

2021 Winter Topical Meeting on VBS

January 28th, 2021

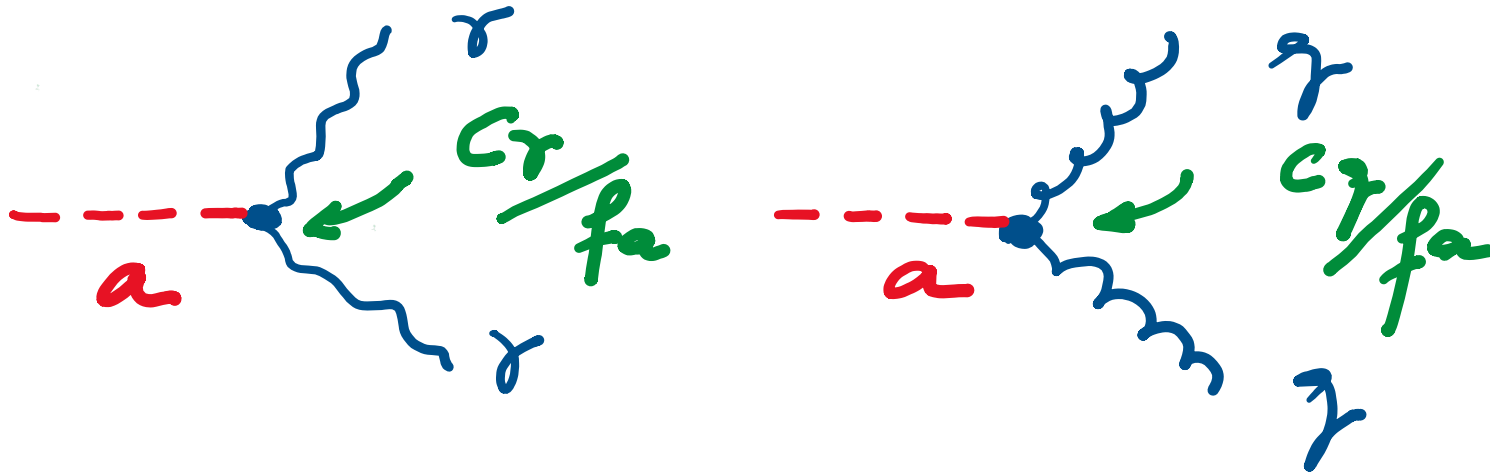
Non-Resonant ALP Searches at the LHC: Implications for VBS

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Theoretical / Phenomenological Background

- **ALPs (Axion-like Particles)** are well motivated theoretically as neutral pseudo-scalar Pseudo-Goldstone Bosons (PGB) of a new spontaneously broken symmetry. Examples: axions, technipions.
- ALP interactions parameterized with a general **Effective Field Theory Lagrangian**, consistent with SM gauge symmetries and CP. Two implementations of EFTs: linear (related to weakly coupled new physics models, minimal) and chiral (related to strongly coupled new physics models, more parameters). In this talk we focus on the linear EFT.



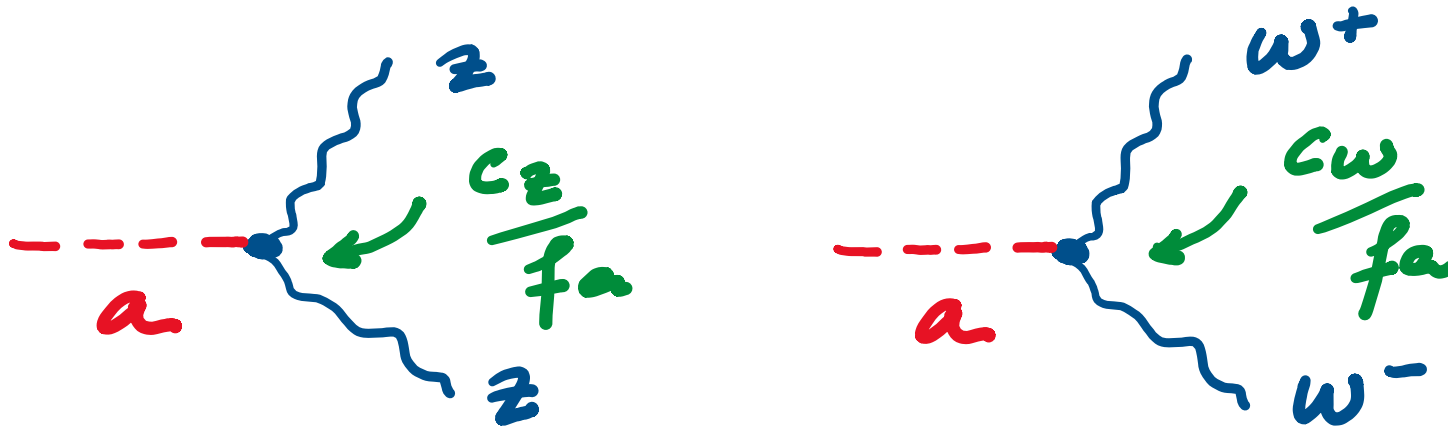
- **ALP interactions are derivative:** they grow with momentum; couplings are proportional to Wilson coefficient c_i and inversely proportional to new physics energy scale f_a .

Theoretical / Phenomenological Background

- Colliders allow searches in a wide range of ALP masses and couplings. We can explore ALP masses beyond astrophysical constraints, and even there, provide important crosschecks. **At the LHC, natural sensitivity is to f_a scales in the TeV region.**

Classical Searches

- Classical searches for ALPs at colliders consider its **couplings to photons and gluons: c_γ and c_g** . More recently, interest in this area has extended to consider ALP **couplings to EWK-bosons: ZZ, WW and Z gamma**. At LO all these and the coupling to photons are related by gauge invariance to two basic EWK couplings: **c_W and c_B** .

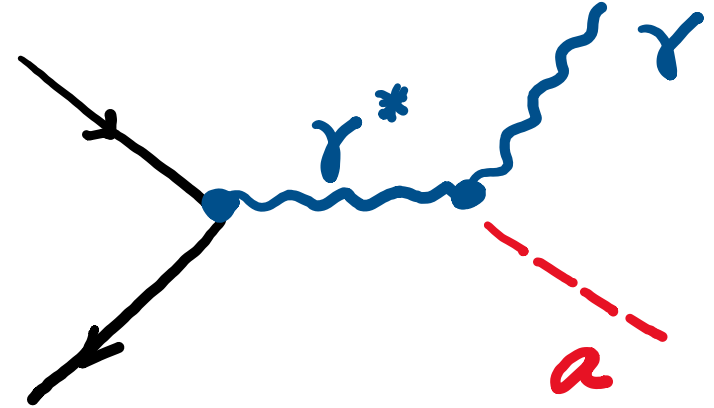


- For ALPs with light masses, they can be considered stable (i.e., can go undetected => MET), use the mono-X signature. Mono-X signatures allow setting limits to ALP couplings “one-by-one”.

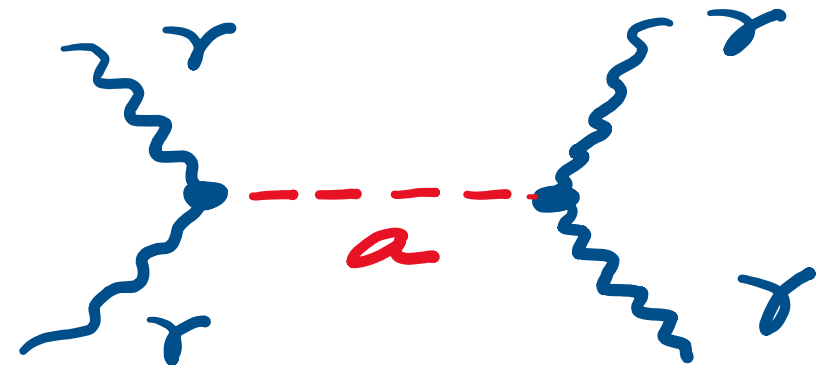
Classical Searches

- Examples:
 - **Mono-photon**, gamma + (a => MET);
 - **Mono-jet**, jet + (a => MET);
 - **Mono-Z, mono-W, mono-H.**

At the LHC, these channels are covered in ATLAS/CMS EXO searches for DM signals.



