

Portal Effective Theories

A framework for the model independent description of
light hidden sector interactions

Philipp Klose

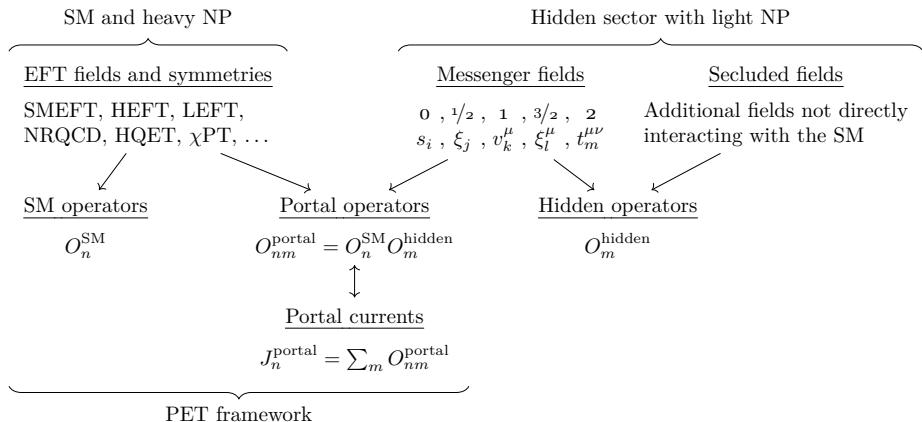
Albert Einstein Center for Fundamental Physics, Institute for Theoretical Physics,
Universität Bern, Sidlerstraße 5, CH-3012 Bern, Switzerland

Searching for long-lived particles at the LHC and beyond: Ninth
workshop of the LLP Community

arXiv:2105.06477

Collaborators: Chiara Arina, Jan Hajer

Portal Effective Theory (PET) Framework



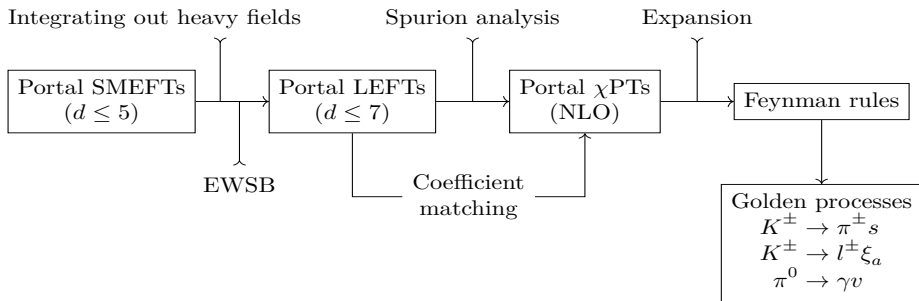
Construct and link multiple PETs

\Rightarrow Translate EW scale \leftrightarrow strong scale \leftrightarrow mesons

Proof of concept: portal SMEFT \rightarrow portal χ PT

Physics Focus: (1) Light mesons (E.g. at NA62, MaTHUSLA, KOTO)
(2) Electroweak interactions
(Needed for e.g. $\pi^0 \rightarrow \gamma X$, $K^+ \rightarrow \pi^+ X$)

- We consider single messenger s_i , ξ_i , ν^μ \Rightarrow **3 sets of PETs**
(SM gauge singlet; spin 0, 1/2, 1)



Strong Scale QCD & χ PT Portal Interactions

SM: **4 currents** capture quark masses, photons, θ angle:

$$\delta\mathcal{L}_{\text{QCD}} = -q^\dagger \bar{\sigma}_\mu \mathbf{l}^\mu q - \bar{q} \sigma_\mu \mathbf{r}^\mu q^\dagger - \theta G_{\mu\nu} \tilde{G}^{\mu\nu} - [\bar{q} \mathbf{m} q + \text{h.c.}]$$

QCD Flavour symmetry determines meson to current coupling
 \Rightarrow spurion analysis

Our work: **10 portal currents** parametrize new physics:

$$\delta\mathcal{L}_{\text{QCD}} \rightarrow -q^\dagger \bar{\sigma}_\mu \mathbf{L}^\mu q - \bar{q} \sigma_\mu \mathbf{R}^\mu q^\dagger - \Omega G_{\mu\nu} G^{\mu\nu} - \Theta G_{\mu\nu} \tilde{G}^{\mu\nu} \\ - [\bar{q} \mathbf{M} q + \bar{q} \bar{\sigma}^{\mu\nu} \mathbf{T}_{\mu\nu} q - \bar{q} \mathbf{\Gamma} \bar{\sigma}^{\mu\nu} G_{\mu\nu} q + \text{h.c.}] + \mathcal{L}_{\bar{q}q\bar{q}q} [H_l, H_r, H_s]$$

