis significantly increases cut-&-count sensitivity

 $h \to WW^* \to \ell \nu \ell \nu$

Neural Network (NN) multivariate analysis



HL-LHC sensitivity (3 ab⁻¹):

 $ert \kappa_c ert < 11.8$ $ert \kappa_u ert < 1930$ (95% C.L.)

 $h \to WW^* \to \ell \nu \ell \nu$

Neural Network (NN) multivariate analysis



HL-LHC sensitivity (3 ab^{-1}) :

$$ert \kappa_c ert < 11.8 \ ert \kappa_u ert < 1930 \$$
 (95%C.L.)

Possible to look for $h + \gamma$ in other final states?

Besides its Yukawa sensitivity, $h + \gamma$ interesting in its own right: Unobserved Higgs production mode!

Tri-boson

Triple massive gauge boson production recently observed for the first time @ LHC!



Triple threat: The first observation of three massive gauge bosons produced in proton-proton collisions

by Ingrid Fadelli , Phys.org



(deviation in $h\bar{q}q$ coupling from SM leads to quadratic growth with c.o.m energy for $q\bar{q} \rightarrow VVV$ XS)



Key in controlling high-energy behaviour of VVV amplitude

(deviation in $h\bar{q}q$ coupling from SM leads to quadratic growth with c.o.m energy for $q\bar{q} \rightarrow VVV$ XS)

VVV can be used to constrain Higgs Yukawa couplings

... following idea of "measuring Higgs couplings without Higgs bosons" Henning, Lombardo, Riembau, Riva. PRL 123 (2019) 181801



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Deviations of Higgs Yukawas from SM: EFT description

Add D = 6 operators to SM:

$$\mathcal{L}_{\text{SMEFT}} \supset \frac{Y_u |H|^2}{v^2} \bar{u}_R Q_{1,L} H + \frac{Y_d |H|^2}{v^2} \bar{d}_R H^{\dagger} Q_{1,L} + \frac{Y_s |H|^2}{v^2} \bar{s}_R H^{\dagger} Q_{2,L} + \text{h.c.}$$

- \Rightarrow We focus on u, d, s
- \Rightarrow For $q = c, t, pp \rightarrow WWqj$ more sensitive

Brooijmans, Buckley, Caron, Falkowski, Fuks, Gilbert, Murray, Nardecchia, No, Torre, You, Zevi della Porta. PhysTeV 2019. New Physics WG, 2002.12220 (Contribution 12)



Key in controlling high-energy behaviour of VVV amplitude

(deviation in $h\bar{q}q$ coupling from SM leads to quadratic growth with c.o.m energy for $q\bar{q} \rightarrow VVV$ XS)

VVV can be used to constrain Higgs Yukawa couplings

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Relation between mass and Yukawa after EWSB:

By **Equivalence Theorem**:

$$\mathcal{M}(q\bar{q} \to V_L V_L V_L) \xleftarrow{\sqrt{\hat{s}} \gg m_Z} \mathcal{M}(q\bar{q} \to GGG)$$

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} i\sqrt{2}G_+\\ v+h+iG_z \end{pmatrix}$$

$$\mathcal{L}_{\text{SMEFT}} \supset \frac{Y_u |H|^2}{v^2} \bar{u}_R Q_{1,L} H + \frac{Y_d |H|^2}{v^2} \bar{d}_R H^{\dagger} Q_{1,L} + \frac{Y_s |H|^2}{v^2} \bar{s}_R H^{\dagger} Q_{2,L} + \text{h.c.}$$

Contact interaction (2 quarks + 3 Goldstone bosons)

$$\mathcal{L} \supset \frac{1}{v^2} \left(G_+ G_- + \frac{1}{2} G_z^2 \right) \left\{ i y_u^{\mathrm{SM}} \delta y_u \left(\sum_{q'=d,s} \bar{u}_R q'_L G_+ - \bar{u}_R u_L \frac{G_z}{\sqrt{2}} \right) + i \sum_{q'=d,s} y_{q'}^{\mathrm{SM}} \delta y_{q'} \left(\bar{q}'_R u_L G_- + \bar{q}'_R q'_L \frac{G_z}{\sqrt{2}} \right) + \text{h.c.} \right\}.$$



