

# Exploring Multilepton Signatures From Dark Matter at the LHC

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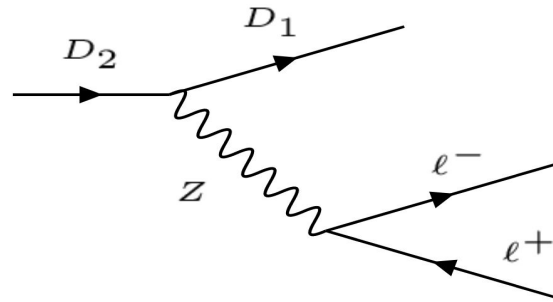
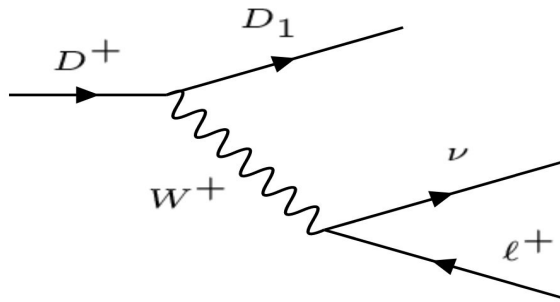
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6th Workshop of the BSM Re-interpretation Forum  
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# Outline

- Motivations for beyond mono-X searches, e.g. multilepton+missing ET
- Two representative examples of Minimal consistent Dark Matter Models (MCDM)
- Multilepton+missing ET is complementary to LLP signature from weak multiplet DM
- We validate our own analysis code against previous analyses
- New limits for I2HDM and MFDM and suggestion for a better parameterisation
- The map of 2&3 lepton channel cross-section limits and efficiencies using CheckMATE for a wider DM re-interpretation analysis for our community



# Inert 2 Higgs Doublet Model (I2HDM)

$$\mathcal{L}_\phi = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - V(\phi_1, \phi_2)$$

$$V = -m_1^2(\phi_1^\dagger \phi_1) - m_2^2(\phi_2^\dagger \phi_2) + \lambda_1(\phi_1^\dagger \phi_1)^2 + \lambda_2(\phi_2^\dagger \phi_2)^2 + \lambda_3(\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) + \lambda_4(\phi_2^\dagger \phi_1)(\phi_1^\dagger \phi_2) + \frac{\lambda_5}{2}[(\phi_1^\dagger \phi_2)^2 + (\phi_2^\dagger \phi_1)^2]$$

$$M_H^2 = 2\lambda_1 v^2 = 2m_1^2 \quad M_{D+}^2 = \frac{1}{2}\lambda_3 v^2 - m_2^2$$

$$M_{D_1}^2 = \frac{1}{2}(\lambda_3 + \lambda_4 - |\lambda_5|)v^2 - m_2^2 \quad M_{D_2}^2 = \frac{1}{2}(\lambda_3 + \lambda_4 + |\lambda_5|)v^2 - m_2^2 > M_{D_1}^2$$

1.  $\lambda_2$  is quartic inert doublet self-coupling    2.  $\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5$  is Higgs-DM coupling:  $HD_1D_1$   
 3.  $M_{D_1}$  is DM mass    4.  $M_{D_2}$  is second lightest, neutral Higgs mass    5.  $M_{D+}$  is charged Higgs mass

Relevant parameters for our study:  $[M_{D_1}, M_{D_2}, M_{D+}, \lambda_2, \lambda_{345}] \longrightarrow [M_{D_1}, M_{D_2}, M_{D+}]$

Parameterisations which are more physical for our analysis:  $\Delta M_{D_2} = M_{D_2} - M_{D+}$   
 $[M_{D_1}, \Delta M_{DP}, \Delta M_{D_2}] \quad \Delta M_{DP} = M_{D+} - M_{D_1}$

# Minimal Fermion Dark Matter (MFDM)

$$\mathcal{L}_{FDM} = \mathcal{L}_{SM} + \bar{\psi}(i\not{D} - m_{\psi})\psi + \frac{1}{2}\bar{\chi}_s^0(i\not{\partial} - m_s)\chi_s^0 - (Y(\bar{\psi}\Phi\chi_s^0) + h.c.)$$

- Minimal model with an EW fermion DM doublet
- To provide the correct amount of relic density, suppress DM scattering through intermediate Z/Higgs boson:

$$\psi = \begin{pmatrix} \chi^+ \\ \frac{1}{\sqrt{2}}(\chi_1^0 + i\chi_2^0) \end{pmatrix}$$

- Majorana neutral D-odd particles  $\chi_1^0, \chi_2^0$
- additional Majorana singlet fermion  $\chi_s^0$

- $\chi_1^0$  and  $\chi_s^0$  mix via Yukawa coupling,  $\chi_2^0$  and  $\chi^+$  are mass degenerate  $Y_{DM} = \frac{\sqrt{(M_{D3} - M_{DP})(M_{DP} - M_{D1})}}{v}$

**1.**  $M_{D1}$  is DM mass **2.**  $M_{DP} = M_{D2}$  is chargino mass **3.**  $M_{D3}$  is third lightest, neutralino mass

$$M_{D3} > M_{DP} = M_{D2} > M_{D1}$$

Parameterisations which are more physical for our analysis:

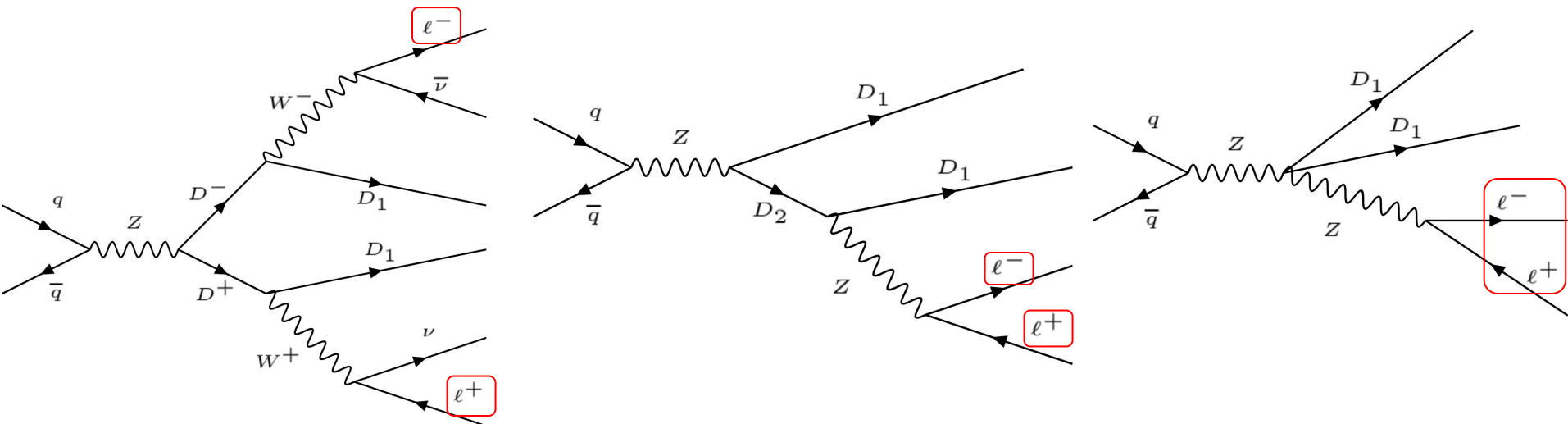
$$[M_{D1}, \Delta M_{DP}, \Delta M_{D3}]$$

$$\Delta M_{D3} = M_{D3} - M_{DP}$$

$$\Delta M_{DP} = M_{DP} - M_{D1}$$

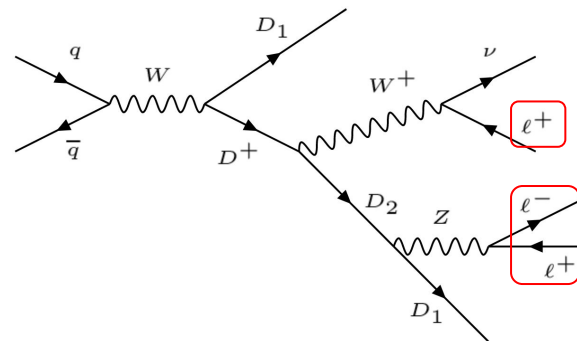
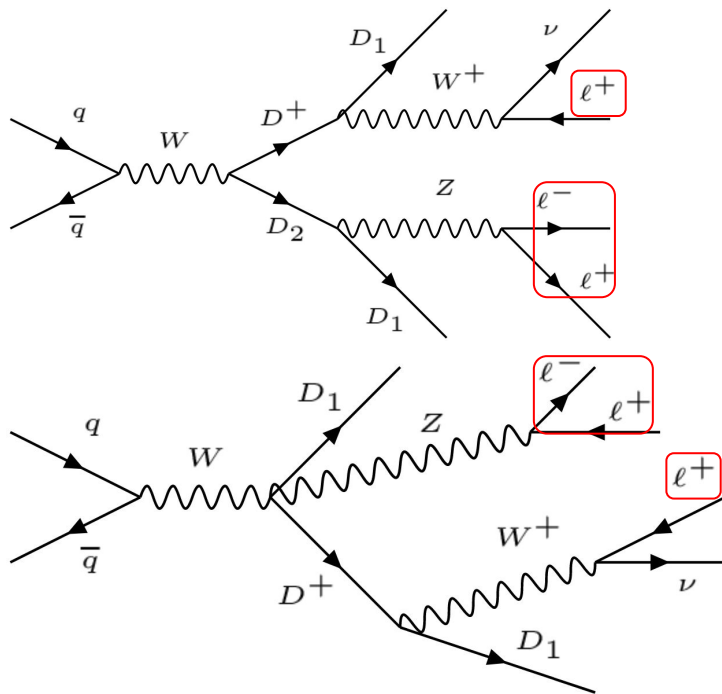
$$Y_{DM} = \frac{\sqrt{\Delta M_{D3}\Delta M_{DP}}}{v}$$

