

# Heavy Neutrinos from Heavy Bosons with MadGraph5\_aMC@NLO

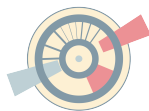
#LHCLLP9 - Zoom National Accelerator Laboratory

Richard Ruiz

Institute of Nuclear Physics – Polish Academy of Science (IFJ PAN)



26 May 2021



## the case for $\nu$ physics

$m_\nu \neq 0$  + renormalizability + SM gauge inv.  $\implies$  new particles!

[Ma ('98)]

Incredibly powerful but also incredibly vague since new particles:

- ... can be light 😊 or heavy ☹️
- ... can be short-lived 😊 or long-lived 😊
- ... can have SM gauge interactions, e.g.,  $H^{\pm\pm}$  in Type II Seesaw
- ... can have new gauge interactions, e.g.  $\nu_R$  and  $Z_{B-L}$  in  $U(1)_{B-L}$
- ... must couple to  $\Phi_{SM}$  and  $L$ , often inducing collider processes that do not conserve **lepton number (LNV)** and/or **lepton flavor (LFV)**

## the case for sterile neutrinos

# The case for sterile neutrinos

To generate Dirac masses for  $\nu$  like other SM fermions, we need  $\nu_R$

$$\begin{aligned}\mathcal{L}_{\nu \text{ Yuk.}} &= -y_\nu \bar{L} \tilde{\Phi} \nu_R + H.c. = -y_\nu (\bar{\nu}_L \quad \bar{\ell}_L) \begin{pmatrix} \langle \Phi \rangle + h \\ 0 \end{pmatrix} \nu_R + H.c. \\ &= \underbrace{-y_\nu \langle \Phi \rangle}_{=m_D} \bar{\nu}_L \nu_R + H.c. + \dots\end{aligned}$$

However,  $\nu_R^k$  do not exist in the SM, implying  $m_D = 0$

# The case for sterile neutrinos

To generate Dirac masses for  $\nu$  like other SM fermions, we need  $\nu_R$

$$\begin{aligned}\mathcal{L}_{\nu \text{ Yuk.}} &= -y_\nu \bar{L} \tilde{\Phi} \nu_R + H.c. = -y_\nu (\bar{\nu}_L \quad \bar{\ell}_L) \begin{pmatrix} \langle \Phi \rangle + h \\ 0 \end{pmatrix} \nu_R + H.c. \\ &= \underbrace{-y_\nu \langle \Phi \rangle \bar{\nu}_L \nu_R}_{=m_D} + H.c. + \dots\end{aligned}$$

However,  $\nu_R^k$  do not exist in the SM, implying  $m_D = 0$

**A solution: The Type I Seesaw mechanisms** generate  $m_\nu$  by hypothesizing right-handed neutrinos  $\nu_R^k$  with Majorana mass  $m_R^k = \Lambda_{LNV}$

- depending on assumptions,  $m_\nu \sim \Lambda_{LNV}$  or  $v^2/\Lambda_{LNV}$

After EWSB,  $\nu_\ell$  and  $\nu_R$  have same quantum numbers  $\implies$  mixing!

**In practice?**

For **discovery purposes**, take agnostic/pheno. approach with **generic mixing**  $V_{\ell N}$  and **Dirac or Majorana**  $N$  mass eigenstates

Atre, Han, Pascoli, Zhang [0901.3589]

$$\underbrace{\nu_{\ell L}}_{\text{flavor basis}} \approx \underbrace{\sum_{m=1}^3 U_{\ell m} \nu_m + V_{\ell m'=4} N_{m'=4}}_{\text{mass basis}}$$



For **discovery purposes**, take agnostic/pheno. approach with **generic mixing**  $V_{\ell N}$  and **Dirac or Majorana**  $N$  mass eigenstates

Atre, Han, Pascoli, Zhang [0901.3589]

$$\underbrace{\nu_{\ell L}}_{\text{flavor basis}} \approx \underbrace{\sum_{m=1}^3 U_{\ell m} \nu_m + V_{\ell m'=4} N_{m'=4}}_{\text{mass basis}}$$

