

Quantum enhanced methods for ultralight dark matter searches



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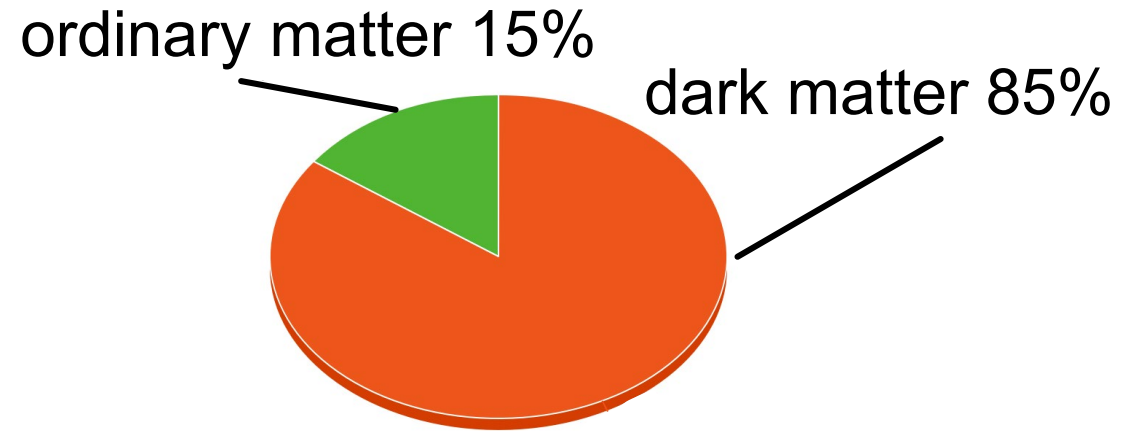


search for ultralight dark matter: a quantum metrology problem

quantum squeezing is already speeding up the search for axionic dark matter

quantum technology development may dramatically accelerate an axion search

Particle nature of dark matter remains unknown



dark matter:

85% of the matter in the universe

detected only gravitationally

cold (gravitationally bound)

mean density: $\sim 0.4 \text{ GeV/cm}^3$

Ultralight dark matter must be bosonic



ultralight dark matter (e.g. axions):

quantum degenerate Bose gas

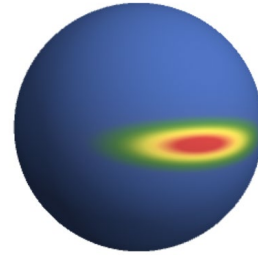
more wave-like than particle-like

tone: feeble, persistent, unknown frequency

Laboratory based searches for ultralight dark matter

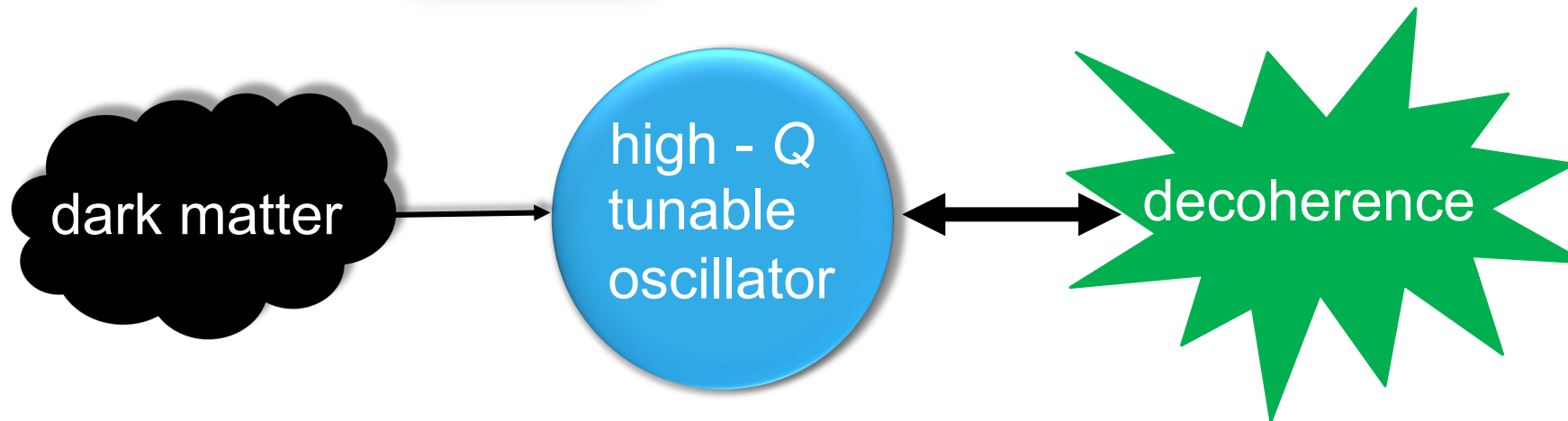
classical signal acting on a quantum harmonic oscillator

microwave
cavities



spin
ensembles

ADMX, HAYSTAC,
CASPEr, DMradio



evolution in rotating frame

null hypothesis

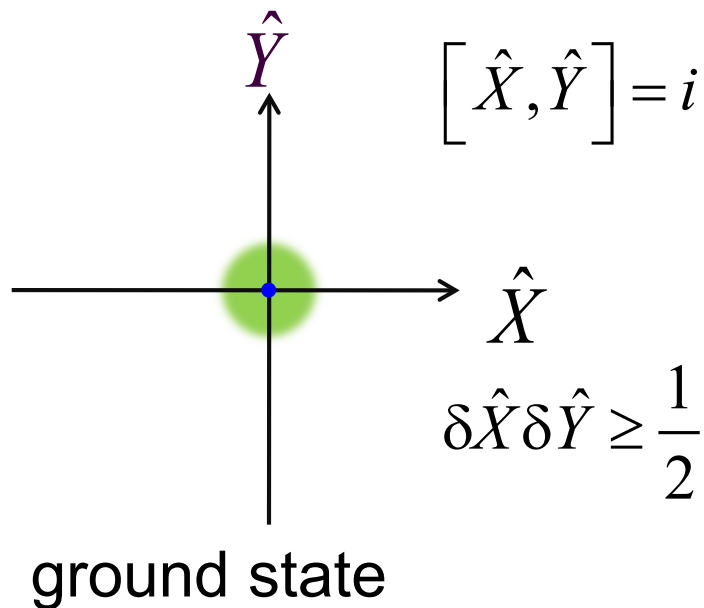
$$\hat{H}_r = 0$$

dark matter hypothesis

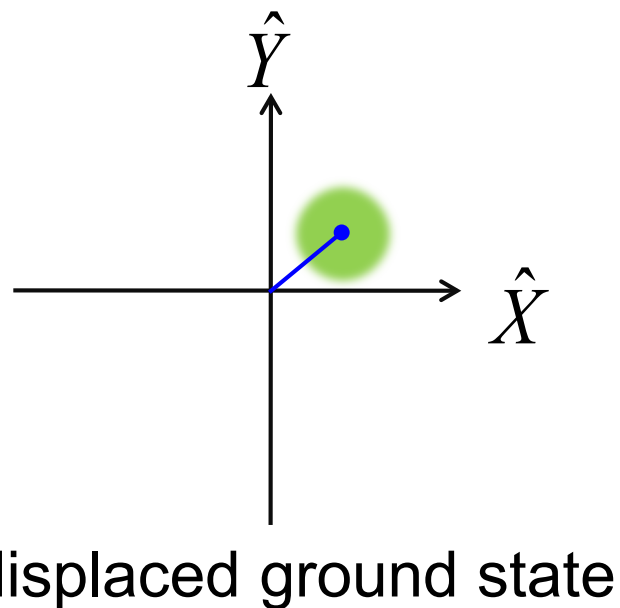
$$\hat{H}_r = F_y(t)\hat{X} + F_x(t)\hat{Y}$$

Quantum noise pollutes inference of classical force

null hypothesis



dark matter hypothesis

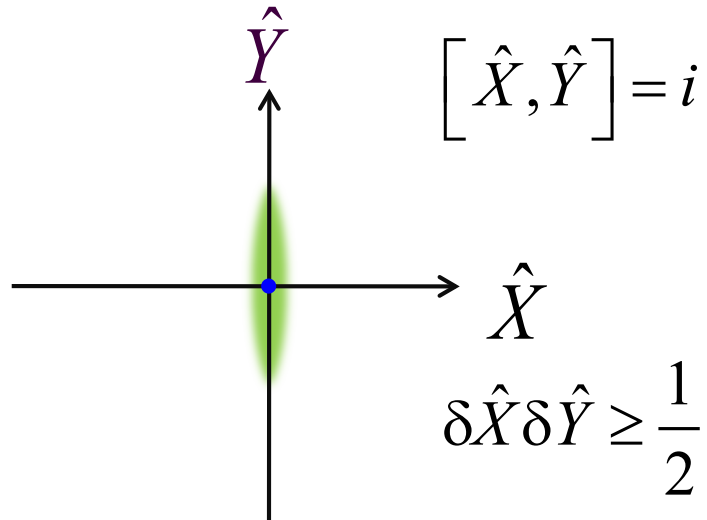


“quantum limit”--- coherent state limit (CSL)