Portal χ PT: Couple mesons to **all portal currents**

- We extend spurion approach to include new currents
- We estimate resulting new χ PT coefficients
(Using scale anomaly, QCD condensates, large n_c limit)

Application: Kaon Decays into Hidden Particles

- 1 Fermionic Messengers: $\mathcal{A}(K^+ \rightarrow \ell_a^+ \xi_b)$ depends on **1** portal current
 - Reproduces standard result for HNL
 - Captures models with additional **secluded fields**
- 2 Scalar Messengers: $\mathcal{A}(K^+ \rightarrow \pi^+ s_j)$ depends on **7** portal currents

$$\mathcal{A} = \mathcal{A}_M + \mathcal{A}_{S_w} + \mathcal{A}_{S_\theta} + \mathcal{A}_\Gamma + \mathcal{A}_{H_l} + \mathcal{A}_{H_r} + \mathcal{A}_{H_s}$$

- Captures production of **general spin 0 messengers** (ALPs, light Higgs, complex scalars, sGoldstinos, etc.)
- Includes i) mixing with mesons, ii) full flavour dependence
- Encompasses prior model-dependent results for ALPs, light Higgs, etc.

Summary

- 1 We developed the **PET Framework**
⇒ couple EFTs to generic light hidden sectors
- 2 We constructed PETs with single messenger of spin 0, 1/2, 1 :
 - portal SMEFTs containing all $d \leq 5$ operators
 - portal LEFTs also containing flavour violating $d = 6, 7$ operators
⇒ **10 portal currents**
 - portal χ PT couples mesons to portal currents
- 3 We computed **generic decay amplitudes**:
 - $\mathcal{A}(K^+ \rightarrow \pi^+ s_i)$ (spin 0: e.g. ALP, Higgs portal, etc.)
 - $\mathcal{A}(K^+ \rightarrow \ell^+ \xi_i)$ (spin 1/2: e.g. HNL, etc.)
 - $\mathcal{A}(\pi^0 \rightarrow \gamma \nu_i)$ (Spin 1 messenger: E.g. Dark photon)

Thank you for your attention!

EW Scale PETs: Portal Operators

	d	Higgs	Yukawa + h.c.	Fermions	Gauge bosons
	3	$s_i H ^2$			
	4	$s_i s_j H ^2$			
s_i	5	$s_i s_j s_k H ^2$	$s_i q_a \bar{u}_b \tilde{H}^\dagger$		$s_i G_{\mu\nu}^a G_a^{\mu\nu}$
		$s_i D^\mu H^\dagger D_\mu H$	$s_i q_a \bar{d}_b H^\dagger$		$s_i W_{\mu\nu}^a W_a^{\mu\nu}$
	5	$s_i H ^4$	$s_i \ell_a \bar{e}_b H^\dagger$		$s_i B_{\mu\nu} B^{\mu\nu}$
					$s_i G_{\mu\nu}^a \tilde{G}_a^{\mu\nu}$
					$s_i W_{\mu\nu}^a \tilde{W}_a^{\mu\nu}$
					$s_i B_{\mu\nu} \tilde{B}^{\mu\nu}$
ξ_a	4		$\xi_a \ell_b \tilde{H}^\dagger$		
† h.c.	5	$\xi_a \xi_b H ^2$	$\xi_a^\dagger \bar{\sigma}^\mu \ell_b D_\mu \tilde{H}^\dagger$		$\xi_a \sigma^{\mu\nu} \xi_b B_{\mu\nu}$
v^μ	4	$v_\mu v^\mu H ^2$		$v^\mu q_a^\dagger \bar{\sigma}_\mu q_b$	
		$\partial_\mu v^\mu H ^2$		$v^\mu \bar{u}_a^\dagger \sigma_\mu \bar{u}_b$	
		$v^\mu H^\dagger \overset{\leftrightarrow}{D}_\mu H$		$v^\mu \bar{d}_a^\dagger \sigma_\mu \bar{d}_b$	
				$v^\mu \ell_a^\dagger \bar{\sigma}_\mu \ell_b$	
				$v^\mu \bar{e}_a^\dagger \sigma_\mu \bar{e}_b$	