The SM  $W$  chiral coupling to **leptons** in **flavor basis** is

$$\mathcal{L}_{\text{Int.}} = -\frac{g_W}{\sqrt{2}} W_{\mu}^{-} \sum_{\ell=e}^{\tau} [\bar{\ell} \gamma^{\mu} P_L \nu_{\ell}] + \text{H.c.}, \quad \text{where } P_L = \frac{1}{2}(1 - \gamma^5)$$

For **discovery purposes**, take agnostic/pheno. approach with **generic mixing**  $V_{\ell N}$  and **Dirac or Majorana**  $N$  mass eigenstates

Atre, Han, Pascoli, Zhang [0901.3589]

$$\underbrace{\nu_{\ell L}}_{\text{flavor basis}} \approx \underbrace{\sum_{m=1}^3 U_{\ell m} \nu_m + V_{\ell m'=4} N_{m'=4}}_{\text{mass basis}}$$

The SM  $W$  chiral coupling to **leptons** in **flavor basis** is

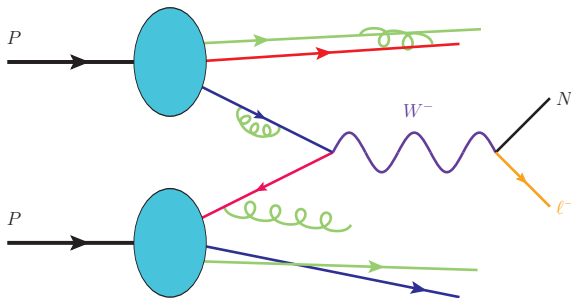
$$\mathcal{L}_{\text{Int.}} = -\frac{g_W}{\sqrt{2}} W_{\mu}^{-} \sum_{\ell=e}^{\tau} [\bar{\ell} \gamma^{\mu} P_L \nu_{\ell}] + \text{H.c.}, \quad \text{where } P_L = \frac{1}{2}(1 - \gamma^5)$$

$\implies$  SM  $W$  coupling to  $N$  and charged **leptons** in the **mass basis** is

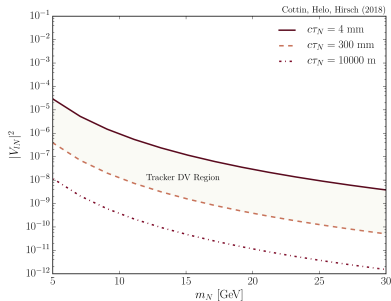
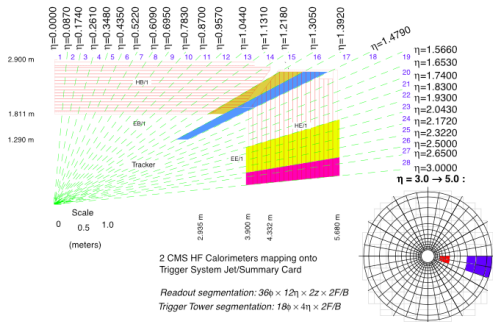
$$\mathcal{L}_{\text{Int.}} = -\frac{g_W}{\sqrt{2}} W_{\mu}^{-} \sum_{\ell=e}^{\tau} \left[ \bar{\ell} \gamma^{\mu} P_L \left( \sum_{m=1}^3 U_{\ell m} \nu_m + V_{\ell N} N \right) \right] + \text{H.c.}$$

$\implies$   $N$  is **accessible through**  $W/Z/H$  currents

## searching for long-lived heavy neutrinos at colliders<sup>1</sup>



<sup>1</sup>for reviews, see w/ Y. Cai, T. Li, T. Han [1711.02180], and w/ S. Pascoli, C. Weiland [1812.08750]



Decays of light  $N$  through SM weak currents can be very long-lived:

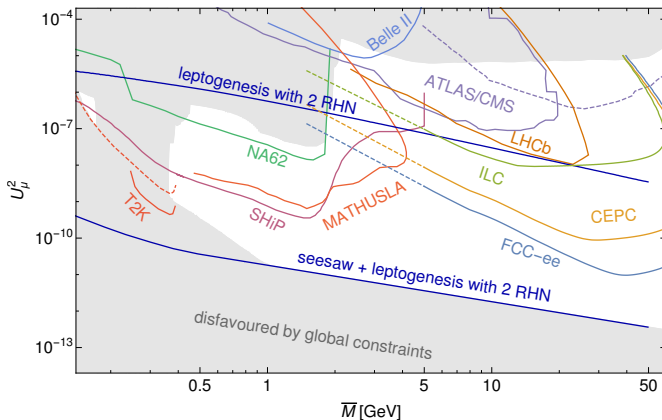
$$\Gamma_{\text{Tot.}} \sim G_F^2 m_N^5 \sum |V|^2 \quad (\text{small } |V| \implies \text{long lifetime!})$$

$$\implies d_0 = \beta c\tau = \frac{\beta c\hbar}{\Gamma_{\text{Tot.}}} \sim \frac{1.45 \text{ m}}{\sum |V|^2} \left( \frac{1 \text{ GeV}}{m_N} \right)$$

(Near) detectors have *finite* detector volume, with radius  $< \mathcal{O}(10) \text{ m}$

- $N$  may decay in **ECAL** (1-2m), **HCAL** (2-3m),  $\mu$ Chamber ( $>5\text{m}$ )

**Community Message:** Current and next-gen. facilities can directly test *simplest* resonant leptogenesis scenarios with high-scale Type I Seesaw



Update of Drewes, et al [[1609.09069](#)]

**Note:** LHC picture evolving with better strategies and add'l channels

Cottin, Helo, et al [[1806.05191](#)]; Abada, Bernal, et al [[1812.01720](#)]; Cheung, Ishida, et al [[2004.11537](#)]

**“Just how can I do this great physics?”**  
**- Anonymous audience member**

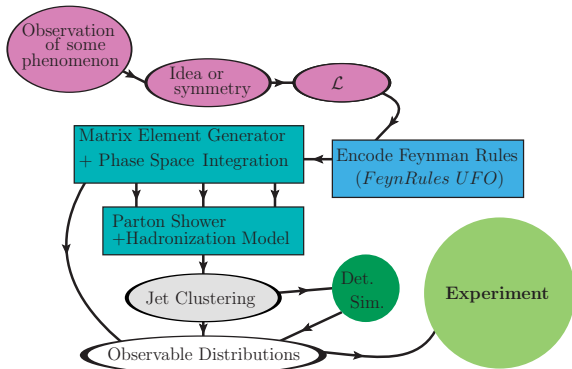
Searching for heavy Majorana and (pseudo)Dirac neutrinos  $N$  at LHC experiments follows any other analysis chain

- BIG PROBLEM (neutrino masses!) ✓
- BIG SOLUTION (sterile neutrinos!) / model / Lagrangian ✓
- Signal simulation with favorite generator ☹  
e.g., MadGraph5, Whizard, SHERPA
- Background simulation with favorite generator ✓
- Design long-lived particle analysis ☹
- Collect data, control regions, unblind, do some statistics ☹

# Major effort over past five-six years to make it possible to **simulation** Seesaw models with your favorite generator

- **Universal FeynRules Object (UFO) libraries** encode Feynman rules (.py) that plug into mainstream event generators, e.g., MadGraph5, Whizard, SHERPA

Alloul, Christensen, Duhr, Degrande, and Fuks [feynrules.irmp.ucl.ac.be](http://feynrules.irmp.ucl.ac.be)





# HeavyN: feynrules.irmp.ucl.ac.be/wiki/HeavyN

TypeIISeesaw, EffLRSM, WZPrime, SMWeinberg also available from [feynrules.irmp.ucl.ac.be/wiki/NLOModels!](https://feynrules.irmp.ucl.ac.be/wiki/NLOModels)

## HeavyN : The Standard Model + Heavy Neutrinos at NLO in QCD

### Contact Author

Richard Ruiz

- Universite Catholique de Louvain
- richard.ruiz AT uclouvain.be

In collaboration with: Daniel Alva and Tao Han [ 1 ]; Celine Degrande, Olivier Mattelear, and Jessica Turner [ 2 ]; and Silvia Pascoli and Cedric Weiland [ 3, 4 ].

