

LISA: Lifetime measurements with Solid Active targets

Kathrin Wimmer

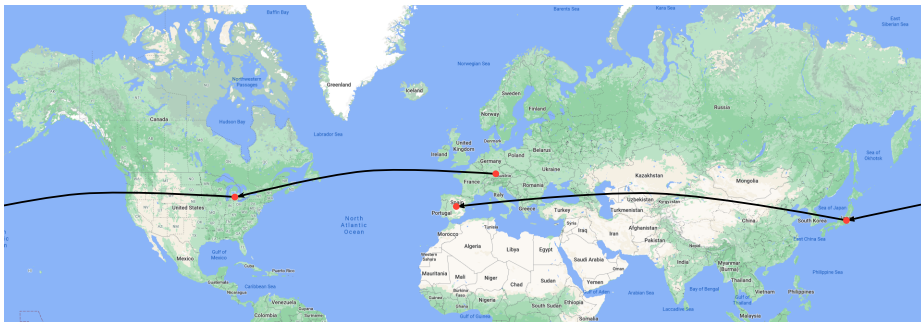
Instituto de Estructura de la Materia - CSIC Madrid

10. May 2021



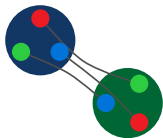
CSIC
CONSEJO SUPERIOR DE INVESTIGACIONES CIENTÍFICAS

- studied at Ludwigs-Maximilians Universität (Munich) and Università di Pisa (Italy)
- PhD at TU Munich (2010), work done at REX-ISOLDE, CERN (Switzerland)
- moved to the National Superconducting Cyclotron Laboratory (USA) as postdoc (2010 - 2012)
- Assistant Professor at Central Michigan University (2012 - 2014)
- Faculty (Kōshi) at The University of Tokyo (2014 - 2019)
- since June 2019 at CSIC
- main research theme: structure and reactions with exotic nuclei

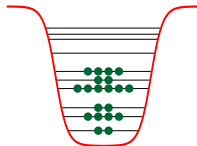
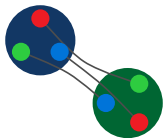


- 1 Introduction and motivation
- 2 In-beam γ -ray Spectroscopy of Exotic Nuclei
- 3 Measuring lifetimes of excited states
- 4 LISA: Lifetime measurements with Solid Active targets

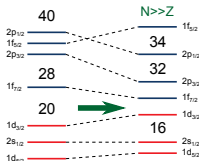
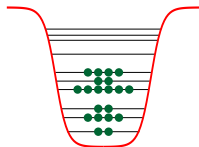
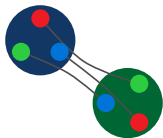
- nuclei govern the Universe at many length (fm to km) and time scales (μs to Gy)
- dual quantum liquid consisting of protons and neutrons which are interacting
- leads to the well-known shell structure of atomic nuclei
- strong nuclear interaction and its isospin dependence leads to changes along isotopic and isotonic chains
- shell evolution can change the spacing and ordering of the single-particle orbitals
 - degenerate levels in a nuclear Jahn-Teller effect
 - spontaneous symmetry breaking and quantum phase transitions
- excitations associated with collective quantum mechanical rotations or vibrations



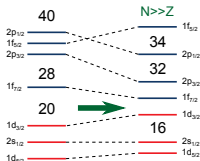
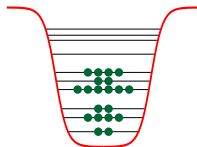
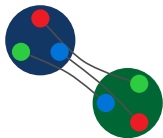
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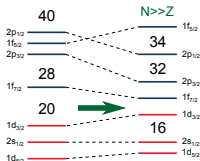
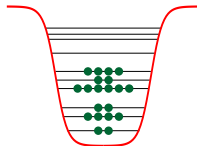
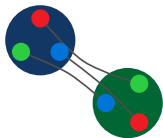
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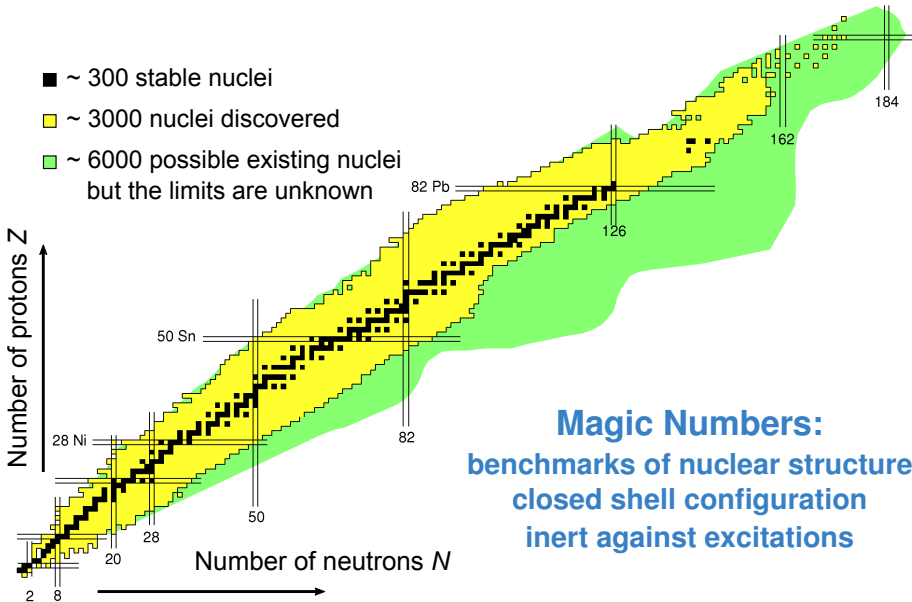
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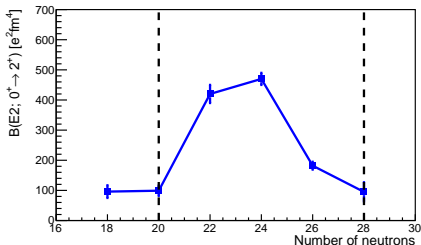
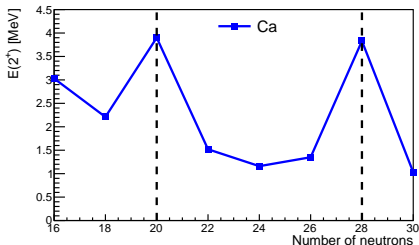
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where and how does collective motion and deformation of nuclei emerge from the single-particle degrees of freedom of the protons and neutrons?



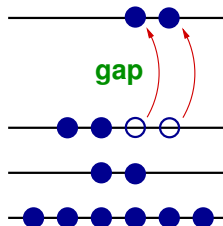
Example: disappearance of shell closures



<http://www.nndc.bnl.gov/ensdf/>, K. Wimmer and P. Doornenbal, Prog. Part. Nucl. Phys. (2021)

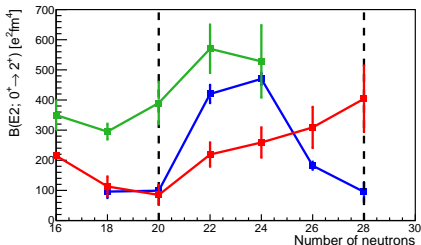
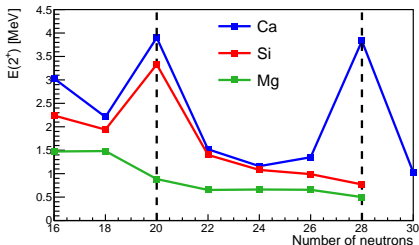
³⁶ Ca	³⁸ Ca	⁴⁰ Ca	⁴² Ca	⁴⁴ Ca	⁴⁶ Ca	⁴⁸ Ca	⁵⁰ Ca
³⁴ Ar	³⁶ Ar	³⁸ Ar	⁴⁰ Ar	⁴² Ar	⁴⁴ Ar	⁴⁶ Ar	⁴⁸ Ar
³² S	³⁴ S	³⁶ S	³⁸ S	⁴⁰ S	⁴² S	⁴⁴ S	⁴⁶ S
³⁰ Si	³² Si	³⁴ Si	³⁶ Si	³⁸ Si	⁴⁰ Si	⁴² Si	⁴⁴ Si
²⁸ Mg	³⁰ Mg	³² Mg	³⁴ Mg	³⁶ Mg	³⁸ Mg	⁴⁰ Mg	

- high energy of the first excited state
- small B(E2) value, difficult to excite



■ many exotic nuclei do not show these signatures

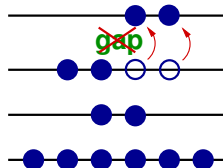
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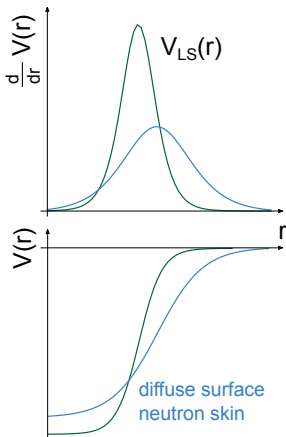
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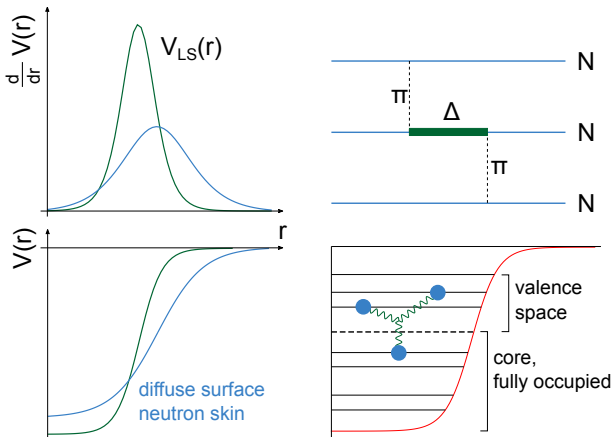