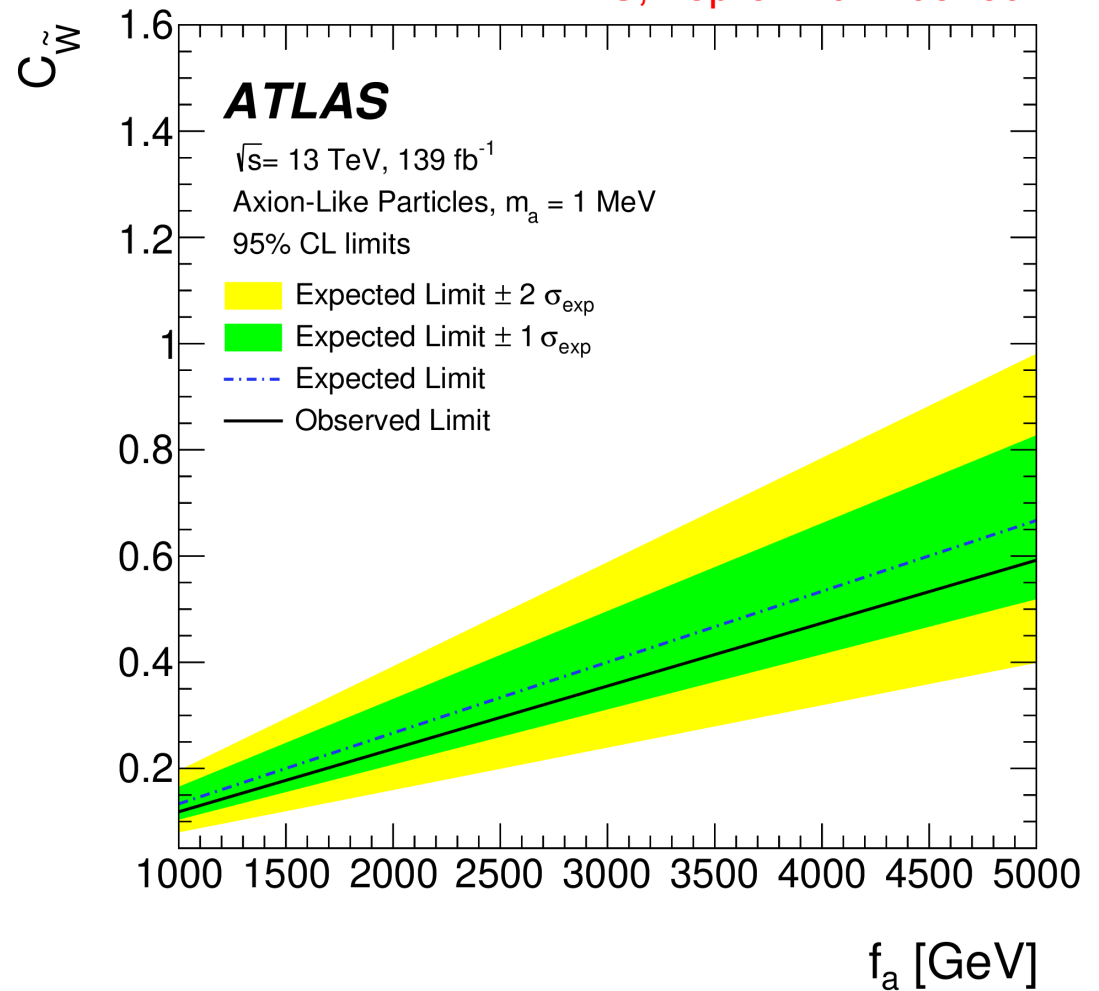
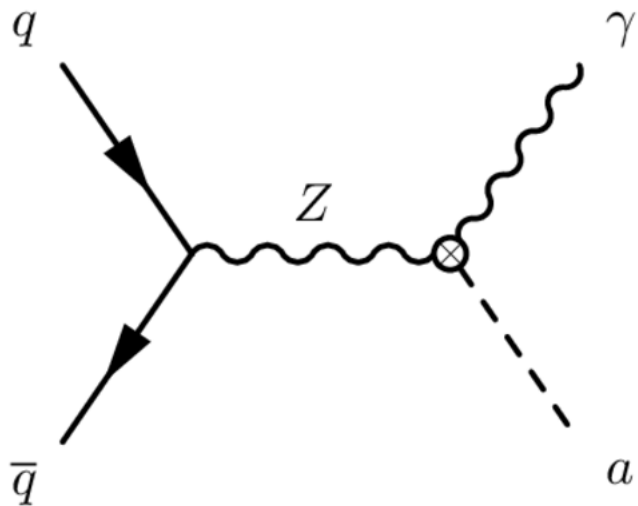
- Alternatively, ALPs can have large masses, where they are assumed to have widths wide enough to decay promptly in the detector; yet narrow enough with respect to experimental resolution.
- Examples:
 - **Dijet and diphoton** resonances.
 - **Light-by-Light scattering**: on-shell ALP production, a => 2 photons.



ATLAS Mono-Photon Limits

ATLAS; hep-ex 2011.05259

- ATLAS, 139 fb⁻¹
- ALP "detected" as MET
- 1 MeV < M_{ax} < 1 GeV
- photophobic



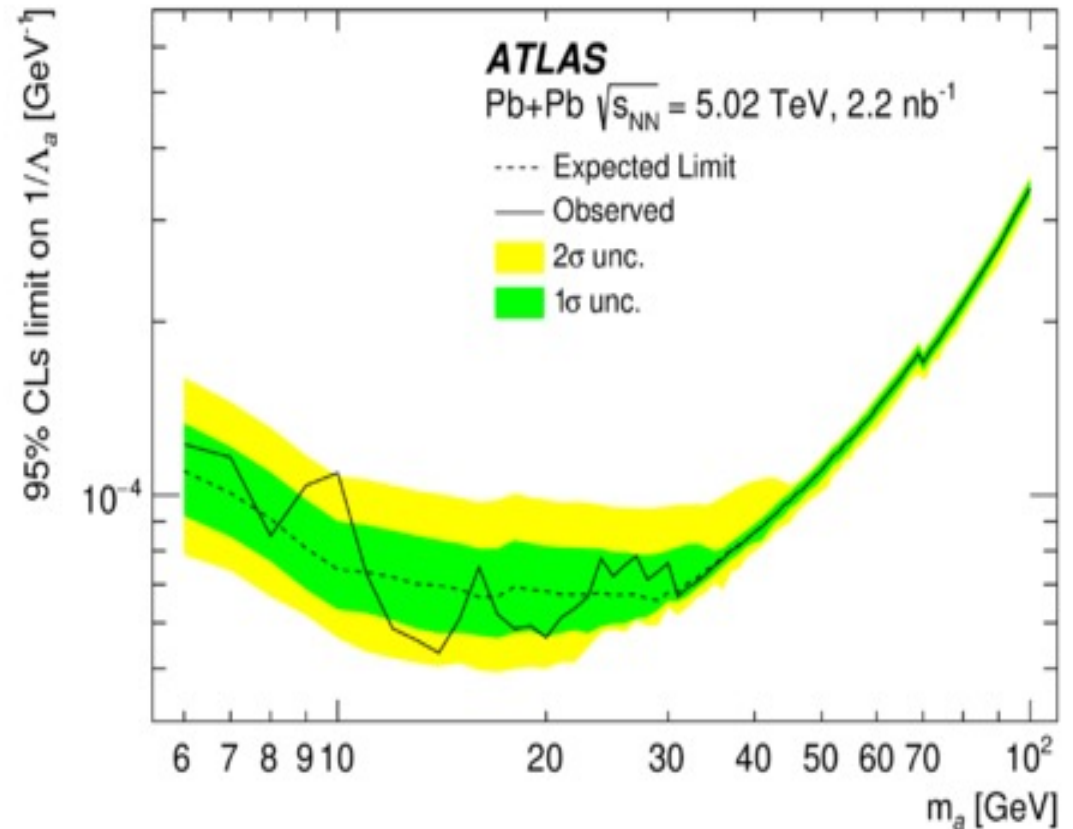
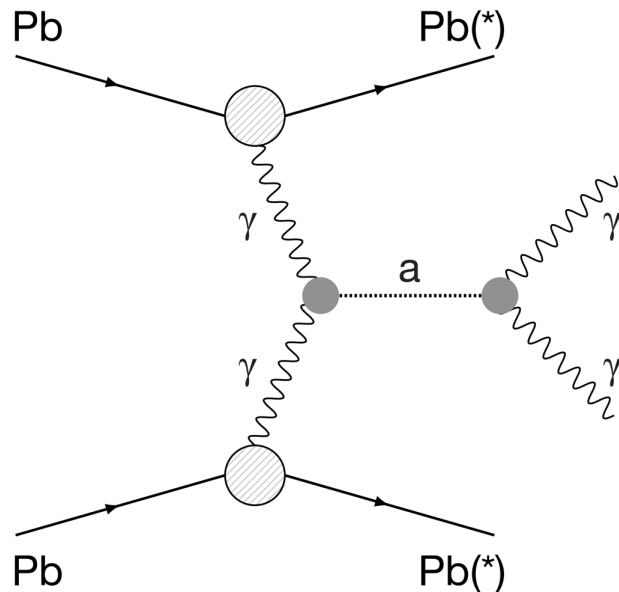
$f_a / c_W > 8$ TeV @ 95% CL,