For detailed instructions and examples on using the HeavyN UFO libraries, see C. Degrande, et al, ⇨ [arXiv:1602.06957](https://arxiv.org/abs/1602.06957) and S. Pascoli, et al, ⇨ [arXiv:1812.08750](https://arxiv.org/abs/1812.08750)

- For studies of heavy Majorana neutrinos, please consider citing [ 5 ] for the Lagrangian and [ 1, 2 ] for the Majorana FR/UFO files.
- For studies of heavy Dirac neutrinos, please consider citing [ 2, 4 ].

### Model Description

#### Majorana

This effective/simplified model extends the Standard Model (SM) field content by introducing three right-handed (RH) neutrinos, which are singlets under the SM gauge or weak hypercharge charges). Each RH neutrino possesses one RH Majorana mass. After electroweak symmetry breaking, the Lagrangian with three heavy Majorana [ 5 ]

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_N + \mathcal{L}_{N \text{ Int.}} \quad (1)$$

The first term is the Standard Model Lagrangian. In the mass basis, i.e., after mixing with active neutrinos, the heavy Majorana neutrinos' kinetic and mass terms are

$$\mathcal{L}_N = \frac{1}{2} \overline{N_{ki}} \not{\partial} N_k - \frac{1}{2} m_{N_k} \overline{N}_k N_k, \quad k = 1, \dots, 3, \quad (1)$$

and its interactions with the Weak gauge and Higgs bosons are given by

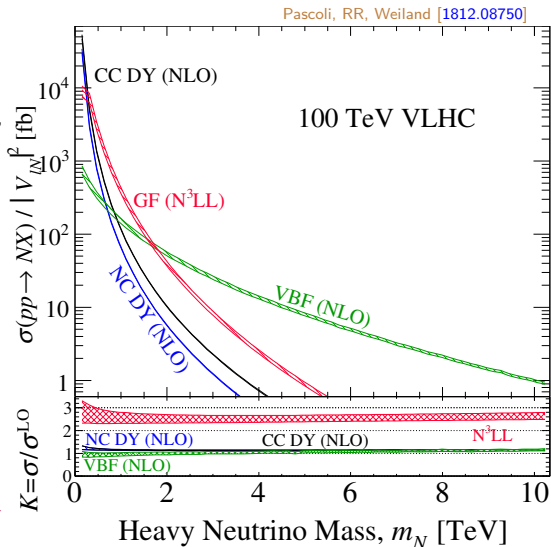
**time for some plots**

# FeynRules to MadGraph5aMC@NLO

Given a *Universal FeynRules Object* (UFO) file, run mg5amc out of the box

```
$ ./bin/mg5_aMC
> import model SM_HeavyN_NLO
> define p = g u c d s u~ c~ d~
s~ a
> define ell = mu+ mu-
> generate p p > n2 ell [QCD]
> output PP_Nmu_NLO
> launch PP_Nmu_NLO
> order=NLO
> fixed_order=ON
> set LHC 100
> set vmun2 1.0
> set mn2 scan:range(5,1001,25)
> set wn2 auto
```

$\mathcal{O}(10)$  lines to get each curve  $\rightarrow$



what about  $e^+e^-$  or  $\mu^+\mu^-$  collisions?

# Making LO(+PS) events for planned lepton colliders as easy as for LHC!

output from cmds on the left!

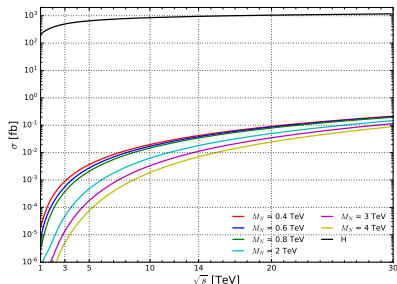
Results in the SM\_HeavyN\_NLO for  $e^+ e^- \rightarrow z > N1 \nu\nu$  QCD=0 QED=2,  $n1 > ell j j$

```

$ ./bin/mg5_aMC
import model SM_HeavyN_NLO
define vv = ve ve~
define ell = e+ e-
generate e+ e- > z > N1 vv
          QCD=0 QED=2, n1 > ell j j
output Test_FCee_ee_Nv_XLOPS
launch
shower=py8
set mz 92.0
set VeN1 1e-4
set mN1 scan:range(10,61,10)
set wn1 auto
set nevents 1k
set ebeam 46.0
set time_of_flight 0.0
set no_parton_cut
done
    
```