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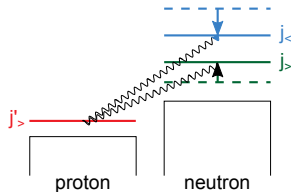
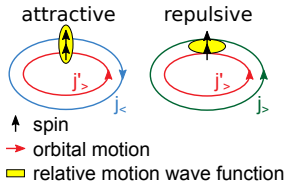
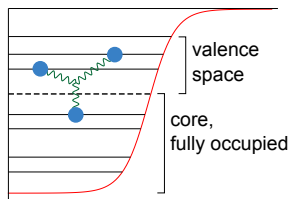
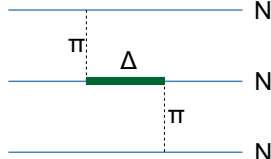
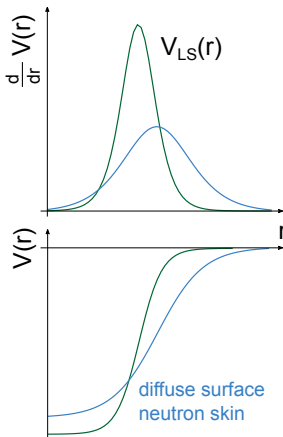
Figures adapted from T. Otsuka et al., Phys. Rev. Lett. 105, 032501 (2010), Phys. Rev. Lett. 87, 082502 (2001), Phys. Rev. Lett. 95, 232502 (2005).

- spin-orbit interaction, excess neutron distribution leads to weaker spin-orbit effect
- three-body interactions, needed to explain binding of light nuclei
- tensor interaction between protons and neutrons
- all affect the level spacing and ordering



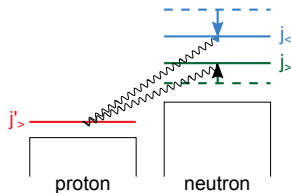
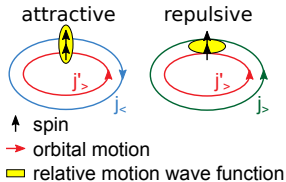
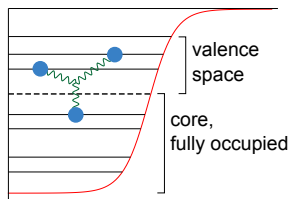
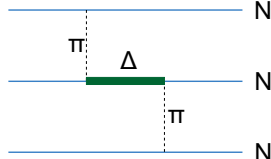
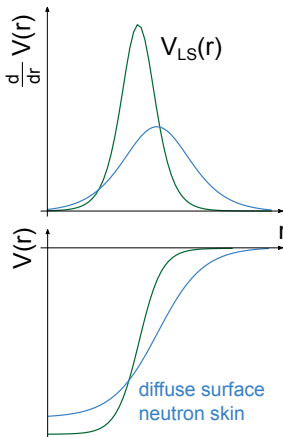
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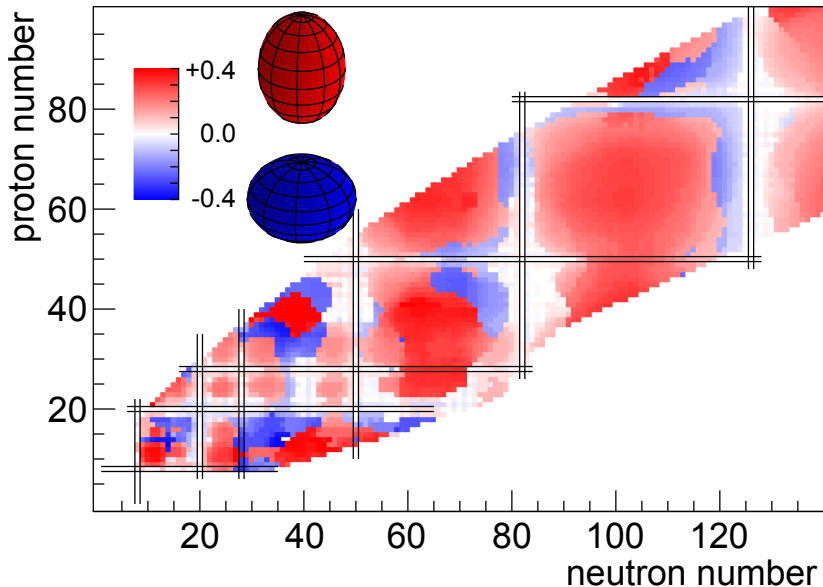
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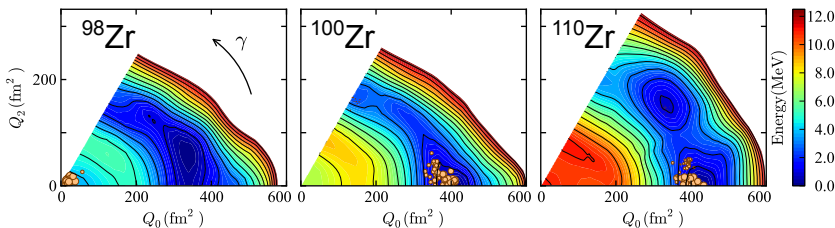
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- predicted deformation parameters using finite-range droplet macroscopic model



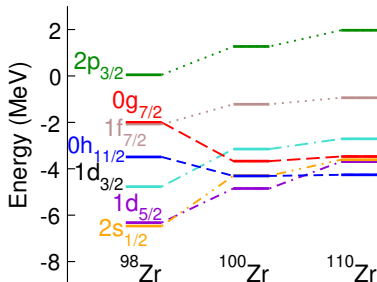
P. Möller et al., At. Data Nucl. Data Tables **109** (2016) 1.

- a deformed nucleus is characterized by collective and coherent quantum mechanical rotations or vibrations



- shapes change as a function of proton or neutron number
- proton-neutron interaction changes the ordering and spacing of levels
- (near-) degeneracy triggers symmetry breaking and deformation
- calculations predict deformation and triaxiality for ^{110}Zr

T. Togashi et al., Phys. Rev. Lett. **117** (2016) 172502.



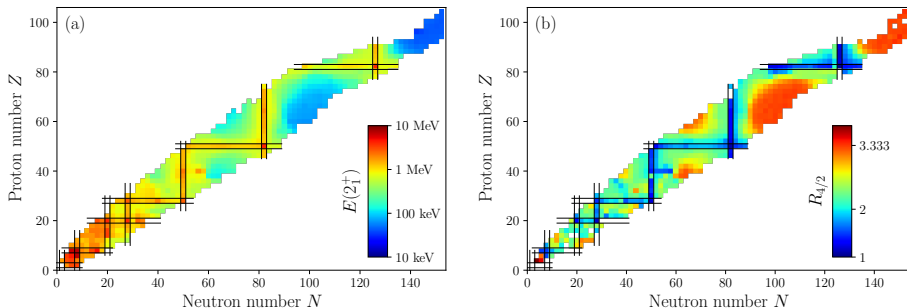


Figure from K. Wimmer and P. Doornenbal, Prog. Part. Nucl. Phys. (2021).

- excitation energy of the first 2_1^+ state
- ratio $R_{4/2} = E(4_1^+)/E(2_1^+)$ mirrors the $E(2_1^+)$ behavior
- magic nuclei: pair breaking \rightarrow similar energies of 2_1^+ and $4_1^+ \rightarrow R_{4/2} \sim 1$
- vibrational nuclei: 2_1^+ state one phonon, 4_1^+ two phonon excitation $\rightarrow R_{4/2} = 2$
- rotational nuclei: $E(J) \propto J(J+1) \rightarrow R_{4/2} = 3.33$

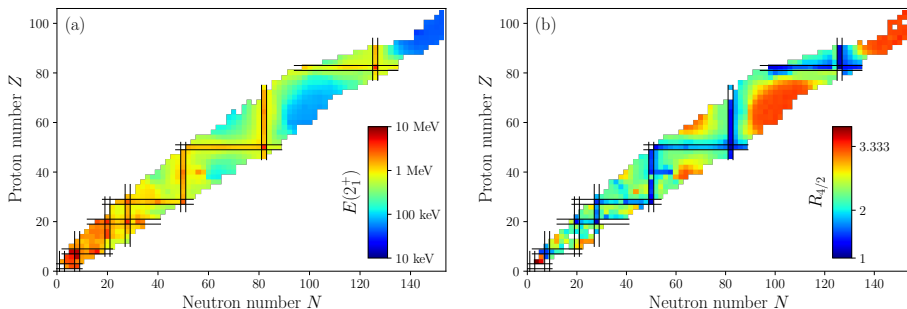


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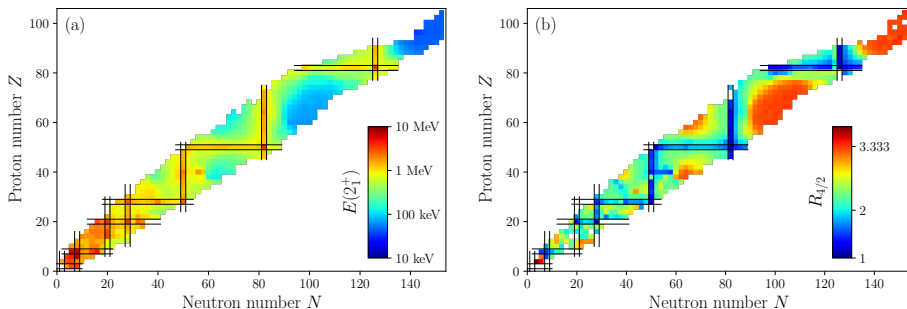


