



# Structure of the Higgs potential in Gauge-Higgs Unification with a Flat Extra Dimension

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## Introduction (Mystery in the Higgs Sector)

Hierarchy problem suggest the new physics at TeV scale

- Supersymmetry
- Composite Models
- **Gauge-Higgs Unification**

$$A_M = (A_\mu, A_m)$$

4D gauge     4D Higgs

$$\Rightarrow V_{\text{eff}}^{\text{GHU}} = \text{const.} + \sum_i \frac{n_i}{2} \int \frac{d^4 p_E}{(2\pi)^4} \ln(p_E^2 + m_i^2(\phi)) \quad \text{one-loop correction}$$

We analyze the Higgs potential and triple coupling in the SU(3) model with 5D Lorentz symmetry relaxed

**SU(3) Model** [Panico, Serone, Wulzer, NPB 739 (2006)]

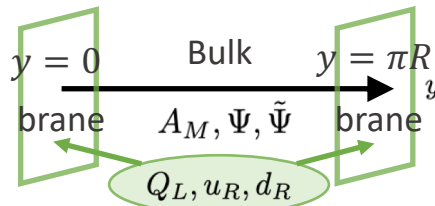
**Space-time** : Flat  $M^4 \times S^1/Z_2$  with compactification scale  $1/R$

**Symmetry Breaking** :  $SU(3)_w \times U(1)'$

$$\xrightarrow{\text{BC}} SU(2)_L \times U(1)_Y \times U(1)_X$$

**Field contents**

- Bulk gauge field :  $A_M$
- Bulk fermion pair :  $\{\Psi, \tilde{\Psi}\}$
- Brane fermion :  $Q_L, u_R, d_R$



## Effective Potential in the Flat SU(3) Model

Conventional notation :  $\frac{g_4}{2} \phi = \frac{\alpha}{R}$

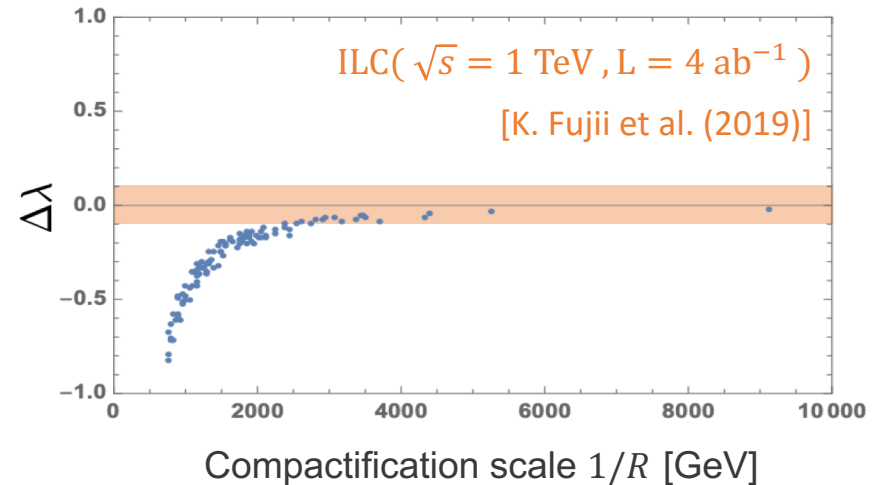
$$V_{\text{eff}}(\alpha) = \text{const.} + \sum V_A(q\alpha) + \sum V_\Psi(q\alpha) + \sum V_a$$

$$\xrightarrow{\alpha \ll 1} \left(\frac{1}{R}\right)^2 A\phi^2 + B\phi^4 + C\phi^4 \ln \frac{\phi^2}{v^2} + \left(\frac{1}{R}\right)^{-2} D\phi^6$$

$$V_{\text{SM}}(\phi) = A'\phi^2 + B'\phi^4 + C\phi^4 \ln \frac{\phi^2}{v^2}$$

## Deviation of the Triple Higgs Boson Coupling

$$\Delta\lambda \equiv \frac{\lambda_{hhh} - \lambda_{hhh}^{\text{SM}}}{\lambda_{hhh}^{\text{SM}}} \sim \left(\frac{1}{R}\right)^{-2} : \left. \frac{\partial V}{\partial \phi} \right|_{\phi=v} = 0, \quad \left. \frac{\partial^2 V}{\partial \phi^2} \right|_{\phi=v} = m_h^2$$



# SU(3) Model [Panico, Serone, Wulzer, NPB 739 (2006)]

space-time : Flat  $M^4 \times S^1/Z_2$  with compactification scale  $1/R$



Boundary conditions ( gauge field )

$$\begin{cases} S^1 : A_M(y + 2\pi R) = A_M(y) \\ Z_2 : A_\mu(-y) = P^\dagger A_\mu(y) P, A_5(-y) = -P^\dagger A_5(y) P \end{cases}$$

symmetry :  $SU(3)_c \times SU(3)_w \times U(1)' \xrightarrow{BC} SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_X$

$$A_\mu = \begin{pmatrix} (+,+) & (+,+) & (-,-) \\ (+,+) & (+,+) & (-,-) \\ (-,-) & (-,-) & (+,+) \end{pmatrix}, A_5 = \begin{pmatrix} (-,-) & (-,-) & (+,+) \\ (-,-) & (-,-) & (+,+) \\ (+,+) & (+,+) & (-,-) \end{pmatrix} \rightarrow \begin{matrix} \text{SM gauge} \\ \text{Higgs doublet} \end{matrix}$$