for $c_{\text{gamma}} = 0, c_g = 0$

ATLAS / CMS LbyL On-shell ALP Limits

- For instance: ATLAS,
Pb+Pb, 2.2 nb⁻¹
- On-shell ALP production,
BR(gamma gamma) = 100%
- 6 GeV < M_ax < 100 GeV

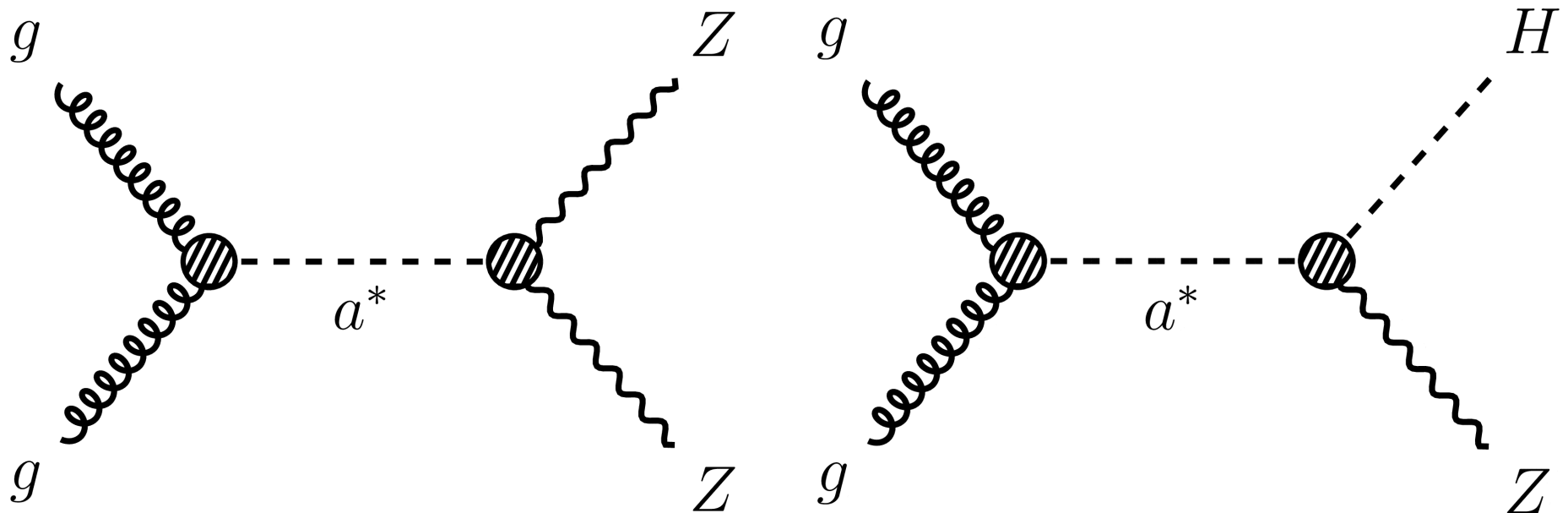
ATLAS; hep-ex 2008.05355



Limits on $4 c_{\text{gamma}} / f_a$ @ 95% CL,
for $c_g = 0$

GGF ALP-Mediated Processes

- Gluon-initiated ALP-mediated processes provide more possibilities to test the ALP universe.
- These channels are sensitive to the product of the ALP coupling to gluons times the coupling to EWK dibosons.



GGF Non-Resonant ALP-Mediated Processes

- **Off-shell ALP production.** This is very promising because the cross-sections are large enough to constraint significantly the theoretical models using Run 2 data.
- ALPs are **s-channel mediators** in $gg \Rightarrow VV$ production with $\hat{s} \gg M_a + \text{many } \Gamma_a$. The size of \hat{s} is enhanced by the mass threshold of the on-shell diboson system in the final state; but most importantly by the hard p_T -spectrum provided by the derivative couplings.
- The analysis uses the ZV, WW, ZH searches looking for **high- p_T / high-mass deviations** in the tails of the transverse momentum / mass spectra with respect to SM expectations.
- This works for ALPs light enough, and cross-sections and limits are independent of M_a **from the very-light limit, up to masses of order of 100 GeV.**

Cross-Section Estimates @ 13 TeV

□ Linear EFT: photophobic, $c_g/f_a = c_W/f_a = 1/\text{TeV}$, $M_a = 1 \text{ MeV}$

- $q\bar{q} \Rightarrow Z a$ 5 pb
- $gg \Rightarrow ZZ$ 40 pb (80 pb for $c_B = c_W$)
- $gg \Rightarrow WW$ 180 pb
- $gg \Rightarrow Z \gamma$ 60 pb
- $gg \Rightarrow \gamma\gamma$ 50 pb $c_B = c_W$; $M(\gamma\gamma) > 0.5 \text{ TeV}$