By Equivalence Theorem:

$$\mathcal{M}(q\bar{q} \to V_L V_L V_L) \xleftarrow{\sqrt{\hat{s}} \gg m_Z} \mathcal{M}(q\bar{q} \to GGG)$$

$$\mathcal{L}_{\text{SMEFT}} \supset \frac{Y_u |H|^2}{v^2} \bar{u}_R Q_{1,L} H + \frac{Y_d |H|^2}{v^2} \bar{d}_R H^{\dagger} Q_{1,L} + \frac{Y_s |H|^2}{v^2} \bar{s}_R H^{\dagger} Q_{2,L} + \text{h.c.}$$

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 $H = \frac{1}{\sqrt{2}} \begin{pmatrix} i\sqrt{2}G_+ \\ v + h + iG_z \end{pmatrix}$

Leading contribution

(other tree-level diagrams more suppressed from extra internal propagators)

$$\mathcal{M}(q\bar{q} \to GGG) \sim \mathcal{O}(\delta y_q E/v^2)$$

$$\begin{aligned} \sigma(q\bar{q} \to G_z G_+ G_-) &= (y_q^{\rm SM} \delta y_q)^2 I(\hat{s}) \\ \sigma(q\bar{q} \to 3G_z) &= \frac{3}{2} (y_q^{\rm SM} \delta y_q)^2 I(\hat{s}) \\ \sigma(u\bar{q}' \to G_+ G_z G_z) + \sigma(q'\bar{u} \to G_- G_z G_z) &= \frac{1}{2} \left[(y_u^{\rm SM} \delta y_u)^2 + (y_{q'}^{\rm SM} \delta y_{q'})^2 \right] I(\hat{s}) \\ \sigma(u\bar{q}' \to G_+ G_+ G_-) + \sigma(q'\bar{u} \to G_- G_- G_+) &= 2 \left[(y_u^{\rm SM} \delta y_u)^2 + (y_{q'}^{\rm SM} \delta y_{q'})^2 \right] I(\hat{s}) \end{aligned}$$

$$\left(\begin{array}{c} q = u, d, s \\ q' = d, s \end{array}\right)$$

$$I(\hat{s}) \equiv \frac{\hat{s}}{6144\pi^3 v^4}$$

c.o.m. energy of partonic collision $(\sqrt{\hat{s}})$
 \mathbf{E}^2 growth of partonic cross section

• For charged (±1) final states ($W^{\pm}W^{\pm}W^{\mp}$, $W^{\pm}ZZ$) same cross section enhancement for δy_u and δy_d

• For **neutral** final states (ZW^+W^- , ZZZ) **different** cross section enhancement for δy_u and δy_d

Break degeneracies combining several tri-boson channels

HL-LHC	SM	BSM $(Y_d = 1)$	BSM $(Y_u = 1)$	BSM $(Y_s = 1)$
$W^+W^-W^+$	152 fb	3.6 pb	3.6 pb	109 fb
$W^+W^-W^-$	87 fb	1.5 pb	1.5 pb	109 fb
ZZW^+	40 fb	1.0 pb	1.0 pb	31 fb
ZZW ⁻	23 fb	0.43 pb	0.43 pb	31 fb
ZW^+W^-	191 fb	1.5 pb	2.4 pb	115 fb
ZZZ	16 fb	0.99 pb	1.7 pb	66 fb

(SM: NLO with MadGraph)

(BSM: LO with MadGraph)

 $W^{\pm}W^{\pm}W^{\mp}$ Largest cross section (+ BR into leptons larger for W than for Z) ZZZ Largest cross section enhancement w.r.t. SM