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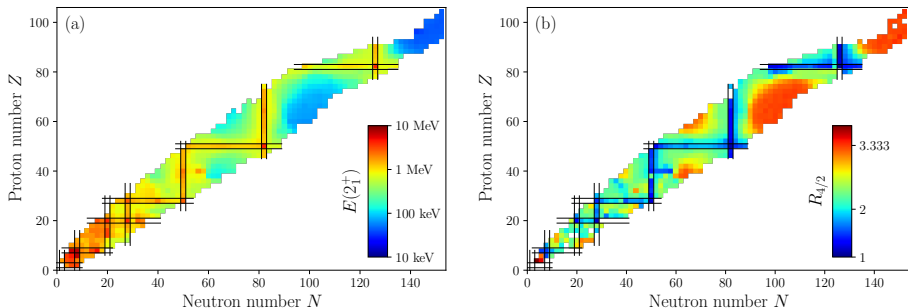


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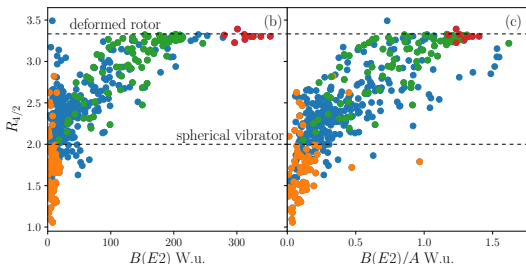
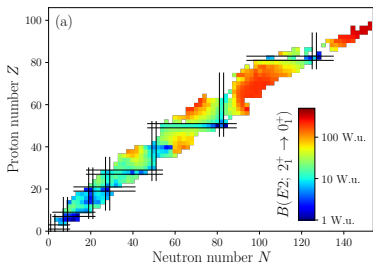
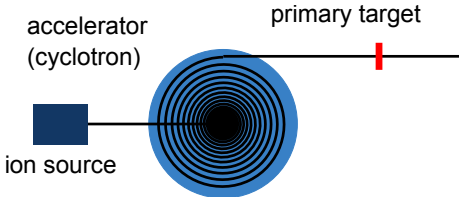


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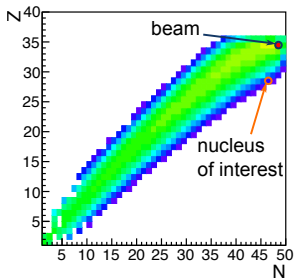
- reduced transition probability $B(E2)$
- direct measure of nuclear collectivity
- in Weisskopf units: how many nucleons participate in the excitation
- experimental observable:
 - excitation cross section or probability, lifetime and decay branching
- main techniques (for exotic nuclei): Coulomb excitation and direct lifetime measurements (this talk)

In-beam γ -ray Spectroscopy of Exotic Nuclei

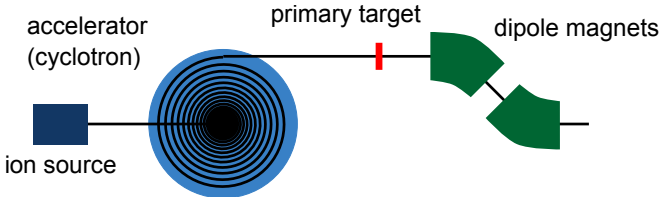
- projectile fragmentation of heavy beams on a light primary reaction target



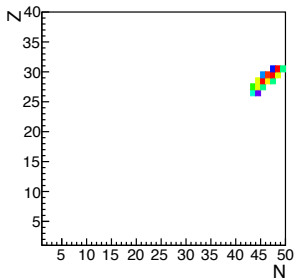
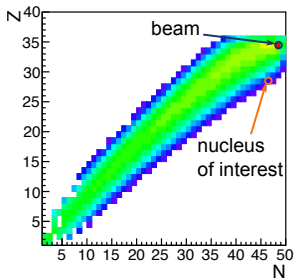
- need to select the ion of interest
 - using dipole magnets and slits
 - Z specific energy loss in a wedge shaped degrader



- projectile fragmentation of heavy beams on a light primary reaction target

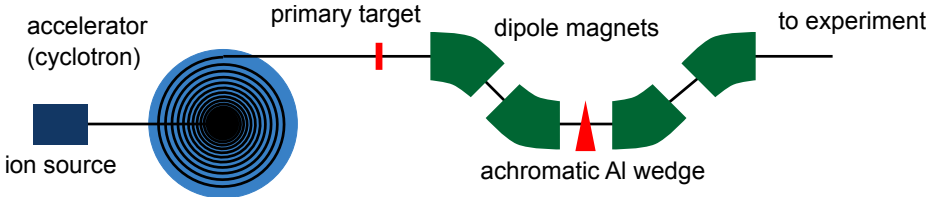


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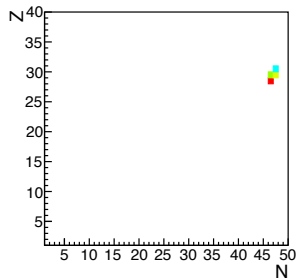
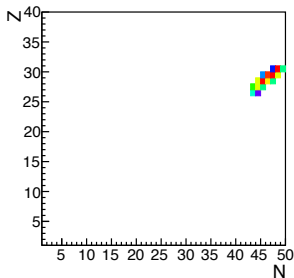
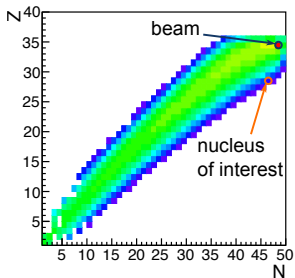


Making radioactive nuclei

- projectile fragmentation of heavy beams on a light primary reaction target



- need to select the ion of interest
- using dipole magnets and slits
- Z specific energy loss in a wedge shaped degrader



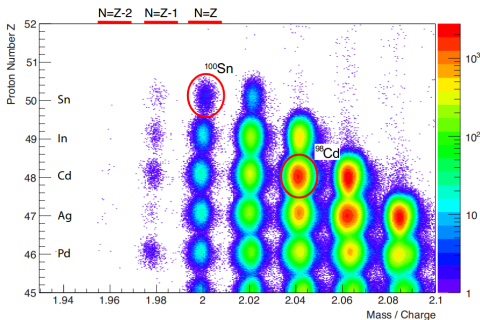
■ $B\rho - \Delta E - TOF$ method

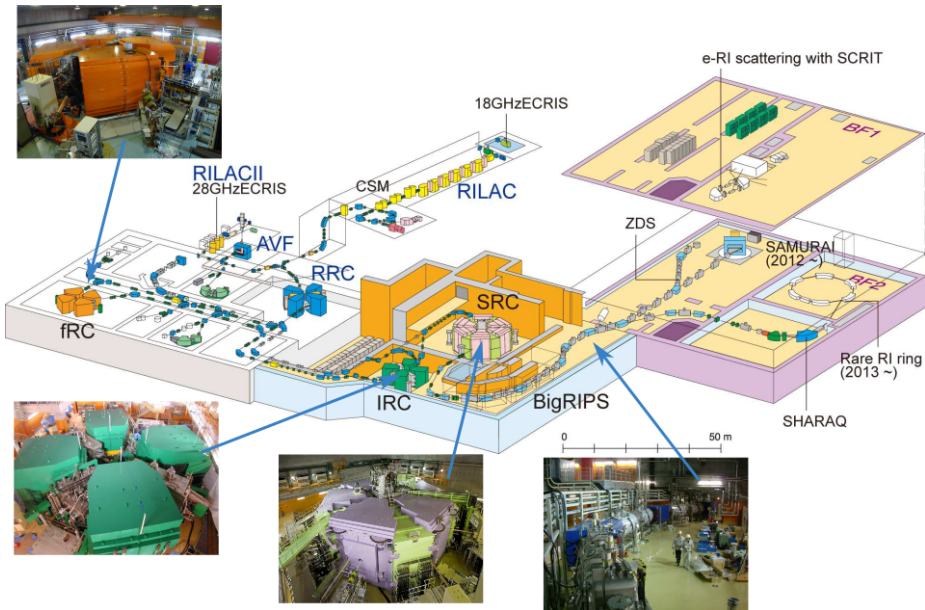
$$\frac{dE}{dx} = \frac{4\pi e^4 Z^2}{m_e v^2} N z \left[\ln \left(\frac{2m_e v^2}{I} \right) - \ln(1 - \beta^2) - \beta^2 \right]$$

$$TOF = \frac{L}{\beta c}$$

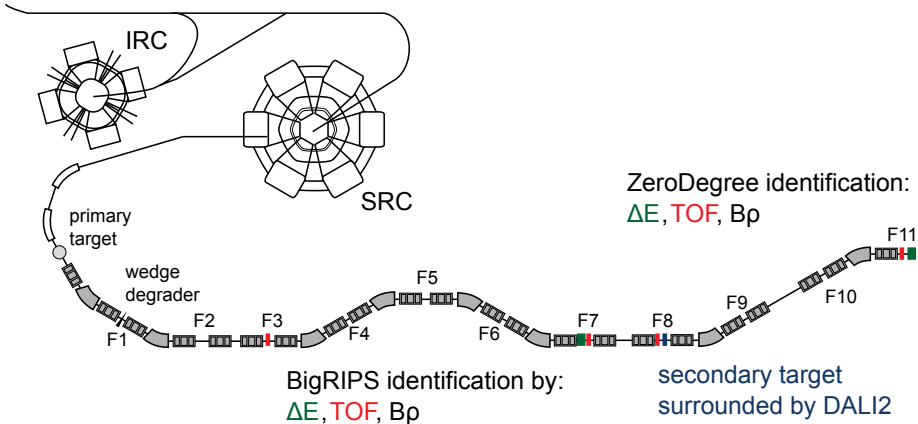
$$\frac{A}{Q} = \frac{B\rho}{\beta\gamma} \frac{c}{m_u}$$