Available Results

Run	Collider	Banner	Cross section (pb)	Events	Data	Output	Action
run_01	e+ e- 46.0 x 46.0 GeV	ban_1	5.136e-05 ± 2.8e-07	1000	parton madevent	LHE	remove run launch detector simulation
					pythia8	LOG HEPMC	remove run launch detector simulation
run_02	e+ e- 46.0 x 46.0 GeV	ban_1	4.855e-05 ± 3.7e-07	1000	parton madevent	LHE	remove run launch detector simulation
					pythia8	LOG HEPMC	remove run launch detector simulation
run_03	e+ e- 46.0 x 46.0 GeV	ban_1	4.38e-05 ± 3.6e-07	1000	parton madevent	LHE	remove run launch detector simulation
					pythia8	LOG HEPMC	remove run launch detector simulation
run_04	e+ e- 46.0 x 46.0 GeV	ban_1	3.776e-05 ± 3e-07	1000	parton madevent	LHE	remove run launch detector simulation
					pythia8	LOG HEPMC	remove run launch detector simulation
run_05	e+ e- 46.0 x 46.0 GeV	ban_1	3.037e-05 ± 1.8e-07	1000	parton madevent	LHE	remove run launch detector simulation
					pythia8	LOG HEPMC	remove run launch detector simulation
run_06	e+ e- 46.0 x 46.0 GeV	ban_1	2.137e-05 ± 1.3e-07	1000	parton madevent	LHE	remove run launch detector simulation
					pythia8	LOG HEPMC	remove run launch detector simulation

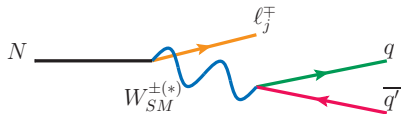


$W^+ W^- \rightarrow e_i^+ e_j^-$  possible at muon colliders! [2005.10289]

# mg5amc+MadSpin+Parton Shower

If the **narrow width approximation** is justified ( $\Gamma_N/m_N \ll 1$ ), efficient generation of  $e^+e^- \rightarrow Z \rightarrow \nu N \rightarrow \nu \ell^\pm q \bar{q}'$  possible with MadSpin:

Spin-correlation fully treated, RR [2008.01092]



In madspin\_card.dat, write:

```
set spinmode onshell
define q = u c d s u~ c~ d~ s~
define ee = e+ e-
decay n1 > ee q q
launch
```

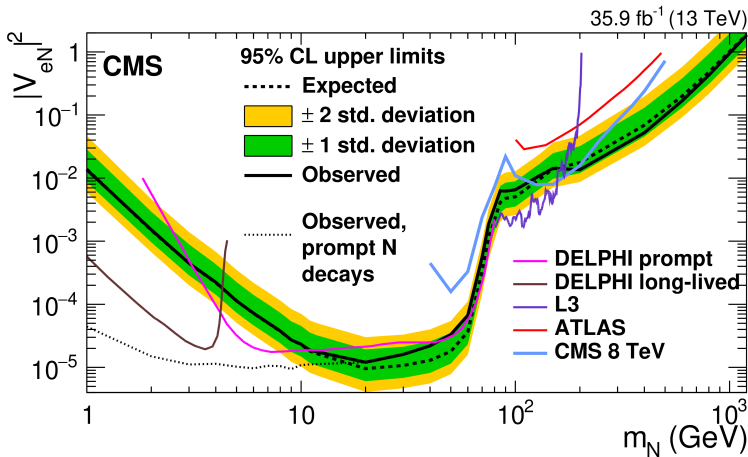
Parton showering with PY8 or HERWIG straightforward