# $2\ell + E_T^{miss}$ I2HDM Final States



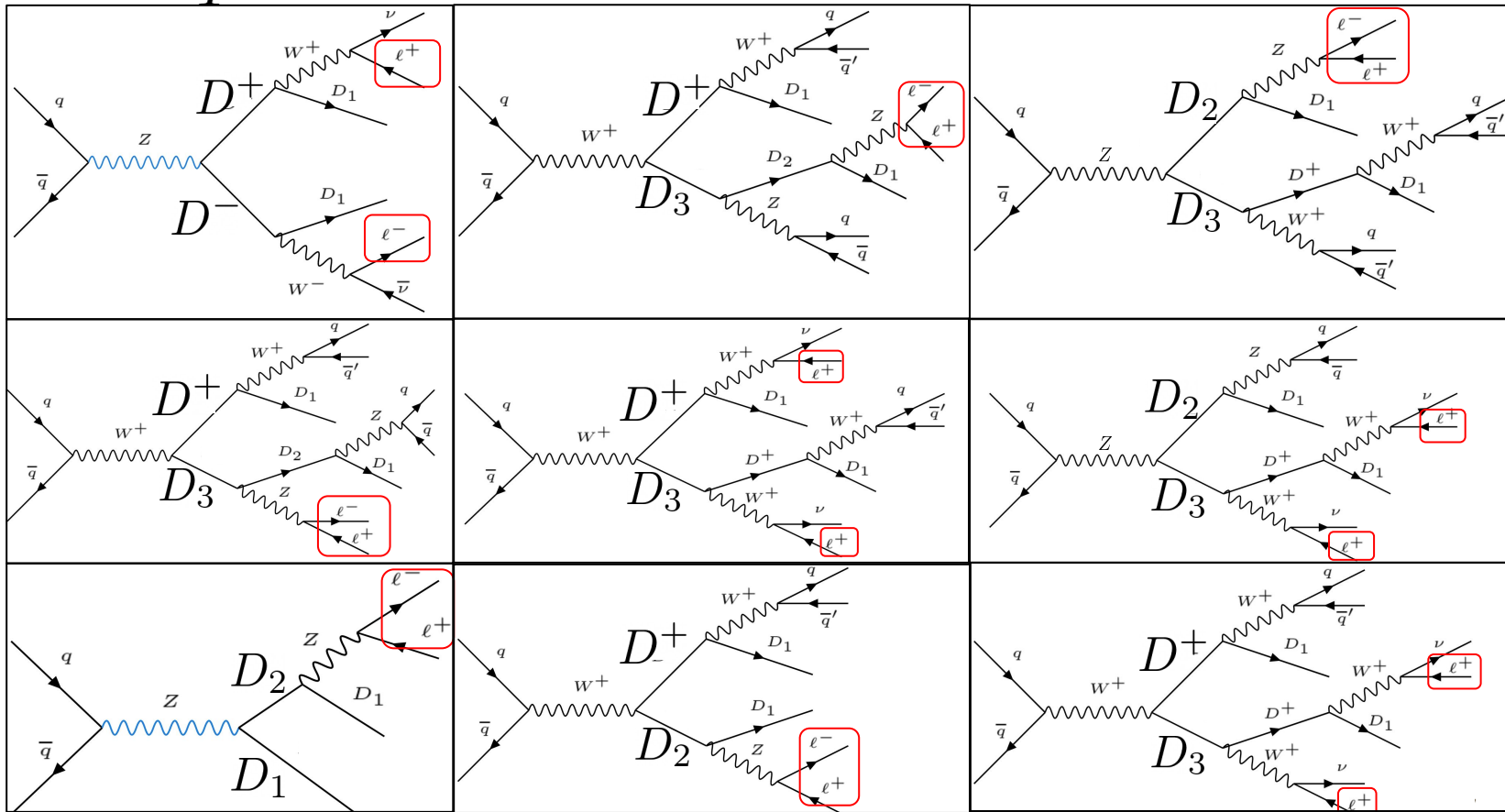
- DM decays via  $Z$  production
- Looking at Higgs funnel:  $\lambda_{345} \sim 0$ , and  $\lambda_2$  not relevant

# $3\ell + E_T^{miss}$ I2HDM Final States

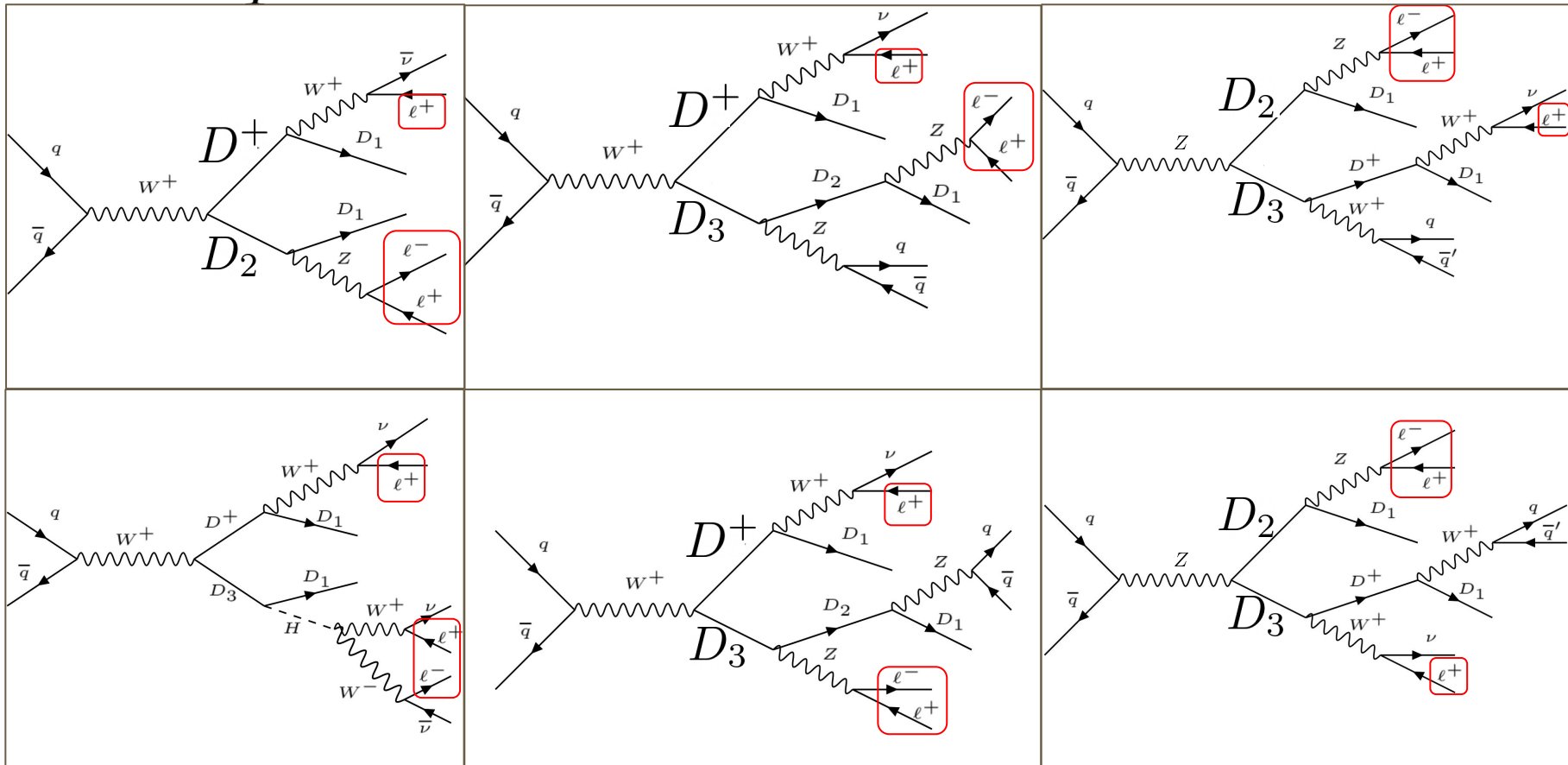


- DM decays via  $W$  production, x2 for the  $+/-$  processes

# $2\ell + E_T^{miss}$ MFDM Final States



# $3\ell + E_T^{miss}$ MFDM Final States





# HEP Tools

- CalcHEP: Parton-level event production and decays: LHE file output
- CheckMATE (+ Pythia + Delphes): Decays, parton-showers, detector effects and analysis checks
- 8 TeV: written new analysis for final states with  $2\ell$  and  $E_T^{miss}$
- 13 TeV: Check any available ATLAS and CMS analyses, lists  $2\ell$  and  $3\ell$  channels

$$\mu = \frac{\sigma_{DM}}{\sigma_{95}} \quad \begin{array}{l} \text{Cross-section of DM events produced} \\ \text{Cross-section required to exclude point at 95\% confidence level} \end{array}$$

- Point excluded if  $\mu \geq 1$

# 8 TeV Analysis Cuts

- 8 TeV ATLAS SUSY analysis [arXiv:1403.5294](https://arxiv.org/abs/1403.5294)

cutflows for dilepton+MET finals states,  
implemented in CheckMATE

Global Cut	
$E_T^{miss}$	$> 0$ GeV
Base leptons	2
$e + e^-$ trigger	97%
$\mu^+ \mu^-$ trigger	89%
$e\mu$ trigger	75%
Signal leptons	2
Leading lepton $p_T$	$> 35$ GeV
sub-leading lepton $p_T$	$> 20$ GeV
$M_{\ell\ell}$	$> 20$ GeV
jets	0
$ M_{\ell\ell} - M_Z $	$> 10$ GeV

SR	$m_{T2}^{90}$	$m_{T2}^{120}$	$m_{T2}^{150}$	WWa	WWb	WWc	Zjets
$M_{\ell\ell}$				$< 120$	$< 170$		
$p_T(\ell\ell)$				$> 80$			$> 80$
$E_T^{miss,rel}$				$> 80$			$> 80$
$m_{T2}$	$> 90$	$> 120$	$> 150$		$> 90$	$> 100$	

best for  
these results

- 8 TeV ATLAS Higgs analysis [arXiv:1402.3244](https://arxiv.org/abs/1402.3244) cutflows for dilepton+MET