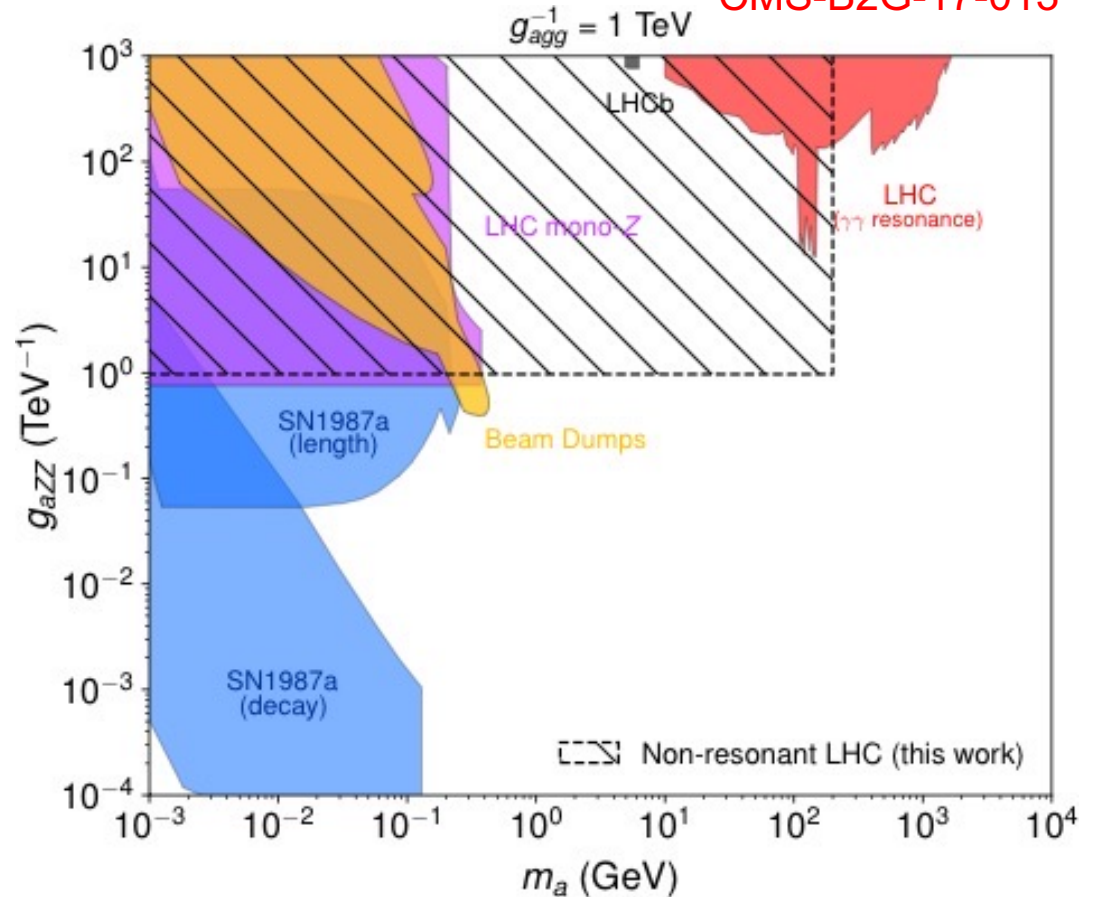
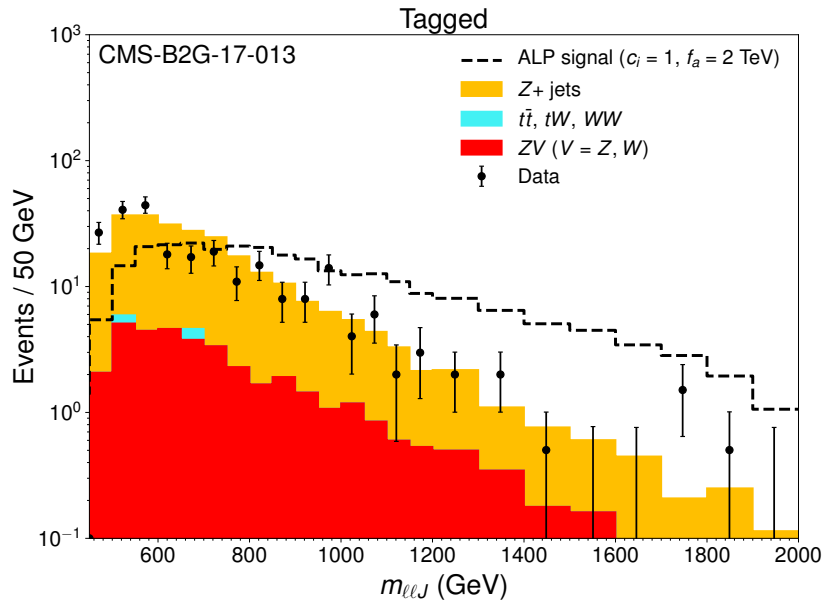
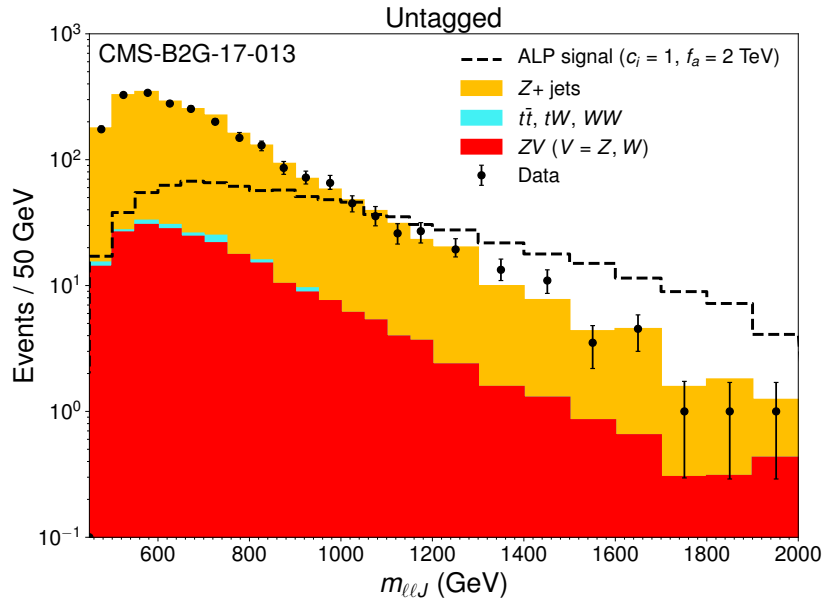
□ Chiral EFT: $a_{\tilde{2D}} = 1$, $c_g/f_a = 1/\text{TeV}$, $M_{ax} = 1 \text{ MeV}$