_					
	BSM $(Y_s = 1)$	BSM $(Y_u = 1)$	BSM $(Y_d = 1)$	SM	HL-LHC
	109 fb	3.6 pb	3.6 pb	152 fb	$W^+W^-W^+$
(SM: NLO with	109 fb	1.5 pb	1.5 pb	87 fb	$W^+W^-W^-$
	31 fb	1.0 pb	1.0 pb	40 fb	ZZW^+
(BSM: LO with	31 fb	0.43 pb	0.43 pb	23 fb	ZZW ⁻
MadGraph)	115 fb	2.4 pb	1.5 pb	191 fb	ZW^+W^-
]	66 fb	1.7 pb	0.99 pb	16 fb	ZZZ

Can readily use CMS observation of SM tri-boson production to constrain δy_q

Sirunyan et al (CMS). PRL 125 (2020) 151802 [CMS-SMP-19-014]

Selection cuts: $pp \to W^{\pm}W^{\pm}W^{\mp} \to \ell^{\pm}\ell^{\pm}\nu\nu jj$ $p_T^{\ell_{1,2}} > 25 \,\text{GeV}, \ m_{\ell\ell} > 20 \,\text{GeV}, \ m_{jj} \in [65, 95] \,\text{GeV} \,(\text{``m_{jj} in''})$ $E_T^{\text{miss}} > 45 \,\text{GeV}, m_T^{\text{max}}(\ell) > 90 \,\text{GeV}$



]	BSM $(Y_s = 1)$	BSM $(Y_u = 1)$	BSM $(Y_d = 1)$	SM	HL-LHC
	109 fb	3.6 pb	3.6 pb	152 fb	$W^+W^-W^+$
(SM: NLO with	109 fb	1.5 pb	1.5 pb	$87~{\rm fb}$	$W^+W^-W^-$
	31 fb	1.0 pb	1.0 pb	$40~{\rm fb}$	ZZW^+
(BSM: LO with	31 fb	0.43 pb	0.43 pb	23 fb	ZZW ⁻
MadGraph)	115 fb	2.4 pb	$1.5 \mathrm{~pb}$	$191~{\rm fb}$	ZW^+W^-
	66 fb	1.7 pb	0.99 pb	16 fb	ZZZ

Can readily use CMS observation of SM tri-boson production to constrain δy_a

Sirunyan et al (CMS). PRL 125 (2020) 151802 [CMS-SMP-19-014]



Selection cuts: $pp \to W^{\pm}W^{\pm}W^{\mp} \to \ell^{\pm}\ell^{\pm}\nu\nu jj$ $p_T^{\ell_{1,2}} > 25 \,\text{GeV}, \ m_{\ell\ell} > 20 \,\text{GeV}, \ m_{jj} \in [65, 95] \,\text{GeV} (``m_{jj} \,\text{in''})$ $E_T^{\text{miss}} > 45 \,\text{GeV}, m_T^{\text{max}}(\ell) > 90 \,\text{GeV}$

> These cuts are not optimized for BSM!

				I
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ZZW^+	$40~{\rm fb}$	$1.0 \ \mathrm{pb}$	$1.0 \ \mathrm{pb}$	31 fb
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(SM: NLO with MadGraph)

(BSM: LO with MadGraph)



BSM much harder kinematics than SM

Improved selection cuts:

 $p_T^{\ell_{1,2}} > 60 \, {\rm GeV} \;, \quad E_T^{\rm miss} > 120 \, {\rm GeV} \;, \quad p_T^{jj} > 120 \, {\rm GeV} \;, \quad |\Delta \eta(\ell_1,\ell_2)| < 2$

+ binned likelyhood analysis:

 $\delta y_d \lesssim 430 \; (\text{HL-LHC})$ $\delta y_u \lesssim 850 \; (\text{HL-LHC})$ $\delta y_s \lesssim 150 \; (\text{HL-LHC})$

HL-LHC	SM	BSM $(Y_d = 1)$	BSM $(Y_u = 1)$	BSM $(Y_s = 1)$
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Cross sections

(BSM: LO with

(SM: NLO with

MadGraph)

MadGraph)

