

Mapping the viable parameter space for testable leptogenesis

Yannis Georis

based on work in collaboration with M. Drewes and J. Klaric
[arXiv:2106.16226]

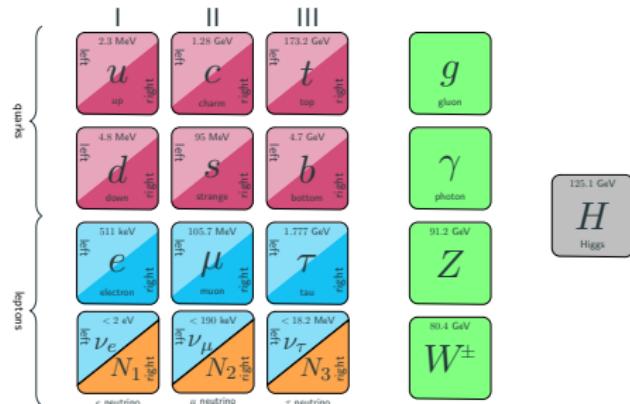
Searching for long-lived particles at the LHC and beyond:
Ninth workshop of the LLP Community
CERN, May 28, 2021



Heavy neutral lepton (HNLs)

Heavy neutrinos can solve three problems at once:

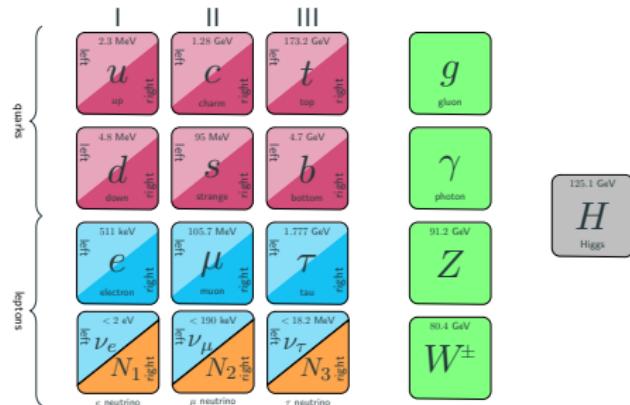
- Non-zero mass of the Standard Model (SM) neutrinos via the **type-I seesaw** mechanism.
(P.Minkowski, 1977)



Heavy neutral lepton (HNLs)

Heavy neutrinos can solve three problems at once:

- Non-zero mass of the Standard Model (SM) neutrinos via the **type-I seesaw** mechanism.
(P. Minkowski, 1977)
- Overabundance of matter with respect to antimatter through **leptogenesis**.
(M. Fukugita, T. Yanagida, 1986)



Heavy neutral lepton (HNLs)

Heavy neutrinos can solve three problems at once:

- Non-zero mass of the Standard Model (SM) neutrinos via the **type-I seesaw** mechanism.
(P. Minkowski, 1977)

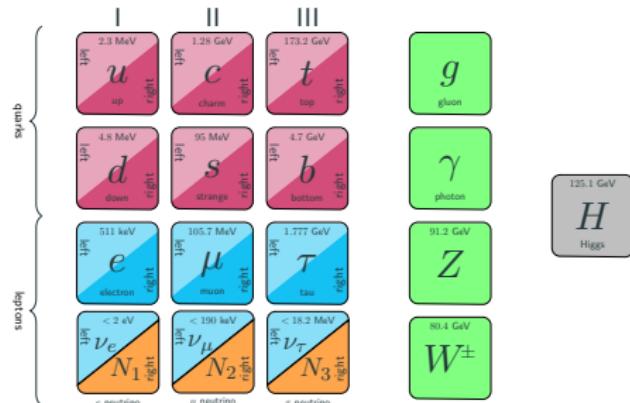
- Overabundance of matter with respect to antimatter through **leptogenesis**.

(M. Fukugita, T. Yanagida, 1986)

- Can be a **Dark Matter** candidate

(Dodelson/Widrow, hep-ph/9303287,

Asaka/Shaposhnikov, hep-ph/0505013)



Leptogenesis from neutrinos oscillations

- "Standard" leptogenesis requires $M_N \gtrsim 10^9$ GeV (Davidson-Ibarra bound) if mass hierarchy ($M_1 \ll M_2 \ll M_3$).

Leptogenesis from neutrinos oscillations

- "Standard" leptogenesis requires $M_N \gtrsim 10^9$ GeV (Davidson-Ibarra bound) if mass hierarchy ($M_1 \ll M_2 \ll M_3$).
→ direct detection difficult at LHC !