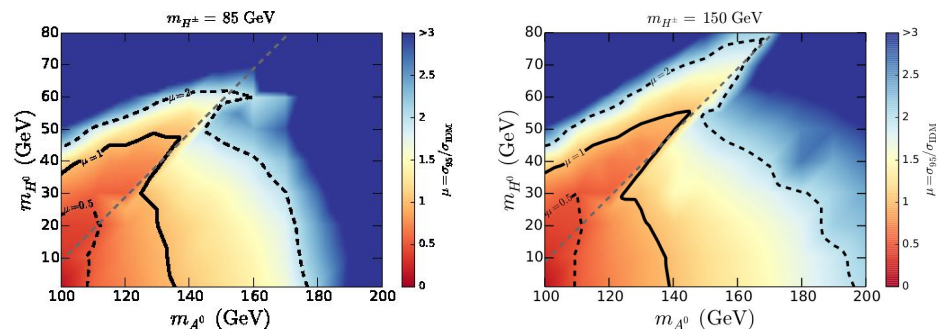
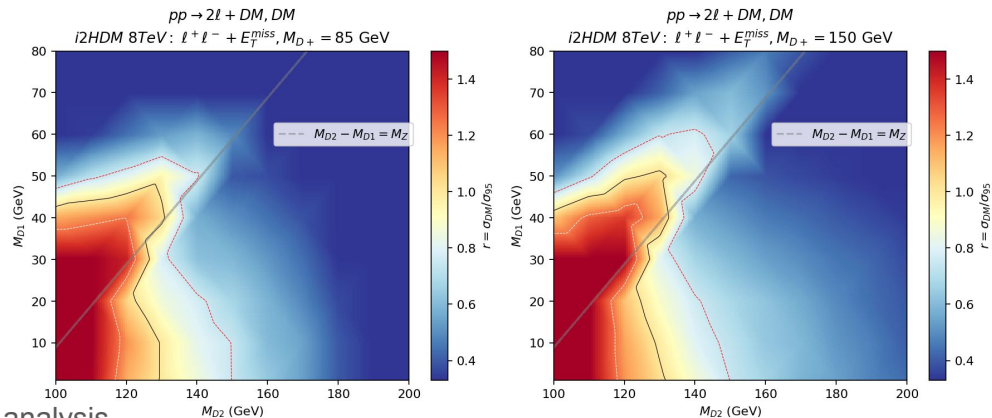
finals states, implemented in CheckMATE

Global Cut	
Base leptons	2
Lepton $p_T$	$> 20$ GeV
Z-window	$76 < M_{\ell\ell} < 106$ GeV
$E_T^{miss}$	$> 90$ GeV
$d\phi(E_T^{miss}, p_T^{miss})$	$< 0.2$
$\Delta\phi(p_T(\ell\ell), E_T^{miss})$	$> 2.6$
$\Delta\phi(\ell, \ell)$	$< 1.7$
$ \frac{E_T^{miss} - p_T(\ell\ell)}{p_T(\ell\ell)} $	$> 0.2$
jets	0

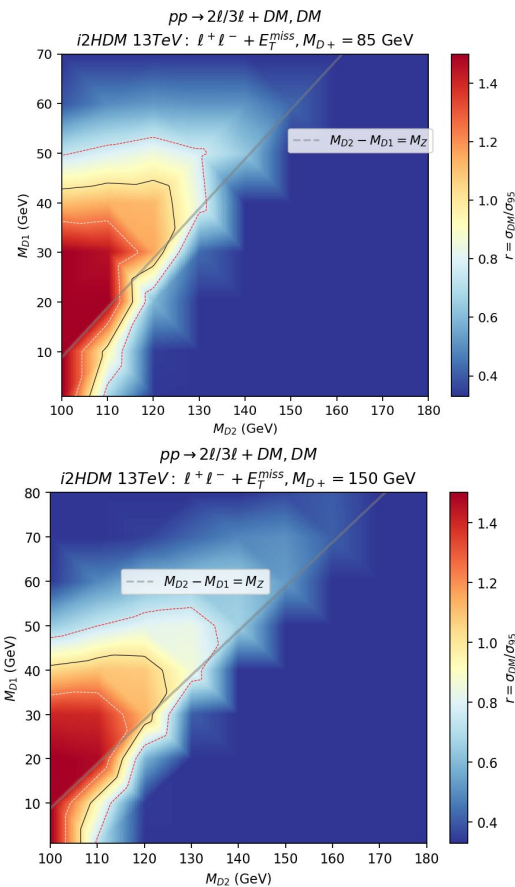
- Validated against MadAnalysis (Belanger et.al paper [arXiv:1503.07367](https://arxiv.org/abs/1503.07367))

# I2HDM Validations

8 TeV



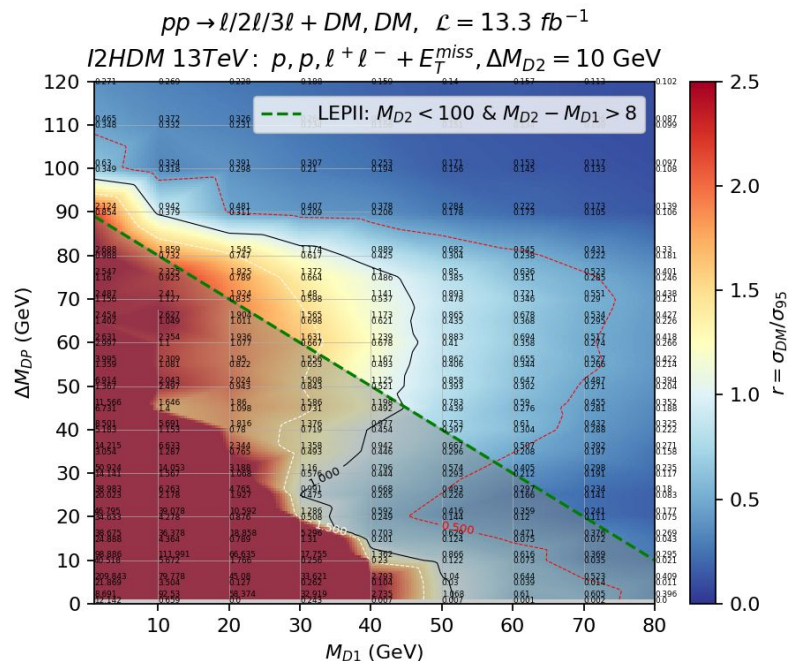
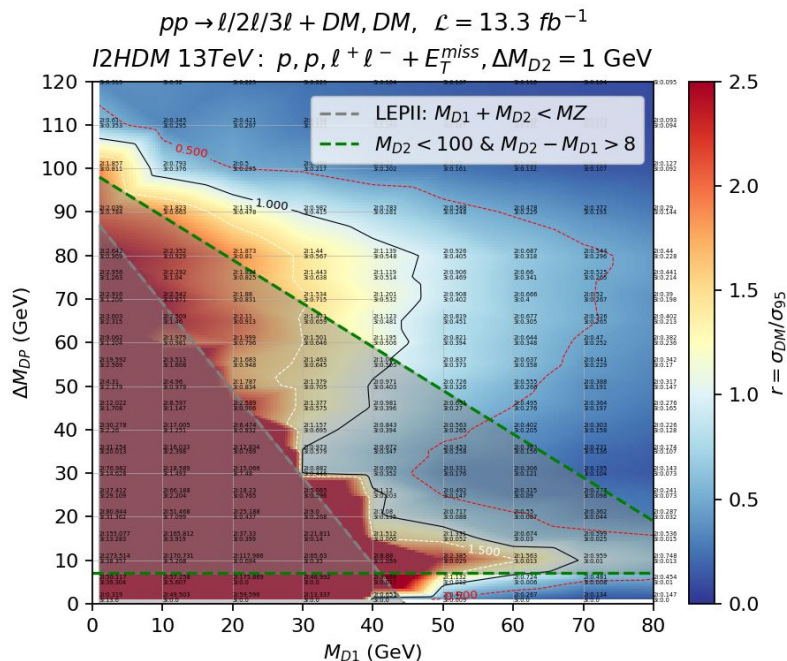
13 TeV



# New I2HDM Results

$$\Delta M_{D2} = M_{D2} - M_{D+}$$

$$\Delta M_{DP} = M_{D+} - M_{D1}$$

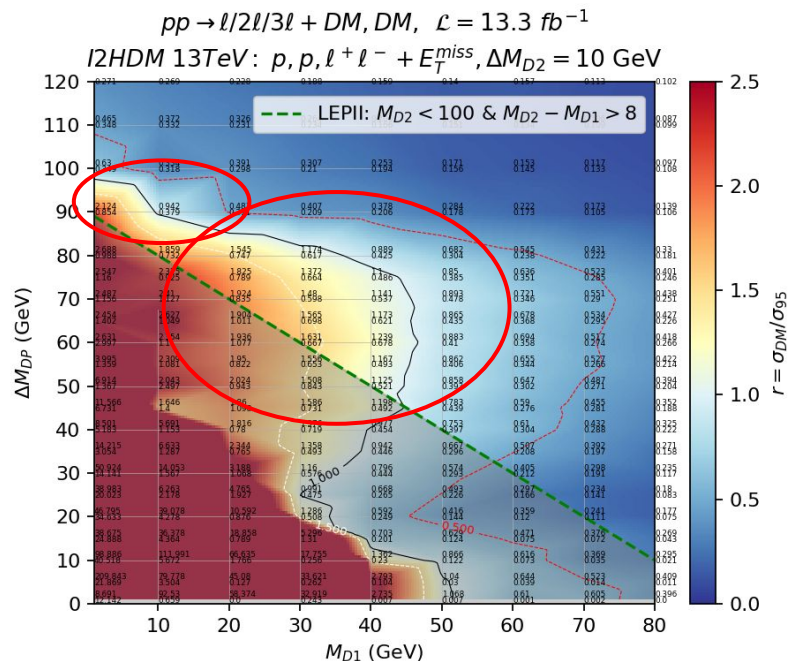
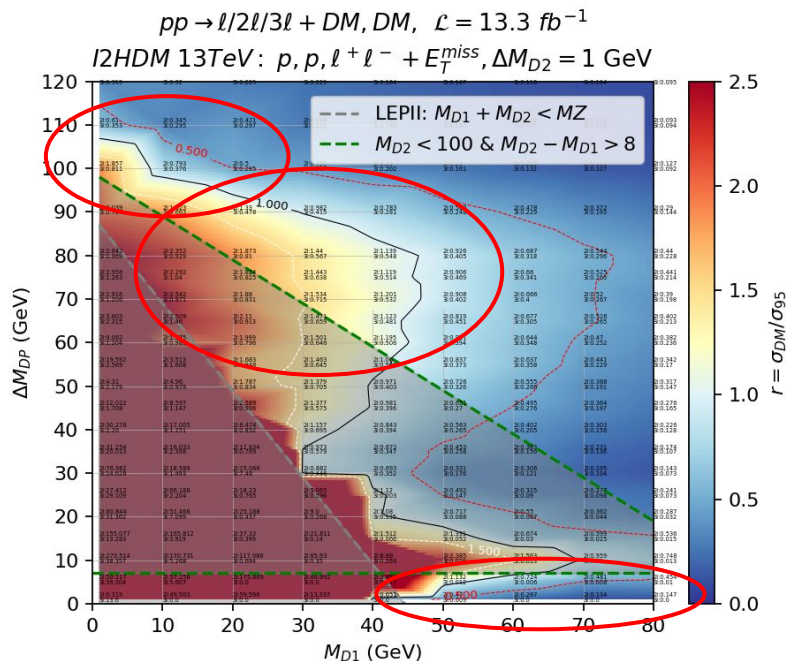