**Fun Fact:** possible to steer entire process with a script →

```
rruiz@mac-1R0-359:~/Scripts/MG5aMC$ more runEffLRSMnlo_pp_Ne_Update.txt
launch EffLRSMnlo_pp_wr_Ne_NLO
order=NLO
shower=PY8
madspin=ON
done
set mwr 4000
set mn1 100
compute_widths wr+
compute_widths n1
set no_parton_cut
set nevents 100k
set LHC 13
set shower_card nsplit_jobs 100
set shower_card ue_enabled true

launch EffLRSMnlo_pp_wr_Ne_NLO
order=LO
shower=ON
```

Plotted: LHC 13 limits in search for  $pp \rightarrow 3\ell + MET$  ( $\ell_X = e, \mu$ )



- HeavyN used in CMS trilepton [1802.02965] and dilepton [1806.10905] searches
- ATLAS is currently adopting software (JobOptions already available!)

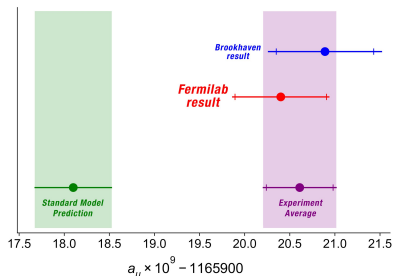
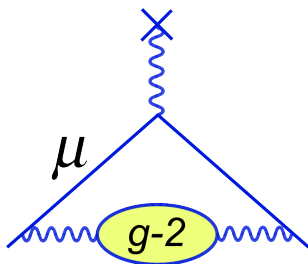
**NEW:**<sup>2</sup> sterile neutrinos  $N$  and  $\Delta a_\mu$

---

<sup>2</sup>V. Cirigliano, W. Dekens, J. de Vries, K. Fuyuto, E. Mereghetti [2105.11462] 



# Anomalous magnetic moment of the $\mu$



Fermilab's Muon  $g-2$  has *confirmed* that  $a_\mu = (g_\mu - 2)/2$  is *\*a bit\** large

[2104.03281]

$$a_\mu^{\text{average}} = (116\,592\,061 \pm 41) \cdot 10^{-11}$$

$$a_\mu^{\text{SM}} = (116\,591\,810 \pm 43) \cdot 10^{-11}$$

The difference? Large enough to start taking BSM solutions seriously.

$$\Delta a_\mu = (251 \pm 59) \cdot 10^{-11} \text{ or about } 4.2\sigma!$$

Can new  $N$  interactions account for this?

Yes, in a surprisingly succinct manner.