Leptogenesis from neutrinos oscillations

- "Standard" leptogenesis requires $M_N \gtrsim 10^9$ GeV (Davidson-Ibarra bound) if mass hierarchy ($M_1 \ll M_2 \ll M_3$).
→ direct detection difficult at LHC !

Low-scale leptogenesis

- "Standard" leptogenesis requires $M_N \gtrsim 10^9$ GeV (Davidson-Ibarra bound) if mass hierarchy ($M_1 \ll M_2 \ll M_3$).
→ direct detection difficult at LHC !

Low-scale leptogenesis

1) Freeze-out mechanism

- TeV-scale HNLs allowed if small HNLs mass splittings
- resonant leptogenesis ([Pilaftsis, hep-ph/9707235](#))

- "Standard" leptogenesis requires $M_N \gtrsim 10^9$ GeV (Davidson-Ibarra bound) if mass hierarchy ($M_1 \ll M_2 \ll M_3$).
→ direct detection difficult at LHC !

Low-scale leptogenesis

1) Freeze-out mechanism

- TeV-scale HNLs allowed if small HNLs mass splittings
- resonant leptogenesis ([Pilaftsis, hep-ph/9707235](#))

2) Freeze-in mechanism

- Allow for GeV-scale HNLs.
- *ARS mechanism* ([Akhmedov/Rubakov/Smirnov, hep-ph/9803255](#))

- "Standard" leptogenesis requires $M_N \gtrsim 10^9$ GeV (Davidson-Ibarra bound) if mass hierarchy ($M_1 \ll M_2 \ll M_3$).
→ direct detection difficult at LHC !

Low-scale leptogenesis

1) Freeze-out mechanism

- TeV-scale HNLs allowed if small HNLs mass splittings
- resonant leptogenesis (Pilafitsis, hep-ph/9707235)

2) Freeze-in mechanism

- Allow for GeV-scale HNLs.
- ARS mechanism (Akhmedov/Rubakov/Smirnov, hep-ph/9803255)

→ Represented by the same set of equations !

(cfr. B.Garbrecht 1812.02651)

Quantum kinetic equations

$$i \frac{d n_{\Delta_\alpha}}{dt} = -2i \frac{\mu_\alpha}{T} \int \frac{d^3 k}{(2\pi)^3} \text{Tr}[\Gamma_\alpha] f_N (1 - f_N) + i \int \frac{d^3 k}{(2\pi)^3} \text{Tr}[\bar{\Gamma}_\alpha (\rho_N - \rho_{\bar{N}})]$$

$$i \frac{d \rho_N}{dt} = [H, \rho_N] - \frac{i}{2} \{ \Gamma, \rho_N - \rho_N^{eq} \} - \frac{i}{2} \sum_\alpha \bar{\Gamma}_\alpha \left[\frac{2\mu_\alpha}{T} f_N (1 - f_N) \right]$$

$$i \frac{d \rho_{\bar{N}}}{dt} = -[H^*, \rho_{\bar{N}}] - \frac{i}{2} \{ \Gamma^*, \rho_{\bar{N}} - \rho_{\bar{N}}^{eq} \} + \frac{i}{2} \sum_\alpha \bar{\Gamma}_\alpha^* \left[\frac{2\mu_\alpha}{T} f_N (1 - f_N) \right]$$

- Rate equations for density matrices $\rho_N, \rho_{\bar{N}}$ and chemical potential μ_α .
- Allow to cover the whole mass range 100 MeV- 70 TeV

Seesaw type-I

- Seesaw mass term

$$\begin{aligned}\mathcal{L} &\supset \frac{1}{2} (\bar{\nu}_L \quad \bar{\nu}_R^c) \cdot \begin{pmatrix} 0 & m_D \\ m_D^t & M_M \end{pmatrix} \cdot \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} \\ &\longrightarrow \quad m_\nu = -m_D \cdot M_M^{-1} \cdot m_D^t, \quad m_D = v F.\end{aligned}$$

Seesaw type-I

- Seesaw mass term

$$\mathcal{L} \supset \frac{1}{2} (\bar{\nu}_L \quad \bar{\nu}_R^c) \cdot \begin{pmatrix} 0 & m_D \\ m_D^t & M_M \end{pmatrix} \cdot \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

$$\longrightarrow \quad m_\nu = -m_D \cdot M_M^{-1} \cdot m_D^t, \quad m_D = v F.$$