- $\Delta M_{D2} = 1 \text{ GeV}$ : Small wedge above  $M_{D1} > 50 \text{ GeV}$  and below  $\Delta M_{DP} < 8 \text{ GeV}$  still allowed by LEP
- $M_{DP}$  is a better variable, results not dependent on  $M_{D2}$ , only require plane of 2 variables
- Important contributions from 3-lepton (up to 70%) which could be combined with 2-lepton

# New I2HDM Results

$$\Delta M_{D2} = M_{D2} - M_{D+}$$

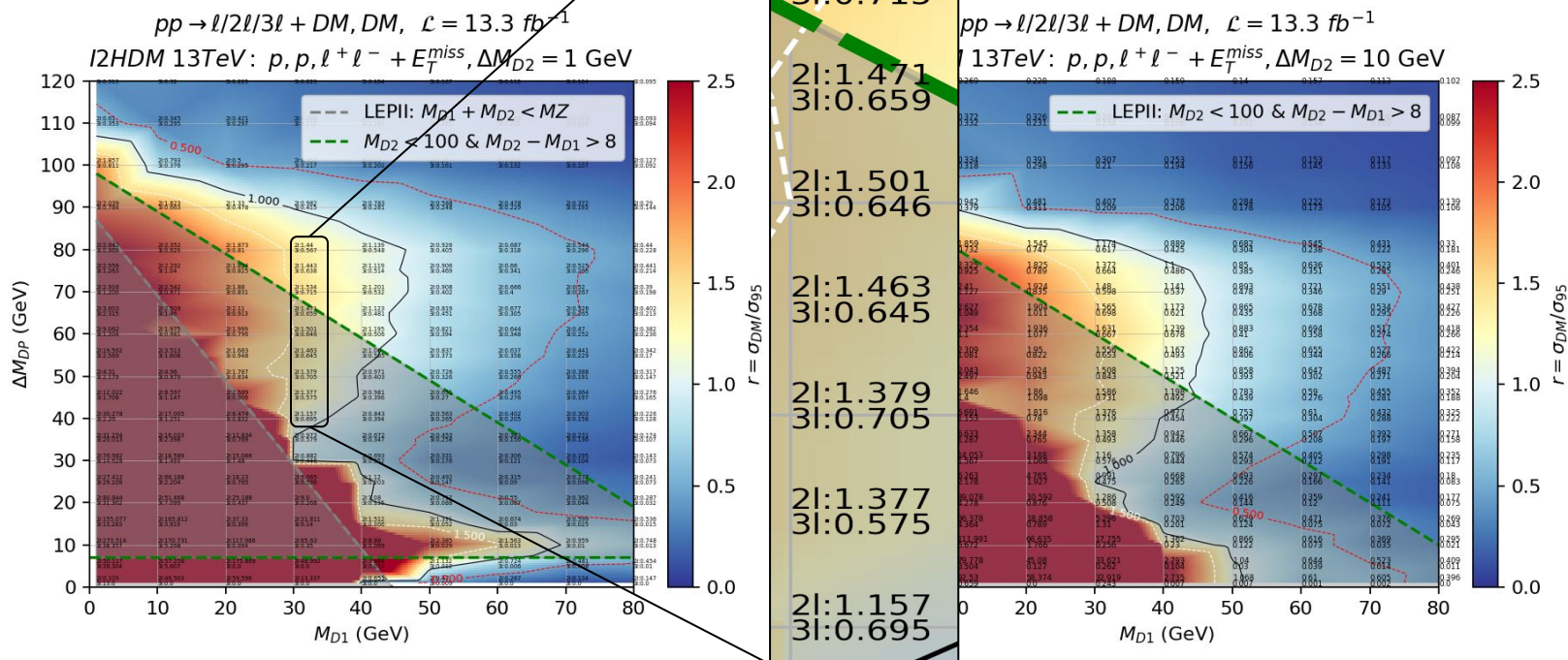
$$\Delta M_{DP} = M_{D+} - M_{D1}$$



- $\Delta M_{D2} = 1 \text{ GeV}$ : Small wedge above  $M_{D1} > 50 \text{ GeV}$  and below  $\Delta M_{DP} < 8 \text{ GeV}$  still allowed by LEP
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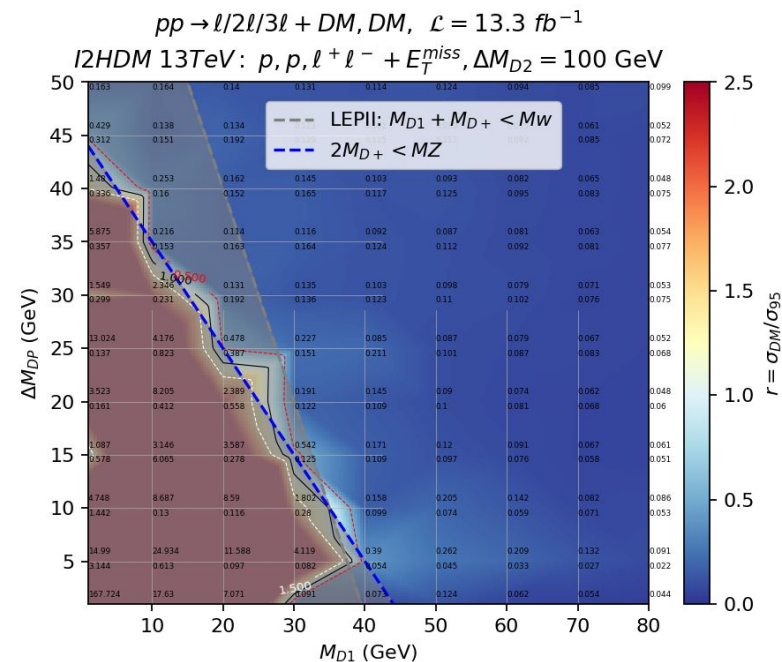
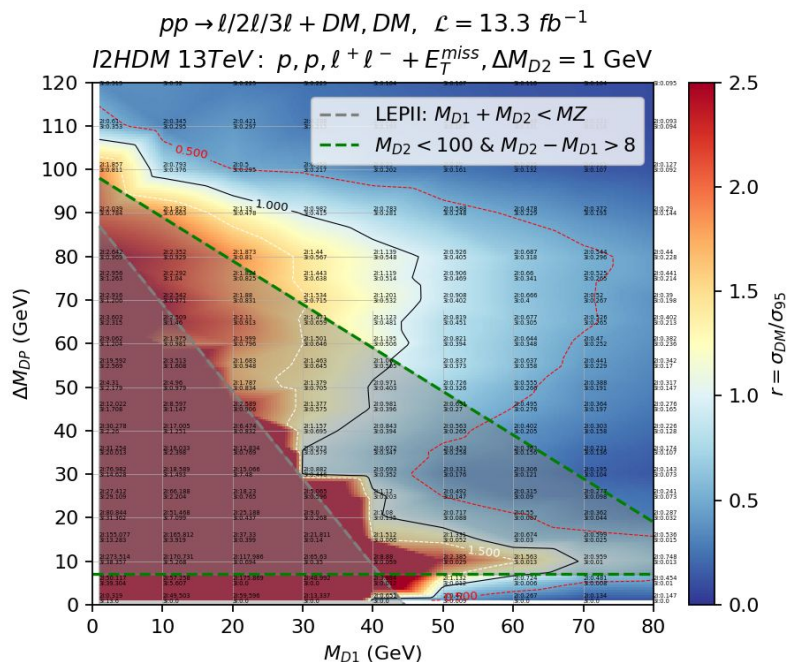


# New I2HDM Results



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- $M_{DP}$  is a better variable, results not dependent on  $M_{D2}$ , only require plane of 2 variables
- Important contributions from 3-lepton (up to 70%) which could be combined with 2-lepton

# New I2HDM Results



- Increasing  $\Delta M_{D2}$  to 200 GeV means the Z veto  $M_{\ell\ell} > 100 \text{ GeV}$  requirement can no longer be fulfilled as production cross-section of the heavier states has fallen

# Re-interpretation: Providing Cross-section Limits

## I2HDM

	Sample A	Sample B	Sample C
No# Events:	50,000	150,000	100,000
Production:	$pp \rightarrow D^+ D^-$ $pp \rightarrow D_2 D_1$	$pp \rightarrow D^\pm D_2$	$pp \rightarrow Z D_1 D_1$
Decays:	$D^\pm \rightarrow (W^\pm \rightarrow \ell^\pm \nu) D_1$ $D_2 \rightarrow (Z \rightarrow \ell^+ \ell^-) D_1$	$D_2 \rightarrow (Z \rightarrow \ell^+ \ell^-) D_1$	$Z \rightarrow \ell^+ \ell^-$

- While the genuine 2-2 process  $pp \rightarrow D_2 D_1$  is separate to 3-body decay  $pp \rightarrow Z D_1 D_1$ , width of  $D_2$  is small, so expected interference between these diagrams is small



# Re-interpretation: Providing Cross-section Limits

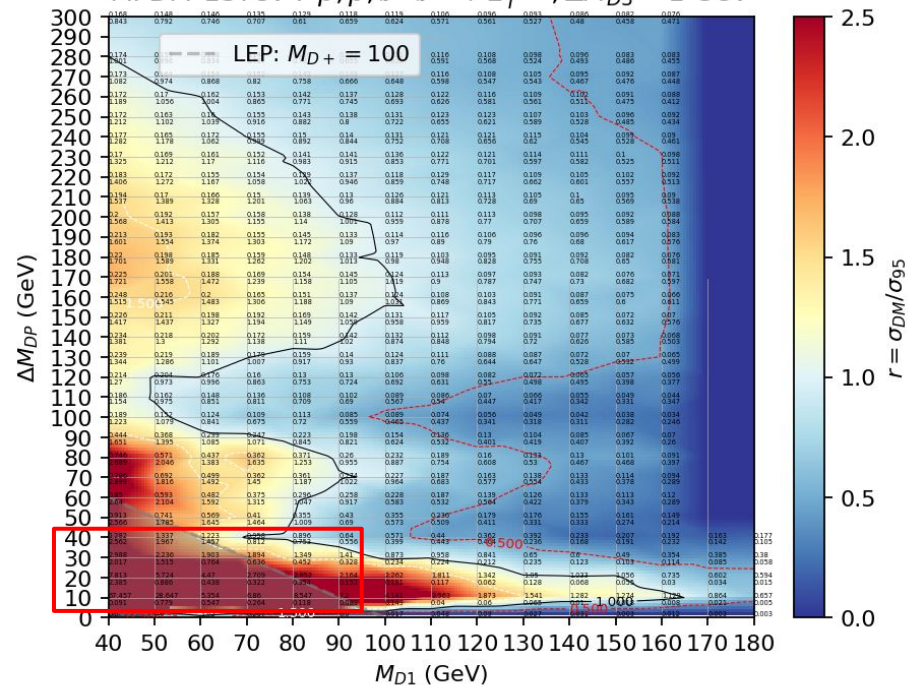
## MFDM

	Sample A	Sample B	Sample C
No# Events:	50,000	150,000	100,000
Production:	$pp \rightarrow D^+ D^-$ $pp \rightarrow D_2 D_1$	$pp \rightarrow D_2 D_3$	$pp \rightarrow D^\pm D_2$ $pp \rightarrow D^\pm D_3$
Decays:	$D^\pm \rightarrow (W^\pm \rightarrow \ell^\pm \nu) D_1$ $D_2 \rightarrow (Z \rightarrow \ell^+ \ell^-) D_1$	Any	$D_2 \rightarrow (Z \rightarrow \ell^+ \ell^-) D_1$ $D_3 \rightarrow (W^\pm \rightarrow \ell^\pm \nu) D^\pm$ $D_3 \rightarrow (Z \rightarrow \ell^+ \ell^-) D_2$