- $q\bar{q} \Rightarrow H a$ 0.2 pb
- $gg \Rightarrow ZH$ 70 pb

High-Mass ZZ 2L2Q Limits

Gavela, No, Sanz, JFT; hep-ph 1905.12953

CMS-B2G-17-013



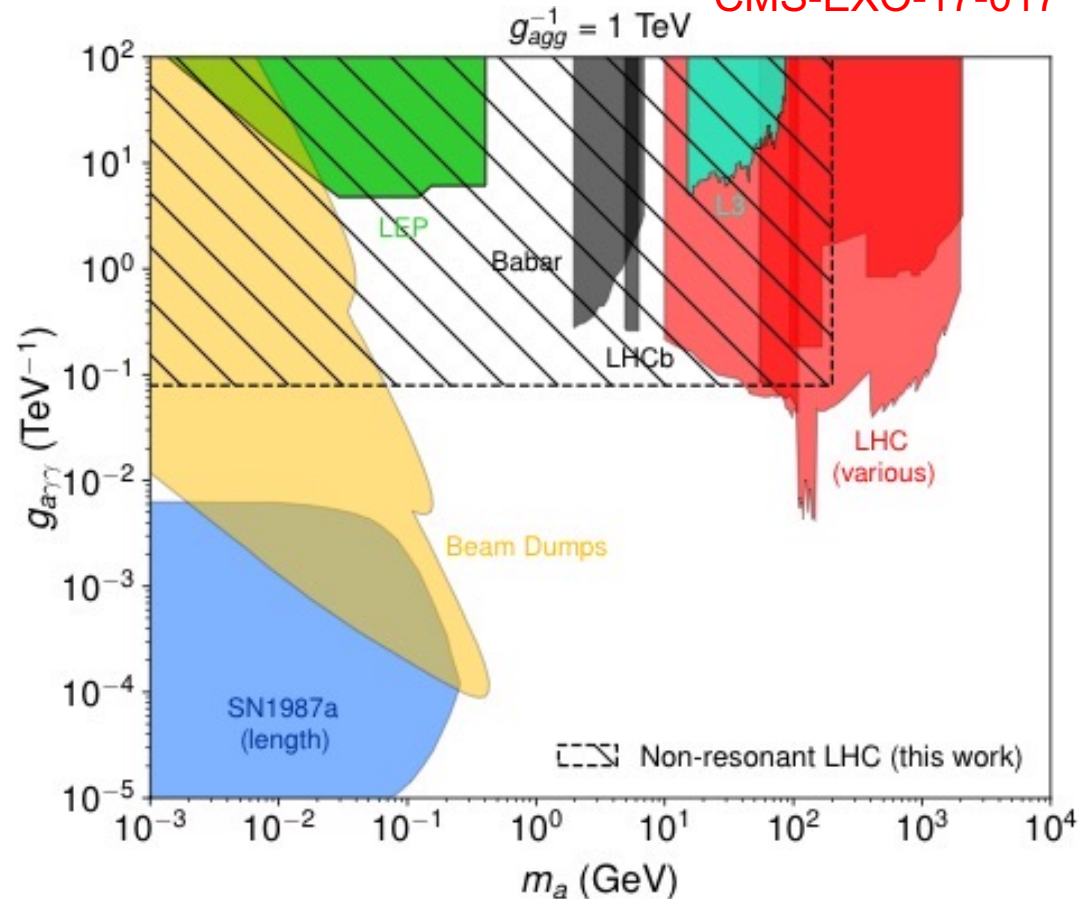
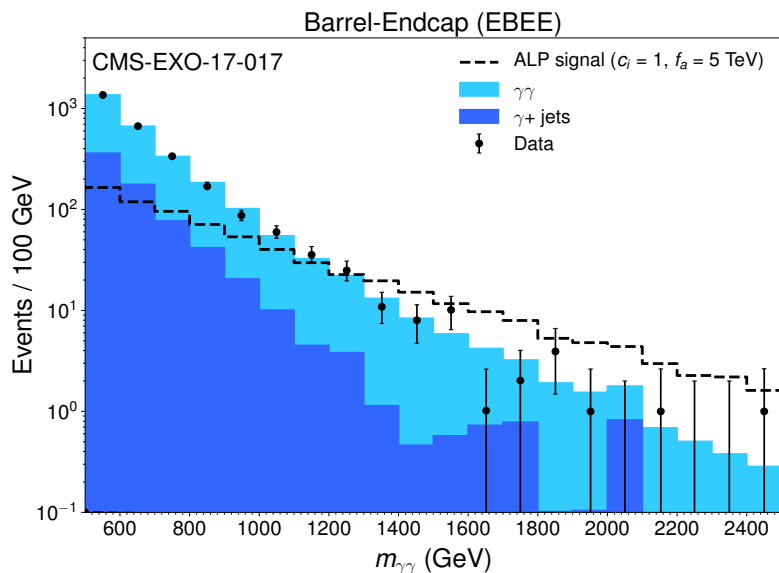
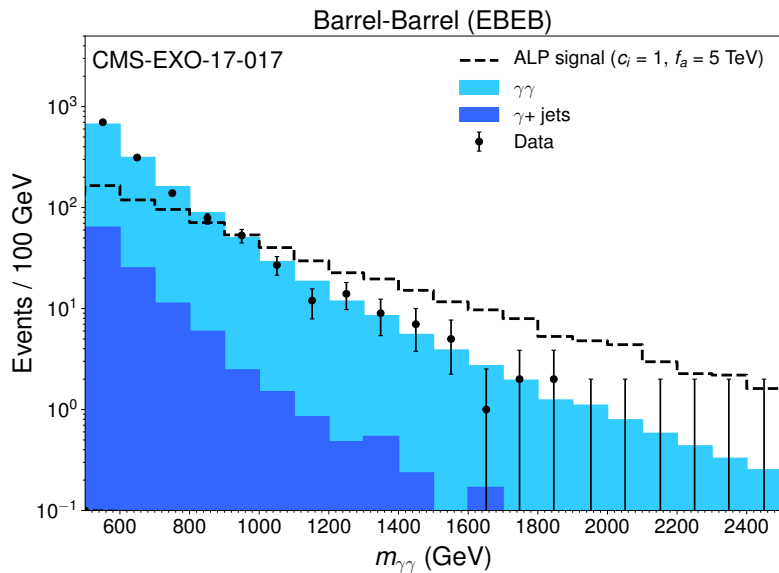
$$f_a / c_Z > 17 \text{ TeV @ 95\% CL,}$$

$$\text{for } f_a / c_g < 1 \text{ TeV}$$

High-Mass Diphoton Limits

Gavela, No, Sanz, JFT; hep-ph 1905.12953

CMS-EXO-17-017

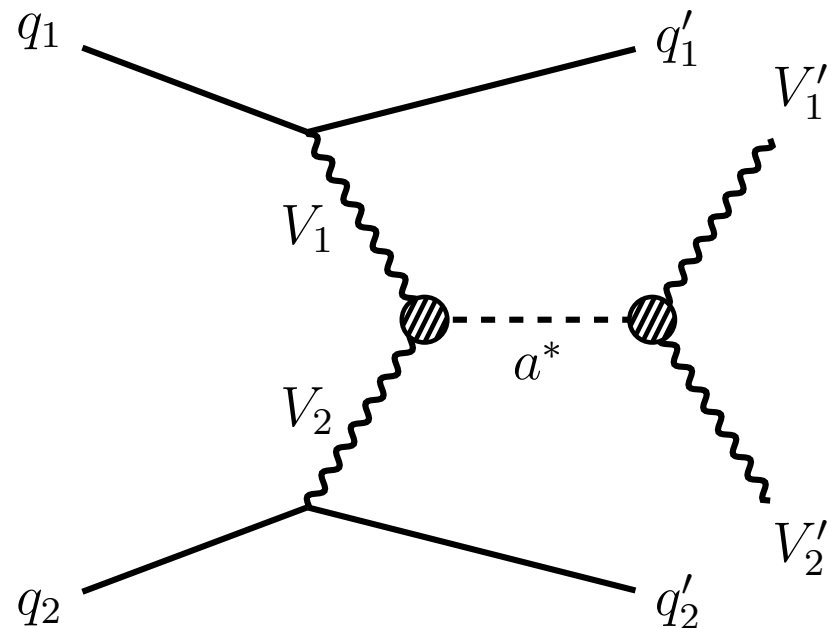
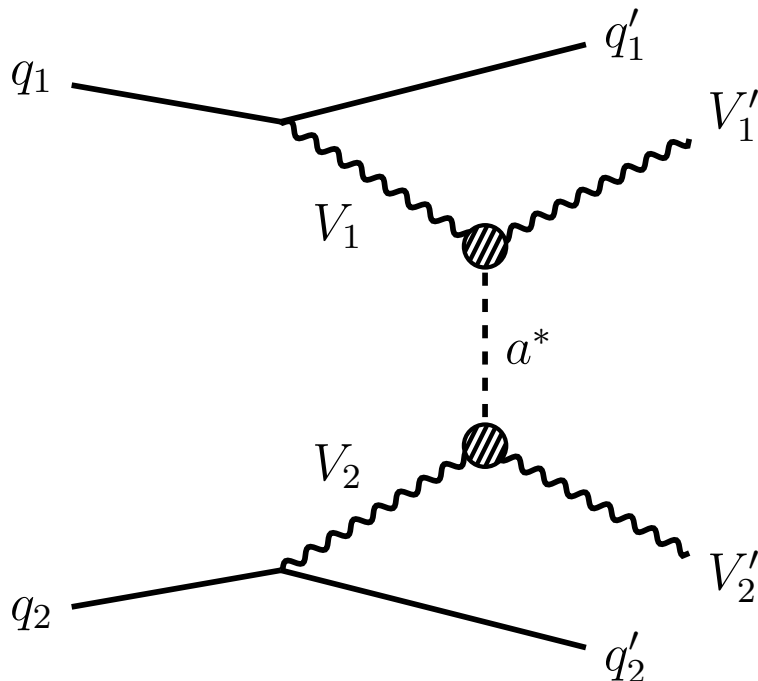


$$f_a / c_{\text{gamma}} > 200 \text{ TeV @ 95\% CL,}$$

$$\text{for } f_a / c_g < 1 \text{ TeV}$$

Non-Resonant ALPs in VBS

- Besides its intrinsic interest, **non-resonant ALP-mediated processes in VBS channels are useful because they provide limits on ALP couplings to vector bosons independently of the gluon coupling.**
- The VBS diagrams contain an **off-shell ALP** interchanged in the **t-channel**; **s-channel** is relevant for the ZZ and the Z gamma final states.



Non-Resonant ALPs in VBS

- Cross-sections and limits still independent of M_a from the **very-light limit, up to masses of order of 100 GeV.**
- Last but not least, ATLAS / CMS have recently published Run 2 measurements. Allows first comparison to the data, calibration of the simulation tools, calculation of educated predictions for higher luminosities.