prepare in ground state, evolve, measure X noiselessly

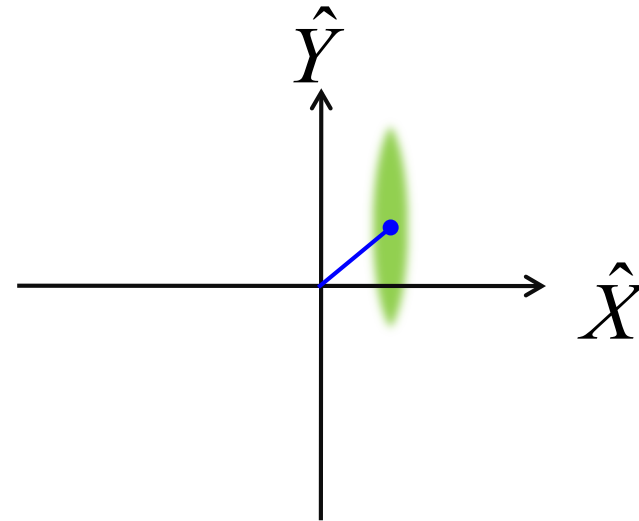
Squeezing and noiseless measurement circumvent quantum limit ⁷

null hypothesis



squeezed state

dark matter hypothesis



displaced squeezed state

beat the quantum limit

prepare in squeezed state, evolve, measure X noiselessly

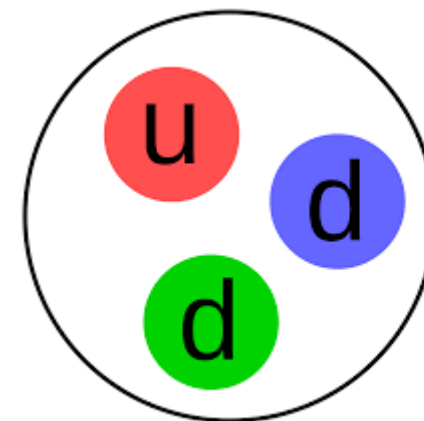
quantum enhanced axion search

Hypothetical QCD axion: a light particle that is cold and dense enough to contribute to dark matter

mechanism to resolve strong-CP problem (Peccei and Quinn)

Why is CP symmetry well-preserved in QCD?

Why is the neutron so round?



axion mass range

$$2 \mu\text{eV} < m_a c^2 < 2000 \mu\text{eV}$$

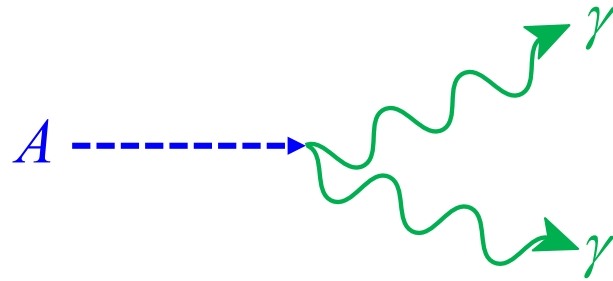
as a frequency

$$500 \text{ MHz}^* < m_a c^2 / h < 500 \text{ GHz}$$

*post-inflation scenario

Axion field couples to electromagnetism

modified QCD Lagrangian $\Rightarrow g_{A\gamma\gamma} A (\vec{E} \cdot \vec{B})$

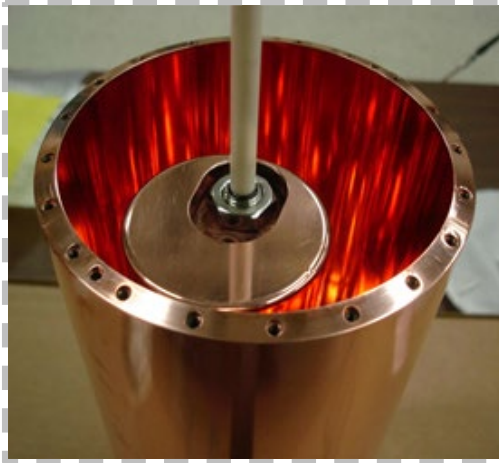


linearize coupling around static \vec{B} -field

$$A \text{---} \text{---} \text{---} \gamma \quad E_A = E_\gamma$$

dark matter axions create microwave photons

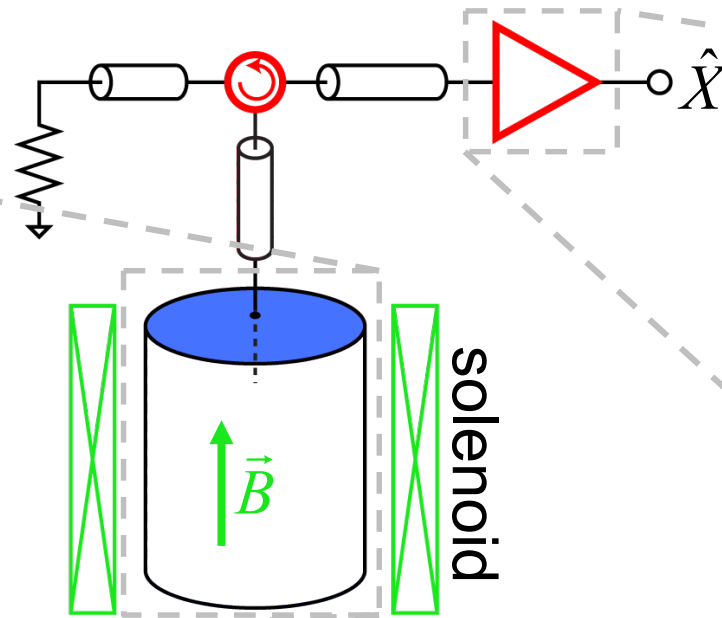
Scan cavity to search for resonant axion to photon conversion



1 L volume

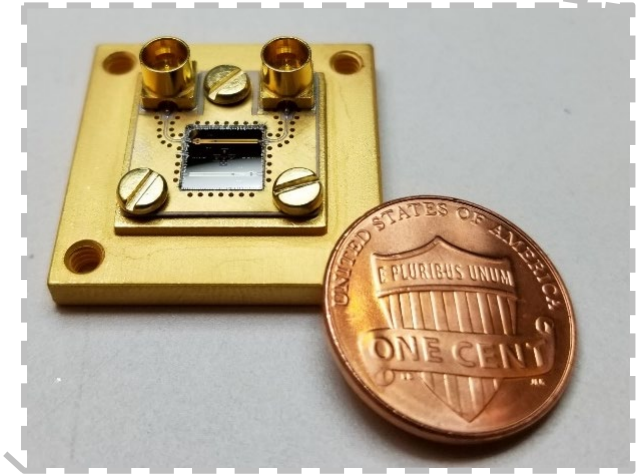
$$Q = 10^4$$

$$\omega_{\text{cav}} \sim 2\pi \times 5 \text{ GHz}$$



tunable cavity

$$T \ll \frac{\hbar \omega_{\text{cav}}}{k_B}$$

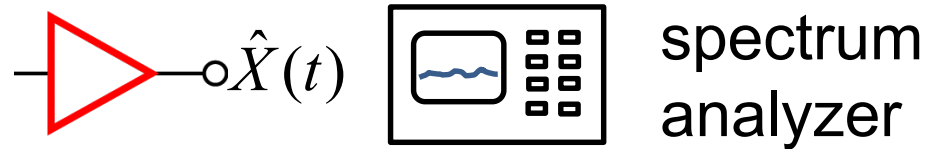


Josephson parametric
amplifier (JPA)

measure \hat{X} noiselessly

haloscope (Sikivie 1983) at the quantum limit (HAYSTAC 2017)

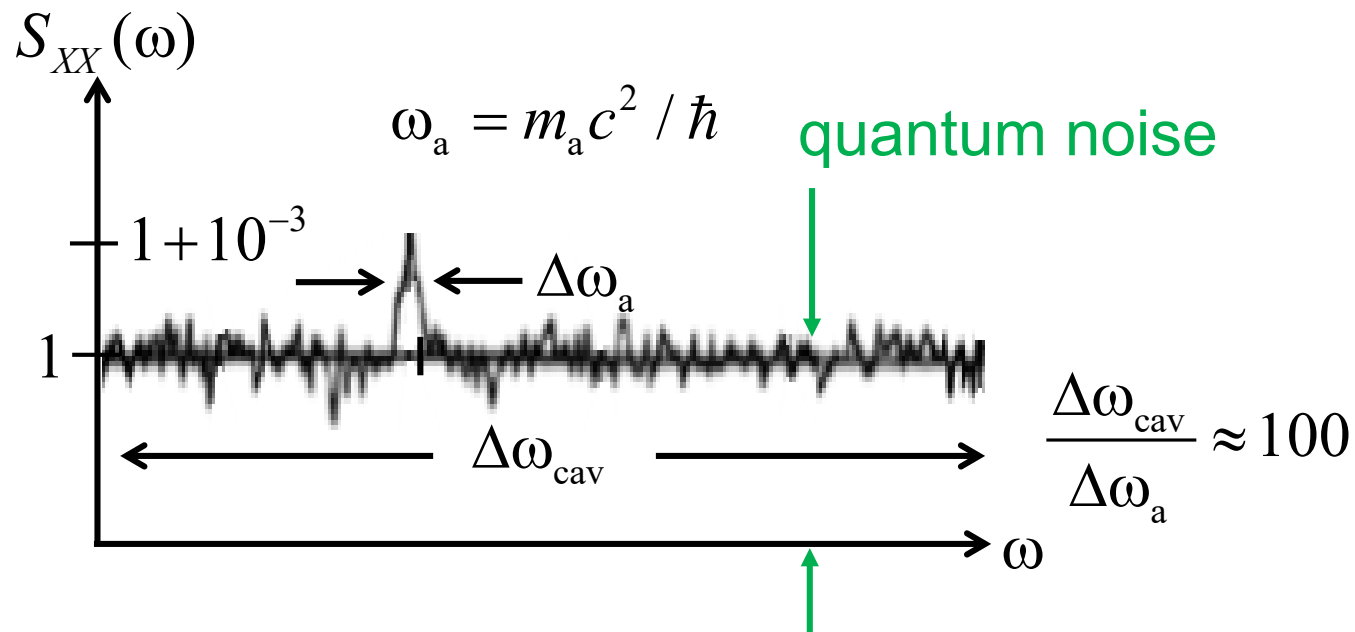
Average noise to resolve tiny excess axion power