$\nu$ SMEFT is the Standard Model Effective Field Theory extended by  $\nu_R$

$\psi^2 H^3$		$\psi^2 H^2 D$		$\psi^2 HX(+\text{H.c.})$	
$\mathcal{O}_{L\nu H}(+\text{H.c.})$	$(\bar{L}\nu_R)\tilde{H}(H^\dagger H)$	$\mathcal{O}_{H\nu}$	$(\bar{\nu}_R\gamma^\mu\nu_R)(H^\dagger i\overleftrightarrow{D}_\mu H)$	$\mathcal{O}_{\nu B}$	$(\bar{L}\sigma_{\mu\nu}\nu_R)\tilde{H}B^{\mu\nu}$
		$\mathcal{O}_{H\nu e}(+\text{H.c.})$	$(\bar{\nu}_R\gamma^\mu e)(H^\dagger i\overleftrightarrow{D}_\mu H)$	$\mathcal{O}_{\nu W}$	$(\bar{L}\sigma_{\mu\nu}\nu_R)\tau^I\tilde{H}W^{I\mu\nu}$
$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$		$(\bar{L}R)(\bar{L}R)(+\text{H.c.})$	
$\mathcal{O}_{\nu\nu}$	$(\bar{\nu}_R\gamma^\mu\nu_R)(\bar{\nu}_R\gamma_\mu\nu_R)$	$\mathcal{O}_{L\nu}$	$(\bar{L}\gamma^\mu L)(\bar{\nu}_R\gamma_\mu\nu_R)$	$\mathcal{O}_{L\nu Le}$	$(\bar{L}\nu_R)\epsilon(\bar{L}e)$
$\mathcal{O}_{e\nu}$	$(\bar{e}\gamma^\mu e)(\bar{\nu}_R\gamma_\mu\nu_R)$	$\mathcal{O}_{Q\nu}$	$(\bar{Q}\gamma^\mu Q)(\bar{\nu}_R\gamma_\mu\nu_R)$	$\mathcal{O}_{L\nu Qd}$	$(\bar{L}\nu_R)\epsilon(\bar{Q}d)$
$\mathcal{O}_w$	$(\bar{u}\gamma^\mu u)(\bar{\nu}_R\gamma_\mu\nu_R)$			$\mathcal{O}_{LdQ\nu}$	$(\bar{L}d)\epsilon(\bar{Q}\nu_R)$
$\mathcal{O}_{d\nu}$	$(\bar{d}\gamma^\mu d)(\bar{\nu}_R\gamma_\mu\nu_R)$				
$\mathcal{O}_{d\nu e}(+\text{H.c.})$	$(\bar{d}\gamma^\mu u)(\bar{\nu}_R\gamma_\mu e)$				
$(\bar{L}R)(\bar{R}L)$		$(\bar{L}\cap B)(+\text{H.c.})$		$(\bar{L}\cap \mathcal{B})(+\text{H.c.})$	
$\mathcal{O}_{Q\nu L}(+\text{H.c.})$	$(\bar{Q}u)(\bar{\nu}_R L)$	$\mathcal{O}_{\nu\nu\nu}$	$(\bar{\nu}_R^i\nu_R)(\bar{\nu}_R^j\nu_R)$	$\mathcal{O}_{QQd\nu}$	$\epsilon_{ij}\epsilon_{\alpha\beta\sigma}(Q_\alpha^i C Q_\beta^j)(d_\sigma C\nu_R)$
				$\mathcal{O}_{udd\nu}$	$\epsilon_{\alpha\beta\sigma}(u_\alpha C d_\beta)(d_\sigma C\nu_R)$

Table 1: The complete basis of dimension-six operators involving  $\nu_R$  taken from Ref. [24]. The operators are expressed in terms of a column vector of  $n$  gauge singlet fields,  $\nu_R$ , and of SM fields, the lepton and Higgs doublets,  $L$  and  $H$ , the quark left-handed doublet  $Q = (u_L, d_L)^T$ , and the right-handed fields  $e$ ,  $u$ , and  $d$ .

Unexpectedly, only one  $\nu$ SMEFT can generate the right  $\Delta a_\mu$

$$\mathcal{L}_{H\nu e} \approx \frac{g\nu^2}{2\sqrt{2}\Lambda^2} \sum_{k=1}^3 [\bar{C}_{H\nu e}]_{k\ell} (\bar{N}_k \gamma^\mu P_R \ell_R) W_\mu^+ \left(1 + \frac{h}{v}\right)^2 + \text{H.c.}$$

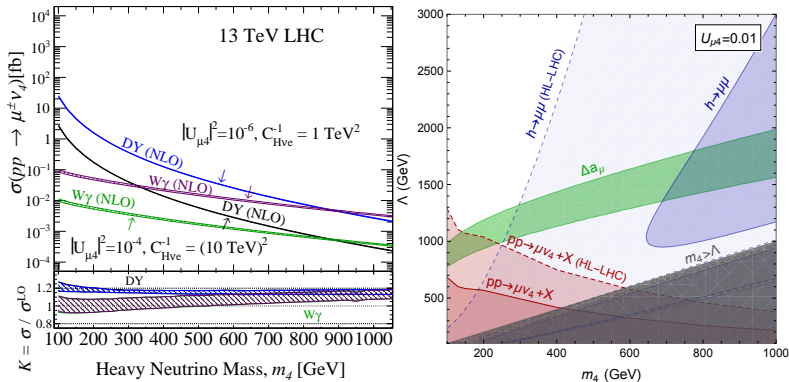
This generates  $\Delta a_\mu$  of the form

$$\Delta a_\mu \sim -\frac{2m_\mu m_N}{(4\pi)^2 \Lambda^2} \text{Re} \left( V_{\mu N} [\bar{C}_{H\nu e}]_{N\mu} \right) \quad (\text{see [2105.11462] for exact formula!})$$