# MFDM Results

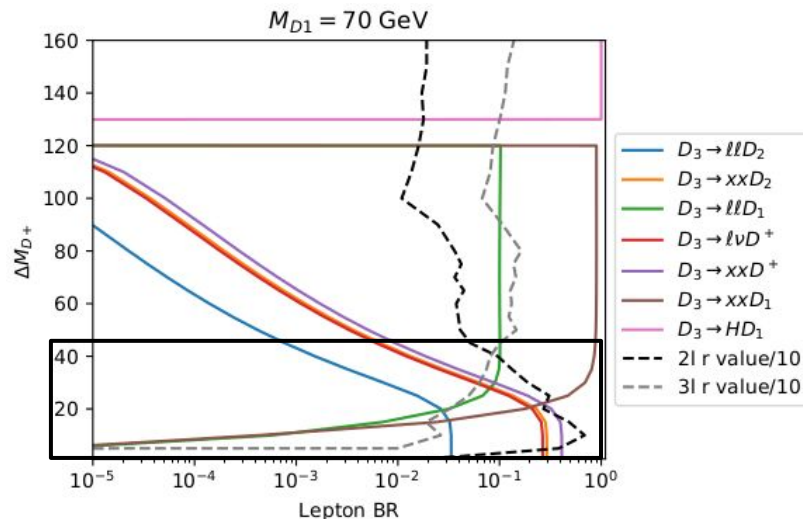
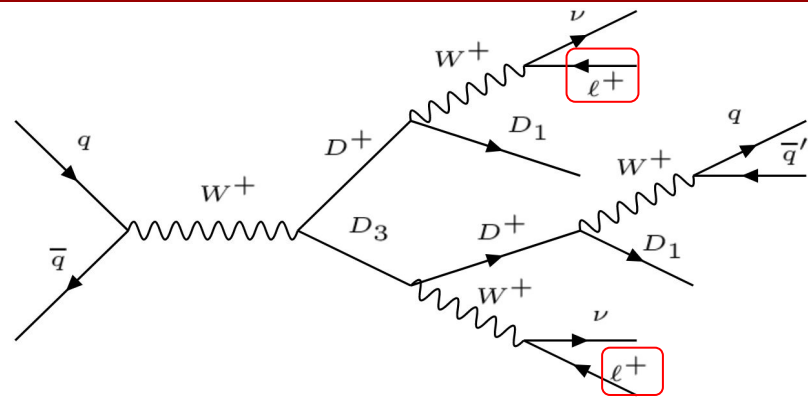
$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$

MFDM 13TeV:  $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 1 \text{ GeV}$



- Similar shapes to I2HDM, but with 3-lepton channel sensitivity due to crossing between

$$D_3 \rightarrow \ell\nu D_1 \quad D_3 \rightarrow Z(\rightarrow \ell\ell)D_1$$

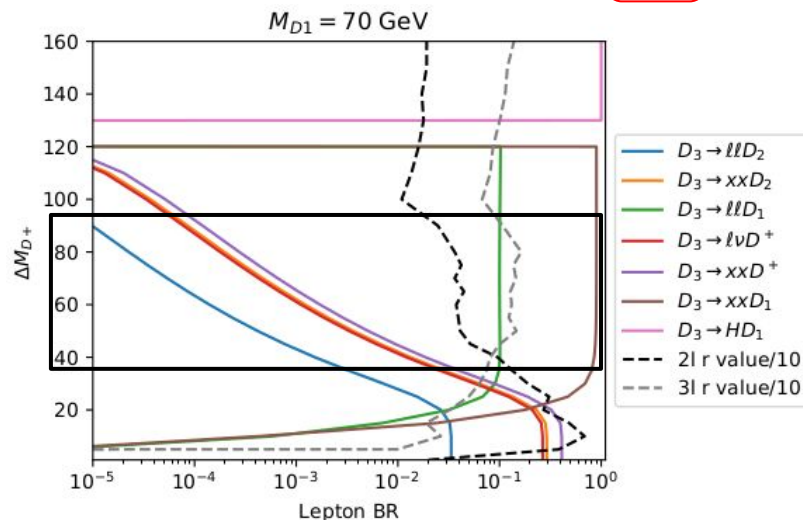
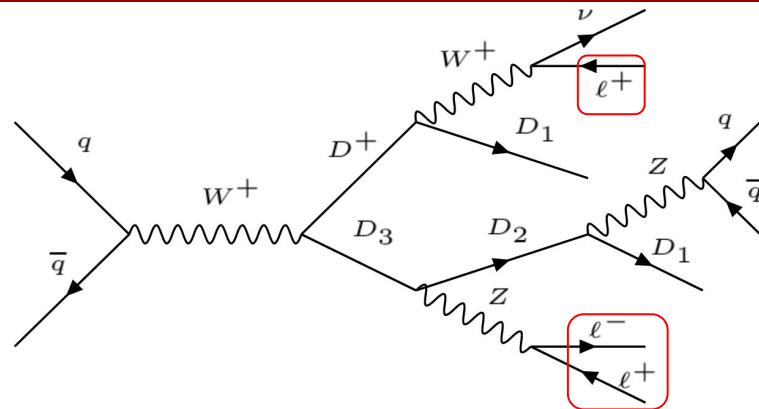
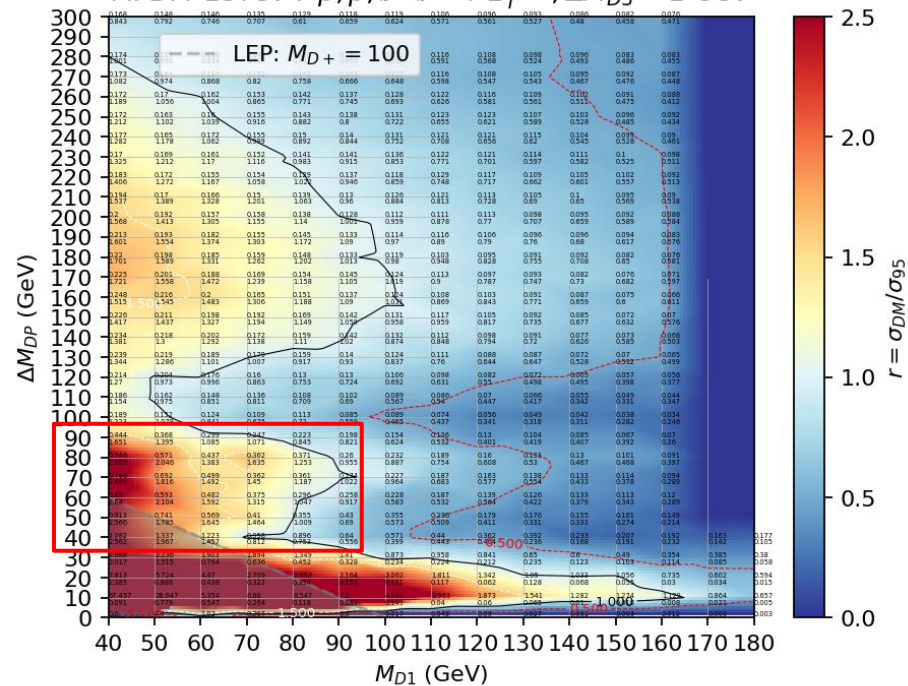


$$D^\pm(\rightarrow \ell\nu D_1)D_3 \quad \Delta M_{DP} = 45$$

# MFDM Results

$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$

MFDM 13TeV:  $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 1 \text{ GeV}$



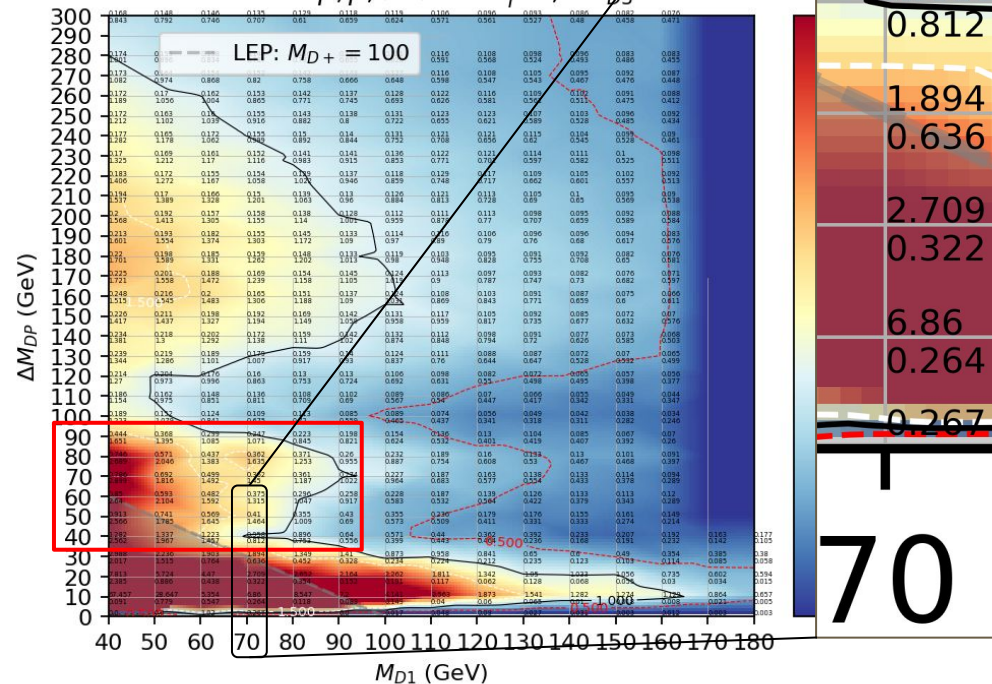
- Similar shapes to I2HDM, but 3-lepton channel sensitivity begins to dominate due to crossing between  $D_3 \rightarrow \ell\nu D_1$  and  $D_3 \rightarrow Z(\rightarrow \ell\ell)D_1$