10^{-3} noise precision $\Rightarrow 10^6$ measurements

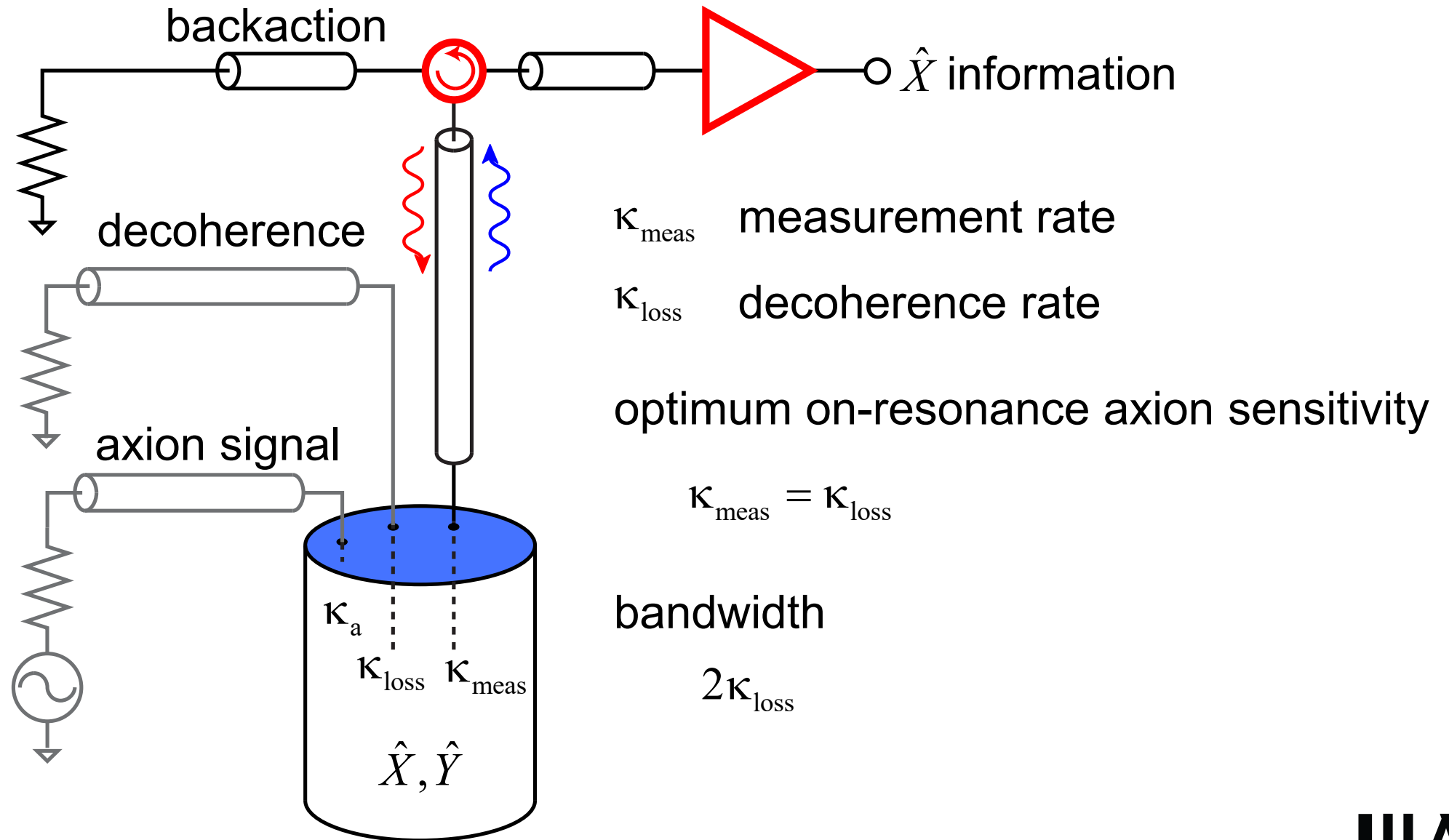
tune \rightarrow wait and average \rightarrow tune