Strong Scale PETs: $d = 6, 7$ and $|\Delta F| = 1$ Portal Operators

d	Two quarks	Quark dipole	Four fermions
6	$s_i s_j s_k \bar{d} d$	$s_i F^{\mu\nu} \bar{d} \sigma_{\mu\nu} d$	
	$\partial^2 s_i \bar{d} d$	$s_i G^{\mu\nu} \bar{d} \sigma_{\mu\nu} d$	
	$s_i \partial_\mu s_j d^\dagger \bar{\sigma}^\mu d$		
s_i	$s_i s_j s_k s_l \bar{d} d$		$s_i d^\dagger \bar{q}^\dagger \bar{q} d$
			$s_i q^\dagger \bar{\sigma}^\mu q q^\dagger \bar{\sigma}_\mu q$
			$s_i d^\dagger \bar{\sigma}^\mu d \bar{q} \sigma_\mu \bar{q}^\dagger$
			$s_i e^\dagger \bar{\sigma}_\mu \nu u^\dagger \bar{\sigma}^\mu d$
			$s_i \nu^\dagger \bar{\sigma}_\mu \nu d^\dagger \bar{\sigma}^\mu d$
ξ_a h.c.	6	$\xi_a^\dagger \bar{\sigma}_\mu e d^\dagger \bar{\sigma}^\mu u$	
		$\xi_a^\dagger \bar{\sigma}_\mu \nu d^\dagger \bar{\sigma}^\mu d$	

$K^+ \rightarrow \pi^+ s_i$ Interactions

Currents:

$$\begin{aligned} \Omega &\supset \frac{\epsilon_{UV}}{v} c_i^{S_\omega} s_i, & \mathbf{M} &\supset \epsilon_{UV} \left(c_i^{S_m} + c_{\partial^2 i}^{S_m} \frac{1}{v^2} \partial^2 \right) s_i, & H_x &\supset h_{xi} \frac{\epsilon_{UV}}{v} s_i, \\ \Theta &\supset \frac{\epsilon_{UV}}{v} c_i^{S_\theta} s_i, & \mathbf{\Gamma} &\supset \epsilon_{UV} \left(\lambda_d^s c_{i\bar{s}d}^\gamma + \lambda_s^d c_{i\bar{d}s}^\gamma \right) s_i \end{aligned}$$

Scalar current coupling to Kaons:

$$\mathcal{L}_{\text{portal}} \supset -\frac{b}{2} \epsilon_{UV} K^+ \pi^- \left(c_{K\pi s_i} + \text{Re } c_{\partial^2 i s}^{S_m d} \frac{\partial^2}{v^2} \right) s_i,$$

where

$$\begin{aligned} c_{K\pi s_i} = & \text{Re } c_i^{S_m d} + \frac{\epsilon_{EW}}{2} \left((m_K^2 - m_\pi^2) \text{Re } c_i^{S_m u} + m_K^2 \text{Re } c_i^{S_m d} - m_\pi^2 \text{Re } c_i^{S_m s} \right) \theta_{K^\pm \pi^\mp} \\ & - \frac{\epsilon_{EW}}{2} \left(2v h_b \left(\frac{m_{ud}}{m_s} \left(c_i^{S_m d} - c_i^{S_m s^\dagger} \right) + c_i^{S_m d^\dagger} \right) + \frac{m_{ud} + m_s}{v} h_{bi} - \kappa_\gamma \left(c_{i\bar{d}s}^\gamma + c_{i\bar{s}d}^{\gamma^\dagger} \right) \right) \end{aligned}$$

$\theta_{K^\pm \pi^\mp}$ is kaon to pion mixing angle, m_{ud} is light quark mass

$K^+ \rightarrow \pi^+ s_i$ Amplitude

$$\mathcal{A}(K^+ \rightarrow \pi^+ s_i) = \mathcal{A}_{\text{Re } S_m}(K^+ \rightarrow \pi^+ s_i) + \mathcal{A}_{\text{Im } S_m}(K^+ \rightarrow \pi^+ s_i) \\ + \mathcal{A}_\omega(K^+ \rightarrow \pi^+ s_i) + \mathcal{A}_\theta(K^+ \rightarrow \pi^+ s_i) + \mathcal{A}_{8+27}(K^+ \rightarrow \pi^+ s_i)$$

where

$$\mathcal{A}_{\text{Re } S_m}(K^+ \rightarrow \pi^+ s_i) = -\frac{\epsilon_{UV} b}{2} \left(c_{K\pi s_i} - \text{Re } \mathbf{c}_{\partial^2 i s}^{S_m d} \frac{m_s^2}{v^2} \right)$$