$$D^\pm(\rightarrow \ell\nu D_1)D_3 \quad \Delta M_{DP} = 45$$

# MFDM Results

$$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$$

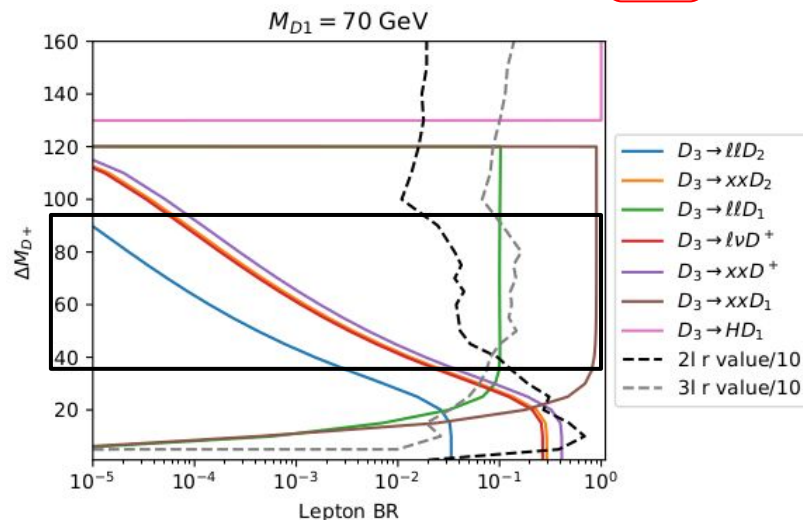
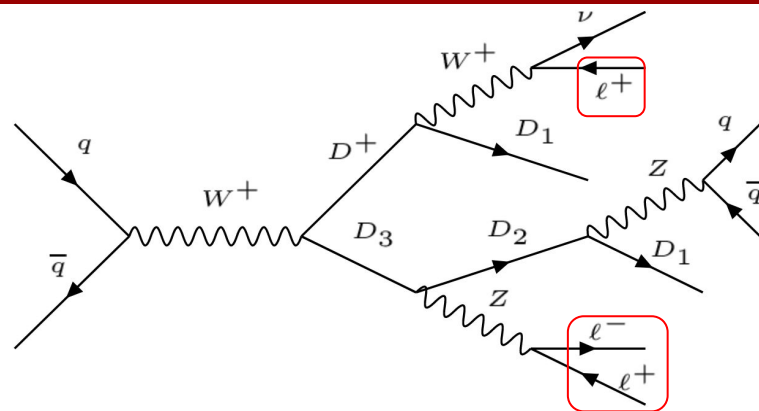
MFDM 13TeV:  $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{DP} = 1 \text{ GeV}$



- Similar shapes to I2HDM, but 3-lepton channel sensitivity begins to dominate due to crossing between

$$D_3 \rightarrow \ell \nu D_1 \quad D_3 \rightarrow Z(\rightarrow \ell \ell) D_1$$

$$D^\pm(\rightarrow \ell \nu D_1) D_3 \quad \Delta M_{DP} = 45$$

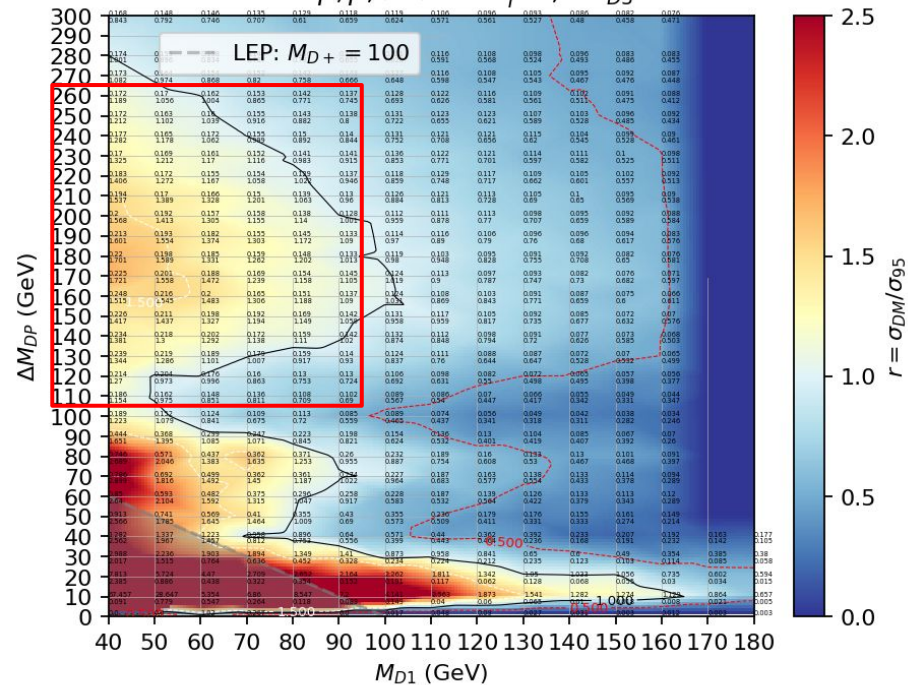




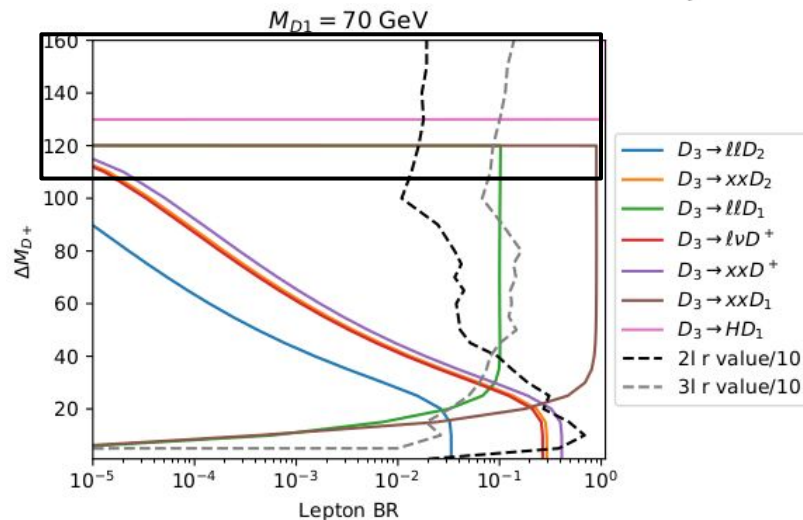
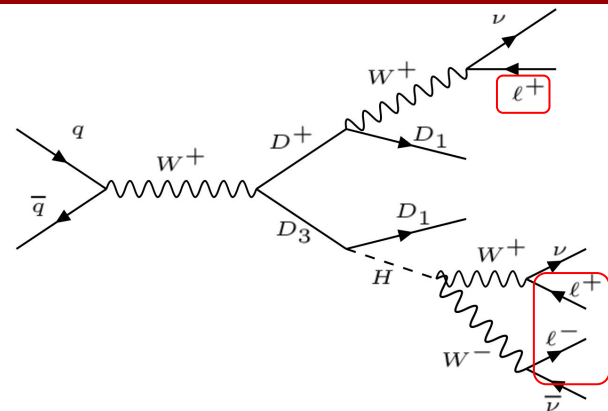
# MFDM Results

$$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$$

MFDM 13TeV:  $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 1 \text{ GeV}$



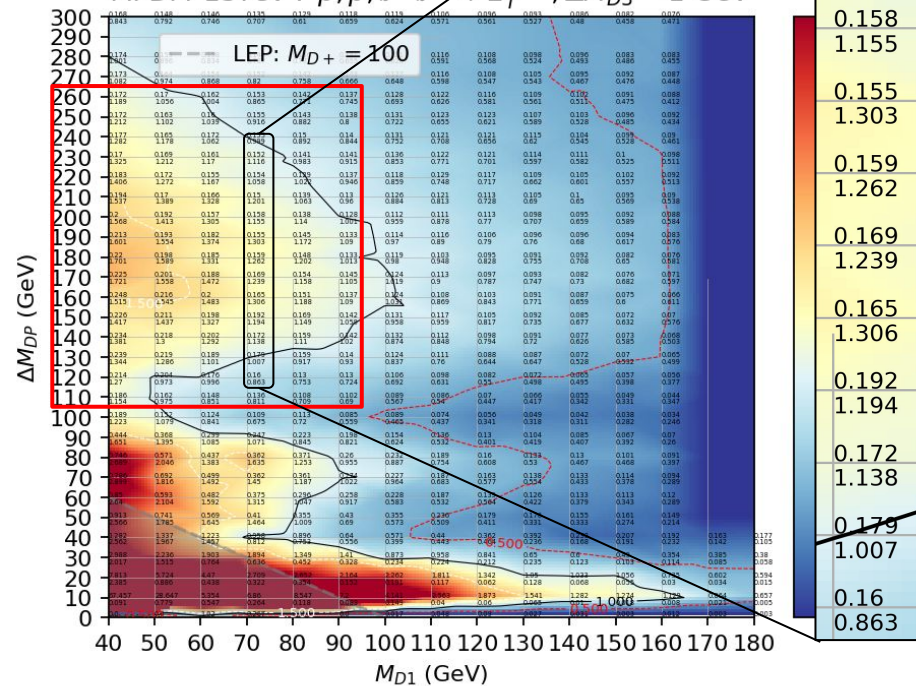
- Second shape due to 3-lepton channel sensitivity due to Higgs decay  $D_3 \rightarrow Z(\rightarrow \ell\ell)D_1$  to  $D_3 \rightarrow H(\rightarrow W^+W^-)D_1$  with production of  $D^\pm(\rightarrow \ell\nu D_1)D_3$ , at  $\Delta M_{DP} = 125 \text{ GeV}$