100 axion bands per cavity tuning

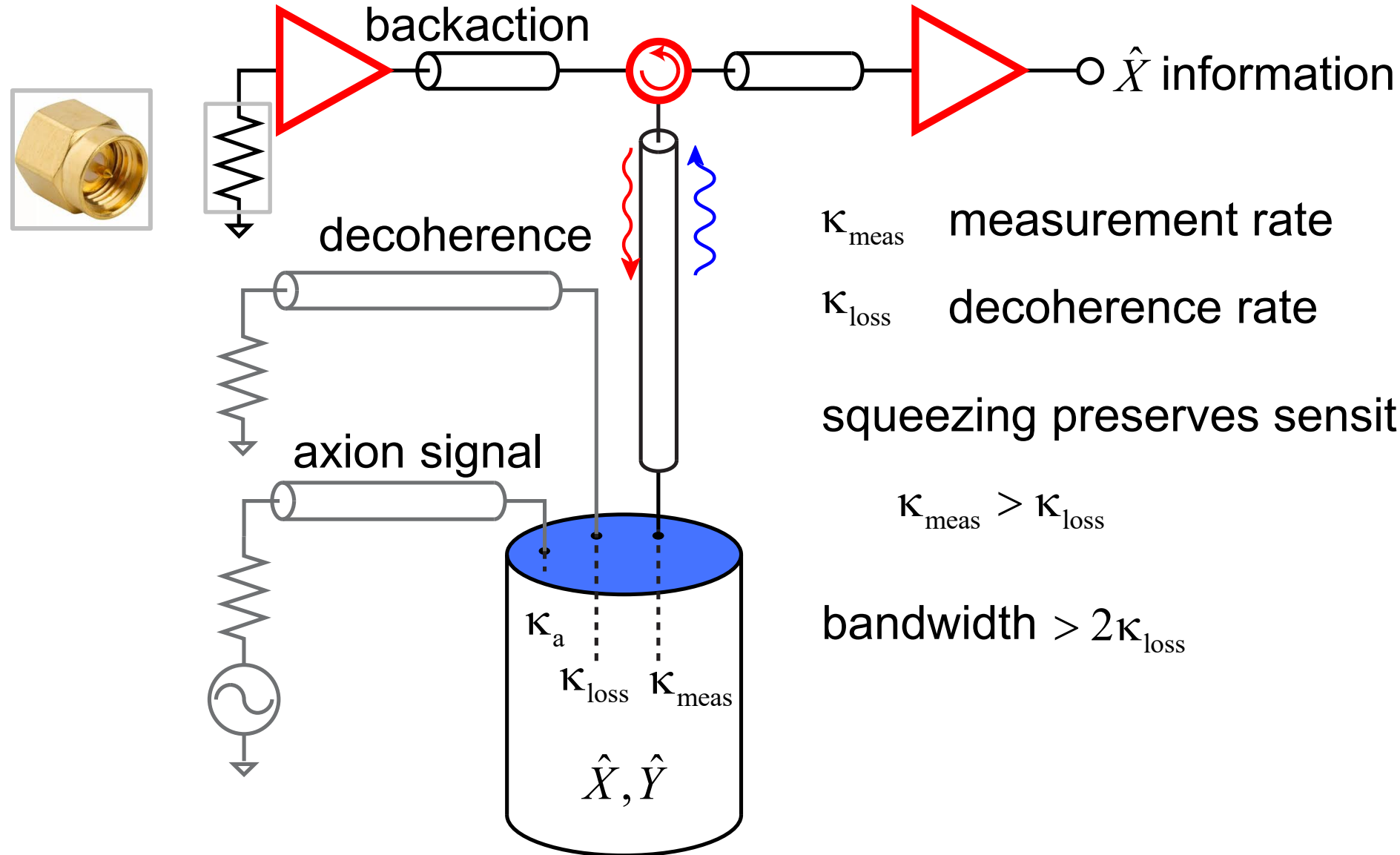


scan rate proportional to bandwidth

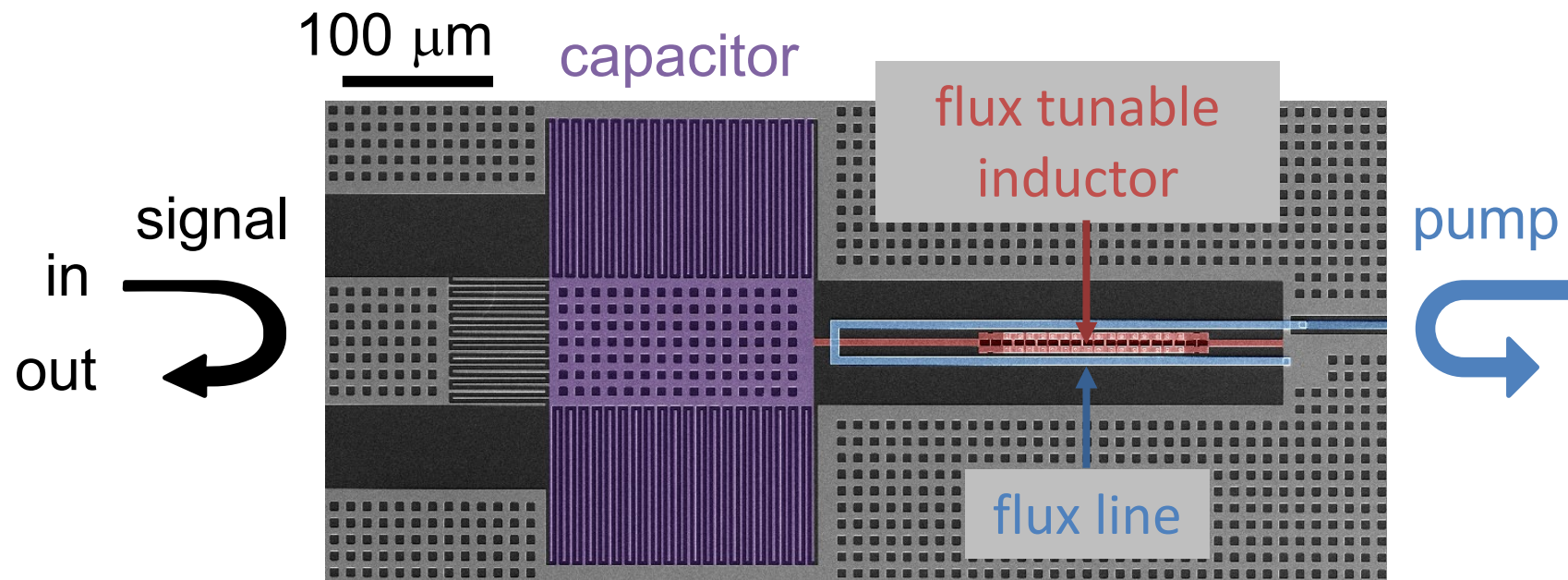
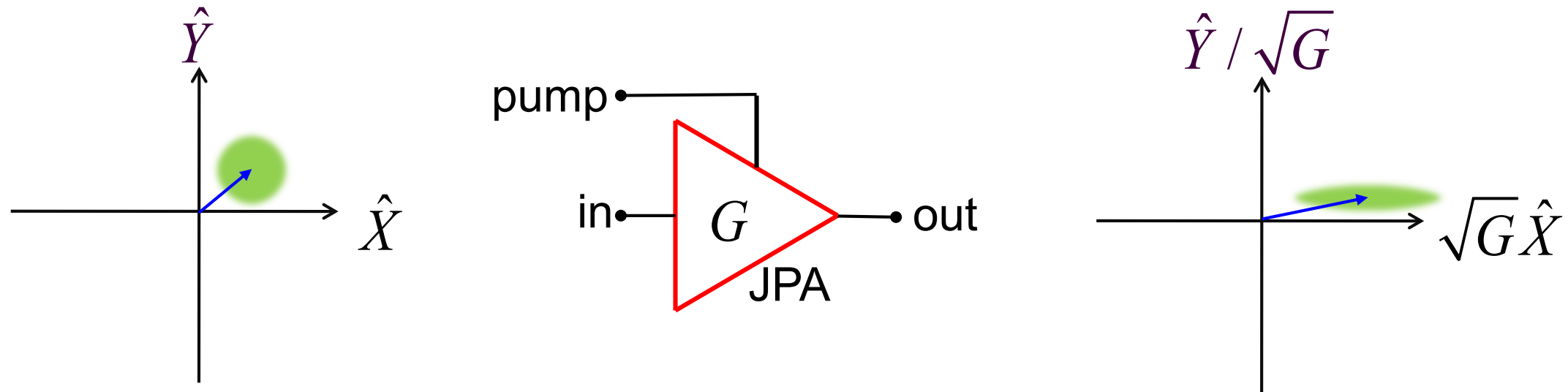
Measurement backaction limits haloscope bandwidth



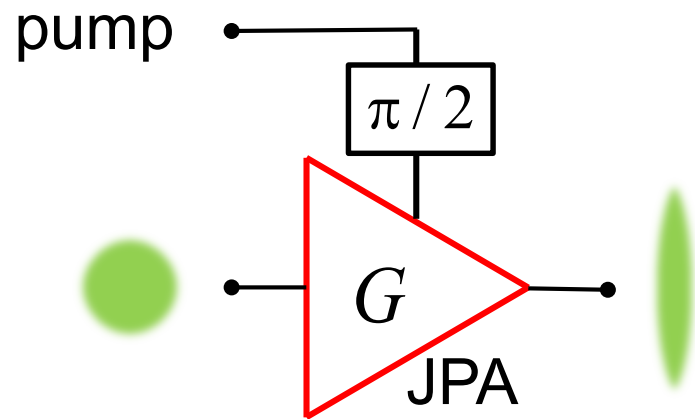
Squeezing yields larger bandwidth through backaction evasion