- measure trajectory (position and angles at focal planes), time-of-flight, and energy loss
 → unique A/Q and Z identification

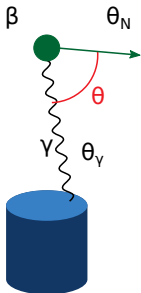




- 345 AMeV primary beams fragmented on a Be target
- 2-stage separation and identification
- decay and in-beam γ -ray spectroscopy

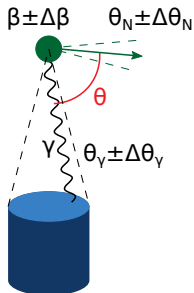


- γ rays are emitted in flight from a moving nucleus with velocity β
- Doppler shift depending on the emission angle in the lab system



$$E_{\text{lab}} = E_0 \frac{\sqrt{1 - \beta^2}}{1 - \beta \cos \theta}$$

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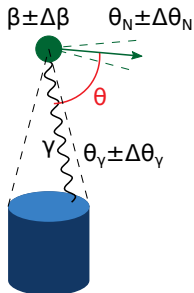


$$E_{\text{lab}} = E_0 \frac{\sqrt{1 - \beta^2}}{1 - \beta \cos \theta}$$

uncertainties lead to Doppler broadening

- velocity spread $\Delta\beta$
- angle of particles $\Delta\theta_N$
- detector opening angle $\Delta\theta_Y$

- γ rays are emitted in flight from a moving nucleus with velocity β
- Doppler shift depending on the emission angle in the lab system

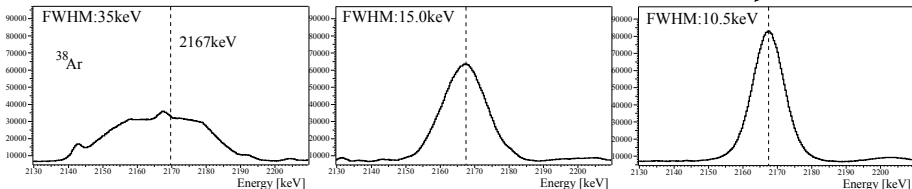


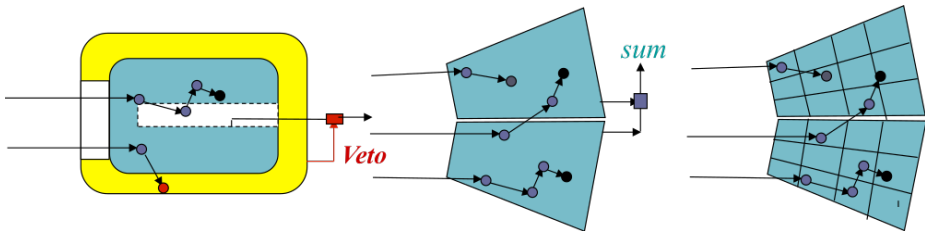
$$E_{\text{lab}} = E_0 \frac{\sqrt{1 - \beta^2}}{1 - \beta \cos \theta}$$

uncertainties lead to Doppler broadening

- velocity spread $\Delta\beta$
 - angle of particles $\Delta\theta_N$
 - detector opening angle $\Delta\theta_\gamma$
- ideally small detectors to minimize $\Delta\theta_\gamma$

decreasing effective detector size





$N \sim 100$

$N\Omega\epsilon \sim 0.1$

limited efficiency,
 suppressor cover space

$N \sim 1000$ (small)

$N\Omega\epsilon \sim 0.6$

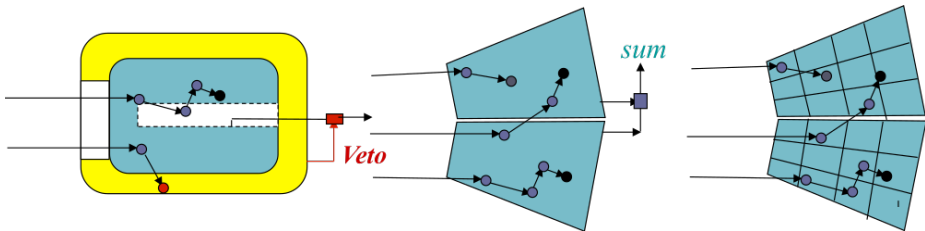
many detectors

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with segmentation to
 reconstruct position

- ideal: a 4π sphere of highly segmented Ge detectors
- allows to sum Compton events (instead of rejecting)
- signal decomposition for sub-segment position resolution
 → a set of interaction points with (E_i, x_i, y_i, z_i)
- tracking: use the Compton and Klein-Nishina formulas to determine order
 → background rejection, Doppler-correction, polarization



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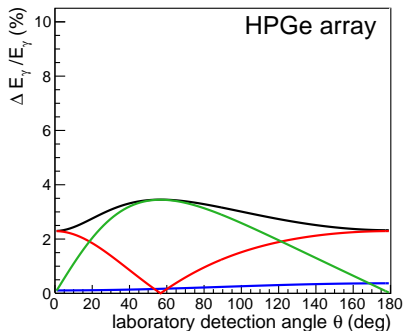
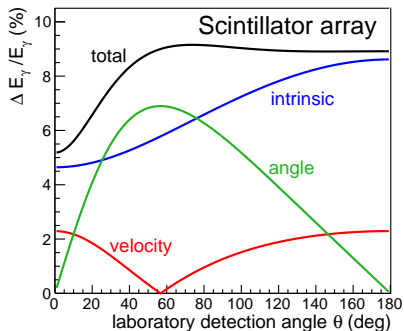
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- depends on intrinsic resolution
- uncertainty from velocity (where in the target happened the reaction)

$$\left(\frac{\Delta E}{E}\right)_{\beta} = \frac{\cos \theta - \beta}{(1 - \beta \cos \theta)(1 - \beta^2)} \cdot \Delta \beta$$

- position resolution \rightarrow signal decomposition and tracking

$$\left(\frac{\Delta E}{E}\right)_{\theta} = \frac{\beta \sin \theta}{(1 - \beta \cos \theta)} \cdot \Delta \theta$$

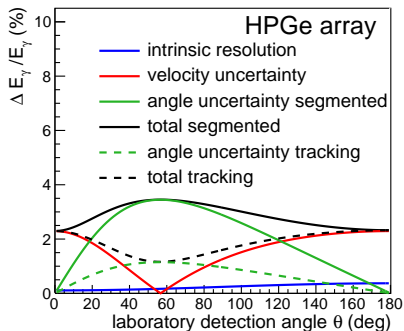
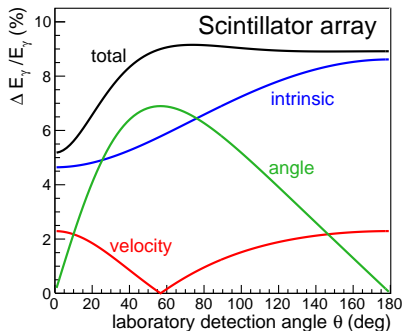


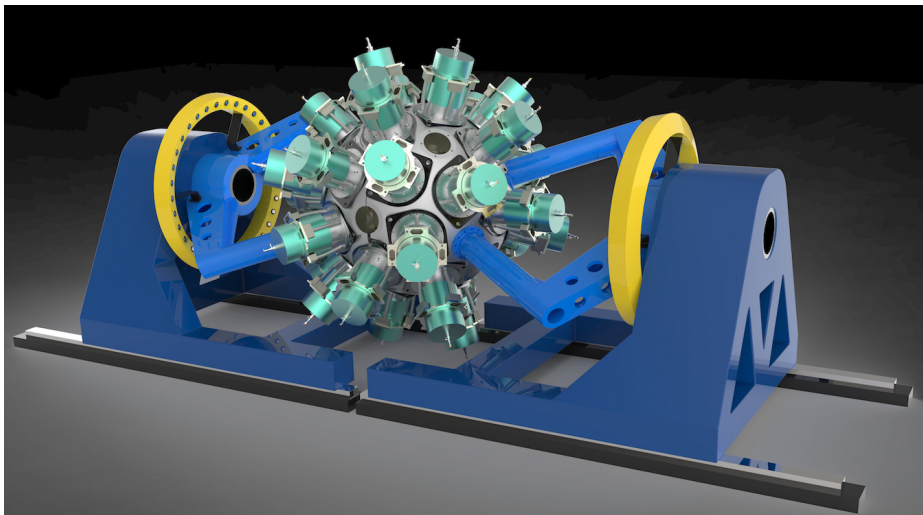
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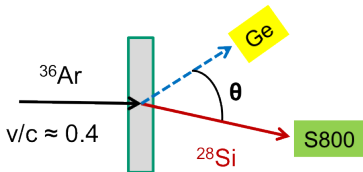
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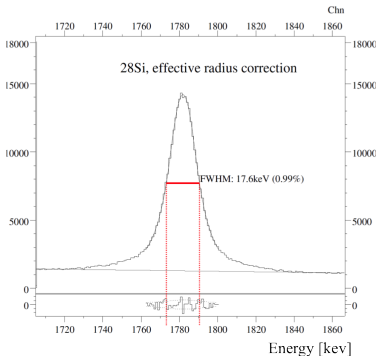




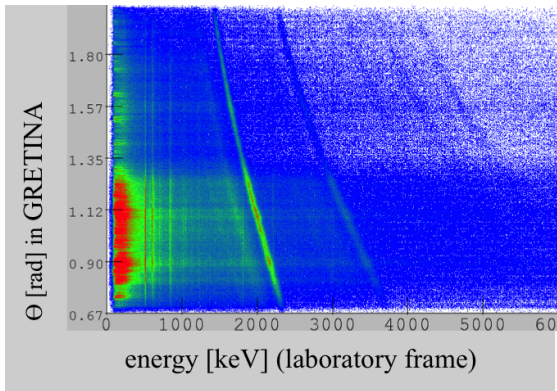
- GRETA: 4π array of 120 HPGe detectors with 36 segments each (USA)
- AGATA: Advanced GAMMA Tracking Array in Europe



- ^{28}Si from ^{36}Ar on a 47 mg/cm^2 Be target
- $v = 0.38 c$
- measured FWHM 1.00 % meets expectations
- 2 mm position resolution



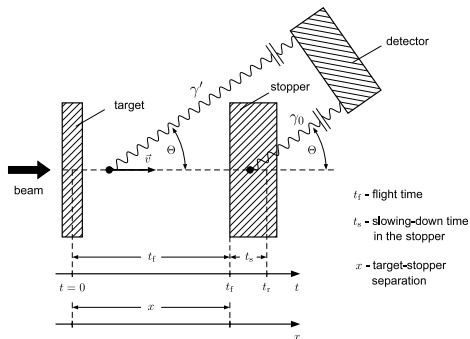
from D. Weisshaar et al.