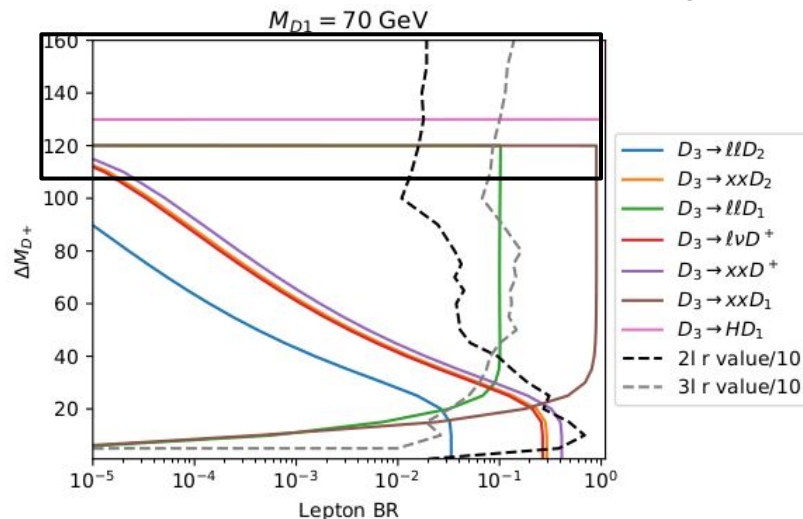
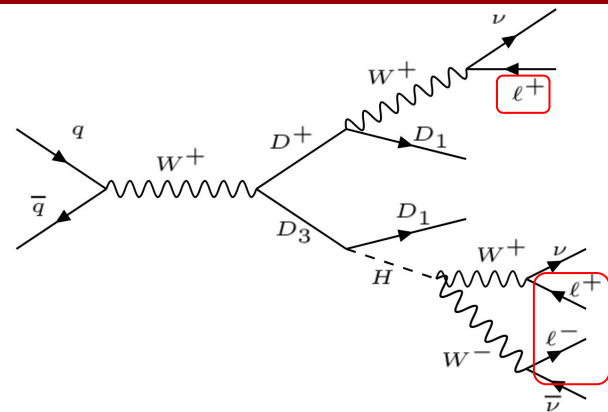
# MFDM Results

$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mu = 13.3 \text{ fb}^{-1}$

MFDM 13TeV:  $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 1 \text{ GeV}$



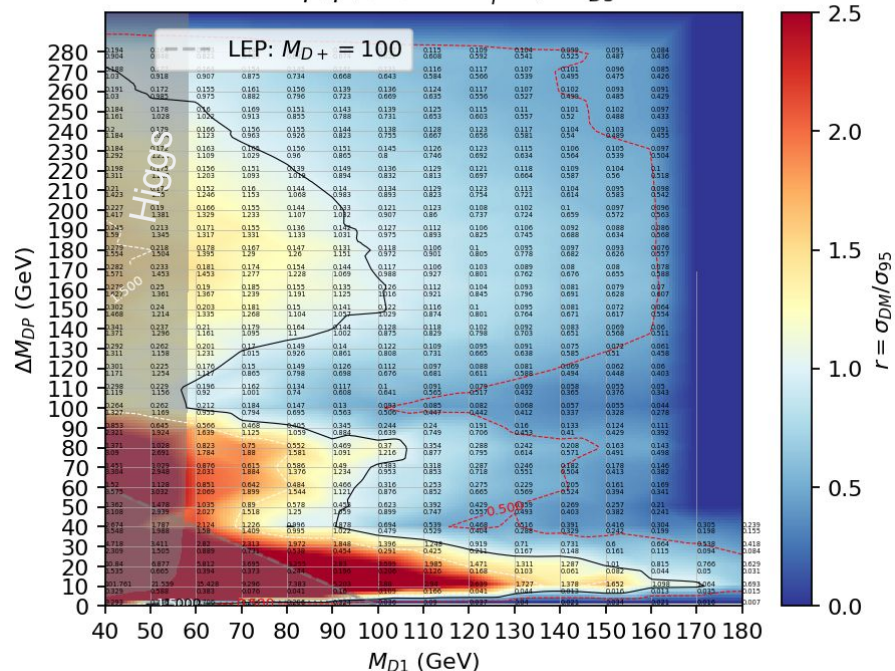
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# MFDM Results

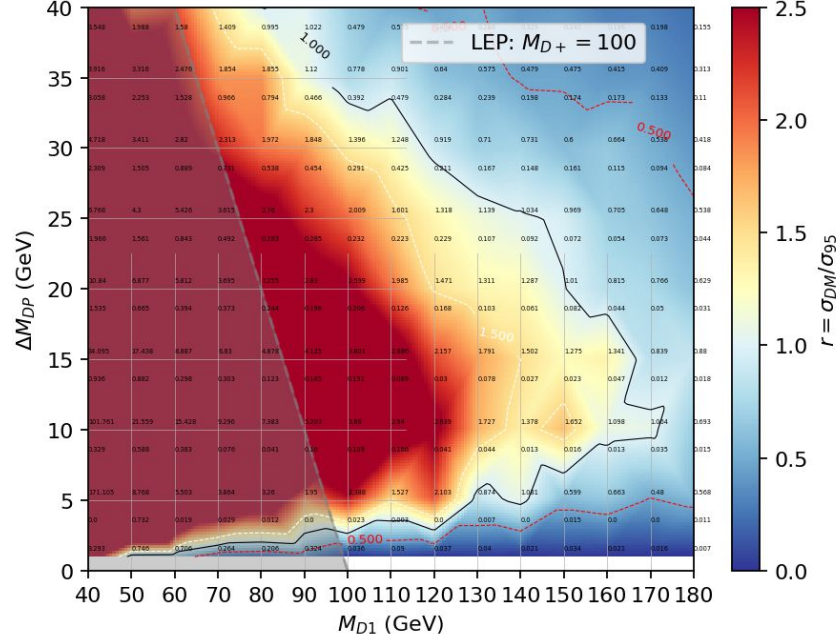
$$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$$

MFDM 13TeV:  $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 10 \text{ GeV}$



$$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$$

MFDM 13TeV:  $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 10 \text{ GeV}$



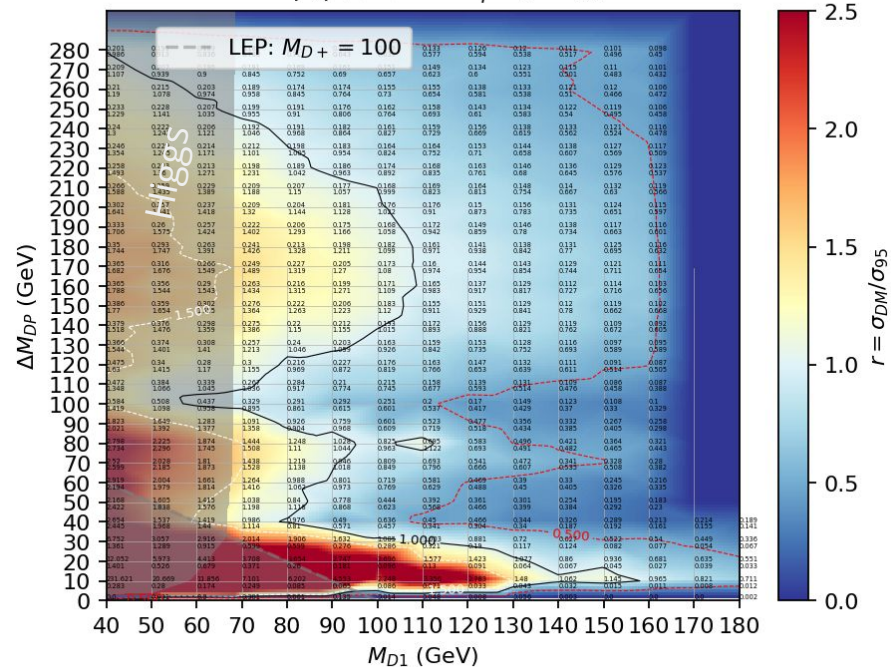
- As  $\Delta M_{D3}$  increases, coupling between  $D_1 - D^\pm$  increases, while heavy  $D_3$  leads to suppressed production cross-section - 'no-lose' theorem



# MFDM Results

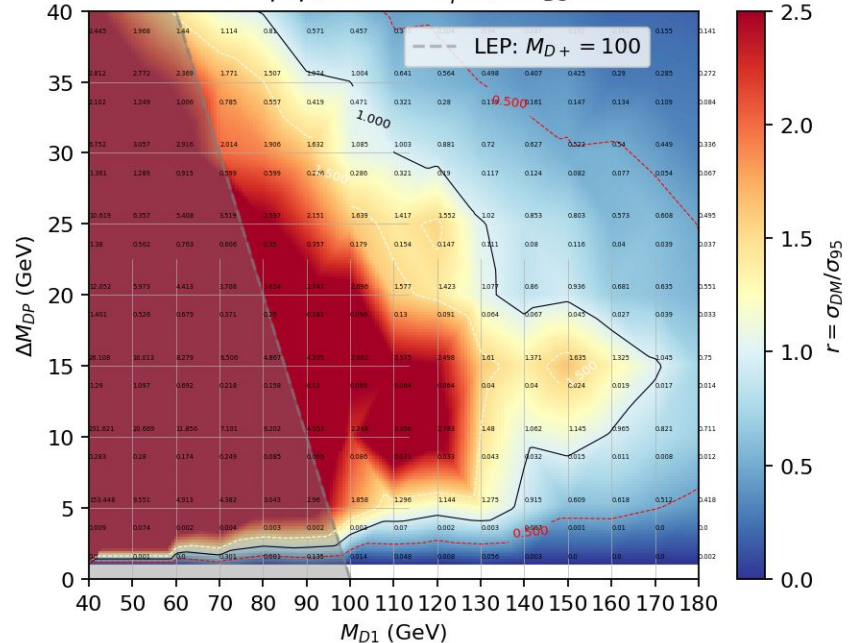
$$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$$

$$\text{MFDM 13TeV: } p, p, \ell + \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 100 \text{ GeV}$$



$$pp \rightarrow \ell/2\ell/3\ell + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$$

$$\text{MFDM 13TeV: } p, p, \ell + \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 100 \text{ GeV}$$



- With increasing  $\Delta M_{D3}$ , Higgs to invisible limit covers larger  $M_{D1}$