Matter Lagrangian

5D Lorentz invariance relaxed

$$\begin{aligned} \mathcal{L}_{\text{mat}}^{5D} = & \sum_j \left\{ \bar{\Psi}_j (i\not{D}_4 - \mathbf{k}_j D_5 \gamma^5) \Psi_j + \bar{\tilde{\Psi}}_j (i\not{D}_4 - \tilde{\mathbf{k}}_j D_5 \gamma^5) \tilde{\Psi}_j + \frac{1}{\pi R} (\bar{\Psi}_j \lambda_j \tilde{\Psi}_j + \text{h.c.}) \right\} \textcircled{1} \\ & + \delta(y - y_1) \left\{ \bar{Q}_L i\not{D}_4 Q_L + \sqrt{\frac{2}{\pi R}} (\epsilon_1^d \bar{Q}_L \psi_d + \epsilon_1^u \bar{Q}_R^c \psi_u + \text{h.c.}) \right\} \textcircled{2} \\ & + \delta(y - y_2) \left\{ \bar{u}_R i\not{D}_4 u_R + \bar{d}_R i\not{D}_4 d_R + \sqrt{\frac{2}{\pi R}} (\epsilon_2^d \bar{d}_R \chi_d + \epsilon_2^u \bar{u}_L^c \chi_u + \text{h.c.}) \right\} \textcircled{2} \end{aligned}$$

① TeV scale bulk mass term

② Bulk-Brane mixing : Source of the brane fermion mass

→ Reproduce the mass scale of the SM fermions

1-loop effective potential

$$V_{\text{eff}}(\alpha) = \text{const.} + \sum V_g(q\alpha) + \sum V_\Psi^\pm(q\alpha) + \sum V_a(\alpha)$$

bulk gauge    bulk fermion    brane fermion

$$V_g(q\alpha) = -\frac{9}{64\pi^6 R^4} \sum_{n=1}^{\infty} \frac{1}{n^5} \cos(2\pi n q\alpha)$$

$$V_\Psi^\pm(q\alpha) = \frac{3k^4}{8\pi^6 R^4} \sum_{n=1}^{\infty} \frac{(\pm 1)^n}{n^5} \left( 1 + 2n \frac{\lambda}{k} + \frac{4}{3} n^2 \frac{\lambda^2}{k^2} \right) e^{-2n \frac{\lambda}{k}} \cos(2\pi n q\alpha)$$

$$V_a(\alpha) = \frac{c_a}{2} \int \frac{d^4 p_E}{(2\pi)^4} \ln \{ Z_1^a(\alpha) Z_2^a(\alpha) p_E^2 + m_a(\alpha)^2 \}$$

Matter configuration

	$SU(3)_c \times SU(3)_w$	$\eta$	Contribution to $V_{\text{eff}}(\alpha)$
$\{\Psi_t, \tilde{\Psi}_t\}$	$(\mathbf{3}, \bar{\mathbf{6}})$	0	$3V_{\Psi_t}^+(2\alpha) + 3V_{\Psi_t}^+(\alpha)$
$\{\Psi_b, \tilde{\Psi}_b\}$	$(\mathbf{3}, \mathbf{3})$	0	$3V_{\Psi_b}^+(\alpha)$
$\{\Psi_A, \tilde{\Psi}_A\}$	$(\mathbf{1}, \mathbf{6})$	1/2	$V_{\Psi_A}^-(2\alpha) + V_{\Psi_A}^-(\alpha)$

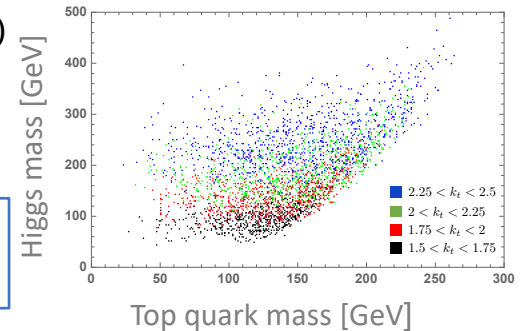
Brane fermion :  $Q_L = (t_L, b_L)^T \rightarrow y = 0$  ;  $t_R, b_R \rightarrow y = \pi R$