# $\Delta a_\mu$ at the LHC

We created a new UFO `HeavyN_vSMEFTdim6` containing  $\mathcal{O}_{H\nu e}$

Already available from the FeynRules UFO database: [feynrules.irmp.ucl.ac.be/wiki/HeavyN](https://feynrules.irmp.ucl.ac.be/wiki/HeavyN)



**Conclusion:** If  $N$  are involved in  $\Delta a_\mu$ , then expect something in

$$pp \rightarrow N\mu^\pm + X \text{ and } H \rightarrow \mu^+\mu^-$$

in Run III data and at the HL-LHC (see the paper for more details! [2105.11462])

## summary

- Available: UFOs with Majorana or Dirac  $N$
- Available: UFOs with full CKM matrix or all quark masses
- Available: LO (XL0) and LO+NLO (NLO) Feynman rules
  
- NEW: (some)  $\nu$ SMEFT operators at dimension six
- In development: decays to/from vector mesons for  $m_N \lesssim 5$  GeV
- In development: more studies on spin-correlation
- In development: more studies on off-shell mediators

#### HeavyN : The Standard Model + Heavy Neutrinos at NLO in QCD

##### Contact Author

Richard Ruiz

- Universite Catholique de Louvain
- richard.ruiz AT uclouvain.be

In collaboration with: Daniel Alva and Tao Han [ 1 ]; Celine Degrande, Olivier Mattelear, and Jessica Turner [ 2 ]; and Silvia Pascoli and Cedric Weiland [ 3, 4 ].

For detailed instructions and examples on using the HeavyN UFO libraries, see C. Degrande, et al, [arXiv:1602.06957](https://arxiv.org/abs/1602.06957) and S. Pascoli, et al, [arXiv:1812.08750](https://arxiv.org/abs/1812.08750)

- For studies of heavy Majorana neutrinos, please consider citing [ 5 ] for the Lagrangian and [ 1, 2 ] for the Majorana FR/UFO files.
- For studies of heavy Dirac neutrinos, please consider citing [ 2, 4 ].

##### Model Description

###### Majorana

This effective/simplified model extends the Standard Model (SM) field content by introducing three right-handed (RH) neutrinos, which are singlets under the SM gauge symmetry (no color, or weak hypercharge charges). Each RH neutrino possesses one RH Majorana mass. After electroweak symmetry breaking, the Lagrangian with three heavy Majorana neutrinos  $N_i$  (for  $i=1,2,3$ ):

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_N + \mathcal{L}_{N \text{ Int}}. \quad (1)$$

The first term is the Standard Model Lagrangian. In the mass basis, i.e., after mixing with active neutrinos, the heavy Majorana neutrinos' kinetic and mass terms are

$$\mathcal{L}_N = \frac{1}{2} \overline{N_k} i \not{\partial} N_k - \frac{1}{2} m N_k \overline{N_k} N_k, \quad k = 1, \dots, 3, \quad (1)$$

and its interactions with the Weak gauge and Higgs bosons are given by

UFO	NLO	Spin	M or D	#	$V'$	LN $V$	LF $V$	arXiv
HeavyN	✓	✓	M	3		✓	✓	[1602.06957]
HeavyN_Dirac	✓	✓	D	3			✓	[1812.08750]
HeavyN_vSMEFTdim6	✓	✓	M	3		✓	✓	[2105.11462]
EffLRSM	✓	✓	M	3	✓	✓	✓	[1610.08985]
WZPrime	✓	✓			✓		✓	[1701.05263]
TypeIISeesaw	✓	✓			✓	✓	✓	[1912.08975]
SMWeinberg	✓	✓				✓	✓	[2012.09882]

## Legend:

- “NLO” = simulations at LO and NLO in QCD possible
- “Spin” = spin correlation fully described for details, see RR [2008.01092]
- “M or D” = Majorana or Dirac  $N$
- “#” = number of  $N$  in the model file
- “ $V'$ ” = new gauge, scalar, or pseudoscalar bosons

(other new particles can be long-lived!)

Happy to make more public, just ask!



**Thank you.**