# Re-interpretation: Providing Cross-section Limits


Mass parameter points



$M_{D1}$	$\Delta M_{D+}$	$\Delta M_{D2}$	$2\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$
1	5	1	$3.26 \times 10^3$	71	-	100	$6.51 \times 10^4$	71	-	-	$1.21 \times 10^3$	24	-	-
1	10	1	97.0	41	-	100	-	-	-	-	$1.21 \times 10^3$	24	-	100
1	20	1	$1.47 \times 10^3$	58	$6.63 \times 10^3$	71	-	100	-	-	933	21	-	-
1	40	1	$1.02 \times 10^5$	35	$8.17 \times 10^4$	58	$8.17 \times 10^4$	71	-	-	$1.2 \times 10^3$	8	-	-
1	60	1	$8.84 \times 10^3$	45	$5.3 \times 10^3$	20	$2.94 \times 10^4$	58	-	-	220	6	-	100
1	80	1	783	11	326	4	$1.15 \times 10^3$	9	-	-	93.0	6	-	-
10	5	1	698	58	$3.14 \times 10^3$	71	-	100	-	-	-	-	-	-
10	10	1	161	38	674	45	-	-	-	-	-	-	-	-
10	20	1	287	45	-	100	$1.43 \times 10^4$	71	-	-	$1.87 \times 10^3$	30	-	100
10	40	1	$1.40 \times 10^4$	50	$1.29 \times 10^4$	28	$2.23 \times 10^4$	45	-	-	531	5	$6.82 \times 10^4$	71
10	60	1	$4.44 \times 10^3$	26	507	5	604	7	-	-	165	5	-	-
10	80	1	150	5	248	4	630	7	-	-	80.0	5	-	-
10	120	1	281	6	$1.32 \times 10^3$	8	411	6	-	-	62.0	4	-	-
20	5	1	97.0	41	877	71	-	-	-	-	-	-	-	-
20	10	1	140	35	562	41	-	-	-	-	-	-	-	-
20	20	1	$4.78 \times 10^3$	58	$1.08 \times 10^4$	50	-	-	-	-	$9.32 \times 10^3$	21	-	-
20	40	1	$6.31 \times 10^3$	38	$6.02 \times 10^3$	21	$1.76 \times 10^4$	45	-	-	366	7	-	-
20	60	1	247	6	377	4	438	6	-	-	148	5	-	-
20	80	1	91.0	4	230	3	534	6	-	-	62.0	5	-	-
20	120	1	247	6	$1.50 \times 10^3$	9	321	5	-	100	58.0	4	$9.40 \times 10^3$	58

# Re-interpretation: Providing Cross-section Limits


Cross-section limit (95% cl) for 2  
lepton channel of sample A,B,C



$M_{D1}$	$\Delta M_{D+}$	$\Delta M_{D2}$	$2\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$
1	5	1	$3.26 \times 10^3$	71	-	100	$6.51 \times 10^4$	71	-	-	$1.21 \times 10^3$	24	-	-
1	10	1	97.0	41	-	100	-	-	-	-	$1.21 \times 10^3$	24	-	100
1	20	1	$1.47 \times 10^3$	58	$6.63 \times 10^3$	71	-	100	-	-	933	21	-	-
1	40	1	$1.02 \times 10^5$	35	$8.17 \times 10^4$	58	$8.17 \times 10^4$	71	-	-	$1.2 \times 10^3$	8	-	-
1	60	1	$8.84 \times 10^3$	45	$5.3 \times 10^3$	20	$2.94 \times 10^4$	58	-	-	220	6	-	100
1	80	1	783	11	326	4	$1.15 \times 10^3$	9	-	-	93.0	6	-	-
10	5	1	698	58	$3.14 \times 10^3$	71	-	100	-	-	-	-	-	-
10	10	1	161	38	674	45	-	-	-	-	-	-	-	-
10	20	1	287	45	-	100	$1.43 \times 10^4$	71	-	-	$1.87 \times 10^3$	30	-	100
10	40	1	$1.40 \times 10^4$	50	$1.29 \times 10^4$	28	$2.23 \times 10^4$	45	-	-	531	5	$6.82 \times 10^4$	71
10	60	1	$4.44 \times 10^3$	26	507	5	604	7	-	-	165	5	-	-
10	80	1	150	5	248	4	630	7	-	-	80.0	5	-	-
10	120	1	281	6	$1.32 \times 10^3$	8	411	6	-	-	62.0	4	-	-
20	5	1	97.0	41	877	71	-	-	-	-	-	-	-	-
20	10	1	140	35	562	41	-	-	-	-	-	-	-	-
20	20	1	$4.78 \times 10^3$	58	$1.08 \times 10^4$	50	-	-	-	-	$9.32 \times 10^3$	21	-	-
20	40	1	$6.31 \times 10^3$	38	$6.02 \times 10^3$	21	$1.76 \times 10^4$	45	-	-	366	7	-	-
20	60	1	247	6	377	4	438	6	-	-	148	5	-	-
20	80	1	91.0	4	230	3	534	6	-	-	62.0	5	-	-
20	120	1	247	6	$1.50 \times 10^3$	9	321	5	-	100	58.0	4	$9.40 \times 10^3$	58

# Re-interpretation: Providing Cross-section Limits

Cross-section limit (95% cl) for 3  
lepton channel of sample A,B,C




$M_{D1}$	$\Delta M_{D+}$	$\Delta M_{D2}$	$2\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$
1	5	1	$3.26 \times 10^3$	71	-	100	$6.51 \times 10^4$	71	-	-	$1.21 \times 10^3$	24	-	-
1	10	1	97.0	41	-	100	-	-	-	-	$1.21 \times 10^3$	24	-	100
1	20	1	$1.47 \times 10^3$	58	$6.63 \times 10^3$	71	-	100	-	-	933	21	-	-
1	40	1	$1.02 \times 10^5$	35	$8.17 \times 10^4$	58	$8.17 \times 10^4$	71	-	-	$1.2 \times 10^3$	8	-	-
1	60	1	$8.84 \times 10^3$	45	$5.3 \times 10^3$	20	$2.94 \times 10^4$	58	-	-	220	6	-	100
1	80	1	783	11	326	4	$1.15 \times 10^3$	9	-	-	93.0	6	-	-
10	5	1	698	58	$3.14 \times 10^3$	71	-	100	-	-	-	-	-	-
10	10	1	161	38	674	45	-	-	-	-	-	-	-	-
10	20	1	287	45	-	100	$1.43 \times 10^4$	71	-	-	$1.87 \times 10^3$	30	-	100
10	40	1	$1.40 \times 10^4$	50	$1.29 \times 10^4$	28	$2.23 \times 10^4$	45	-	-	531	5	$6.82 \times 10^4$	71
10	60	1	$4.44 \times 10^3$	26	507	5	604	7	-	-	165	5	-	-
10	80	1	150	5	248	4	630	7	-	-	80.0	5	-	-
10	120	1	281	6	$1.32 \times 10^3$	8	411	6	-	-	62.0	4	-	-
20	5	1	97.0	41	877	71	-	-	-	-	-	-	-	-
20	10	1	140	35	562	41	-	-	-	-	-	-	-	-
20	20	1	$4.78 \times 10^3$	58	$1.08 \times 10^4$	50	-	-	-	-	$9.32 \times 10^3$	21	-	-
20	40	1	$6.31 \times 10^3$	38	$6.02 \times 10^3$	21	$1.76 \times 10^4$	45	-	-	366	7	-	-
20	60	1	247	6	377	4	438	6	-	-	148	5	-	-
20	80	1	91.0	4	230	3	534	6	-	-	62.0	5	-	-
20	120	1	247	6	$1.50 \times 10^3$	9	321	5	-	100	58.0	4	$9.40 \times 10^3$	58

# Re-interpretation: Providing Cross-section Limits

$$\frac{100}{\sqrt{\text{Number of Monte Carlo events survived}}} \%$$

• Gives a percentage uncertainty



$M_{D1}$	$\Delta M_{D+}$	$\Delta M_{D2}$	$2\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$2\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_A^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_B^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$	$3\ell \sigma_C^{95} \text{ (fb)}$	$\frac{100}{\sqrt{N_{MC}}}$
1	5	1	$3.26 \times 10^3$	71	-	100	$6.51 \times 10^4$	71	-	-	$1.21 \times 10^3$	24	-	-
1	10	1	97.0	41	-	100	-	-	-	-	$1.21 \times 10^3$	24	-	100
1	20	1	$1.47 \times 10^3$	58	$6.63 \times 10^3$	71	-	100	-	-	933	21	-	-
1	40	1	$1.02 \times 10^5$	35	$8.17 \times 10^4$	58	$8.17 \times 10^4$	71	-	-	$1.2 \times 10^3$	8	-	-
1	60	1	$8.84 \times 10^3$	45	$5.3 \times 10^3$	20	$2.94 \times 10^4$	58	-	-	220	6	-	100
1	80	1	783	11	326	4	$1.15 \times 10^3$	9	-	-	93.0	6	-	-
10	5	1	698	58	$3.14 \times 10^3$	71	-	100	-	-	-	-	-	-
10	10	1	161	38	674	45	-	-	-	-	-	-	-	-
10	20	1	287	45	-	100	$1.43 \times 10^4$	71	-	-	$1.87 \times 10^3$	30	-	100
10	40	1	$1.40 \times 10^4$	50	$1.29 \times 10^4$	28	$2.23 \times 10^4$	45	-	-	531	5	$6.82 \times 10^4$	71
10	60	1	$4.44 \times 10^3$	26	507	5	604	7	-	-	165	5	-	-
10	80	1	150	5	248	4	630	7	-	-	80.0	5	-	-
10	120	1	281	6	$1.32 \times 10^3$	8	411	6	-	-	62.0	4	-	-
20	5	1	97.0	41	877	71	-	-	-	-	-	-	-	-
20	10	1	140	35	562	41	-	-	-	-	-	-	-	-
20	20	1	$4.78 \times 10^3$	58	$1.08 \times 10^4$	50	-	-	-	-	$9.32 \times 10^3$	21	-	-
20	40	1	$6.31 \times 10^3$	38	$6.02 \times 10^3$	21	$1.76 \times 10^4$	45	-	-	366	7	-	-
20	60	1	247	6	377	4	438	6	-	-	148	5	-	-
20	80	1	91.0	4	230	3	534	6	-	-	62.0	5	-	-
20	120	1	247	6	$1.50 \times 10^3$	9	321	5	-	100	58.0	4	$9.40 \times 10^3$	58

# Conclusions

- Validated previous 8 TeV results from MadAnalysis, in CheckMATE
- New 13 TeV limits on DM mass splittings, beyond mono-jet for:
  - I2HDM: Important contributions from 3-lepton final states
  - MFDM: 3-lepton channel becomes leading role via  $D_3 \rightarrow Z(\rightarrow \ell\ell)D_1$  and Higgs becomes important for  $D_3 \rightarrow H(\rightarrow W^+W^-)D_1$
- Map of 2&3 lepton channel cross-section limits and efficiencies using CheckMATE for a wider DM re-interpretation analysis for our community