JPAAs measure one quadrature noiselessly



JPAAs prepare microwave squeezed states

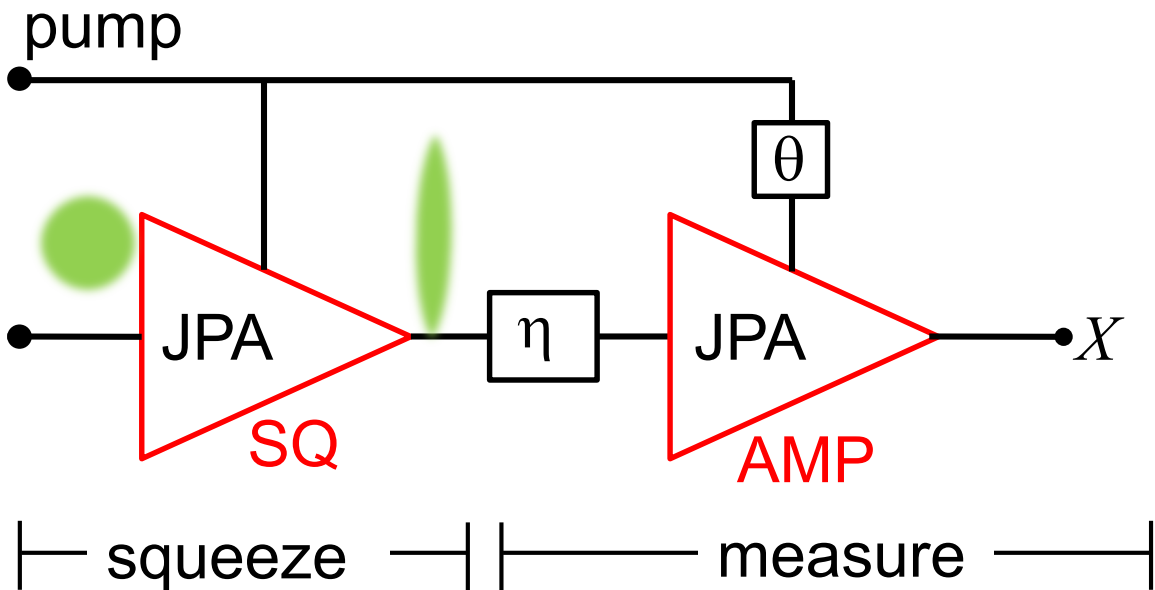


$$\delta^2 \hat{Y} = \frac{G}{2}$$

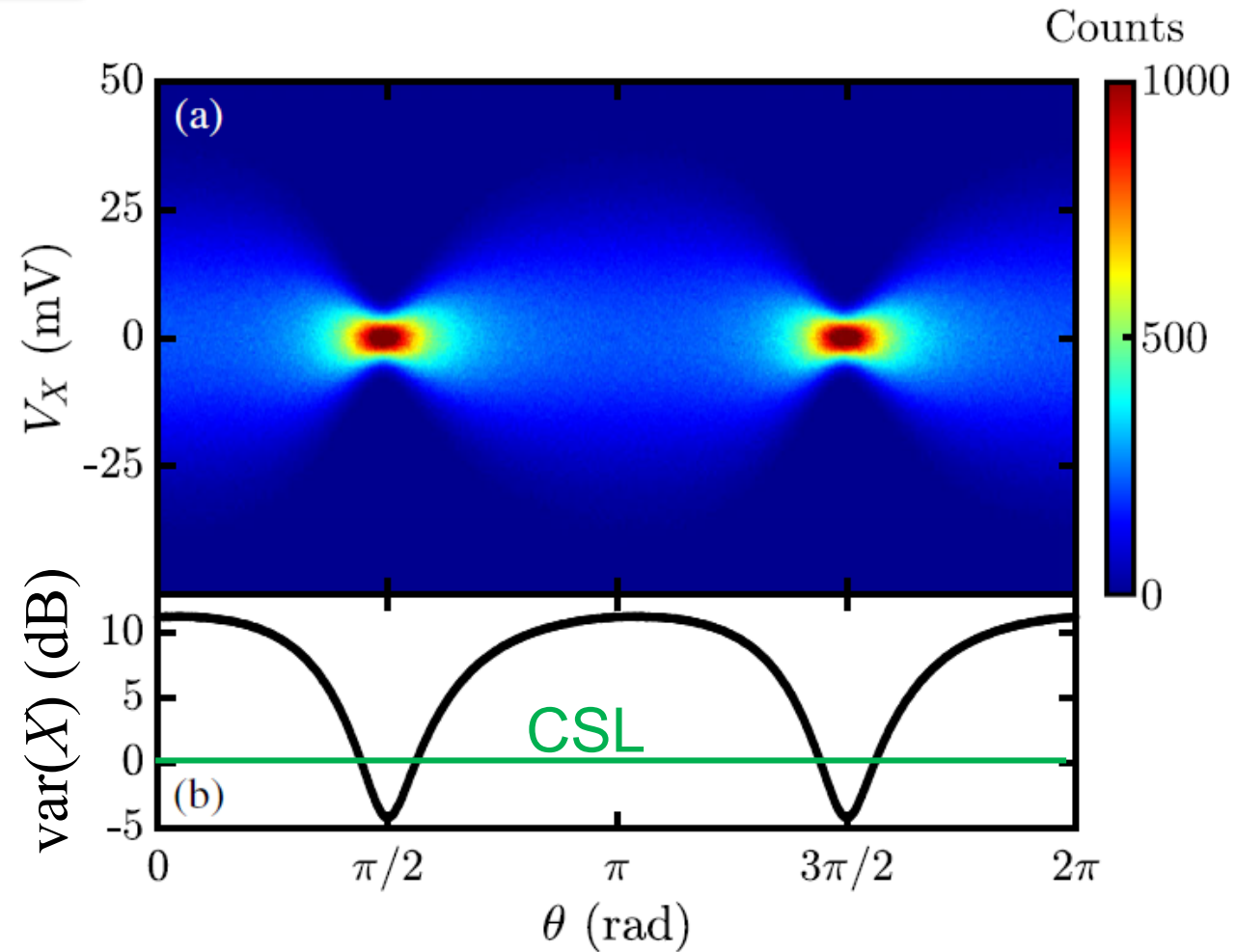
$$\delta^2 \hat{X} = \frac{1}{2G}$$

squeezed quadrature: quantum noise suppressed

Demonstration of squeezing and determination of loss: second JPA analyzes squeezed state created by first



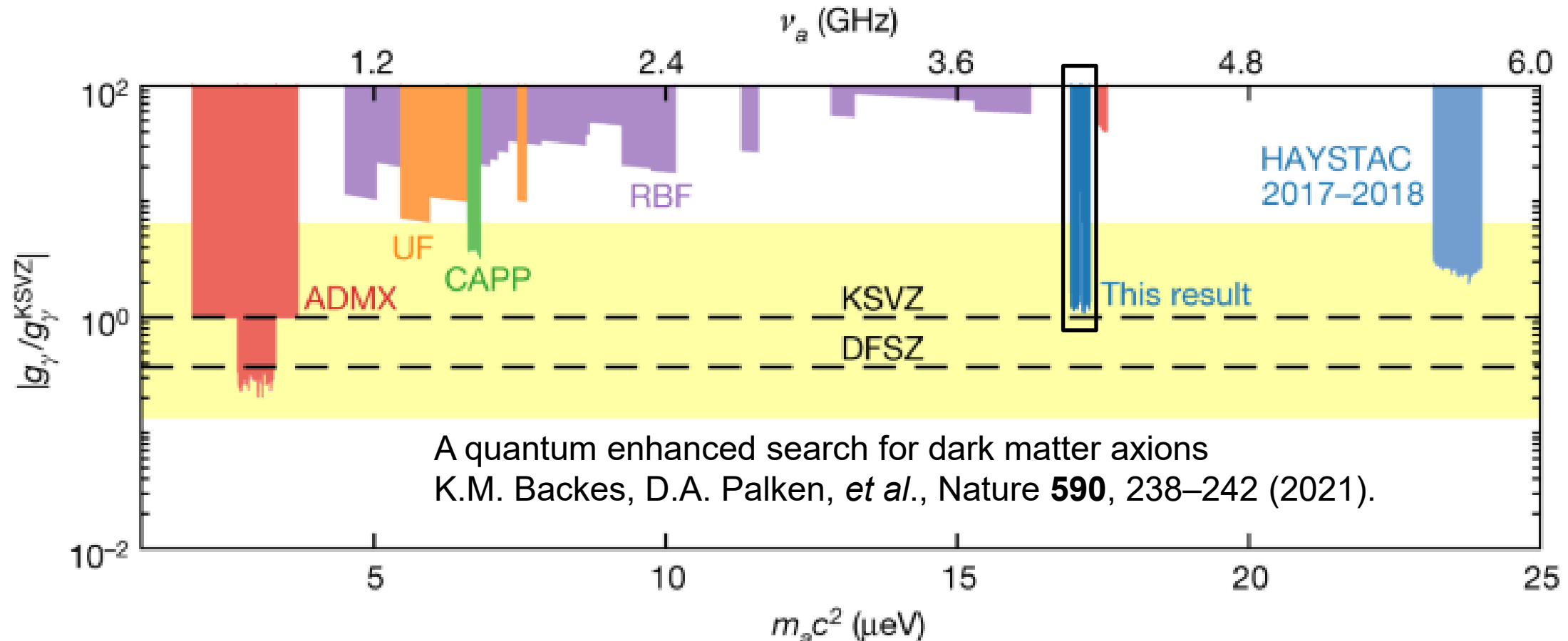
phase θ chooses measured quadrature relative to squeezed quadrature



loss η : 1.63 dB

First quantum enhanced dark matter search

double the CSL scan rate



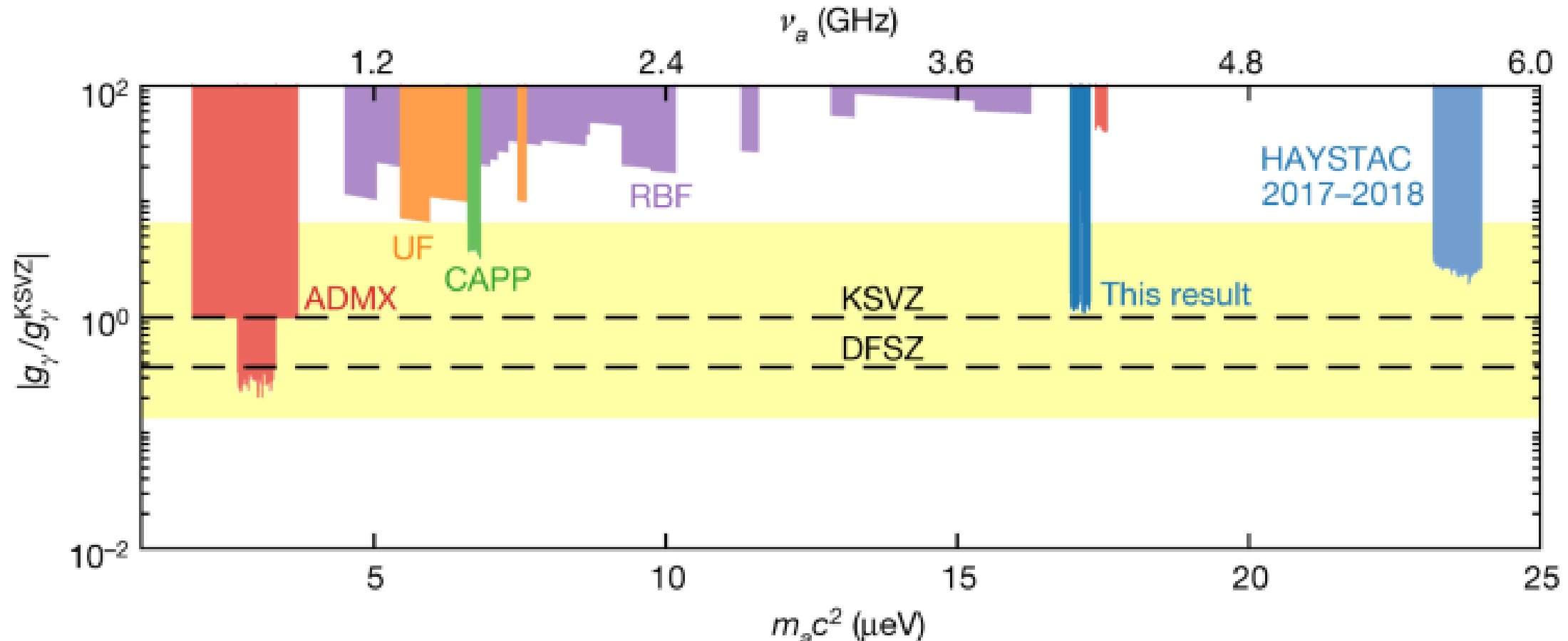
Prohibitive resources required to scan one decade at CSL