$$\mathcal{A}_\omega(K^+ \rightarrow \pi^+ s_i) = \frac{\epsilon_{UV} \epsilon_{EW} c_i^{S_\omega}}{\beta_0 v} \left(h'_b m_K^2 - \frac{1}{2} (h_8 + (3-1)h_{27}) (m_K^2 + m_\pi^2 - m_s^2) \right)$$

$$\mathcal{A}_{8+27}(K^+ \rightarrow \pi^+ s_i) = -\frac{\epsilon_{UV} \epsilon_{EW}}{4v} (h_{8i} + (3-1)h_{27i}) (m_K^2 + m_\pi^2 - m_s^2)$$

$$\mathcal{A}_{\text{Im } S_m}(K^+ \rightarrow \pi^+ s_i) = -i \epsilon_{UV} \epsilon_{EW} b \left(c_{s_i \pi} \frac{V_{K\pi\pi}}{m_s^2 - m_\pi^2} + c_{s_i \eta} \frac{V_{K\pi\eta}}{m_s^2 - m_\eta^2} + c_{s_i \eta'} \frac{V_{K\pi\eta'}}{m_s^2 - m_{\eta'}^2} \right)$$

$$\mathcal{A}_\theta(K^+ \rightarrow \pi^+ s_i) = i \frac{\epsilon_{UV} \epsilon_{EW} c_i^{S_\theta} m_0^2}{v} \left(c_\eta \frac{V_{K\pi\eta'}}{m_s^2 - m_{\eta'}^2} - s_\eta \frac{V_{K\pi\eta}}{m_s^2 - m_\eta^2} \right).$$

- c_η, s_η are (co-)sine of $\eta - \eta'$ mixing angle, $V_{\phi_a \phi_b \phi_c}$ are SM 3 meson vertices
- $c_{K\pi s_i}, c_{s_i \phi_a}$ parametrize coupling to hidden sectors

$K^+ \rightarrow \nu_a \xi_b$ Width

Interaction:

$$\mathcal{L}_{\text{portal}} \supset -\epsilon_{UV} v c_{ba}^\nu \nu_b \xi_a - \frac{\epsilon_{UV} f}{v^2} c_{\bar{u}s,ba}^{L\dagger} \xi_a^\dagger \bar{\sigma}_\mu e_b \partial^\mu K^+ + \text{h.c.}$$

Width:

$$\Gamma = |\theta'_{ba}|^2 \Gamma(K^+ \rightarrow \ell_b^+ \nu_b) \frac{\rho(x_\ell, x_\xi)}{\rho(x_\ell, 0)}, \quad \theta'_{ba} = \epsilon_{UV} \left(\frac{c_{ba}^\nu v}{m_\xi} + \frac{c_{\bar{u}s,ba}^L}{V_{us}} \right)$$

Phasespace factor:

$$\rho(x_\ell, x_\xi) = \left(x_\ell + x_\xi - (x_\ell - x_\xi)^2 \right) \sqrt{\left(\frac{1 - x_\ell - x_\xi}{2} \right)^2 - x_\ell x_\xi}, \quad x_i = \frac{m_i^2}{m_K^2}$$

$\pi^0 \rightarrow \gamma \nu_i$ Width

Interaction:

$$\mathcal{L}_{\text{portal}} \supset \frac{2}{(4\pi)^2 f} (2\partial^\mu \mathbf{V}_{\nu u}^\nu + \partial^\mu \mathbf{V}_{\nu d}^\nu) \frac{\pi^0}{\sqrt{2}} e \tilde{F}_{\mu\nu}, \quad \mathbf{V}_\nu^\mu = \epsilon_{UV} (\mathbf{c}_\nu^L + \mathbf{c}_\nu^R) \nu^\mu$$

Width:

$$\Gamma(\pi^0 \rightarrow \gamma \nu_i) = 2\epsilon_{\text{eff}}^2 \Gamma_{\pi^0 \rightarrow \gamma\gamma} \left(1 - \frac{m_\nu^2}{m_\pi^2}\right)^3, \quad \epsilon_{\text{eff}} = \epsilon_{UV} \frac{2(\mathbf{c}_\nu^R + \mathbf{c}_\nu^L)_u + (\mathbf{c}_\nu^R + \mathbf{c}_\nu^L)_d}{2e(2q_u + q_d)}$$