- Strength of the mixing with SM neutrinos parametrised by

$$U^2 = \sum_{a,i} |(m_D \cdot M_M^{-1})_{ai}|^2 \equiv \sum_a U_a^2.$$

Seesaw type-I

- Seesaw mass term

$$\begin{aligned}\mathcal{L} &\supset \frac{1}{2} (\bar{\nu}_L \quad \bar{\nu}_R^c) \cdot \begin{pmatrix} 0 & m_D \\ m_D^t & M_M \end{pmatrix} \cdot \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} \\ &\longrightarrow \quad m_\nu = -m_D \cdot M_M^{-1} \cdot m_D^t, \quad m_D = v F.\end{aligned}$$

- Strength of the mixing with SM neutrinos parametrised by

$$U^2 = \sum_{a,i} |(m_D \cdot M_M^{-1})_{ai}|^2 \equiv \sum_a U_a^2.$$

Question: Leptogenesis possible for high U^2 ?

Seesaw type-I

- Seesaw mass term

$$\begin{aligned}\mathcal{L} &\supset \frac{1}{2} (\bar{\nu}_L \quad \bar{\nu}_R^c) \cdot \begin{pmatrix} 0 & m_D \\ m_D^t & M_M \end{pmatrix} \cdot \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} \\ &\longrightarrow \quad m_\nu = -m_D \cdot M_M^{-1} \cdot m_D^t, \quad m_D = v F.\end{aligned}$$

- Strength of the mixing with SM neutrinos parametrised by

$$U^2 = \sum_{a,i} |(m_D \cdot M_M^{-1})_{ai}|^2 \equiv \sum_a U_a^2.$$

Question: Leptogenesis possible for high U^2 ?

- B-L approximate symmetry allows to avoid the naive seesaw bound

$$U_i^2 \sim \frac{\sqrt{\Delta m_{atm}^2 + m_{light}^2}}{M} \lesssim 10^{-10} \frac{\text{GeV}}{M_i}$$

$n=2$ leptogenesis

Large U^2

→ large washout

→ asymmetric Yukawa coupling
needed to hide BAU from the
washout

Asymmetry

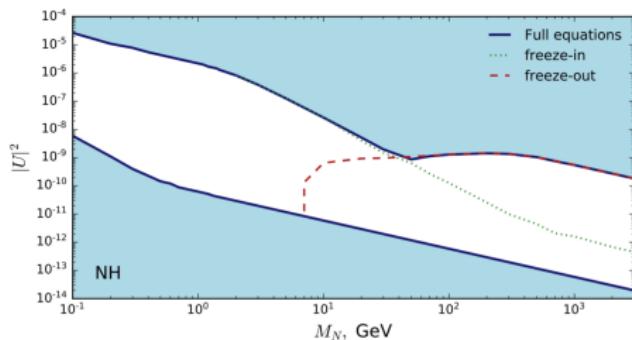
$$f \simeq \frac{\min|F_a|}{\max|F_a|}$$

constrained by neutrino oscillation
data

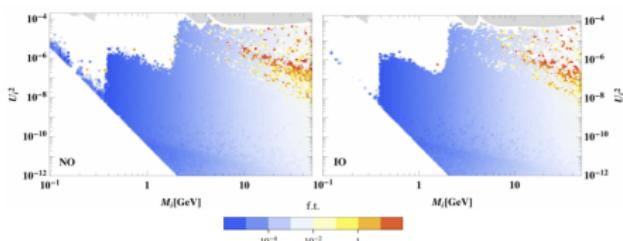
$$f > 5 \cdot 10^{-3}$$

→ high U^2 not possible.

For $n=3$, lepton number violating
processes can play a significant role.