Non-Resonant ALPs in VBS

Bonilla, Brivio, Machado, JFT, preprint in prep.

- We have started a re-interpretation of five CMS VBS papers with lepton / photon final states. Results reported here should be considered preliminary.
- Simulation of ALP signals based on ALP EFT linear model implementation in MadGraph, by Brivio, Gavela, No, Sanz, et al.
- Expected numbers of events in the next slides use the selections and integrated luminosities in the CMS publications.
- Numbers of events reduced $\sim 85\%$ due to EFT consistency condition $M(VV) < f_a \sim 2 \text{ TeV}$.
- Delphes simulation efficiencies “calibrated” using SM EWK VBS channel.

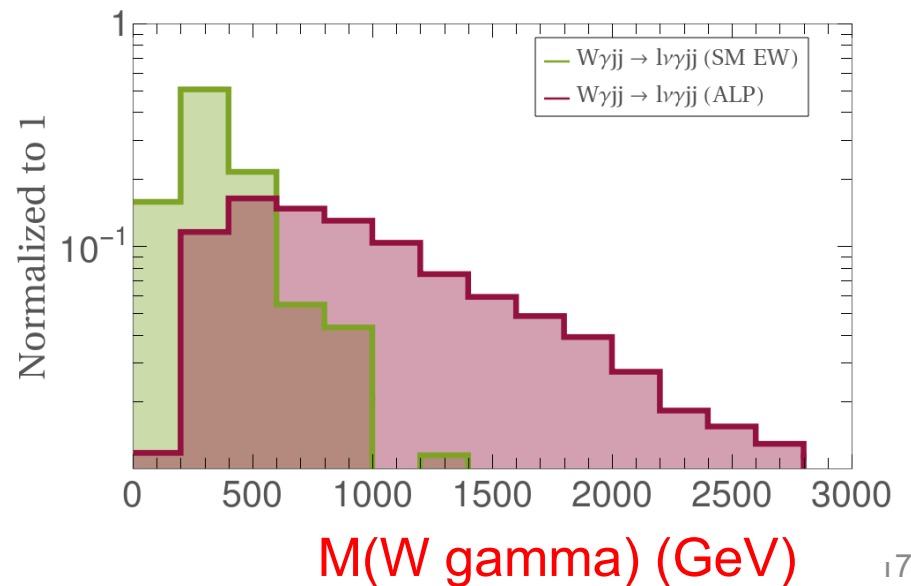
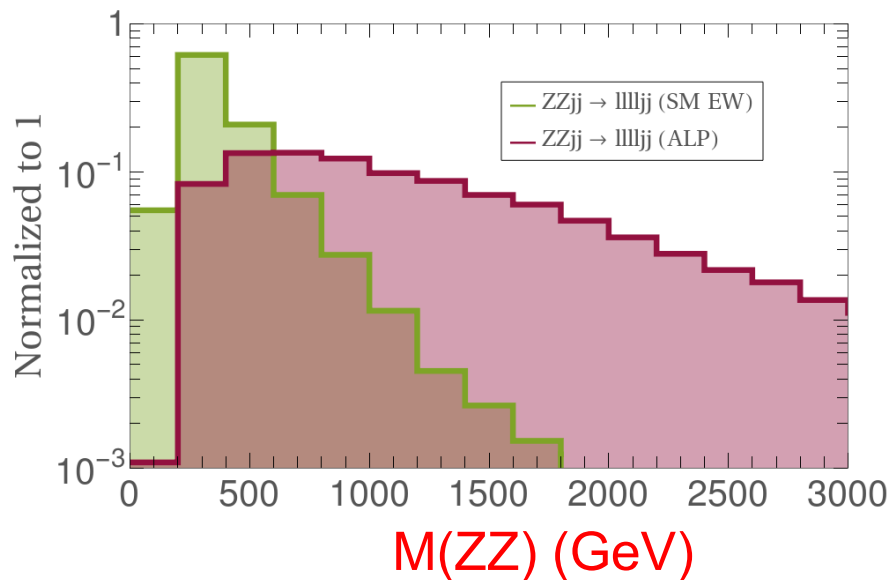
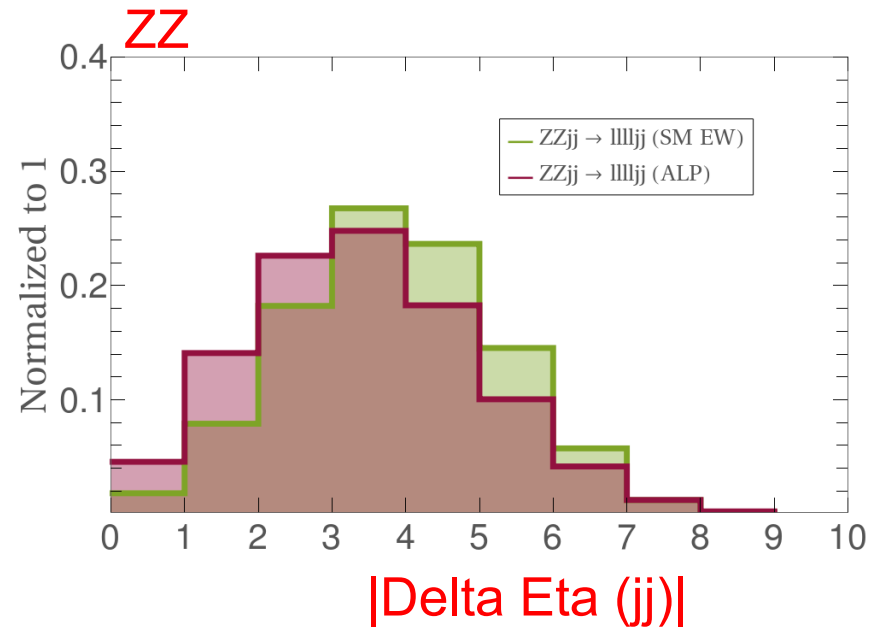
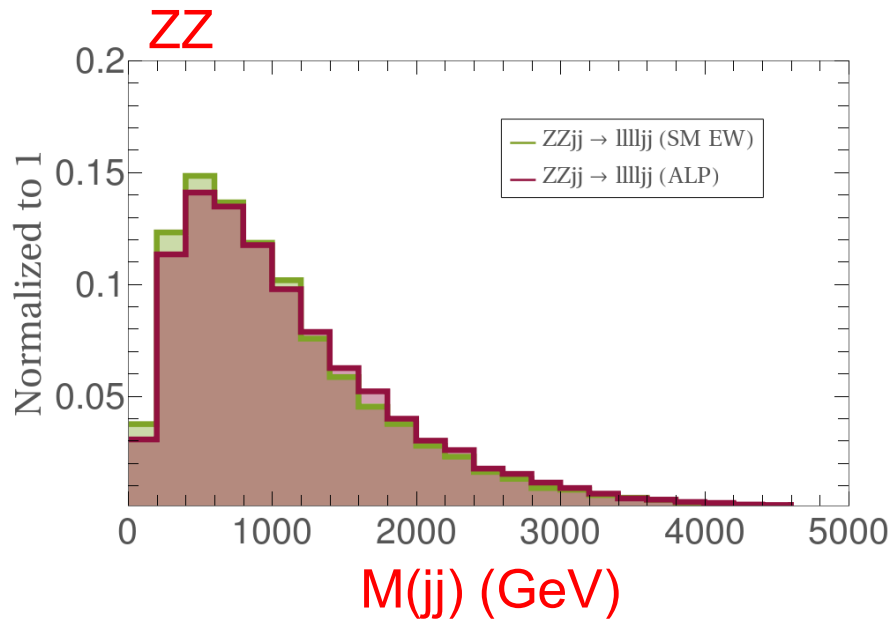
ALP VBS Cross-Sections @ 13 TeV

- Linear EFT Benchmark Case (1): $c_B/fa = c_W/fa = 1/TeV$, $M_a = 1 MeV$;
Case (2): photophobic, $c_W/fa = 1/TeV$.