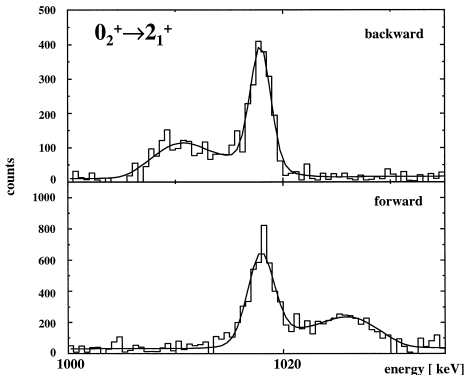
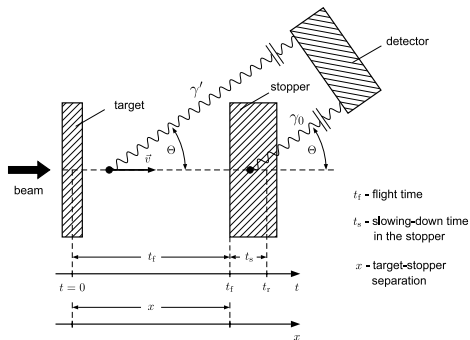
How to measure the lifetime of an excited nuclear state

- excited states produced in the target decay in flight
- measure distance instead of time
- place a stopper a certain distance after the target
- two components to the spectrum: shifted (in-flight) and stopped

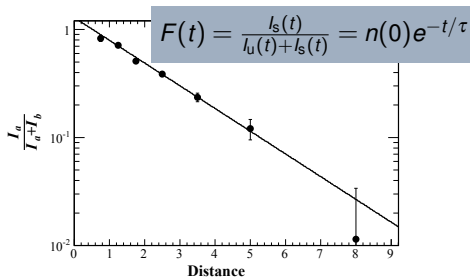
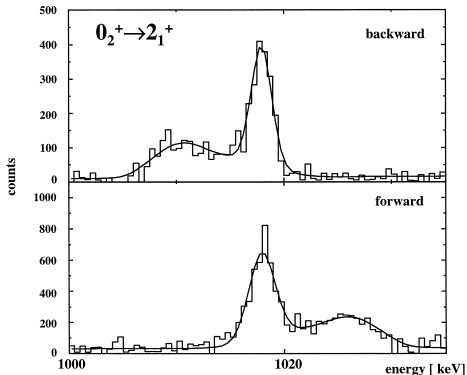
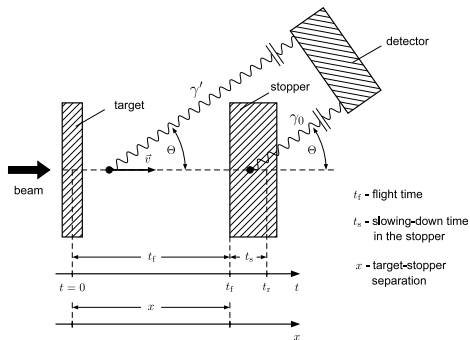
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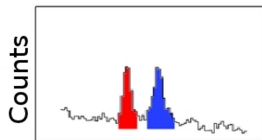
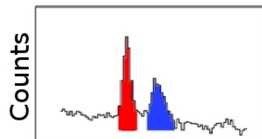
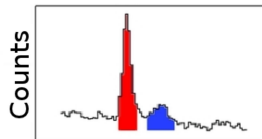
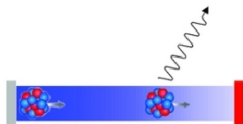
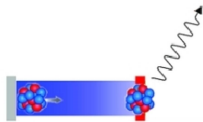
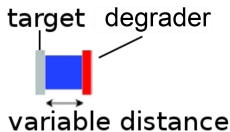
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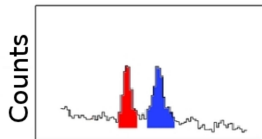
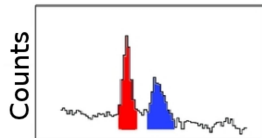
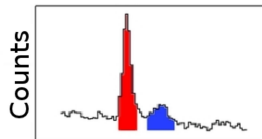
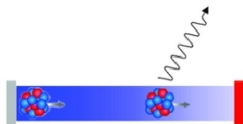
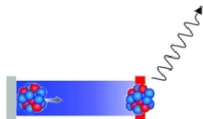
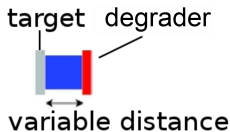


- you do not want to stop the beam in your target chamber, this will create a lot of background from decay
- use a degrader instead of a stopper
 - only change velocity β
- two different emission velocities, two peaks in spectrum
- vary distance and measure the relative yields of the two peaks
 - decay curve → lifetime τ



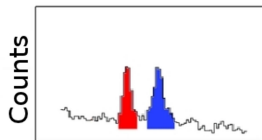
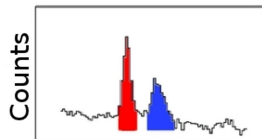
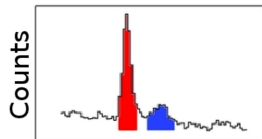
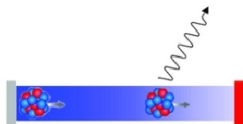
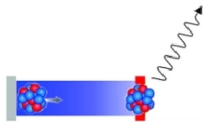
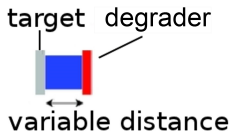
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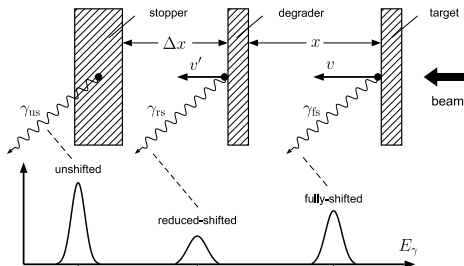
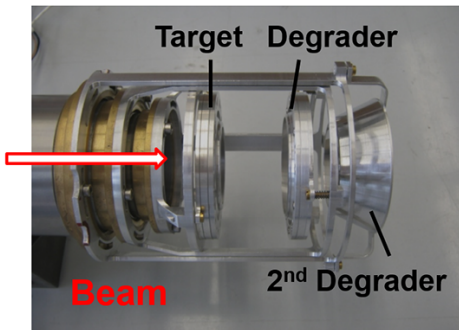
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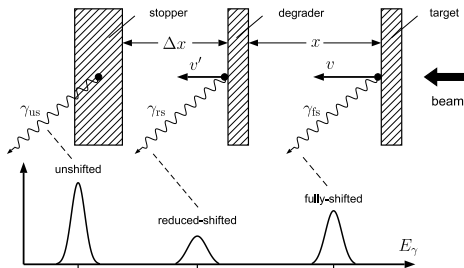
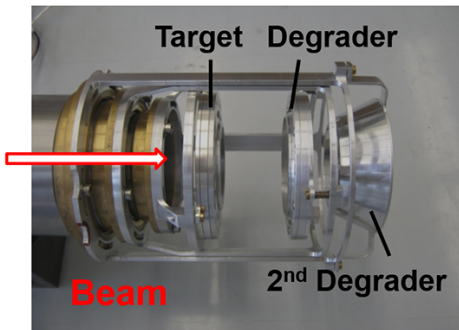
Energy

- the beam intensity is low, beam time is scarce
- long measurements with many distances are expensive
- using target and two degraders
- fast, reduced and slow (or stopped) peaks
- measure derivative of decay curve