1 – 10 GHz at CSL for DFSZ

1 x 9T magnet
~20,000 yrs

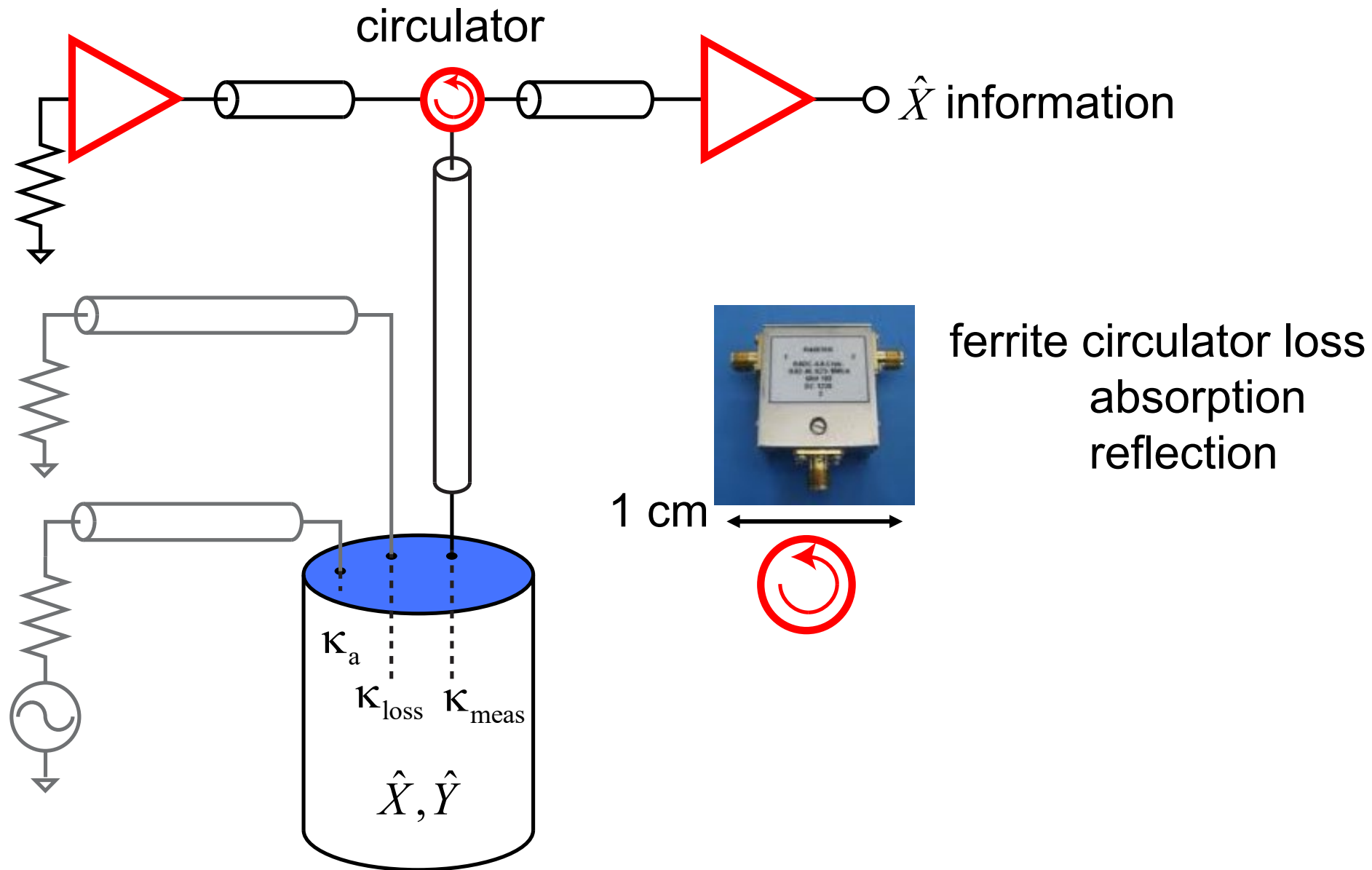
or

16 x 30T magnets
~10 yrs + 1.2 B\$



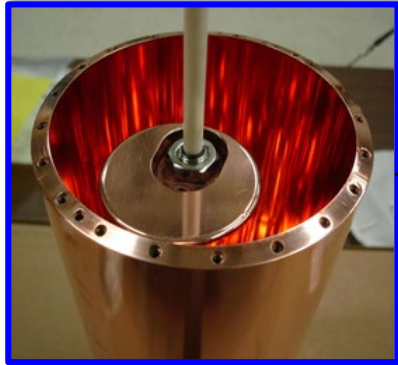
engineering more quantum enhancement

Squeezing diminished by circulator and cascaded signal path



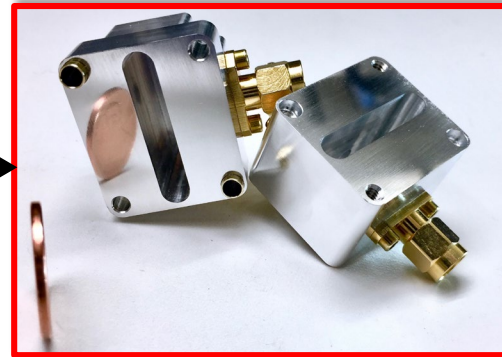
A better approach to backaction evading measurements

axion cavity



ω_A

readout cavity



ω_R

no ferrite circulators
no loss from reflection

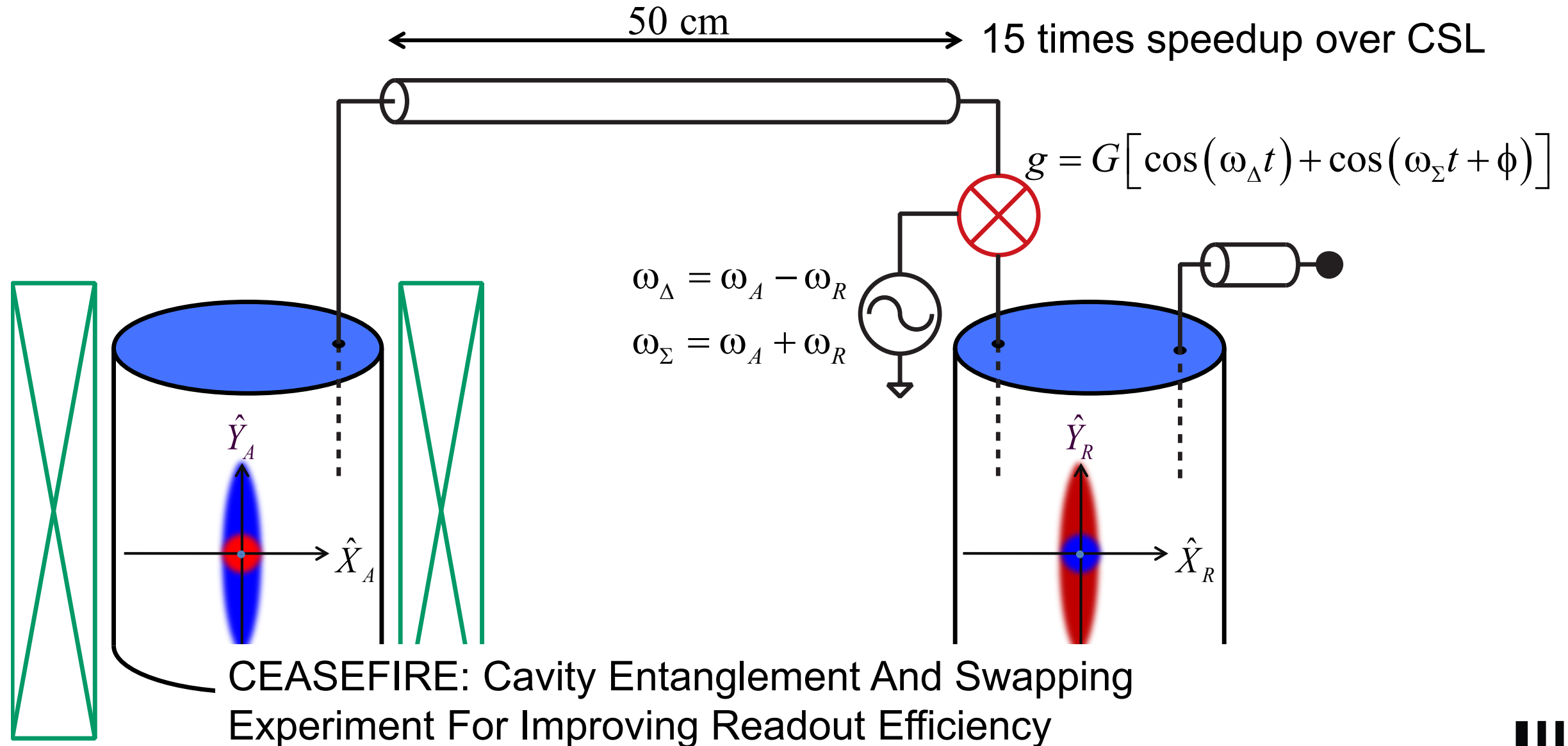
engineering QND interaction between two cavities

$$\hat{H}_I = \hbar G \hat{X}_R \hat{X}_A \quad \frac{d}{dt} \hat{Y}_R = -G \hat{X}_A$$