 $pp \to ZZZ \to 4\ell + 2\nu$

$$\begin{split} p_T^{\ell_{1,2}} > 25 ~{\rm GeV} \,, \, p_T^{\ell_{3,4}} > 10 ~{\rm GeV} \,, \, |\eta_\ell| < 2.5 \,, \, \Delta R_{\ell\ell} > 0.1 \,, \, |m_Z - m_{\ell\ell}| < 10 ~{\rm GeV} \,, \, E_T^{\rm miss} \, > \, 200 ~{\rm GeV} \\ (\Delta R_{\ell\ell} > 0.01 ~{\rm FCC-hh}) & (E_T^{\rm miss} > 500 ~{\rm GeV} ~{\rm FCC-hh}) \end{split}$$





HL-LHC	SM	BSM $(Y_d = 1)$	BSM $(Y_u = 1)$	BSM $(Y_s = 1)$
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Cross sections

MadGraph)

(SM: NLO with

(BSM: LO with MadGraph)

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Sensitivity Summary HL-LHC (FCC-hh)

	WWW			ZZZ		
	$\ell^{\pm}\ell^{\pm} + 2\nu + 2j$	$\ell^{\pm}\ell^{\pm}\ell^{\mp} + 3\nu$	Comb.	$4\ell + 2\nu$	$4\ell + 2j$	Comb.
δy_d	430 (36)	840(54)	420(34)	1500 (65)	1300 (93)	1100~(60)
δy_u	850 (71)	1700 (110)	830(68)	2300(100)	1800(140)	1600 (92)
δy_s	150(13)	230(33)	140(13)	300(12)	290(16)	250(11)



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	WWW			ZZZ		
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δy_s	150(13)	230(33)	140(13)	300(12)	290(16)	250(11)



Other EFT operators may also impact VVV (e.g. anomalous TGCs)

[Global Fit]

Complementarity among probes

 $[h\,\gamma]\,,\,\,[VVV]\,,\,\,[h\!+\!j]\,(ggf)\,,\,\,[h\,h]\,,\,\,[h\rightarrow J/\psi\!+\!\gamma]\ldots$



Thank you!

TTT TT		CN	DCM(V 1)	DCM(V 1)	DOM (V 1)	1
HL-Lf	IC	SM	$BSM (Y_d = 1)$	$BSM(Y_u = 1)$	$BSM (Y_s = 1)$	
W^+W^-	W^+	152 fb	3.6 pb	3.6 pb	109 fb	
W^+W^-	W^-	$87~{\rm fb}$	1.5 pb	1.5 pb	109 fb	(SM: NLO with
ZZW	·+	40 fb	1.0 pb	1.0 pb	31 fb	
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ZW^+V	V^{-}	191 fb	1.5 pb	2.4 pb	115 fb	MadGraph)
	7	16 fb	0.99 pb	1.7 pb	66 fb]

Cross sections

 $pp \to W^{\pm}W^{\pm}W^{\mp} \to \ell^{\pm}\ell^{\pm}\nu\nu jj$ $\sigma(Y_d) = 7.5 \,\text{fb} + Y_d^2 \times 205 \,\text{fb}$ Large BSM cross section enhancement

Limits – No Reducible Bck $\delta y_d \lesssim 430 \; (\text{HL-LHC})$ $\delta y_u \lesssim 850 \; (\text{HL-LHC})$ $\delta y_s \lesssim 150 \; (\text{HL-LHC})$ These projected limits get very little affected by reducible SM backgrounds:

We simulate

$$t\bar{t}W^{\pm}, t\bar{t}Z$$
 (NLO in QCD)

 $W^{\pm}Z\,jj$ (LO)

Together they yield (after cuts!) ~ 20% of SM tri-boson (LHC)

ZZ reducible background MET fit for $pp \rightarrow ZZZ \rightarrow 4\ell + 2\nu$



Auxiliary material of:

Sirunyan et al (CMS). PRL 125 (2020) 151802 [CMS-SMP-19-014]