Model parameters  $\sim O(1)$

$$\epsilon_1^t, \epsilon_2^t, \epsilon_1^b, \epsilon_2^b$$

$$\lambda_t, \lambda_b, \lambda_A \quad k_t, k_b, k_A$$

By tuning them, we can reproduce the  $m_h, m_t, m_{KK}$



# Analyze of the Higgs potential

GHU :  $V_{\text{eff}}^{\text{GHU}} = \text{const.} + \sum_i \frac{n_i}{2} \int \frac{d^4 p_E}{(2\pi)^4} \ln(p_E^2 + m_i^2(\phi))$

$$\xrightarrow{1/R \rightarrow \infty} \left(\frac{1}{R}\right)^2 A\phi^2 + B\phi^4 + C\phi^4 \ln \frac{\phi^2}{v^2} + \left(\frac{1}{R}\right)^{-2} D\phi^6$$

SM :  $V_{\text{SM}}(\phi) = A'\phi^2 + B'\phi^4 + C\phi^4 \ln \frac{\phi^2}{v^2}$

## Tadpole and mass conditions

$$\left. \frac{\partial V_{\text{eff}}}{\partial \phi} \right|_{\phi=v} = 0 \quad \rightarrow \quad B' = -\frac{2}{C} - \frac{A'}{2v^2} - \frac{3}{2} D \frac{v^4}{(1/R)^2}$$

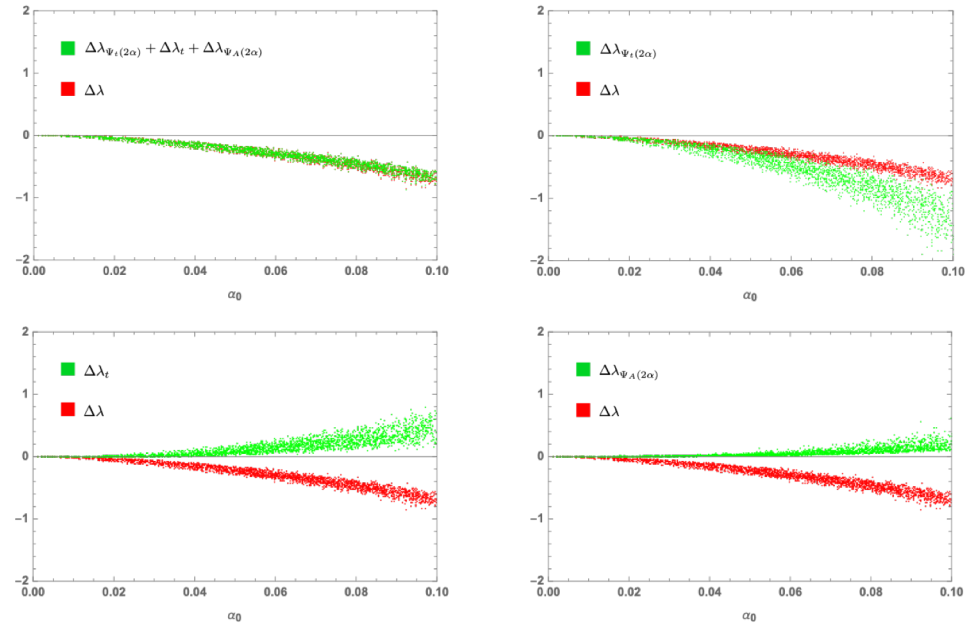
$$\left. \frac{\partial^2 V_{\text{eff}}}{\partial \phi^2} \right|_{\phi=v} = m_h^2 \quad \rightarrow \quad m_h^2 = -4A' + 8Cv^2 + 12 \frac{v^4}{(1/R)^2}$$

## Triple Higgs boson coupling in $1/R \rightarrow \infty$ limit

$$\lambda_{hhh} = \frac{3}{v} \left\{ m_h^2 + \frac{16}{3} \left( C + 3D \frac{v^2}{(1/R)^2} \right) v^2 \right\} \rightarrow \Delta\lambda \sim \left(\frac{1}{R}\right)^{-2}$$

This feature is not only found in this model but also models with bulk-brane mixing on flat  $S^1/Z_2$  orbifold

# Dominant contribution of $\Delta\lambda$

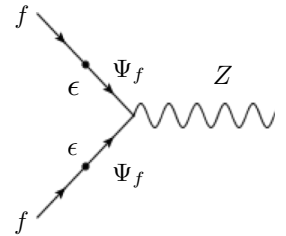


## limitation of the compactification scale

G. Panico, M. Serone, A. Wulzer, NPB 739 (2006) , NPB 762 (2007)

### Direct correction

- $Z b_L \bar{b}_L$ :  $\frac{\delta g_b}{g_b} \simeq \frac{1}{1 - \frac{2}{3} \sin^2 \theta_W} \frac{\epsilon_1^{t2} k_t^2}{\lambda_t^2 Z_1} \left(\frac{m_W}{M_t}\right)^2$
- FCNC :  $g_{\text{FCNC}} \sim \frac{\epsilon_1^u \epsilon_1^c}{\lambda_u \lambda_c} \left(\frac{m_W}{M_t}\right)^2$



### Oblique correction

- $T = \hat{T}/\alpha_e = \alpha_0^2 \pi^2 / \alpha_e \Rightarrow 1/R \geq 5 \text{ TeV}$