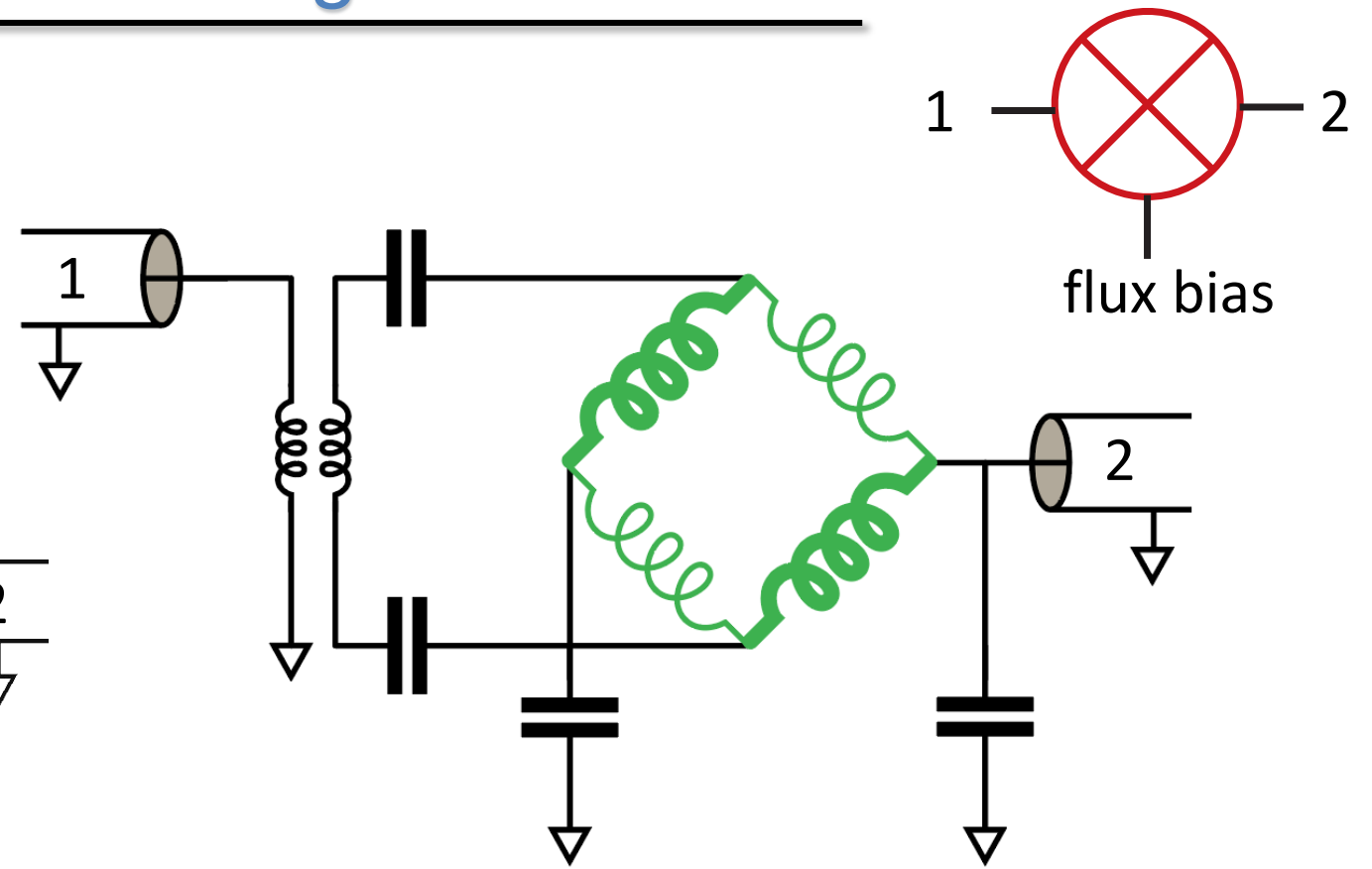
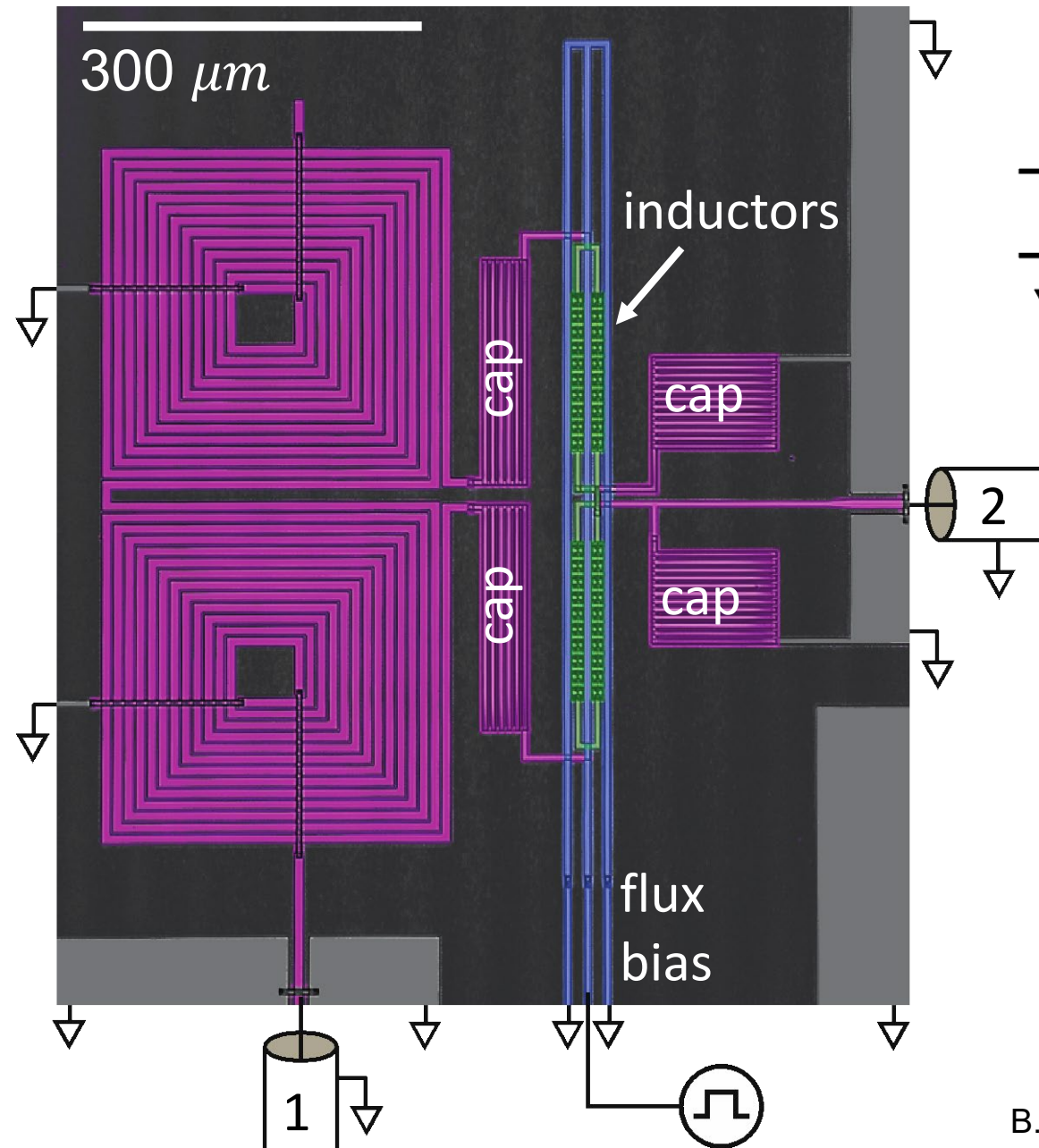
monitor \hat{Y}_R infer \hat{X}_A

backaction deposited in unmonitored quadrature

Two mode squeezing, state swapping yield QND Hamiltonian 23



Tunable coupler from an inductor bridge



tune bridge balance/imbalance with
one on-chip bias line

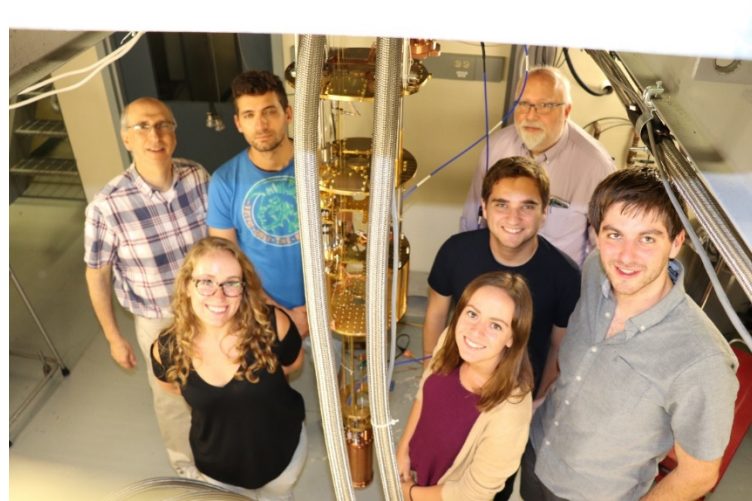
Conclusions

superconducting quantum technology
accelerates the search for dark matter

quantum squeezing doubles search rate in
HAYSTAC apparatus

greater speedup possible with circulator-less
concepts

Yale, August 2018,
SSR commissioning



CEASEFIRE

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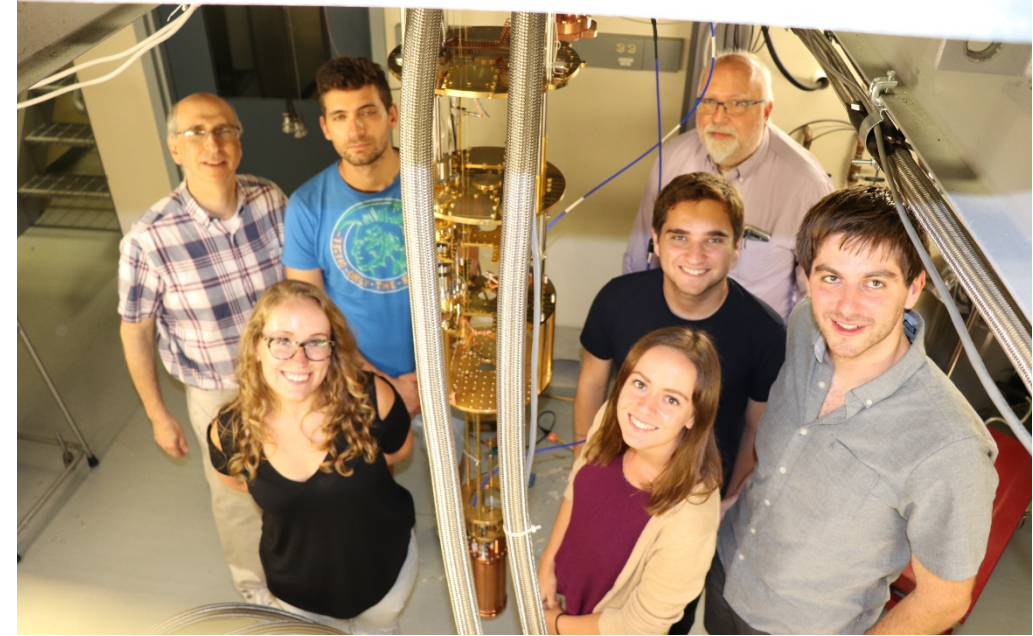
Conclusions