Klaric/Shaposhnikov/Timirsyasov 2008.13771



Abada/Arcadi/Domcke/Drewes/Klaric/Lucente hep-ph/1810.12463

- Consistency with ν -oscillation data
- Casas-Ibarra parametrisation

$$F = \frac{i}{\nu} U_\nu \sqrt{m_\nu^{\text{diag}}} R \sqrt{M_M}$$

- 3 light neutrino masses (2)
3 complex angles (1)
3 Majorana masses (2)
3 CP-violating phases (2)

- Consistency with ν -oscillation data
- Casas-Ibarra parametrisation

$$F = \frac{i}{\nu} U_\nu \sqrt{m_\nu^{\text{diag}}} R \sqrt{M_M}$$

- 3 light neutrino masses (2)
3 complex angles (1)
3 Majorana masses (2)
3 CP-violating phases (2)
- Theoretical constraints:

$n=3$ leptogenesis

- Consistency with ν -oscillation data
- Casas-Ibarra parametrisation

$$F = \frac{i}{\nu} U_\nu \sqrt{m_\nu^{\text{diag}}} R \sqrt{M_M}$$

- 3 light neutrino masses (2)
3 complex angles (1)
3 Majorana masses (2)
3 CP-violating phases (2)
- Theoretical constraints:
 - ➊ Perturbative unitarity $\Gamma < \frac{M}{2}$

$n=3$ leptogenesis

- Consistency with ν -oscillation data
- Casas-Ibarra parametrisation

$$F = \frac{i}{\nu} U_\nu \sqrt{m_\nu^{\text{diag}}} R \sqrt{M_M}$$

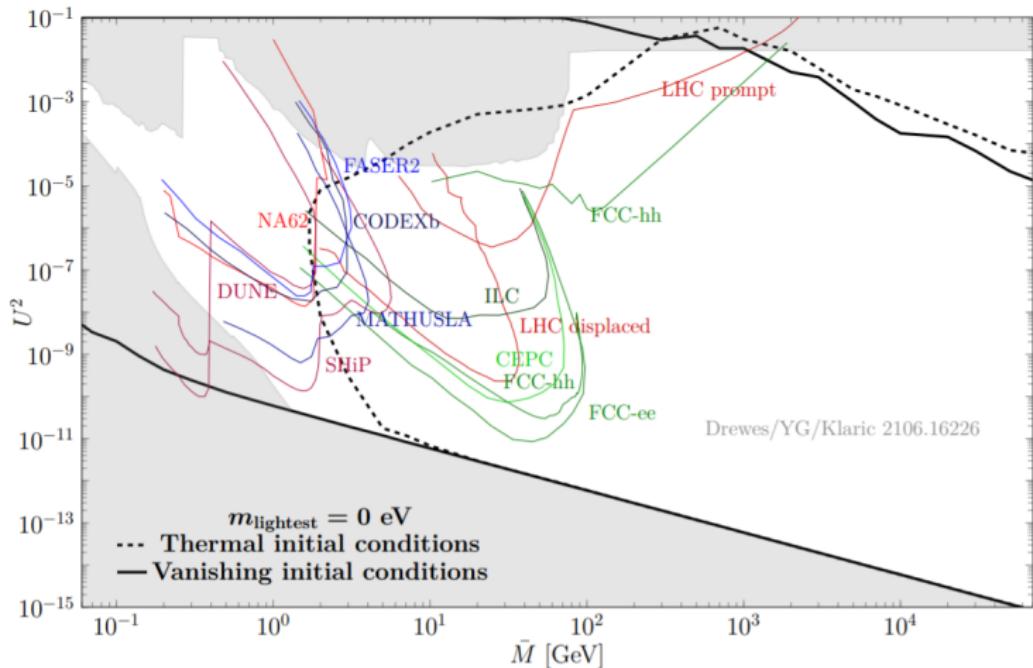
- 3 light neutrino masses (2)
3 complex angles (1)
3 Majorana masses (2)
3 CP-violating phases (2)
- Theoretical constraints:
 - ① Perturbative unitarity $\Gamma < \frac{M}{2}$
 - ② Seesaw expansion $U^2 < 0.1$

- Consistency with ν -oscillation data
- Casas-Ibarra parametrisation

$$F = \frac{i}{\nu} U_\nu \sqrt{m_\nu^{\text{diag}}} R \sqrt{M_M}$$

- 3 light neutrino masses (2)
3 complex angles (1)
3 Majorana masses (2)
3 CP-violating phases (2)
- Theoretical constraints:
 - ① Perturbative unitarity $\Gamma < \frac{M}{2}$
 - ② Seesaw expansion $U^2 < 0.1$
 - ③ No large radiative corrections $(1 - \frac{m_{\text{tree}}}{m_{\text{loop}}})^2 < \frac{1}{4}$.

Results and conclusions



- Reaches theoretical constraint at low masses.
- Parameter space much larger than for the $n=2$ scenario.
- Seesaw line reached as for $n=2$.