	Case (1) sgnl. / interf.	Case (2) sgnl. / interf.	# dilepton events at CMS	int. lum. (fb^{-1})
• ZZ	42 / -13 fb	18 / -9 fb	9.3 / -3.2	137
• WZ	18 / 1.7 fb	24 / -0.1 fb	4.2 / 0.05	137
• ssWW	16 / -4.0 fb		18 / -5.5	137
• W gamma	29 / 4.3 fb	5.4 / 1.7 fb	3.6 / -0.04	36
• Z gamma	11 / 0.3 fb	21 / -9 fb	3.8 / 0.02	36

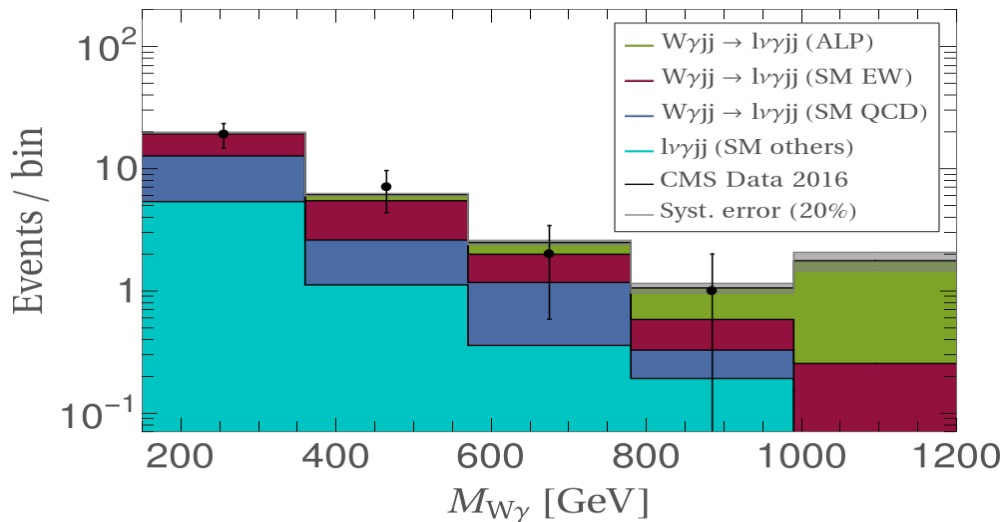
- Selection efficiencies range from 5% to 30%, depending on the dijet selection cuts. Typical efficiency is $\sim 15\%$.