superconducting quantum technology can accelerate search for ultralight dark matter

quantum squeezing doubles search rate in demonstration

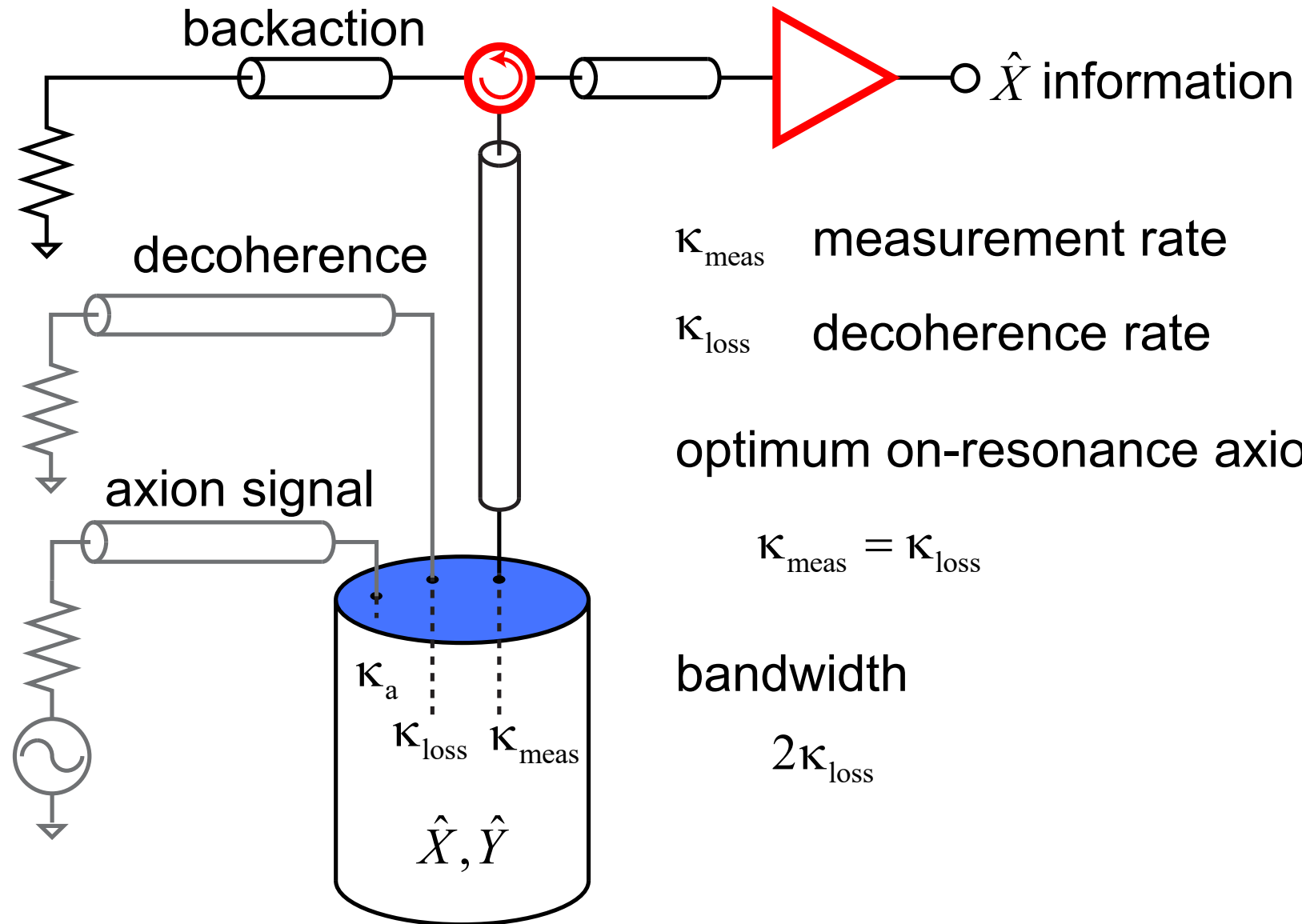
squeezed state receiver now operating in HAYSTAC apparatus

greater speedup possible with circulator-less concepts

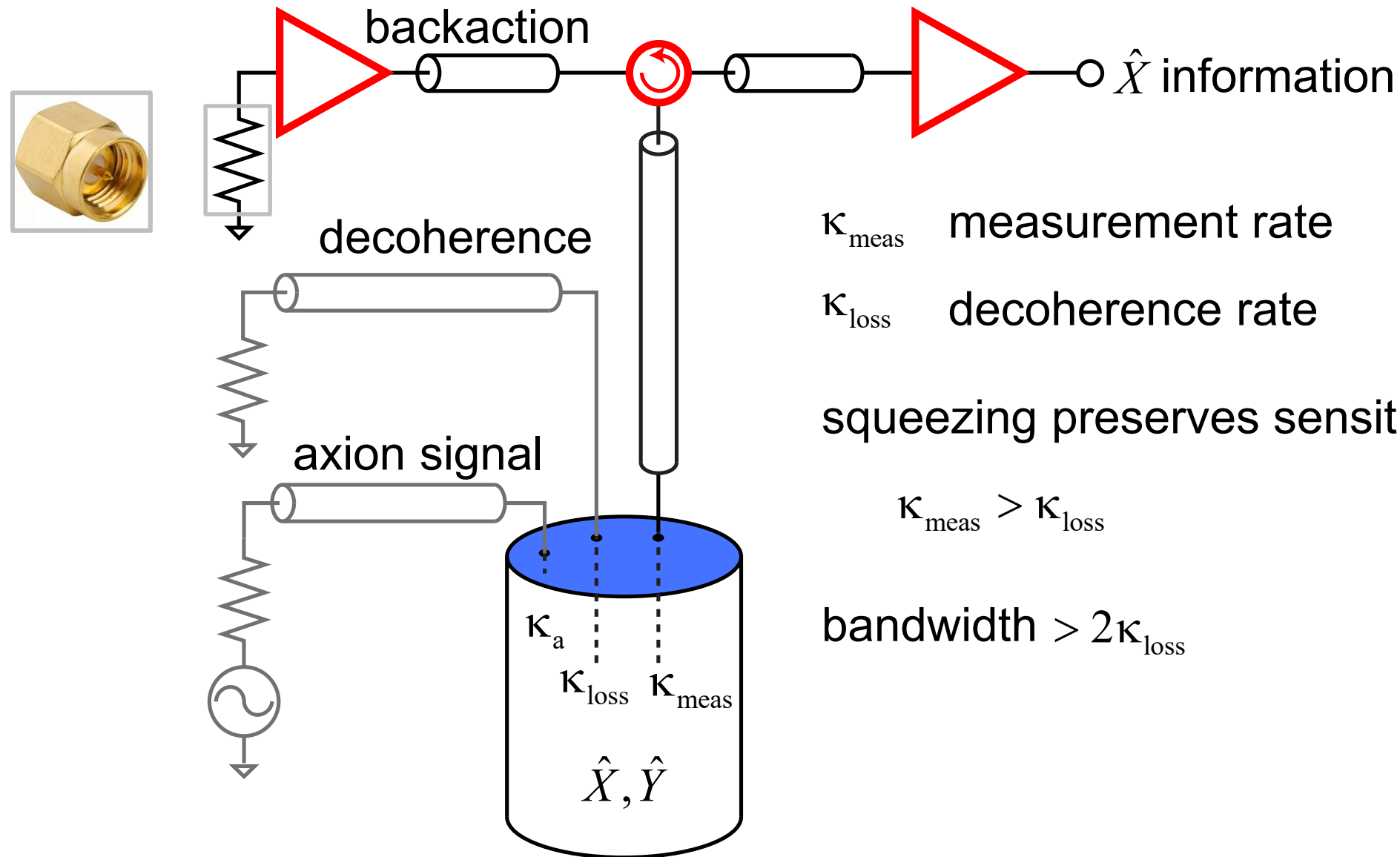


Yale, August 2018. SSR commissioning

Measurement backaction limits haloscope bandwidth



Squeezing yields larger bandwidth through backaction evasion



Questions? (you might ask)

Don't you lose half your signal measuring one quadrature?

(yes, but also half the noise power: [arxiv:1809.06470](#): App. C)

Should you count photons instead?

(maybe, more so at higher frequencies: [arXiv:1607.02529](#))

Is it possible to measure both axion quadratures noiselessly?

(in principle, yes! [arXiv:1607.02529](#))

What would become possible if the cavity were more coherent than the signal?