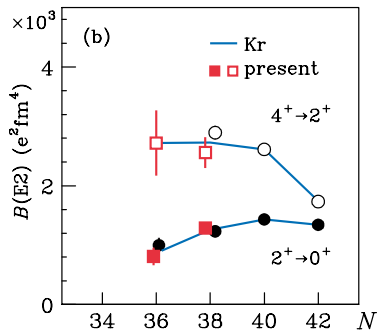
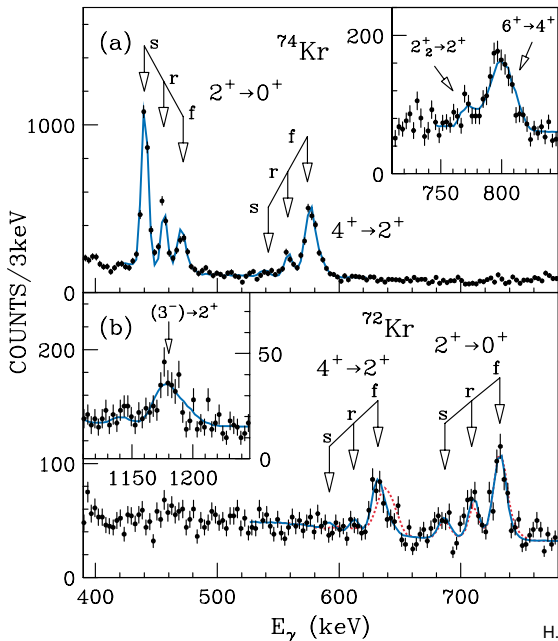
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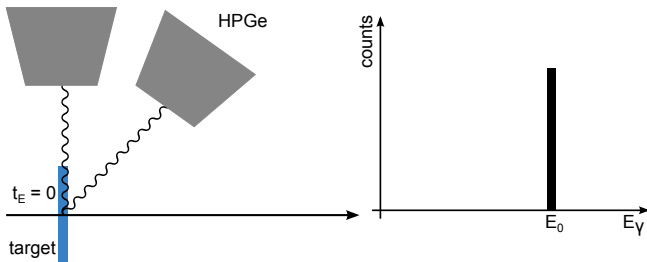


$$\tau = \left| \frac{\exp(-t/\tau)}{\frac{d}{dt} \exp(-t/\tau)} \right| \approx \frac{\Delta x l_s}{v l_r}$$



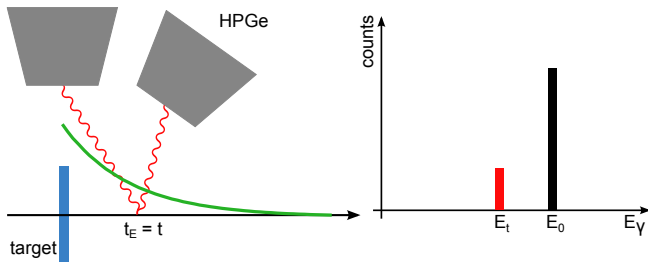
- short lifetime of 4^+ state in ^{72}Kr
- large $B(E2; 4^+ \rightarrow 2^+)$
- indicates shape transition in the yrast band
- oblate ground state, prolate for higher spins

- method to extract level lifetimes without a degrader
- emission at target position \rightarrow Doppler correction gives correct energy E_0



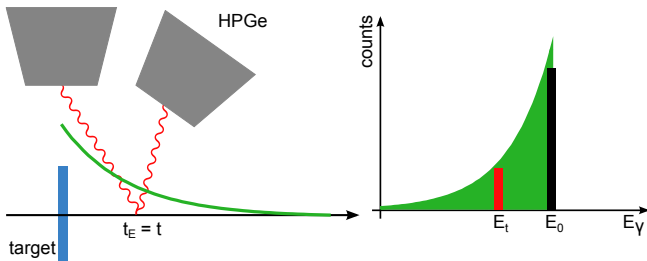
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- warning: this is simplified, some nuclei decay in the target while the nucleus is slowing down

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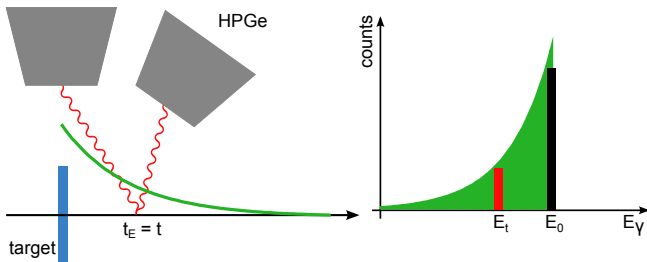
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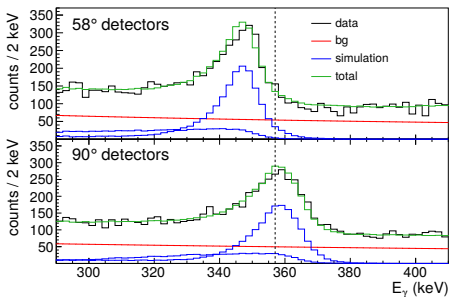
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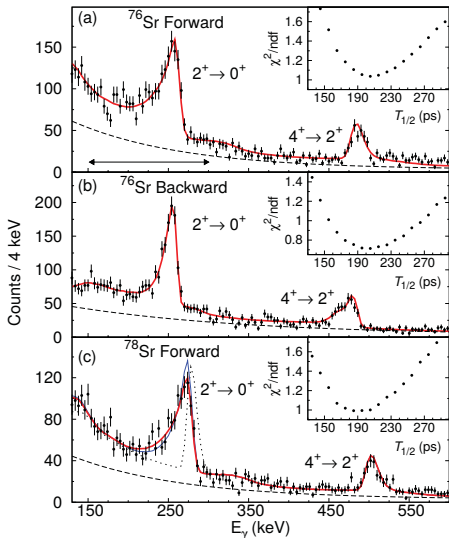
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 - shift in peak position (short lifetimes) and shape (long lifetimes)



K. Wimmer et al., NSCL experiment.

- lifetime determination through χ^2 fit with simulation
- lifetime measurement at $N = Z = 38$: ^{76}Sr
- $B(E2)$ values in $N = Z$ nuclei
- mutual enhanced collectivity in self-conjugate ^{76}Sr

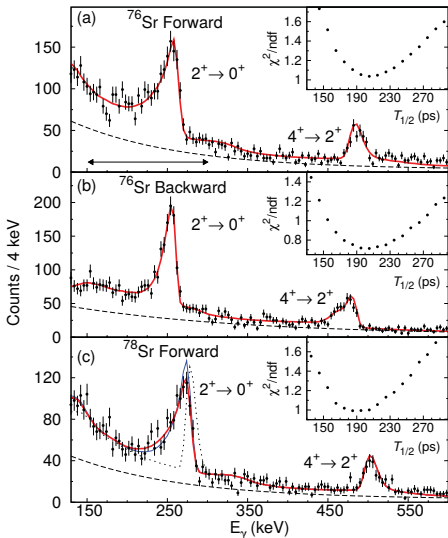
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A. Lemasson et al., Phys. Rev. C **85** (2012) 041303.

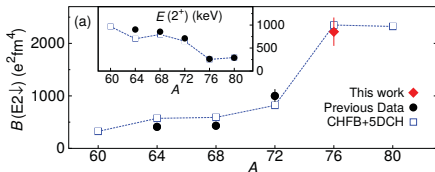
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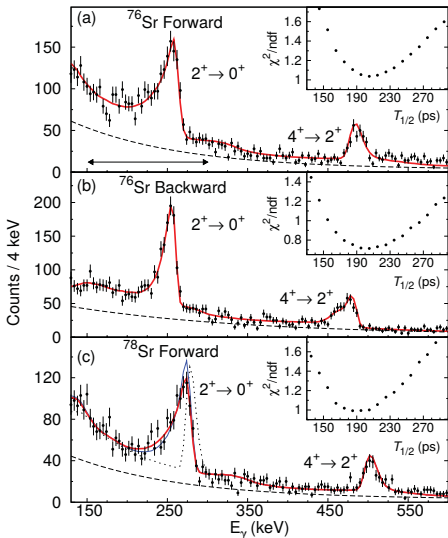


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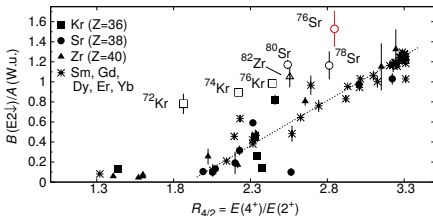


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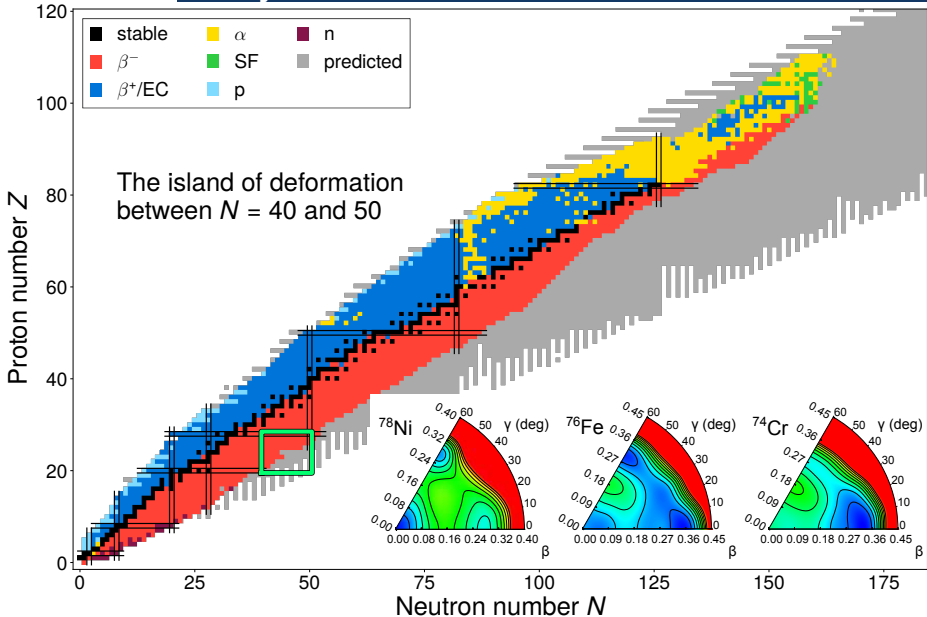
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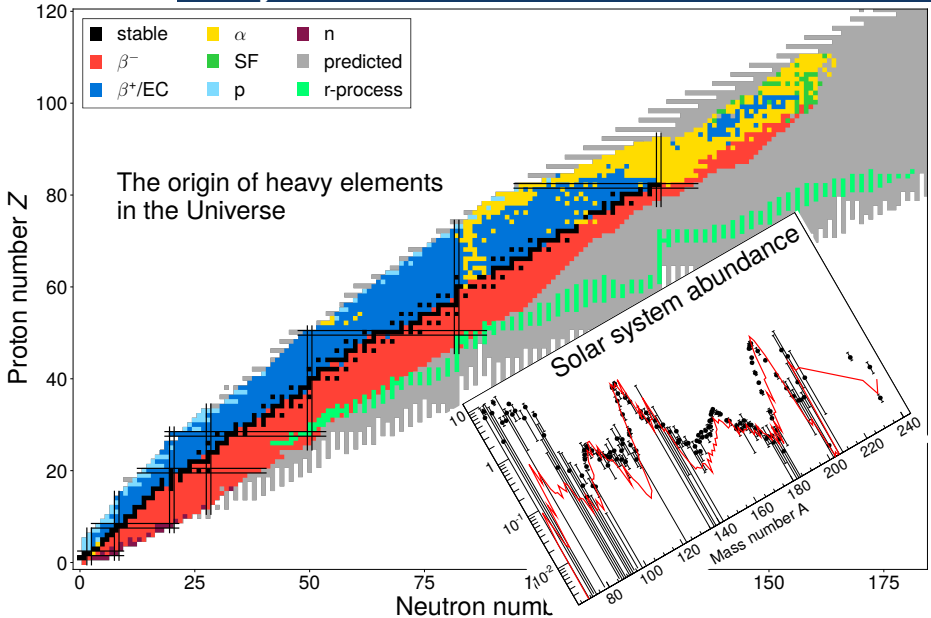
LISA:

Lifetime measurements with Solid Active targets

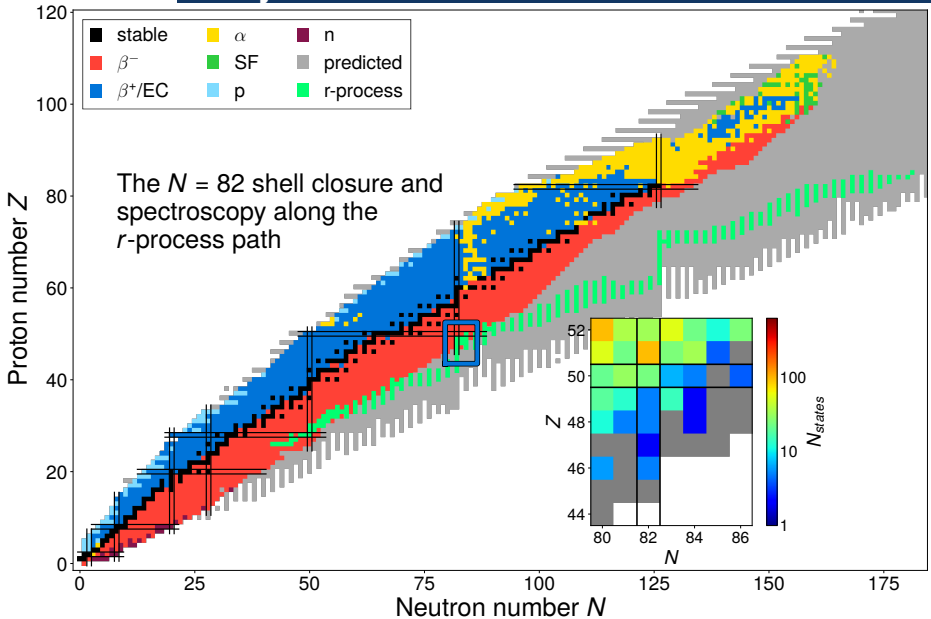


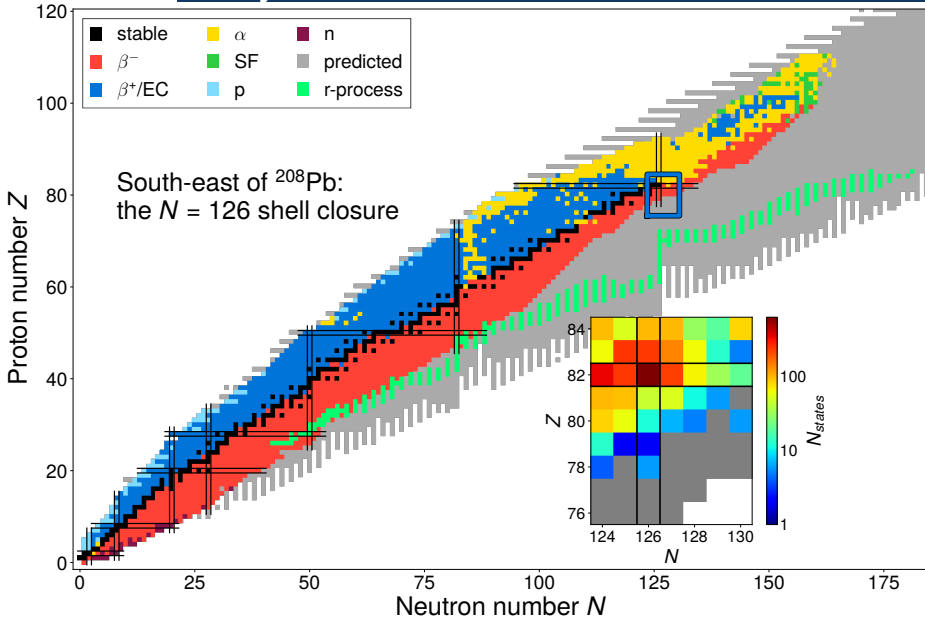


F. Nowacki et al., Phys. Rev. Lett. **117** (2016) 272501.



G. Lorusso et al., Phys. Rev. Lett. **114** (2015) 192591.





- γ rays emitted at a velocity β are subject to the Doppler effect

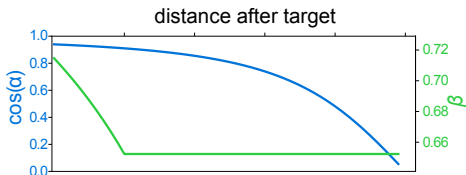
$$E_{\text{lab}} = E_0 \cdot \frac{\sqrt{1 - \beta^2}}{1 - \beta \cos \alpha}$$

- need to know velocity β and emission angle α to reconstruct transition energy E_0 from measured E_{lab}

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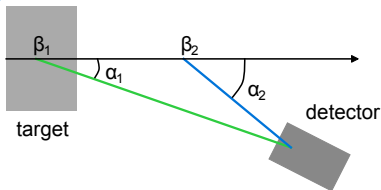
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a finite lifetime affects both

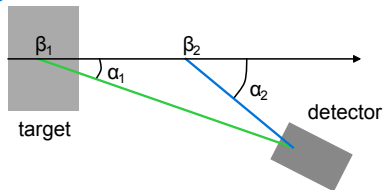
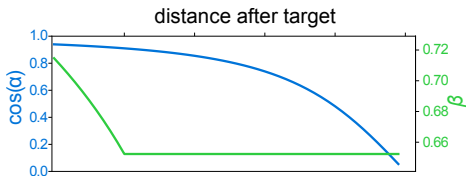
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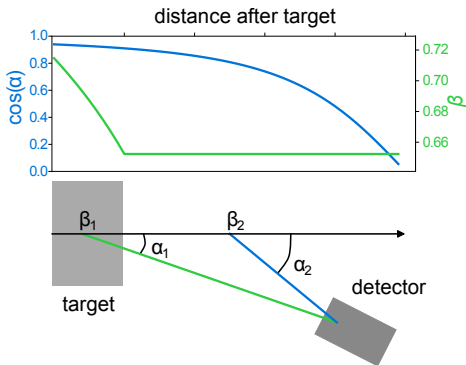
a finite lifetime affects both

- ejectile loses energy in target, smaller velocity β
 - decay further downstream, emission at larger angle α
- change in Doppler reconstructed energy
 → sensitive to the lifetime of states

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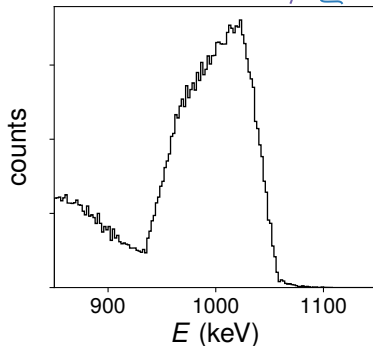
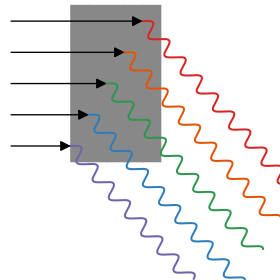
a finite lifetime affects both

- ejectile loses energy in target, smaller velocity β
 - decay further downstream, emission at larger angle α
- change in Doppler reconstructed energy
 → sensitive to the lifetime of states

need very good energy and position resolution for γ rays

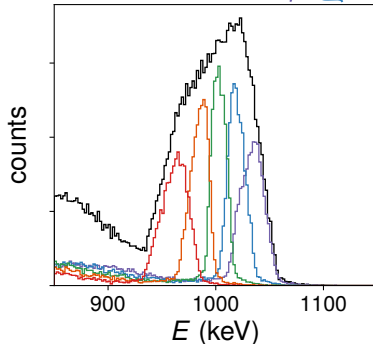
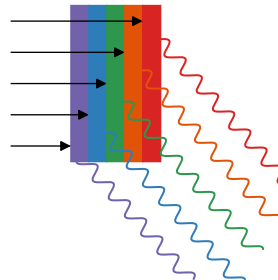
with the Advanced Gamma-ray Tracking Array AGATA, the achievable resolution is now given by the target ($\sim\text{g}/\text{cm}^2$ or few mm thick)