VBS Observables: ALP vs. SM EWK

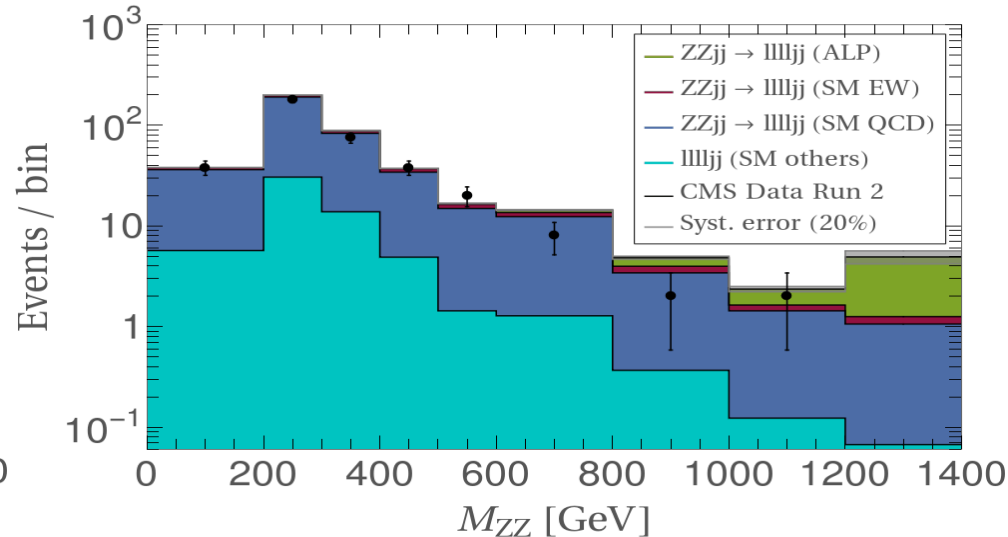


ALP $M(VV)$ in CMS Leptonic Analyses

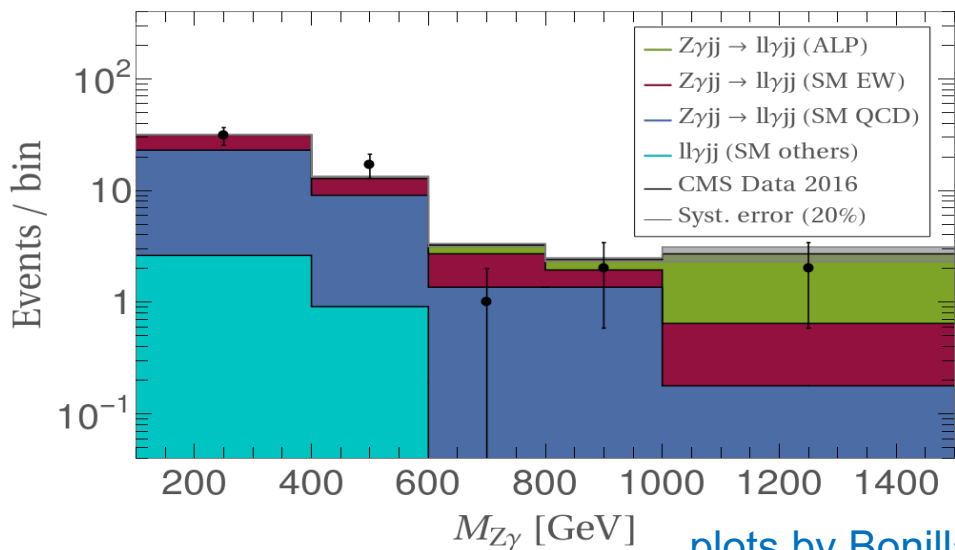
W gamma CMS-SMP-19-008



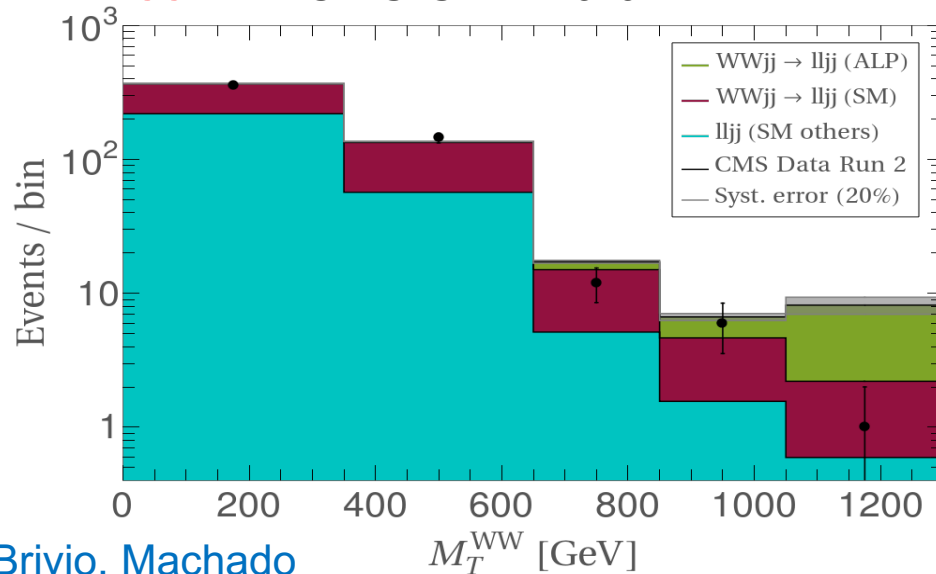
ZZ CMS-SMP-20-001



Z gamma CMS-SMP-18-007

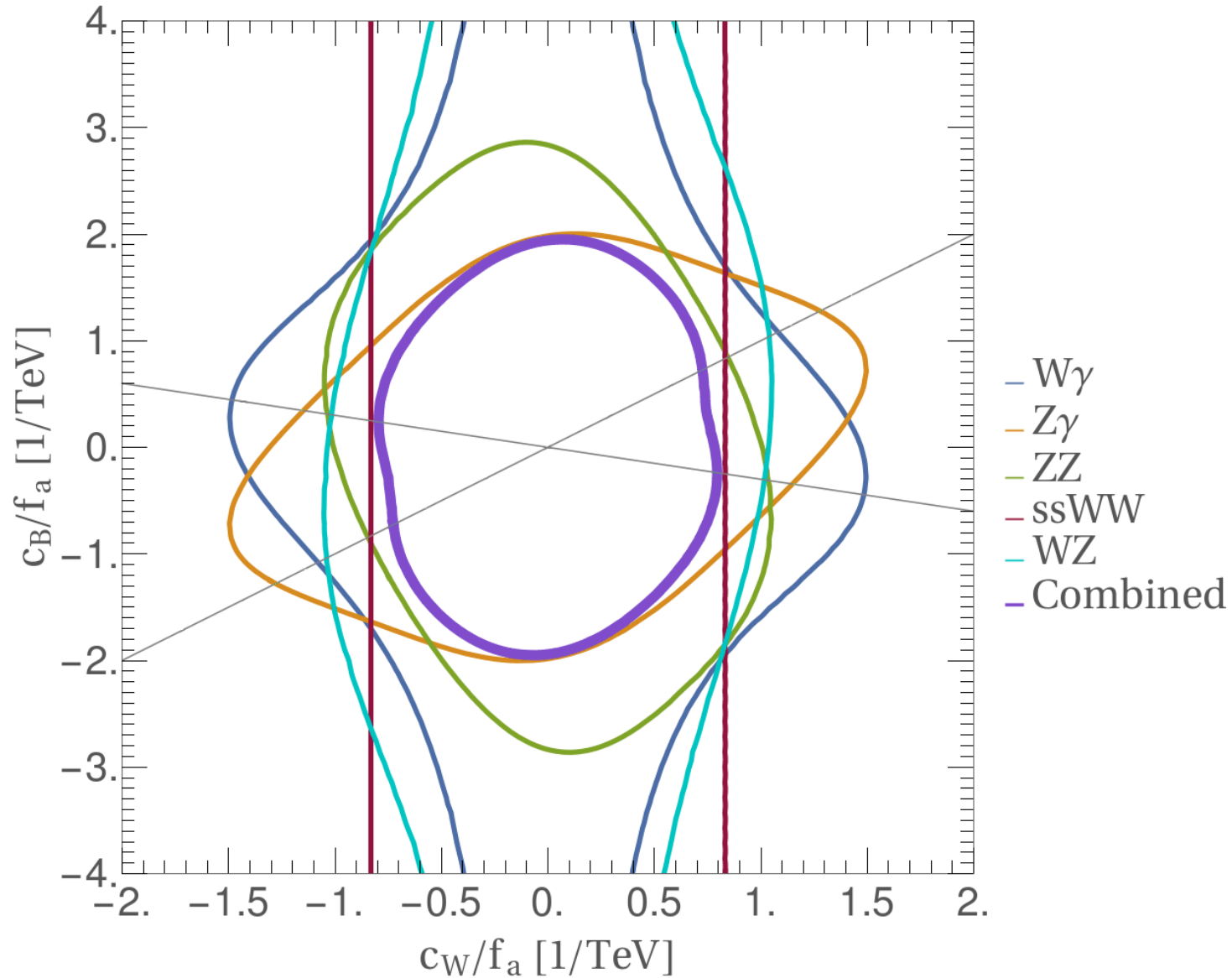


ssWW CMS-SMP-19-012

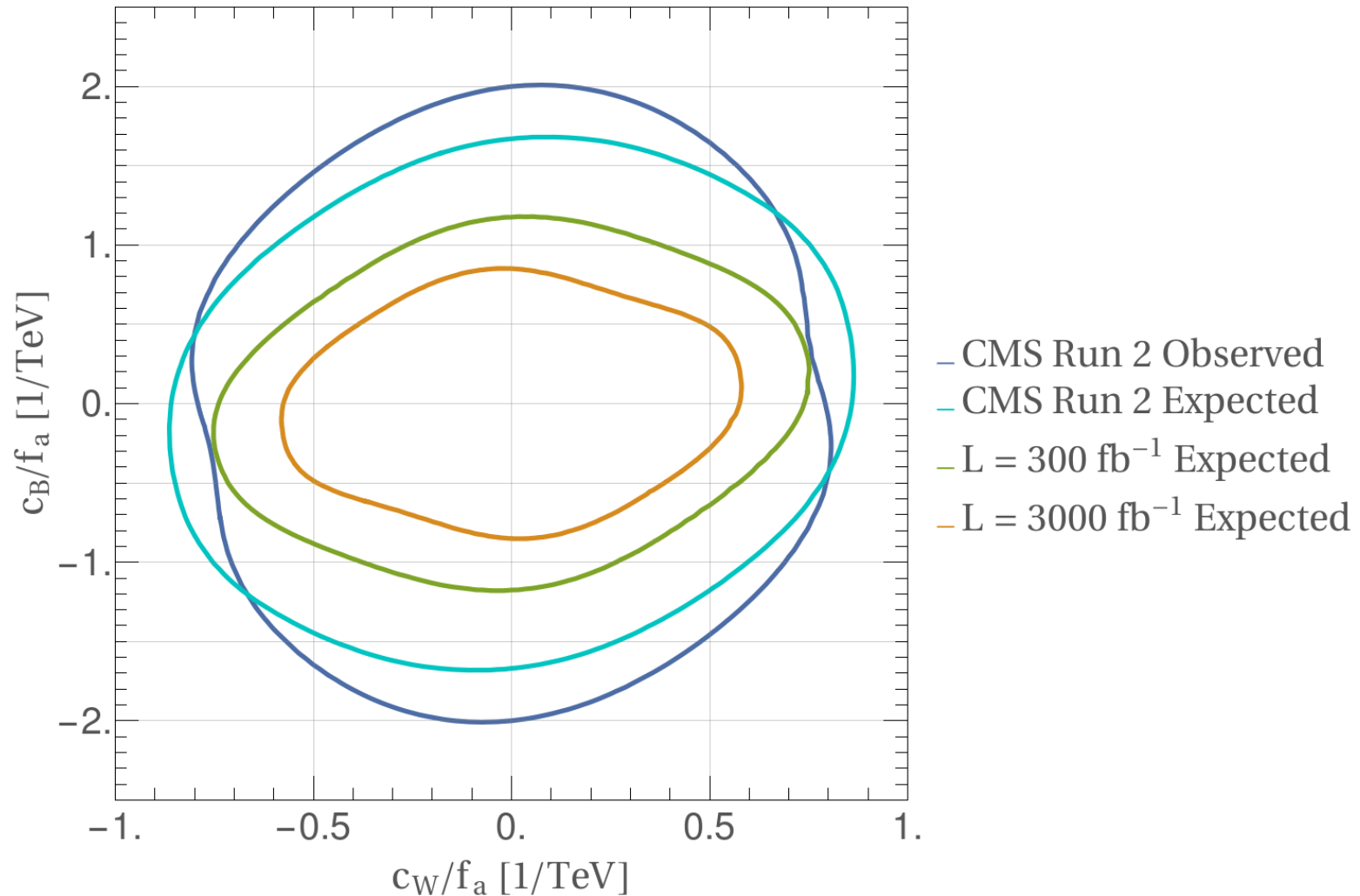


plots by Bonilla, Brivio, Machado

Current Limits w/ CMS Run 2 Data



Projected Limits at Run3 and HL-LHC



Expected Results and Conclusions

- Current analysis at the limit of (small) statistics.
- Expected limits on f_a / c_W in the 1 – 2 TeV region.
- Next exercise, calculate expectations at larger LHC luminosities.