- reaction and emission at different velocities
- (angle dependent) spread in Doppler reconstructed spectrum
- different mean decay velocities and different depths in the target
- with an active target resolution can be greatly improved
- use proper position and velocity for Doppler reconstruction



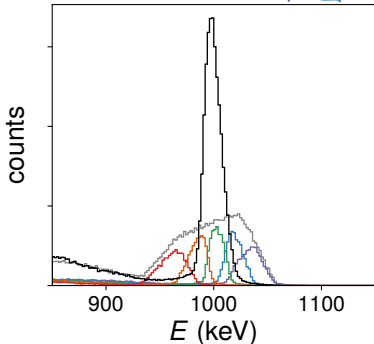
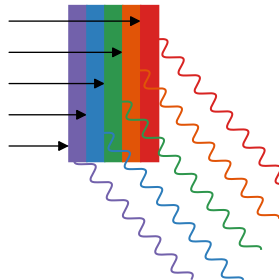
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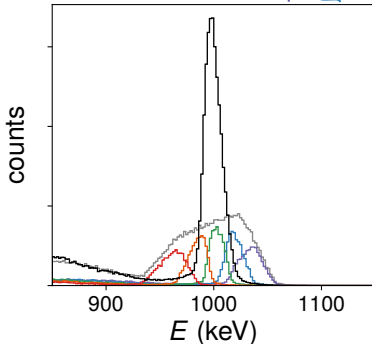
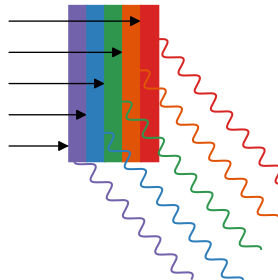
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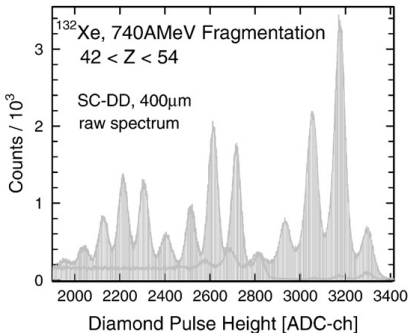
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Lifetime measurements with Solid Active targets

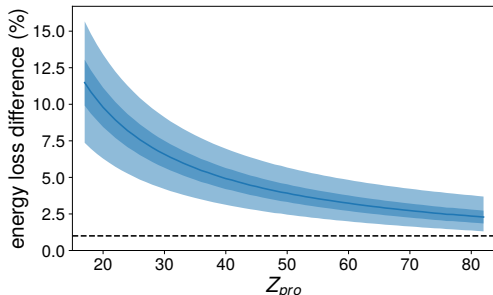


- proton removal, change in $Z \rightarrow$ different energy loss
- with single-crystalline diamond a unique Z identification is possible



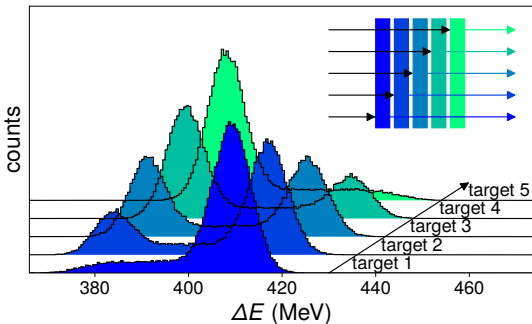
E. Berdermann et al., Dia. Rel. Mat. 17 (2008) 1159.

differential energy loss for proton removal at 200 and 800 AMeV



- required resolution for Z identification is few %
- new technology based on poly- and single-crystalline CVD diamond detectors

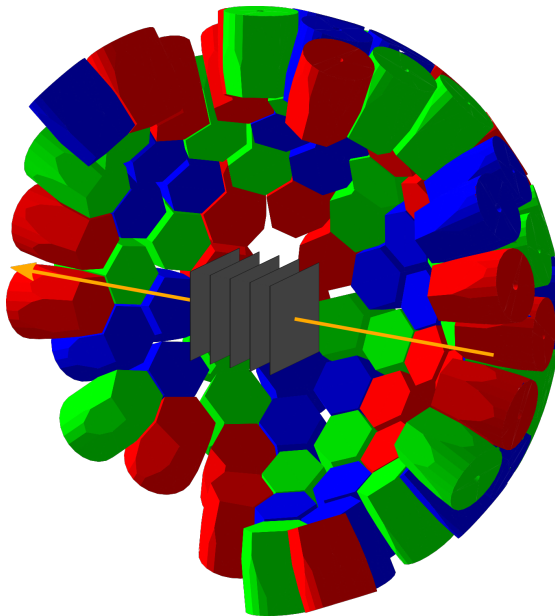
- each layer measures the Z -dependent energy loss of the beam-like particles
- event-by-event Z identification in each layer

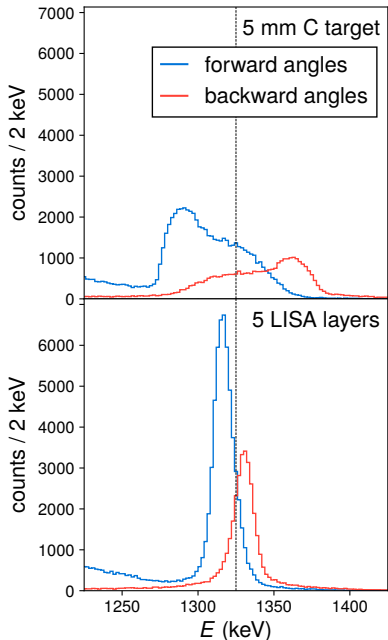


- proton removal reaction
 → change in Z and energy loss ΔE
- analysis with machine learning techniques
- simulation study:
 improved accuracy for layer identification from 70 to 90% for challenging cases compared to simple gating

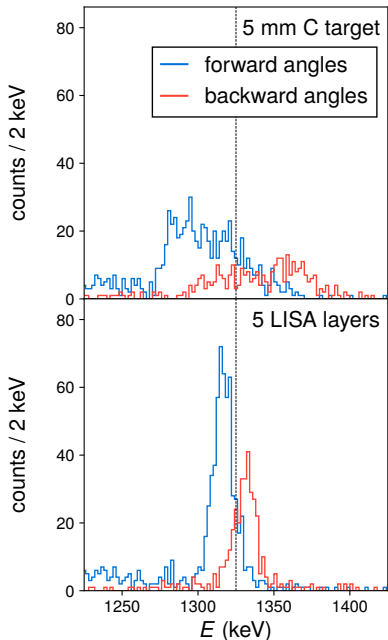
- Doppler-correction with position of reaction layer and corresponding velocity

match target position resolution to resolution of modern γ -ray tracking detectors for improved lifetime measurements

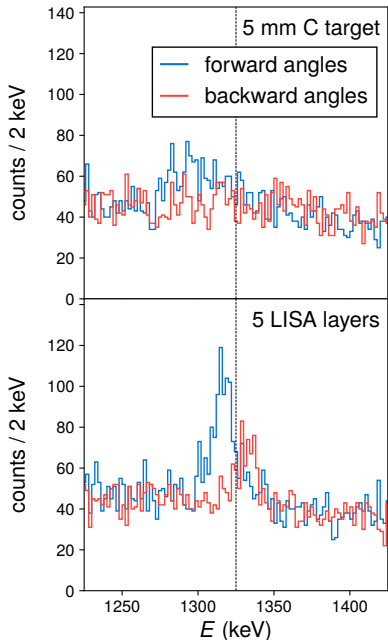




- Coulomb excitation challenging: low beam intensity, small cross section
- knockout from high-intensity ^{132}Sn beam
- assumption: 1 W.u. $\rightarrow \tau = 5$ ps
- decay in target: shift in energy for forward and backward laboratory angles
- LISA improves resolution significantly (5 targets, equivalent thickness)
 - with realistic statistics measurement of shift becomes challenging
 - including a typical background, only LISA has sufficient peak-to-background ratio to measure the lifetime



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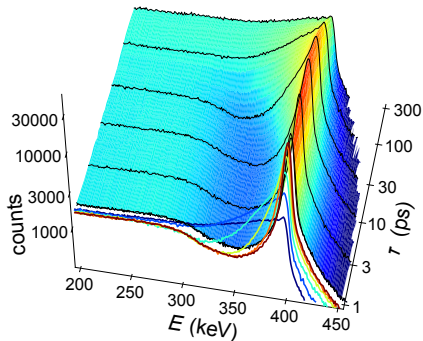


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- probe underlying symmetries and dynamics of nuclei

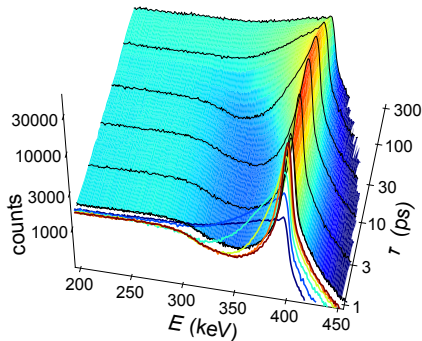
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- LISA is an active target system that matches the excellent energy resolution given by tracking spectrometers
- using solid active targets to determine the reaction point
- lifetime measurements to explore collectivity in Islands of Inversion and shell closures in exotic nuclei



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Lots of opportunities for future experiments